

Buoyancy Effects  $\rightarrow l_b$  or  $l_o = (\epsilon/N^3)^{1/2}$

We reviewed B as a source of the examples of convectively driven mixing in daytime atmospheric boundary layer  
nighttime ocean boundary layer.

↑  
at least the upper boundary layer

Notes on Evolution of Turbulence in the Diurnal Warm layer Moulin et al 2019

neglecting transport terms:

$$\frac{\partial}{\partial t} \left( \frac{1}{2} \overline{u_i u_i} \right) = - \overline{u_i u_j} \frac{\partial U_i}{\partial x_j} - \frac{g}{\rho} \overline{u_3 \rho'} - \epsilon$$

$\uparrow$   $\uparrow$   
 $\rho$   $B$

Scale  $\frac{1}{2} \overline{u_i u_i} \sim \frac{3}{2} u^2$

Re stress  $\overline{u_i u_j} \sim C_{uv} u^2$

$\frac{\partial U_i}{\partial x_j} \sim S$  (for shear)

B is a the sink

lab.  
empirical  
result  
for shear  
flow  
 $C_{uv} \sim 0.3$

Important aside: how to quantify  $B$ ?

Consider steady state  $\mathcal{P} = B + \epsilon \leftarrow B$  is a sink

define flux Richardson #  $R_f = \frac{B}{\mathcal{P}}$

= ratio of PE gained by mixing  
mechanical prod. of tke

tells us how much of tke  
goes to  $\uparrow$  system PE

where  $B > 0$ ,  $R_f > 0$  since  $\mathcal{P} > 0$

In a lab or DNS we can compute changes  
in KE & PE due to mixing events

- this is impossible in natural flows

$\rightarrow u_i(z), \rho_i(z)$  initial  $u, \rho$  profiles

$\rightarrow u_f(z), \rho_f(z)$  final state.

$$\begin{aligned} \delta KE &= \int_0^h \rho_f(z) u_f^2(z) dz - \int_0^h \rho_i(z) u_i^2(z) dz \\ &\sim \rho_0 \int (u_f^2 - u_i^2) dz \end{aligned}$$

$$\delta PE = \int_0^h \rho_f(z) g z dz - \int_0^h \rho_i(z) g z dz$$

lab studies suggest that turbulence cannot be  
sustained if  $R_f \approx 0.2$

that is turbulence decays if  $\approx 20\%$  of the  
is lost to B.

let's rewrite  $B = -g_p \overline{w'p'}$  also  $N^2 = -\frac{g}{\rho} \frac{dp}{dz}$

define eddy diffusivity  $\overline{w'p'} = K_p \frac{dp}{dz}$

$$g/p \overline{w'p'} = g/p K_p \frac{dp}{dz} = -K_p N^2$$

$$B = K_p N^2$$

$$P = \epsilon + B$$

$$P/B = \epsilon/B + 1$$

$$1/R_f = \epsilon/K_p N^2 + 1$$

$$K_p = \frac{R_f}{(1-R_f)} \frac{\epsilon}{N^2}$$

let  $\Gamma = \frac{R_f}{1-R_f}$ ,  $R_f$  define by lab measurements

$$K_p \downarrow = \Gamma \epsilon / N^2 \quad \Gamma \approx 0.2$$

$$B \approx \Gamma \epsilon$$

$$\frac{d}{dt} (\frac{3}{2} u^2) = C_{uv} u^2 S - \epsilon (\Gamma + 1)$$

inviscid estimate of  $\epsilon$   $\epsilon = C_\epsilon u^3 / l$   $C_\epsilon \approx 0.5$   
empirical

replace  $l$  with  $l_b$

$$\frac{d\epsilon}{dt} = \frac{2}{3} \epsilon \left[ C_{uv} S - C_\epsilon^{2/3} N(1+r) \right]$$

assume that  $N(t)$  changes slowly relative to  $\epsilon$   
 $\uparrow$  constant

general solution  $\epsilon(t) = C_0 e^{t/\delta}$   
 $\uparrow$  const. of integration

$$\delta = \left[ \frac{2}{3} (C_{uv} S - \alpha N) \right]^{-1} \quad \alpha = C_\epsilon^{2/3} (1+r) \approx .75$$

when  $\delta < 0$ , or  $C_{uv} S < \alpha N$ ,  $\epsilon(t)$  decays.

when  $\delta > 0$ , or  $C_{uv} S > \alpha N$ ,  $\epsilon(t)$  grows.

$\rightarrow$  let's go back to paper Fig 7.

### Dynamics of Temperature (Scalar) Fluctuations

temperature variance eq'n.  $\leftarrow$  related to the

$$\frac{\partial}{\partial t} \left( \frac{1}{2} \overline{\theta^2} \right) + U_j \frac{\partial}{\partial x_j} \left( \frac{1}{2} \overline{\theta^2} \right) = - \frac{\partial}{\partial x_j} \left[ \frac{1}{2} \overline{u_j \theta^2} - \delta \frac{\partial}{\partial x_j} \left( \frac{1}{2} \overline{\theta^2} \right) \right]$$

$\uparrow$   
time

$\uparrow$   
adv

$$- \overline{u_j \theta} \frac{\partial \overline{\theta}}{\partial x_j}$$

$\uparrow$   
 $\rho_\theta$

$$- \delta \frac{\partial \overline{\theta}}{\partial x_j} \frac{\partial \overline{\theta}}{\partial x_j}$$

$\uparrow$   
diffusive  
molecular diffusion  
of  $\theta^2$

$\nwarrow$   
divergence

consider steady-state ( $\partial_t = 0$ ), homogeneous.

stratified  $\frac{\partial \Theta}{\partial z} \neq 0$ , neglect transport

$$-\overline{w\Theta} \frac{\partial \Theta}{\partial z} = \gamma \frac{\partial \Theta}{\partial x_j} \frac{\partial \Theta}{\partial x_j}$$

gradient  $\partial_\Theta$  = molecular diffusion of  $\overline{\Theta^2}$

$$\chi = 2\gamma \frac{\partial \Theta}{\partial x_j} \frac{\partial \Theta}{\partial x_j} \leftarrow \text{something that can be measured using fast T sensors.}$$

$$\overline{w\Theta} \frac{\partial \Theta}{\partial z} = \frac{\chi}{2}$$

$$\overline{w\Theta} = -K_T \frac{\partial \Theta}{\partial z}$$

$$K_T = \frac{\chi}{2\overline{\Theta_z^2}} \leftarrow K_p = \frac{\Gamma \epsilon}{N^2}$$

$$K_T = K_p \Rightarrow \Gamma = \frac{\chi N^2}{2\epsilon \overline{\Theta_z^2}}$$

extensive comparisons  $\Rightarrow \Gamma \approx 0.2$