

From random walk to diffusion equation

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where (in two dimensions):

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- The *Laplacian*, Δu , is defined by $\Delta u = u_{xx} + u_{yy}$.
- One can also consider the diffusion equation *with drift and reaction*:

$$u_t = k\Delta u + \mathbf{b} \cdot \nabla u - \gamma u$$

Symmetric multidimensional random walk - set up

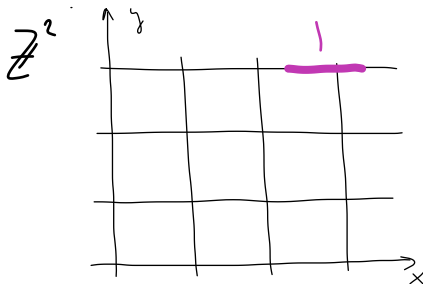
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- Lattice $h\mathbb{Z}^2$ is the set of points whose coordinates are integers multiplied by h .



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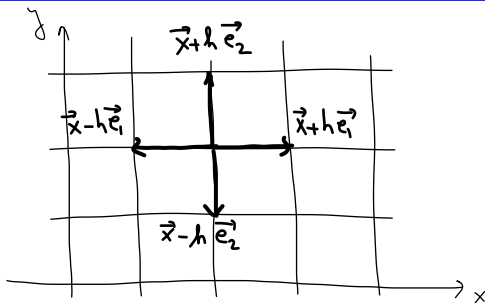
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If we denote by \mathbf{e}_1 and \mathbf{e}_2 the basis of \mathbb{R}^2 , every point $\mathbf{x} \in h\mathbb{Z}^2$, has a discrete neighborhood of 4 points at distance h , given by:

$$\mathbf{x} + h\mathbf{e}_1, \quad \mathbf{x} - h\mathbf{e}_1, \quad \mathbf{x} + h\mathbf{e}_2, \quad \mathbf{x} - h\mathbf{e}_2.$$

Symmetric multidimensional random walk -rules



The particle moves according to the following rules:

- 1 It starts from $\mathbf{x} = \mathbf{0}$.
- 2 If at time t , the particle is located at \mathbf{x} , then at time $t + \tau$ it is at one of the four neighboring points $\mathbf{x} \pm h\mathbf{e}_1$, $\mathbf{x} \pm h\mathbf{e}_2$, with probability $\frac{1}{4}$.
- 3 Each step is independent of the previous one.

Symmetric multidimensional random walk - probability

We now calculate the probability $p(\mathbf{x}, t)$ of finding the particle at position \mathbf{x} at time t . This probability has the following initial conditions: $p(\mathbf{0}, 0) = 1$ and $p(\mathbf{x}, 0) = 0$ if $\mathbf{x} \neq \mathbf{0}$.

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$$p(\mathbf{x}, t + \tau) = \frac{1}{4} \sum_{j=1}^2 \left(p(\mathbf{x} + h\mathbf{e}_j, t) + p(\mathbf{x} - h\mathbf{e}_j, t) \right).$$

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For fixed \mathbf{x} and t , we want to find the limit as $h \rightarrow 0$ and $\tau \rightarrow 0$. Assuming p is defined and smooth, by Taylor's formula we have:

$$p(\mathbf{x}, t + \tau) = p(\mathbf{x}, t) + p_t(\mathbf{x}, t)\tau + o(\tau)$$

$$p(\mathbf{x} + h\mathbf{e}_j, t) = p(\mathbf{x}, t) + p_{x_j}(\mathbf{x}, t)h + \frac{1}{2}p_{x_j x_j}(\mathbf{x}, t)h^2 + o(h^2)$$

$$p(\mathbf{x} - h\mathbf{e}_j, t) = p(\mathbf{x}, t) - p_{x_j}(\mathbf{x}, t)h + \frac{1}{2}p_{x_j x_j}(\mathbf{x}, t)h^2 + o(h^2)$$

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Combining the last four formulas, we have:

$$p(\mathbf{x}, t) + p_t(\mathbf{x}, t)\tau + o(\tau) = \frac{1}{4} \sum_{j=1}^2 \left(2p(\mathbf{x}, t) + p_{x_j x_j}(\mathbf{x}, t)h^2 + o(h^2) \right)$$

Symmetric multidimensional random walk - probability

Therefore, after canceling $p(\mathbf{x}, t)$ and dividing by τ :

$$p_t + o(1) = \frac{1}{4} \frac{h^2}{\tau} \Delta p + o\left(\frac{h^2}{\tau}\right).$$

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To obtain nontrivial limit, $\frac{h^2}{\tau}$ must have a finite and positive limit. The simplest choice is to set the ratio to be equal to a constant:

$$\frac{h^2}{\tau} = 4k, \quad k > 0.$$

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$$p_t = k \Delta p$$

with the initial condition

$$\lim_{t \rightarrow 0+} p(\mathbf{x}, t) = \delta.$$

In this limit, the random walk has become a continuous walk. If $k = 1$, it is called (2-dimensional) Brownian motion.

Walks with different space steps

One variant of a symmetric random walk is to choose a different space step h_j for every direction \mathbf{e}_j , i.e. space step h_1 for direction \mathbf{e}_1 and space step h_2 for direction \mathbf{e}_2 .

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For fixed \mathbf{x} and t , we want to find the limit as $h_j \rightarrow 0$ and $\tau \rightarrow 0$. By Taylor's formula we have:

$$\begin{aligned} p(\mathbf{x}, t + \tau) &= p(\mathbf{x}, t) + p_t(\mathbf{x}, t)\tau + o(\tau) \\ p(\mathbf{x} + h_j \mathbf{e}_j, t) &= p(\mathbf{x}, t) + p_{x_j}(\mathbf{x}, t)h_j + \frac{1}{2}p_{x_j x_j}(\mathbf{x}, t)h_j^2 + o(h_j^2) \\ p(\mathbf{x} - h_j \mathbf{e}_j, t) &= p(\mathbf{x}, t) - p_{x_j}(\mathbf{x}, t)h_j + \frac{1}{2}p_{x_j x_j}(\mathbf{x}, t)h_j^2 + o(h_j^2) \end{aligned}$$

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$$\begin{aligned} p(\mathbf{x}, t) + p_t(\mathbf{x}, t)\tau + o(\tau) &= \frac{1}{4} \sum_{j=1}^2 \left(2p(\mathbf{x}, t) + p_{x_j x_j}(\mathbf{x}, t)h_j^2 + o(h_j^2) \right) \\ &= p(\mathbf{x}, t) + \frac{1}{4} \sum_{j=1}^2 \left(h_j^2 p_{x_j x_j} + o(h_j^2) \right). \end{aligned}$$

Walks with different space steps

Canceling $p(\mathbf{x}, t)$ and dividing by τ yields:

$$p_t + o(1) = \frac{1}{4} \sum_{j=1}^2 \left(\frac{h_j^2}{\tau} p_{x_j x_j} + o\left(\frac{h_j^2}{\tau}\right) \right).$$

This time, to obtain nontrivial limit, we ask that:

$$\frac{h_j^2}{\tau} = 4k_j, \quad k_j > 0, \quad j = 1, 2.$$

This results in a **diffusion equation**

$$p_t = \sum_{j=1}^2 k_j p_{x_j x_j} = k_1 p_{x_1 x_1} + k_2 p_{x_2 x_2}.$$

Walks with drift

We can also set different probabilities in different directions. For example, if at time t , the particle is located at \mathbf{x} , then at time $t + \tau$ it is

at $\mathbf{x} + h_1 \mathbf{e}_1$ with probability p_1 ,

at $\mathbf{x} - h_1 \mathbf{e}_1$ with probability q_1 ,

at $\mathbf{x} + h_2 \mathbf{e}_2$, with probability p_2 ,

at $\mathbf{x} - h_2 \mathbf{e}_2$, with probability q_2 ,

with $p_1 + q_1 = \frac{1}{2}$ and $p_2 + q_2 = \frac{1}{2}$.

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Then:

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Using the Taylor formulas from before:

$$p(\mathbf{x}, t + \tau) = p(\mathbf{x}, t) + p_t(\mathbf{x}, t)\tau + o(\tau)$$

$$q_j p(\mathbf{x} + h_j \mathbf{e}_j, t) = q_j p(\mathbf{x}, t) + q_j p_{x_j}(\mathbf{x}, t) h_j + \frac{q_j}{2} p_{x_j x_j}(\mathbf{x}, t) h_j^2 + o(h_j^2)$$

$$p_j p(\mathbf{x} - h_j \mathbf{e}_j, t) = p_j p(\mathbf{x}, t) - p_j p_{x_j}(\mathbf{x}, t) h_j + \frac{p_j}{2} p_{x_j x_j}(\mathbf{x}, t) h_j^2 + o(h_j^2)$$

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... we get:

$$p_t(\mathbf{x}, t)\tau + o(\tau) = \sum_{j=1}^2 \left(\frac{1}{4} p_{x_j x_j}(\mathbf{x}, t) h_j^2 + (q_j - p_j) p_{x_j}(\mathbf{x}, t) h_j + o(h_j^2) \right)$$

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If, in addition to $\frac{h_j^2}{\tau} = 4k_j$, we have

$$\frac{q_j - p_j}{h_j} \rightarrow \beta_j \quad \text{and} \quad b_j = 4k_j \beta_j,$$

the limit leads to a **drift-diffusion equation**:

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$$p_t = \sum_{j=1}^2 k_j p_{x_j x_j} + \sum_{j=1}^2 b_j p_{x_j}$$

Walks with drift and reaction

One can also assume that the particle loses mass at the constant rate γ . That is, from time t to time $t + \tau$, the following part of the particle disappears:

$$Q(\mathbf{x}, t) = \tau\gamma p(\mathbf{x}, t).$$

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Therefore,

$$p(\mathbf{x}, t + \tau) = \sum_{j=1}^2 \left[q_j \left(p(\mathbf{x} + h_j \mathbf{e}_j, t) - Q(\mathbf{x} + h_j \mathbf{e}_j, t) \right) + p_j \left(p(\mathbf{x} - h_j \mathbf{e}_j, t) - Q(\mathbf{x} - h_j \mathbf{e}_j, t) \right) \right].$$

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Note that:

$$\begin{aligned} \sum_{j=1}^2 \left(q_j Q(\mathbf{x} + h_j \mathbf{e}_j, t) + p_j Q(\mathbf{x} - h_j \mathbf{e}_j, t) \right) &= Q(\mathbf{x}, t) + \sum_{j=1}^2 (q_j - p_j) h_j Q_x(\mathbf{x}, t) + \dots \\ &= \tau\gamma p(\mathbf{x}, t) + O(\tau h_j), \end{aligned}$$

Walks with drift and reaction

Therefore,

$$p_t(\mathbf{x}, t)\tau + o(\tau) = \sum_{j=1}^2 \left(\frac{1}{4} p_{x_j x_j}(\mathbf{x}, t) h_j^2 + (q_j - p_j) p_{x_j}(\mathbf{x}, t) h_j + o(h_j^2) - \tau \gamma p(\mathbf{x}, t) + O(\tau h_j) \right)$$

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Assuming, as before, that $\frac{h_j^2}{\tau} = 4k_j$ and $\frac{q_j - p_j}{h_j} \rightarrow \beta_j$ and $b_j = 4k_j\beta_j$, we obtain a **drift-diffusion-reaction equation**:

$$p_t = \sum_{j=1}^2 k_j p_{x_j x_j} + \sum_{j=1}^2 b_j p_{x_j} - \gamma p.$$