

# One step reconstruction of wave speed and optical properties for PAT

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# Photoacoustic Tomography(PAT)

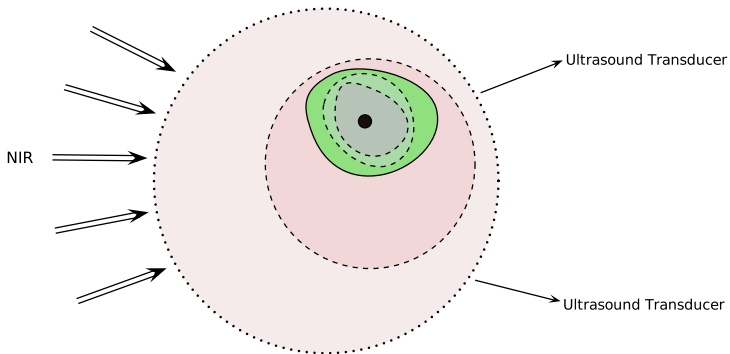
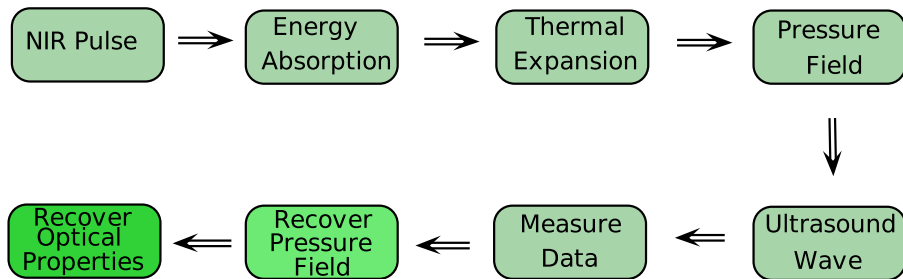


Figure : Photoacoustic Tomography: To recover **scattering**, **absorption** and **photoacoustic efficiency** properties of tissues from boundary measurement of acoustic signal generated with the photoacoustic effect. Two processes: **propagation of NIR radiation** and **propagation of ultrasound**. There is a (time) **scale separation** between the two processes.

# Regular PAT Workflow



**Figure :** Workflow chart of both physical process during PAT and regular inversion process

# Photon Diffusion Process

Let  $\Omega \subset \mathbb{R}^d$  be the domain of medium. Then the equation of photon density  $u$  at position  $\mathbf{x} \in \Omega$  satisfies

$$-\nabla \cdot D \nabla u(\mathbf{x}) + \sigma u(\mathbf{x}) = 0, \quad (1)$$

with boundary condition

$$\psi|_{\partial\Omega} = g(\mathbf{x}), \quad (2)$$

where  $D(\mathbf{x})$  is diffusion coefficient,  $\sigma(\mathbf{x})$  is absorption coefficient.

# Photoacoustic Effect

The medium absorbs part of the energy of NIR photons, generate initial pressure field  $H(\mathbf{x})$  through photoacoustic effect.

$$H(x) = \Gamma(\mathbf{x})\sigma(\mathbf{x})u(\mathbf{x}), \quad (3)$$

where  $\Gamma(x)$  is Grüneisen coefficient, measures the efficiency of photoacoustic effect(energy  $\rightarrow$  pressure).

# Acoustic Wave Propagation

The initial pressure field generates acoustic wave(ultrasound),

$$\begin{aligned}\frac{1}{c^2(\mathbf{x})}p_{tt} - \Delta p &= 0, \\ p(0, \mathbf{x}) &= H(x) = \Gamma(\mathbf{x})\sigma_a(\mathbf{x})u(\mathbf{x}), \\ p_t(0, \mathbf{x}) &= 0,\end{aligned}\tag{4}$$

during the photoacoustic process, wave speed  $c(\mathbf{x})$  is assumed to be unchanged. And we measure the pressure field(ultrasound signal) on the surface  $\Sigma = \partial\Omega$  of domain for sufficient long time  $T \gg 1$ ,

$$\mathcal{M}(t, \mathbf{x}) = p(t, \mathbf{x})|_{[0, T] \times \Sigma}\tag{5}$$

# Two Step Reconstruction

- 1 The first step is solve an inverse source problem with (4) and (5). In practice,  $c(\mathbf{x})$  is variable, and time reversal method works well when  $c(\mathbf{x})$  is known.
- 2 The second step is nonlinear and there are plenty of literatures, it is known that we can reconstruct any two of  $D(\mathbf{x})$ ,  $\sigma(\mathbf{x})$  and  $\Gamma(\mathbf{x})$  if the third is known. e.g., if  $D$  is given, we have a chain like

$$\mathcal{M}(t, \mathbf{x}) \rightarrow H \rightarrow (\Gamma, \sigma) \quad (6)$$

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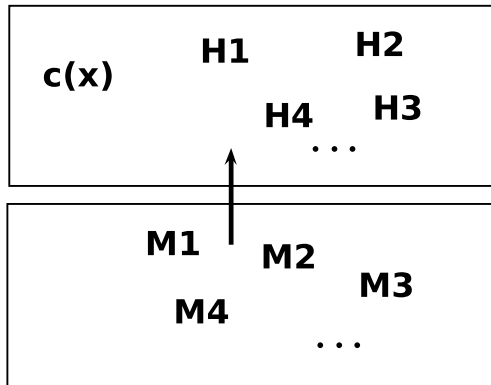
$$\mathcal{M}(t, \mathbf{x}) \rightarrow H \rightarrow (\Gamma, \sigma) \quad (7)$$



# Two Step Reconstruction

It is proved that when  $c(\mathbf{x})$  is unknown, the reconstruction in first step will be very unstable.

Why? Intuitively, we are using less data to recover more unknowns.



# Two Step Reconstruction

No matter how many measurements are taken, the first step always requires recovery of the new  $H$ , does not reduce the number of unknowns at all. The fact is:

$H$  is not important.

# One Step Reconstruction

We combine the two process as one, try to use the measurement to directly recover the properties of interest. Since we introduced a new unknown  $c(\mathbf{x})$ , we need to know one more in  $(D, \sigma, \Gamma)$ , e.g.,  $D$  and  $\Gamma$  are known. And the problem turns to be a simple chain,

$$\mathcal{M}(t, \mathbf{x}) \rightarrow (c, \sigma) \quad (8)$$

# One Step Reconstruction

Current main results on uniqueness and stability:

- 1 When  $c(\mathbf{x})$  is constant and unknown, we can uniquely recover  $c(x)$ .
- 2 When  $c(\mathbf{x}) = \alpha f(\mathbf{x})$  where  $\alpha$  is unknown real number and  $f$  is known, we can uniquely recover  $c(\mathbf{x})$ .
- 3 There is no Sobolev stability estimate. (no stability involving finite order of derivatives).

In other words, it seems if  $c(\mathbf{x})$  only depends on finite (one ?) variables, there is no problem on uniqueness.

# One Step Reconstruction

Our results on uniqueness:

- 1 In 1D, **use one measurement**,  $c(\mathbf{x})$  can be recovered uniquely.
- 2 In 3D, **use one measurement**, if  $c(\mathbf{x})$  is radial (only depends on  $|\mathbf{x}|$ ), then  $c(\mathbf{x})$  can be determined uniquely.
- 3 In 3D, **use infinite measurements**,  $c(\mathbf{x})$  can be determined uniquely.

In other words, if  $c(\mathbf{x})$  is one dimensional, single measurement is enough.

# The Idea

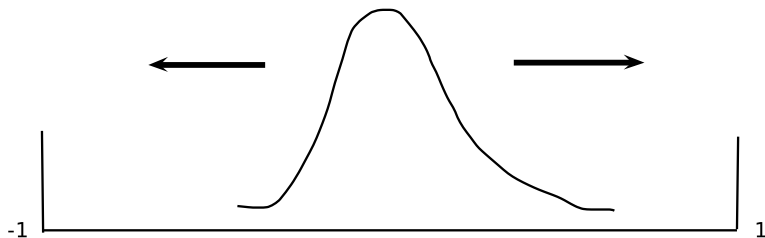
Use linearized model, say  $c(x) = c_0 + \delta c$ ,  $\sigma = \sigma_0 + \delta \sigma$ . Then the perturbed wave equation is

$$\begin{aligned}\partial_{tt}\delta p - c_0^2\delta p &= \frac{2\delta c}{c_0}\partial_{tt}p_0 \\ \delta p(x, 0) &= \delta H \\ \partial_t\delta p(x, 0) &= 0\end{aligned}\tag{9}$$

perturbed measurement is  $\delta\mathcal{M} = \delta p(\mathbb{R}^+ \times \Sigma)$ .

# The Idea

For example, in 1D.



**Figure :** Compare the measurements on both sides from different direction of time (forward and backward).

# The Idea

We somehow can cancel the term with perturbation from  $\sigma(x)$  (or  $\delta H$ ). And the remaining will only depend on perturbed  $\delta c(x)$ . From here to show uniqueness requires some tricks.

More important thing: boundary measurement is enough to determine both wave speed and initial condition.



# Conclusion

- 1 Made some progress on PAT.
- 2 Some uniqueness and stability results, very likely to be logarithm stability.
- 3 In 2D, our method does not work.