

Derivation of tri-diagonal matrix for a General Implicit Heat Flow problem

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1 Introduction

In the lecture slides you can find a derivation of the implicit tridiagonal matrix for the heat flow problem using a Dirichlet problem. A small problem you face with this approach is that this solution is not very easy to develop in to a general implicit solver where we can use so-called higher order implicit solvers such as the Crank Nicholson approach based on the trapezoidal rule, which is second order accurate compared to the first order implicit Euler method.

A general approach is possible when we use the Robin boundary condition instead for the lower and the top boundary conditions which for the top boundary can be written as:

$$q_{h_{top}} = -k_{rob_{top}}(T_{top} - T)$$

and for the bottom boundary as:

2 General solution for 1D flow

The general heat flow equation is

$$\mathbf{q}_h = -\lambda \nabla T$$

which combined with the heat balance equation gives

$$\zeta \frac{\partial T}{\partial t} = -\nabla \cdot \mathbf{q}_H.$$

For a 1D vertical implementation on a staggered grid we can discretize these equations as follows.

$$q_{h_{i-\frac{1}{2}}} = -\lambda_{i-\frac{1}{2}} \frac{\Delta T_{i-1}}{\Delta z_{i-1}} = -\lambda_{i-\frac{1}{2}} \frac{T_i - T_{i-1}}{\Delta z_{i-1}}$$

and

$$\begin{aligned}
\zeta_i \frac{T_i^{m+1} - T_i^m}{\Delta t} &= -\frac{q_{h_{i+\frac{1}{2}}} - q_{h_{i-\frac{1}{2}}}}{\Delta z_{i-\frac{1}{2}}} = \\
&= -\left(\frac{\frac{-\lambda_{i+\frac{1}{2}}(T_{i+1}-T_i)}{\Delta z_i} - \frac{-\lambda_{i-\frac{1}{2}}(T_i-T_{i-1})}{\Delta z_{i-1}}}{\Delta z_{i-\frac{1}{2}}} \right) = \\
&= \frac{\lambda_{i-\frac{1}{2}}}{\Delta z_{i-1}\Delta z_{i-\frac{1}{2}}} T_{i-1} - \left(\frac{\lambda_{i-\frac{1}{2}}}{\Delta z_{i-1}\Delta z_{i-\frac{1}{2}}} + \frac{\lambda_{i+\frac{1}{2}}}{\Delta z_i\Delta z_{i-\frac{1}{2}}} \right) T_i + \frac{\lambda_{i+\frac{1}{2}}}{\Delta z_i\Delta z_{i-\frac{1}{2}}} T_{i+1} \\
&= \alpha_i T_{i-1} + \beta_i T_i + \gamma_i T_{i+1}
\end{aligned}$$

vectorizing this for all inner nodes leads to

$$\mathbf{M}(\mathbf{T}^{m+1} - \mathbf{T}^m) = \mathbf{K}\mathbf{T}.$$

3 Top and bottom boundary condition

For the top and the bottom boundary condition we use the Robin condition defined above. Discretizing for the bottom condition we obtain

$$\begin{aligned}
\zeta_1 \frac{T_1^{m+1} - T_1^m}{\Delta t} &= -\frac{q_{H_{\frac{1}{2}}} - q_{H_{-\frac{1}{2}}}}{\Delta z_{\frac{1}{2}}} = -\left(\frac{\frac{-\lambda_{\frac{1}{2}}(T_2-T_1)}{\Delta z_1} + k_{rob_{bot}}(T_1 - T_{bnd_{bot}})}{\Delta z_{\frac{1}{2}}} \right) = \\
&= -\left(\frac{k_{rob_{bot}}}{\Delta z_{\frac{1}{2}}} + \frac{\lambda_{\frac{1}{2}}}{\Delta z_1\Delta z_{\frac{1}{2}}} \right) T_1 + \frac{\lambda_{\frac{1}{2}}}{\Delta z_1\Delta z_{\frac{1}{2}}} T_2 + \frac{k_{rob_{bot}}}{\Delta z_{\frac{1}{2}}} T_{bnd_{bot}}
\end{aligned}$$

which can be written as

$$\zeta_1 \frac{T_1^{m+1} - T_1^m}{\Delta t} = \beta_1 T_1 + \gamma_1 T_2 + y_1$$

for a zero flow boundary condition we can set $k_{rob_{bot}} = 0$.

A similar derivation occurs for the top boundary condition leading to

$$\begin{aligned}
\zeta_n \frac{T_n^{m+1} - T_n^m}{\Delta t} &= -\left(\frac{-k_{rob_{top}}(T_{bnd_{top}} - T_n) - \frac{-\lambda_{n-\frac{1}{2}}(T_n - T_{n-1})}{\Delta z_{n-1}}}{\Delta z_{n-\frac{1}{2}}} \right) = \\
&= \left(\frac{\lambda_{n-\frac{1}{2}}}{\Delta z_{n-1}\Delta z_{n-\frac{1}{2}}} T_{n-1} - \left(\frac{\lambda_{n-\frac{1}{2}}}{\Delta z_{n-1}\Delta z_{n-\frac{1}{2}}} + \frac{k_{rob_{top}}}{\Delta z_{n-\frac{1}{2}}} \right) T_n + \frac{k_{rob_{top}}}{\Delta z_{n-\frac{1}{2}}} T_{bnd_{top}} \right) \\
&= \alpha_n T_{n-1} + \beta_n T_n + y_n.
\end{aligned}$$

4 Explicit, implicit and semi-implicit solutions

It is clear that the general solution including top and boundary conditions is the following matrix equation

$$M(\mathbf{T}^{m+1} - \mathbf{T}^m) = \mathbf{K}\mathbf{T} + \mathbf{y}$$

which lead to the following explicit Euler solution:

$$\mathbf{T}^{m+1} = \mathbf{T}^m + \frac{\mathbf{K}\mathbf{T}^m}{M} + \mathbf{y}^m.$$

We also can develop an implicit solution as follows

$$(M + \mathbf{K})\mathbf{T}^{m+1} = M\mathbf{T}^m + \mathbf{y}^{m+1}.$$

Finally the half implicit (also known as the Crank Nicholson solution) which is based on the trapezoidal rule can be developed as:

$$M(\mathbf{T}^{m+1} - \mathbf{T}^m) = \mathbf{K} \frac{\mathbf{T}^{m+1} + \mathbf{T}^m}{2} + \frac{\mathbf{y}^{m+1} + \mathbf{y}^m}{2}$$

leading to

$$(M - \frac{\mathbf{K}}{2})\mathbf{T}^{m+1} = (M + \frac{\mathbf{K}}{2})\mathbf{T}^m + \frac{\mathbf{y}^{m+1} + \mathbf{y}^m}{2}$$

which is considered to be the most efficient and stable solution for this problem. A challenge however is that for large grid problems the matrices tend to become very large. Luckily they are sparse matrices for which most modern programming languages have efficient tools available.