

Course Outline:
Based on Lectures from Winter, 2024 by Prof. Antony Humphries.

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1 Introduction

1.1 Definitions

↪ Definition 1.1: Differential equation

A *differential equation* (DE) is an equation with derivatives. *Ordinary* DE's (ODE) will be covered in this course; other types (PDE's, SDE's, DDE's, FDE's, etc.) exist as well but won't be discussed. ODE's only have one independent variable (typically, $y = f(x)$ or $y = f(t)$).

⊗ Example 1.1: A Trivial Example

$\frac{dy}{dx} = 6x$. Integrating both sides:

$$\int \frac{dy}{dx} dx = \int 6x dx \implies y(x) = 3x^2 + C.$$

⊗ Example 1.2: Another One

$$\frac{d^2u}{dt^2} = 0 \implies y = at + b.$$

↪ Definition 1.2: Order

The order of a differential equation is defined as the order of the highest derivative in the equation.

1.2 Initial Values

Remark 1.1. Note the existence of arbitrary constants in the previous examples, indicating infinite solutions. We often desire unique solutions by fixing these coefficients. For first order ODEs, we simply specify a single initial condition (say, some $y(x_0) = \alpha_0$). For higher order ODEs of degree n , we can either specify $n - 1$ initial conditions for $n - 1$ derivatives (say, $y(x_0) = \alpha_0, y'(x_0) = \beta_0$), or boundary conditions (say, $y(x_0) = \alpha_0, y(x_1) = \alpha_1$) where values for the solution itself are specified.

⊗ Example 1.3: A Less Trivial Example

$\frac{dy}{dx} = y$. We cannot simply integrate both sides as before, as we have no way to know what $\int y dx$ (the RHS) is equal to. We can fairly easily guess that $y = e^x$ is a solution; its derivative is equal to itself, hence it does indeed solve the equation. This is not the

only solution; indeed, given $y = ce^x$, we have

$$\frac{dy}{dx} = ce^x = y = ce^x.$$

Luckily, we were rather limited in how many places constants could appear; this doesn't always hold.

1.3 Physical Applications

⊛ Example 1.4: Simple Pendulum

Let θ be the angle of a pendulum of mass m from vertical and length l . Then, we have the equation of motion

$$ml\ddot{\theta} = -mg \sin \theta \implies \ddot{\theta} + \frac{g}{l} \sin \theta = 0 \implies \ddot{\theta} + \omega^2 \sin \theta = 0.$$

Take θ small, then, $\sin \theta \approx \theta$. Then, $\ddot{\theta} + \omega^2 \theta = 0$. This is linear simple harmonic motion, and has periodic solutions; how do we know this is a valid solution to the non-linear model?

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⊛ Example 1.5: Lorenz Equations

$$\begin{aligned}\frac{dx}{dt} &= \sigma(y - x) \\ \frac{dy}{dt} &= rx - y - xz \\ \frac{dz}{dt} &= xy - bz\end{aligned}$$

These are a famous set of equations originally derived from atmospheric modeling, known for its chaotic behavior for particular parameters. This is a nonlinear system of de's, and beyond the scope of this class (indeed, it is not solvable exactly).

1.4 Uniqueness

Given an ODE of the general form $y^{(n)} = f(t, y, y', \dots, y^{(n-1)})$, if we wish to determine $y^{(n)}(t_0)$ uniquely, we need to specify the initial conditions

$$y(t_0), y'(t_0), \dots, y^{(n-1)}(t_0).$$

Moreover, this not only determines uniqueness of $y^{(n)}(t_0)$, but the uniqueness of solution y for $t \in I$ for some "interval of validity" I .

↪ **Definition 1.3: Autonomous/Nonautonomous**

An ODE of the form

$$y^{(n)} = f(y, y', \dots, y^{(n-1)})$$

is called *autonomous*; that is, if it has no explicit dependence on the independent variable. Otherwise, the system is called *nonautonomous*.

↪ **Definition 1.4: Linear/Nonlinear**

Linear ODEs of dimension n have a solution space which is a vector space of dimension n . As a result, solutions can be written as a linear combination of n basis solutions (or “fundamental set of solutions”). Solutions to nonlinear ODEs cannot be written this way (except locally).

Alternatively (but equivalently), if we can write an n th order ODE in the form

$$a_n(t)y^n(t) + \dots + a_1(t)y'(t) + a_0(t)y(t) = g(t),$$

or equivalently,

$$\sum_{i=0}^n a_i(t)y^i(t) = g(t), \quad \circledast$$

where each $a_i(t)$ and $g(t)$ are known functions of t , then we say that the ODE is linear. Otherwise, it is nonlinear.

⊛ Example 1.6

The pendulum

$$\ddot{\theta} + \omega^2 \sin \theta = 0$$

is autonomous and linear;

$$\ddot{\theta} + \omega^2 \sin \theta = 0$$

is autonomous and nonlinear, due to the $\sin \theta$ term (indeed, this is a nonlinear oscillator equation); a damped-forced oscillator

$$\ddot{\theta} + k^2 \dot{\theta} + \omega^2 \theta = A \sin(\mu t)$$

is nonautonomous and linear.

Remark 1.2. Note that the following definitions apply only to linear ODEs.

↪ **Definition 1.5: Homogeneous/Nonhomogeneous**

A linear ODE of the form \circledast is *homogeneous* if $g(t) = 0$; otherwise it is *nonhomogeneous*.

↪ **Definition 1.6: Constant/Variable**

A linear ODE of the form \ast is *constant coefficient* if $a_j(t) = \text{constant} \ \forall j$; if at least one a_j not constant, it is *non-constant* or *variable coefficient*.

Remark 1.3. Note that while we define linearity of ODEs in terms of the form of $y^{(n)} = f(t, y, \dots)$, this more “helpfully” relates to the form of the solution of such an ODE, which is indeed linear.

1.5 Solutions

Given an n order ODE $y^{(n)} = f(t, y, \dots)$, and assuming f continuous, then for $y(t)$ to be a solution, we need y to be n -times differentiable; hence, $y, \dots, y^{(n-1)}$ must all exist and be continuous. Then, $y^{(n)}$, being a continuous function of continuous functions, is, itself, continuous.

↪ **Definition 1.7: Solution**

The function $y(t) : I \rightarrow \mathbb{R}$ is a solution to an ODE on an interval $I \subseteq \mathbb{R}$ if it is n -times differentiable on I , and satisfies the ODE on this interval.

Given an well-defined IVP with $n - 1$ initial values defined at t_0 , then $y(t)$ is a solution if $t_0 \in I$, y satisfies the initial values, and $y(t)$ is a solution on the interval.

↪ **Definition 1.8: Interval of Validity**

The largest I on which $y(t) : I \rightarrow \mathbb{R}$ solves an ODE is called the *interval of validity* of the problem.

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