

Enhancing Accuracy of FV3 finite-volume operators

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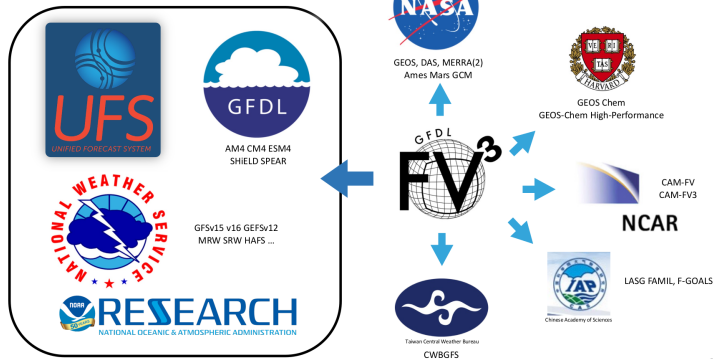
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Numerical methods for GFD seminar

October 6, 2023

The Global FV3 Community Past, present, future earth and beyond



6

www.gfdl.noaa.gov/fv3
github.com/NOAA-GFDL/GFDL_atmos_cubed_sphere

- To discuss strategies to reduce larger errors at the edges (grid imprinting).
- Introduce the duogrid to compute stencils near to the edges.
- Improve the accuracy of the discrete operators from FV3 dynamical core.

The gnomonic equiangular cubed-sphere from Ronchi et al. (1996)

Six maps from a cube panel p to the sphere $\Phi_p : [-\alpha, \alpha]^2 \rightarrow \mathbb{S}_R^2$,
 $p = 1, \dots, 6$, $\alpha = \frac{\pi}{4}$.

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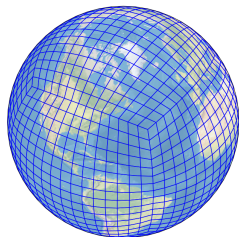
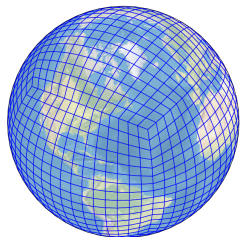


Figure: CS grid with $N = 15$ cells along a coordinate axis

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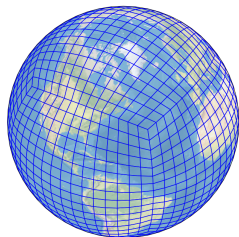


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- $\sqrt{g}(x, y) = \sqrt{\det(D\Phi_p^T D\Phi_p)}$;
- Point on the sphere is represented by (x, y, p) ;
- Gridpoints use indexes (i, j, p) ;
- Edge values are represented using half integers.

Linear advection equation on the cubed-sphere

$$\begin{cases} \partial_t(\sqrt{g}q) + \partial_x(u\sqrt{g}q) + \partial_y(v\sqrt{g}q) = 0, \\ q_0(x, y, p) = q(x, y, p, 0), \end{cases}$$

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- ① Metric tensor term \sqrt{g} ;
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- ③ Coordinate system discontinuity at the cube edges.

Dimension splitting scheme

Lie-Trotter splitting: firstly we solve

$$[\partial_t q^x + \partial_x(uq^x)](x, y_j, p, t) = 0,$$

and then solve

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$$\begin{aligned} Q^{n+1} &= \frac{(Q^{xy} + Q^{yx})}{2}, \quad \text{or} \\ Q^{n+1} &= Q^n - \Delta t \mathbb{D}_{ijp}^n, \end{aligned}$$

where

$$\begin{aligned} \mathbb{D}_{ijp}^n &= \frac{-1}{\Delta t \sqrt{g_{ijp}}} \left[\mathbf{F} \left(Q^n + \frac{1}{2} \mathbf{G}(Q^n, \tilde{v}^n), \tilde{u}^n \right) \right. \\ &\quad \left. + \mathbf{G} \left(Q^n + \frac{1}{2} \mathbf{F}(Q^n, \tilde{u}^n), \tilde{v}^n \right) \right]. \end{aligned}$$

is the discrete divergence.

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- 3 Integrate the parabolas to obtain the fluxes.

Main references: Carpenter et al. (1990) and Colella and Woodward (1984)

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This approach is referred to as **MT-PL07**, which ensures that $\mathbf{F}(Q^n, \tilde{u}) = -\bar{q} \frac{\Delta t}{\Delta x} \delta_i (\sqrt{g} \tilde{u})_{ijp}^n$ when $Q^n = \bar{q}$ is constant.

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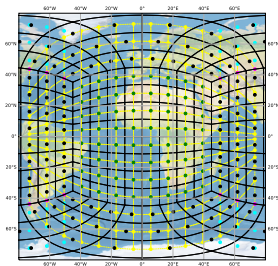
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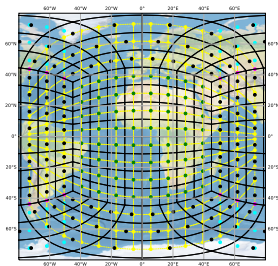
- This first-order error is eliminated by reconstructing $\sqrt{g}q$ (**MT-0**), but allows splitting error.

Duogrid and mass conservation

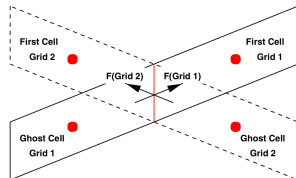


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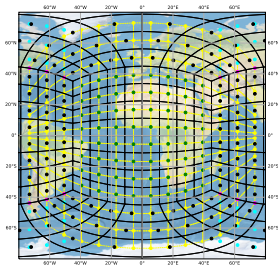


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The flux being computed twice on the cube edge breaks the total mass conservation. Figure taken from Rossmanith (2006).

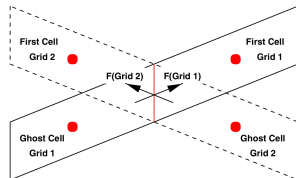
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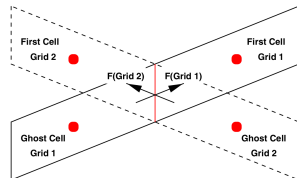
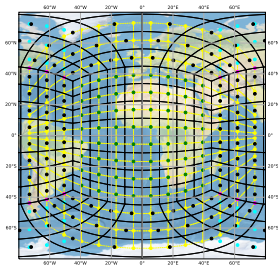
Mass conservation may be achieved by:

- 1 Average the flux (**AF**) at the cube interfaces (it can be shown that this approach reduces the \mathbb{D}_{ijp}^n accuracy by one only at the cube edges);



The flux being computed twice on the cube edge breaks the total mass conservation. Figure taken from Rossmanith (2006).

Duogrid and mass conservation



The flux being computed twice on the cube edge breaks the total mass conservation. Figure taken from Rossmanith (2006).

Duogrid (named by Chen (2021))

Mass conservation may be achieved by:

- 1 Average the flux (**AF**) at the cube interfaces (it can be shown that this approach reduces the \mathbb{D}_{ijp}^n accuracy by one only at the cube edges);
- 2 Compute the L_2 projection (**PR**) of \mathbb{D}_{ijp}^n onto $V_0 = \{\delta \in \mathcal{CS}_N : \sum_{p=1}^6 \sum_{i,j=1}^N \delta_{ijp} \sqrt{g_{ijp}} \Delta x \Delta y = 0\}.$

Numerical results

Scheme	Splitting	DP	ET	Mass fixer	Metric tensor
PL07	PL07	RK1	PL07	none	MT-PL07
PL07-PR	PL07	RK1	Duogrid	PR	MT-PL07
AVLT-AF	AVLT	RK2	Duogrid	AF	MT-0
AVLT-PR	AVLT	RK2	Duogrid	PR	MT-0

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Steady test case from Williamson et al. (1992).

$$\begin{cases} u_\lambda(\lambda, \phi) &= u_0(\cos(\phi)\cos(\alpha) + \sin(\phi)\cos(\lambda)\sin(\alpha)), \\ v_\phi(\lambda, \phi) &= -u_0\sin(\lambda)\sin(\alpha), \end{cases}$$

where $u_0 = \frac{2\pi}{5}$, $\alpha = -\frac{45\pi}{180}$.

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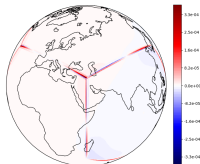
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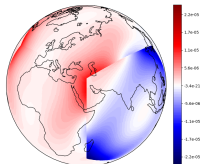
where $u_0 = \frac{2\pi}{5}$, $\alpha = -\frac{45\pi}{180}$. The advected scalar field is defined in such way that $\nabla \cdot (\mathbf{u}q) = 0$.

$\nabla \cdot (\mathbf{u}q)$ computation

$N = 128$, CFL = $6.94e-01$, $\kappa = 6$, $\nu f = 5$, $sp = p07$
recon = hyppen, dp = rk3, rnt = p07, rnf = none, et = p07, mass variation = $1.56e-17$

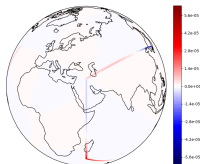


$N = 128$, CFL = $6.94e-01$, $\kappa = 6$, $\nu f = 5$, $sp = p07$
recon = hyppen, dp = rk3, rnt = p07, rnf = gpr, et = diagrid, mass variation = $5.05e-20$



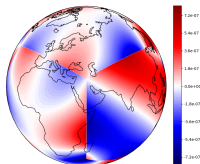
PL07

$N = 128$, CFL = $6.94e-01$, $\kappa = 6$, $\nu f = 5$, $sp = avt$
recon = hyppen, dp = rk2, rnt = rnt0, rnf = af, et = diagrid, mass variation = $-0.03e-21$



PL07-PR

$N = 128$, CFL = $6.94e-01$, $\kappa = 6$, $\nu f = 5$, $sp = avt$
recon = hyppen, dp = rk2, rnt = rnt0, rnf = gpr, et = diagrid, mass variation = $1.34e-21$

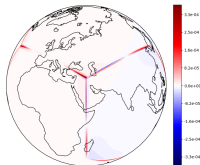


AVLT-AF

AVLT-PR

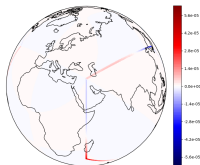
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recon = hyppen, dp = rk3, rnt = p07, rnf = none, rnt = p07, mass variation = $1.56e-17$



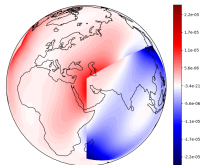
PL07

$N = 128$, CFL = $6.94e-01$, $\kappa = 6$, $\nu f = 5$, $sp = avlt$
recon = hyppen, dp = rk2, rnt = rnt0, rnf = avlt, rnt = diagrid, mass variation = $-0.83e-21$



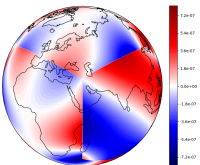
AVLT-AF

$N = 128$, CFL = $6.94e-01$, $\kappa = 6$, $\nu f = 5$, $sp = p07$
recon = hyppen, dp = rk3, rnt = p07, rnf = gpr, rnt = diagrid, mass variation = $5.05e-20$

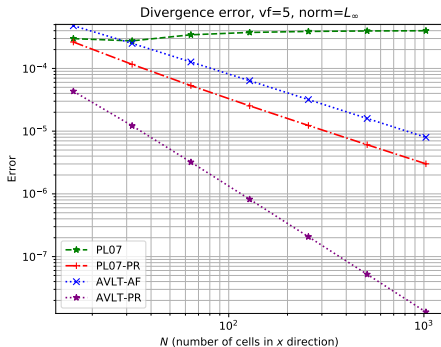


PL07-PR

$N = 128$, CFL = $6.94e-01$, $\kappa = 6$, $\nu f = 5$, $sp = avlt$
recon = hyppen, dp = rk2, rnt = rnt0, rnf = gpr, rnt = diagrid, mass variation = $1.34e-21$

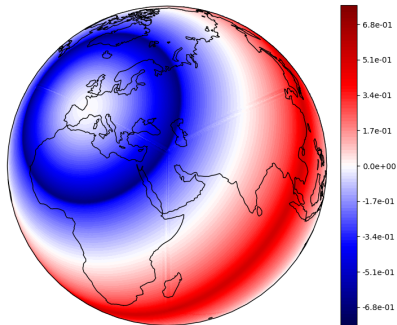


AVLT-PR



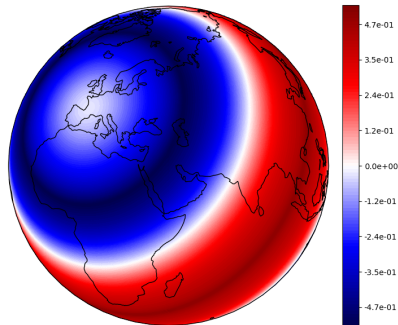
Steady test case error

Error - N = 512, Time = 5.00e+00, CFL = 6.98e-01, ic = 6, vf = 5
sp = pl07, recon = hyppm, dp = rk1, mt = pl07, mf = none, et = pl07



PL07

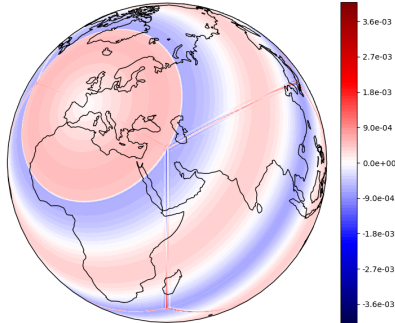
Error - N = 512, Time = 5.00e+00, CFL = 6.98e-01, ic = 6, vf = 5
sp = pl07, recon = hyppm, dp = rk1, mt = pl07, mf = gpr, et = duogrid



PL07-PR

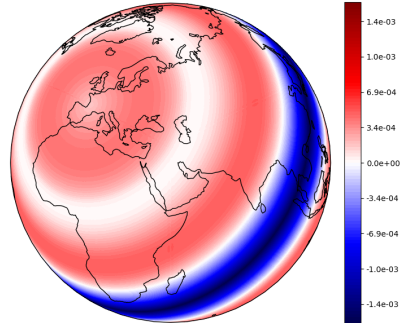
Steady test case error

Error - N = 512, Time = 5.00e+00, CFL = 6.98e-01, ic = 6, vf = 5
sp = avlt, recon = hyppm, dp = rk2, mt = mt0, mf = af, et = duogrid



AVLT-AF

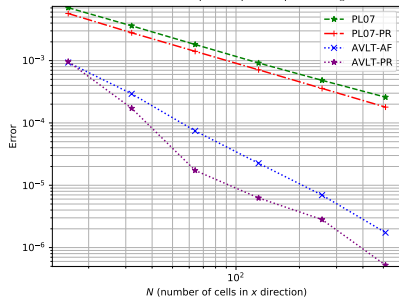
Error - N = 512, Time = 5.00e+00, CFL = 6.98e-01, ic = 6, vf = 5
sp = avlt, recon = hyppm, dp = rk2, mt = mt0, mf = gpr, et = duogrid



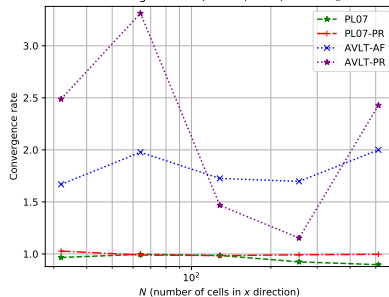
AVLT-PR

Steady test case error

Advection error, $ic = 6$, $vf = 5$, $norm=L_\infty$

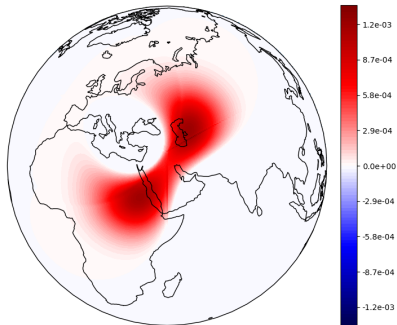


Convergence rate, $ic = 6$, $vf=5$, $norm=L_\infty$



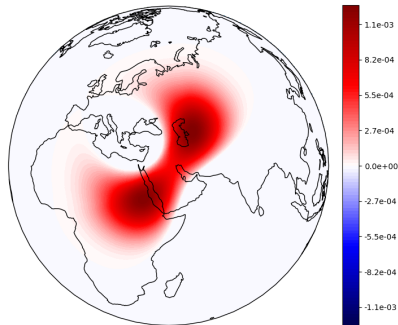
One Gaussian hill advection error

Error - N = 512, Time = 5.00e+00, CFL = 4.52e-01, ic = 2, vf = 1
sp = pl07, recon = hyppm, dp = rk1, mt = pl07, mf = none, et = pl07



PL07

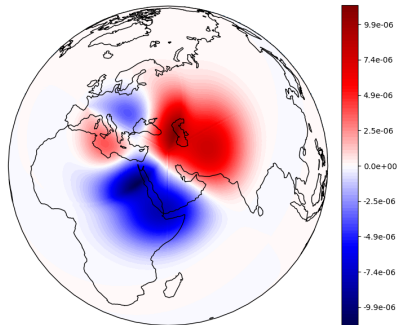
Error - N = 512, Time = 5.00e+00, CFL = 4.52e-01, ic = 2, vf = 1
sp = pl07, recon = hyppm, dp = rk1, mt = pl07, mf = gpr, et = duogrid



PL07-PR

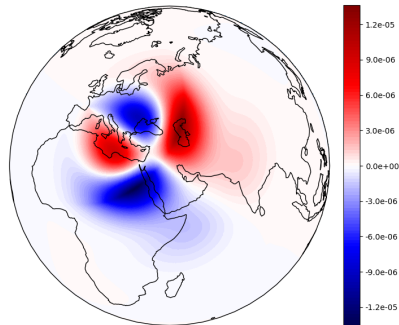
One Gaussian hill advection error

Error - N = 512, Time = 5.00e+00, CFL = 4.52e-01, ic = 2, vf = 1
sp = avlt, recon = hyppm, dp = rk2, mt = mt0, mf = af, et = duogrid



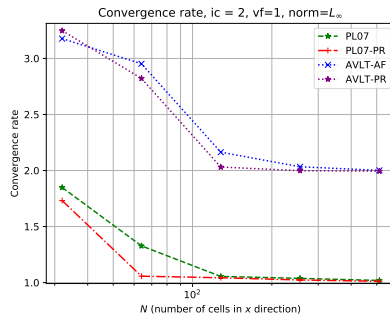
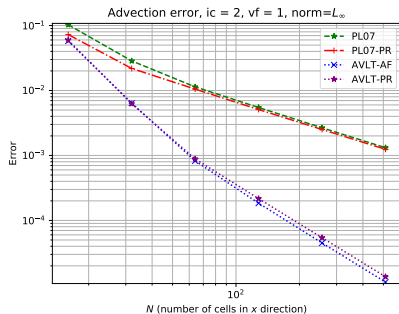
AVLT-AF

Error - N = 512, Time = 5.00e+00, CFL = 4.52e-01, ic = 2, vf = 1
sp = avlt, recon = hyppm, dp = rk2, mt = mt0, mf = gpr, et = duogrid



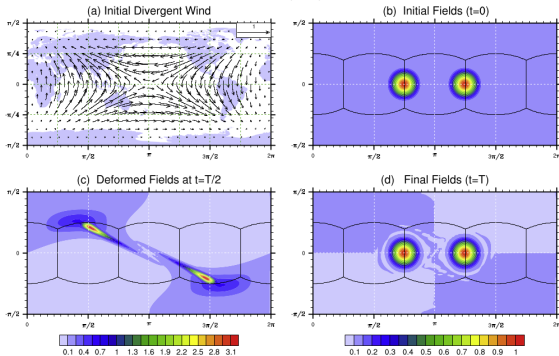
AVLT-PR

One Gaussian hill advection error



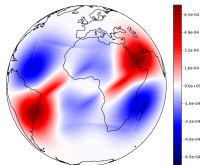
Non-divergent flow from Nair and Lauritzen (2010)

Deformational Flow (DG3): Case-3, CB



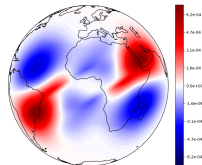
Non-divergent flow from Nair and Lauritzen (2010)

Error - N = 512, Time = 5.00e+00, CFL = 0.90e-02, ic = 3, vf = 3
 up = p07, recon = hyperv, dp = r01, r02 = p07, r03 = none, r04 = p07



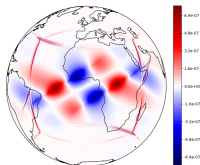
PL07

Error - N = 512, Time = 5.00e+00, CFL = 0.90e-02, ic = 3, vf = 3
 up = p07, recon = hyperv, dp = r01, r02 = p07, r03 = ggr, r04 = diagr02



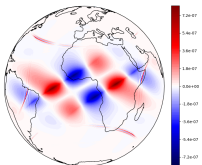
PL07-PR

Error - N = 512, Time = 5.00e+00, CFL = 0.90e-02, ic = 3, vf = 3
 up = avlt, recon = hyperv, dp = r02, r03 = r03, r04 = av, r05 = diagr02

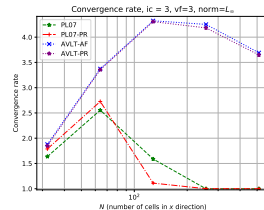
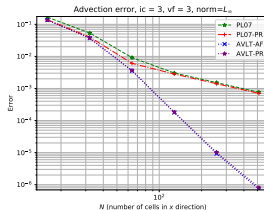


AVLT-AF

Error - N = 512, Time = 5.00e+00, CFL = 0.90e-02, ic = 3, vf = 3
 up = avlt, recon = hyperv, dp = r02, r03 = r03, r04 = ggr, r05 = diagr02



AVLT-PR



Conclusions

Take away message

- Duogrid reduces the grid imprinting but we need to use a mass fixer.
- Flux averaging leads to numerical inconsistency and generates grid imprinting.
- Divergence projection preserves the consistency order and is more efficient to reduce grid imprinting.

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- Flux averaging leads to numerical inconsistency and generates grid imprinting.
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On going work

- Modify the PL07 splitting to ensure the elimination of splitting errors when employing the most accurate metric tensor treatment (MT-0).
- Propose a local mass fixer that preserves the scheme order.
- Implementation of some of these ideas on the FV3 shallow water core.

References I

- Carpenter, R. L., Droegemeier, K. K., Woodward, P. R., & Hane, C. E. (1990). Application of the piecewise parabolic method (ppm) to meteorological modeling. *Monthly Weather Review*, 118(3), 586–612. [https://doi.org/10.1175/1520-0493\(1990\)118<0586:AOTPPM>2.0.CO;2](https://doi.org/10.1175/1520-0493(1990)118<0586:AOTPPM>2.0.CO;2) (cit. on pp. 24–29).
- Chen, X. (2021). The Imars based shallow-water dynamical core on generic gnomonic cubed-sphere geometry [e2020MS002280 2020MS002280]. *Journal of Advances in Modeling Earth Systems*, 13(1), e2020MS002280. <https://doi.org/https://doi.org/10.1029/2020MS002280> (cit. on pp. 35–38).

References II

- Colella, P., & Woodward, P. R. (1984). The piecewise parabolic method (ppm) for gas-dynamical simulations. *Journal of Computational Physics*, 54(1), 174–201.
[https://doi.org/https://doi.org/10.1016/0021-9991\(84\)90143-8](https://doi.org/https://doi.org/10.1016/0021-9991(84)90143-8)
(cit. on pp. 24–29).
- Lin, S.-J. (2004). A “vertically lagrangian” finite-volume dynamical core for global models. *Monthly Weather Review*, 132(10), 2293–2307. [https://doi.org/10.1175/1520-0493\(2004\)132<2293:AVLFDC>2.0.CO;2](https://doi.org/10.1175/1520-0493(2004)132<2293:AVLFDC>2.0.CO;2)
(cit. on pp. 30–34).
- Lin, S.-J., & Rood, R. B. (1996). Multidimensional flux-form semi-lagrangian transport schemes. *Monthly Weather Review*, 124(9), 2046–2070.
[https://doi.org/10.1175/1520-0493\(1996\)124<2046:MFFSLT>2.0.CO;2](https://doi.org/10.1175/1520-0493(1996)124<2046:MFFSLT>2.0.CO;2) (cit. on pp. 19–23).

References III

- Nair, R. D., & Lauritzen, P. H. (2010). A class of deformational flow test cases for linear transport problems on the sphere. *Journal of Computational Physics*, 229(23), 8868–8887. <https://doi.org/https://doi.org/10.1016/j.jcp.2010.08.014> (cit. on pp. 50, 51).
- Putman, W. M., & Lin, S.-J. (2007). Finite-volume transport on various cubed-sphere grids. *Journal of Computational Physics*, 227(1), 55–78. <https://doi.org/https://doi.org/10.1016/j.jcp.2007.07.022> (cit. on pp. 19–23, 30–34).
- Ronchi, C., Iacono, R., & Paolucci, P. (1996). The “cubed sphere”: A new method for the solution of partial differential equations in spherical geometry. *Journal of Computational Physics*, 124(1), 93–114. <https://doi.org/https://doi.org/10.1006/jcph.1996.0047> (cit. on pp. 4–7).

- Rossmannith, J. A. (2006). A wave propagation method for hyperbolic systems on the sphere. *Journal of Computational Physics*, 213(2), 629–658. [https://doi.org/https://doi.org/10.1016/j.jcp.2005.08.027](https://doi.org/10.1016/j.jcp.2005.08.027) (cit. on pp. 35–38).
- Williamson, D., Drake, J., Hack, J., Jakob, R., & Swarztrauber, P. (1992). A standard test set for numerical approximations to the shallow water equations in spherical geometry. *J. Comput. Phys.*, 102, 211–224. [https://doi.org/10.1016/S0021-9991\(05\)80016-6](https://doi.org/10.1016/S0021-9991(05)80016-6) (cit. on pp. 39–41).

Code availability

- Code: <https://github.com/luanfs/cubed-sphere>
- Report: <https://github.com/luanfs/doc-thesis>

Work under progress!

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Work under progress!

Thank you!

Contact: luan.santos@usp.br
<https://luanfs.github.io/>