Show that $\left(\frac{\partial x}{\partial y}\right)_f \left(\frac{\partial y}{\partial f}\right)_x \left(\frac{\partial f}{\partial x}\right)_y = -1$

Show that
$$\left(\frac{\partial x}{\partial y}\right)_f \left(\frac{\partial y}{\partial f}\right)_x \left(\frac{\partial f}{\partial x}\right)_y = -1$$

$$df = \left(\frac{\partial f}{\partial x}\right)_{y} dx + \left(\frac{\partial f}{\partial y}\right)_{x} dy$$

$$dx = \left(\frac{\partial x}{\partial f}\right)_{y} df + \left(\frac{\partial x}{\partial y}\right)_{f} dy$$

$$dy = \left(\frac{\partial y}{\partial f}\right)_{x} df + \left(\frac{\partial y}{\partial x}\right)_{f} dx$$

$$dx = \left(\frac{\partial x}{\partial f}\right)_{y} df + \left(\frac{\partial x}{\partial y}\right)_{f} \left[\left(\frac{\partial y}{\partial f}\right)_{x} df + \left(\frac{\partial y}{\partial x}\right)_{f} dx\right]$$

$$dx = \left[\left(\frac{\partial x}{\partial f}\right)_{y} + \left(\frac{\partial x}{\partial y}\right)_{f} \left(\frac{\partial y}{\partial f}\right)_{x}\right] df + \left(\frac{\partial x}{\partial y}\right)_{f} \left(\frac{\partial y}{\partial x}\right)_{f} dx$$

matching the coefficients for dx and df

$$0 = \left(\frac{\partial x}{\partial f}\right)_y + \left(\frac{\partial x}{\partial y}\right)_f \left(\frac{\partial y}{\partial f}\right)_x \quad ; \quad 1 = \left(\frac{\partial x}{\partial y}\right)_f \left(\frac{\partial y}{\partial x}\right)_f$$

and final results

$$\left(\frac{\partial f}{\partial x}\right)_y \left(\frac{\partial x}{\partial y}\right)_f \left(\frac{\partial y}{\partial f}\right)_x = -1 \quad ; \quad \frac{1}{\left(\frac{\partial x}{\partial y}\right)_f} = \left(\frac{\partial y}{\partial x}\right)_f$$

Description of the state of a system

Energy: *U*

Entropy and temperature: S, T

Generalized forces and displacements:

pressure, volume: p, V

tension, length: τ, L

chemical potential, particle number: μ, N

magnetic field, magnetization: B, M

Response functions:
$$\kappa_x = -\frac{1}{V} \left(\frac{\partial V}{\partial p} \right)_x; \quad x = \{T, S\}$$

$$C_x = T \left(\frac{\partial S}{\partial T}\right)_x; \quad x = \{V, p\}$$

What constitutes a complete description of a system?

How many variables can be independent chosen?

How many EOSs do we need to know everything?

Are some EOSs more complete than others?

Internal energy as a function of S, V, N

$$U = U(S, V, N)$$

First order (partial) derivatives:

$$\left(\frac{\partial U}{\partial S}\right)_{V,N} = T; \quad \left(\frac{\partial U}{\partial V}\right)_{S,N} = -p; \quad \left(\frac{\partial U}{\partial N}\right)_{S,V} = \mu$$

The thermodynamic relations involving first order derivatives of the energy are sometimes called equations of state in a strict sense because they are experimentally determined. U(S,V,N) can be calculated from them by integration up to a (irrelevant) constant.

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Second order derivatives: $\left(\frac{\partial^2 U}{\partial S^2}\right)_{N,V}$

$$\left(\frac{\partial^2 U}{\partial S^2}\right), \quad \left(\frac{\partial^2 U}{\partial V^2}\right), \quad \left(\frac{\partial^2 U}{\partial N^2}\right)$$

$$\left(\frac{\partial^2 U}{\partial S \partial V}\right), \quad \left(\frac{\partial^2 U}{\partial V \partial S}\right), \quad \left(\frac{\partial^2 U}{\partial S \partial N}\right), \quad \left(\frac{\partial^2 U}{\partial N \partial S}\right) \text{ etc. (9 total)}$$

They are not all independent: because a state function must be a conservative function, the order of the partial derivatives in the second order partial derivatives can be switched without changing the value:

$$\left(\frac{\partial^2 U}{\partial S \partial V}\right) = \left(\frac{\partial^2 U}{\partial V \partial S}\right), \quad \left(\frac{\partial^2 U}{\partial S \partial N}\right) = \left(\frac{\partial^2 U}{\partial N \partial S}\right), \quad \left(\frac{\partial^2 U}{\partial S \partial N}\right) = \left(\frac{\partial^2 U}{\partial N \partial S}\right)$$

Maxwell relations

$$\left(\frac{\partial^2 U}{\partial S \partial V}\right) = \left(\frac{\partial^2 U}{\partial V \partial S}\right)$$
$$\left(\frac{\partial T}{\partial V}\right)_{S,N} = -\left(\frac{\partial p}{\partial S}\right)_{V,N}$$

For 3 independent state functions, there are 3 Maxwell relations for the internal energy U.

Response functions are second order partial derivatives. How many should there be?

Maxwell relations

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For 3 independent state functions, there are 3 Maxwell relations for the internal energy U.

Response functions are second order partial derivatives. How many should there be?

The ones we have defined so far all are for constant N:

$$\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_{p,N} \qquad \kappa_T = -\frac{1}{V} \left(\frac{\partial V}{\partial p} \right)_{T,N} \qquad \kappa_S = -\frac{1}{V} \left(\frac{\partial V}{\partial p} \right)_{S,N}$$

$$1 \quad (\partial S)$$

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$$c_p = \frac{1}{N}T\left(\frac{\partial S}{\partial T}\right)_{p,N}$$
 $c_V = \frac{1}{N}T\left(\frac{\partial S}{\partial T}\right)_{V,N}$

For constant N: U(S, V, N) = U(S, V)

For constant N: U(S, V, N) = U(S, V) there are 3 independent second order partial derivatives:

$$\left(\frac{\partial^2 U}{\partial S^2}\right), \quad \left(\frac{\partial^2 U}{\partial S \partial V}\right) = \left(\frac{\partial^2 U}{\partial V \partial S}\right), \quad \left(\frac{\partial^2 U}{\partial V^2}\right)$$

and hence there can only be 3 independent response functions.

Out 5 standard response functions we only need to know 3 and can calculate the other 2. This explains the existence of the well known thermodynamic identities:

$$c_p = c_V + \frac{\alpha^2 TV}{N\kappa_T}$$
$$\kappa_T = \kappa_S + \frac{\alpha^2 TV}{Nc_p}$$

What about the well known relation $\frac{C_p}{C_V} = \frac{\kappa_T}{\kappa_S}$?

Internal energy as a function of T, V, N

Common: we want to control the temperature: U = U(T, V, N)

First order (partial) derivatives:

$$\left(\frac{\partial U}{\partial T}\right)_{V,N} = ? \; ; \; \left(\frac{\partial U}{\partial V}\right)_{T,N} = ? \; ; \; \left(\frac{\partial U}{\partial N}\right)_{T,V} = ?$$

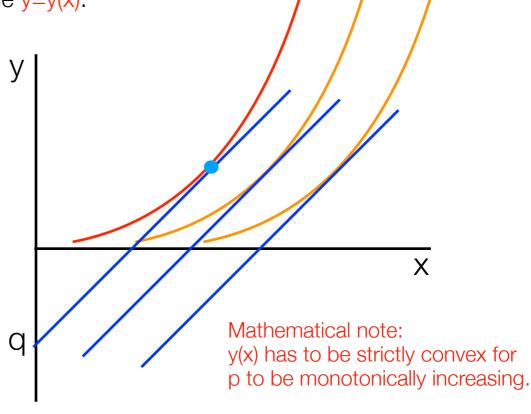
Loss of information!

Simple example: Consider a function of one variable y=y(x).

Any point (x,y=y(x)) defines a state.

The derivative is p=dy/dx = p(x) and inverting this relationship x = x(p) we find y=y(x(p)) = y(p).

Unfortunately any function y(x-a) gives the same y(p) and we have lost all information about a. To describe a state given by points (x,y) with the slope p we need to replace y with y-intercept of the slope $q(x,y) \rightarrow (p,q)$:



y=px+q or q=y-px which defines the Legendre transform y[p] of y(x).

Helmholtz Free Energy

$$T = \left(\frac{\partial U}{\partial S}\right)_{V,N}$$

$$F(T, V, N) \equiv U[T] = U - TS$$

Note on notation: Many textbooks use F to denote the Helmholtz Free Energy, but some use A (German for labor (Arbeit), including our textbook. I prefer F.

$$dF = \left(\frac{\partial F}{\partial T}\right)_{V,N} dS + \left(\frac{\partial F}{\partial V}\right)_{T,N} dV + \left(\frac{\partial F}{\partial N}\right)_{T,V} dN$$

To identify the partial derivatives of F start with total differential of F: dF

$$dF = dU - TdS - SdT$$

and use the thermodynamic identity,

$$dU = TdS - pdV + \mu dN$$

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To identify the partial derivatives of F start with total differential of F: dF

$$dF = dU - TdS - SdT$$

and using the thermodynamic identity,

$$dU = TdS - pdV + \mu dN$$

we arrive at

$$dF = -SdT - pdV + \mu dN$$

and
$$S = -\left(\frac{\partial F}{\partial T}\right)_{V,N}; \quad p = -\left(\frac{\partial F}{\partial V}\right)_{T,N}; \quad \mu = \left(\frac{\partial F}{\partial N}\right)_{T,V}$$

Enthalpy

$$H(S, p, N) \equiv U[p] = U + pV$$

Gibbs Free Energy

$$G(T, p, N) \equiv U[T, p] = U - ST + pV$$

Maxwell relations

Fundamental relations