

High-pressure effects on nonfluorinated BiS₂-based superconductors La_{1-x}M_xOBiS₂ ($M = \text{Ti}$ and Th)

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Layered LnOBiS₂ (where Ln represents a lanthanide element) compounds with Ln = La, Ce, Pr, Nd, and Yb can be rendered conducting and superconducting via two routes, substitution of F for O or the tetravalent ions Ti, Zr, Hf, and Th for trivalent Ln ions. Electrical resistivity measurements on nonfluorinated La_{0.80}Ti_{0.20}OBiS₂ and La_{0.85}Th_{0.15}OBiS₂ superconductors were performed between ~ 1.5 and 300 K and under pressure up to 2.4 GPa. For both compounds, the superconducting transition temperature T_c , which is ~ 2.9 K at ambient pressure, gradually increases with pressure to 3.2–3.7 K at ~ 1 GPa, above which it is suppressed and the superconducting transitions become very broad. Measurements of the normal-state electrical resistivity of the two compounds reveal discontinuous changes in the resistivity as a function of pressure at ~ 0.6 GPa. Surprisingly, above 1.3 GPa, semiconductinglike behavior reappears in La_{0.80}Ti_{0.20}OBiS₂. This paper reveals a new high-pressure phase of La_{1-x}M_xOBiS₂ containing the tetravalent ions $M = \text{Ti}, \text{Th}$ which does not favor superconductivity. In contrast, application of pressure to fluorinated LaO_{0.5}F_{0.5}BiS₂ produces an abrupt tetragonal-monoclinic transition to a metallic phase with an enhanced T_c . These results demonstrate that the response of the normal and superconducting properties of LaOBiS₂-based compounds depends strongly on the atomic site where the electron donor ions are substituted.

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I. INTRODUCTION

Following the discovery of superconductivity in the layered bismuth oxysulfide compound Bi₄O₄S₃, an intense effort to find new BiS₂-based superconductors quickly ensued as a result of the remarkable flexibility of the BiS₂ layers in forming compounds with new chemical compositions [1–4]. By doping the BiS₂ layers with electrons via appropriate chemical substitutions, superconductivity was found in Ln(O, F)BiS₂ (Ln = La–Nd, Yb), (Sr, La)FBiS₂, (La, M)OBiS₂ ($M = \text{Ti}, \text{Zr}, \text{Hf}, \text{Th}$), Bi₃O₂S₃, La(O, F)BiSe₂, EuBiS₂F, and Eu₃Bi₂S₄F₄ with values of the superconducting transition temperature T_c ranging from 2.7 to 10.6 K [2,5–12]. One of the fascinating features in the behavior of several of the BiS₂-based superconductors, including LnO_{0.5}F_{0.5}BiS₂ (Ln = La–Nd, Yb), (La, Sm)O_{0.5}F_{0.5}BiS₂, Sr_{0.5}Re_{0.5}FBiS₂ (Re = La–Pr, Sm) (where Re represents a rare earth element), LaO_{0.5}F_{0.5}BiSe₂, and EuBiS₂F, is the pressure-induced enhancement of T_c that occurs at about 0.5–2.2 GPa, revealing the existence of two distinct superconducting phases [13–19]. This phenomenon is believed to be associated with a structural phase transition, which has been demonstrated in the case of LaO_{0.5}F_{0.5}BiS₂, from a tetragonal structure to a monoclinic structure with a sudden change in unit-cell volume at an external pressure (P) of ~ 0.5 GPa [20].

In addition to the requisite superconducting BiS₂ layers, there are other similarities in the compounds reported to have a second superconducting phase under high pressure: (1) All of these compounds contain F, which resides in blocking layers and donates conduction electrons to the BiS₂ layers, and (2) the high-pressure phase has a significantly larger value of T_c than

the low-pressure phase. These similarities raise the issue of whether F is essential for the formation of the pressure-induced high- T_c phase. Recently, superconductivity was also found in single-crystalline CeOBiS₂ at 1.3 K, and the value of T_c could gradually be enhanced under an external pressure [21]. This phenomenon was suggested to be related to Ce valence fluctuations, which can be tuned via the application of pressure; however, it is still not clear if there are concomitant phase transitions. On the other hand, the compounds La_{1-x}M_xOBiS₂ ($M = \text{Ti}, \text{Zr}, \text{Hf}, \text{Th}$) were reported to have the same crystal structure, very similar lattice parameters, and values of T_c (~ 3 K) similar to that of LaO_{0.5}F_{0.5}BiS₂ [8]. The M ions in La_{1-x}M_xOBiS₂ are tetravalent and contribute conduction electrons to the BiS₂ layers in the absence of F; hence, this system is ideal for investigating whether non-F-substituted BiS₂ compounds have a second pressure-induced high- T_c phase. High-pressure studies of La_{1-x}M_xOBiS₂ also have the potential for uncovering new superconducting phases in BiS₂-based superconductors and will be very helpful in developing an understanding of the nature of the pressure-induced low- T_c to high- T_c phase transition in these materials.

II. EXPERIMENTAL DETAILS

Polycrystalline samples of La_{0.80}Ti_{0.20}OBiS₂ and La_{0.85}Th_{0.15}OBiS₂ were synthesized by solid-state reaction as described elsewhere [8]. Geometric factors used in determining the electrical resistivity for each sample were measured before applying pressure. Hydrostatic pressure was generated by using a clamped piston-cylinder cell in which a 1:1 mixture of *n*-pentane and isoamyl alcohol by volume was used as the pressure-transmitting medium. The pressures applied to the samples were estimated by measuring the T_c

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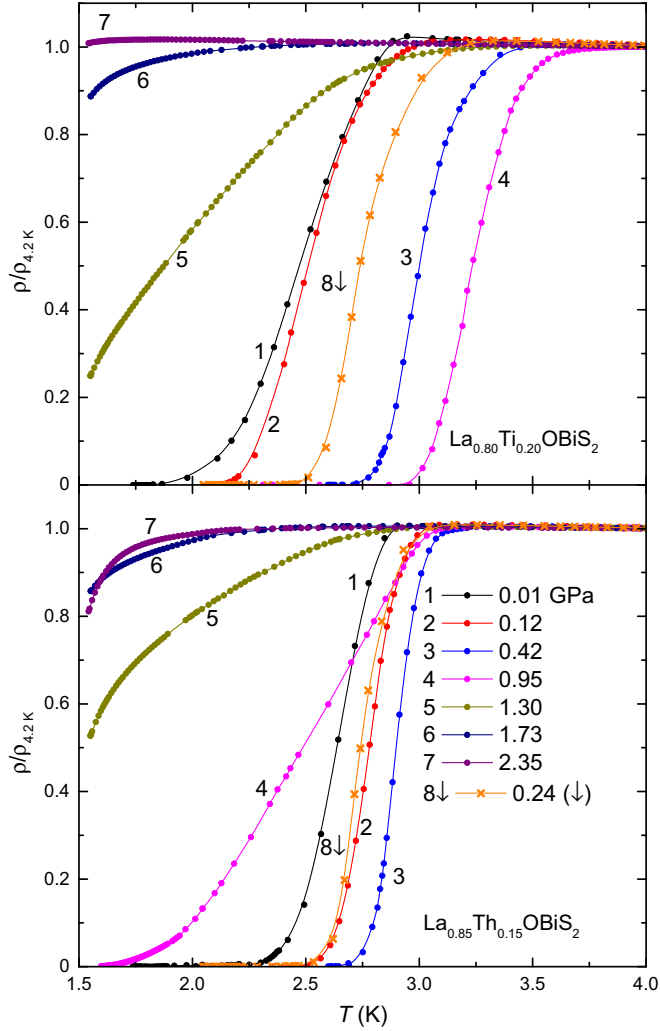


FIG. 1. Temperature T dependence of electrical resistivity ρ , normalized to its normal-state value at 4.2 K, below 4 K for $\text{La}_{0.80}\text{Ti}_{0.20}\text{OBiS}_2$ (upper panel) and $\text{La}_{0.85}\text{Th}_{0.15}\text{OBiS}_2$ (lower panel). The numbers and symbols with different colors in both panels are used to indicate the pressures listed in the lower panel. Most of the $\rho(T)$ data (represented by circles) were obtained upon increasing pressure. Crosses and the \downarrow on the right side of the number 8 in this figure as well as in the following figures throughout this paper are used to represent data taken upon decreasing pressure.

of a high-purity ($> 99.99\%$ -purity) Sn disk inside the sample chamber of the cell and comparing the measured values with the well-determined $T_c(P)$ of high-purity Sn [22]. Electrical resistance measurements under high pressure up to ~ 2.4 GPa were performed in a temperature range between 1.5 and 280 K using a standard four-probe method in a pumped ^4He dewar.

III. RESULTS AND DISCUSSION

The resistive superconducting transition curves $\rho(T)$, normalized to the value of ρ in the normal state at 4.2 K for polycrystalline $\text{La}_{0.80}\text{Ti}_{0.20}\text{OBiS}_2$ and $\text{La}_{0.85}\text{Th}_{0.15}\text{OBiS}_2$ samples are shown in the upper and lower panels of Fig. 1, respectively. The superconducting transitions are very sharp for pressures below 0.95 GPa for both samples; however, the

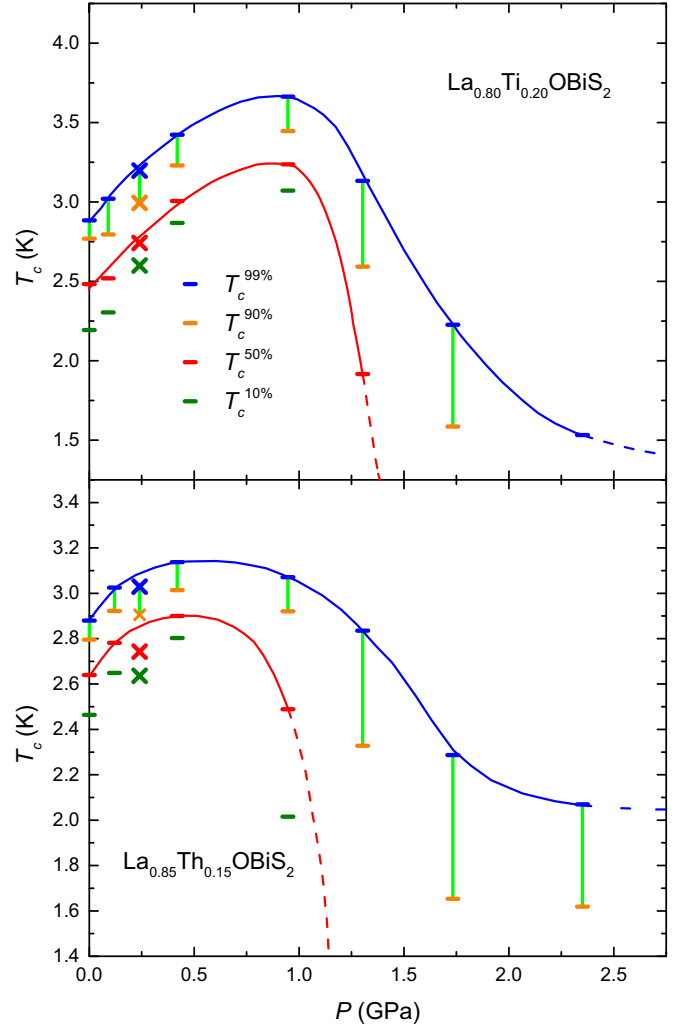


FIG. 2. Superconducting transition temperature T_c vs pressure for $\text{La}_{0.80}\text{Ti}_{0.20}\text{OBiS}_2$ (upper panel) and $\text{La}_{0.85}\text{Th}_{0.15}\text{OBiS}_2$ (lower panel). The temperature at which ρ drops to 99%, 90%, 50%, and 10% of its normal-state value above the superconducting transition is indicated by symbols $T_c^{99\%}$, $T_c^{90\%}$, $T_c^{50\%}$, and $T_c^{10\%}$, respectively. The lengths of the vertical lines represent the difference between $T_c^{99\%}$ and $T_c^{90\%}$ and are measures of the broadening of the superconducting transition under pressure. The crosses represent data taken upon decreasing pressure. The solid and dashed lines are guides to the eye for the experimental data and the corresponding extrapolations, respectively.

transitions become quite broad at higher pressures, and the values of ρ remain finite down to 1.5 K. The relationship between T_c and pressure for the two compounds can be seen in Fig. 2. $T_c^{99\%}$, $T_c^{90\%}$, $T_c^{50\%}$, and $T_c^{10\%}$ represent the temperatures where ρ falls to 99%, 90%, 50%, and 10% of its normal-state value, respectively. The value of $T_c^{99\%}$ is 3.5 K for $\text{La}_{0.80}\text{Ti}_{0.20}\text{OBiS}_2$ at ambient pressure and increases slightly with increasing pressure to 3.7 K at 0.95 GPa above which $T_c^{99\%}$ starts decreasing gradually. A similar dome-shaped T_c vs P superconducting phase boundary can also be observed in the case of $\text{La}_{0.85}\text{Th}_{0.15}\text{OBiS}_2$, although the values of T_c are slightly different than those of $\text{La}_{0.80}\text{Ti}_{0.20}\text{OBiS}_2$. For both compounds, the difference between $T_c^{99\%}$ and $T_c^{90\%}$ (indicated by the green vertical lines in Fig. 2), which is related to

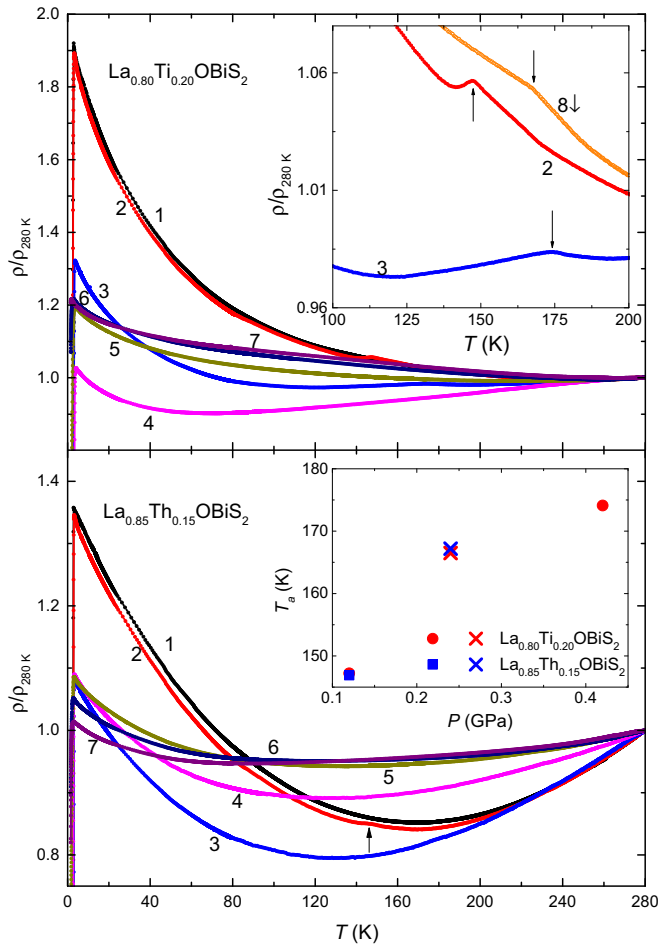


FIG. 3. Temperature T dependence of the electrical resistivity ρ , normalized to the value of ρ at 280 K, ($\rho/\rho_{280\text{ K}}$) of $\text{La}_{0.80}\text{Ti}_{0.20}\text{OBiS}_2$ (upper panel) and $\text{La}_{0.85}\text{Th}_{0.15}\text{OBiS}_2$ (lower panel) at various pressures. The pressure corresponding to each ($\rho/\rho_{280\text{ K}}$) curve is indicated in Fig. 1 with the same color. The upper inset: enlargement of the T dependence of $\rho/\rho_{280\text{ K}}$ where the anomalies in $\rho(T)$ can be observed. The orange curve represents the data obtained at 0.24 GPa during the pressure releasing cycle. The lower inset: temperature at which the anomaly in $\rho(T)$ occurs, (T_a) vs pressure (P). The black arrows in the figure indicate the location of the anomalies in the resistivity.

the width of the superconducting transition, becomes much larger at pressures at which T_c starts decreasing (0.95 GPa), reflecting dramatic changes in the superconductivity. During the pressure releasing cycle, the values of T_c are reversible, and the superconducting transitions become sharp again when the samples have returned to low pressures as can be seen in Figs. 1 and 2 (data represented by the crosses). These features in the T_c vs P data indicate that there are significant differences in the superconducting behavior below and above 0.95 GPa.

To obtain further information about the origin of the suppression of superconductivity at high pressure, we investigated the normal-state resistivity of the two compounds under pressure. Upon warming, $\rho(T)$ of $\text{La}_{0.85}\text{Th}_{0.15}\text{OBiS}_2$ first decreases significantly, passes through a minimum, and then increases nearly linearly at higher temperatures as shown in Fig. 3. Such minima in the normal-state resistivity are also

observed for $\text{La}_{0.80}\text{Ti}_{0.20}\text{OBiS}_2$ under pressure and other BiS₂-based superconductors [8,11,13,23,24]. In addition, subtle anomalies in resistivity, indicated by the black arrows in Fig. 3, were observed in both compounds when pressure was applied to the samples. Upon a further increase in pressure, the temperature at which the anomaly displays a kink (T_a) increases with increasing pressure (see the lower inset of Fig. 3); however, at pressures above 0.42 GPa, such anomalies in $\rho(T)$ can no longer be observed. Very recently, a similar kink in $\rho(T)$ was reported in nonsuperconducting LaOBiSe_2 at ambient pressure, which may related to a charge-density wave transition and can be suppressed by F substitution [25]. In this paper, the appearance of these anomalies is possibly due to a temperature-driven phase transition; however, the nature of the transition will require further investigation.

In this paper, we use the parameter $a(T)$, which represents the slope of $\rho(T)$, ($d\rho/dT$), at a certain temperature to characterize the tendency of the compound to exhibit metallic- or semiconductinglike behavior. For simple metals at high temperatures, a is related closely to the scattering of conduction electrons by phonons. Although the value of a does not have a clear physical significance for semiconductors, a positive value of a suggests that the sample is metalliclike and a negative value of a means it is semiconductinglike. In this paper, the quantity $a(T)$ represents the tendency of the compounds to exhibit metallic- or semiconductinglike behavior at different pressures. As shown in Fig. 4(a), the values of a at ~ 10 K ($a_{10\text{ K}}$) for the two compounds are negative at ambient pressure, suggesting semiconductinglike behavior at low temperatures. Upon application of an external pressures, $a_{10\text{ K}}$ first increases rapidly at low pressures but relatively moderately at high pressures, resulting in a kink in the pressure dependence of $a_{10\text{ K}}$ at ~ 0.7 GPa. Although the values of $a_{10\text{ K}}$ are all negative at pressures up to 2.4 GPa, the increase in $a_{10\text{ K}}$ with pressure clearly indicates suppression of the semiconductinglike behavior. The kinks in $a_{10\text{ K}}(P)$ found at ~ 0.7 GPa suggests there are some fundamental differences between behaviors of the two compounds at low and high pressures, which are consistent with the decrease in T_c and broadening of the superconducting transitions at high pressure discussed previously.

Although the pressure dependences of a at low temperatures are quite similar for the two compounds, the values of a at 250 K ($a_{250\text{ K}}$) for $\text{La}_{0.85}\text{Th}_{0.15}\text{OBiS}_2$ are positive and almost the same at low pressures, in contrast to the remarkable enhancement of $a_{250\text{ K}}$ from negative to positive values in the case of $\text{La}_{0.80}\text{Ti}_{0.20}\text{OBiS}_2$ as shown in Fig. 4(b). As indicated by the dashed lines in Fig. 4(b), the pressure dependence of $a_{250\text{ K}}$ dramatically decreases at higher pressures for both compounds, which suggests not only pressure-induced suppression of metalliclike behavior at high temperatures for the two compounds, but also indicates the appearance of high-pressure phases that exhibit different types of electronic conduction compared to the corresponding low-pressure phases. This behavior can be observed (Fig. 3) at high temperatures above ~ 60 K for the Ti-substituted sample and above ~ 140 K for the Th-substituted sample. It should be noted that the values of $a_{250\text{ K}}$ for $\text{La}_{0.80}\text{Ti}_{0.20}\text{OBiS}_2$ are even negative above 1.3 GPa, indicating the reappearance of semiconductinglike behavior at high temperatures.

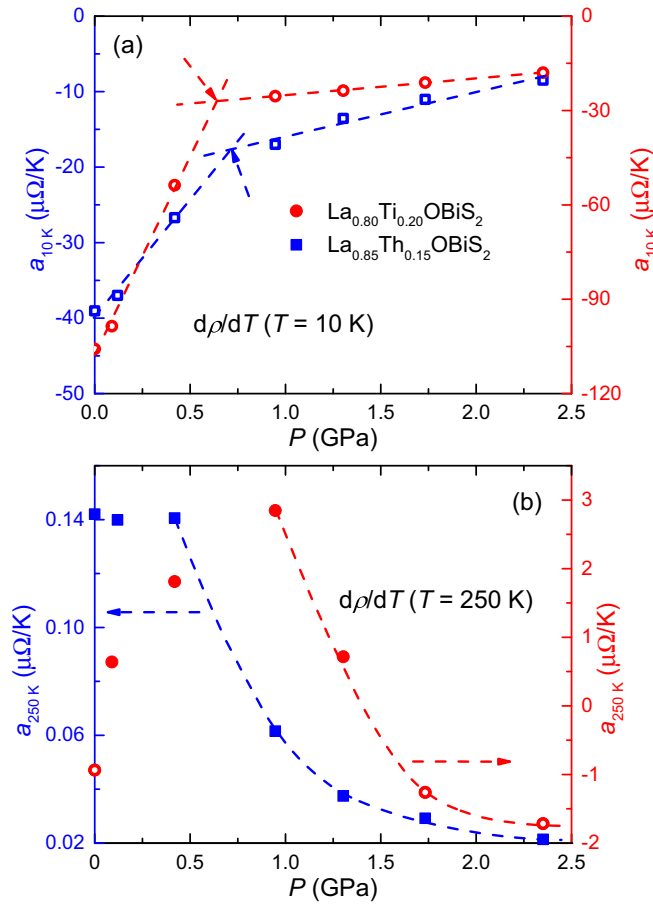


FIG. 4. Evolution of the slopes of $\rho(T)$ at (a) ~ 10 K and at (b) ~ 250 K with increasing pressure for $\text{La}_{0.80}\text{Ti}_{0.20}\text{OBiS}_2$ (red circles) and $\text{La}_{0.85}\text{Th}_{0.15}\text{OBiS}_2$ (blue squares), respectively. The open circles indicate that the values are negative, suggesting semiconductinglike behavior. The dashed lines are guides to the eye.

Either the suppression of metallic conductivity or the reappearance of semiconductinglike behavior is unusual for solids under pressure since applying external pressures generally decreases the distance between neighboring atoms and results in an increase in the carrier density as well as a suppression of the semiconducting energy gap. For a better understanding of the normal-state resistivity, we plotted representative ρ at 4.2 K ($\rho_{4.2\text{K}}$) vs pressure data for the two compounds in Fig. 5. It can be observed that $\rho_{4.2\text{K}}$ for both compounds is significantly suppressed when pressure is applied first, which can easily be understood due to the increase in the electron density under pressures. With a further increase in pressure, however, $\rho_{4.2\text{K}}$ of $\text{La}_{0.80}\text{Ti}_{0.20}\text{OBiS}_2$ increases slowly, whereas $\rho_{4.2\text{K}}$ of $\text{La}_{0.85}\text{Th}_{0.15}\text{OBiS}_2$ decreases very slightly, resulting in kinks in $\rho_{4.2\text{K}}(P)$ at ~ 0.6 GPa for both compounds, as indicated by the arrows in Fig. 5. Consistent with previous discussions about the behavior of superconductivity and a under pressure, a different phase whose normal-state resistivity exhibits an unusual response to pressure appears at pressures above ~ 0.6 GPa.

As mentioned in the Introduction, all of the BiS_2 -based superconductors which were reported to show pressure-induced low- T_c to high- T_c phase transitions contain F. The previous

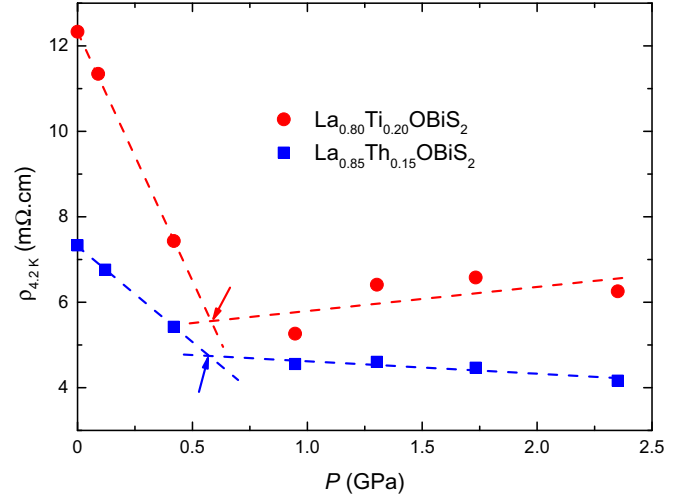


FIG. 5. Normal-state electrical resistivity ρ at 4.2 K for $\text{La}_{0.80}\text{Ti}_{0.20}\text{OBiS}_2$ (red circles) and $\text{La}_{0.85}\text{Th}_{0.15}\text{OBiS}_2$ (blue squares) under pressure. The dashed lines are linear fits of $\rho_{4.2\text{K}}$ vs P data, whose slopes change at ~ 0.6 GPa (indicated by the arrows) for both compounds.

study of Sm-substituted $\text{LaO}_{0.5}\text{F}_{0.5}\text{BiS}_2$ suggested that the pressure dependences of T_c in BiS_2 -based superconductors are closely related to the lattice parameter a rather than the chemical composition of the blocking layers [13]. In this paper, although phase transitions were clearly observed in $\text{La}_{0.80}\text{Ti}_{0.20}\text{OBiS}_2$ and $\text{La}_{0.85}\text{Th}_{0.15}\text{OBiS}_2$ via investigations of the behavior of both the superconducting transitions and the normal-state electrical resistivity under pressure, the high-pressure phases of the two compounds are apparently unfavorable for superconductivity. The results indicate that F in the blocking layer of BiS_2 -based compounds plays a critical role of not only providing conduction electrons, but also providing the sudden pressure-induced enhancement of superconductivity. This is presumably associated with a structural phase transition from tetragonal to orthorhombic symmetry as has been shown for $\text{LaO}_{0.5}\text{F}_{0.5}\text{BiS}_2$.

Although superconductivity is suppressed and ρ remains finite to 1.5 K, the decrease in ρ signaling the appearance of superconducting currents can still be observed up to the highest pressure obtained in this paper. For homogeneous samples, in general, pressure-induced phase transitions (either structural or electronic, first or second order) usually occur in a rather narrow hydrostatic pressure range. The abrupt changes in normal-state resistivity presented in this paper also suggest the phase transition should be complete before 1 GPa; this raises the question of why the superconducting transitions of the two compounds remain broad at pressures above 1 GPa. Considering the fact that bulk superconductivity of the samples at ambient pressure has been verified by means of specific heat and magnetization measurements [8], it is possible that the high-pressure phase is not superconducting and the superconducting currents observed above ~ 1 GPa are associated with small regions of the low-pressure phase that remains at high pressure due to slight inhomogeneities in the samples. However, the possibility that the high-pressure state of the compounds is still superconducting at lower tempera-

tures cannot be ruled out. Further studies need to be performed to determine whether the compounds are superconducting at lower temperatures.

IV. CONCLUDING REMARKS

To summarize, the observed phenomena including suppression and broadening of the superconducting transition, changes in the normal-state electrical resistivity with pressure, and the dramatic changes in the behavior of the normal-state electronic conduction between 0.42 and 0.95 GPa strongly suggest a pressure-induced phase transition occurs in the $\text{La}_{1-x}\text{M}_x\text{OBiS}_2$ compounds. In contrast to the pressure-induced low- T_c to high- T_c phase transitions reported in F-containing BiS_2 superconductors, the high-pressure phase observed in this paper does not favor superconductivity. The results suggest that fluorine plays a critical role for the pressure-induced enhancement of T_c for BiS_2 -based compounds. For the low-pressure phase, anomalies in the normal-state resistivity,

which seem to be pressure dependent, were observed in the pressure range from 0.24 to 0.45 GPa. Another unusual phenomenon observed in this paper is the tendency toward more semiconducting behavior with increasing pressure in the high-pressure phase at temperatures above ~ 140 K. Further studies of these unusual phenomena are expected to enrich our understanding of the physics of solids at high pressure.

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- [1] Y. Mizuguchi, H. Fujihisa, Y. Gotoh, K. Suzuki, H. Usui, K. Kuroki, S. Demura, Y. Takano, H. Izawa, and O. Miura, *Phys. Rev. B* **86**, 220510 (2012).
 - [2] D. Yazici, K. Huang, B. D. White, A. H. Chang, A. J. Friedman, and M. B. Maple, *Philos. Mag.* **93**, 673 (2013).
 - [3] Y. Fang, C. T. Wolowiec, D. Yazici, and M. B. Maple, *Nov. Supercond. Mater.* **1**, 79 (2015).
 - [4] D. Yazici, I. Jeon, B. D. White, and M. B. Maple, *Physica C* **514**, 218 (2015).
 - [5] Y. Fang, D. Yazici, B. D. White, and M. B. Maple, *Phys. Rev. B* **91**, 064510 (2015).
 - [6] I. Jeon, D. Yazici, B. D. White, A. J. Friedman, and M. B. Maple, *Phys. Rev. B* **90**, 054510 (2014).
 - [7] X. Lin, X. X. Ni, B. Chen, X. F. Xu, X. X. Yang, J. H. Dai, Y. K. Li, X. J. Yang, Y. K. Luo, Q. Tao, G. H. Cao, and Z. Xu, *Phys. Rev. B* **87**, 020504 (2013).
 - [8] D. Yazici, K. Huang, B. D. White, I. Jeon, V. W. Burnett, A. J. Friedman, I. K. Lum, M. Nallaiyan, S. Spagna, and M. B. Maple, *Phys. Rev. B* **87**, 174512 (2013).
 - [9] Y. Yu, J. F. Shao, S. Tan, C. J. Zhang, and Y. H. Zhang, *J. Phys. Soc. Jpn.* **82**, 034718 (2013).
 - [10] M. Tanaka, M. Nagao, Y. Matsushita, M. Fujioka, S. J. Denholme, T. Yamaguchi, H. Takeya, and Y. Takano, *J. Solid State Chem.* **219**, 168 (2014).
 - [11] H. F. Zhai, Z. T. Tang, H. Jiang, K. Xu, K. Zhang, P. Zhang, J. K. Bao, Y. L. Sun, W. H. Jiao, I. Nowik, I. Felner, Y. K. Li, X. F. Xu, Q. Tao, C. M. Feng, Z. A. Xu, and G. H. Cao, *Phys. Rev. B* **90**, 064518 (2014).
 - [12] H. F. Zhai, P. Zhang, S. Q. Wu, C. Y. He, Z. T. Tang, H. Jiang, Y. L. Sun, J. K. Bao, I. Nowik, I. Felner, Y. W. Zeng, Y. K. Li, X. F. Xu, Q. Tao, Z. A. Xu, and G. H. Cao, *J. Am. Chem. Soc.* **136**, 15386 (2014).
 - [13] Y. Fang, D. Yazici, B. D. White, and M. B. Maple, *Phys. Rev. B* **92**, 094507 (2015).
 - [14] C. T. Wolowiec, B. D. White, I. Jeon, D. Yazici, K. Huang, and M. B. Maple, *J. Phys.: Condens. Matter* **25**, 422201 (2013).
 - [15] Y. Fang, C. T. Wolowiec, A. J. Breindel, D. Yazici, P.-C. Ho, and M. B. Maple, *Supercond. Sci. Technol.* **30**, 115004 (2017).
 - [16] R. Jha, B. Tiwari, and V. P. S. Awana, *J. Phys. Soc. Jpn.* **83**, 063707 (2014).
 - [17] R. Jha, B. Tiwari, and V. P. S. Awana, *J. Appl. Phys.* **117**, 013901 (2015).
 - [18] M. Fujioka, M. Tanaka, S. J. Denholme, T. Yamaki, H. Takeya, T. Yamaguchi, and Y. Takano, *Europhys. Lett.* **108**, 47007 (2014).
 - [19] C. Y. Guo, Y. Chen, M. Smidman, S. A. Chen, W. B. Jiang, H. F. Zhai, Y. F. Wang, G. H. Cao, J. M. Chen, X. Lu, and H. Q. Yuan, *Phys. Rev. B* **91**, 214512 (2015).
 - [20] T. Tomita, M. Ebata, H. Soeda, H. Takahashi, H. Fujihisa, Y. Gotoh, Y. Mizuguchi, H. Izawa, O. Miura, S. Demura, K. Deguchi, and Y. Takano, *J. Phys. Soc. Jpn.* **83**, 063704 (2014).
 - [21] M. Tanaka, M. Nagao, R. Matsumoto, N. Kataoka, I. Ueta, H. Tanaka, S. Watauchi, I. Tanaka, and Y. Takano, *J. Alloys Compd.* **722**, 467 (2017).
 - [22] T. F. Smith, C. W. Chu, and M. B. Maple, *Cryogenics* **9**, 53 (1969).
 - [23] S. Demura, Y. Mizuguchi, K. Deguchi, H. Okazaki, H. Hara, T. Watanabe, S. J. Denholme, M. Fujioka, T. Ozaki, H. Fujihisa, G. Yoshito, M. Osuke, Y. Takahide, T. Hiroyuki, and T. Yoshihiko, *J. Phys. Soc. Jpn.* **82**, 033708 (2013).
 - [24] G. C. Kim, M. Cheon, Y. C. Kim, and R.-K. Ko, *arXiv:1512.05070*.
 - [25] S. Wu, Z. Sun, F. Chiang, C. Ma, H. Tian, R. Zhang, B. Zhang, J. Li, and H. Yang, *Solid State Commun.* **205**, 14 (2015).