# Distributing Multipartite Entanglement over Noisy Quantum Networks

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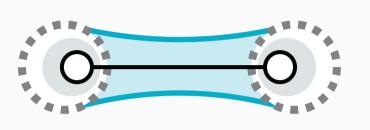


#### Quantum Networks

A future quantum internet provides support for several quantum technologies, ranging from quantum secure communications to distributed quantum computation and distributed quantum metrology and sensing.

Some applications require the distribution of multipartite entanglement, such as quantum conference key agreement (CKA) [1], atomic clocks synchronization [2] and quantum sensor networks [3]. Optimizing its distribution, naturally impacts the applications that utilize this important resource.

Our quantum network model is a graph constituted by a set of nodes and a set of links. Each network link, connecting two neighbouring nodes, contains:



- Noisy quantum channel Imperfect quantum memory Bipartite entangled state

The protocols for entanglement generation and entanglement swapping are probabilistic, and each entangled pair of the network is modelled by a Werner state:

$$\rho_W = \gamma |\phi^+\rangle \langle \phi^+| + (1-\gamma)\mathbb{I}$$

where  $\gamma = (4F-1)/3$  , and F is the fidelity of the state.

Optimising bipartite entanglement distribution over a quantum network is essentially finding the shortest-path in a graph. The definition of shortest-path depends on what we want to optimise for.

For bipartite entanglement we ought to maximise the fidelity of the state and probability of success and minimise the communications time.

However, since the fidelity of an entangled pair will depend on the generation fidelity and the decoherence, we must first optimise four parameters individually and then map them to only three:

$$[p_{m:n}, t_{m:n}, \gamma_{m:n}, \sigma_{m:n}] \mapsto [p_{m:n}, t_{m:n}, F_{m:n}]$$

$$\frac{4F_{m:n} - 1}{3} = \gamma_{m:n}e^{-t_{m:n}/\sigma_{m:n}}$$

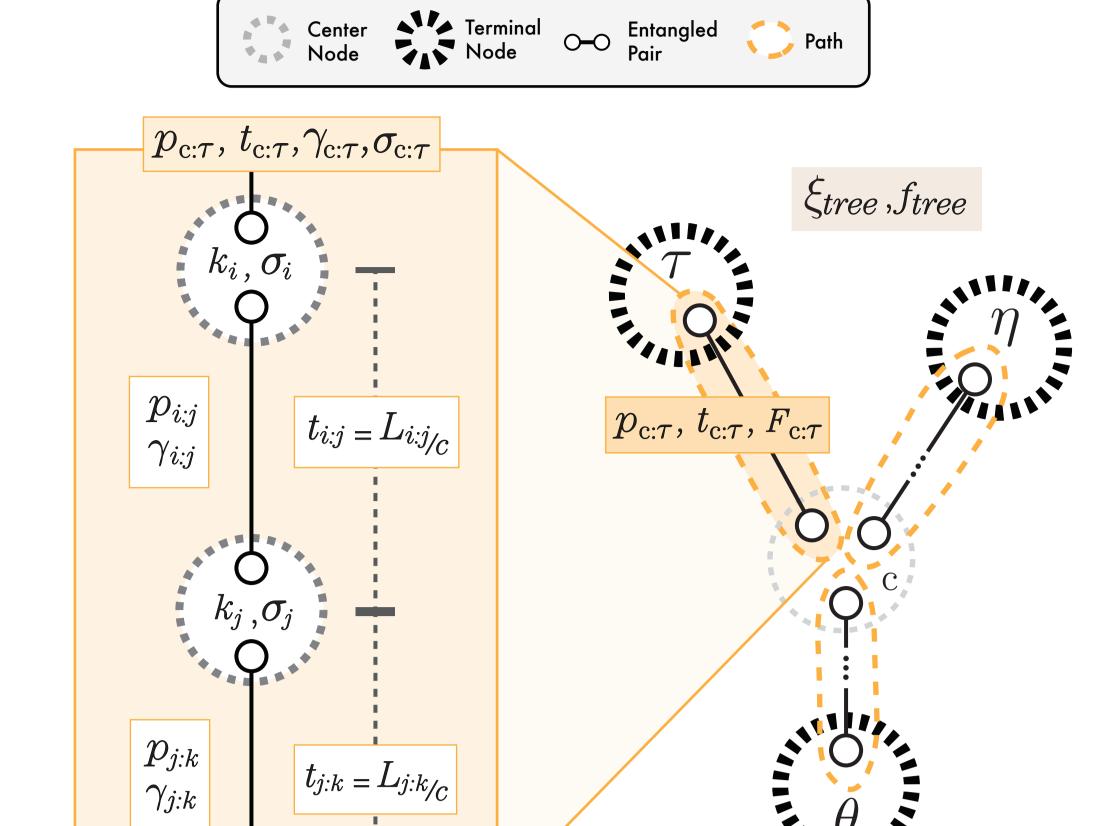


Figure 1. Branch composed by an arbitrary number of links. Each link has a fidelity  $\gamma_{i:j}$  , a communications time  $t_{i:j}$  and probability of successful generation  $p_{i:j}$ . Moreover, each node has a memory coherence time  $\sigma_i$  and a probability of successful entanglement swapping  $k_i$ .

Figure 2. Star composed by 3 different paths, connecting terminal nodes  $\mathcal{T}$ ,  $\eta$  and  $\theta$  to the center node c. Notice that the rate  $\xi_{tree}$  depends on each branch probability of success  $p_{c: au}$  and communications time  $t_{c: au}$ . Moreover, the fidelity of the final multipartite state  $f_{tree}$  depends on each branch fidelity  $F_{c:\tau}$ .

Figure 4. Simulations of the MOSP algorithm and the Star Algorithm for 3 and 4 qubits GHZ state distribution in Erdös-Rényi networks with average degree 3. Complexity scaling with the number of network nodes.

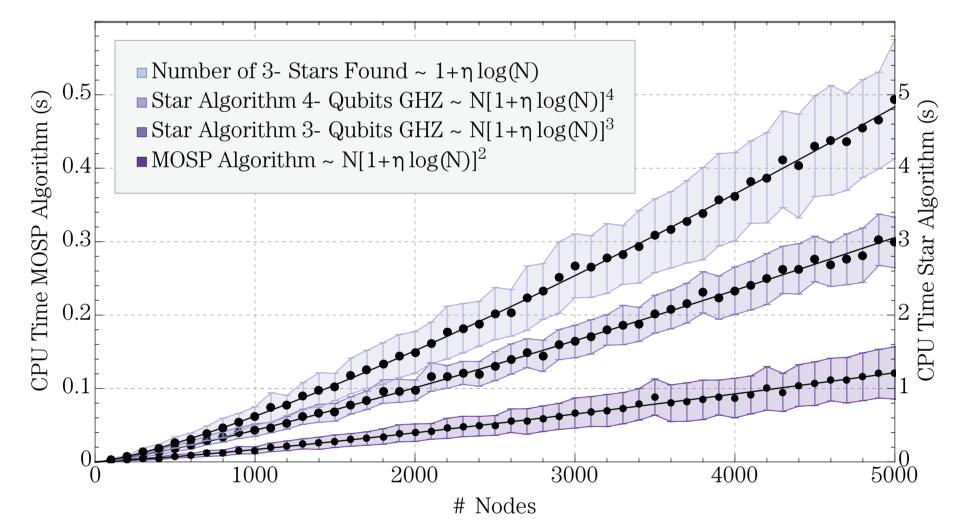
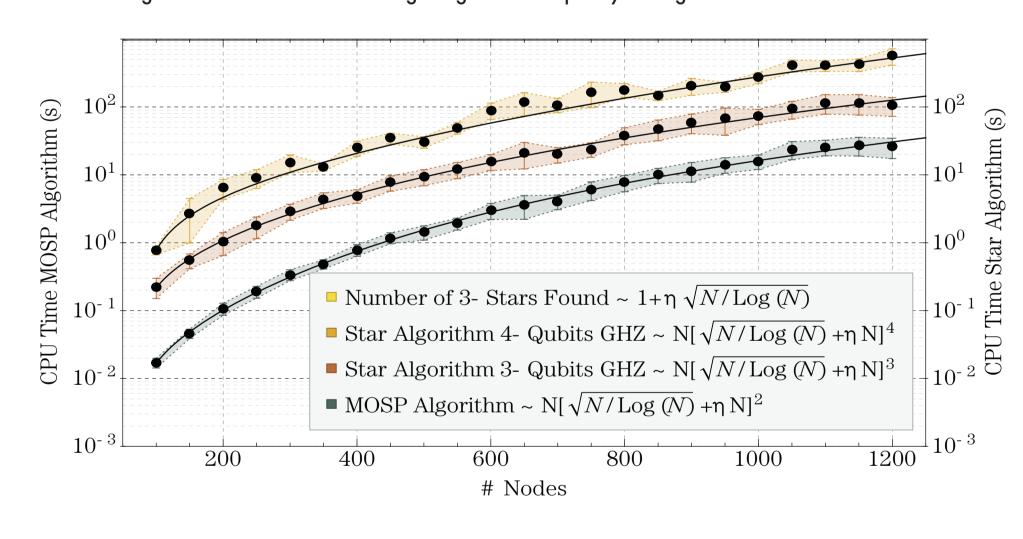


Figure 5. Simulations of the MOSP algorithm and the Star Algorithm for 3 and 4 qubits GHZ state distribution in random geometric networks with average degree 8. Complexity scaling with the number of network nodes.



### Multipartite Entanglement Distribution

Similarly to [4], to distribute a 3-qubit GHZ state, first we have to find the shortest-tree connecting the set of terminal nodes, and then perform a set of measurements on some qubits. For 3 qubits, the shortest-tree is a star-graph, and the scheme is equivalent to first distributing bipartite entanglement between every terminal and a center node and then performing a set of operations on the center node.

Thus, the central idea is finding the center node for which this results in the optimal solution. Optimality is defined concerning the metrics for the distribution of the final multipartite entangled state, namely the rate of distribution and the fidelity.

The average time for distribution depends on the maximum of each path's communications time multiplied by the average number of attempts before the first success. The average rate would then be the inverse of:

$$T_{tree} = 2 \cdot \frac{\max_{\tau \in \mathcal{T}} \{t_{c:\tau}\}}{\prod_{\tau \in \mathcal{T}} p_{c:\tau}} , \quad \xi_{tree} = \frac{1}{T_{tree}}$$

The form of the final distributed state is equivalent to applying a depolarising channel to each qubit of the GHZ state, correspondent to each branch connecting the center node to a terminal. The final fidelity of this state is described by:

$$f_{tree} = \frac{1}{2} \left[ \prod_{\tau \in \mathcal{T}} \frac{1 + 2F_{\tau}}{3} + \prod_{\tau \in \mathcal{T}} \frac{2(1 - F_{\tau})}{3} + \prod_{\tau \in \mathcal{T}} \frac{4F_{\tau} - 1}{3} \right]$$

# Multi-objective Routing Algorithm

In our problem, we want to optimise two distinct parameters. For that, we require a multi-objective approach. In 1984, Martins [5] introduced multi-objective shortest-paths (MOSP) algorithms, defining the dominance relation and using the Pareto optimality definition - optimal paths are non-dominated paths. This results in a set of optimal solutions that is usually larger than one the Pareto front.

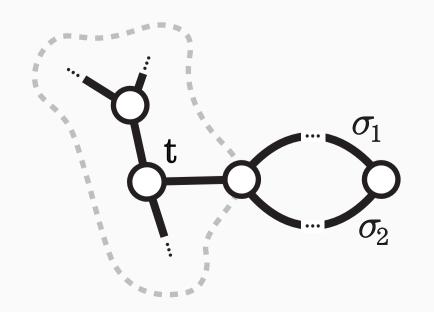
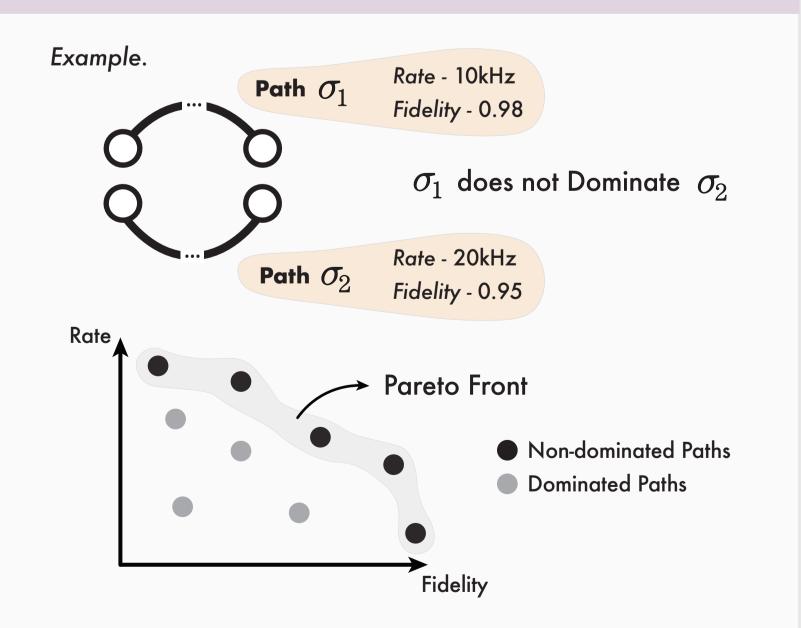


Figure 3. Label-Isotonicity: for every tree and path, when extending a tree with a path, choosing a better path results in a better tree.



In [6] Sobrinho introduces a more fundamental definition for these metrics involved in routing problems - algebras for routing. These algebras incorporate the parameters of the links, paths and how paths are extended with links - the metrics -, allowing for a bigger level of abstraction.

Multi-objective routing - one algebra for each parameter (rate, fidelity,...). To prove optimality, we expanded these algebras for the case of trees (instead of paths) and introduced label-isotonicity (Figure 3).

[5] E. Q. V. Martins, Eur. J. Oper. Res., vol. 16, no. 2, pp. 236–245 (1984)

# Results and Conclusions

- > We were able to calculate the effect of imperfect entanglement throughout quantum networks on the distribution of GHZ and W states, importante resources for distributed applications, such as communication and sensing;
- > We defined a metric for the fidelity and realise an algorithm with polinomial runtime for distributing these states optimally, maximising simultaneously the fidelity and the rate of distribution;
- > Using classical techniques and tools, we can prove the optimality and create the necessary formalism to expand the methodology, while also providing the necessary conditions to attain this optimality.

[4] C. Meignant, D. Markham, and F. Grosshans, Phys. Rev. A, vol. 100, no. 5, p. 052333 (2019)

About the Authors:





