

Group theory, Topology and Spin-1/2 Particles

From Dirac's belt to fermions

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Brussels Summer School of Mathematics 2022

1. Dirac's belt trick and rotations
2. Homotopy theory
3. Quantum spin and $SU(2)$
4. Conclusion

Dirac's belt trick and rotations

You need:

- a belt (not necessarily Dirac's)
- a heavy book

Rules:

1. you can only move the end of the belt
2. you cannot twist or rotate it

Goal: untwist a 2π -twist.

\Rightarrow it turns out to be impossible ! One turn negates the twist: $2\pi \rightarrow -2\pi$.

Therefore, possible for a 4π twist ...

Why is that ?

Space of rotations: $\text{SO}(3)$ as a group

Rotations in 3-dimensional space: matrices that acts on \mathbb{R}^3 s.t.

1. preserve the **scalar product**: $O^T O = \mathbb{1}$ ($\Leftrightarrow O$ is orthogonal)
2. preserve the **orientation**: $\det O = 1$

Special orthogonal group

$\text{SO}(3)$ is the set of 3×3 real matrices such that $O^T O = \mathbb{1}$ and $\det O = 1$.

Three “fundamental” rotations:

$$x : \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos \theta & -\sin \theta \\ 0 & \sin \theta & \cos \theta \end{bmatrix} \quad y : \begin{bmatrix} \cos \theta & 0 & -\sin \theta \\ 0 & 1 & 0 \\ \sin \theta & 0 & \cos \theta \end{bmatrix} \quad z : \begin{bmatrix} \cos \theta & -\sin \theta & 0 \\ \sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

\Rightarrow It forms a **group**.

Space of rotations: $\text{SO}(3)$ as a topological space

Fundamental data that describes a rotation:

- an **axis** of rotation, i.e. a unit vector \vec{n} $\rightarrow 2$ parameters
- an **angle** of rotation $\theta \in [-\pi, \pi]$ (with $-\pi \sim \pi$) $\rightarrow 1$ parameter

The space of rotations can then alternately be defined as a **3-sphere of radius π and its antipodal points identified**:

$$\boxed{\text{SO}(3) \cong B^3(\pi) / \sim}$$

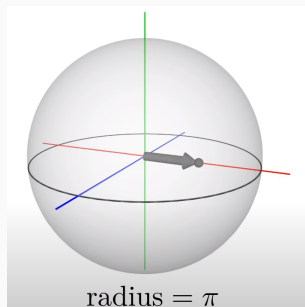
and for each point:

direction \leftrightarrow axis

norm \leftrightarrow angle

\Rightarrow It forms a **topological space**.

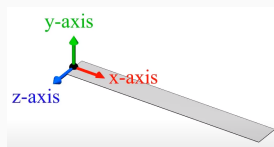
(group + topological space = Lie group)



Back to the belt

Mathematical description of the belt ?

- ▷ a belt is a strip, which is just a **path** + an **orientation**.
- ▷ given axis on the middle line along the belt, each set of axis is related by a rotation
- ▷ a belt configuration is equivalent to a continuous set of axis and therefore to a continuous set of translations, i.e. a **path in $SO(3)$**

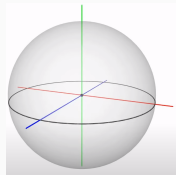


There is a bijection:

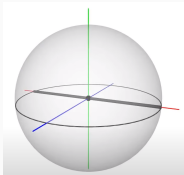
$$\text{belt configuration} \Leftrightarrow \text{path in } SO(3)$$

This gives us a new language to analyze the problem !

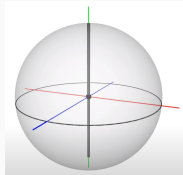
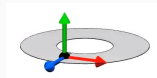
trivial rotation



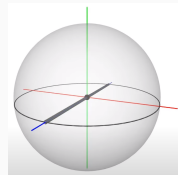
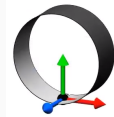
2π *x*-rotation



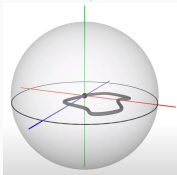
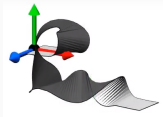
2π *y*-rotation



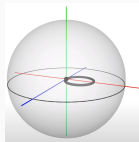
2π *z*-rotation



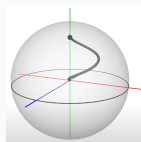
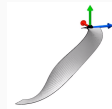
random rotation



closed



open path



<u>Belt</u>		<u>Path</u>
specific configuration	\longleftrightarrow	specific path
moving the ends	\longleftrightarrow	continuous deformation
ends have same orientation	\longleftrightarrow	closed path (loop)
can be flattened	\longleftrightarrow	contractible

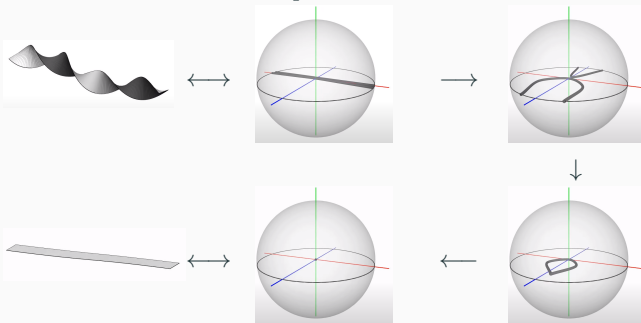
Back to Dirac's belt trick:

1. ends of the belt have same orientation \rightarrow we consider loops
(passing through the origin)
2. moving the ends of the belt \rightarrow continuous deformation
3. belt in original (flat) position \rightarrow trivial path

The question then becomes: **which loops are contractible ?**

Problem solved ?

- 4π -twist: We saw in the beginning the the 4π -twist can be flattened, how can we see this in terms of paths ?



\Rightarrow the 4π -twist is **contractible** ! Great.

- 2π -twist: we “clearly” see that is not contractible... no ?! Great..?..

Wierd aftertaste: our “proof” is good to show contractibility but bad to show non-contractibility and it only works for simple examples.

\Rightarrow We want a consistent and general way of studying paths in topological spaces.

Homotopy theory

Starting observation: depending on the topological space, all loops might not be contractible. Moreover, some loops are “fundamentally different” from each other.

Examples: \mathbb{R}^3 , S^2 , \mathbb{T}^2 , etc.

Paths and homotopies

For a topological space X :

- *Path* in X : continuous map $\gamma : [0, 1] \rightarrow X$, *loop* if closed
- γ_1 and γ_2 are *homotopically equivalent* if one can be deformed into the other: there exists $H : [0, 1] \times [0, 1] \rightarrow X$ such that

$$H(0, t) = \gamma_1(t) \quad \text{and} \quad H(1, t) = \gamma_2(t).$$

This is an equivalence relation.

For each $x_0 \in X$, we define

$$\pi_1(X, x_0) = \{\text{all loops based at } x_0\} / \sim_{\text{hom.}},$$

it is the set of “fundamentally different” loops passing through x_0 .

Fundamental group

Group structure:

- **Product** of paths: $\gamma_1 \cdot \gamma_2 = \gamma_1$ then γ_2
- **Inverse** path: $\gamma^{-1} = \gamma$ traversed in the opposite direction
- **Neutral** path: $e = \text{constant path at the identity}$
- For equivalence classes: $[\gamma_1] \cdot [\gamma_2] = [\gamma_1 \cdot \gamma_2]$ and $[\gamma]^{-1} = [\gamma^{-1}]$

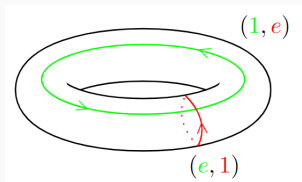
Important fact: up to isomorphism, $\pi_1(X, x_0)$ does not depend on x_0
 \Rightarrow we denote it as $\pi_1(X)$, it is called the **fundamental group** of X .

Contractible loops are the ones in $[e]$ (\sim to a point).

How to compute the fundamental group ? Difficult task, not discussed here.

Examples:

- $\pi_1(\mathbb{R}^2) = \{e\}$
- $\pi_1(S^2) = \{e\}$
- $\pi_1(\mathbb{T}^2) = \mathbb{Z} \times \mathbb{Z}$
- $\pi_1(\mathbb{R}^2 \setminus \{pt\}) = \mathbb{Z}$

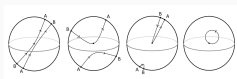


Question we had: **are all loops in $\mathrm{SO}(3)$ contractible ?**

In homotopy language: **is $\pi_1(\mathrm{SO}(3))$ trivial ?**

The belt trick is a way of physically demonstrating that the fundamental group of $\mathrm{SO}(3)$ is \mathbb{Z}_2 .

In quantum mechanics



Quantum spin and $SU(2)$

Conclusion

Three different block environments are pre-defined and may be styled with an optional background color.

Some text.

Default

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Alert

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Example

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Sometimes, it is useful to add slides at the end of your presentation to refer to during audience questions.

The best way to do this is to include the `appendixnumberbeamer` package in your preamble and call `\appendix` before your backup slides.

metropolis will automatically turn off slide numbering and progress bars for slides in the appendix. [?]



A. Hatcher.

Algebraic Topology.

Algebraic Topology. Cambridge University Press, 2002.



N. Miller.

Representation theory and quantum mechanics, 2018.