Specification testing for errors-in-variables models[*]

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December 6, 2019

Abstract

This paper considers specification testing for regression models with errors-in-variables and proposes a test statistic comparing the distance between the parametric and nonparametric fits based on deconvolution techniques. In contrast to the methods proposed by Hall and Ma (2007) and Song (2008), our test allows general nonlinear regression models and possesses complementary local power properties. We establish the asymptotic properties of our test statistic for the ordinary and supersmooth measurement error densities. Simulation results endorse our theoretical findings: our test has advantages in detecting high-frequency alternatives and dominates the existing tests under certain specifications.

1 Introduction

In this paper, we propose a specification, or goodness-of-fit test, for (possibly nonlinear) regression models with errors-in-variables by comparing the distance between the parametric and nonparametric fits based on deconvolution techniques. We establish asymptotic properties of the test statistic and use the bootstrap procedure of Hall and Ma (2007) to obtain critical values. Compared to existing methods, our test allows nonlinear regression models and possesses complementary local power properties.

In the enormous literature on specification testing, relatively little attention has been given to the issue of measurement error despite its obvious need. Papers such as Zhu, Song and Cui (2003), Zhu and Cui (2005), and Cheng and Kukush (2004) proposed χ^2 statistics based on moment conditions of observables implied from errors-in-variables regression models. However,

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as is the case without measurement error, these tests are generally inconsistent for some fixed alternatives. Song (2008) proposed a consistent specification test for linear errors-in-variables regression models by comparing nonparametric and model-based estimators on the conditional mean function of the dependent variable Y given the *mismeasured* observable covariates W, that is $E[Y|W]$. As we clarify at the end of the next section, this approach may not have sensible local power for the original hypothesis on $E[Y|X]$, where X is a vector of error-free unobservable covariates. Hall and Ma (2007) proposed a nonsmoothing specification test for polynomial regression models with errors-in-variables, which is able to detect local alternatives at the $\sqrt{n}$-rate. We propose a smoothing specification test that complements Hall and Ma's (2007) test (see further discussion below). Other papers that study specification testing under measurement error include Butucea (2007), Holzmann and Boysen (2006) and Holzmann, Bissantz and Munk (2007) (for testing probability densities), Koul and Song (2009, 2010) (for Berkson measurement error models), and Song (2009) and Xu and Zhu (2015) (for errors-in-variables models with validation data).

Consistent specification tests can be broadly split into those that use a nonparametric estimator (called smoothing tests) and those that do not (called nonsmoothing or integral-transform tests). In contrast to Hall and Ma (2007) which adopted the nonsmoothing approach, we propose a kernel-based smoothing test for the goodness-of-fit of parametric regression models with errors-in-variables. There are two important features of our test. First, our smoothing test is not restricted to polynomial models and allows testing of general nonlinear regression models. Second, analogous to the literature on conventional specification testing, our smoothing test complements Hall and Ma's (2007) test (if applied to polynomial models) due to its distinct power properties. Rosenblatt (1975) explained that although local power properties of nonsmoothing tests suggest they are more powerful than smoothing tests, 'there are other types of local alternatives for which tests based on density estimates are more powerful'. Fan and Li (2000) showed that in the non-measurement error case, smoothing tests are generally more powerful for high-frequency alternatives and nonsmoothing tests are more powerful for low-frequency alternatives. Thus, smoothing tests 'should be viewed as complements to, rather than substitutes for, [nonsmoothing tests].' Our simulation results suggest that this phenomenon extends to errors-in-variables models.

In contrast to the above papers and our own, Ma *et al.* (2011) moved away from Wald-type tests where restricted and unrestricted estimates are compared. They proposed a local test that is more analogous to the score test where only the model under the null hypothesis must be estimated, and they then extended this to an omnibus test using a system of different basis

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functions. However, the test does not allow the number of basis functions to increase with the sample size and so is not strictly a nonparametric test. They also discuss that as a result of the way the test is constructed, it has low power against high-frequency alternatives - similar to the tests of Song (2008) and Hall and Ma (2007).

The rest of this paper is organized as follows. Section 2 describes the setup in detail and introduces the test statistic and its motivation. Section 3 outlines the main asymptotic properties of the test statistic, outlines the bootstrap procedure used, and discusses how to implement the test, including in the case where the distribution of the measurement error is unknown but repeated measurements are available. Finally, Section 4 analyses the small sample properties of the test through a Monte Carlo simulation. All mathematical proofs are deferred to the Appendix.

2   Setup and test statistic

Consider the nonparametric regression model

$$Y = m(X) + U \quad \text{with } E[U|X] = 0,$$

where $Y \in \mathbb{R}$ is a response variable, $X \in \mathbb{R}^d$ is a vector of covariates, and $U \in \mathbb{R}$ is the error term. In this paper, we focus on the situation where X is not directly observable due to the measurement mechanism or nature of the environment. Instead a vector of variables W is observed through

$$W = X + \epsilon,$$

where $\epsilon \in \mathbb{R}^d$ is a vector of measurement errors that has a known density f_ϵ and is independent of (Y, X). The case of unknown density f_ϵ will be discussed in Section 3.2. We are interested in specification, or goodness-of-fit, testing of a parametric functional form of the regression function m. More precisely, for a parametric model m_θ, we wish to test the hypothesis

$$\text{H}_0 \quad : \quad \text{there exists some } \theta \in \mathbb{R}^k \text{ such that } m(x) = m_\theta(x) \text{ for almost every } x \in \mathbb{R}^d,$$

$$\text{H}_1 \quad : \quad \text{H}_0 \text{ is false},$$

based on the random sample $\{Y_i, W_i\}_{i=1}^n$ of observables (while X_i is unobservable).

To test the null H_0, we adapt the approach of Härdle and Mammen (1993), which compares nonparametric and parametric regression fits, to the errors-in-variables model. As a nonpara-

metric estimator of m, we use the deconvolution kernel estimator (see, e.g., Fan and Truong, 1993, and Meister, 2009, for a review)

$$\hat{m}(x) \;=\; \frac{\sum_{i=1}^{n} Y_i \mathcal{K}_b(x - W_i)}{\sum_{i=1}^{n} \mathcal{K}_b(x - W_i)},$$

where

$$\mathcal{K}_b(a) = \frac{1}{(2\pi)^d} \int e^{-\mathrm{i}t \cdot a} \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)} dt,$$

is the so-called deconvolution kernel, $\mathrm{i} = \sqrt{-1}$, b is a bandwidth, and K^{ft} and f_ϵ^{ft} are the Fourier transforms of a kernel function K and the measurement error density f_ϵ, respectively.

To simplify the exposition, we concentrate on the case where all elements of X are mismeasured. If X contains both correctly measured and mismeasured covariates (denoted by X_1 and X_2, respectively), then the kernel estimator is modified as $\hat{m}(x) = \frac{\sum_{i=1}^{n} Y_i K_{1b}(x_1 - X_{1i}) \mathcal{K}_b(x_2 - W_i)}{\sum_{i=1}^{n} K_{1b}(x_1 - X_{1i}) \mathcal{K}_b(x_2 - W_i)}$, where $K_{1b}(a) = \frac{1}{b^{d_1}} K_1\left(\frac{a}{b}\right)$ and K_1 is a conventional kernel function for X_1, and analogous results can be established.

Throughout the paper we assume $f_\epsilon^{\mathrm{ft}}(t) \neq 0$ for all $t \in \mathbb{R}^d$ and K^{ft} has compact support so that the above integral is well-defined.

For the parametric functional form, m_θ, several methods are available to estimate θ under certain regularity conditions. For example, based on Butucea and Taupin (2008), we can estimate θ by the (weighted) least squares regression of Y on the implied conditional mean function $E[m_\theta(X)|W]$. In this paper, we do not specify the construction of the estimator $\hat{\theta}$ for θ except for assuming consistency.

In order to construct a test statistic for H_0, we compare the nonparametric and parametric estimators of the regression function based on the L_2-distance,

$$D_n = n \int \left| \hat{m}(x)\hat{f}(x) - m_{\hat{\theta}}(x)\hat{f}(x) \right|^2 dx \tag{1}$$

where $|\cdot|$ is the Euclidean norm, and $\hat{f}(x) = \frac{1}{n}\sum_{i=1}^{n} \mathcal{K}_b(x - W_i)$ is the deconvolution kernel density estimator for X.

We close this section with a remark on an alternative testing approach. To test the null hypothesis H_0, one may consider testing an alternative formulation of H_0 based on the conditional mean $E[Y|W]$ of observables, i.e., consider $\mathrm{H}_0' : f_W(w)E[Y|W = w] = \int m_\theta(w - u)f_X(w - u)f_\epsilon(u)du$ for almost every w, and test H_0' by a conventional method, such as Härdle and Mammen (1993). This approach was employed by Song (2008), and we argue that our testing approach can be a useful complement to this method since it may have non-trivial local power against

local alternatives which may not be detected by Song's (2008) approach.

To illustrate this point, consider the following local alternative hypothesis for the regression function

$$m_n(x) = m_\theta(x) + 2a_n \cos(A_n x) \left(\frac{\sin x}{x}\right),$$

where $a_n \to 0$ and $A_n \to \infty$ as $n \to \infty$. In this case, m_n converges to m_θ at the rate of a_n under the L_2-norm, and the test based on D_n will have non-trivial power for a certain rate of a_n. On the other hand, local power of the test based on the implied null H_0' is determined by the L_2-norm of the convolution $\{(m_n - m_\theta)f_X\} * f_\epsilon$. By Parseval's identity and the Fourier shift formula, we have

$$\|\{(m_n - m_\theta)f_X\} * f_\epsilon\|^2 = a_n^2 \left\|\{q^{\mathrm{ft}}(\cdot - A_n) + q^{\mathrm{ft}}(\cdot + A_n)\}f_\epsilon^{\mathrm{ft}}\right\|^2,$$

where $q(x) = \left(\frac{\sin x}{x}\right)f_X(x)$. For example, if f_ϵ is Laplace with $f_\epsilon^{\mathrm{ft}}(t) = 1/(1 + t^2)$, then we have that the L_2-norm $\|\{(m_n - m_\theta)f_X\} * f_\epsilon\|$ is of order a_n/A_n^2. By letting A_n diverge at an arbitrarily fast rate, the rate a_n/A_n^2 becomes arbitrarily fast so that any conventional test for H_0' fails to detect deviations from this null.

We emphasize that our test should be treated as a complement to the existing approach by Song (2008). One may construct examples where Song's (2008) approach yields better power properties. This point is illustrated in our simulation study in Section 4.

3 Asymptotic properties

In this section, we present asymptotic properties of the test statistic D_n. We first derive the limiting distribution under the null hypothesis H_0. To this end, we impose the following assumptions.

Assumption D.

(i) $\{Y_i, X_i, \epsilon_i\}_{i=1}^n$ *are i.i.d.* ϵ *is independent of* (Y, X) *and has a known density* f_ϵ. *Elements of* ϵ *are mutually independent.*

(ii) m_θ *is twice differentiable in a neighborhood of* θ. *Each of the following are in* L_1 *and* L_2: $m^2 f_X$, $m_2 f_X$, $\left(\frac{dm_\theta}{d\theta}\right)^2 f_X$, *and* $\frac{d^2 m_\theta}{d\theta^2} f_X$, *where* $m_2(x) = E[Y^2|X = x]$.

(iii) $K^{\mathrm{ft}}(t)$ *is compactly supported on* $[-1, 1]^d$, *is symmetric around zero (i.e.,* $K^{\mathrm{ft}}(t) = K^{\mathrm{ft}}(-t)$), *and is bounded.*

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(iv) *As $n \to \infty$, it holds that $b \to 0$ and $nb^d \to \infty$.*

(v) $\hat{\theta} - \theta = o_p(1)$ *under* H_0.

Assumption D (i) is common in the literature of classical measurement error. Extensions to the case of unknown f_ϵ are discussed in Section 3.2. Assumption D (ii) contains boundedness conditions on f_X, the regression functions m and m_θ, and m_2. Assumption D (iii) and (iv) contain standard conditions on the kernel function K and bandwidth b, respectively. A popular choice for the kernel function in the context of deconvolution methods is the sinc kernel $K(x) = \frac{\sin x}{\pi x}$ whose Fourier transform is equal to $K^{\text{ft}}(t) = \mathbb{I}\{-1 \le t \le 1\}$. Assumption (v) stipulates that we only require some consistent estimator of θ. When the regression model under the null hypothesis is linear (i.e., $m_\theta(x) = x'\theta$), we can employ the methods in, for example, Gleser (1981), Bickel and Ritov (1987), or van der Vaart (1988). For nonlinear regression, we may choose the estimators by e.g., Taupin (2001) or Butucea and Taupin (2008) under certain regularity conditions. It is interesting to note that in contrast to the no measurement error case as in Härdle and Mammen (1993), the limiting distribution of the estimation error $(\hat{\theta} - \theta)$ does not influence the first-order asymptotic properties of the test statistic D_n. This is due to the measurement error slowing down the convergence rate of the dominant term of D_n.

For simplicity, we use product kernels in the following way. Following Masry (1993), let $\tilde{K}(\cdot)$ be a univariate kernel and $\tilde{K}^{\text{ft}}(\cdot)$ denote its Fourier transform. Define the univariate deconvolution kernel as

$$\tilde{\mathcal{K}}_b(x_j) = \frac{1}{2\pi b} \int e^{-ita} \frac{\tilde{K}^{\text{ft}}(t)}{\tilde{f}_{\epsilon_j}^{\text{ft}}(t/b)} dt.$$

where $\tilde{f}_{\epsilon_j}^{\text{ft}}(\cdot)$ is the Fourier transform of ϵ_j. Finally, set $K(x) = \prod_{j=1}^{\dim(x)} \tilde{K}(x_j)$ and $\mathcal{K}_b(x) = \prod_{j=1}^{\dim(x)} \tilde{\mathcal{K}}_b(x_j)$. Since we assume that ϵ is vector valued with independent elements, we can write $f_\epsilon^{\text{ft}}(t) = \prod_{j=1}^{\dim(t)} \tilde{f}_{\epsilon_j}^{\text{ft}}(t_j)$. Combining these facts, we have $K^{\text{ft}}(x) = \prod_{j=1}^{\dim(x)} \tilde{K}^{\text{ft}}(x_j)$.

We impose additional assumptions based on the bounds of the rate of decay of the tail of the characteristic function of the measurement error, f_ϵ^{ft}. Let $\sigma^2(x) = E[U^2|X = x]$ be the conditional variance of the error term. The first case, known as ordinary smooth measurement error, is characterised as follows.

Assumption O.

$f_\epsilon^{\text{ft}}(t) \ne 0$ *for all $t \in \mathbb{R}^d$ and there exist positive constants c, C, and α such that*

$$c|s|^{-\alpha} \le |\tilde{f}_{\epsilon_j}^{\text{ft}}(s)| \le C|s|^{-\alpha},$$

for all $1 \leq j \leq d$ *as* $|s| \to \infty$.

Assumption O requires that the Fourier transform $\tilde{f}^{\text{ft}}_{\epsilon_j}$ decays in some finite power. A popular example of an ordinary smooth density is the Laplace density. Although it is beyond the scope of this paper, isolated zeros in $f^{\text{ft}}_{\epsilon}(t)$ may be allowed by introducing the spectral method (Carrasco and Florens, 2011) or ridge approach (Hall and Meister, 2007).

For the second case, known as supersmooth measurement error, we impose the following assumptions.

Assumption S. (i) $f^{\text{ft}}_{\epsilon}(t) \neq 0$ *for all* $t \in \mathbb{R}^d$ *and there exist positive constants* C_{ϵ}, μ, γ_0, *and* $\gamma > 1$ *such that*

$$\tilde{f}^{\text{ft}}_{\epsilon_j}(s) \sim C_{\epsilon}|s|^{\gamma_0} e^{-|s|^{\gamma}/\mu},$$

as $|s| \to \infty$.

(ii) *There exist constants* $A > 0$ *and* $\beta \geq 0$ *such that*

$$\tilde{K}^{\text{ft}}(1 - s) = As^{\beta} + o(s^{\beta}),$$

as $s \to 0$.

(iii) $E[Y^4] < \infty$, $E[W^4] < \infty$.

Assumption S (i) is adopted from Holzmann and Boysen (2006). This assumption requires that the Fourier transform $\tilde{f}^{\text{ft}}_{\epsilon}$ decays at an exponential rate. An example of a supersmooth density satisfying this assumption is the normal density, where $C_{\epsilon} = 1$, $\gamma_0 = 0$, $\gamma = 2$, and $\mu = 2$. However, due to the requirement $\gamma > 1$, the Cauchy density is excluded. As is clarified in the proof of Theorem 1 (iii), the condition $\gamma > 1$ is imposed to make a bias term negligible. Assumption S (ii) contains an additional condition on the kernel function. The sinc kernel, $K(x) = \frac{\sin x}{\pi x}$, for example, satisfies this assumption with $A = 1$ and $\beta = 0$. Assumption S (iii) requires bounded fourth moments for our observable random variables.

Under these assumptions, the null distribution of D_n is given as follows.

Theorem 1.

(i) *Suppose that Assumptions D and O hold true. Then under* H_0,

$$C_{V,b}^{-1/2} D_n \overset{d}{\to} N\left(0, \frac{2}{(2\pi)^{2d}}\right),$$

where $C_{V,b} = O(b^{-d(1+4\alpha)})$ *is defined in (2) in the Appendix.*

(ii) *Suppose that Assumptions D and S hold true and ϵ is multivariate normal. Then under* H_0,

$$\varphi(b)D_n \xrightarrow{d} \sum_{k=1}^{\infty} \lambda_k(Z_k^2 - 1),$$

where $\varphi(b) = \frac{(2\pi)^d}{b^{d(1+4\beta)}e^{d/b^2}A^{2d}\Gamma(1+2d\beta))}$ *with the gamma function* Γ, $\{Z_k\}$ *is an independent sequence of standard normal random variables and* $\{\lambda_k\}$ *is defined in (6) in the Appendix.*

(iii) *Suppose that Assumptions D and S hold true. Then under* H_0,

$$\phi(b)D_n \xrightarrow{d} \sum_{k=1}^{\infty} \lambda_k(Z_k^2 - 1),$$

where $\phi(b) = \frac{(2\pi)^d \gamma^{1+2d\beta} C_\epsilon^{2d}}{\mu^{1+2d\beta}b^{d(\gamma-1+2\gamma\beta+2\gamma_0)}e^{2d/(\mu b^\gamma)}A^2\Gamma(1+2d\beta)}.$

Theorem 1 (i) says that for the ordinary smooth case, the test statistic D_n is asymptotically normal. The normalizing term $C_{V,b}$ comes from the variance of the U-statistic of the leading term in D_n. Note that the convergence rate $C_{V,b}^{-1/2} = O(b^{d(\frac{1}{2}+2\alpha)})$ of the statistic D_n is slower than the rate $O(b^{d/2})$ of Härdle and Mammen's (1993) statistic for the no measurement error case. As the dimension d of X or the decay rate α of f_ϵ^{ft} increases, the convergence rate of D_n becomes slower. This theorem can also be used to show the 'rule-of-thumb' rate for the bandwidth is of the order $b = n^{-\left(\frac{1}{4+d+2\alpha}\right)}$.

Theorem 1 (ii) focuses on the case of normal measurement error and shows that the test statistic converges to a weighted sum of chi-squared random variables. The normalizing term $\varphi(b)$ is characterized by the shape of the kernel function specified in Assumption S (i). For example, if we employ the sinc kernel (i.e., $A = 1$ and $\beta = 0$), the normalization becomes $\varphi(b) = \frac{(2\pi)^d}{b^d e^{d/b^2}\Gamma(d)}$. In this supersmooth case, the non-normal limiting distribution emerges because the leading term of the statistic D_n is characterized by a degenerate U-statistic with a fixed kernel (see, e.g., Serfling, 1980, Theorem 5.5.2). In contrast, for the ordinary smooth case in Part (i) of this theorem, the leading term is characterized by a U-statistic with a varying kernel so that the central limit theorem in Hall (1984) applies. An analogous result is obtained in Holzmann and Boysen (2006) for the integrated squared error of the deconvolution density estimator.

Theorem 1 (iii) presents the limiting null distribution of the test statistic for the case of general supersmooth measurement error. In this case, after normalization by $\varphi(b)$, the test statistic obeys the same limiting distribution as the normal case in Part (ii) of this theorem. Thus, similar comments to Part (ii) apply. The normalization term $\varphi(b)$ is characterized by the shapes of the kernel function and Fourier transform $f_\epsilon^{\mathrm{ft}}(t)$ of the measurement error specified in

Assumption S (i).

In order to investigate the power properties of the test based on T_n, we consider a local alternative hypothesis of the form

$$H_{1n} : m(x) = m_\theta(x) + c_n \Delta(x), \text{ for almost every } x \in \mathbb{R}^d$$

where $c_n \to 0$ and $\Delta(x)$ is a non-zero function satisfying $\Delta(\cdot)^2 f_X(\cdot) \in L_2$ and such that the limit $\lim_{n\to\infty} \Delta_n$ defined in (7) in the Appendix exists. The local power properties are obtained as follows.

Theorem 2.

(i) *Suppose that Assumptions D and O hold true. Then under* H_{1n} *with* $c_n = n^{-1/2} b^{-d(\frac{1}{4}+\alpha)}$,

$$C_{V,b}^{-1/2} D_n \xrightarrow{d} N\left(\lim_{n\to\infty} \Delta_n, \frac{2}{(2\pi)^{2d}}\right).$$

(ii) *Suppose that Assumptions D and S hold true and* ϵ *is multivariate normal. Then under* H_{1n} *with* $c_n = n^{-1/2} b^{d(1/2+2\beta)} e^{d/(2b^2)}$,

$$\varphi(b) D_n \xrightarrow{d} \lim_{n\to\infty} \Delta_n + \sum_{k=1}^{\infty} \lambda_k (Z_k^2 - 1).$$

(iii) *Suppose that Assumptions D and S hold true. Then under* H_{1n} *with* $c_n = b^{d((\lambda-1)/2+\lambda\beta+\lambda_0)} e^{d/(\mu b^\lambda)}$,

$$\varphi(b) D_n \xrightarrow{d} \lim_{n\to\infty} \Delta_n + \sum_{k=1}^{\infty} \lambda_k (Z_k^2 - 1),$$

Theorem 2 (i) says that in the ordinary smooth case, our test has non-trivial power against local alternatives drifting with the rate of $c_n = n^{-1/2} b^{-d(\frac{1}{4}+\alpha)}$. This is a nonparametric rate, with the test becoming less powerful as the dimension d of X or the decay rate α of f_ϵ^{ft} increases. For the no measurement error case, Härdle and Mammen's (1993) statistic has non-trivial power for local alternatives with the rate $n^{-1/2} b^{-d/4}$. Therefore, as expected, the test becomes less powerful in the presence of measurement error. Theorem 2 (ii) and (iii) present local power properties of our test for the normal and general supersmooth measurement error cases, respectively. Except for the normalizing constants, the test statistic has the same limiting distribution. Also, for $c_n \to 0$, the bandwidth b should decay at a log rate. As an example, consider the case of $\epsilon \sim N(0, 1)$. In this case, if we choose $b \sim \left(\frac{d}{\log(n)}\right)^{1/2}$, then the rate for the local alternative will

be $c_n \sim \left(\frac{d}{log(n)}\right)^{d(1/4+\beta)}$. So the rate at which we can detect local alternatives is typically a log rate in the supersmooth measurement error case.

3.1 Bootstrap and implementation

As is typically the case in testing problems which involve nonparametric components, the asymptotic distribution is relatively complex and can be difficult to estimate in practice. As such, we suggest a bootstrap procedure to calculate critical values for this test. We follow the approach of Hall and Ma (2007) who propose a type of moment matching bootstrap that is adapted to the setting of mismeasured regressors. Measurement error poses many problems for the implementation of bootstrap procedures, resulting in a relative scarcity of bootstraps in the literature. We do not have access to the true regressors, so cannot compute any residual based bootstrap, we also do not have access to the measurement errors which provides further issues. Hall and Ma (2007) suggest the following approach in answer to these issues. Compute $\hat{m}$, $\hat{f}$, $m_{\hat{\theta}}$, and $\hat{F}_X$ from the original data using standard deconvolution techniques and a parametric estimator for $\hat{\theta}$. Estimate $\omega_2 = E[U^2]$ (Hall and Ma, 2007, provide a method to do this in their polynomial setting. For more general models, one can use the approach of Delaigle, Hall and Jamshidi, 2015) and choose a distribution, $G_\epsilon(\cdot|\hat{\omega}_2)$, which has first and second moments equal to 0 and $\hat{\omega}_2$, respectively (in Section 4 we use Mammen's 2-point distribution for G_ϵ). Now, draw data $\{X_i^*\}_{i=1}^n$ from $\hat{F}_X$, $\{\epsilon_i^*\}_{i=1}^n$ from $F_\epsilon(\cdot)$, and $\{U_i^*\}_{i=1}^n$ from $G_\epsilon(\cdot|\hat{\omega}_2)$. Construct $Y_i^* = m_{\hat{\theta}}(X_i^*) + U_i^*$ and $W_i^* = X_i^* + \epsilon_i^*$. Using the pairs (Y_i^*, W_i^*) compute D_n^* in the same manner as D_n. The distribution of D_n^* can be approximated by repeatedly carrying out this procedure.

In order to implement our test, we also need to choose a suitable bandwidth. Unfortunately, to the best of our knowledge, there is currently no theory on the optimal choice of bandwidth in a measurement error testing situation. It is beyond the scope of this paper to provide such a choice, however, we recommend the estimated plug-in approach of Delaigle and Gijbels (2004) as a good starting point; this is the approach we use in our Monte Carlo study. Although this procedure is designed to select the optimal bandwidth in a regression context - rather than for a testing problem - it provides good results in terms of power and size control (as shown in Section 4) and tends to produce greater power than other bandwidth choices around the selected value.

In summary, we suggest the following procedure for implementation of our specification test.

1. Based on the deconvolution regression and density estimators $\hat{m}$ and $\hat{f}$, and the parameter estimator $\hat{\theta}$, compute the test statistic D_n in (1). The bandwidths b in $\hat{m}$ and $\hat{f}$ are chosen as outlined in the 'two-stage selection plug-in' method of Delaigle and Gijbels

(2004). That is, start with a normal-reference assumption to calculate $\int f_X^{(4)}(x)^2 dx$ (where $f_X^{(4)}(x)$ denotes the 4th order derivative) which can be used to select an 'optimal' plug-in bandwidth for estimating $\int f_X^{(3)}(x)^2 dx$ (required for the estimation of the asymptotic mean integrated squared error (AMISE) for $f_X^{(1)}$). In turn this bandwidth can be used to estimate $\int f_X^{(2)}(x)^2 dx$ which is plugged into the AMISE for f_X. The final bandwidth is given as the minimiser of this estimated AMISE.

2. Compute the cdf $\hat{F}_X$ by integrating $\hat{f}$ obtained in Step 1, and compute $\hat{\omega}_2$ using the method of Delaigle, Hall and Jamshidi (2015), that is, $\hat{\omega}_2 = \max\{\hat{\mu}_2 - \hat{\xi}_2, 0\}$, where $\hat{\mu}_2 = \frac{1}{n} \sum_{i=1}^n Y_i^2$ and $\hat{\xi}_2 = \int \hat{m}(x)^2 \hat{f}(x) dx$.

3. Draw $\{X_i^*\}_{i=1}^n$ from $\hat{F}_X$, $\{\epsilon_i^*\}_{i=1}^n$ from F_ϵ, and $\{U_i^*\}_{i=1}^n$ from $G_\epsilon(\cdot|\hat{\omega}_2)$ (some distribution with variance $\hat{\omega}_2$, such as Mammen's 2 point distribution). Construct $Y_i^* = m_{\hat{\theta}}(X_i^*) + U_i^*$ and $W_i^* = X_i^* + \epsilon_i^*$ for $i = 1, \ldots, n$. Using the pairs $\{Y_i^*, W_i^*\}_{i=1}^n$ compute the bootstrap counterpart D_n^* in the same manner as D_n in (1).

4. Repeat Step 3 many times and use the $(1-\alpha)$-th quantile of the distribution of D_n^* as the critical value.

3.2 Case of unknown f_ϵ

In practical applications, it is unrealistic to assume that the density of the measurement error, f_ϵ, is known to the researcher. In the literature on nonparametric deconvolution, several estimation methods for f_ϵ are available: these are typically based on additional data (see, e.g., Section 2.6 of Meister (2009) for a review). Although the analysis of the asymptotic properties is different, we can modify the test statistic D_n by inserting the estimated Fourier transform of the measurement error density, $\hat{f}_\epsilon^{\text{ft}}$.

For example, suppose the researcher has access to repeated measurements on X in the form of $W = X + \epsilon$ and $W^r = X + \epsilon^r$, where ϵ and ϵ^r are identically distributed and $(X, \epsilon, \epsilon^r)$ are mutually independent, see Delaigle, Hall and Meister (2008) for a list of examples of such repeated measurements. If we further assume that the Fourier transform f_ϵ^{ft} is real-valued (i.e. f_ϵ is symmetric around zero), then we can employ the estimator proposed by Delaigle, Hall and Meister (2008)

$$\hat{f}_\epsilon^{\text{ft}}(t) = \left| \frac{1}{n} \sum_{i=1}^n \cos\{t(W_i - W_i^r)\} \right|^{1/2}.$$

Delaigle, Hall and Meister (2008) studied the asymptotic properties of the deconvolution density estimator and the regression estimator using $\hat{f}_\epsilon^{\text{ft}}$ and found conditions to guarantee that the

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differences between the estimators with known f_ϵ and those with unknown f_ϵ are asymptotically negligible. Under similar conditions, we can expect that the asymptotic distributions of the test statistic D_n obtained above remain unchanged when we replace f_ϵ^{ft} with $\hat{f}_\epsilon^{\mathrm{ft}}$. If the researcher wishes to remove the assumption that f_ϵ^{ft} is real-valued and ϵ and ϵ^r are identically distributed, it may be possible to employ the estimator by Li and Vuong (1998) based on Kotlarski's identity.

4 Simulation

We evaluate the small sample performance of our test through a Monte Carlo experiment. We also compare our test to Hall and Ma (2007) and Song (2008). Recall that although Hall and Ma's (2007) test is confined to a polynomial regression model and Song's (2008) test places restrictions on the type of model that can be tested, our test allows any nonlinear model. We take the true unobservable regressor $\{X_i\}_{i=1}^n$ to be distributed as $U[-1,1]$ and $Y_i = X_i + 0.5\sin(\delta X_i) + U_i$, where $U_i \sim N(0, 1/4)$ and δ is a constant to be varied. The contaminated regressor is given by $W_i = X_i + \epsilon_i$. We consider two distributions for ϵ_i to be drawn from. For the ordinary smooth case, we use the Laplace distribution with variance of $1/12$. For the supersmooth case, we use $N(0, 1/12)$. So the signal-to-noise ratio is $\sigma_X^2/\sigma_\epsilon^2 = 4$ in both cases. We use the following kernel for all of our simulations (Fan, 1992)

$$K(x) = \frac{48\cos(x)}{\pi x^4}\left(1 - \frac{15}{x^2}\right) - \frac{144\sin(x)}{\pi x^5}\left(2 - \frac{5}{x^2}\right).$$

We report results for a small (250) and a medium (500) sample size. As discussed in Section 3.1, we use the estimated plug-in approach of Delaigle and Gijbels (2004) to select our bandwidth. For the moment matching bootstrap, we use Mammen's 2-point distribution for $G(\cdot|\omega_2)$.

For the parametric estimator, we use the polynomial estimator of degree 1 proposed by Cheng and Schneeweiss (1998) so as to remain consistent with the experiment conducted by Hall and Ma (2007). For the test of Song (2008) we use the same kernel as for our test and also choose bandwidths by cross-validation. All results are based on 1000 Monte Carlo replications.

Table 1 shows the results for the power and the size of three tests. The column labelled D_n refers to the test proposed in this paper, the column labelled 'S' corresponds to the test put forward in Song (2008), and the column labelled 'HM' corresponds to the test proposed by Hall and Ma (2007). When $\delta = 0$ we have the null model. We also consider $\delta = \{1, 2, 3\}$ corresponding to departures from the null of varying frequencies which we term low, medium, and high frequency alternatives, respectively.

The results are encouraging and seem to be consistent with the theory. The first column ($\delta = 0$) indicates that all three tests track the nominal level closely. This indicates that the bootstrap procedure discussed in Section 3.1 works well in practice.

As we expected, across all parameter settings, our test performs more poorly for supersmooth measurement error in comparison to ordinary smooth error. This reflects the slower convergence rates derived in Section 3. Somewhat surprisingly, the test of Hall and Ma (2007) lags behind the other two tests in all cases, although it has the benefit of not requiring a tuning parameter.

As conjectured, in the low frequency setting ($\delta = 1$) our test is generally less powerful than the test of Song (2008). Although, in the ordinary smooth case for a sample size of 250, our test displays higher power. This reflects that the test of Song (2008) is able to detect local alternatives at the rate $\sqrt{nb^{d/2}}$ for both ordinary and supersmooth measurement error distributions. However, our test achieves a slightly slower polynomial rate in the ordinary smooth case and only a $\log(n)$-rate in the super smooth case. Thus it is not surprising to see our test underperform when faced with Gaussian measurement error.

On the other hand, as previously mentioned, we suspect that our test is better suited to detecting high-frequency alternatives than Song (2008). This is confirmed in the final column of Table 1 ($\delta = 3$). We find that the test of Song (2008) and Hall and Ma (2007) have trivial power in this setting for both sample sizes, while our test still enjoys reasonable power which increases with the sample size.

The dominance over the test of Hall and Ma (2007) in this high-frequency setting is intuitive and is explained in Fan and Li (2000). Nonsmoothing tests can be thought of as smoothing tests but with a fixed bandwidth. Thus, it is the asymptotically vanishing nature of the bandwidth in smoothing tests that allows for the superior detection of high-frequency alternatives over nonsmoothing tests. Furthermore, as discussed at the end of Section 2, the test of Song (2008) will have poor power properties for some high-frequency alternatives due to testing the hypothesis based on $E[Y|W]$ rather than $E[Y|X]$.

Table 1: $Y = X + 0.5\sin(\delta X) + U$

Ordinary Smooth										δ			
		0			**1**			**2**			**3**		
n	**Level**	D_n	S	HM	D_n	S	HM	D_n	S	HM	D_n	S	HM
250	**5%**	5.5	4.4	4.9	72.2	44.3	16.1	92.8	83.6	31.4	16.2	4.9	4.1
	10%	10.4	9.3	9.9	80.9	63.4	32.8	94.8	93.3	50.7	24.3	9.6	8.8

500	**5%**	5.1	4.2	4.7	92.1	86.1	37.1	97.9	99.8	64.6	37.4	4.1	4.0
	10%	10.2	8.3	9.8	93.8	93.9	58.3	98.1	99.9	81.0	45.5	9.2	8.1

Supersmooth

250	**5%**	5.2	4.7	5.1	25.5	70.0	15.9	55.6	89.5	23.3	15.4	5.4	5.2
	10%	10.9	9.5	9.8	33.6	82.3	30.3	62.5	95.1	41.2	27.8	11.0	10.6
500	**5%**	5.0	4.1	4.2	51.1	98.6	29.4	85.9	99.8	46.2	27.7	5.2	4.0
	10%	11.7	8.9	8.8	59.9	99.4	49.1	90.0	100	65.5	41.6	10.2	10.2

A   Mathematical Appendix

Hereafter, $f(b) \sim g(b)$ means $f(b)/g(b) \to 1$ as $b \to 0$.

A.1   Proof of Theorem 1

A.1.1   Proof of (i)

We start be decomposing D_n as

$$D_n = \frac{1}{n(2\pi)^d} \sum_{i=1}^n \int |\zeta_i(x)|^2 dx + \frac{1}{n(2\pi)^d} \sum_{i \neq j} \int \zeta_i(x)\overline{\zeta_j(x)} dx \equiv B_n + T_n,$$

where $\zeta_i(x) = \int \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)} \{Y_i - m_{\hat{\theta}}(x)\} e^{-\mathrm{i}t \cdot (x-W_i)} dt$ and $\overline{\zeta_j(x)}$ is the complex conjugate of $\zeta_j(x)$. We will show that B_n is of smaller asymptotic order than T_n: as such, T_n will characterise the asymptotic distribution of the test.

First, we define the normalization term $C_{V,b}$ and characterize its asymptotic order. Let

$$\xi_i(x) \equiv \int \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)} \{Y_i - m_\theta(x)\} e^{-\mathrm{i}t \cdot (x-W_i)} dt, \qquad H_{i,j} \equiv \int \xi_i(x)\overline{\xi_j(x)} dx.$$

Then $C_{V,b}$ is defined as

$$C_{V,b} \equiv E[H_{1,2}^2] = \int \int |E[\xi_1(x)\xi_1(z)]|^2 dx dz \equiv Q_1 + Q_2 - 2Q_3, \qquad (2)$$

where

$$
\begin{aligned}
Q_1 &= \int \int \left| \int \int E\left[Y_i^2 e^{-\mathrm{i}t \cdot (x-W_i)} e^{-\mathrm{i}s \cdot (z-W_i)}\right] \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)} \frac{K^{\mathrm{ft}}(sb)}{f_\epsilon^{\mathrm{ft}}(s)} dt ds \right|^2 dx dz, \\
Q_2 &= \int \int \left| \int \int E\left[e^{-\mathrm{i}t \cdot (x-W_i)} e^{-\mathrm{i}s \cdot (z-W_i)}\right] m_\theta(x) m_\theta(z) \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)} \frac{K^{\mathrm{ft}}(sb)}{f_\epsilon^{\mathrm{ft}}(s)} dt ds \right|^2 dx dz, \\
Q_3 &= \int \int \left| \int \int E\left[Y_i e^{-\mathrm{i}t \cdot (x-W_i)} e^{-\mathrm{i}s \cdot (z-W_i)}\right] m_\theta(z) \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)} \frac{K^{\mathrm{ft}}(sb)}{f_\epsilon^{\mathrm{ft}}(s)} dt ds \right|^2 dx dz.
\end{aligned}
$$

For Q_1, we have

$$
\begin{aligned}
Q_1 &= \int \int \left| \int \int f_\epsilon^{\mathrm{ft}}(t+s) \int m_2(v) f_X(v) e^{-\mathrm{i}t \cdot (x-v)} e^{-\mathrm{i}s \cdot (z-v)} dv \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)} \frac{K^{\mathrm{ft}}(sb)}{f_\epsilon^{\mathrm{ft}}(s)} dt ds \right|^2 dx dz \\
&= \int \int \left| f_\epsilon^{\mathrm{ft}}(t+s) \int m_2(v) f_X(v) e^{\mathrm{i}(t+s) \cdot v} dv \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)} \frac{K^{\mathrm{ft}}(sb)}{f_\epsilon^{\mathrm{ft}}(s)} \right|^2 dt ds \\
&= \int \int \left| [m_2 f_X]^{\mathrm{ft}}(t+s) f_\epsilon^{\mathrm{ft}}(t+s) \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)} \frac{K^{\mathrm{ft}}(sb)}{f_\epsilon^{\mathrm{ft}}(s)} \right|^2 dt ds \\
&= b^{-d} \int \int \left| [m_2 f_X]^{\mathrm{ft}}(a) f_\epsilon^{\mathrm{ft}}(a) \frac{K^{\mathrm{ft}}(ab-r)}{f_\epsilon^{\mathrm{ft}}(a-r/b)} \frac{K^{\mathrm{ft}}(r)}{f_\epsilon^{\mathrm{ft}}(r/b)} \right|^2 da dr \\
&= O(b^{-d(1+4\alpha)}),
\end{aligned}
$$

where we have used Parseval's identity in the second equality, the change of variables $(t+s = a, s = r/b)$ in the penultimate equality, and for the final equality we have used Assumption D (ii) $(f_X m_2 \in L_2)$ which implies $\sup_a |[m_2 f_X]^{\mathrm{ft}}(a)|^2 < \infty$, Assumption D (iii) (compactness of K^{ft}), Assumption O (i) (rate of decay of f_ϵ^{ft}), and that a characteristic function is bounded by 1.

For Q_2, we can use similar arguments as for Q_1 to write

$$
\begin{aligned}
Q_2 &\leq \int \int \int \left| \int \int f_\epsilon^{\mathrm{ft}}(t+s) e^{-\mathrm{i}t \cdot (x-v)} e^{-\mathrm{i}s \cdot (z-v)} f_X(v) m_\theta(x) m_\theta(z) \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)} \frac{K^{\mathrm{ft}}(sb)}{f_\epsilon^{\mathrm{ft}}(s)} dt ds \right|^2 dv dx dz \\
&= b^{2d} \int \int \int \left| \int \int f_\epsilon^{\mathrm{ft}}(t+s) e^{-\mathrm{i}t \cdot ub} e^{-\mathrm{i}s \cdot rb} f_X(v) m_\theta(v+ub) m_\theta(v+rb) \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)} \frac{K^{\mathrm{ft}}(sb)}{f_\epsilon^{\mathrm{ft}}(s)} dt ds \right|^2 dv du dr \\
&= b^{-2d} \int \int \int \left| \int \int f_\epsilon^{\mathrm{ft}}\left(\frac{t+s}{b}\right) e^{-\mathrm{i}t \cdot u} e^{-\mathrm{i}s \cdot r} f_X(v) m_\theta(v+ub) m_\theta(v+rb) \frac{K^{\mathrm{ft}}(t)}{f_\epsilon^{\mathrm{ft}}(t/b)} \frac{K^{\mathrm{ft}}(s)}{f_\epsilon^{\mathrm{ft}}(s/b)} dt ds \right|^2 dv du dr \\
&\sim b^{-2d} \int \int \int \left| \int \int f_\epsilon^{\mathrm{ft}}\left(\frac{t+s}{b}\right) e^{-\mathrm{i}t \cdot u} e^{-\mathrm{i}s \cdot r} f_X(v) m_\theta(v) m_\theta(v) \frac{K^{\mathrm{ft}}(t)}{f_\epsilon^{\mathrm{ft}}(t/b)} \frac{K^{\mathrm{ft}}(s)}{f_\epsilon^{\mathrm{ft}}(s/b)} dt ds \right|^2 dv du dr \\
&= b^{-2d} \int |f_X(v) m_\theta(v) m_\theta(v)|^2 dv \int \int \left| f_\epsilon^{\mathrm{ft}}\left(\frac{t+s}{b}\right) \frac{K^{\mathrm{ft}}(t)}{f_\epsilon^{\mathrm{ft}}(t/b)} \frac{K^{\mathrm{ft}}(s)}{f_\epsilon^{\mathrm{ft}}(s/b)} \right|^2 dt ds \\
&= b^{-2d} \int |[f_X m_\theta^2]^{\mathrm{ft}}(\zeta)|^2 d\zeta \int \int \left| f_\epsilon^{\mathrm{ft}}\left(\frac{t+s}{b}\right) \frac{K^{\mathrm{ft}}(t)}{f_\epsilon^{\mathrm{ft}}(t/b)} \frac{K^{\mathrm{ft}}(s)}{f_\epsilon^{\mathrm{ft}}(s/b)} \right|^2 dt ds \\
&= O(b^{-d(1+4\alpha)}),
\end{aligned}
$$

where we have used the change of variables $(x = ub+v, z = rb+v)$ in the first equality, the wave relation follows from two Taylor expansions in each m_θ around v, Parseval's theorem is used in the third and fourth equalities, and the same arguments as used for bounding Q_1 are used in the final step.

We can similarly show $Q_3 = O(b^{-d(1+4\alpha)})$ and conclude $C_{V,b} = O(b^{-d(1+4\alpha)})$.

Second, we show that the estimation error of θ is negligible for the limiting distribution of

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T_n. Decompose $\zeta_i(x) = \xi_i(x) + \rho_i(x)$, where

$$\rho_i(x) = \int \frac{K^{\text{ft}}(tb)}{f_\epsilon^{\text{ft}}(t)} \{m_\theta(x) - m_{\hat{\theta}}(x)\} e^{-it\cdot(x-W_i)} dt.$$

Then T_n is written as

$$
\begin{aligned}
T_n &= \frac{1}{n}\sum_{i\neq j}\int \xi_i(x)\overline{\xi_j(x)}dx + \frac{1}{n}\sum_{i\neq j}\int \rho_i(x)\overline{\rho_j(x)}dx \\
&\quad + \frac{1}{n}\sum_{i\neq j}\int \rho_i(x)\overline{\xi_j(x)}dx + \frac{1}{n}\sum_{i\neq j}\int \xi_i(x)\overline{\rho_j(x)}dx \\
&\equiv \tilde{T}_n + R_{1n} + R_{2n} + R_{3n}.
\end{aligned}
\tag{3}
$$

By an expansion around $\hat{\theta} = \theta$ and Assumption O (iii) ($\hat{\theta} - \theta = o_p(1)$), the term R_{1n} satisfies

$$R_{1n} = o_p(1)\left|\frac{1}{n}\sum_{i\neq j}\int \rho_{1i}(x)\overline{\rho_{1j}(x)}dx\right|,$$

where $\rho_{1i}(x) = \int \frac{\partial m_\theta(x)}{\partial\theta}\frac{K^{\text{ft}}(tb)}{f_\epsilon^{\text{ft}}(t)}e^{-it\cdot(x-W_i)}dt$. Using similar arguments as those used to bound Q_2, together with Assumption D (ii) ($\frac{dm_\theta}{d\theta}f_X \in L_2$), we have

$$
\begin{aligned}
&E\left[\int \rho_{1i}(x)\overline{\rho_{1j}(x)}dx\right] \\
&\leq \int\int\left|\int \frac{dm_\theta(x)}{d\theta}f_X(v)K^{\text{ft}}(tb)e^{-it\cdot(x-v)}dt\right|^2 dvdx \\
&= b^{-d}\int\int\left|\int \frac{dm_\theta(v+ub)}{d\theta}f_X(v)K^{\text{ft}}(t)e^{-it\cdot u}dt\right|^2 dvdu \\
&\sim b^{-d}\int\int\left|\int \frac{dm_\theta(v)}{d\theta}f_X(v)K^{\text{ft}}(t)e^{-it\cdot u}dt\right|^2 dvdu \\
&= b^{-d}\int\left|\left[\frac{dm_\theta}{d\theta}f_X\right]^{\text{ft}}(\zeta)\right|^2 d\zeta\int|K^{\text{ft}}(t)|^2 dt = \int\left|\left[\frac{dm_\theta}{d\theta}f_X\right]^{\text{ft}}(\zeta)\right|^2 d\zeta\int|K^{\text{ft}}(tb)|^2 dt \\
&= O(1).
\end{aligned}
$$

Also, by applying the same argument as used to derive the bound of $C_{V,b}$, we have

$$E\left[\left(\int \rho_{1i}(x)\overline{\rho_{1j}(x)}dx\right)^2\right] = O(b^{-d(1+4\alpha)}).$$

Combining these two results, along with $\hat{\theta} - \theta = o_p(1)$ and $C_{V,b} = O(b^{-d(1+4\alpha)})$, we obtain $C_{V,b}^{-1/2}R_{1n} = o_p(1)$. In the same manner we can show $C_{V,b}^{-1/2}R_{2n} = o_p(1)$ and $C_{V,b}^{-1/2}R_{3n} = o_p(1)$

and thus $C_{V,b}^{-1/2}T_n = C_{V,b}^{-1/2}\tilde{T}_n + o_p(1)$.

Thirdly, we derive the limiting distribution of $C_{V,b}^{-1/2}\tilde{T}_n$. Note that $\tilde{T}_n$ is written as $\tilde{T}_n = \frac{2}{n(2\pi)^d}\sum_{i<j}H_{i,j}$ and is a U-statistic with zero mean under H_0. To prove the asymptotic normality of $\tilde{T}_n$, we apply the central limit theorem in Hall (1984, Theorem 1). To this end, it is enough to show

$$\frac{E[H_{1,2}^4]}{n(E[H_{1,2}^2])^2} \to 0, \quad \text{and} \quad \frac{E[G_{1,2}^2]}{(E[H_{1,2}^2])^2} \to 0, \tag{4}$$

where $G_{i,j} = E[H_{1,i}H_{1,j}|Y_1, W_1]$. Recall that $C_{V,b} = E[H_{1,2}^2]$ satisfies $C_{V,b} = O(b^{-d(1+4\alpha)})$. By a similar argument used to bound $E[H_{1,2}^2]$, we can show

$$E[H_{1,2}^4] = E\left[\int \cdots \int \prod_{k=1}^4 \xi_1(x_k)\overline{\xi_2(x_k)}dx_1\cdots dx_4\right] = O(b^{-3d(1+8\alpha)}).$$

For $E[G_{1,2}^2]$, we can equivalently look at (see, p. 5 of Hall, 1984)

$$\begin{aligned}
&E[H_{1,3}H_{1,4}H_{2,3}H_{2,4}] \\
=\ &\int \cdots \int \prod_{k=1}^4 \xi_1(x_1)\overline{\xi_3(x_1)}\xi_1(x_2)\overline{\xi_4(x_2)}\xi_2(x_3)\overline{\xi_3(x_3)}\xi_2(x_4)\overline{\xi_4(x_4)}dx_1\cdots dx_4 \\
=\ &O(b^{-d(1+8\alpha)}).
\end{aligned}$$

These results combined with Assumption D (iv) guarantee the conditions in (4). Thus, Hall (1984, Theorem 1) implies

$$C_{V,b}^{-1/2}T_n \xrightarrow{d} N\left(0, \frac{4}{(2\pi)^{2d}}\right).$$

Finally, the B_n term is simply a sample mean of independent and identically distributed random variables. The second moment can be calculated in the same manner as T_n to give $C_{V,b}^{-1/2}B_n = O_p(n^{-1})$. Combining these results, the conclusion follows.

A.1.2   Proof of (ii)

A similar argument to the proof of Part (i) guarantees $\varphi(b)T_n = \varphi(b)\tilde{T}_n + o_p(1)$, where $\varphi(b) = \frac{(2\pi)^d}{b^{d(2+4\beta)}e^{d/b^2}A^{2d}\Gamma(1+2d\beta)}$. Thus, we hereafter derive the limiting distribution of $\tilde{T}_n$. Decompose

$\tilde{T}_n = \bar{T}_n + r_{1n} - 2r_{2n}$, where

$$\bar{T}_n = \frac{1}{n}\sum_{i\neq j}\int\int\int \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)}\frac{K^{\mathrm{ft}}(sb)}{f_\epsilon^{\mathrm{ft}}(s)}Y_iY_j e^{-\mathrm{i}t\cdot(x-W_i)}e^{-\mathrm{i}s\cdot(x-W_j)}dtdsdx \qquad (5)$$

$$r_{1n} = \frac{1}{n}\sum_{i\neq j}\int\int\int \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)}\frac{K^{\mathrm{ft}}(sb)}{f_\epsilon^{\mathrm{ft}}(s)}m_\theta(x)^2 e^{-\mathrm{i}t\cdot(x-W_i)}e^{-\mathrm{i}s\cdot(x-W_j)}dtdsdx$$

$$r_{2n} = \frac{1}{n}\sum_{i\neq j}\int\int\int \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)}\frac{K^{\mathrm{ft}}(sb)}{f_\epsilon^{\mathrm{ft}}(s)}Y_i m_\theta(x) e^{-\mathrm{i}t\cdot(x-W_i)}e^{-\mathrm{i}s\cdot(x-W_j)}dtdsdx$$

First, we derive the limiting distribution of $\bar{T}_n$. Observe that

$$\bar{T}_n = \frac{1}{n}\sum_{i\neq j}\int\left(\int \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)}Y_i e^{-\mathrm{i}t\cdot W_i}dt\right)\left(\int \frac{K^{\mathrm{ft}}(sb)}{f_\epsilon^{\mathrm{ft}}(s)}Y_j e^{-\mathrm{i}s\cdot W_j}ds\right)$$

$$= \frac{1}{n}\sum_{i\neq j}\frac{1}{(2\pi)^{2d}}\int \frac{|K^{\mathrm{ft}}(tb)|}{f_\epsilon^{ft}(t)}\frac{|K^{\mathrm{ft}}(sb)|}{f_\epsilon^{ft}(s)}Y_iY_j\{\cos(tW_i)\cos(sW_j)+\sin(tW_i)\sin(sW_j)\}dtds$$

$$= \frac{1}{(2\pi)^{2d}}\left|\int \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{ft}(t)}dt\right|^2 \frac{1}{n}\sum_{i\neq j}Y_iY_j\left\{\cos\left(\frac{W_i}{b}\right)\cos\left(\frac{W_j}{b}\right)+\sin\left(\frac{W_i}{b}\right)\sin\left(\frac{W_j}{b}\right)\right\}$$

$$\quad + O_p\left(b^{d(2+4\beta)}e^{d/b^2}\right)$$

$$\equiv \left(\frac{1}{(2\pi)^{2d}}\left|\int \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{ft}(t)}dt\right|^2\right)\dot{T}_n + O_p\left(b^{d(2+4\beta)}e^{d/b^2}\right)$$

where the first equality follows from Plancheral's theorem, the second equality follows from $f_\epsilon^{\mathrm{ft}}(t) = e^{-dt^2/2}$ and $e^{\mathrm{i}tW_i} = \cos(tW_i) + \mathrm{i}\sin(tW_i)$, the fourth equality follows from the change of variables $(t = v/b, s = t + r)$ and from a simple multivariate extension of Holzmann and Boysen (2006, Theorem 1) based on Assumption S (ii). Note that

$$\dot{T}_n = \frac{1}{n}\sum_{i\neq j}Y_iY_j\left[\begin{array}{l}\left\{\cos\left(\frac{X_i}{b}\right)\cos\left(\frac{\epsilon_i}{b}\right)-\sin\left(\frac{X_i}{b}\right)\sin\left(\frac{\epsilon_i}{b}\right)\right\}\left\{\cos\left(\frac{X_j}{b}\right)\cos\left(\frac{\epsilon_j}{b}\right)-\sin\left(\frac{X_j}{b}\right)\sin\left(\frac{\epsilon_j}{b}\right)\right\}\\ +\left\{\sin\left(\frac{X_i}{b}\right)\cos\left(\frac{\epsilon_i}{b}\right)+\cos\left(\frac{X_i}{b}\right)\sin\left(\frac{\epsilon_i}{b}\right)\right\}\left\{\sin\left(\frac{X_j}{b}\right)\cos\left(\frac{\epsilon_j}{b}\right)+\cos\left(\frac{X_j}{b}\right)\sin\left(\frac{\epsilon_j}{b}\right)\right\}\end{array}\right].$$

From van Es and Uh (2005, proof of Lemma 6), it holds $\left(\frac{X_i}{b}\bmod 2\pi\right)\overset{d}{\to} V_i^X \sim U[0, 2\pi]$ and $\left(\frac{\epsilon_i}{b}\bmod 2\pi\right)\overset{d}{\to} V_i^\epsilon \sim U[0, 2\pi]$ as $b\to 0$ for each i, where V_i^ϵ is independent from (Y_i, V_i^X). Thus by applying Holzmann and Boysen (2006, Lemma 1), $\dot{T}_n$ has the same limiting distribution as $\tilde{T}_n^V = \frac{1}{n}\sum_{i\neq j}h(Q_i, Q_j)$, where $Q_i = (Y_i, V_i^X, V_i^\epsilon)$ and

$$h(Q_i, Q_j) = Y_iY_j\left[\begin{array}{l}\left\{\cos(V_i^X)\cos(V_i^\epsilon)-\sin(V_i^X)\sin(V_i^\epsilon)\right\}\left\{\cos(V_j^X)\cos(V_j^\epsilon)-\sin(V_j^X)\sin(V_j^\epsilon)\right\}\\ +\left\{\sin(V_i^X)\cos(V_i^\epsilon)+\cos(V_i^X)\sin(V_i^\epsilon)\right\}\left\{\sin(V_j^X)\cos(V_j^\epsilon)+\cos(V_j^X)\sin(V_j^\epsilon)\right\}\end{array}\right].$$

Observe that $\mathrm{Cov}\,(h(Q_1, Q_2), h(Q_1, Q_3)) = 0$ because $E[\cos(V_i^\epsilon)] = E[\sin(V_i^\epsilon)] = 0$. Therefore, by applying the limit theorem for degenerate U-statistics with a fixed kernel h (Serfling, 1980, Theorem 5.5.2), we obtain

$$\tilde{T}_n^V \xrightarrow{d} \sum_{k=1}^\infty \lambda_k(Z_k^2 - 1), \qquad (6)$$

where $\{Z_k\}$ is an independent sequence of standard normal random variables and $\{\lambda_k\}$ are the eigenvalues of the integral operator

$$(\Lambda g)(Q_1) = \lambda g(Q_1).$$

where $(\Lambda g)(Q_1) = E[h(Q_1, Q_2)g(Q_2)|Q_1]$. Also, we have

$$\left| \int \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)} dt \right|^2 \leq \int \left| \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)} \right|^2 dt = \frac{1}{b^d} \int \left| \frac{K^{\mathrm{ft}}(t)}{f_\epsilon^{\mathrm{ft}}(t/b)} \right|^2 dt$$

$$= \frac{1}{b^d} \int \left| K^{\mathrm{ft}}(t) e^{\left(\frac{t}{2b}\right)^2} \right|^2 dt \sim \frac{1}{\varphi(b)},$$

where the inequality follows from Cauchy's inequality, the second step follows from a simple change of variables, the third step follows from the definition of the characteristic function of the Gaussian distribution, and the final wave relation comes from van Es and Uh (2005, Lemma 5). Combining these results gives

$$\varphi(b)\bar{T}_n \xrightarrow{d} \sum_{k=1}^\infty \lambda_k(Z_k^2 - 1).$$

Next, we show that r_{1n} is negligible. Notice that $E[r_{1n}^2]$ is equal to Q_2 from the proof of Theorem 1 (i), hence we can write

$$
\begin{aligned}
E[r_{1n}^2] &= b^{-2d} \int |[f_X m_\theta^2]^{\mathrm{ft}}(\zeta)|^2 d\zeta \int \int \left| f_\epsilon^{\mathrm{ft}}\left(\frac{t+s}{b}\right) \frac{K^{\mathrm{ft}}(t)}{f_\epsilon^{\mathrm{ft}}(t/b)} \frac{K^{\mathrm{ft}}(s)}{f_\epsilon^{\mathrm{ft}}(s/b)} \right|^2 dt ds \\
&= b^{-d} \int |[f_X m_\theta^2]^{\mathrm{ft}}(\zeta)|^2 d\zeta \int \int \left| f_\epsilon^{\mathrm{ft}}(v) \frac{K^{\mathrm{ft}}(vb-s)}{f_\epsilon^{\mathrm{ft}}(v-s/b)} \frac{K^{\mathrm{ft}}(s)}{f_\epsilon^{\mathrm{ft}}(s/b)} \right|^2 dv ds \\
&\sim b^{-d} \int |[f_X m_\theta^2]^{\mathrm{ft}}(\zeta)|^2 d\zeta \int |f_\epsilon^{\mathrm{ft}}(v)|^2 dv \int \left| \frac{K^{\mathrm{ft}}(s)}{f_\epsilon^{\mathrm{ft}}(s/b)} \right|^4 ds \\
&= O(b^{-d}) \int \left| \frac{K^{\mathrm{ft}}(s)}{f_\epsilon^{\mathrm{ft}}(s/b)} \right|^4 ds \sim \frac{1}{\varphi(2b)},
\end{aligned}
$$

where the second equality follows from a change of variables $t = vb - s$, the penultimate equality uses Assumption D (ii), and the final relation uses van Es and Uh (2005, Lemma 5) again. The first moment of r_{1n} is bounded in a similar manner.

By the same argument used to show the asymptotic behaviour of $\dot{T}_n$, it holds

$$\frac{1}{n}\sum_{i\neq j}\left\{\cos\left(\frac{W_i}{b}\right)\cos\left(\frac{W_j}{b}\right)+\sin\left(\frac{W_i}{b}\right)\sin\left(\frac{W_j}{b}\right)\right\}=O_p(1).$$

Combining these results, we have

$$\varphi(b)r_{1n}=O_p(b^{d(1+2\beta)}),$$

and thus r_{1n} is negligible. Similar arguments imply that the terms r_{2n} and r_{3n} are also asymptotically negligible, and the conclusion follows.

A.1.3   Proof of (iii)

The proof for the general supersmooth case follows the same steps as in the proof of Part (ii) for the normal case. As the proof is similar, only essential details are provided in what follows. Hereafter we show why the condition $\gamma > 1$ is imposed in this case. The dominant term $\bar{T}_n$ defined in (5) satisfies

$$\bar{T}_n \sim \frac{b^{2d(\gamma_0-1)}}{(2\pi)^d C_\epsilon^{2d}}\int |K^{\mathrm{ft}}(t)||K^{\mathrm{ft}}(s)||t|^{-d\gamma_0}|s|^{-d\gamma_0}e^{\frac{d}{\mu b^\gamma}(|t|^\gamma+|s|^\gamma)}$$
$$\times\frac{1}{n}\sum_{i\neq j}Y_iY_j\left\{\cos\left(\frac{tW_i}{b}\right)\cos\left(\frac{sW_j}{b}\right)+\sin\left(\frac{tW_i}{b}\right)\sin\left(\frac{sW_j}{b}\right)\right\}dtds$$

We now show that

$$D_{\cos} \equiv \frac{b^{2d(\gamma_0-1)}}{(2\pi)^d C_\epsilon^{2d}}\int |K^{\mathrm{ft}}(t)||K^{\mathrm{ft}}(s)||t|^{-d\gamma_0}|s|^{-d\gamma_0}e^{\frac{d}{\mu b^\gamma}(|t|^\gamma+|s|^\gamma)}$$
$$\times\frac{1}{n}\sum_{i\neq j}Y_iY_j\left\{\cos\left(\frac{tW_i}{b}\right)\cos\left(\frac{sW_j}{b}\right)-\cos\left(\frac{W_i}{b}\right)\cos\left(\frac{W_j}{b}\right)\right\}dtds$$

is asymptotically negligible, as well as the correspondingly defined $D_{\sin}$. We have seen in the proof of part (ii) that each term is zero mean. Following the proof of Holzmann and Boysen (2006, Theorem 1), we obtain

$$\left|\cos\left(\frac{tW_i}{b}\right)\cos\left(\frac{sW_j}{b}\right)-\cos\left(\frac{W_i}{b}\right)\cos\left(\frac{W_j}{b}\right)\right|\leq\frac{(1-t)|W_i|}{b}+\frac{(1-s)|W_k|}{b}.$$

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Similar arguments to van Es and Uh (2005, Lemmas 1 and 5) and Assumption S (ii), imply

$$
\begin{aligned}
\mathrm{Var}(D_{\cos}) &= O(n^{-2}b^{4d\gamma_0-6d})\sum_{i\neq j}E\left[|Y_i||Y_j|\int|K^{\mathrm{ft}}(t)||K^{\mathrm{ft}}(s)||t|^{-d\gamma_0}|s|^{-d\gamma_0}e^{\frac{d}{\mu b^\gamma}(|t|^\gamma+|s|^\gamma)}\right.\\
&\qquad\left.\times\{(1-t)|W_i|+(1-s)|W_k|\}\,dtds\right]^2\\
&= O(b^{4d\gamma_0-6d})\left[\int(1-t)|K^{\mathrm{ft}}(t)||K^{\mathrm{ft}}(s)||t|^{-d\gamma_0}|s|^{-d\gamma_0}e^{\frac{d}{\mu b^\gamma}(|t|^\gamma+|s|^\gamma)}dtds\right]^2\\
&= O(b^{4d(\gamma_0-1)})\left[\int(1-t)|K^{\mathrm{ft}}(t)||K^{\mathrm{ft}}(t+vb)||t|^{-d\gamma_0}|t+vb|^{-d\gamma_0}e^{\frac{d}{\mu b^\gamma}(|t|^\gamma+|t+vb|^\gamma)}dtdv\right]^2\\
&\sim O(b^{4d(\gamma_0-1)})\left[\int(1-t)|K^{\mathrm{ft}}(t)|^2|t|^{-2d\gamma_0}e^{\frac{2d|t|^\gamma}{\mu b^\gamma}}dt\right]^2\\
&= O(b^{4d(\gamma_0-1)})\left(b^{2\gamma}\int v|K^{\mathrm{ft}}(1-b^\gamma v)|^2|1-b^\gamma v|^{-2\gamma_0}e^{\frac{2|1-b^\gamma v|^\gamma}{\mu b^\gamma}}dv\right)^2\\
&= O\left(b^{4d(\gamma_0-1)}\left(b^{\gamma(2+2\beta)}e^{\frac{2}{\mu b^\gamma}}\right)^2\right),
\end{aligned}
$$

where the first two steps follow from standard arguments, the third step follows from the change of variable $s = vb + t$, the fourth step results from the definition of the wave relation, the fifth step uses the change of variable $t = 1 - b^\gamma v$, and the final step uses van Es and Uh (2005, Lemma 5). Using similar arguments as in the proof of Theorem 1 (ii), we obtain $D_{\cos} = O_p\left(b^{2(\gamma-1)+2\gamma\beta+2\gamma_0}e^{\frac{2}{\mu b^\gamma}}\right)$. The same argument applies to $D_{\sin}$. Note that

$$
\begin{aligned}
\bar{T}_n &= \frac{b^{2d(\gamma_0-1)}}{(2\pi)^dC_\epsilon^{2d}}\int|K^{\mathrm{ft}}(t)||K^{\mathrm{ft}}(s)||t|^{-d\gamma_0}|s|^{-d\gamma_0}e^{\frac{d}{\mu b^\gamma}(|t|^\gamma+|s|^\gamma)}\\
&\qquad\times\frac{1}{n}\sum_{i\neq j}Y_iY_j\left\{\cos\left(\frac{W_i}{b}\right)\cos\left(\frac{W_j}{b}\right)+\sin\left(\frac{W_i}{b}\right)\sin\left(\frac{W_j}{b}\right)\right\}dtds\\
&\qquad+O\left(b^{d\{2(\gamma-1)+2\gamma\beta+2\gamma_0\}}e^{\frac{2d}{\mu b^\gamma}}\right)\\
&= \frac{A^{2d}\mu^{1+2d\beta}b^{d(\gamma-1+2\gamma\beta+2\gamma_0)}e^{\frac{2d}{\mu b^\gamma}}\Gamma(1+2d\beta)}{\lambda^{1+2d\beta}(2\pi)^dC_\epsilon^{2d}}\tilde{T}_n+O\left(b^{d\{2(\gamma-1)+2\gamma\beta+2\gamma_0\}}e^{\frac{2d}{\mu b^\gamma}}\right)\\
&\equiv \frac{\tilde{T}_n}{\varphi(b)}+O\left(b^{d\{2(\gamma-1)+2\gamma\beta+2\gamma_0\}}e^{\frac{2d}{\mu b^\gamma}}\right),
\end{aligned}
$$

where the second equality follows from the definition of $\tilde{T}_n$ in (3) and a modification of van Es and Uh (2005, Lemma 5). Therefore, we obtain

$$
\varphi(b)T_n = \tilde{T}_n + O(b^{d(\gamma-1)}).
$$

The limiting distribution of $\tilde{T}_n$ - and hence D_n - is obtained in the proof of Theorem 1 (ii). The remainder term becomes negligible if we impose $\gamma > 1$.

A.2 Proof of Theorem 2

A.2.1 Proof of (i)

By a similar argument to the proof of Theorem 1 (i), the estimation error $\hat{\theta} - \theta$ is negligible for the asymptotic properties of T_n and thus it is written as

$$
\begin{aligned}
T_n &= \frac{1}{n}\sum_{i\neq j}\frac{1}{(2\pi)^d}\int \xi_i(x)\overline{\xi_j(x)}dx + \frac{1}{n}\sum_{i\neq j}\frac{1}{(2\pi)^d}\int \eta_i(x)\overline{\eta_j(x)}dx \\
&\quad + \frac{1}{n}\sum_{i\neq j}\frac{1}{(2\pi)^d}\int \xi_i(x)\overline{\eta_j(x)}dx + \frac{1}{n}\sum_{i\neq j}\frac{1}{(2\pi)^d}\int \eta_i(x)\overline{\xi_j(x)}dx + o_p(C_{V,b}^{1/2}) \\
&\equiv \tilde{T}_n + R_{1n}^* + R_{2n}^* + R_{3n}^* + o_p(C_{V,b}^{1/2}),
\end{aligned}
$$

where

$$
\begin{aligned}
\eta_i(x) &= \int \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)}\{m(x) - m_\theta(x)\}e^{-\mathrm{i}t\cdot(x-W_i)}dt \\
&= c_n\int \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)}\Delta(x)e^{-\mathrm{i}t\cdot(x-W_i)}dt,
\end{aligned}
$$

under H_{1n}. By Theorem 1 (i), we have seen $C_{V,b}^{-1/2}\tilde{T}_n \xrightarrow{d} N\left(0, \frac{2}{(2\pi)^{2d}}\right)$. For R_{1n}^*, observe that $\eta_i(x)$ takes the same form as $\rho_{1i}(x)$ in the proof of 1 (i), but with $\frac{\partial m_\theta(x)}{\partial\theta}$ replaced by $\Delta(x)$ and multiplied by c_n. As such, we can reuse the same arguments as used in that proof. First,

$$
\begin{aligned}
E[C_{V,b}^{-1/2}R_{1n}^*] &= E\left[\frac{c_n^2}{(2\pi)^d C_{V,b}^{1/2}}\frac{1}{n}\sum_{i\neq j}\int\int \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)}\Delta(x)e^{-\mathrm{i}t\cdot(x-W_i)}dt\int \frac{K^{\mathrm{ft}}(sb)}{f_\epsilon^{\mathrm{ft}}(s)}\Delta(x)e^{-\mathrm{i}s\cdot(x-W_j)}ds\,dx\right] \\
&= E\left[\int \frac{c_n^2}{(2\pi)^d C_{V,b}^{1/2}}\frac{1}{n}\sum_{i\neq j}\int \frac{K^{\mathrm{ft}}(tb)}{f_\epsilon^{\mathrm{ft}}(t)}e^{-\mathrm{i}t\cdot W_i}dt\int \frac{K^{\mathrm{ft}}(sb)}{f_\epsilon^{\mathrm{ft}}(s)}\Delta(x)e^{-\mathrm{i}s\cdot W_j}ds\right] \\
&= \int \Delta^2(x)dx\frac{(n-1)c_n^2}{(2\pi)^d C_{V,b}^{1/2}}\left|\int K^{\mathrm{ft}}(tb)f_X^{ft}(t)dt\right|^2 \equiv \Delta_n.
\end{aligned}
\tag{7}
$$

By the definition of c_n, the fact that $C_{V,b} = O(b^{-d(1+4\alpha)})$, and Assumption D (ii), it holds that $E[C_{V,b}^{-1/2}R_{1n}^*] = O(1)$ and the limit of Δ_n exists. Again, using a similar argument used to bound the variance of $C_{V,b}$, we obtain

$$
E[R_{1n}^{*2}] = O(b^{-d(1+4\alpha)})c_n^4 = O(n^{-2}b^{-2d(1+4\alpha)})
$$

Therefore, $\mathrm{Var}(C_{V,b}^{-1/2}R_{1n}^*) \to 0$ and we obtain $C_{V,b}^{-1/2}R_{1n}^* \xrightarrow{p} \lim_{n\to\infty}\Delta_n$. Finally, using similar arguments, combined with $E[\xi_i(t)] = 0$, we can show that $C_{V,b}^{-1/2}R_{2n}^* \xrightarrow{p} 0$ and $C_{V,b}^{-1/2}R_{3n}^* \xrightarrow{p} 0$.

Taken together, the conclusion follows.

A.2.2 Proof of (ii)

Similar to the proof of Theorem 2 part (i), we can decompose

$$T_n \;\;=\;\; \tilde{T}_n + R_{1n}^* + R_{2n}^* + R_{3n}^* + o_p(\varphi(b)^{-1}).$$

Theorem 1 (ii) gives the limiting distribution of $\varphi(b)\tilde{T}_n$. For R_{1n}^*, we have the same result as for part (i) and the limit of Δ_n exists from the definition of c_n and the bound of $C_{V,b}$ in the supersmooth case. The conclusion follows.

A.2.3 Proof of (iii)

The proof is identical to that of Part (ii) after setting $c_n = b^{d\{(\lambda-1)/2 + \lambda\beta + \lambda_0\}} e^{\frac{d}{\mu b^\lambda}}$. Each of the objects in the asymptotic distribution is defined analogously to part (ii).

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