# Statistical Studies: Experiment

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Two energy peaks of Cs-137 decay where observed with a multi-channel analyzer and NaI detector, a Gamma-ray emission line with energy  $679.1^{+2.7}_{-2.7} \text{KeV}$  and a X-ray emission line with energy  $14.88^{0.15}_{-0.15} \text{KeV}$  was observed. The resolution of the Multi-channel analyzer was determened to be  $6.943 \pm 0.039$ .

#### I. INTRODUCTION

The firs component of this experiment focus on the statistical atool sand tequiques for analyzing adata and statistical distribution. Spesifically methods for fitting statistical distributions and tecniques for determining if data is consistent with a theoretical model are investigated. In the second component of theis experiment the statistical method investigated in the first section are applied to the practical problem of calibrating labritory equiment, and arror analysis of the subsequent mesurements. In the corse of this lab we will be investigating the Poisson distribution, the Gaussian distributheon the  $\chi^2$  distribution and the  $\chi^2$  statistic and theoric application to the analysis of counting statistics and the callibration of a multi-chanel analyzer.

#### II. THEORY

The statistic of counting random events is modeled by assuming the process in question can be described as a large number of simple binary events which each have a small probability of leading to a count being mesured, where the number of binary event n is  $n \sim 1/x$  each event having a probability of the binary even occuring is  $p \sim 1/x$  for x << 1. The probability that  $\nu$  events will occur in a given interval is described by the binomial distribution,

$$B_{n,p}(\nu) = \binom{n}{\nu} p^{\nu} (1-p)^{n-\nu}$$

In the limiting case of a continuous process p->0 and  $n->\infty$  the distribution simplifies to the Poisson distribution.

$$P_{\mu}(\nu) = \frac{e^{-\mu}\mu^{\nu}}{\nu!} \tag{1}$$

where the expected number of events  $\bar{\nu}$  is  $\bar{\nu} = \sum_{\nu=0}^{\infty} \nu P_{\mu}(\nu) = \mu$  and the varience is  $\sigma_{\nu}^2 = \sum_{\nu=0}^{\infty} (\nu - \mu)^2 P_{\mu}(\nu) = \mu$ . The Poisson distribution can be simplified in the limiting case where  $\mu$  is large, in that case the

Poisson distribution (1) a proches a gaussian distribution with mean  $\mu$  and varience  $\mu$ . In this limit the distribution becomes,

$$P_{\mu}(\nu) \approx G_{\mu,\sqrt{\mu}}(\nu) = \frac{1}{\sqrt{2\pi\mu}} e^{-(\nu-\mu)^2/2\mu}$$
 (2)

A. 
$$\chi^2$$
 Test

The  $\chi^2$  test test is a method for determining the probability that a set of mesurements is consistent with a given model assuming Gaussian errors. The  $\chi^2$  test uses the  $\chi^2$  statistic which in the case of descreat variables is defined as,

$$\chi^2 = \sum_{i} \left( \frac{O_i - E_i}{\sigma_i} \right)^2$$

where  $O_i$  is the observed value,  $E_i$  is the expected value and  $\sigma_i$  is the standard deviation. Given the expected value is distributed as a Poisson distribution with mean  $\mu = E_i$  and variance  $\sigma_i = \sqrt{E_i}$  then the  $\chi^2$  statistic is,

$$\chi^{2} = \sum_{i} \frac{(O_{i} - E_{i})^{2}}{E_{i}} \tag{3}$$

Assuming the errors are gaussaing then the  $\chi^2$  statistic is described by the associated  $\chi^2$  distribution. The  $\chi^2$  distribution which has the cumulative distribution function (CDF),

$$\operatorname{Prob}_{d}(\chi^{2} \leq \chi_{o}^{2}) = \frac{1}{2^{d/2}\Gamma(d/2)} \int_{0}^{\chi_{o}^{2}} t^{d/2 - 1} e^{-t/2} dt \quad (4)$$

So the probability of measureing a  $\chi^2$  as large or larger than  $\chi_o^2$  is  $\operatorname{Prob}_d(\chi^2 \geq \chi_o^2) = 1 - \operatorname{Prob}_d(\chi^2 \leq \chi_o^2)$ .

## B. Poisson errors and modified $\chi^2$ statistic

The statistic given by equation (3) can be used with the  $\chi^2$  assuming  $O_i$  is described by a gausian distribution but that assumption is invalid when  $E_i$  is not sufficently large that the Poisson distribution can be approximated

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as a Gaussian distribution. If when  $E_i$  is small the  $\chi^2$  statistic must be modified to account for the asymmetry of the Poisson distribution. For the guassian distribution the probability of getting a value below the lower bound of a one sigma confidence interval is  $p = \frac{1}{2} - \text{erf}(\frac{1}{\sqrt{2}})/2$ , for the modified  $\chi^2$  statistic insead of the standard deviation lets use the value  $|\nu_o - \mu|$  where  $\nu_o$  solves the equation  $\text{Prob}_{\mu} (\nu \leq \nu_o) = p$  or  $\text{Prob}_{\mu} (\nu \geq \nu_o) = 1 - p$  depending on whether the observed value is above or below the expected value. To evaluate these probabilities for  $\mu$ . The CDF of the poison distribution is,

$$\operatorname{Prob}_{\mu}\left(\nu \leq \nu_{o}\right) = \sum_{k=0}^{\nu_{o}} \frac{e^{-\mu} \mu^{\nu}}{\nu!}$$

where  $\nu_o$  is a positive integer. When  $\mu \sim 1$  a continuous analouge to the Poisson CDF is needed to evaluate  $\nu_o$  for the adjusted  $\chi^2$  statistic. In the paper Continuous Counterparts of Poisson and Binomial Distributions and their Properites[1] A. Ilienko defines a continuous analogue to the Poisson distribution, the CDF of this distribution is

$$\operatorname{Prob}_{\mu} (\nu \leq \nu_o) = \frac{\Gamma(\nu_o + 1, \mu)}{\Gamma(\nu_o + 1)}$$

where  $\nu_o$  is a positive real number,  $\Gamma(x,\lambda)$  is the incomplete Gamma function  $\Gamma(x,\lambda)=\int_{\lambda}^{\infty}t^{x-1}e^{-t}dt$  and  $\Gamma(x)=\Gamma(x,0)$  is the Gamma function. Note that if  $\nu_o$  is a positive integer then,

$$\operatorname{Prob}_{\mu}(\nu \leq \nu_{o}) = \frac{\Gamma(\nu_{o} + 1, \mu)}{\Gamma(\nu_{o} + 1)} = \sum_{k=0}^{\nu_{o}} \frac{e^{-\mu} \mu^{\nu}}{\nu!}$$

Using the continuous analogue to the Poisson distribution the adjusted  $\chi^2$  statistic is defined as,

$$\chi^2 = \sum_{i} \left( \frac{O_i - E_i}{\nu_{o_i} - E_i} \right)^2$$

$$\frac{\Gamma(\nu_{o_i} + 1, E_i)}{\Gamma(\nu_{o_i} + 1)} = \begin{cases} p_{\text{mid}} - p_{\text{int}}/2 & \text{if } O_i < E_i \\ p_{\text{mid}} + p_{\text{int}}/2 & \text{if } O_i > E_i \end{cases}$$

where  $p_{\rm mid} = \Gamma(E_i+1,E_i)/\Gamma(E_i+1)$  and  $p_{\rm int} = {\rm erf}(\frac{1}{\sqrt{2}})$ . It should be noted that for  $E_i > 30$  the relitive difference between  $|\nu_{o_i}-E_i|$  and  $\sqrt{E_i}$  is less than ten percent and that if  $p_{\rm mid} < p_{\rm int}/2$  the interval must be set to,

$$\frac{\Gamma(\nu_{o_i}+1, E_i)}{\Gamma(\nu_{o_i}+1)} = \begin{cases} 0 & \text{if } O_i < E_i \\ p_{\text{int}} & \text{if } O_i > E_i \end{cases}$$

### C. Energy resolution

The resolution R of a spectromter is,

$$R = \frac{\delta E}{E} \cdot 100 \tag{5}$$

where  $\delta E$  is the Full Width at Half Maximum (FWHM) of an energy peak and where E is the energy at the maximum of the energy peak. The resolution R gauges the detectors ability to resolve spectral features.

### III. DESIGN AND RESULTS

The experiment is split into two components, in the first section the distribution of counts produced by a sintilator(s) is measured and popaired to theoretical distributions, in the second section the methodes investigated for analyzing distributions is used to callibrate a multichannel analyzer and to measure the resolution of the device.

### A. Section I

The distribution produced by one sintilator for the casses of  $\mu \approx 5$  and  $\mu \approx 100$  where measured. We setup the sintilator with a -1.5kV bias voltage. The signal from the sintilarot was passed into an amplifier and then through a discriminarot. the signal from te disescriminator was then passed to a timer/counter. Two sets of counts were colected with this setup each with 500 measurements. In the first set the descriminator cutoff was tuned to produce a low count rate  $\mu \approx 5$ , in the second set the cutoff was tuned to produce a high count rate of  $\mu \approx 100$ .

In the second part two sintilatorst where used to reject any signal that was not corrillated between the devieces, so as to reduce the random noise from the sintelators. The sintelators where stacked verticaly with approximatly 6cm vertical seporation. The signals from the two sintelators where then passed into two amplifieers and then into two descriminators. The two signels where then passed into a corrillator with corrillation setting set to two. The output from the corrillator was then connected to a counter/timer. As in the first part two sets of data was taken with 500 measurements each, one with  $\mu \approx 5$  and the other with  $\mu \approx 100$ .

### B. Section II

using a NaI detector we callibrated a multi-channel analyzer by fitting the energy peaks of Na-22 and Co-60 using the analyssis tecniques investigated in Section I. Once calibrated a gamaray and xray peak of Cs-137 was

TABLE I. Channel number and Energy of peaks

|        | $\mu$ in (channel)          | dof | $\chi^2$ | p-value | energy in (MeV) |
|--------|-----------------------------|-----|----------|---------|-----------------|
|        | $337.159^{+0.016}_{-0.016}$ | 55  | 200.92   | 0.0%    | 0.511           |
|        | 0.071                       | 113 | 121.90   | 26.8%   | 1.274           |
|        | $755.289^{+0.041}_{-0.040}$ | 82  | 97.97    | 11.0%   | 1.17            |
|        | $856.509^{+0.041}_{-0.041}$ | 100 | 104.25   | 36.6%   | 1.33            |
|        | $443.914^{+0.017}_{-0.016}$ | 95  | 550.09   | 0.0%    | 0.6612          |
| Cs-137 | $22.130^{+0.012}_{-0.012}$  | 15  | 2024.62  | 0.0%    | 0.0322          |

TABLE II. Energy of Cs-137 lines

|                            |                                     | expected energy      | percent error |
|----------------------------|-------------------------------------|----------------------|---------------|
|                            | -2.7                                | $661.2 \mathrm{KeV}$ | +2.71%        |
| $22.130^{+0.012}_{-0.012}$ | $14.88^{0.15}_{-0.15} \mathrm{KeV}$ | $32.2 \mathrm{KeV}$  | -53.8%        |

measured and used to find the energy resolution of the multi-channel analyzer.

#### IV. ANALYSIS

#### A. section I

### B. section II

When fitting the energy peaks, it was assumed that peak itself was Gaussian and that the local region around the peak could be described by a straight line. The model of the energy peak used in the fitting was,

$$y = Ae^{-(x-\mu)^2/c} + m(x-\mu) + b$$

The Gamma-ray line energies are from *Scintillation* Spectrometry[2], the energy of the Ba  $K\alpha$  line in the Cs-137 spectrum is from X-ray wavelengths[3].

Using the energy peaks from the Na-22 and Co-60 samples the Multi-channel analyzer was calibrated using linear regression. The calibration equation was found to be,

$$E = (1.5748 \pm 0.0061 \text{KeV/channel}) x - 19.971 \pm 0.066 \text{KeV}$$

Using the calibration equation the energies of the Cs-137 where found to be,

The using the 661.2 KeV line of Cs-137 the parameter  $c = 323.33^{+0.51}_{-0.51}$  in channels<sup>2</sup> the resolution of the Multichannel analyzer was determined to be  $6.943 \pm 0.039$ .

### V. CONCLUTION

Using a multi-channel analyzer to measure the gamma-ray emission of Cs-137 two lines where observed. A Gamma-ray emission line with energy 679.1 $^{+2.7}_{-2.7}$ KeV was observed and found to have a +2.71% deviation from the expected energy of 661.2KeV. A X-ray emission line with energy 14.88 $^{0.15}_{-0.15}$ KeV was observed and found to have a -53.8% deviation from the expected energy of 32.2KeV. The resolution of the Multi-channel analyzer was found to be 6.943  $\pm$  0.039.

<sup>[1]</sup> A. Ilienko, Continuous Counterparts of Poisson and Binomial Distributions and their Properties, (Annales Univ. Sci. Budapest., Sect. Comp. 39(2013) 137-147).

<sup>[2]</sup> R. L. Heath, Scintillation Spectrometry: Gamma-ray Spectrum Catalogue, (2nd Edition, Vol. 1, August 1964).

<sup>[3]</sup> J. A. Bearden, Xray wavelengths, (Review of Modern Physics, January 1967, p 86-99).