

MRI Simulator Notes

Liam

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Tasks

- ▶ I looked at the simulation method for excitation and for free precession using rotations and scalings.
- ▶ I wrote a proposal for the basic structure of our simulation.

Simulating excitation

Bloch equation in rotating frame

Let $R_z(\theta)$ be the rotation matrix about the z-axis by the angle θ . Assume the magnetic field during excitation is given by

$$\mathbf{B}(z, t) = R_z(\omega t) \underbrace{\begin{bmatrix} B_{1,x}(t) \\ B_{1,y}(t) \\ B_0 + G_z(t)z \end{bmatrix}}_{\mathbf{B}_{\text{rot}}(z, t)},$$

which corresponds to an excitation pulse with carrier frequency ω in the transverse plane in the presence of the main field and a time-varying gradient in the longitudinal direction.

Consider the Bloch equation in a frame of reference rotating about the z axis at frequency ω . Define the rotating magnetization vector, $\mathbf{M}_{\text{rot}}(z, t)$, through the relation

$$\mathbf{M}(z, t) = R_z(\omega t) \mathbf{M}_{\text{rot}}(z, t).$$

In this rotating frame of reference,

$$\frac{d\mathbf{M}_{\text{rot}}}{dt} = \mathbf{M}_{\text{rot}} \times \gamma \mathbf{B}_{\text{eff}},$$

where

$$\mathbf{B}_{\text{eff}} = \mathbf{B}_{\text{rot}} - \frac{\omega}{\gamma} \hat{\mathbf{z}}.$$

Assume $\omega = \gamma B_0$. Then

$$\mathbf{B}_{\text{eff}}(z, t) = B_{1,x}(t) \hat{\mathbf{x}} + B_{1,y}(t) \hat{\mathbf{y}} + G_z(t) z \hat{\mathbf{z}}.$$

Analytical solution for constant B_1 and G_z

Assume the RF pulse modulation and gradient are constant, so $B_{1,x}(t) \equiv B_{1,x}$, $B_{1,y}(t) \equiv B_{1,y}$, $G_z(t) \equiv G_z$, and

$$\mathbf{B}_{\text{eff}}(z) = B_{1,x}\hat{\mathbf{x}} + B_{1,y}\hat{\mathbf{y}} + G_z z\hat{\mathbf{z}}.$$

Solving the Bloch equation in the rotating frame under this assumption, we find \mathbf{M}_{rot} precesses about the axis

$$\mathbf{u}(z) := \frac{1}{\|\mathbf{B}_{\text{eff}}(z)\|} \mathbf{B}_{\text{eff}}(z)$$

with rotational frequency

$$\gamma \|\mathbf{B}_{\text{eff}}(z)\| = \gamma \sqrt{B_{1,x}^2 + B_{1,y}^2 + (G_z z)^2}.$$

Numerical solution for time-varying B_1 and G_z

- ▶ We no longer assume the RF pulse modulation and the gradient are constant. Suppose, however, that these signals are well-approximated by piecewise constant functions with time steps Δt . Let t_0 be the time of excitation and let $t_i := t_0 + i\Delta t$ for $i = 0, 1, 2, \dots, N$ be the set of times spanning the RF pulse.
- ▶ During the time interval $[t_i, t_{i+1}]$, \mathbf{M}_{rot} rotates about the axis

$$\mathbf{u}^i(z) := \frac{1}{\|\mathbf{B}_{\text{eff}}(z, t_i)\|} \mathbf{B}_{\text{eff}}(z, t_i)$$

by the angle

$$\begin{aligned}\theta^i(z) &:= \gamma \|\mathbf{B}_{\text{eff}}(z, t_i)\| \Delta t \\ &= \gamma \Delta t \sqrt{(B_{1,x}(t_i))^2 + (B_{1,y}(t_i))^2 + (G_z(t_i)z)^2}.\end{aligned}$$

Therefore, we may compute $\mathbf{M}_{\text{rot}}(z, t_N)$ by applying a sequence of rotations to $\mathbf{M}_{\text{rot}}(z, t_0)$.

Implementation using quaternions

- ▶ For each z position, define an initial orientation quaternion $m_0(z)$ with imaginary part $\mathbf{M}_{\text{rot}}(z, t_0)$ and zero real part.
- ▶ Define the rotation quaternion for each time interval $[t_i, t_{i+1}]$

$$q_i(z) := \exp \left(\frac{\theta_i(z)}{2} (u_x^i(z)\mathbf{i} + u_y^i(z)\mathbf{j} + u_z^i(z)\mathbf{k}) \right) .$$

- ▶ The total effect of the RF pulse is the rotation quaternion q_{RF} given by the product of the individual rotation quaternions:

$$q_{\text{RF}}(z) := q_N(z) \cdots q_0(z) .$$

- ▶ The final magnetization is the imaginary part of the orientation quaternion

$$m_N(z) := q_{\text{RF}}(z)m_0(z)q_{\text{RF}}(z)^{-1} .$$

Simulating free precession

Bloch equation for free precession

During free precession, the Bloch equation in the laboratory frame reads

$$\frac{d\mathbf{M}}{dt} = \mathbf{M} \times \gamma \mathbf{B} - \frac{(M_z - M_0)}{T_1} \hat{\mathbf{z}} - \frac{(M_x \hat{\mathbf{x}} + M_y \hat{\mathbf{y}})}{T_2},$$

where M_0 is the equilibrium magnetization along the z direction. Let $\mathbf{M} = M_x \hat{\mathbf{x}} + M_y \hat{\mathbf{y}} + M_z \hat{\mathbf{z}}$. Assume that the magnetic field only has a component in the z direction so that $\mathbf{B} = B_z \hat{\mathbf{z}}$.

Define the transverse magnetization as the complex number $M := M_x + jM_y$. The Bloch equation divides into two uncoupled equations

$$\begin{aligned}\frac{dM}{dt} &= - \left(\frac{1}{T_2} + i\gamma B_z \right) M \\ \frac{dM_z}{dt} &= - \frac{M_z - M_0}{T_1}.\end{aligned}$$

One may recover the x and y components of the magnetization as $M_x = \text{Re}(M)$ and $M_y = \text{Im}(M)$.

Solving these equations, we find

$$M(\mathbf{r}, t) = M(\mathbf{r}, 0) \exp\left(-\frac{t}{T_2(\mathbf{r})}\right) \exp\left(-i\gamma \int_0^t B_z(\mathbf{r}, t) dt\right)$$

and

$$M_z(\mathbf{r}, t) = M_0 + (M_z(\mathbf{r}, 0) - M_0) \exp\left(-\frac{t}{T_1(\mathbf{r})}\right).$$

The received MR signal is

$$s(t) = \int M(\mathbf{r}, t) d^3\mathbf{r},$$

for t in the readout interval.

Basic structure of our simulation

Basic setup

(Ignore metabolic conversion for now.)

- ▶ Define a set of “em”s.
- ▶ Each em has (at least) two data attributes: a magnetization vector $\boldsymbol{\mu} \in \mathbb{R}^3$ and a position vector $\mathbf{r} \in \mathbb{R}^3$.
- ▶ The simulation instantiates a large number of ems.
- ▶ The simulation reads in a pulse sequence of piecewise constant signals, a set of readout times T_{readout} , and a set of physical parameters $\{T_1(\mathbf{r}), T_2(\mathbf{r})\}$.
- ▶ The simulation steps through time. At each time step, the simulation can be in one of two states: excitation or free precession.

Excitation

- ▶ During excitation, an RF pulse and a z-gradient are applied.
- ▶ Excitation will cause a rotation of each em's magnetization vector in 3D space.
- ▶ I propose it will be most efficient to compute this rotation using quaternions.

- ▶ Assign a rotation quaternion $q_{\text{RF}} := 1 + \mathbf{0}$ to each em.
- ▶ For each time t_i in the excitation interval and each em:
 - ▶ Compute the position $\mathbf{r}(t_i + \Delta t)$ according to the physics (blood flow, Brownian motion, etc.).
 - ▶ Define

$$\mathbf{u} := \frac{1}{\|\mathbf{B}_{\text{eff}}(\mathbf{r}^*, t_i)\|} \mathbf{B}_{\text{eff}}(\mathbf{r}^*, t_i)$$

$$\theta := \gamma \|\mathbf{B}_{\text{eff}}(\mathbf{r}^*, t_i)\| \Delta t .$$

where $\mathbf{r}^* = (\mathbf{r}(t_i) + \mathbf{r}(t_i + \Delta t))/2$.

- ▶ Update the rotation quaternion using

$$q_{\text{RF}} := \exp(\theta/2(u_x \mathbf{i} + u_y \mathbf{j} + u_z \mathbf{k})) q_{\text{RF}} .$$

- ▶ At the end of the excitation interval (of length T seconds) update the rotation quaternion to account for the rotation in the laboratory frame:

$$q_{\text{RF}} := \exp(\omega T/2(1\mathbf{k}))q_{\text{RF}}$$

- ▶ Update the magnetization of each em using

$$\mu := \text{Im} \{ q_{\text{RF}}(0 + \mu)q_{\text{RF}}^{-1} \} .$$

Free precession

- ▶ During free precession, no RF pulse is applied; the only magnetic field present is along the longitudinal direction.
- ▶ During free precession, rotation occurs around the z-axis only (i.e. rotation occurs in the 2D transverse plane alone).
- ▶ I propose it will be most efficient to represent this rotation using complex numbers.

Free precession

- ▶ Assign a transverse magnetization $m := \mu_x + i\mu_y$ to each em.
- ▶ For each time t_i in the free precession interval and each em:
 - ▶ If $t_i \in T_{\text{readout}}$ then compute the MR signal sample

$$s(t_i) := \sum_{k \in (\text{set of ems})} m_k$$

and store in memory.

- ▶ Compute the position $\mathbf{r}(t_i + \Delta t)$ according to the physics (blood flow, Brownian motion, etc.).
- ▶ Update the transverse magnetization as

$$m := m \exp\left(-\frac{\Delta t}{T_2(\mathbf{r}^*)}\right) \exp(-i\gamma B_z(\mathbf{r}^*, t_i)\Delta t),$$

and update the longitudinal magnetization as

$$\mu_z := \mu_0 + (\mu_z - \mu_0) \exp(-\Delta t / T_1(\mathbf{r}^*)),$$

where $\mathbf{r}^* = (\mathbf{r}(t_i) + \mathbf{r}(t_i + \Delta t))/2$.

- At the end of the free precession interval, for each em, assign

$$\mu_x := \operatorname{Re}(m)$$

$$\mu_y := \operatorname{Im}(m).$$