

## A kinetic Poincaré inequality

### Theorem 1 ([1, 2]):

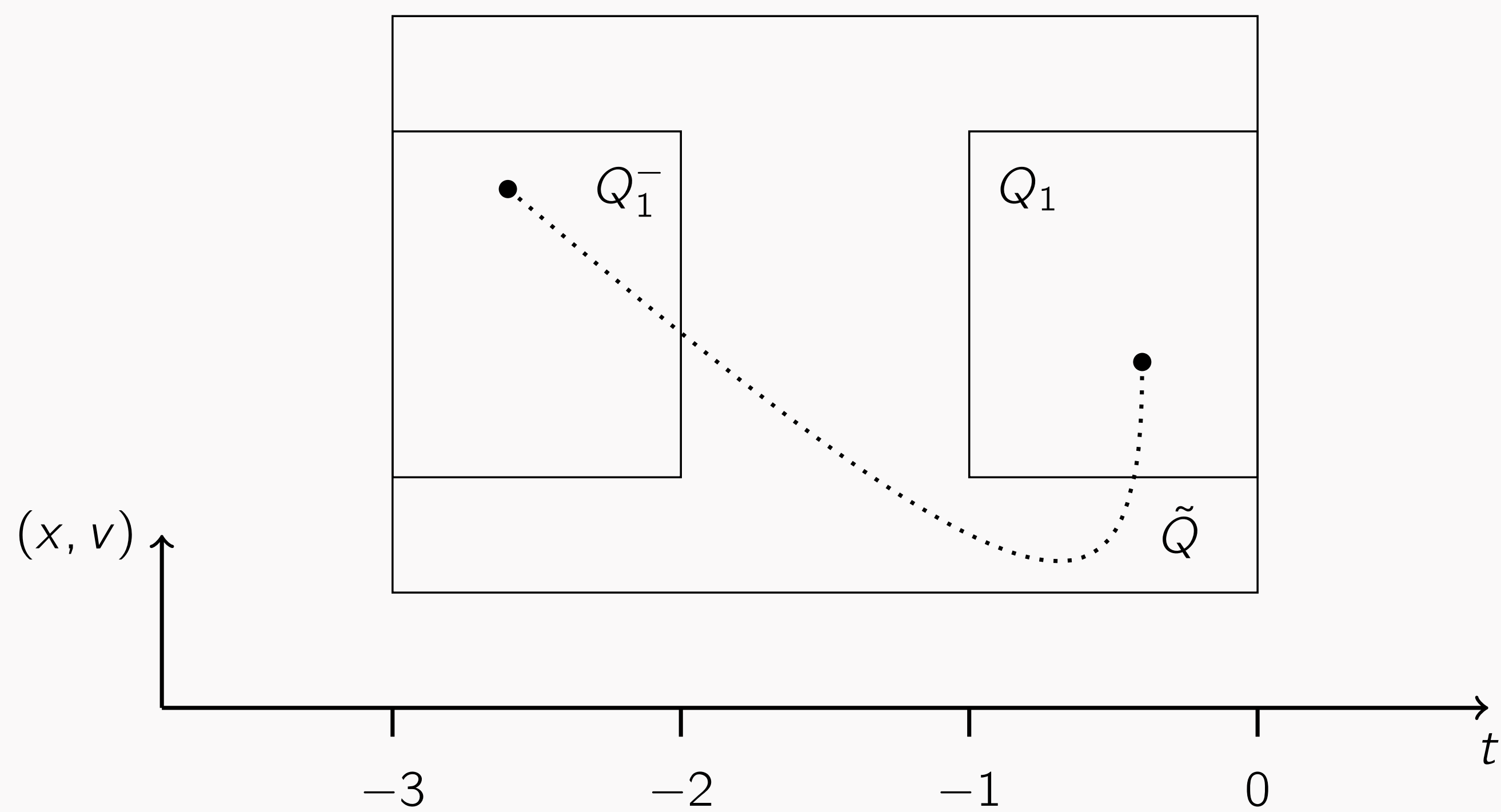
There exists  $R > 1$  such that the following holds for any  $0 \leq \varphi \in C_c^\infty(B_1 \times B_1)$  nonzero. Let  $f \in L^1_{t,x}((-3, 0] \times B_R; W^{1,1}_v(B_R))$  be such that

$$(\partial_t + v \cdot \nabla_x) f \leq \nabla_v \cdot S$$

in the distributional sense for some  $S \in L^1_{t,x,v}((-3, 0] \times B_R \times B_R; \mathbb{R}^d)$ . Then,

$$\left\| \left( f - \frac{1}{c_\varphi} \int_{Q_1^-} f \varphi \, d(s, y, w) \right)_+ \right\|_{L^1(Q_1)} \leq C \left( \|\nabla_v f\|_{L^1(\tilde{Q})} + \|S\|_{L^1(\tilde{Q})} \right),$$

where  $C = C(d, \varphi) > 0$ ,  $c_\varphi = \int_{Q_1^-} \varphi d(s, y, w)$ ,  $Q_1^- = (-3, -2] \times B_1 \times B_1$ ,  $Q_1 = (-1, 0] \times B_1 \times B_1$ , and  $\tilde{Q} = (-3, 0] \times B_R \times B_R$ .



## Kinetic trajectories

### Definition:

Let  $(t_0, x_0, v_0), (t_1, x_1, v_1) \in \mathbb{R}^{1+2d}$  with  $t_0 \neq t_1$ . A kinetic trajectory is a map  $\gamma \in C([0, 1]; \mathbb{R}^{1+2d})$  differentiable in  $(0, 1)$  with  $\gamma(0) = (t_0, x_0, v_0)$  and  $\gamma(1) = (t_1, x_1, v_1)$  satisfying  $\dot{\gamma}_x = \dot{\gamma}_t \gamma_v$  in  $(0, 1)$ .

### Important property:

For  $g: \mathbb{R}^{1+2d} \rightarrow \mathbb{R}$  smooth we have

$$\begin{aligned} \frac{d}{dr} g(\gamma(r)) &= \dot{\gamma}_t(r) [\partial_t g] + \dot{\gamma}_x(r) \cdot [\nabla_x g](\gamma(r)) + \dot{\gamma}_v(r) \cdot [\nabla_v g](\gamma(r)) \\ &= \dot{\gamma}_t(r) [\partial_t g + v \cdot \nabla_x g](\gamma(r)) + \dot{\gamma}_v(r) \cdot [\nabla_v g](\gamma(r)). \end{aligned}$$

### Construction ([2, 3]):

Ansatz: set  $\gamma_t = t_0 + (t_1 - t_0)r$  (kinetic scaling  $r \mapsto (rt, r^{\frac{3}{2}}x, r^{\frac{1}{2}}v)$ ) and

$$\dot{\gamma}_v(r) = \ddot{g}_0(r) \mathbf{m}_0 + \ddot{g}_1(r) \mathbf{m}_1$$

for two forcings  $\ddot{g}_0, \ddot{g}_1: [0, 1] \rightarrow \mathbb{R}$  and vectorial parameters  $\mathbf{m}_0, \mathbf{m}_1 \in \mathbb{R}^d$ .

Integration and  $\dot{\gamma}_x = \dot{\gamma}_t \gamma_v$  yields

$$\begin{cases} \gamma_t(r) = t_0 + (t_1 - t_0)r \\ \gamma_x(r) = (t_1 - t_0)g_0(r) \mathbf{m}_0 + (t_1 - t_0)g_1(r) \mathbf{m}_1 + (t_1 - t_0)r v_0 + x_0 \\ \gamma_v(r) = \dot{g}_0(r) \mathbf{m}_0 + \dot{g}_1(r) \mathbf{m}_1 + v_0. \end{cases}$$

Ansatz:  $g_0(r) \sim r^{p_0}$  and  $g_1(r) \sim r^{p_1}$  for  $p_0 \neq p_1 \in (1, 2)$ . This is needed to have enough linear independence at  $r = 1$  and for good properties of  $\gamma$ .

If this is the case, solving  $\gamma(1) = (t_1, x_1, v_1)$  yields  $\mathbf{m}_0, \mathbf{m}_1$ .

## Proof

Take kinetic trajectory connecting  $(t, x, v) \in Q_1$  with  $(s, y, w) \in Q_1^-$ . Then

$$\begin{aligned} f - \frac{1}{c_\varphi} \int_{Q_1^-} f \varphi \, d(s, y, w) &= \frac{1}{c_\varphi} \int_{Q_1^-} (f(t, x, v) - f(s, y, w)) \varphi(y, w) d(s, y, w) \\ &= -\frac{1}{c_\varphi} \int_{Q_1^-} \int_0^1 \frac{d}{dr} f(\gamma(r)) dr \varphi(y, w) d(s, y, w) \\ &= -\frac{1}{c_\varphi} \int_0^1 \int_{Q_1^-} \dot{\gamma}_t(r) [(\partial_t + v \cdot \nabla_x) f](\gamma(r)) \varphi(y, w) d(s, y, w) dr \\ &\quad - \frac{1}{c_\varphi} \int_0^1 \int_{Q_1^-} \dot{\gamma}_v(r) \cdot [\nabla_v f](\gamma(r)) \varphi(y, w) d(s, y, w) dr =: I_1 + I_2. \end{aligned}$$

As  $\dot{\gamma}_t(r) = (s - t) \leq 0$ , we can use the subsolution property. Substitute

$$\begin{pmatrix} \tilde{y} \\ \tilde{w} \end{pmatrix} = \Phi_{r,t,x,v,s}(y, w) = \begin{pmatrix} \gamma_x(r) \\ \gamma_v(r) \end{pmatrix} = \mathcal{A}_{s-t}(r) \begin{pmatrix} y \\ w \end{pmatrix} + \mathcal{B}_{s-t}(r) \begin{pmatrix} x \\ v \end{pmatrix}$$

to perform a partial integration as follows

$$\begin{aligned} I_1 &\lesssim \int_0^1 \int_{Q_1^-} [\nabla_v \cdot S](\gamma(r)) \varphi(y, w) d(s, y, w) dr \\ &= \int_0^1 \int_{-3}^{-2} \int_{\Phi(B)} [\nabla_v \cdot S](\gamma_t(r), \tilde{y}, \tilde{w}) \varphi(\Phi^{-1}(\tilde{y}, \tilde{w})) |\det \mathcal{A}(r)|^{-1} d(\tilde{y}, \tilde{w}) ds dr \\ &= \int_0^1 \int_{-3}^{-2} \int_{\Phi(B)} S(\gamma_t(r), \tilde{y}, \tilde{w}) \cdot ([\nabla \varphi]^T(\Phi^{-1}(\tilde{y}, \tilde{w}))[D_{\tilde{w}} \Phi^{-1}](\tilde{y}, \tilde{w})) \\ &\quad \cdot |\det \mathcal{A}(r)|^{-1} d(\tilde{y}, \tilde{w}) ds dr \\ &= \int_0^1 \int_{Q_1^-} S(\gamma(r)) \cdot ([\nabla \varphi]^T(y, w) (\mathcal{A}^{-1})_{:,2}) d(s, y, w) dr \\ &\lesssim \int_0^1 \int_{Q_1^-} r^{1-\max\{p_0, p_1\}} |S|(\gamma(r)) d(s, y, w) dr. \end{aligned}$$

Moreover,

$$\begin{aligned} I_2 &\lesssim \int_0^1 \int_{Q_1^-} |\dot{\gamma}_v(r)| |\nabla_v u|(\gamma(r)) \varphi(y, w) d(s, y, w) dr \\ &\lesssim \int_0^1 \int_{Q_1^-} r^{\min\{p_0, p_1\}-2} |\nabla_v u|(\gamma(r)) d(s, y, w) dr. \end{aligned}$$

The singularities are integrable due to the careful construction. Next, integrate on  $Q_1$  split the  $r$  integral in  $(0, \frac{1}{2})$  and  $(\frac{1}{2}, 1)$  and change coordinates with the opposite end of the trajectory.

## Literature

- Theorem 1 was proven first in [1] using piecewise trajectories along the vector fields  $\partial_t + v \cdot \nabla_x$ ,  $\nabla_v$ , and  $\nabla_x$  where control on the transfer of regularity is needed.
- A similar Poincaré inequality for weak solutions is proven in Albritton-Armstrong-Mourrat-Novack '24, see also Guerand-Imbert-Mouhot '24.
- Follow-up work [3] studying a Kolmogorov equation with more commutators.
- Theorem 1 is an important tool in the context of the De Giorgi-Nash-Moser theory for the Kolmogorov equation with rough coefficients, see [1].
- Trajectories and Moser's approach to the regularity of weak solutions to elliptic/parabolic equations with rough coefficients is studied in [4].

## References

- [1] Jessica Guerand and Clément Mouhot. Quantitative De Giorgi methods in kinetic theory. *Journal de l'École polytechnique. Mathématiques*, 2022.
- [2] Lukas Niebel and Rico Zacher. On a kinetic Poincaré inequality and beyond. *Journal of Functional Analysis*, 2025.
- [3] Francesca Anceschi, Helge Dietert, Jessica Guerand, Amélie Loher, Clément Mouhot, and Annalaura Rebusci. Poincaré inequality and quantitative De Giorgi method for hypoelliptic operators. arXiv:2401.12194, 2024.
- [4] Lukas Niebel and Rico Zacher. A trajectorial interpretation of Moser's proof of the Harnack inequality. *To appear in Annali della Scuola Normale Superiore di Pisa. Classe di Scienze. Serie V*, 2025.

