Possible overdoping in TlSrLaCuO₅ and superconductivity in TlBaLaCuO₅ by isovalent-cation substitution

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It is known that superconductivity in TlSrLaCuO₅ is due to the overlap of empty Tl 6s band with the antibonding Cu $3d_{x^2-y^2}$ band at the Fermi level. This motivated us to investigate the effects of isovalent-cation substitution in superconducting TlSrLaCuO₅ and semiconducting TlBaLaCuO₅. We have succeeded in inducing superconductivity in TlBaLaCuO₅ by introducing Ca²⁺ at the Ba sites.

I. INTRODUCTION

in single Tl-O layered compounds $T1A_2Ca_{n-1}Cu_nO_{2n+3}$ (where A = Ba or Sr and n = 1-5) Tl-O double layered compounds $Tl_2Ba_2Ca_{n-1}Cu_nO_{2n+4}$ (where n=1-4) are created via two different mechanisms. 1,2 In the former class, itinerant charge carriers are created either by introducing higher-valence cations²⁻⁴ in place of lower-valence cations or by creating anion nonstoichiometry.⁵ In the latter class, holes are largely created by the preexisting redox mechanism, 6,7 i.e., due to the overlap of the empty Tl 6s band with the antibonding Cu $3d_{x^2-y^2}$ band [see Fig. 1(a)], electrons are transferred from CuO₂ planes to Tl-O layers $Tl^{3+} + Cu^{2+} \rightleftharpoons Tl^{(3-\delta)+}$ and vice versa: $+Cu^{(2+\delta)+}$. Recently, it has been proposed that the redox mechanism is responsible for hole generation in a derivative of the single Tl-O layered compound with n=1, TlSrLaCuO₅ ($T_c=32$ K). This conclusion follows from tight-binding electronic band-structure calculations⁸ and has been established experimentally by x-ray-photoelectron spectroscopic studies on TlBa_{1-x}Sr_xLaCuO₅.9 It should be pointed out that the isostructural compound TlBaLaCuO₅ is not a superconductor, in fact it is a semiconductor.³ In this case, the Fermi level lies below the bottom of the Tl 6s band, as shown in Fig. 1(b), rendering the redox mechanism nonoperative.10

In order to make the redox mechanism in TlBaLaCuO₅ operative, the lower Cu $3d_{x^2-y^2}$ Hubbard band must be raised (in energy) so as to overlap with the empty Tl 6s band. This is possible if one is able to decrease the inplane Cu-O bond length. As the Cu-O bond length decreases, the Fermi level rises and moves toward the empty Tl 6s band. Eventually, the Cu $3d_{x^2-y^2}$ band overlaps

with the Tl 6s band [see Fig. 1(c)]. This causes electrons to flow from Cu ions to Tl ions and holes are introduced in the Cu band.

Experimentally, a decrease of the Cu-O bond length may be accomplished by substituting the bigger Ba^{2+} ions in $TlBaLaCuO_5$ with smaller isovalent cations, 9,11 for example, Sr^{2+} . Both the lattice parameters a and c of the tetragonal structure of the system $TlBa_{1-x}Sr_xLaCuO_5$ decrease as a function of x. The decrease in a implies a decrease of the in-plane Cu-O bond length because it has been observed that in all thallium-based high- T_c superconductors the in-plane Cu-O bond length is one-half of a. The concentration of holes may be increased by increasing x in $TlBa_{1-x}Sr_xLaCuO_5$. As a

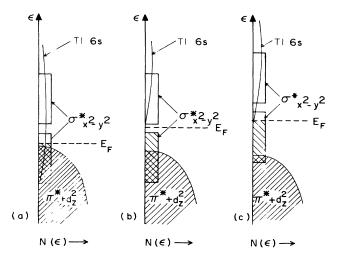


FIG. 1. Energy density of states $N(\epsilon)$ vs energy ϵ for nominal (a) Tl₂Ba₂CuO₆, (b) TlBaLaCuO₅, and (c) TlSrLaCuO₅. [(a) and (c) adapted from Ref. 10].

result, T_c also increases. Interestingly, the maximum T_c is observed in TlSrLaCuO₅ (32-46 K depending upon the preparative conditions^{9,12,13}). One may expect that the hole concentration in TlSrLaCuO₅ may increase (leading to overdoping of holes) if the degree of overlap is further increased by substituting Sr2+ ions by still smaller isovalent cations (Ca²⁺). This should induce a variation of T_c (T_c should further increase, or if the compound is already optimally hole doped T_c should decrease due to overdoping) in the system TlSr_{1-v}Ca_vLaCuO₅ in a manner similar to that in other high- T_c superconductors, such as La_{2-x}Sr_xCuO₄. ¹⁴ These considerations led us to examine the effect of hole doping in TlBaLaCuO₅ by introducing the smaller Ca²⁺ ions. We have, in fact, succeeded in inducing superconductivity in this material following this approach. In this paper, we report the results of these investigations.

II. EXPERIMENTAL

Samples of the type $TlSr_{1-y}Ca_yLaCuO_5$ ($0 \le y \le 1.0$) and $TlBa_{1-z}Ca_zLaCuO_5$ ($0 \le z \le 1.0$) were prepared by the method followed in our preparation of the materials $TlBa_{1-x}Sr_xLaCuO_5$. The phase purity of the samples was checked by x-ray powder diffraction. The superconducting properties were examined by the ac mutual inductance (frequency = 313 Hz) technique and four-probe resistance measurements. In order to check reproducibility of the results, three independent batches of the above-mentioned samples were made and investigated.

III. RESULTS

X-ray powder diffraction (XRD) showed that essentialsingle-phase samples of the composition $TlSr_{1-y}Ca_yLaCuO_5$ are formed for $0.0 \le y \le 0.6$ and of the composition $TlBa_{1-z}Ca_zLaCuO_5$ for $0.0 \le z \le 0.8$. The powder-diffraction pattern of TlBa₁₋₂Ca₂LaCuO₅ (z = 0.0 - 0.8) is shown in Fig. 2. Traces of some impurity phases in materials of both systems were seen. Diffraction lines could be indexed on the basis of the 1:2:0:1 phase, with the space group P4/mmm. Lattice parameters a and c, obtained using a least-squares fitting procedure, are given in Table I. ac susceptibility meaof two samples of the system surements TlSr_{1-v}Ca_vLaCuO₅, namely, those corresponding to y = 0.3 and 0.6, showed that the superconducting onset temperature, respectively, is 32 and 26 K. Figure 3 shows the normalized resistance of these samples as a function of temperature. Zero-resistance temperatures T_{c0} for the two samples are 30 and 20 K, respectively.

Superconductivity was observed in the system $TlBa_{1-z}Ca_zLaCuO_5$ with $12 < T_c < 16$ K for $0.3 \le z \le 0.8$. Figure 4 shows the resistance of the material $TlBa_{0.4}Ca_{0.6}LaCuO_5$ as a function of temperature; $T_{c0}=13$ K. The inset of Fig. 4 shows the diamagnetic response of this material as a function of temperature, with onset temperature ~ 16 K. In view of the fact that the effect of Ca doping in $TlBaLaCuO_5$ is of particular interest, we synthesized two independent samples of

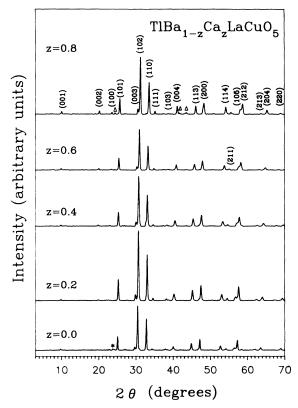


FIG. 2. X-ray powder-diffraction pattern of $TlBa_{1-z}Ca_zLaCuO_5$ (z=0.0, 0.2, 0.4, 0.6, and 0.8). The lines marked (*) and (\triangle) indicate the presence of traces of $BaCO_3$ (only in the higher-Ba-content samples) and La_2CuO_4 (only in the higher-Ca-content sample), respectively. All the other lines can be indexed in terms of the 1:2:0:1 phase (space group P4/mmm.)

TlBa_{0.4}Ca_{0.6}LaCuO₅ and both the samples gave similar results. The strength of the ac χ signal at the lowest temperature of our measurements (4.2 K) suggests clearly the TlBa_{0.4}Ca_{0.6}LaCuO₅ is a bulk superconductor (the superconducting volume fraction is estimated to be in excess of 60%).

TABLE I. Compounds and their corresponding lattice parameters a and c, tolerance factor t, and diamagnetic onset temperature.

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Compound	a (Å)	c (Å)	t	T_c (K)
TlBaLaCuO ₅	3.852	8.9928	0.860	
TlBa _{0.6} Sr _{0.4} LaCuO ₅	3.821	8.901	0.852	20
TlBa _{0.3} Sr _{0.7} LaCuO ₅	3.794	8.885	0.848	30
TlSrLaCuO ₅	3.779	8.864	0.842	32
TlSr _{0.7} Ca _{0.3} LaCuO ₅	3.761	8.807	0.838	32
TlSr _{0.4} Ca _{0.6} LaCuO ₅	3.753	8.737	0.834	26
TlBa _{0.8} Ca _{0.2} LaCuO ₅	3.827	8.948	0.854	
TlBa _{0.6} Ca _{0.4} LaCuO ₅	3.817	8.885	0.847	14
TlBa _{0.4} Ca _{0.6} LaCuO ₅	3.790	8.815	0.841	16
TlBa _{0.2} Ca _{0.8} LaCuO ₅	3.766	8.750	0.834	13

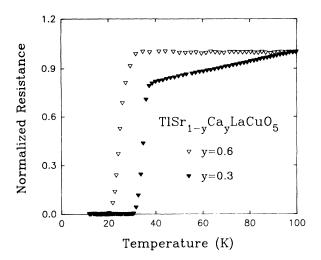


FIG. 3. Normalized resistance as a function of temperature in $TlSr_{1-y}Ca_yLaCuO_5$ (x=0.3,0.6).

IV. DISCUSSION

A. Structural stability and role of thallium

Careful analysis of x-ray powder-diffraction patterns revealed the presence of a trace of La_2CuO_4 as impurity phase in the materials $TlSr_{1-y}Ca_yLaCuO_5$ and $TlBa_{1-z}Ca_zLaCuO_5$. The intensity of the impurity lines increases with increasing y in $TlSr_{1-y}Ca_yLaCuO_5$. In the system $TlBa_{1-z}Ca_zLaCuO_5$ the impurity phase appears with z=0.8, as shown in Fig. 2. The stability of the tetragonal 1:2:0:1 phase of these systems could be understood in terms of a geometrical factor, the so-called tolerance factor, as discussed below.

All the known copper oxide superconductors are built up of oxygen-deficient perovskite layers intergrown with

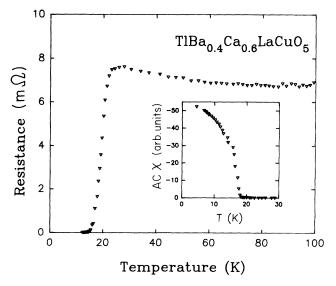


FIG. 4. Resistance as a function of temperature of TlBa_{0.4}Ca_{0.6}LaCuO₅. The inset shows the diamagnetic response, as a function of temperature; the estimated superconducting volume fraction is in excess of 60%.

rocksalt or fluorite-type layers. In general, perovskite layers are made up of copper and oxygen, and, particularly in thallium compounds, the rocksalt or fluorite-type layers are Tl, Ba or Sr, Ca or rare earths, and oxygen. Stabilization of the intergrowth structure requires bondlength matching across the intergrowth interface. A measure of bond-length matching across the interface is the tolerance factor, 15 $t=r_{A-O}/\sqrt{2}r_{\text{Cu-O}}$, where A represents the ions in the rocksalt or fluorite-type layers and r_{A-O} is the average length of the A-O bonds in the structure. An ideal matching occurs when t=1. t<1 means that the CuO₂ planes are under compression and the A-O layers are under tension. 16 The empirically observed limits for the stability of the intergrowth structure are $0.86 \le t \le 1.02.17$

The tolerance factor for $TlBa_{1-x}Sr_xLaCuO_5$ $(0 \le x \le 1)$, $TlSr_{1-y}Ca_yLaCuO_5$ $(0 \le y \le 0.6)$, and $TlBa_{1-z}Ca_zLaCuO_5$ $(0 \le z \le 0.8)$ has been calculated according to the above relation using the average radii of the A ions, taken from Ref. 18, for the appropriate coordination number. Values of t for various materials are shown in Table I.

The tolerance factor for the semiconducting TlBaLaCuO₅ is 0.860, which lies right at the lower limit for the stability of the intergrowth structure. Upon substituting Sr at the Ba site, and considering that Tl remains in the 3+ state (so that its effective radius does not vary as a function of doping), the tolerance factor decreases below the lower limit (see Table I), due to the smaller ionic radius of Sr. which means that the structure should not form. The fact that the structure forms (and therefore the tolerance factor must be within the stability limits), and the phase is stabilized, however, suggests that the valence of the Tl ions does change; it moves toward 1 +and the ions expand in size (radius=0.885 Å for T1³⁺ and 1.50 Å for T1⁺ for six coordination). These studies, therefore, clearly throw light on the dual role of thallium in these materials, i.e., it stabilizes the structure by varying its size and allows the generation of holes, via the redox mechanism, by changing its valence. Of course, at the same time, the decrease in $r_{\text{Cu-O}}$ (see below) due to the creation of holes in the CuO2 planes also contributes to the increase in the effective value of the tolerance factor.

The solubility of Ca is limited by the tolerance factor t=0.834, which is the limiting value for Ca in $TlSr_{1-y}Ca_yLaCuO_5$ and $TlBa_{1-z}Ca_zLaCuO_5$. Beyond the value of t=0.834 the impurity phase dominates. Large concentrations of such ions correspond to a large fraction of monovalent Tl, which the lattice finds too difficult to accommodate because of its bigger ionic radius (1.50 Å).

B. Overdoping in $TlSr_{1-\nu}Ca_{\nu}LaCuO_5$

The lattice parameters a and c decrease with increasing y (see Table I). Decrease of c with increase of y implies that the Ca^{2+} ions substitute essentially for Sr^{2+} ions. Decrease of lattice parameter a with increase of y shows that the in-plane Cu-O bond length r_{Cu-O} decreases, which should increase the overlap of the Tl 6s and Cu $3d_{x^2-y^2}$ bands and thereby increase the hole concentra-

tion. The decrease of T_c for y=0.6 reveals that the material is in the overdoped state. The variation of T_c with hole concentration in the systems $\text{TlBa}_{1-x}\text{Sr}_x\text{LaCuO}_5$ and $\text{TlSr}_{1-y}\text{Ca}_y\text{LaCuO}_5$, as a whole, is similar to what has been observed in the other high- T_c superconductors.

C. Superconductivity in TlBa_{1-z}Ca_zLaCuO₅

In order to check the validity of the creation of the redox mechanism and thus the semiconductor-metal transition in TlBaLaCuO₅ by substituting smaller alkaline rare earths (Sr²⁺) at the Ba site and further Ca²⁺ at the Sr site in TlSrLaCuO₅, we have substituted Ca²⁺ at the Ba site in TlBaLaCuO₅. Decrease of the lattice parameters a and c with increase of z in TlBa_{1-z}Ca_zLaCuO₅ suggests that the in-plane Cu-O bond length decreases, which should favor the overlap of Tl 6s and Cu $3d_{x^2-y^2}$ bands at the Fermi level, and Ca essentially substitutes at the Ba site. The observation of superconductivity in TlBa_{1-z}Ca_zLaCuO₅ 0.3 < z < 0.8 proves the existence or creation of the redox mechanism when the lattice parameter a is small or made smaller. It is important to note that the compound TlCaLaCuO₅ does not form, presumably again in accordance with the consideration of the tolerance factor. It is instructive to note that our attempt to prepare single-phase TlBaLaCuO₅, substituted with smaller rare earths at the bigger La site, which is also expected to decrease the in-plane Cu-O bond length and hence favor overlap of Tl 6s and Cu $3d_{x^2-y^2}$ bands, was not successful. This may be due to the increase of the mismatch between Ba²⁺ and the smaller rare earths.

V. CONCLUSIONS

We have shown here that, using the concept of the redox mechanism, the hole concentration can be increased in TlSrLaCuO₅ by substituting a smaller isovalent cation (Ca²⁺) at the bigger Sr site, leading to an overdoped region. Based on this model, we have also been able to induce superconductivity in the semiconducting TlBaLaCuO₅ by substituting a smaller isovalent cation (Ca²⁺) at the bigger Ba site. Structural-stability considerations give useful information with respect to the solubility limit of the substituent and the valence state of thallium.

ACKNOWLEDGMENTS

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