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Computer Simulations in Statistical Physics

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Problem set 3

Proseminar

Problem 3.1 Ising model

The Ising model is a cornerstone of classical statistical mechanics. It is a mathematical model of ferromagnetism that can be solved exactly in one and two dimensions. In two and higher dimensions the model shows a phase transition.

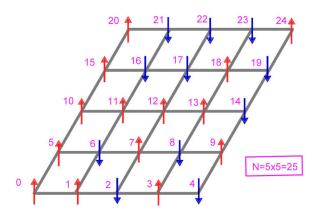
The Ising model it is defined on a lattice in d dimensions, whose sites are occupied by a spin taking two values $s = \pm 1$. The energy of a configuration $\{s_k\}$ is given by

$$E\left(\left\{s_{k}\right\}\right) = -J\sum_{\langle i,j\rangle} s_{i}s_{j} , \qquad (1)$$

where J > 0 is a coupling constant and the sum is extended to nearest neighbors. In the following, we set J = 1.

- a) Which is the energy of a completely ordered configuration with all spin values equal to1? What do you expect for the energy of a random configuration?
- b) Generate a random configuration of N spins displaced over a square lattice with periodic boundary conditions (take $N=10\times 10$) and compute its energy. Programming Hint: think object oriented, i.e. define a class lattice with its own relevant functions. When you add a function to a class, allways test if it is working!

Also remember: it is faster to deal with arrays of size N than with a matrix $\sqrt{N} \times \sqrt{N}$!



- c) Repeat the previous point for M = 100 different random configurations. Plot the energy E_k with k = 1, ..., M and compare the result with the energy of a configuration in which all the spins are equal to 1.
- d) use the code snippet "plot_Ising_configuration.py" to visualize the configurations.

Problem 3.2 Ising model - Metropolis algorithm

We discuss here the Metropolis algorithm for the Ising model. In this algorithm spin configurations are sampled by a Markov chain, which is characterized by a transition probability $P(\mu \to \nu)$ where μ and ν are two spin configurations. In the algorithm $P(\mu \to \nu) \neq 0$ only if the two configurations differ by a single spin. We initialize the system by selecting a random configuration μ of a $N = L \times L$ lattice with periodic boundary conditions. The practical implementation of the algorithm is as follows:

- 1. Select one of the N spins at random, say spin s_i .
- 2. Calculate $\Delta E = E_{fin} E_{in}$, where E_{in} is the energy of the present configuration and E_{fin} the energy of the configuration obtained by flipping the spin $s_i \to -s_i$.
- 3. Flip the spin $s_i \to -s_i$ if $\Delta E \le 0$ or if $r \le \exp(-\Delta E/k_B T)$, where r is a uniform random number in [0, 1].
- 4. Go to (1).

Typical observables one computes in the Ising model are the energy, as defined by Eq. (1), or the magnetization per spin defined as

$$m = \frac{1}{N} \sum_{i=1}^{N} s_i \,, \tag{2}$$

or spin-spin autocorrelations functions. These quantities reach their equilibrium value after some equilibration time τ_{eq} . Another quantity we compute is the (temporal) autocorrelation function which is defined for a system in which time is a discrete variable

$$\chi(t) = \frac{1}{T} \sum_{s=\tau_{eq}}^{\tau_{eq}+T} \left[m(s) - \langle m \rangle \right] \left[m(s+t) - \langle m \rangle \right] , \qquad (3)$$

where m(t) is the magnetization at timestep t, and $\langle m \rangle$ is the average equilibrium magnetization. This equation is only valid when the Monte Carlo run has reached an equilibrium configuration $(t > \tau_{eq})$. This autocorrelation function decreases exponentially like

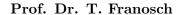
$$\chi(t) \sim \exp\left(-t/\tau\right) ,$$
(4)

where τ defines the autocorrelation time of the simulation. This is the timescale of the simulation, and it indicates how much time it takes to arrive at a new uncorrelated configuration of the model.

a) For a temperature T=2 and N=2500 and starting from a random configuration, use the Metropolis algorithm to compute and plot the energy and the magnetization per spin as a function of the Monte Carlo time step. (For making the plot use the code snippet "Plot_Ising_trajectory.py") Compare the average magnetization with the exact value (valid in the ferromagnetic phase):

$$m(T) = \left[1 - \sinh^{-4}(2J/T)\right]^{1/8}$$
, (5)

- b) From your plot estimate the equilibration time τ_{eq} .
- c) Use the code snippet "Plot_Ising_trajectory.py" to visualize time evolution.
- d) Determine the correlation time τ (For making the fit and plotting use the code snippet "Plot_Ising_trajectory.py")





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Solution to Problem 3.1 Ising model

We start by defining a class "lattice". In this class are contained all relevant functions and subroutines that we can implement on a lattice configurations.

file = lattice.h

```
class Lattice {
1
2
     public:
3
       Lattice(int, double [], double []);
                                                   // Constructor
       int get_spin_index(int, int);
4
5
       double get_spin_value(double [], int, int);
6
       int get_right(int);
7
       int get_left(int);
       int get_above(int);
8
9
       int get_below(int);
       double get_energy(double []);
10
11
12
     private:
13
       int L;
14 };
```

file = lattice.cpp

```
#include <math.h>
#include "lattice.h"

#include "nrutil.h"

Lattice::Lattice(int N, double s[], double s_initial_values[]){
    // call constructur make a lattice with periodic boundary,
```

```
7
    // spins initially setted equal to the array "s_initial_values"
     L = sqrt(N);
8
9
     for (int i =0; i<N; i++){</pre>
10
       s[i] = s_initial_values[i];
11
     }
12
  }
13
   int Lattice::get_spin_index(int i, int j){
14
15
     // i (j) = row (column) index from 0 to L-1
     // Output: the spin index corresponding to (i,j)
16
17
     int k = i*L + j;
18
     return k;
  }
19
20
21
  double Lattice::get_spin_value(double s[], int i, int j){
     // i (j) = row (column) index from 0 to L-1
22
23
     // Output: the spin value corresponding to (i,j)
24
     int k = i*L + j;
25
     return s[k];
26 }
27
28
  int Lattice::get_left(int i){
29
     // return the index of the site at the left of site i
30
     int j=-1;
31
     if (i%L == 0){
32
       j = i+L-1;
33
     } else {
34
       j = i-1;
35
36
     return j;
  }
37
38
39
   int Lattice::get_right(int i){
40
     // return the index of the site at the right of site i
41
     int j=-1;
42
     if ((i+1)%L == 0){
       j = i+1-L;
43
     } else {
44
45
       j = i+1;
     }
46
47
     return j;
48
   }
49
   int Lattice::get_below(int i){
50
     // return the index of the site above of site i
51
52
     int j=-1;
53
     if (i<L){</pre>
       j = i + (L-1)*L;
54
55
     } else {
       j = i-L;
56
     }
57
58
     return j;
59 }
```

```
60
61
62
   int Lattice::get_above(int i){
63
     // return the index of the site below of site i
64
     int j=-1;
     if (i>=L*(L-1)){
65
66
       j = (i+L) \% (L*L);
     } else {
67
       j = i+L;
68
     }
69
70
     return j;
   }
71
72
73
  double Lattice::get_energy(double s[]){
     \ensuremath{//} return the energy E of a given configuration
74
     int N = L*L;
75
76
     double E = 0.0;
     for (int i =0; i<N; i++){</pre>
77
       E -= s[i]*s[get_right(i)] + s[i]*s[get_below(i)];
78
79
     }
80
     return E;
81 }
```

Then we write a very compact "main" file.

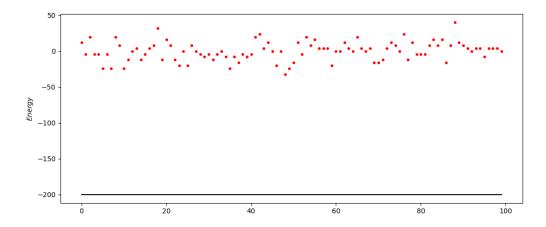
file = Ising.h

```
1 | int L;
2
   int N;
3
4
  double * s;
  double * s_initial_values;
  double * s0;
6
7
  double * s1;
  double * s2;
8
9
   double * s3;
  double * s4;
10
  double * s5;
11
12
13 | double * mvector;
14 double mean (double x[], int NN);
15 double variance (double x[], int NN);
16 double autocorrelation (double x[], int NN, int tau);
```

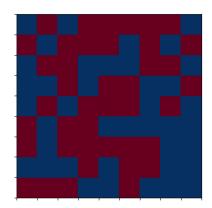
file = Ising.cpp

```
// compile with: g++ Ising.cpp lattice.cpp nrutil.cpp -o test -02 -w
!include <stdio.h>
#include <stdlib.h>
#include <math.h>
#include <sstream>
#include "Ising.h"
#include "lattice.h"
#include "nrutil.h"
```

```
10 | int main (){
     // set parameters
11
12
     L = 10;
     N = L*L;
13
14
15
     // array of length N
16
     s = dvector(0, N-1);
17
18
     // array for a completely ordered configuration
     s_{initial\_values} = dvector(0,N-1);
19
20
     for (int i =0; i<N; i++){</pre>
21
         s_initial_values[i] = 1.0;
22
23
24
     /st define lattice as an element of the Lattice class in the
         initial configuration defined by the array "s_initial_values" */
25
26
     class Lattice *lattice;
27
     lattice = new Lattice(N, s, s_initial_values);
28
     // get the energy of this configuration
29
     double Energy_order = lattice->get_energy(s);
30
     printf("%f \n", Energy_order);
31
     // open output file
32
33
     FILE * out;
34
     out = fopen("output1.dat","w");
35
36
     // random configurations
37
     for (int j=0; j<100; j++){</pre>
38
       for (int i =0; i<N; i++){</pre>
39
          double x = rand()/((double) RAND_MAX)-0.5;
40
          if (x>0) {
41
           s[i]=1.0;
42
          } else {
43
            s[i] = -1.0;
         }
44
45
46
       double Energy = lattice->get_energy(s);
       fprintf(out, "%d %.4f %.4f \n", j, Energy_order, Energy);
47
48
49
     fclose(out);
50
     out = fopen("output2.dat","w");
51
52
     for (int i =0; i<N; i++){</pre>
53
       fprintf(out,"%.1f \n",s[i]);
54
     fclose(out);
55
56
57
     return 0;
58 }
```



Energy of different realizations of random configuration (red points) vs. the energy of a completely ordered configuration (black line).



Visualization of a random configuration.

Solution to Problem 3.2 Ising model - Metropolis algorithm

Add the following modules to the lattice class:

file = lattice.cpp

```
double Lattice::get_magnetization(double s[]){
1
2
     // return the magnetization m of a given configuration
3
     int N = L*L;
4
     double m =0.0;
         (int i =0; i<N; i++){
5
         += s[i];
6
7
8
     m = fabs(m)/N;
9
     return m;
10 | }
```

```
11
12
   void Lattice::Metropolis (double s[], double T, int nmcstep){
     // perform nmcstep Monte Carlo time steps
13
     // using the Metropolis Algorithm
14
15
     // T = temperature
     int N = L*L;
16
17
     double beta = 1./T;
18
     double * precalc;
     precalc = dvector(-4,4);
19
     for (int i=-4; i<=4; i++){</pre>
20
21
       precalc[i] = exp(-beta*2*i);
22
23
     for (int t=0; t<nmcstep; t++){</pre>
24
       int i = rand() % N;
       double sum_neigbours = s[get_below(i)] + s[get_above(i)] +
25
                                s[get_right(i)] + s[get_left(i)];
26
27
       double E_in = -s[i] * sum_neigbours;
       double E_fin = s[i] * sum_neigbours;
28
29
       int j = int((E_fin-E_in)/2);
30
       if (j<=0) {
         s[i] = -s[i];
31
       } else {
32
          if (precalc[j] > rand()/((double) RAND_MAX)) {
33
34
            s[i] = -s[i];
35
36
       }
37
     }
```

Obviously also the *lattice.h* should be modified accordingly.

And wrap up a main file:

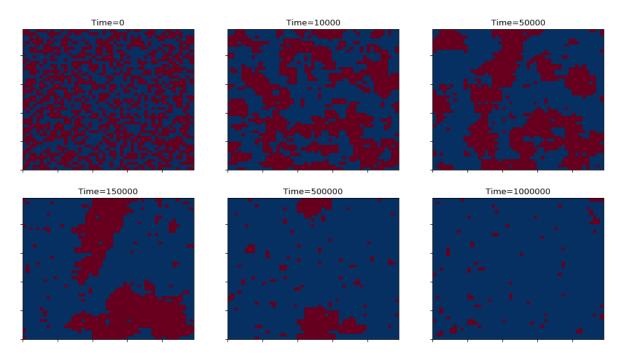
 $file = Ising_Metropolis.cpp$

```
1 // compile with:
2 |// g++ Ising_Metropolis.cpp lattice.cpp nrutil.cpp -o test -02 -w
3 | #include <stdio.h>
4 | #include <stdlib.h>
  | #include <math.h>
5
  #include <sstream>
6
   #include "Ising.h"
7
   #include "lattice.h"
8
9
   #include "nrutil.h"
10
  int main (){
11
12
    // set parameters
    L = 50;
13
     N = L*L;
14
     double T = 2.0;
15
16
     int nmcstep = 10000;
17
     printf("critical temperature = %.4f \n", 2./log(1.+sqrt(2.)));
18
     double ma = pow((1.-pow(sinh(2./T),-4.)), 0.125);
```

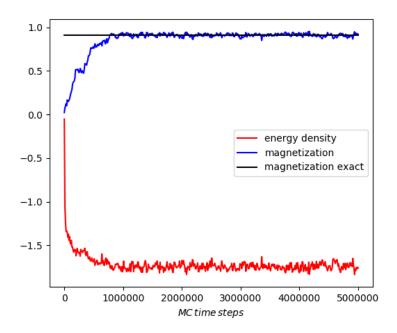
```
20
21
     // array of length N
22
     s = dvector(0,N-1);
     s_{initial\_values} = dvector(0,N-1);
23
24
     s0 = dvector(0, N-1);
     s1 = dvector(0,N-1);
25
26
     s2 = dvector(0,N-1);
     s3 = dvector(0, N-1);
27
28
     s4 = dvector(0, N-1);
     s5 = dvector(0, N-1);
29
30
     // random configuration
31
32
     for (int i =0; i<N; i++){</pre>
33
       double x = rand()/((double) RAND_MAX)-0.5;
34
       if (x>0) {
35
         s_initial_values[i]=1.0;
36
       } else {
37
         s_initial_values[i]=-1.0;
       }
38
     }
39
40
     class Lattice *lattice;
41
     lattice = new Lattice(N, s, s_initial_values);
42
43
     // open output file
44
     FILE * out;
     out = fopen("output3.dat","w");
45
46
     double m = lattice->get_magnetization(s);
47
     fprintf(out, "%d %.1f %.4f %.4f \n",0,lattice->get_energy(s), m, ma);
48
49
     for (int k =0; k<N; k++) {s0[k]=s[k];}</pre>
50
51
     // Monte Carlo - Metropolis rule - equilibration
52
     int 1=0;
53
     int j=1;
54
     int i = 0;
     while (i<3*j && l==1 || l==0){
55
56
57
       lattice -> Metropolis (s, T, nmcstep);
58
       m = lattice->get_magnetization(s);
       fprintf(out, "%d %.1f %.4f %.4f \n",
59
60
                              (i+1)*nmcstep,lattice->get_energy(s), m, ma);
       if (i==1) {
61
62
         for (int k =0; k<N; k++) {s1[k]=s[k];}</pre>
       } else if (i==5) {
63
         for (int k = 0; k < N; k++) {s2[k] = s[k];}
64
       } else if (i==15) {
65
          for (int k =0; k<N; k++) {s3[k]=s[k];}</pre>
66
       } else if (i==50) {
67
68
         for (int k =0; k<N; k++) {s4[k]=s[k];}</pre>
69
       } else if (i==100) {
70
          for (int k =0; k<N; k++) {s5[k]=s[k];}</pre>
71
72
```

```
73
        if (m>ma && l==0) {
74
          printf("relaxation time = %d \n", (i+1)*nmcstep);
75
          j=i;
76
          1=1;
        }
77
      }
78
79
      fclose(out);
80
81
      out = fopen("output4.dat","w");
      for (int i =0; i<N; i++) {</pre>
82
83
        fprintf(out, "%.1f %.1f %.1f %.1f %.1f \n",
                              s0[i],s1[i],s2[i],s3[i],s4[i],s5[i]);
84
85
      fclose(out);
86
87
      // Magnetization autocorrelation function
88
89
      //~ nmcstep=1000;
90
      //~ int nsave = 500000;
      //~ mvector = dvector(0,nsave-1);
91
      //\sim for (int k=0; k<nsave; k++){
92
93
        //~ lattice->Metropolis (s, T, nmcstep);
        //~ mvector[k] = lattice->get_magnetization(s);
94
      //~ }
95
      //~ out = fopen("output5.dat","w");
96
97
      //~ for (int k =0; k<nsave-10; k++){
98
        //\sim fprintf(out,"%d %f\n",k*nmcstep, autocorrelation (mvector,nsave,k));
      //~ }
99
100
      //~ fclose(out);
101
102
103
      return 0;
   }
104
105
    double mean (double x[],int NN){
106
107
      // define the mean function
108
      double xm=0.0;
      for (int i =0; i<NN; i++){</pre>
109
        xm += x[i];
110
111
      }
112
      xm = xm/NN;
113
      return xm;
114 }
115
   double variance (double x[],int NN){
116
117
      // define the variance function
118
      double xm=0.0;
      for (int i =0; i < NN; i++){</pre>
119
120
        xm += x[i]*x[i];
121
122
      xm = xm/NN - mean(x,NN)*mean(x,NN);
123
      return xm;
124 | }
125
```

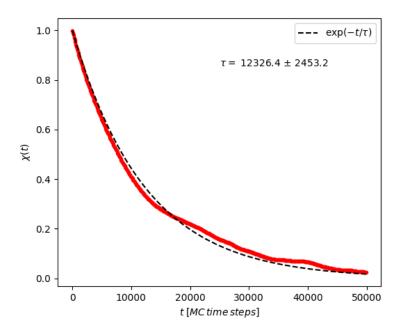
```
double autocorrelation (double x[], int NN, int tau){
126
127
      // define the autocorrelation function
128
      // input:
      //
129
         x = 1D array of data
      //
130
         tau = time lapse
131
      // output:
           c_n = autocorrelation of xs at time lapse n
132
      //
133
      double mu1 = mean (x,NN);
      double sigma2 = variance (x,NN);
134
135
      double c = 0.0;
      for (int i =0; i<NN-tau; i++){</pre>
136
137
        c+=(x[i]-mu1)*(x[i+tau]-mu1);
138
139
      c = c/(NN-tau);
      c = c/sigma2;
140
141
      return c;
142 | }
```



Visualization of the evolution starting with a random configuration.



Magnetization and energy density as a function of time.



Autocorralation function with a fit to the exponential decay.

Some interesting tricks: A Monte Carlo code consists of a central core nucleus which is repeated a large number of times. Therefore we should try to speed up this core to have a very efficient program (with a fast program we can sample a large number of different configurations

and thus get an accurate result).

- A common mistake is to compute the exponentials at each Metropolis step. The $\exp(x)$ is an operation which requires a certain amount of polynomial approximation on a computer and it is very slow. Since in the Ising model there are only few discrete values for $\Delta E > 0$ (how many?). It is much faster to calculate these once and store these numbers in a small vector.
- To calculate ΔE we only need to know the 4 spins up, down, right and left of the given spin, because all the other spins are unchanged.
- Often one performs simulations at different temperatures, say $T_1 < T_2 < T_3 \dots$ To avoid long runs to reach equilibration one can take as begin configuration for the simulation at T_2 the last spin configuration of the run at T_1 , and so on.