

1 **Measurement of the Higgs boson
2 transverse momentum spectrum in
3 the WW decay channel at 8 TeV and
4 first results at 13 TeV**

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7 PhD Thesis

8

9

Abstract

10 The cross section for Higgs boson production in pp collisions is studied
11 using the $H \rightarrow W^+W^-$ decay mode, followed by leptonic decays of the
12 W bosons, leading to an oppositely charged electron-muon pair in the
13 final state. The measurements are performed using data collected by the
14 CMS experiment at the LHC with pp collisions at a centre-of-mass energy
15 of 8 TeV, corresponding to an integrated luminosity of 19.4 fb^{-1} . The
16 Higgs boson transverse momentum (p_T) is reconstructed using the lepton
17 pair p_T and missing p_T . The differential cross section times branching
18 fraction is measured as a function of the Higgs boson p_T in a fiducial
19 phase space defined to match the experimental acceptance in terms of
20 the lepton kinematics and event topology. The production cross section
21 times branching fraction in the fiducial phase space is measured to be
22 $39 \pm 8\text{ (stat)} \pm 9\text{ (syst)}\text{fb}$. The measurements are compared to theoretical
23 calculations based on the standard model to which they agree within
24 experimental uncertainties.

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Chapter 1.

¹⁰⁷ Electroweak and QCD physics at
¹⁰⁸ LHC

Chapter 2.

¹⁰⁹ The CMS experiment at the LHC

¹¹⁰ In this chapter, the main characteristics of the Large Hadron Collider (LHC particle
¹¹¹ accelerator and Compact Muon Solenoid (CMS) experiment are described.

¹¹² 2.1. The Large Hadron Collider

¹¹³ The LHC [1–4] at CERN, officially inaugurated on 21st October 2008, is the largest and
¹¹⁴ most powerful hadron collider ever built. Installed in the underground tunnel which hosted
¹¹⁵ the Large Electron Positron Collider (LEP) [5–7], the leptonic accelerator in operation until
¹¹⁶ 2nd November 2000, the LHC accelerator has the shape of a circle with a length of about
¹¹⁷ 27 km and is located underground at a depth varying between 50 m to 175 m, straddling the
¹¹⁸ Franco-Swiss border near Geneva. It is designed to collide two 7 TeV counter-circulating
¹¹⁹ beams of protons resulting in a center-of-mass energy of 14 TeV, or two beams of heavy
¹²⁰ ions, in particular lead nuclei at an energy of 2.76 TeV/nucleon in the center-of-mass frame.

¹²¹ The transition from a leptonic collider to a hadronic collider entailed the following
¹²² advantages: first, it has been possible to build a machine that having the same size of the
¹²³ previous one (and therefore accommodated in the same LEP tunnel, substantially reducing
¹²⁴ the cost and time of construction), could reach a higher energy in the center-of-mass
¹²⁵ frame. This is due to the much lower amount of energy loss through synchrotron radiation
¹²⁶ emitted by the accelerated particles, that is proportional to the fourth power of the ratio
¹²⁷ E/m between their energy and their mass. Secondly, the composite structure of protons
¹²⁸ compared to the elementary structure of electrons allows LHC to be able to simultaneously
¹²⁹ access a wider energy spectrum, despite the production of many low energies particles in a
¹³⁰ complex environment. This is a particularly important feature for a machine dedicated to
¹³¹ the discovery of “new” physics.

¹³² In Fig. 2.1 a schematic description of the accelerator complex installed at CERN is shown.
¹³³ The acceleration is performed in several stages [4]. The protons source is a *Duoplasmatron*:
¹³⁴ the protons are obtained by removing electrons from a source of hydrogen gas and then
¹³⁵ sent to the LINAC2, a 36 m long linear accelerator which generates a pulsed beam with

an energy of 50 MeV using Radio Frequency Quadrupoles (RFQ) and focusing quadrupole magnets. The beam is subsequently sent to the Proton Synchrotron Booster (PSB), a circular accelerator consisting of four superimposed synchrotron rings with a circumference of about 160 m, which increases the proton energy up to 1.4 GeV. Then, protons are injected into the Proton Synchrotron (PS), a single synchrotron ring with a circumference of about 600 m where the energy is increased to 25 GeV. The sequential combination of these two synchrotrons also allows to create a series of protons bunches interspersed by 25 ns as required for the correct operation of LHC. The final proton injection stage is the Super Proton Synchrotron (SPS), a synchrotron with a circumference of approximately 7 km where protons reach an energy value of 450 GeV. Subsequently, protons are extracted and injected into the LHC ring via two transmission lines, to generate two beams running in opposite directions in two parallel pipes and which are accelerated up to the energy of interest. In the two pipes an ultrahigh vacuum condition is maintained (about 10^{-10} Torr) to avoid the spurious proton interactions with the gas remnants. At full intensity, each proton beam consists of 2808 bunches and each bunch contains around 10^{11} protons. The beams are squeezed and collide for a length of about 130 m at four interaction points where the four main experiments (ALICE, ATLAS, CMS and LHCb) are placed:

- CMS (Compact Muon Solenoid) [8] and ATLAS (A Toroidal LHC ApparatuS) [9] are two general-purpose detectors designed to investigate the largest possible spectrum of physics. In particular, they have been devoted to the detection of particles produced by a Higgs boson decay and to look for any evidence of possible new physics. The use of two detectors chasing the same objectives but designed independently is crucial for a cross-check of any possible new discovery;
- LHCb (LHC beauty) [10] is an experiment primarily designed to study CP (combined Charge conjugation and Parity symmetry) violation in electroweak interactions and to study asymmetries between matter and antimatter through the analysis of rare decays of hadrons containing b quarks. The detector is also able to perform measurements in the forward region, at small polar angles with respect to the beam line;
- ALICE (A Large Ion Collider Experiment) [11] is an experiment studying heavy ions collisions, through the production of a new state of matter called quark-gluon plasma.

Two other smaller experiments are located along the circumference of the LHC accelerator, TOTEM and LHCf, which focus on particles emitted in the forward direction. TOTEM (TOTal Elastic and diffractive cross section Measurement) [12] measures the proton-proton interaction cross section and accurately monitors the luminosity of the LHC using detectors positioned on either side of the CMS interaction point. LHCf (LHC forward) [13] is made up of two detectors which sit along the LHC beamline, at 140 m either side of the ATLAS collision point. It makes use of neutral particles thrown forward by LHC collisions as a source to simulate the interaction with the atmosphere of very high energy cosmic rays (between 10^{17} TeV and 10^{20} TeV) in laboratory conditions.

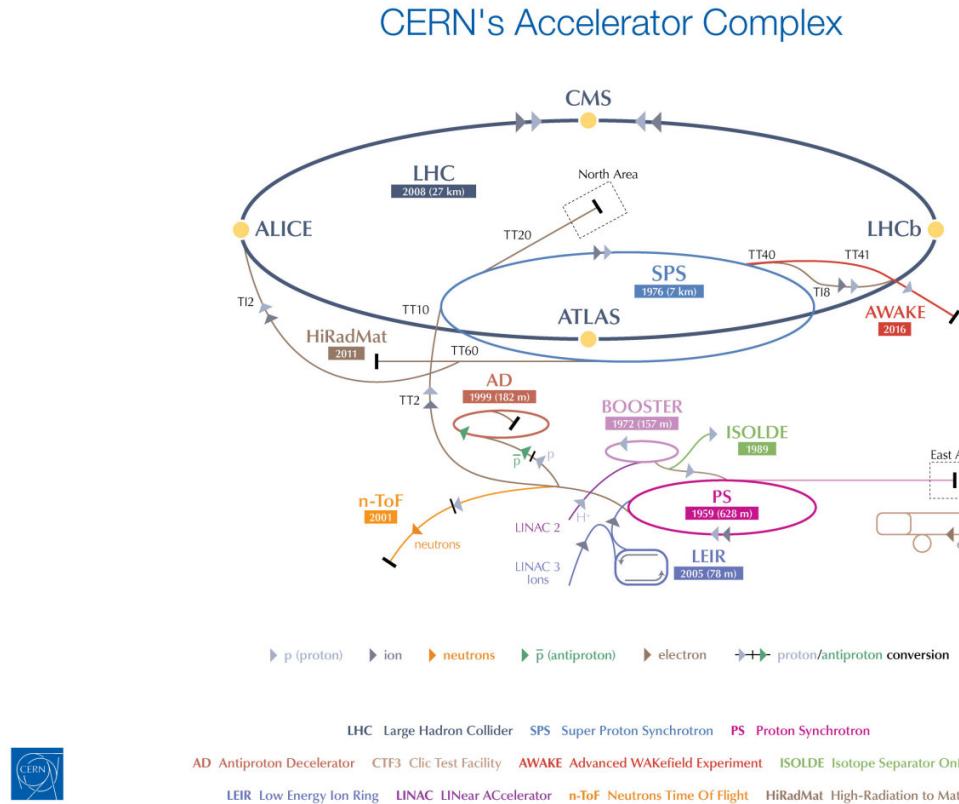


Figure 2.1.: Schematic description of the accelerator complex installed at CERN.

A series of about 1200 magnetic dipoles bend the beams along the accelerator ring. They are located along the “arc” structures of the circumference. The ring, in fact, can be subdivided into octants, with eight curve regions (the “arcs”) separated by rectilinear regions. In these straight regions, instead, almost 400 focusing and defocusing quadrupoles are located, which maintain the beam stable along the orbit, and some other small multipolar magnets (sextupoles and octupoles) are used to make additional minor corrections to the beam direction. A radio frequency acceleration system, consisting of 16 superconducting radio-frequency resonant cavities, is used to increase the proton energy by 0.5 MeV with each beam revolution. The 7 TeV per-beam-energy limit on the LHC is not determined by the electric field generated by the radiofrequency cavity but by the magnetic field necessary to maintain the protons in orbit, given the current technology for the superconducting magnets, which is about 5.4 T on average.

One of the most important parameters of an accelerator is the instantaneous luminosity \mathcal{L} , which gives a measure of the rate of events one can expect given the process cross section. In fact, for a given physics process with cross section σ , producing N events for unit of

time, the instantaneous luminosity is defined by the following equation:

$$N = \sigma \mathcal{L} . \quad (2.1)$$

The LHC design luminosity is $\mathcal{L} = 10^{34} \text{cm}^{-2}\text{s}^{-1}$, leading to around 1 billion proton interactions per second.

The instantaneous luminosity is a parameter which depends on the construction characteristics of the accelerator, and can be expressed by the following approximated formula:

$$\mathcal{L} = f \frac{n_1 n_2}{4\pi\sigma_x\sigma_y} , \quad (2.2)$$

where n_1 and n_2 are the number of particles contained in the two bunches colliding at a frequency f , and σ_x and σ_y are the beam sizes in the transverse plane. At LHC, the bunches collide with $f = 40 \text{MHz}$ and the transverse size of the beam can be squeezed down to around $15 \mu\text{m}$. Then, the integrated luminosity L is defined as the time integral of the instantaneous luminosity:

$$L = \int \mathcal{L} dt . \quad (2.3)$$

The main parameters of the LHC machine are listed in Table 2.1.

The LHC started to be operative in September 2008 but, due to a faulty interconnection between two magnets which caused a helium leakage in the tunnel, the operation was stopped and restarted in March 2010. During 2010 and 2011 LHC ran successfully and provided proton proton collisions at a center-of-mass energy of 7TeV , delivering a total integrated luminosity of about 6.1fb^{-1} . The encouraging results in the Higgs boson search provided by the ATLAS and CMS Collaborations led to the decision of extending the data taking period to the end of 2012, and to increase the center-of-mass energy up to 8TeV . During 2012, LHC delivered to the experiments an integrated luminosity of 23.3fb^{-1} . After the first long shutdown (LS1), a two years period started in the early 2013 where the LHC operation stopped for maintenance and upgrade, the LHC started again delivering proton proton collisions on 3rd June 2015, at the new record center-of-mass energy of 13TeV . During the 2015 the LHC delivered an integrated luminosity of 4.2fb^{-1} . Nowadays, LHC is still colliding bunches of protons at $\sqrt{s} = 13 \text{TeV}$, reaching unprecedented instantaneous luminosities and delivering a total integrated luminosity of 31fb^{-1} . The cumulative delivered luminosity versus time for the different LHC data taking periods is shown in Fig. 2.2.

As the instantaneous luminosity increases, the probability of multiple proton proton interactions to occur in a single bunch crossing grows higher as well. In this instance, the main goal is the identification and reconstruction of a single primary collision where the physics event of interest occurs among the background of the additional proton proton interactions. Such backgrounds are due to processes occurring with very high probability,

Table 2.1.: LHC technical parameters for proton proton collisions.

Parameter	Value
Maximum dipole magnetic field	8.33 T
Dipole operating temperature	1.9 K
Beam energy at injection	450 GeV
Beam energy at collision (nominal)	7 TeV
Beam energy at collision (2012)	4 TeV
Beam energy at collision (2015–2016)	6.5 TeV
Maximum instantaneous luminosity (nominal)	$10^{34} \text{ cm}^{-2}\text{s}^{-1}$
Maximum instantaneous luminosity (2012)	$7.7 \cdot 10^{33} \text{ cm}^{-2}\text{s}^{-1}$
Maximum instantaneous luminosity (2015–2016)	$1.2 \cdot 10^{34} \text{ cm}^{-2}\text{s}^{-1}$
Number of bunches per proton beam (nominal)	2808
Number of bunches per proton beam (2012)	1380
Number of bunches per proton beam (2015–2016)	2220
Maximum number of protons per bunch	$1.69 \cdot 10^{11}$
Bunch separation in time (nominal)	25 ns
Bunch separation in time (2012)	50 ns
Bunch separation in time (2015–2016)	25 ns
Collision frequency (nominal)	40 MHz
Collision frequency (2012)	20 MHz
Collision frequency (2015–2016)	40 MHz
Energy loss per turn at 14 TeV	7 keV

²²¹ like the production of low- p_T jets. These additional collisions are known as pile up (PU).²²² During the LHC current run the average number of pile up events is 23, with some event ²²³ exhibiting over 45 pile up collisions.

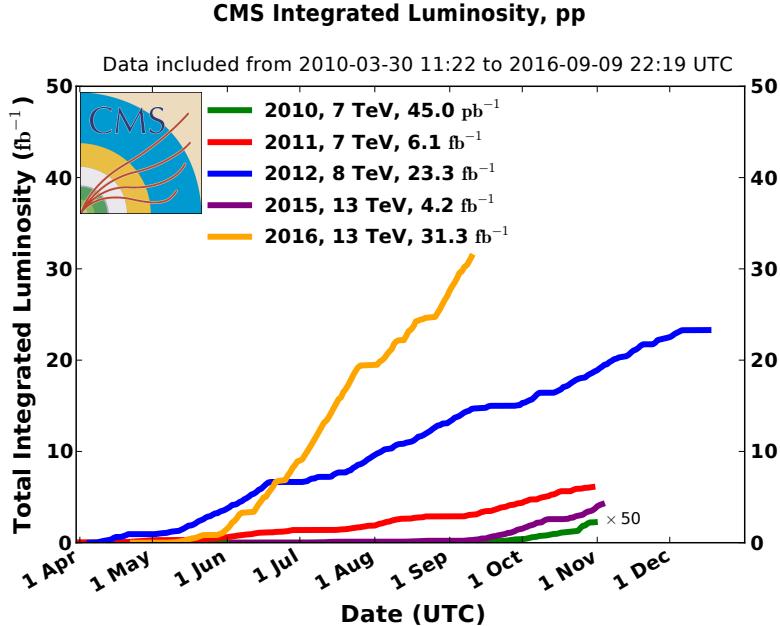


Figure 2.2.: Cumulative luminosity versus day delivered to CMS during proton proton collisions.

224 2.2. The *Compact Muon Solenoid* experiment

225 The CMS apparatus is a general purpose detector situated in one of the four LHC interaction
 226 points¹. The detector is designed to investigate a wide range of physics, from the search of
 227 the Higgs boson, to SM measurements and BSM physics searches. To achieve this goal, the
 228 detector is able to identify and reconstruct all the physics objects that may be produced
 229 in the proton proton collisions: electrons, muons, photons and jets. The main feature of
 230 the CMS detector is a superconducting solenoidal magnet which is capable to produce a
 231 3.8 T magnetic field. Such a strong magnetic field is the key aspect which permits to have
 232 a compact design of the detector. The detector has a cylindrical structure, which is typical
 233 of general purpose detectors, which consists of several cylindrical detecting layers, coaxial
 234 with the beam direction (*barrel* region), closed at both ends with detecting disks (*endcap*
 235 region), in such a way to ensure the hermetic closure of the apparatus.

236 The coordinate system used by CMS is a right-handed Cartesian system, with the
 237 origin in the nominal beam collision point inside the detector. The x -axis is chose to
 238 point radially towards the center of the LHC circumference and the y -axis is directed
 239 upwards along the vertical. The z -axis is oriented along the beam direction, according to
 240 the counter-clockwise direction of the LHC ring if seen from above. The CMS cylindrical
 241 symmetry and the Lorentz invariant description of the proton proton collisions, suggest the

¹The CMS detector is placed in a cavern 100 m underground in the area called Point 5, near the village of Cessy, in France.

²⁴² use of a pseudo-angular reference frame, described by the triplet of coordinates (r, ϕ, η) ,
²⁴³ where r is the distance from the z -axis, ϕ is the azimuthal angle, measured starting from
²⁴⁴ the x -axis positive direction, and η is the pseudorapidity, defined by the following equation:

$$\eta = -\ln \left(\tan \frac{\theta}{2} \right) , \quad (2.4)$$

²⁴⁵ where θ is the polar angle. The used of pseudorapidity is preferred over the polar angle
²⁴⁶ because differences in pseudorapidity are Lorentz invariant under boosts along the z -axis.
²⁴⁷ In the limit of ultrarelativistic particles the pseudorapidity coincides with the rapidity y :

$$y = \frac{1}{2} \ln \left(\frac{E + p_z}{E - p_z} \right) , \quad (2.5)$$

²⁴⁸ where E is the particle energy and p_z is the momentum projection along the z -axis.

²⁴⁹ The schematic view of the CMS detector, which has a length of 21.5 m, a diameter of
²⁵⁰ 15 m and a weight of about 14000 tons, is shown in Fig. 2.3. From the inner region to the
²⁵¹ outer one, the various CMS sub-detectors are:

- **Silicon tracker:** it occupies the region $r < 1.2$ m and $|\eta| < 2.5$. It is composed of an inner silicon pixel vertex detector and a surrounding silicon microstrip detector, with a total active area of about 215 m². It is used to reconstruct charged particle tracks and vertices;
- **Electromagnetic calorimeter (ECAL):** placed in the region $1.2 \text{ m} < r < 1.8$ m and $|\eta| < 3$, it consists of many scintillating crystals of lead tungstate (PbWO_4). It is used for the measurement of the trajectory and the energy released by electrons and photons;
- **Hadronic calorimeter (HCAL):** it is placed in the region $1.8 \text{ m} < r < 2.9$ m and $|\eta| < 5$. It is made up of brass layers alternated with plastic scintillators and it is used to measure the direction and energy deposited by the hadrons produced in the interactions;
- **Superconducting solenoidal magnet:** it occupies the region $2.9 \text{ m} < r < 3.8$ m and $|\eta| < 1.5$ and generates an internal uniform magnetic field with an intensity of 3.8 T, pointing along the direction of the beams. The magnetic field is necessary to bend the trajectories of charged particles, in order to allow the measurement of their momentum through the curvature observed in the tracking system. The magnetic field lines are closed by an external 21 m long iron yoke, that has a diameter of 14 m. Outside the return yoke, a residual 1.8 T magnetic field is present, pointing in the opposite direction with respect to the internal field;
- **Muon system:** the outermost system, which is placed in the region $4 \text{ m} < r < 7.4$ m and $|\eta| < 2.4$, has the purpose of reconstructing the tracks of muons passing through

274 it. It consists of Drift Tubes (DT) in the barrel region and Cathode Strip Chambers
 275 (CSC) in the endcaps. A complementary system of Resistive Plate Chambers (RPC)
 276 is used both in the barrel and endcaps. The muon chambers are housed inside the
 277 iron structure of the return yoke.

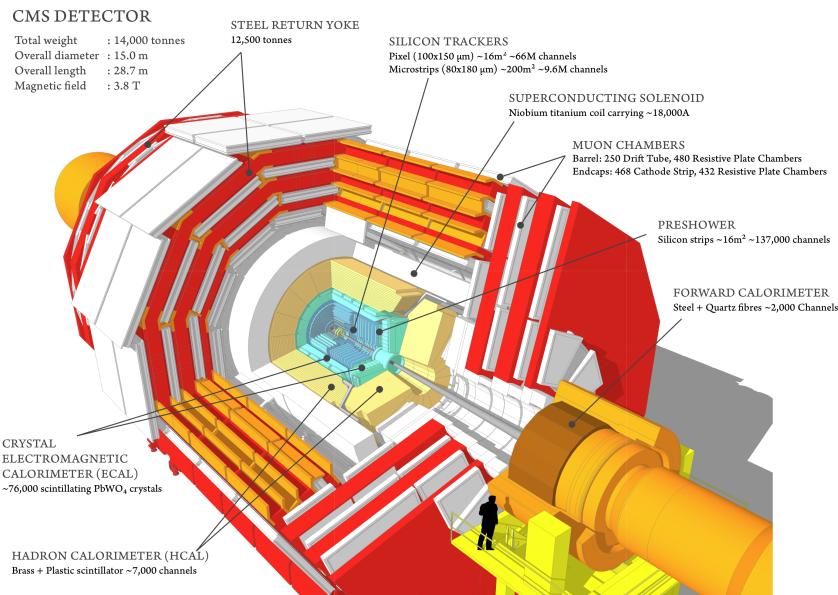


Figure 2.3.: Schematic view of the CMS detector showing its sub-detectors.

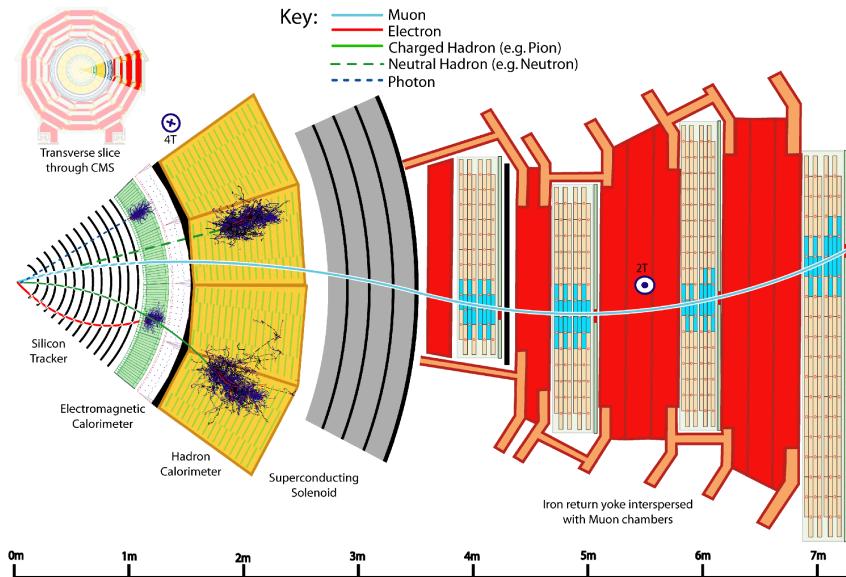


Figure 2.4.: Schematic view of a slice of the CMS detector, showing the sub-detectors response to the passage of different types of particles.

²⁷⁸ In Fig. ?? the response of the various CMS sub-detectors to the passage of different types
²⁷⁹ of particles is sketched. In the following sections a brief description of each sub-detector is
²⁸⁰ given.

281 2.2.1. The tracker**282 The pixel detector****283 The microstrip detector****284 2.2.2. The electromagnetic calorimeter (ECAL)****285 2.2.3. The hadronic calorimeter (HCAL)****286 2.2.4. The solenoid****287 2.2.5. The muon system****288 2.3. The CMS trigger system****289 2.3.1. The Level-1 trigger****290 2.3.2. The high level trigger (HLT)****291 2.4. Objects definition and event reconstruction****292 2.4.1. The Particle Flow technique****293 2.4.2. Leptons reconstruction and identification****294 Electrons****295 Muons****296 2.4.3. Jets reconstruction and identification****297 Jet b tagging****298 2.5. The CMS framework**

Chapter 3.

²⁹⁹ Higgs boson properties in the ³⁰⁰ $H \rightarrow WW$ decay channel

³⁰¹ 3.1. Higgs boson phenomenology at LHC

³⁰² The discovery of a new boson consistent with the standard model (SM) Higgs boson has
³⁰³ been reported by ATLAS and CMS Collaborations in 2012. The discovery has been followed
³⁰⁴ by a comprehensive set of studies of properties of this new boson in several production and
³⁰⁵ decay channels and no evidence of deviation from the SM expectation has been found so
³⁰⁶ far.

³⁰⁷ 3.1.1. Higgs boson production mechanisms

³⁰⁸ 3.1.2. Higgs boson decay channels

³⁰⁹ 3.2. The $H \rightarrow WW$ decay channel

³¹⁰ The CMS studies in the $H \rightarrow WW \rightarrow 2\ell 2\nu$ decay channel include the measurement of the
³¹¹ Higgs properties, as well as constraints on the Higgs total decay width and gauge bosons
³¹² anomalous couplings.

³¹³ **3.2.1. The H → WW → 2ℓ2ν signature**

³¹⁴ **3.2.2. Physics objects definition**

³¹⁵ **3.2.3. Background processes**

³¹⁶ **3.3. Experimental results of the Higgs boson measurements**

Chapter 4.

Measurement of the Higgs boson transverse momentum at 8 TeV using $H \rightarrow WW \rightarrow 2\ell 2\nu$ decays

4.1. Introduction

The Higgs boson production at hadron colliders is characterized by p_T^H and η . The η distribution is essentially driven by the PDF of the partons in the colliding hadrons, and it is only mildly sensitive to radiative corrections. The p_T^H distribution is instead sensitive to QCD radiative corrections. Considering the ggH production mode, at LO in perturbation theory, $\mathcal{O}(\alpha_s^2)$, the Higgs boson is always produced with p_T^H equal to zero. Indeed in order to have p_T different from zero, the Higgs boson has to recoil at least against one parton. Higher order corrections to the ggH process are numerically large and are known at NLO including full top quark mass dependence [14, 15], and at NNLO using the so-called large- m_t approximation [16–18], in which the top quark mass is assumed to be very large and the fermionic loop is replaced by an effective vertex of interaction. Starting from the NLO, the Higgs boson can be produced recoiling against other final state partons, resulting in a finite p_T^H . For this reason the LO process for Higgs production at $p_T \neq 0$ is at $\mathcal{O}(\alpha_s^3)$, and the counting of perturbative orders differs between inclusive Higgs boson production and p_T^H distribution. Also, NNLO QCD corrections in the p_T^H observable have recently been shown [19].

When $p_T^H \sim m_H$ the QCD radiative corrections to p_T^H differential cross section are theoretically evaluated using fixed-order calculations. When $p_T^H \ll m_H$ the perturbative expansion does not converge due to the presence of large logarithmic terms of the form $\alpha_s^n \ln^{2n} m_H^2/p_T^2$, leading to a divergence of $d\sigma/dp_T$ in the limit of $p_T \rightarrow 0$. For computing the p_T^H spectrum in this region soft-gluon resummation techniques are used, and matched to the fixed-order calculation in the $p_T^H \sim m_H$ region. For the p_T^H differential cross section the large- m_t calculation is a crude approximation, since it is known that the top quark mass has a non-negligible effect on the shape of the spectrum. Moreover the inclusion of the

³⁴⁵ bottom quark contribution in the fermionic loop can significantly modify the p_T^H shape [20],
³⁴⁶ as shown in Fig. 4.1. Hence, a precise experimental measurement of the p_T^H spectrum is
³⁴⁷ important to test the existing SM calculations.

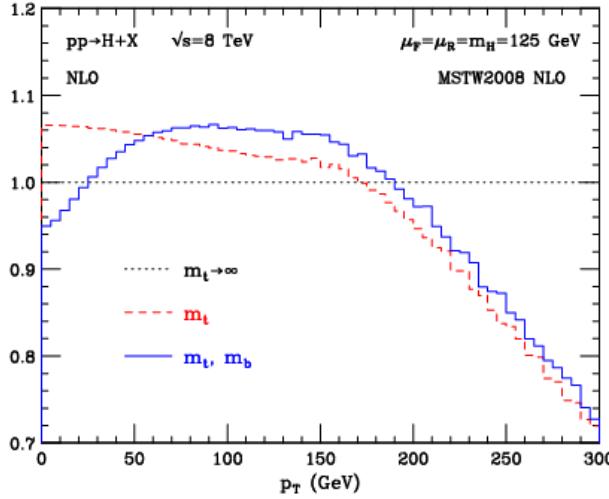


Figure 4.1.: p_T^H distribution computed at NLO (α_s^3) and normalized to the calculation obtained in the large- m_t approximation. The red dashed line corresponds to the calculation including the top quark mass while the blue line refers to the calculation including also the bottom quark effects.

³⁴⁸ Possible extensions of the SM predict a modification of the Higgs boson couplings to
³⁴⁹ gluons and to the top quark. Many of these models actually predict the existence of new
³⁵⁰ states that interact with the SM Higgs boson, but are beyond the direct production reach at
³⁵¹ the actual LHC energies. The effect of these new states could however show up as a deviation
³⁵² of the Higgs boson couplings with respect to the SM expectation. The modification of
³⁵³ the couplings, as shown in Refs. [21, 22], can change the kinematics of the Higgs boson
³⁵⁴ production and the effect can be particularly sizeable in the tail of the p_T^H distribution.
³⁵⁵ Other models, such as Composite Higgs [23], predict the existence of top-partners, which
³⁵⁶ are heavy resonances with the same quantum numbers as the top quark, that can interact
³⁵⁷ with the Higgs boson in the ggH fermionic loop, changing the p_T^H shape with respect to
³⁵⁸ what the SM predicts [24]. The measurement of the p_T^H spectrum is thus a useful tool for
³⁵⁹ indirect searches of new particles predicted by theories beyond the SM.

³⁶⁰ Measurements of the fiducial cross sections and of several differential distributions, using
³⁶¹ the $\sqrt{s} = 8$ TeV LHC data, have been reported by ATLAS [25–27] and CMS [28, 29] for
³⁶² the $H \rightarrow ZZ \rightarrow 4\ell$ ($\ell = e, \mu$) and $H \rightarrow \gamma\gamma$ decay channels. In this chapter a measurement of
³⁶³ the fiducial cross section times branching fraction ($\sigma \times \mathcal{B}$) and p_T spectrum for Higgs boson
³⁶⁴ production in $H \rightarrow WW \rightarrow e^\pm \mu^\mp \nu\nu$ decays, based on $\sqrt{s} = 8$ TeV LHC data, is reported.

³⁶⁵ The analysis is performed looking at different flavour leptons in the final state in order
³⁶⁶ to suppress the sizeable contribution of backgrounds containing a same-flavour lepton pair
³⁶⁷ originating from Z boson decay.

368 Although the $H \rightarrow WW \rightarrow 2\ell 2\nu$ channel has lower resolution in the p_T^H measurement
 369 compared to the $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ \rightarrow 4\ell$ channels because of neutrinos in the final
 370 state, the channel has a significantly larger $\sigma \times \mathcal{B}$, exceeding those for $H \rightarrow \gamma\gamma$ by a factor
 371 of 10 and $H \rightarrow ZZ \rightarrow 4\ell$ by a factor of 85 for a Higgs boson mass of 125 GeV [30], and is
 372 characterized by good signal sensitivity. Such sensitivity allowed the observation of a Higgs
 373 boson at the level of 4.3 (5.8 expected) standard deviations for a mass hypothesis of 125.6
 374 GeV using the full LHC data set at 7 and 8 TeV [31].

375 The measurement is performed in a fiducial phase space defined by kinematic require-
 376 ments on the leptons that closely match the experimental event selection.

377 The effect of the limited detector resolution, as well as the selection efficiency with respect
 378 to the fiducial phase space are corrected to particle level with an unfolding procedure [32],
 379 as explained in Sec. 4.7.

380 4.2. Data sets, triggers and MC samples

381 This analysis relies on the published $H \rightarrow WW$ measurements [31] in terms of code, selections
 382 and background estimates for both the ggH and VBF production mechanisms.

383 4.2.1. Data sets and triggers

384 The data sets used for the analysis correspond to 19.4 fb^{-1} at $\sqrt{s} = 8 \text{ TeV}$ of integrated
 385 luminosity composed of the following CMS data taking periods during 2012: 2012A
 386 (892 pb^{-1}), 2012B (4440 pb^{-1}), and 2012C (6898 pb^{-1}) and 2012D (7238 pb^{-1}). Data have
 387 been checked and validated and only data corresponding to good data taking quality are
 388 considered. The $e^\pm \mu^\mp$ final state is considered in this analysis.

389 For the data samples, the events are required to fire one of the unprescaled single-electron,
 390 single-muon or muon-electron triggers. Due to the rather high LHC instantaneous luminosity
 391 the single-lepton triggers must have high HLT p_T thresholds, otherwise the rate of these
 392 triggers would be too large to be sustained. The double-lepton triggers allow to lower
 393 down the p_T thresholds while keeping a sustainable trigger rate, thus maintaining a good
 394 sensitivity to the Higgs boson signal, for which the lepton p_T can be rather small. A brief
 395 overview of the HLT p_T criteria on the leptons is given in Table 4.1. While the HLT lepton
 396 p_T thresholds of 17 and 8 GeV for the double lepton triggers accommodate the offline
 397 lepton p_T selection of 20 and 10 GeV, the higher p_T thresholds in the single lepton triggers
 398 help partially recovering double lepton trigger inefficiencies as a high p_T lepton is on average
 399 expected due to the kinematics of the Higgs decay.

Table 4.1.: Highest transverse momentum thresholds applied in the lepton triggers at the HLT level. Double set of thresholds indicates the thresholds for each leg of the double lepton triggers.

Trigger Path	8 TeV
Single-Electron	$p_T > 27$ GeV
Single-Muon	$p_T > 24$ GeV
Muon-Electron	$p_T > 17$ and 8 GeV
Electron-Muon	$p_T > 17$ and 8 GeV

400 The trigger is not simulated in MC samples but the combined trigger efficiency is
 401 estimated from data and applied as a weight to all simulated events. The trigger efficiency
 402 for single and double lepton triggers is calculated using a Tag and Probe technique separately
 403 for muons and electrons, in bins of η and p_T . The Tag and Probe method uses a known
 404 mass resonance (e.g. J/Ψ , Z) to select particles of the desired type, and probe the efficiency
 405 of a particular selection criterion on these particles. In general the “tag” is an object that
 406 passes a set of very tight selection criteria designed to isolate the required particle type.
 407 Tags are often referred to as a golden electrons or muons and the fake rate for passing
 408 tag selection criteria should be very small. A generic set of the desired particle type (i.e.
 409 with potentially very loose selection criteria) known as “probes” is selected by pairing
 410 these objects with tags such that the invariant mass of the combination is consistent with
 411 the mass of the resonance. Combinatorial backgrounds may be eliminated through any
 412 of a variety of background subtraction methods such as fitting, or sideband subtraction.
 413 The definition of the probe objects depend on the specifics of the selection criterion being
 414 examined. The simple expression to get the efficiency ϵ as a function of p_T and η is given
 415 below:

$$\epsilon(p_T, \eta) = \frac{N_{\text{pass}}^{\text{probe}}}{N_{\text{pass}}^{\text{probe}} + N_{\text{fail}}^{\text{probe}}} \quad (4.1)$$

416 For double lepton triggers the efficiency is calculated separately for each leg of the
 417 trigger and then combined together. In the calculation the efficiencies of the two trigger legs
 418 are considered as independent, given that the correlations are very small. The combined
 419 efficiency is then used as a kinematics-dependent weight to be applied on top of simulated
 420 events.

421 The event efficiency ϵ_{ev} for an event with two leptons to pass the single lepton trigger is
 422 given by the following formula:

$$\epsilon_{\text{ev}} = 1 - (1 - \epsilon_{S,\ell 1}) \cdot (1 - \epsilon_{S,\ell 2}) , \quad (4.2)$$

where $\epsilon_{S,\ell 1}$ and $\epsilon_{S,\ell 2}$ are the efficiencies for the leading and subleading lepton to pass the single lepton trigger. In other words, the dilepton event passes the single lepton trigger if either one of the two leptons passes the single lepton trigger, excluding the cases for which both leptons pass the trigger. For double lepton triggers, the event efficiency can be written as:

$$\epsilon_{\text{ev}} = \epsilon_{D,\ell 1}^{\text{lead}} \cdot \epsilon_{D,\ell 2}^{\text{trail}} + (1 - \epsilon_{D,\ell 1}^{\text{lead}} \cdot \epsilon_{D,\ell 2}^{\text{trail}}) \cdot \epsilon_{D,\ell 1}^{\text{trail}} \cdot \epsilon_{D,\ell 2}^{\text{lead}} , \quad (4.3)$$

where $\epsilon_{D,\ell 1}^{\text{lead(trail)}}$ is the efficiency of the first lepton to pass the leading (trailing) leg of the double lepton trigger, and $\epsilon_{D,\ell 2}^{\text{lead(trail)}}$ is the efficiency of the second lepton to pass the leading (trailing) leg of the double lepton trigger. The final event efficiency applied to reweight the events in simulation is given by the boolean OR of the event efficiencies corresponding to the single and double lepton triggers, which, using Eqs. (4.2) and (4.3), can be written as:

$$\begin{aligned} \epsilon_{\text{ev}} = & 1 - (1 - \epsilon_{S,\ell 1}) \cdot (1 - \epsilon_{S,\ell 2}) + \\ & + (1 - \epsilon_{S,\ell 1}) \cdot (1 - \epsilon_{S,\ell 2}) \cdot \\ & \cdot [\epsilon_{D,\ell 1}^{\text{lead}} \cdot \epsilon_{D,\ell 2}^{\text{trail}} + (1 - \epsilon_{D,\ell 1}^{\text{lead}} \cdot \epsilon_{D,\ell 2}^{\text{trail}}) \cdot \epsilon_{D,\ell 1}^{\text{trail}} \cdot \epsilon_{D,\ell 2}^{\text{lead}}] . \end{aligned} \quad (4.4)$$

The term that multiplies the double lepton trigger event efficiency is needed to ensure that the events passing the double lepton trigger do not pass also the single lepton trigger.

4.2.2. Monte-Carlo samples

Several Monte Carlo event generators are used to simulate the signal and background processes:

- The first version of the POWHEG program [33–37] (POWHEG V1) provides event samples for the $H \rightarrow WW$ signal for the gluon fusion (ggH) and VBF production mechanisms, as well as $t\bar{t}$ and tW processes [38], with NLO accuracy.
- The $qq \rightarrow W^+W^-$, Drell-Yan, ZZ, WZ, $W\gamma$, $W\gamma^*$, tri-bosons and $W+jets$ processes are generated using the MADGRAPH 5.1.3 [39] event generator.

- 443 • The $gg \rightarrow W^+W^-$ process is generated using the GG2WW 3.1 generator [40] and its
 444 cross section is scaled to the approximate NLO prediction [41, 42].

- 445 • The VH process is simulated using PYTHIA 6.426 [43].

446 For leading-order generators samples, the CTEQ6L [44] set of parton distribution functions
 447 (PDF) is used, while CT10 [45] is used for next-to-leading order (NLO) ones. Cross
 448 section calculations [46] at next-to-next-to-leading order (NNLO) are used for the $H \rightarrow WW$
 449 process, while NLO calculations are used for background cross sections. The $H \rightarrow WW$
 450 process simulation is reweighted so that the p_T^H spectrum and inclusive production cross
 451 section closely match the SM calculations that have NNLO+NNLL pQCD accuracy in the
 452 description of the Higgs boson inclusive production, in accordance with the LHC Higgs
 453 Cross Section Working Group recommendations [30]. The reweighting of the p_T^H spectrum
 454 is achieved by tuning the POWHEG generator, as described in detail in Ref. [47]. Cross
 455 sections computed with NLO pQCD accuracy [30] are used for the background processes.
 456 The contribution of the $t\bar{t}H$ production mechanisms is checked to be negligible in each
 457 bin of p_T^H (below 1%) and is not included among the different production mechanisms. In
 458 Fig. 4.2 the relative fraction of the four production mechanisms is shown for each p_T^H bin.

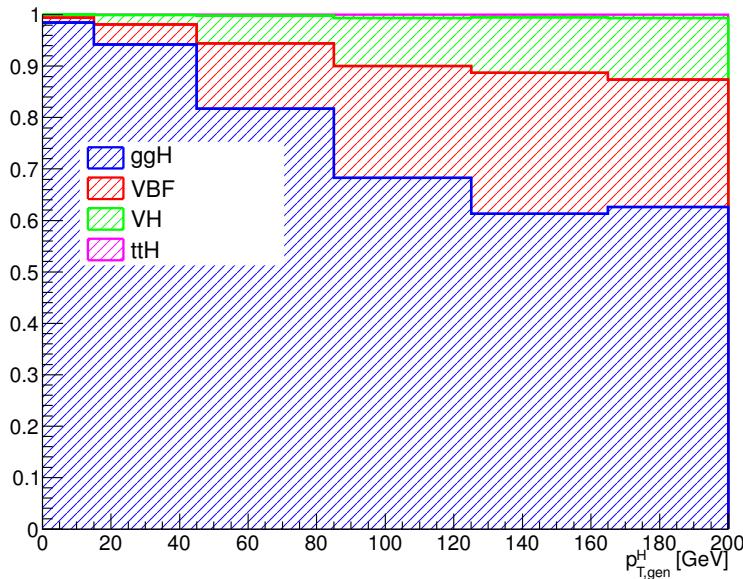


Figure 4.2.: Relative fraction of ggH, VBF, VH and ttH in each bin of the Higgs boson transverse momentum.

459 For all processes, the detector response is simulated using a detailed description of the
 460 CMS detector, based on the GEANT4 package [48].

461 Minimum bias events are superimposed on the simulated events to emulate the additional
 462 pp interactions per bunch crossing (pileup). The number of pile-up events simulated in the

463 MC samples (in the same bunch crossing, in time, or in the previous or following one, out
 464 of time pileup) have been generated poissonianly sampling from a distribution similar to
 465 what is expected from data. For a given range of analyzed runs, the mean number of pileup
 466 interactions per bunch crossing is estimated per luminosity block using the instantaneous
 467 luminosity provided by the LHC, integrated over the entire run range and normalized. The
 468 average number of pileup events per beam crossing in the 2011 data is about 10, and in the
 469 2012 data it is about 20.

470 The simulated events are reweighted to correct for observed differences between data
 471 and simulation in the number of pileup events, as shown in Fig. 4.3, trigger efficiency, and
 472 lepton reconstruction and identification efficiencies [31].

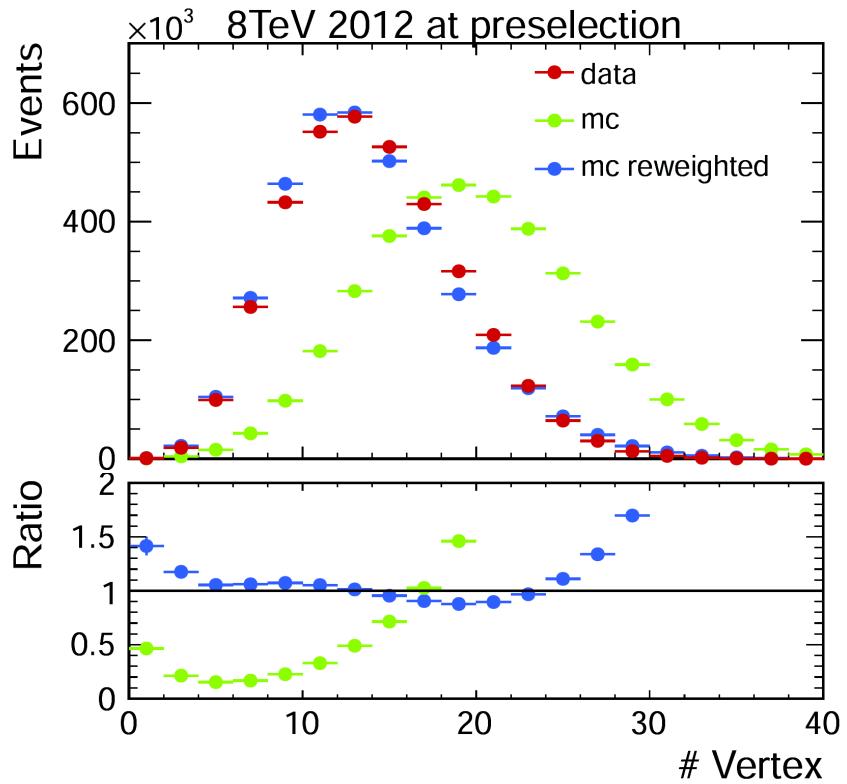


Figure 4.3.: Distribution of the number of vertices in data and in simulation, before and after applying the pile-up reweighting.

473 For the comparison of the measured unfolded spectrum with the theoretical predictions,
 474 two additional MC generators are used for simulating the SM Higgs boson production in
 475 the ggH process: HRES 2.3 [20, 49] and the second version of the POWHEG generator
 476 (POWHEG V2) [50]. HRES is a partonic level MC generator that computes the SM Higgs
 477 boson cross section at NNLO accuracy in pQCD and performs the NNLL resummation
 478 of soft-gluon effects at small p_T . The central predictions of HRES are obtained including
 479 the exact top and bottom quark mass contribution to the gluon fusion loop, fixing the

renormalization and factorization scale central values at a Higgs boson mass of 125 GeV. The cross section normalization is scaled, to take into account electroweak corrections, by a factor of 1.05 and the effects of threshold resummation by a factor of 1.06 [51, 52]. The upper and lower bounds of the uncertainties are obtained by scaling up and down both the renormalization and the factorization scales by a factor of two. The POWHEG V2 generator is a matrix element based generator that provides a NLO description of the ggH process in association with zero jets, taking into account the finite mass of the bottom and top quarks. The POWHEG prediction is tuned using the POWHEG damping factor $hdump$ of 104.17 GeV, in order to match the p_T^H spectrum predicted by HRES in the full phase space. This factor reduces the emission of additional jets in the high p_T regime, and enhances the contribution from the Sudakov form factor in the limit of low p_T . The POWHEG generator is interfaced to the JHUGEN generator version 5.2.5 [53–55] for the decay of the Higgs boson to a W boson pair and interfaced with PYTHIA 8 [56] for the simulation of parton shower and hadronization effects.

4.3. Analysis Strategy

The analysis presented here is based on that used in the previously published $H \rightarrow WW \rightarrow 2\ell 2\nu$ measurements by CMS [31], modified to be inclusive in the number of jets. This modification significantly reduces the uncertainties related to the modelling of the number of jets produced in association with the Higgs boson.

4.3.1. Event reconstruction and selections

The electron selection is based on two multivariate discriminants, one specialised in identifying the electron object and the other for isolation. The cut value for each discriminant is optimised to provide a good fake electron rejection and to improve the signal acceptance.

Muons are reconstructed using the standard CMS selection and are required to be identified both in the tracker (*Tracker Muon*) and in the muon chambers (*Global Muon*). Additionally quality criteria on the muon track are required, such as to have at least 10 hits in the tracker (at least one of which in the pixel detector) and to have $\chi^2/ndf < 10$. Muon isolation is based on the Particle-Flow algorithm. An MVA approach is considered, based on the radial distributions of the Particle-Flow candidates inside a cone of radius 0.5 around the muon direction.

The efficiencies for the identification and isolation of the electrons and muons are measured in data and in simulation selecting a pure sample of leptons coming from the $Z \rightarrow \ell\ell$ decay, and using a Tag and Probe technique very similar to the one described in Sec. 4.2.1 for the trigger efficiency. In this case, the probe lepton is defined by loose isolation and identification requirements and the efficiency to pass the tight analysis selections

515 is measured performing a simultaneous fit of signal plus background in two categories,
 516 corresponding to events in which the probe lepton pass or fail the analysis requirements.
 517 For the electrons, the resonant signal contribution in the fit is modelled as the convolution
 518 of a Breit-Wigner and a Crystal-Ball function. A polynomial function is added to take into
 519 account the tail in the low mass region. For muons the signal is fitted using the sum of two
 520 Voigtian functions. For both electrons and muons the background contribution is modelled
 521 as a third order Bernstein polynomial function. The efficiencies for data and simulation are
 522 extracted as parameters of the fit and are used as scale factors to correct the MC simulation
 523 to precisely model the data.

524 **This part would probably end up in the Object Reconstruction chapter**

525 Jets in this analysis are reconstructed by combining the energy measured in the calorimeters
 526 and tracks from charged particles on basis of the standard CMS particle flow algorithm
 527 and using the anti- k_T clustering algorithm with $R = 0.5$. Events will be classified into zero
 528 jet, one jet and VBF topologies by counting jets within $|\eta| < 4.7$ and for $p_T > 30 \text{ GeV}$.

529 In addition to the standard CMS PF E_T^{miss} , in this analysis a *projected* E_T^{miss} variable
 530 is also used. The *projected* E_T^{miss} is defined as the component of \vec{p}_T^{miss} transverse to the
 531 nearest lepton if the lepton is situated within the azimuthal angular window of $\pm\pi/2$ from
 532 the \vec{p}_T^{miss} direction, or the E_T^{miss} itself otherwise. Since the E_T^{miss} resolution is degraded
 533 by pileup, the minimum of two projected E_T^{miss} variables is used: one constructed from
 534 all identified particles (full projected E_T^{miss}), and another constructed from the charged
 535 particles only (track projected E_T^{miss}).

536 Background events from $t\bar{t}$ and tW production are rejected applying a soft-muon veto
 537 and b-tagging veto. The former selection requires that in the event there are no muons
 538 from b-decays passing the following cuts:

- 539 • the muon is reconstructed as TrackerMuon;
- 540 • the number of hits of the muon in the Silicon Tracker is greater than 10;
- 541 • the transverse impact parameter of the muon is less than 0.2 cm;
- 542 • if $p_T > 20 \text{ GeV}$ then the muon is required to be non-isolated with $ISO/p_T > 0.1$.

543 The latter veto rejects events that contain jets tagged as b-jets using two different
 544 algorithms for high and low p_T jets. For jets with p_T between 10 and 30 GeV, the Track-
 545 Counting-High-Efficiency (TCHE) algorithm, with a cut at 2.1 on the discriminating
 546 variable, is applied. For jets above 30 GeV, a better performing algorithm, Jet-Probability
 547 (JP), is used. Jets are identified as b-jets by the JP algorithm if the discriminating variable
 548 has a value above 1.4. In the following a b-tagged jet is defined as a jet, within $|\eta| < 2.4$
 549 (b-tagging requires the tracker information), with a value of the discriminating variable
 550 above the mentioned thresholds for the two algorithms.

551 The event selection consists of several steps. The first step is to select WW-like events
 552 applying a selection that consists of the following set of cuts:

553 **1. Lepton preselection:**

- 554 • at least two opposite-charge and opposite-flavour ($e\mu$) isolated leptons recon-
 555 structed in the event;
- 556 • $|\eta| < 2.5$ for electrons and $|\eta| < 2.4$ for muons;
- 557 • $p_T > 20$ GeV for the leading lepton. For the trailing lepton, the transverse
 558 momentum is required to be larger than 10 GeV.

559 **2. Extra lepton veto:** the event is required to have two and only two opposite-sign
 560 leptons passing the lepton selection.

561 **3. E_T^{miss} preselection:** particle flow E_T^{miss} is required to be greater than 20 GeV.

562 **4. projected E_T^{miss} selection:** minimum projected E_T^{miss} required to be larger than
 563 20 GeV.

564 **5. Di-lepton mass cut:** $m_{\ell\ell} > 12$ GeV in order to reject low mass resonances and QCD
 565 backgrounds.

566 **6. Di-lepton p_T cut:** $p_T^{\ell\ell} > 30$ GeV.

567 **7. Transverse mass:** $m_T > 60$ GeV to reject Drell-Yan to $\tau\tau$ events.

568 In addition to the WW-like preselection other cuts are applied in order to reduce the top
 569 quark background (both $t\bar{t}$ and tW), which is one of the main backgrounds in this final
 570 state. Two different selections are used depending on the number of jets with $p_T > 30$ GeV
 571 in the event. This is done to suppress the top quark background both in the low p_T^H region,
 572 where 0-jets events have the largest contribution, and for higher p_T^H values where also
 573 larger jet multiplicity events are important. The selection for 0-jets events relies on the
 574 soft-muon veto and on a soft jets (with $p_T < 30$ GeV) anti b-tagging requirement. The
 575 latter requirement exploits the TCHE algorithm to reject soft jets that are likely to come
 576 from b quarks hadronization.

577 For events with a jet multiplicity greater or equal than one, a different selection is
 578 applied. In this case we exploit the good b-tagging performances of the JP tagger to reject
 579 all the jets with $p_T > 30$ GeV that are likely to come from a b quark hadronization. The
 580 analysis selection requires to have no events containing b-tagged jets with $p_T > 30$ GeV.

581 A cut-flow plot is reported in Fig. 4.4, showing the effect of each selection using signal
 582 and background simulations. In the first bin, labelled as “No cut”, no selection is applied
 583 and the bin content corresponds to the total expected number of events with a luminosity
 584 of 19.4 fb^{-1} . All the events in this bin have at least two leptons with a loose transverse
 585 momentum cut of 8 GeV. In the following bin the lepton cuts are applied, including the

586 requirement to have two opposite-sign and opposite-flavour leptons and the extra lepton
 587 veto. Then all the other selections are progressively reported, showing the effect of each cut
 588 on the background and signal yields. For each selection the expected signal over background
 589 ratio is also shown, which, after the full selection requirements, reach a maximum value of
 590 about 3%.

591 4.3.2. Fiducial phase space

592 The Higgs boson transverse momentum is measured in a fiducial phase space, whose
 593 requirements are chosen in order to minimize the dependence of the measurements on the
 594 underlying model of the Higgs boson properties and its production mechanism.

595 The exact requirements are determined by considering the two following correlated
 596 quantities: the reconstruction efficiency for signal events originating from within the fiducial
 597 phase space (fiducial signal efficiency ϵ_{fid}), and the ratio of the number of reconstructed
 598 signal events that are from outside the fiducial phase space (“out-of-fiducial” signal events)
 599 to the number from within the fiducial phase space. The requirement of having a small
 600 fraction of out-of-fiducial signal events, while at the same time preserving a high value of the
 601 fiducial signal efficiency ϵ_{fid} , leads to a loosening of the requirements on the low-resolution
 602 variables, E_T^{miss} and m_T , with respect to the analysis selection.

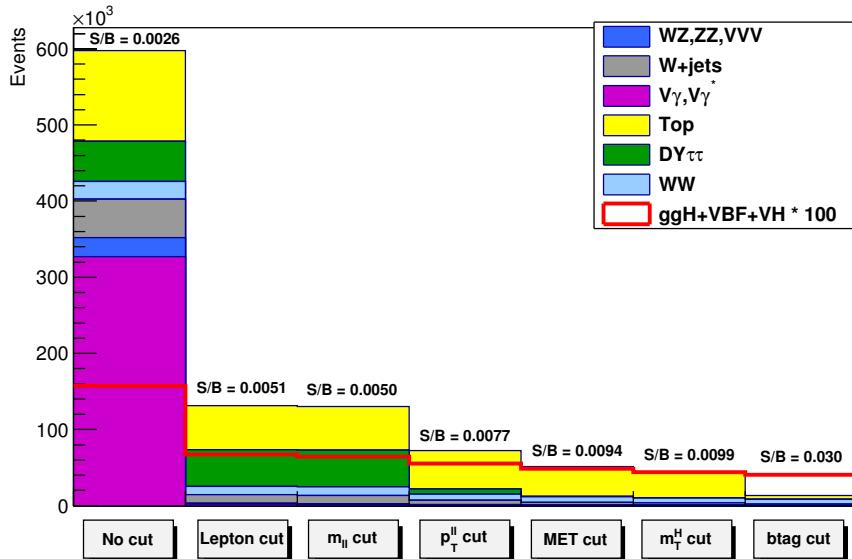


Figure 4.4.: Effect of single selections on MC samples. The signal (red line) is multiplied by 100 and superimposed on stacked backgrounds. In each bin, corresponding to a different selection, is reported the expected number of events in MC at a luminosity of 19.46 fb^{-1} .

603 The fiducial phase space used for the cross section measurements is defined at the
 604 particle level by the requirements given in Table 4.2. The leptons are defined as Born-level
 605 leptons, i.e. before the emission of final-state radiation (FSR), and are required not to
 606 originate from leptonic τ decays. The effect of including FSR is evaluated to be of the
 607 order of 5% in each p_T^H bin. For the VH signal process, the two leptons are required to
 608 originate from the $H \rightarrow WW \rightarrow 2\ell 2\nu$ decays in order to avoid including leptons coming
 609 from the associated W or Z boson.

Table 4.2.: Summary of requirements used in the definition of the fiducial phase space.

Physics quantity	Requirement
Leading lepton p_T	$p_T > 20 \text{ GeV}$
Subleading lepton p_T	$p_T > 10 \text{ GeV}$
Pseudorapidity of electrons and muons	$ \eta < 2.5$
Invariant mass of the two charged leptons	$m_{\ell\ell} > 12 \text{ GeV}$
Charged lepton pair p_T	$p_T^{\ell\ell} > 30 \text{ GeV}$
Invariant mass of the leptonic system in the transverse plane	$m_T^{\ell\ell\nu\nu} > 50 \text{ GeV}$
E_T^{miss}	$E_T^{\text{miss}} > 0$

610 A detailed description of the fiducial region definition and its optimization is given in
 611 appendix A.

612 4.3.3. Binning of the p_T^H distribution

613 Experimentally, the Higgs boson transverse momentum is reconstructed as the vector sum
 614 of the lepton momenta in the transverse plane and E_T^{miss} .

$$\vec{p}_T^H = \vec{p}_T^{\ell\ell} + \vec{p}_T^{\text{miss}} \quad (4.5)$$

615 Compared to other differential analysis of the Higgs cross section, such as those in the ZZ
 616 and $\gamma\gamma$ decay channels, this analysis has to cope with the limited resolution due to the E_T^{miss}
 617 entering the transverse momentum measurement. The effect of the limited E_T^{miss} resolution
 618 has two main implications on the analysis strategy. The first one is that the choice of
 619 the binning in the p_T^H spectrum needs to take into account the detector resolution. The
 620 second implication is that migrations of events across bins are significant and an unfolding
 621 procedure needs to be applied to correct for selection efficiencies and bin migration effects.

622 Given these aspects the criterion that was used to define the p_T^H bin size is devised to
 623 keep under control the bin migrations due to the finite resolution. For any given bin i we
 624 can define the purity P_i on a signal sample as the number events that are generated and

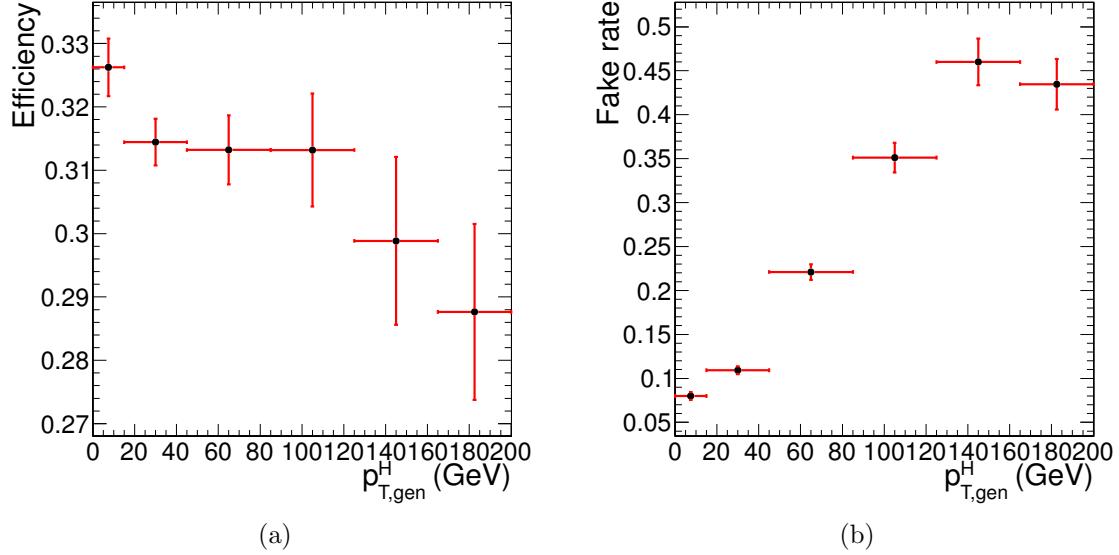


Figure 4.5.: Efficiency of the full selection (a) and fake rate (b) as a function of p_T^H .

625 also reconstructed in that bin, $N_i^{\text{GEN|RECO}}$, divided by the number of events reconstructed
 626 there, N_i^{RECO} :

$$P_i = \frac{N_i^{\text{GEN|RECO}}}{N_i^{\text{RECO}}} \quad . \quad (4.6)$$

627 The bin width is chosen in such a way as to make the smallest bins able to ensure a purity
 628 of about 60% on a ggH signal sample. Following this prescription we have divided the whole
 629 p_T^H range in the following six bins: [0-15 GeV], [15-45 GeV], [45-85 GeV], [85-125 GeV],
 630 [125-165 GeV], [165-∞ GeV].

631 The efficiency of the analysis selection with respect to the fiducial phase space is reported
 632 in Fig. 4.5 (a) for each p_T^H bin. The efficiency denominator is the number of events that
 633 are inside the fiducial phase space, while the numerator is the number of events that pass
 634 both the analysis and the fiducial phase space selections. The fake rate, defined by the
 635 ratio of signal events that pass the analysis selection but are not within the fiducial phase
 636 space, divided by the total number of events passing both the analysis and the fiducial
 637 phase space selections is shown in Fig. 4.5 (b). For both the selection efficiency and the
 638 fake rate, all the signal production mechanisms are included. The overall efficiency and
 639 fake rate are $\epsilon = 0.362 \pm 0.005$ and $\text{fake rate} = 0.126 \pm 0.004$ respectively, where only
 640 statistical uncertainties are taken into account.

641 If a 4π acceptance is defined, requiring just that the Higgs decays to WW and then to
 642 $2\ell 2\nu$, the efficiency becomes $\epsilon = 0.0396 \pm 0.0003$ and the fake rate is zero.

643 4.4. Background estimation

644 4.4.1. Top quark background

645 In this analysis the top quark background is divided into two different categories depending
 646 on the number of jets in the event. In the two categories different selections are applied,
 647 especially concerning the b-tagging requirements.

648 The general strategy for determining the residual top events in the signal region is to
 649 first measure the top tagging efficiencies from an orthogonal region of phase space in data.
 650 The orthogonal phase space is defined inverting the b-veto requirement of the signal region,
 651 in such a way to have a control region enriched in top quark events. Then, using this
 652 efficiency, the number of events with the associated uncertainty is propagated from the
 653 control region to the signal region. The number of surviving top events in the signal region
 654 would then be:

$$N_{b\text{veto}}^{\text{signal}} = N_{b\text{tag}}^{\text{control}} \cdot \frac{1 - \epsilon_{\text{top}}}{\epsilon_{\text{top}}} \quad (4.7)$$

655 where $N_{b\text{tag}}^{\text{control}}$ is the number of events in the control region and ϵ_{top} is the efficiency as
 656 measured in data.

657 The methods to estimate the top background contribution in the two jet categories are
 658 different and are explained below.

659 0-jets category

660 Most of the top background, composed of $t\bar{t}$ and tW processes, is rejected in the 0-jet bin
 661 by the jet veto. The top-tagging efficiency in the zero jet bin, $\epsilon_{\text{tag}}^{0\text{-jet}}$, is the probability for a
 662 top event to fail one of either the b-tagging veto or the soft muon veto, and is defined as:

$$\epsilon_{\text{tag}} = \frac{N_{\text{tag}}^{\text{control}}}{N^{\text{control}}} \quad , \quad (4.8)$$

663 where N^{control} is the number of events in the top control phase space defined requiring
 664 one b-tagged jet with $p_T > 30 \text{ GeV}$, and $N_{\text{tag}}^{\text{control}}$ is the subset of those events that pass
 665 either the soft muon tagging or the low- p_T b jet tagging. The purity of this control sample,
 666 as estimated from simulation, is about 97%. The remaining 3% background contribution is

667 estimated from simulation and subtracted from the numerator and denominator of Eq. (4.9).
668 The efficiency $\epsilon_{\text{top}}^{0-jet}$ can then be estimated using the following formula:

$$\epsilon_{\text{top}}^{0-jet} = f_{t\bar{t}} \cdot \epsilon_{2b} + f_{tW} \cdot (x \cdot \epsilon_{2b} + (1 - x) \cdot \epsilon_{\text{tag}}) \quad , \quad (4.9)$$

$$\epsilon_{2b} = 1 - (1 - \epsilon_{\text{tag}})^2 \quad , \quad (4.10)$$

669 where $f_{t\bar{t}}$ and f_{tW} are the $t\bar{t}$ and tW fractions respectively, x is the fraction of tW events
670 containing 2 b jets, and ϵ_{2b} is the efficiency for a top event with 0 counted jets, i.e. two soft
671 b jets, to pass the top veto. For the ratio of $t\bar{t}$ and tW cross-sections an uncertainty of 17%
672 is assumed. The fraction $f_{t\bar{t}}$ is estimated using MC simulation of the $t\bar{t}$ and tW processes
673 at NLO accuracy.

674 Using this procedure a data/simulation scale factor of 0.98 ± 0.17 is found, and is applied
675 to correct the MC simulation in order to match the data.

676 Category with more than 0 jets

677 The strategy for the estimation of the top background in events with at least one jet with
678 p_T greater than 30 GeV is the following. First of all the efficiency for tagging a b jet is
679 measured both in data and simulation and the values are used to correct the simulation for
680 different b-tagging efficiencies in data and simulation. This evaluation is performed in a
681 control region, called CtrlTP, containing at least two jets, using a Tag&Probe technique.
682 The procedure to extract these scale factors is presented in Sec. 4.4.1. Then a larger
683 statistics control region, CtrlDD, is defined by requiring at least one b-tagged jet and we
684 use the simulation, corrected for the previously computed b-tagging efficiency scale factor,
685 to derive the factor that connects the number of events in CtrlDD to the number of events
686 in the signal region. This second step is explained in detail in Sec. 4.4.1.

687 Tag&Probe

688 The Tag&Probe technique is a method to estimate the efficiency of a selection on data.
689 It can be applied whenever one has two objects in one event, by using one of the two,
690 the *tag*, to identify the process of interest, and using the second, the *probe*, to actually
691 measure the efficiency of the selection being studied. In our case we want to measure the
692 b-tagging efficiency, so what we need is a sample with two b-jets per event. The easiest
693 way to construct such a sample is to select $t\bar{t}$ events.

The CrtlTP control region is defined selecting the events which pass the lepton preselection cuts listed in Sec. 4.3.1, and have at least two jets with p_T greater than 30 GeV. One of the two leading jets is required to have a *JetBProbability* score higher than 0.5. From events in this control region we built *tag-probe* pairs as follows. For each event the two leading jets are considered. If the leading jet passes the *JetBProbability* cut of 0.5, that is considered a *tag*, and the sub-leading jet is the *probe*. In order to avoid any bias that could arise from the probe being always the second jet, the pair is tested also in reverse order, meaning that the sub-leading jet is tested against the *tag* selection, and in case it passes, then the leading jet is used as *probe* in an independent *tag-probe* pair. This means that from each event passing the CrtlTP cuts one can build up to two *tag-probe* pairs.

If the *tag* selection were sufficient to suppress any non top events, one could estimate the efficiency by dividing the number of *tag-probe* pairs in which the *probe* passes the analysis cut $JetBProbability > 1.4$ (*tag-pass-probe*) by the total number of *tag-probe* pairs. However this is not the case. In order to estimate the efficiency in the presence of background a variable that discriminates between true b-jets and other jets in a $t\bar{t}$ sample is chosen. The variable is the p_T of the *probe* jet. For real b-jets this variable has a peak around 60 GeV, while it does not peak for other jets. The idea is to fit simultaneously the p_T spectrum for *probe* jets in *tag-pass-probe* and *tag-fail-probe* pairs, linking together the normalizations of the two samples as follows:

$$N_{TPP} = N_s \epsilon_s + N_b \epsilon_b \quad (4.11)$$

713

$$N_{TFP} = N_s (1 - \epsilon_s) + N_b (1 - \epsilon_b) \quad (4.12)$$

where N_{TPP} is the number of *tag-pass-probe* pairs, N_{TFP} is the number of *tag-fail-probe* pairs, N_s is the number of *tag-probe* pairs in which the probe is a b-jet, N_b is the number of *tag-probe* pairs in which the probe is a not b-jet, ϵ_s is the b-tagging efficiency, ϵ_b is the probability of identifying as b-jet a non-b-jets, i.e. the mistag rate.

A χ^2 simultaneous fit of the *probe* p_T spectrum for *tag-pass-probe* and *tag-fail-probe* pairs is performed, deriving the shapes for true b-jets and non-b-jets from the simulation, and extracting N_s , N_b , ϵ_s and ϵ_b from the fit. The result of the fit on simulation is shown in Fig. 4.6. The relevant efficiencies are:

$$\epsilon_s^{MC} = 0.7663 \pm 0.0072 \quad (4.13)$$

722

$$\epsilon_b^{MC} = 0.208 \pm 0.015 \quad (4.14)$$

We have checked that these values are consistent with the true value for the b-tagging efficiency. The true value is computed by selecting jets that are matched within a cone of $\Delta R < 0.5$ with a generator level b-quark, and counting the fraction of those that pass

the *JetBProbability* cut of 1.4. This means that the *tag-probe* method does not introduce biases within the simulation statistic accuracy.

In order to assess the robustness of the fit, 5000 toy MC samples have been generated with a statistics equivalent to the one expected in data and the same fit is performed. All the 5000 fit succeeded, and the pull distributions for ϵ_s and ϵ_b parameters are shown in Fig. 4.7. The plots show the pull of the efficiencies measured in the fit, where the pull variable for each toy i is defined as:

$$pull(\epsilon_{s(b)}) = \frac{\epsilon_{s(b)}^{\text{true}} - \epsilon_{s(b)}^i}{\sigma(\epsilon_{s(b)}^i)} \quad (4.15)$$

The pulls are centered on 0 and have σ close to 1, as expected.

An example fit for one of the toys is shown in Fig. 4.8

Before running the fit on data, the shapes used in the fit have been validated. To do so, a purer top enriched phase space has been defined by requiring exactly two jets with *JetBProbability* score higher than 1.5 and no additional b-tagged jets, rejecting also jets with p_T smaller than 30 GeV. On this purer sample we have compared data against the shape used to fit the true b-jets in the *tag-pass-probe* distribution. The result is shown in Fig. 4.9 and shows good agreement.

Finally the fit has been performed on data, as shown in Fig. 4.10, providing the following efficiencies:

$$\epsilon_s^{Data} = 0.769 \pm 0.022 \quad (4.16)$$

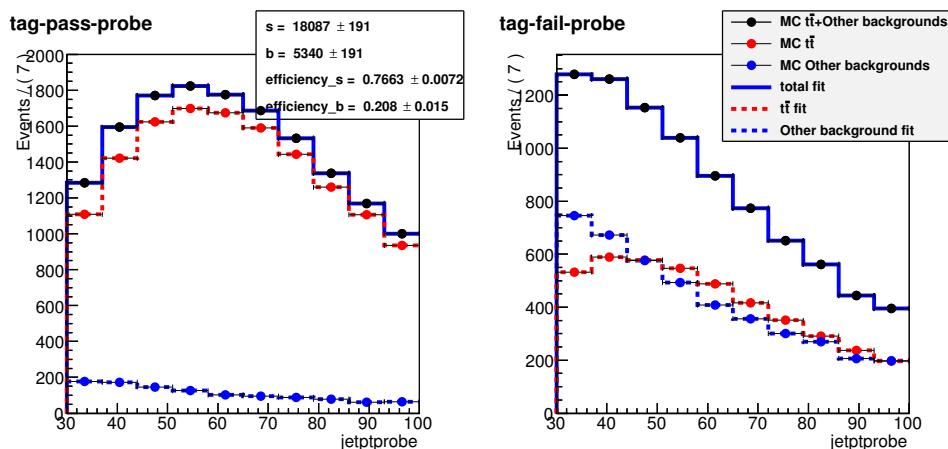


Figure 4.6.: Simultaneous fit of the *tag-pass-probe* and *tag-fail-probe* pairs in the MC.

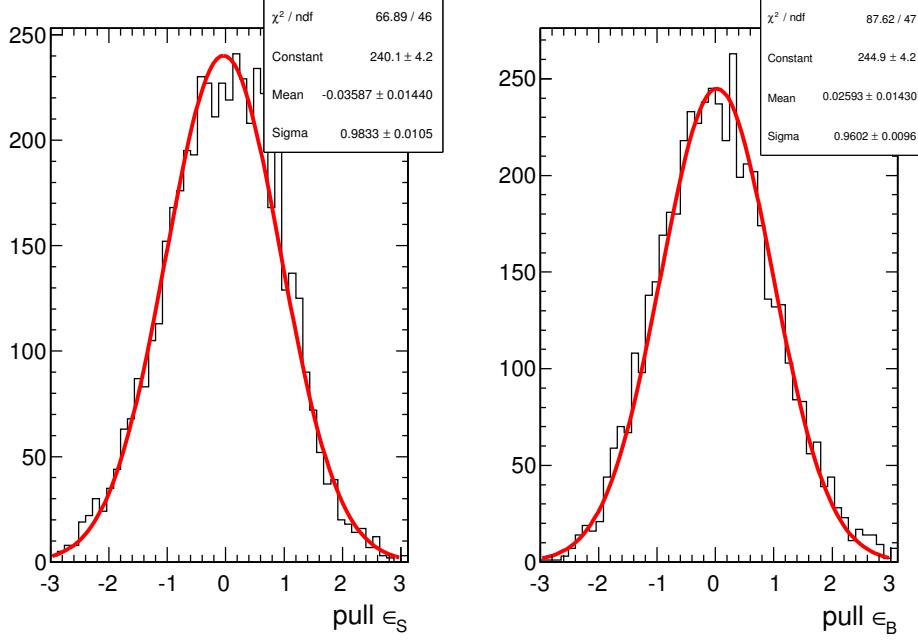


Figure 4.7.: Pulls of the ϵ_s and ϵ_b parameters in 5000 toy MC.

$$\epsilon_b^{Data} = 0.121 \pm 0.054 \quad (4.17)$$

743 Further studies have been performed to assess the effect of the relative uncertainty on
 744 the $t\bar{t}$ and tW event fractions. The same procedure described above has been applied to
 745 different simulation templates obtained varying the $t\bar{t}$ and tW fractions within theoretical
 746 uncertainties, and the effect on the parameters extracted with the fit procedure is found to
 747 be well below the fit uncertainties.

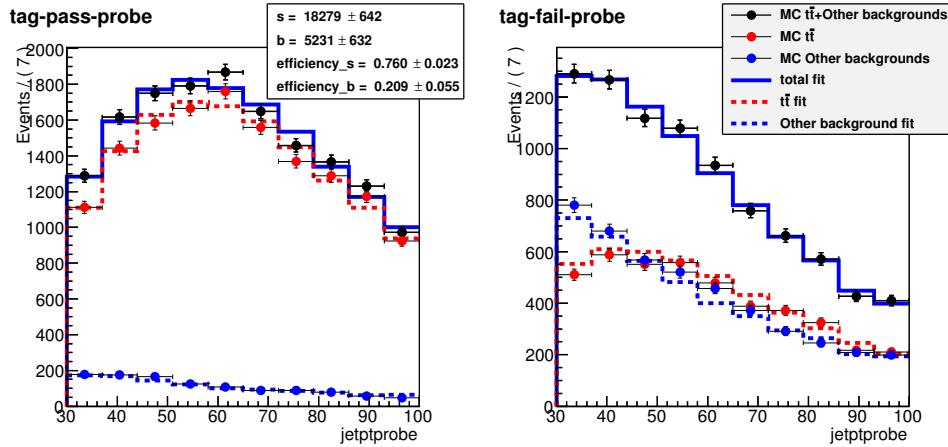


Figure 4.8.: Fit of a toy MC sample.

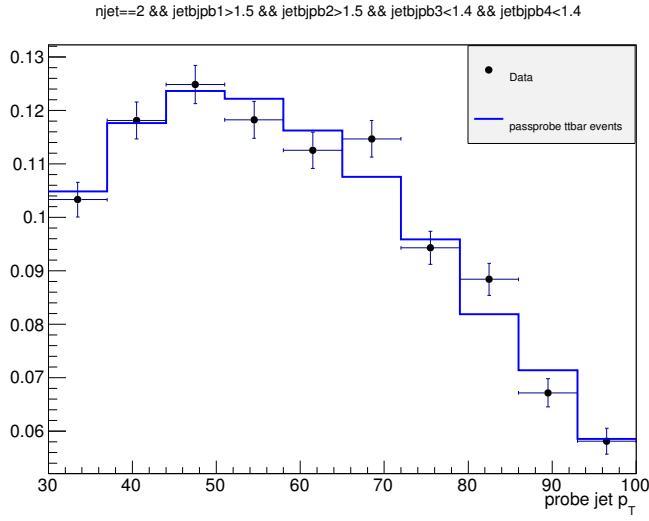


Figure 4.9.: Shape comparison for the *probe p_T* spectrum in data and in MC in a very pure $t\bar{t}$ sample.

748 Data driven estimation

749 In addition to the b-tagging efficiency, the other ingredient to estimate the $t\bar{t}$ background is
 750 the process cross section. The idea is to measure the cross section in a $t\bar{t}$ enriched control
 751 region, that is called CtrlDD. CtrlDD is defined according to the lepton preselection cuts
 752 defined in Sec. 4.3.1, and requiring in addition at least one jet with *JetBProbability* score
 753 higher than 1.4.

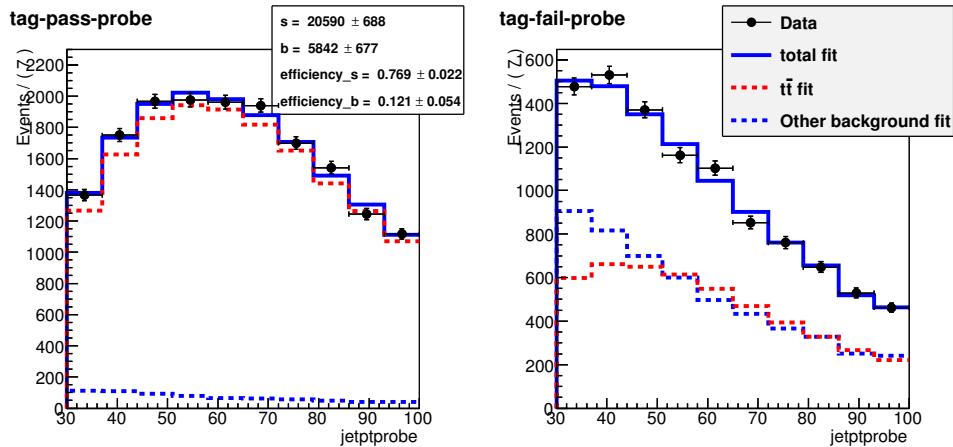


Figure 4.10.: Simultaneous fit of the *tag-pass-probe* and *tag-fail-probe* pairs in data.

From the simulation we derive the factor α that connects CtrlDD to the signal region, calculating the ratio of $t\bar{t}$ events in the two regions:

$$\alpha = \frac{N_{tt\ MC}^{SIG}}{N_{tt\ MC}^{CtrlDD}} . \quad (4.18)$$

The number of events in the CtrlDD region in data is counted, subtracting the expected number of events from non- $t\bar{t}$ backgrounds, and obtaining $N_{tt\ Data}^{CtrlDD}$. Finally the number of expected $t\bar{t}$ events in the signal region ($N_{tt\ Data}^{SIG}$) is obtained as:

$$N_{tt\ Data}^{SIG} = \alpha N_{tt\ Data}^{CtrlDD} . \quad (4.19)$$

In evaluating α and its error the b-tagging efficiencies determined in Sec. 4.4.1 are used. For each event an efficiency scale factor and a mistag rate scale factor are derived, depending on whether the event falls in the signal or CtrlDD region.

$$SF_{SIG} = \left(\frac{1 - \epsilon_s^{Data}}{1 - \epsilon_s^{MC}} \right)^{\min(2, n_{b-jets})} \left(\frac{1 - \epsilon_b^{Data}}{1 - \epsilon_b^{MC}} \right)^{n_{non-b-jets}} \quad (4.20)$$

$$SF_{CtrlDD} = \left(\frac{\epsilon_s^{Data}}{\epsilon_s^{MC}} \right)^{(jet1 == b-jet)} \left(\frac{\epsilon_b^{Data}}{\epsilon_b^{MC}} \right)^{(jet1 == non-b-jets)} \quad (4.21)$$

where n_{b-jets} is the number of true b-jets in the event and $n_{non-b-jets}$ is the number of non-b-jets in the event. The writing $jet1 == b-jet$ ($jet1 == non-b-jets$) is a boolean flag that is true when the leading jet, the one used for the CtrlDD selection, is (not) a true b-jet.

Since the efficiency and mistag rate that have been measured on data are close to the one in the simulation, it was decided to assume a scale factor of 1 for both b-tagging efficiency and mis-tag rate. This means that the central values of the scale factors defined in Eq. 4.20 and Eq. 4.21 is 1, but these numbers have an error that is derived assuming an uncertainty on ϵ_s^{Data} and ϵ_b^{Data} that covers both the statistical error from the fit of the two quantities and the difference with respect to the simulation. This results in an up and a down variation of the scale factors in the signal and CtrlDD regions, that is used to derive an error on α .

p_T^H [GeV]	N_{CTRL}^{DATA}	N_{CTRL}^{TOP}	N_{SIG}^{TOP}	α	$\Delta\alpha$
[0–15]	406.71	358.78	117.83	0.328	0.075
[15–45]	2930.14	2703.44	859.08	0.318	0.071
[45–85]	5481.02	5207.48	1506.05	0.289	0.065
[85–125]	4126.35	4032.56	861.22	0.214	0.052
[125–165]	1612.64	1654.27	304.69	0.184	0.055
[165– ∞]	647.50	760.37	201.70	0.265	0.147

Table 4.3.: Data driven scale factors related to the top quark background estimation.

A data driven estimation of the top quark background with the method described above is performed in each of the p_T^H bins independently. The reason to make this estimation in p_T^H bins, rather than inclusively is explained in Fig. 4.11, where the p_T^H distribution is shown in the CtrlDD region normalized to the cross section measured by a specific CMS analysis [57]. As shown in the ratio plot, an overall normalization factor would not be able to accommodate for the variations of the data/simulation ratio from bin to bin.

The α factors for each bin and the number of events in signal, CtrlDD regions in MC as well as in data are listed in Tab. 4.3.

A comparison of the m_{ll} distribution in the six p_T^H bins used in the analysis in CtrlDD after the data driven correction is shown in Fig. 4.12

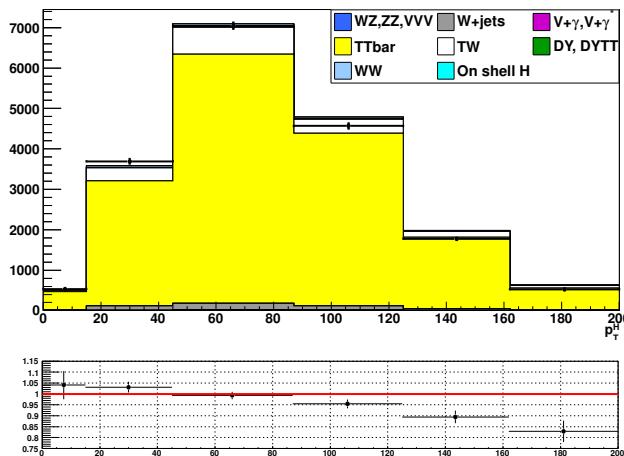


Figure 4.11.: p_T^H distribution in the CtrlDD control region.

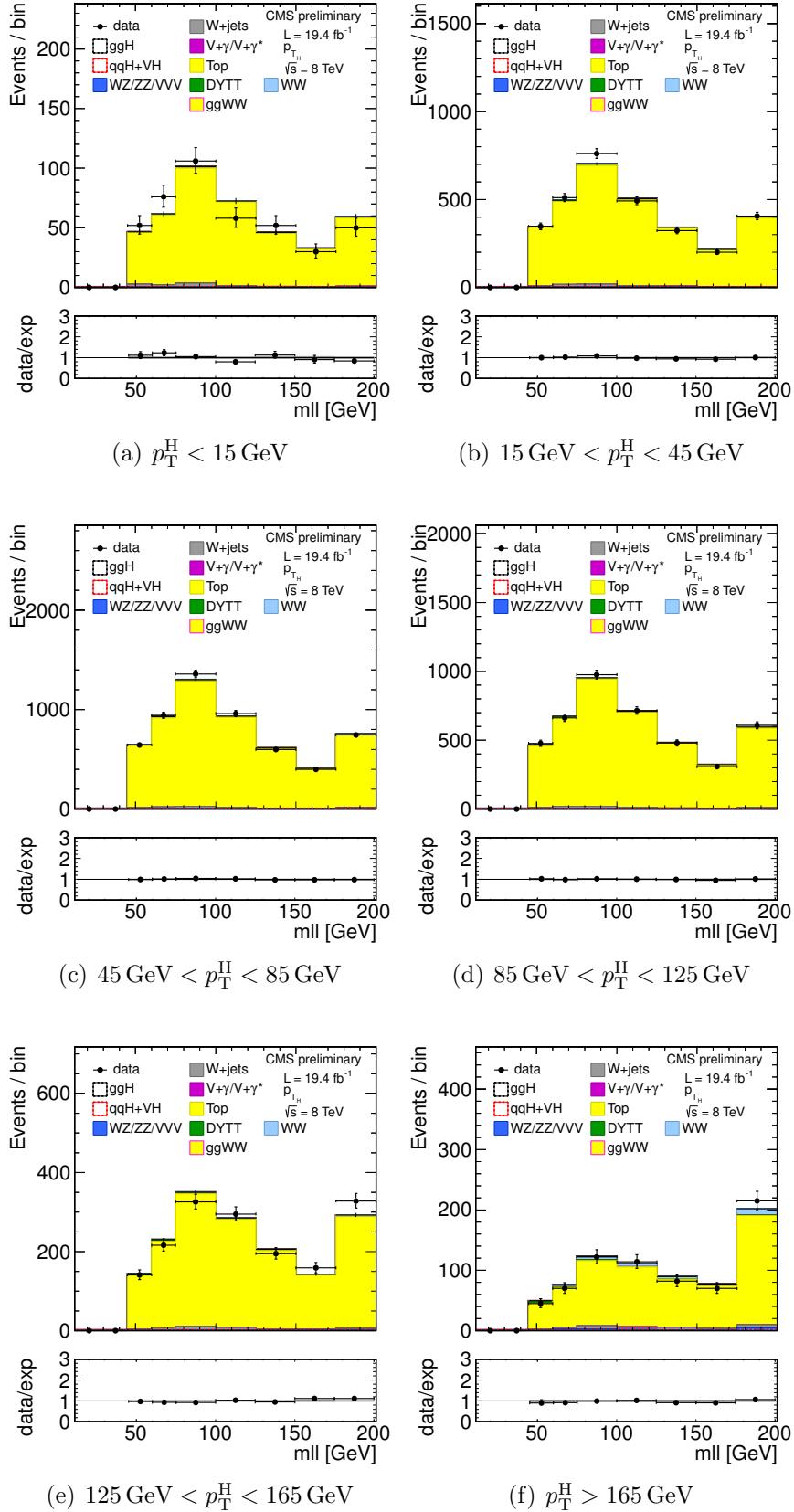


Figure 4.12.: $m_{\ell\ell}$ distributions in the CtrlDD region for the different p_T^H bins.

⁷⁸⁴ **4.4.2. WW background**

⁷⁸⁵ For what the $q\bar{q} \rightarrow W^+W^-$ background shape is concerned, the prediction from the simula-
⁷⁸⁶ tion is used. This background is divided into six different parts, corresponding to the six
⁷⁸⁷ p_T^H bins defined in the analysis. The normalization of the $q\bar{q} \rightarrow W^+W^-$ background is left
⁷⁸⁸ free to float in each bin, in such a way to adjust it in order to match the data during the fit
⁷⁸⁹ procedure. In this way the shape difference between the p_T^{WW} theory prediction and the
⁷⁹⁰ distribution provided by the simulation, which is obtained with the MADGRAPH generator,
⁷⁹¹ is minimized.

⁷⁹² In figure 4.13 a comparison is shown between the p_T^{WW} spectra of two different $q\bar{q} \rightarrow W^+W^-$
⁷⁹³ samples: one obtained with the MADGRAPH generator and the other after applying to the
⁷⁹⁴ same distribution a reweighting in order to match the theoretical prediction at NLO+NNLL
⁷⁹⁵ precision.

⁷⁹⁶ A shape discrepancy can be clearly observed and the effect becomes larger at high
⁷⁹⁷ values of p_T^H . In order to assess the effect of this discrepancy on the shapes of the variables
⁷⁹⁸ used for the signal extraction, $m_{\ell\ell}$ and m_T , the shapes have been checked in all p_T^H bins,
⁷⁹⁹ comparing different MC samples. The MADGRAPH sample used for the nominal shape is
⁸⁰⁰ compared to the MADGRAPH sample with NLO+NNLL reweighting, a POWHEG sample
⁸⁰¹ with NLO accuracy and an AMC@NLO sample. The results of this comparison are shown
⁸⁰² in figures 4.14 and 4.15. The shape discrepancy among the different models is included as
⁸⁰³ an additional systematic uncertainty.

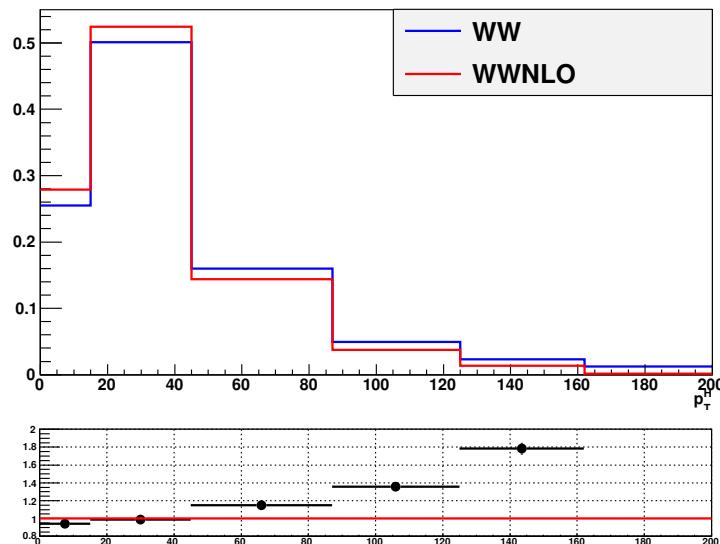


Figure 4.13.: Comparison between the p_T^{WW} distributions obtained with two different MC generators: the blue line corresponds to the MADGRAPH generator and the red line refers to the same sample in which a reweighting has been applied in order to match the theoretical prediction at NLO+NNLL precision.

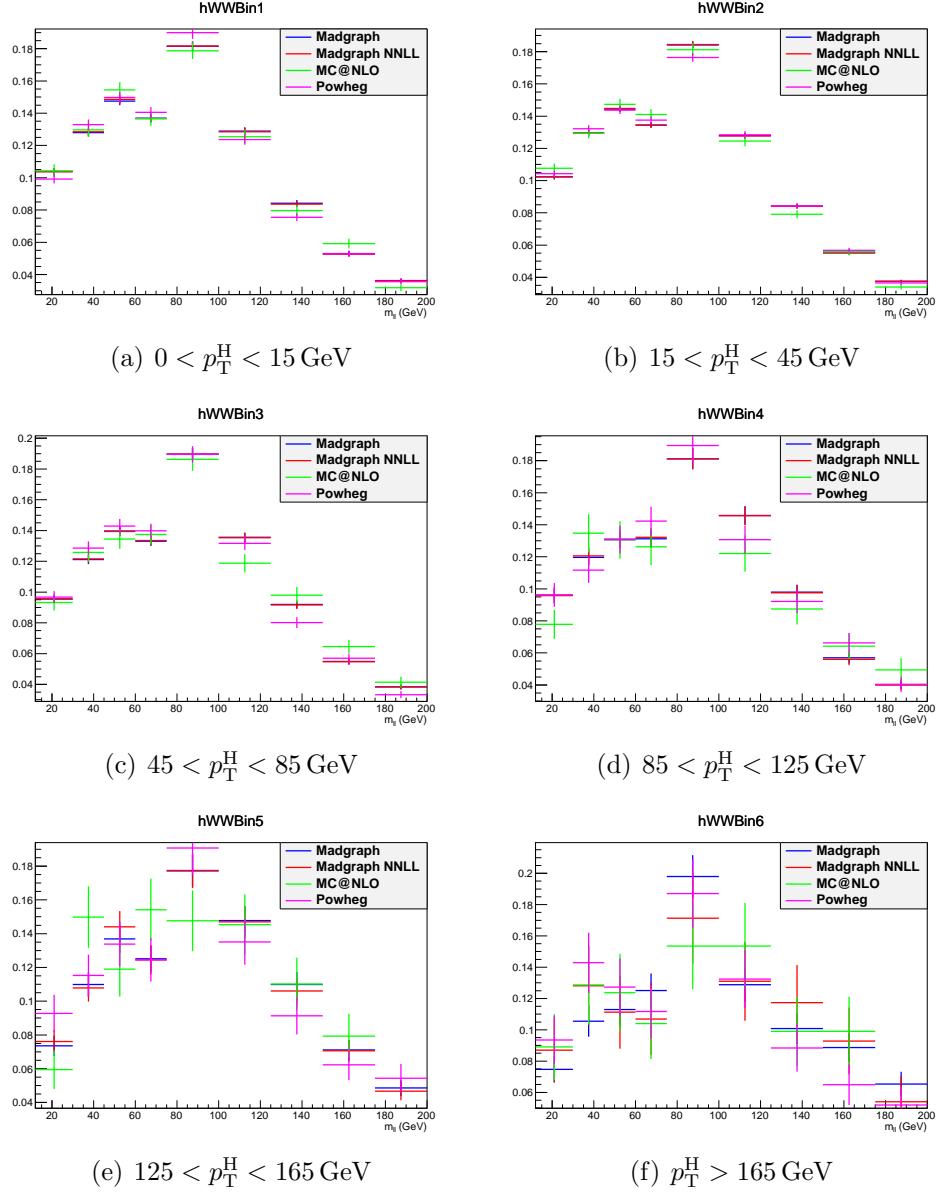


Figure 4.14.: Comparison between the default WW background sample and other theoretical models for the $m_{\ell\ell}$ distributions in every p_T^H bin.

The gluon-induced WW process, i.e. $gg \rightarrow W^+W^-$, has a sub-dominant contribution with respect to the quark-induced process, being the cross section ratio between the two of about 5%. The $m_{\ell\ell}$ and m_T shapes for this background are taken from simulation while the cross section is scaled to the approximate NLO calculation [41, 42].

The agreement of the $m_{\ell\ell}$ and m_T shapes between simulation and data for this background was checked in a signal-free control, defined selecting events with values of $m_{\ell\ell}$

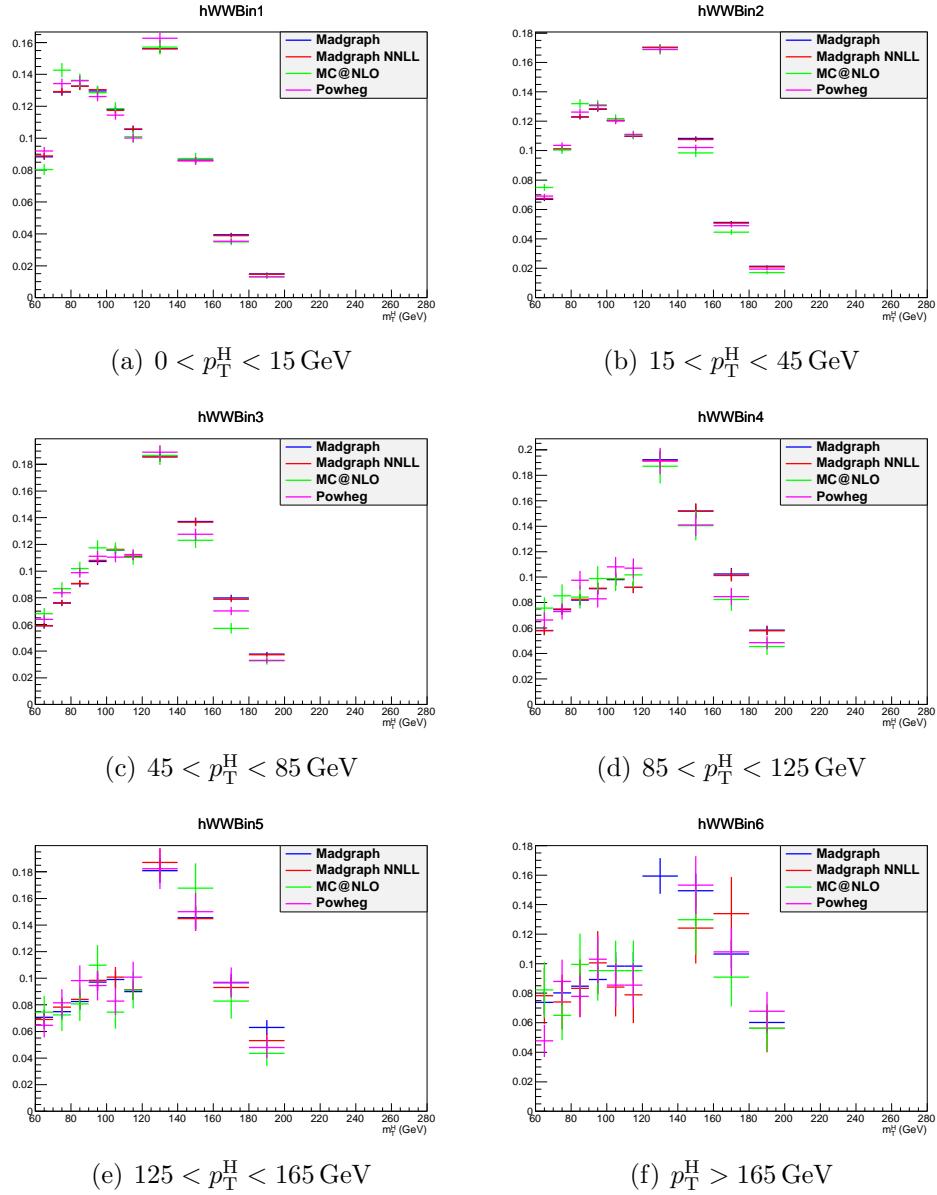


Figure 4.15.: Comparison between the default WW background sample and other theoretical models for the m_T distributions in every p_T^H bin.

810 greater than 70 GeV. A comparison of the $m_{\ell\ell}$ and m_T shapes in data and simulation is
 811 shown in Fig. 4.16 for events containing zero and one jets, inclusive in p_T^H .

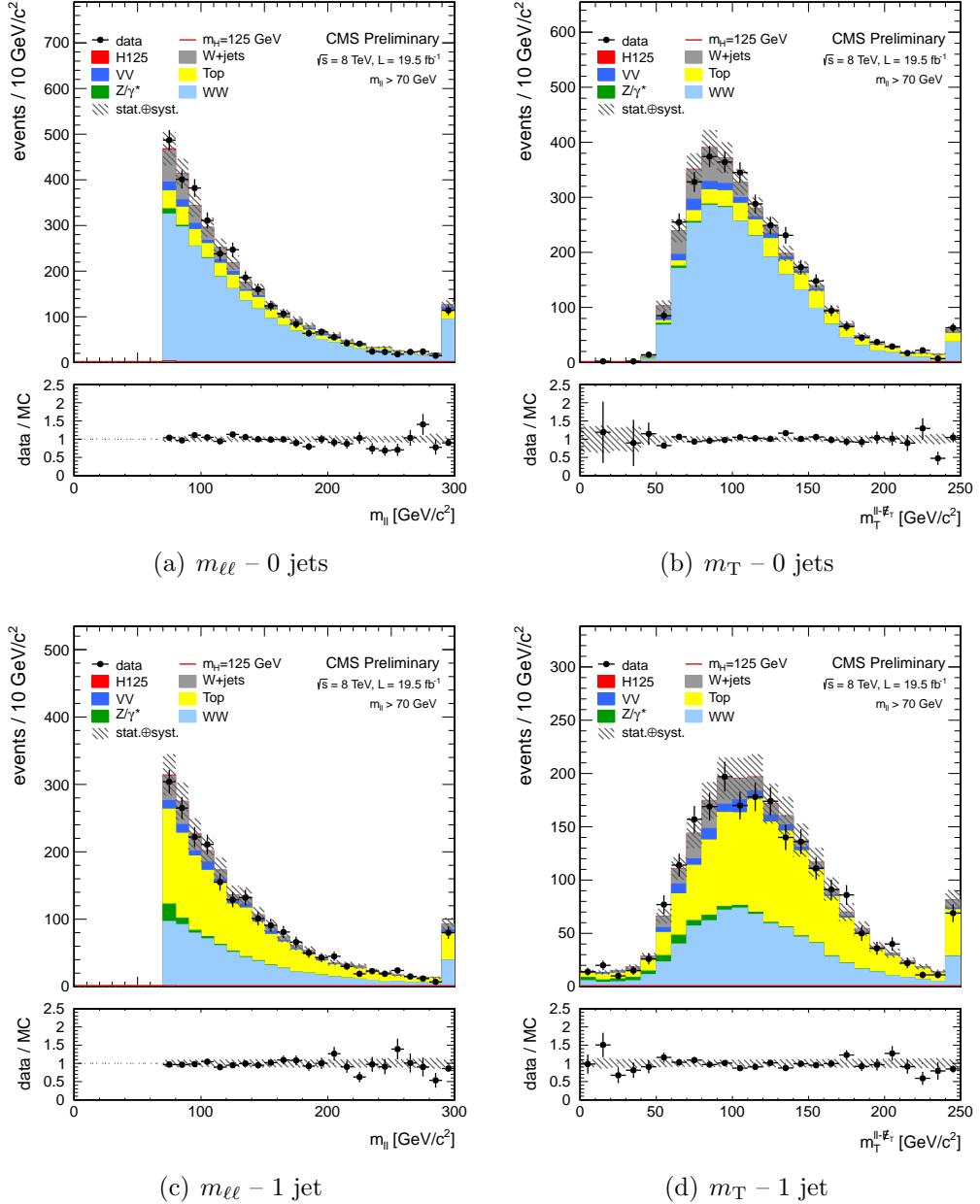


Figure 4.16.: Comparison of the $m_{\ell\ell}$ and m_T shapes in data and simulation for events with zero and 1 jets, inclusive in p_T^H . The events are required to pass the analysis requirements and, in order to define a signal-free control region, to have $m_{\ell\ell} > 70$ GeV.

812 4.4.3. Other backgrounds

813 W+jets background

814 Events in which W bosons are produced in association with jets, as well as multi-jet events,
 815 constitute a background for this analysis, because one or more jets can be misidentified as
 816 leptons. The rate at which jets are misidentified as leptons may be not accurately described
 817 in simulation, hence a data driven method is used to estimate this background.

818 The idea is to estimate the background containing one or two fake leptons selecting
 819 events with relaxed lepton quality criteria, i.e. looser with respect to the selections used at
 820 the analysis level, and computing the efficiencies for real and fake leptons to pass the tight
 821 lepton quality requirements of the analysis. A data-driven approach is pursued to estimate
 822 this background. A set of loosely selected lepton-like objects, referred to as the “fakeable
 823 object” or “denominator” from here on, is defined in a data set of events dominated by
 824 dijet production. To measure the fake rate we count how many fakeable objects pass the
 825 full lepton selection of the analysis, parameterized as a function of the phase space of the
 826 fakeable lepton, therefore it is extracted in bins of η and p_T . The ratio of the fully identified
 827 lepton, referred as “numerator”, to the fakeable objects is taken as the probability for a
 828 fakeable object to fake a lepton:

$$829 \text{Fake Rate} = \frac{\text{\#of fully reconstructed leptons}}{\text{\#of fakeable objects}} \quad (4.22)$$

829 It is then used to extrapolate from the loose leptons sample to a sample of leptons
 830 satisfying the full selection.

831 The definition of the denominator is of large impact in the systematic uncertainties
 832 related to this method. For the 2012 data taking period a summary of the selections used
 833 for the numerator and the denominator of Eq. (4.22) is shown below for electrons and
 834 muons respectively. For electrons the denominator is defined by the following requirements:

- 835 • $\sigma_{\text{inj}\eta} < 0.01(0.03)$ for barrel (endcap);
- 836 • $|\Delta\phi_{\text{in}}| < 0.15(0.10)$ for barrel (endcap);
- 837 • $|\Delta\eta_{\text{in}}| < 0.007(0.009)$ for barrel (endcap);
- 838 • $H/E < 0.12(0.10)$ for barrel (endcap);
- 839 • electron conversion rejection;
- 840 • $|d_0| < 0.02 \text{ cm};$

- 841 • $\frac{\sum_{\text{trk}} E_T}{p_T^{\text{ele}}} < 0.2$;
- 842 • $\frac{\sum_{\text{ECAL}} E_T}{p_T^{\text{ele}}} < 0.2$;
- 843 • $\frac{\sum_{\text{HCAL}} E_T}{p_T^{\text{ele}}} < 0.2$.

844 For muons the selection are loosened with respect to the tight analysis selection requiring
 845 that:

- 846 • $|d_0| < 0.02$ cm;
- 847 • MVA isolation output > -0.6 .

848 The dijet enriched data set used for the fake rate measurement, which is selected using
 849 single lepton triggers with low p_T thresholds, it is not a pure sample containing just fake
 850 leptons, but may still contain prompt leptons coming from the W and Z boson decays.
 851 To reject muons from the W decay, the events are required to have $E_T^{\text{miss}} < 20$ GeV and
 852 a W transverse mass below 20 GeV as well. Muons from the Z decay are instead remove
 853 requiring $m_{\mu\mu} > 20$ GeV and $m_{\mu\mu} \notin [76, 106]$ GeV. For electrons the Z mass peak veto
 854 is enlarged to $m_{ee} \notin [60, 120]$ GeV. Finally both electrons and muons are required to be
 855 isolated from the leading jet in the event, i.e. $\Delta\phi(\ell, j) > 1$. The residual prompt lepton
 856 contamination from EW processes such as W/Z+jets production, which can bias the fake
 857 rate measurement, is estimated using simulation and subtracted from both the numerator
 858 and denominator. The contamination from EW processes is different for the numerator
 859 and denominator and is particularly important for relatively high lepton p_T values.

860 In addition to the fake rate, also a prompt lepton rate is evaluated, defined as the
 861 probability of a prompt lepton passing the loose requirements to also pass the tight analysis
 862 selections. The prompt rate is also measured in data, defining a control region enriched in
 863 $Z \rightarrow \ell\ell$ events, selecting dilepton events with an invariant mass of the two leptons in the Z
 864 peak mass region.

865 Both the fake and prompt rate are used to reweight the data samples used in the analysis
 866 in order to obtain directly from data the contribution of the fake lepton background. The
 867 method to apply those rates is explained below in the simple case of just one lepton in
 868 the data sample, i.e. data selected by single lepton triggers, but can be straightforwardly
 869 generalized to situations with more than one lepton. Suppose that the total number of
 870 leptons passing the loose requirements, N_ℓ , is made up of N_p prompt and N_f fake leptons.
 871 N_p and N_f cannot be directly measured but one can measure the number of events where
 872 no leptons, N_{t0} , or one lepton, N_{t1} , pass the tight analysis requirement. These numbers are
 873 related by the following equations:

$$\begin{aligned}
N_\ell &= N_p + N_f = N_{t0} + N_{t1} \\
N_{t0} &= (1 - p)N_p + (1 - f)N_f \\
N_{t1} &= pN_p + fN_f
\end{aligned} \tag{4.23}$$

where p and f are the prompt and fake rates respectively. Equation (4.23) can be inverted to obtain the number of prompt and fake leptons:

$$\begin{aligned}
N_p &= \frac{1}{p - f} [(1 - f)N_{t1} - fN_{t0}] \\
N_f &= \frac{1}{p - f} [pN_{t0} - (1 - p)N_{t1}]
\end{aligned} \tag{4.24}$$

The number of fake events passing the tight analysis requirement is $N_{\text{fake}} = fN_f$. The fake background contribution is estimated directly from data, applying the kinematics-dependent weights (f and p are estimated in bins of p_T and η) defined in Eq.(4.24).

The prompt and fake rate estimations after the removal of the EW contribution are shown in Tables 4.4 and 4.5 separately for electrons and muons.

The region obtained by reversing the opposite sign lepton requirement in the analysis selection is enriched with W+jets events where one of the jets is misidentified as a lepton. The fake rate procedure can be applied to this same-sign control region to perform a closure test of the method. The results of the closure test on same-sign events gives good agreement with the expectations.

The systematic uncertainty on the prompt and fake rate estimation is evaluated by varying the jet thresholds in the dijet control sample, and an uncertainty on the background normalization is added according to the agreement with data in the same-sign control region. The systematic uncertainty amounts to about 36% of the fake background yield.

890 Drell-Yan to $\tau\tau$ background

The low E_T^{miss} threshold in the $e\mu$ final state requires the consideration of the contribution from $Z/\gamma^* \rightarrow \tau^+\tau^-$, that is estimated from data. This is accomplished by selecting $Z/\gamma^* \rightarrow \mu^+\mu^-$ events in data and replacing both muons with a simulated $\tau \rightarrow \ell\nu_\tau\bar{\nu}_\ell$ decay [31], thus obtaining a “hybrid” event. The Z boson four-momentum is reconstructed in data from the four-momenta of the daughter muons. Then a simulation step allows the replacement

Table 4.4.: Measured prompt rate for electrons and muons in bins of η , p_T . Only the statistical uncertainties are shown.

Electron prompt rate			
p_T range [GeV]	$0 < \eta \leq 1.4442$	$1.4442 < \eta \leq 1.566$	$1.566 < \eta$
$10 < p_T \leq 15$	0.5738 ± 0.0045	0.5366 ± 0.0204	0.2947 ± 0.0047
$15 < p_T \leq 20$	0.7091 ± 0.0020	0.5484 ± 0.0185	0.4477 ± 0.0034
$20 < p_T \leq 25$	0.7175 ± 0.0013	0.6297 ± 0.0067	0.6200 ± 0.0001
$25 < p_T \leq 50$	0.9219 ± 0.0002	0.8404 ± 0.0007	0.8509 ± 0.0001
$p_T > 50$	0.9693 ± 0.0002	0.9398 ± 0.0021	0.9385 ± 0.0005

Muon prompt rate		
p_T range [GeV]	$0 < \eta \leq 1.5$	$1.5 < \eta \leq 2.5$
$10 < p_T \leq 15$	0.7119 ± 0.0003	0.7582 ± 0.0006
$15 < p_T \leq 20$	0.8049 ± 0.0018	0.8495 ± 0.0001
$20 < p_T \leq 25$	0.9027 ± 0.0008	0.8948 ± 0.0012
$25 < p_T \leq 50$	0.9741 ± 0.0001	0.9627 ± 0.0002
$p_T > 50$	0.9900 ± 0.0001	0.9875 ± 0.0003

896 of the muon objects with τ leptons, in such a way to preserve the Z boson momentum
 897 direction is preserved in its rest frame. The $Z/\gamma^* \rightarrow \tau^+\tau^-$ decay is simulated with the
 898 TAUOLA package [58] to correctly describe the τ -polarization effects.

899 After replacing muons from $Z/\gamma^* \rightarrow \mu^+\mu^-$ decays with simulated τ decays, the set
 900 of pseudo- $Z/\gamma^* \rightarrow \tau^+\tau^-$ events undergoes the reconstruction step. Good agreement in
 901 kinematic distributions for this sample and a MC based $Z/\gamma^* \rightarrow \tau^+\tau^-$ sample is found. The
 902 global normalization of pseudo- $Z/\gamma^* \rightarrow \tau^+\tau^-$ events is checked in the low m_T spectrum
 903 where a rather pure $Z/\gamma^* \rightarrow \tau^+\tau^-$ sample is expected.

904 This method allows to avoid the simulation of very large MC samples that would be
 905 needed for an accurate description of this process.

906 ZZ, WZ and $W\gamma$ backgrounds

907 The WZ and ZZ backgrounds are partially estimated from data when the two selected
 908 leptons come from the same Z boson. If the leptons come from different bosons the
 909 contribution is expected to be small. The WZ component is largely rejected by requiring
 910 only two high p_T isolated leptons in the event.

Table 4.5.: Measured electrons and muons fake rates in bins of η and p_T , after the EWK correction. Only statistical uncertainties are shown.

electron fake rate				
p_T range [GeV]	$0 < \eta \leq 1$	$1 < \eta \leq 1.479$	$1.479 < \eta \leq 2$	$2 < \eta \leq 2.5$
$10 < p_T \leq 15$	0.045 ± 0.005	0.033 ± 0.004	0.008 ± 0.002	0.021 ± 0.005
$15 < p_T \leq 20$	0.044 ± 0.003	0.049 ± 0.003	0.017 ± 0.001	0.017 ± 0.002
$20 < p_T \leq 25$	0.041 ± 0.002	0.064 ± 0.003	0.025 ± 0.002	0.025 ± 0.002
$25 < p_T \leq 30$	0.059 ± 0.003	0.101 ± 0.005	0.041 ± 0.003	0.043 ± 0.003
$30 < p_T \leq 35$	0.084 ± 0.006	0.111 ± 0.009	0.058 ± 0.006	0.066 ± 0.005

muon fake rate				
p_T range [GeV]	$0 < \eta \leq 1$	$1 < \eta \leq 1.479$	$1.479 < \eta \leq 2$	$2 < \eta \leq 2.5$
$10 < p_T \leq 15$	0.131 ± 0.002	0.154 ± 0.004	0.194 ± 0.005	0.241 ± 0.009
$15 < p_T \leq 20$	0.143 ± 0.007	0.191 ± 0.012	0.235 ± 0.016	0.308 ± 0.027
$20 < p_T \leq 25$	0.198 ± 0.005	0.239 ± 0.009	0.221 ± 0.011	0.271 ± 0.021
$25 < p_T \leq 30$	0.182 ± 0.011	0.228 ± 0.018	0.195 ± 0.022	0.287 ± 0.045
$30 < p_T \leq 35$	0.170 ± 0.021	0.244 ± 0.036	0.195 ± 0.041	0.289 ± 0.111

911 The $W\gamma^{(*)}$ background, where the photon decays to an electron-positron pair, is expected
 912 to be very small, thanks to the stringent photon conversion requirements. This background
 913 also includes events where a real photon is produced in association with the W boson.
 914 These events constitute a background for this analysis because the photon can interact with
 915 the tracker material converting to an electron-positron pair.

916 Since the WZ simulated sample has a generation level cut on the di-lepton invariant mass
 917 ($m_{\ell\ell} > 12$ GeV) and the cross-section raises quickly with the lowering of this threshold, a
 918 dedicated MADGRAPH sample has been produced with lower momentum cuts on two of
 919 the three leptons ($p_T > 5$ GeV) and no cut on the third one. The surviving contribution
 920 estimated with this sample is still very small, and since the uncertainty on the cross-section
 921 for the covered phase space is large, a conservative 100% uncertainty has been given to it.
 922 A k -factor for $W\gamma^*$ of 1.5 ± 0.5 based on a dedicated measurement of tri-lepton decays,
 923 $W\gamma^* \rightarrow e\mu\mu$ and $W\gamma^* \rightarrow \mu\mu\mu$, is applied [31]. The contribution of $W\gamma^{(*)}$ is also constrained
 924 by a closure test with same sign leptons on data, which reveals a good compatibility of the
 925 data with the expected background.

926 4.5. Systematic uncertainties

927 Systematic uncertainties play an important role in this analysis where no strong mass peak
 928 is expected due to the presence of undetected neutrinos in the final state. One of the most
 929 important sources of systematic uncertainty is the normalization of the backgrounds that
 930 are estimated on data control samples whenever is possible.

931 A summary of the main sources of systematic uncertainty and the corresponding estimate
 932 is reported in Table 4.6. A detailed description of each source of systematic uncertainty is
 933 discussed in the following sections.

Table 4.6.: Main sources of systematic uncertainties and their estimate. The first category reports the uncertainties in the normalization of background contributions. The experimental and theoretical uncertainties refer to the effect on signal yields. A range is specified if the uncertainty varies across the p_T^H bins.

Uncertainties in backgrounds contributions	
Source	Uncertainty
$t\bar{t}$, tW	20–50%
$W + \text{jets}$	40%
WZ, ZZ	4%
$W\gamma^{(*)}$	30%

Effect of the experimental uncertainties on the signal and background yields	
Source	Uncertainty
Integrated luminosity	2.6%
Trigger efficiency	1–2%
Lepton reconstruction and identification	3–4%
Lepton energy scale	2–4%
E_T^{miss} modelling	2%
Jet energy scale	10%
Pileup multiplicity	2%
b mistag modelling	3%

Effect of the theoretical uncertainties on signal yield	
Source	Uncertainty
b jet veto scale factor	1–2%
PDF	1%
WW background shape	1%

934 4.5.1. Background normalization uncertainties

935 The signal extraction is performed subtracting the estimated backgrounds to the event
 936 counts in data. This uncertainty depends on the background:

- 937 • **tt** and tW backgrounds: The efficiency on jets b-tagging is estimated using the
 938 Tag&Probe technique in data and simulation control regions, as explained in 4.4.1. A
 939 per-jet scale factor, which takes into account the possibly different efficiency of the
 940 anti b-tagging selection in data and simulation, is computed by means of the efficiency
 941 measured with the Tag&Probe method. The Tag&Probe method has been used also
 942 to measure the mistag rates in data and simulation, which are the probability to b-tag
 943 a jet that is not produced by the hadronization of a b quark. These factors are used
 944 to reweigh the Top MC samples as explained in 4.4.1. The uncertainties provided by
 945 the Tag&Probe fit are then propagated to the factor α that is used in the top data
 946 driven estimation 4.4.1. These uncertainties are embedded in a systematic error that
 947 affects the shape of the Top background in each p_T^H bin.

948 Provided that the simulated samples include both tt> and tW processes, a systematic
 949 uncertainty related to the tW/tt> fraction has been included. In fact, a relative variation
 950 of the contribution of these two processes could modify the shape of the MC sample,
 951 and is thus included as a shape uncertainty affecting the top quark background shape
 952 in each p_T^H bin in a correlated way.

- 953 • **W+jets background:** It is estimated with data control sample as described in
 954 Sec.4.4.3. With 19.4 fb^{-1} at 8 TeV, the uncertainty receives similar contributions from
 955 statistics and systematic error (mainly jet composition differences between the fake
 956 rate estimation sample and the application sample), the total error being about 40%,
 957 dominated by the closure test of the method on a same-sign control region.
- 958 • **WZ,ZZ,W $\gamma^{(*)}$ backgrounds:** those backgrounds, which are expected to give a
 959 small contribution, are estimated from simulation. Uncertainties on the cross sections
 960 reported in [59, 60] are 4% for WZ and 2.5% for ZZ. A 30% uncertainty is assigned to
 961 the $W\gamma$ [61] yield and another 30% on $W\gamma^{(*)}$ contribution according to the uncertainty
 962 on the normalization study (see Sec. 4.4.3).

963 4.5.2. Experimental uncertainties

964 The following experimental systematic sources have been taken into account:

- 965 • **Luminosity:** Using the online luminosity monitoring CMS reached an uncertainty
 966 on the luminosity of 2.6% at 8 TeV.
- 967 • **Trigger efficiency.** The uncertainties for both electrons and muons are at 1-2% level,
 968 which is added together to the lepton efficiency uncertainty.

- **Lepton reconstruction and identification efficiency:** The lepton reconstruction and identification efficiencies are measured with the Tag&Probe method in data. To correct for the difference in the lepton identification efficiencies between data and MC, a scale factor is applied to MC. The uncertainties resulting from this procedure on the lepton efficiencies are 4% for electrons and 3% for muons.
- **Muon momentum and electron energy scale:** The momentum scale of leptons have relatively large uncertainties due to different detector effects. For electrons a scale uncertainty of 2% for the barrel, and 4% for the endcaps respectively, is assigned. For muons, a momentum scale uncertainty of 1.5%, independent of its pseudorapidity, is assigned.
- **E_T^{miss} modeling:** The E_T^{miss} measurement is affected by the possible mis-measurement of individual particles addressed above, as well as the additional contributions from the pile-up interactions. The effect of the missing transverse momentum resolution on the event selection is studied by applying a Gaussian smearing of 10% on the x - and y -components of the missing transverse momentum. All correlated variables, like the transverse mass, are recalculated.
- **Jet energy scale (JES) uncertainties:** It affects both the jet multiplicity and the jet kinematic variables, such as m_{jj} . We estimate this uncertainty applying variations of the official jet uncertainties on the JES (which depend on η and p_T of the jet [62]) and compute the variation of the selection efficiency.
- **b jets mistag modeling:** A fraction of signal events is rejected because erroneously identified as b jets by the b-tagging algorithms. The mistag rate, as measured with the Tag&Probe technique described in Sec. 4.4.1, comes with an uncertainty due to different modeling of the b-tagging performance in data and simulation.
- **Pileup multiplicity:** Some of the variables used in the analysis are affected by the average number of pileup interactions. The simulated events have been reweighted according the instantaneous luminosity measured on data. The error in the average number of pileup interactions measured in data and the simulation of the modeling and physics aspects of the pileup simulation gives an uncertainty of 5% on the distribution used in the reweighting procedure. This uncertainty is propagated through all the analysis, and the estimated uncertainty on the efficiency is 2%.

4.5.3. Theoretical uncertainties

- **QCD scale uncertainties:** The uncertainties on the total cross sections due to the choice of the renormalization and factorization scale are assigned to MC-driven backgrounds. For the signal processes these uncertainties are separated in two categories: those affecting the selection efficiency and those affecting the jet bin fractions. The effect of renormalization and factorization scale on the selection efficiency is of the

order of 2% for all processes. Although this analysis is inclusive in number of jets, the effect of the QCD scale variation on the jet bin migrations has to be taken into account because of the b-tagging veto efficiency. The efficiency of this selection depends on jet multiplicity and the effect of the QCD scale variation has been evaluated using the Stewart-Tackman method, as explained in 4.5.3.

- **PDFs uncertainties:** The utilization of different PDF sets can affect both the normalization and the shapes of the signal contributions. The uncertainty related due to the variations in the choice of PDFs is considered following the PDF4LHC [63, 64] prescription, using CT10, NNPDF2.1 [65] and MSTW2008 [66] PDF sets.
- **WW:** Due to the fact that the WW shape is entirely taken from simulation, the analysis is strongly relying on theoretical models and can thus be strongly affected by their uncertainties. Especially higher order QCD radiative effects have an influence on the generated WW shape. To study this impact, the shapes of the distributions produced with the MADGRAPH generator (which is the generator for the MC simulation used in the analysis) are compared to the ones produced with MC@NLO. The comparison is performed separately in each bin of p_T^H and the uncertainty includes shape differences originating from the renormalization and factorization scale choice. A comparison of the $m_{\ell\ell}$ and m_T shapes for the WW background using different MC generators is reported in section 4.4.2.

1025 Jet multiplicity uncertainty

The jet bin uncertainty on the ggH production mode has been evaluated using the Stewart-Tackman method, following the recipe proposed in Refs. [30, 67]. Three independent nuisance parameters have to be associated with the inclusive ggH production cross sections $\sigma_{\geq 0}$, $\sigma_{\geq 1}$ and $\sigma_{\geq 2}$, which corresponds to the cross sections with ≥ 0 jets, ≥ 1 jet and ≥ 2 jets respectively. According to the agreement on the treatment of uncertainties in the combination of ATLAS and CMS results [68], these nuisance parameters are labelled as $QCDscale_{ggH}$, $QCDscale_{ggH1in}$ and $QCDscale_{ggH2in}$. However, in case the analysis is split in exclusive jet multiplicity bins, the jet bin uncertainties can be evaluated taking into account the correct correlations among the three nuisances following the Stewart-Tackman prescription. Even though this analysis is inclusive in number of jets, the jet binning uncertainties must be included due to the presence of the b-jet veto, that introduces a dependency of the selection efficiency on the number of jets in the event. The veto efficiency has been evaluated in all the p_T^H bins defined in the analysis and as a function of jets multiplicity. The results are shown in Figs. 4.17(a) and 4.17(b). The drop of the veto efficiency at high values of p_T^H is due to the correlation with jets multiplicity.

The first step of this procedure is to take the inclusive ggH cross section, σ_{ggH} , and to convert the relative QCD up/down scale uncertainties, ϵ_+ and ϵ_- , to a log-normal uncertainty, i.e. $\kappa = \sqrt{\exp(\epsilon_+) \cdot \exp(\epsilon_-)}$. The exclusive cross sections, σ_0 , σ_1 and σ_2 , can

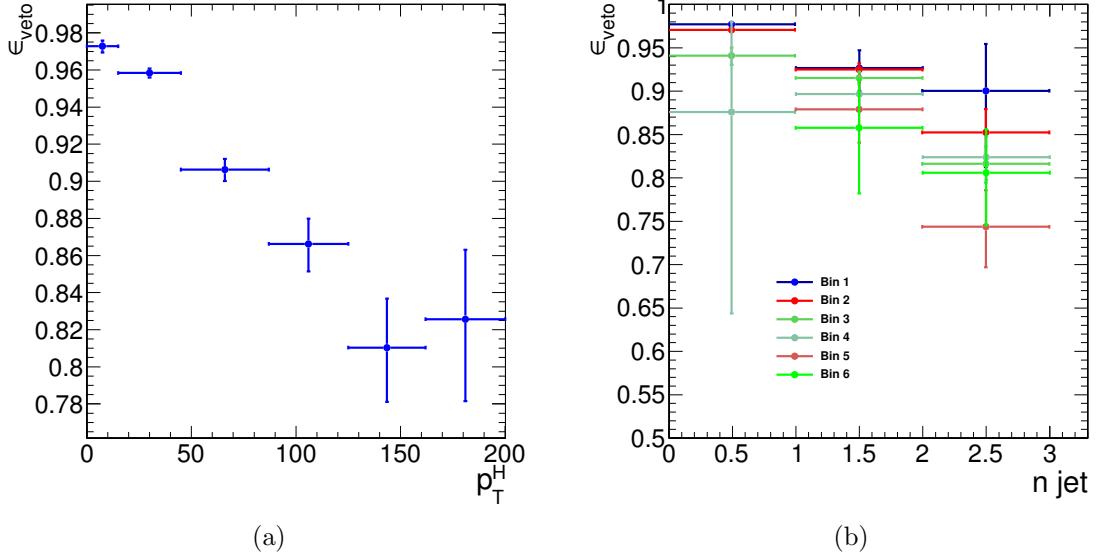


Figure 4.17.: (a) Efficiency of the b-tagging veto in different bins of p_T^H . **(b)** Efficiency of the b-tagging veto in different bins of p_T^H , as a function of number of jets.

be calculated starting from σ_{ggH} and using the selection efficiencies for the three jet bins.
 For every exclusive cross section the corresponding relative uncertainty is computed varying
 the renormalization (μ_R) and factorization (μ_F) scales independently of a factor 2 and
 $1/2$, and taking the cross section value corresponding to half of the maximum variation.
 The inclusive cross sections are then obtained summing the exclusive cross sections and
 propagating the uncertainties, i.e. $\sigma_{\geq 0} = \sigma_0 + \sigma_1 + \sigma_2$, $\sigma_{\geq 1} = \sigma_1 + \sigma_2$, $\sigma_{\geq 2} = \sigma_2$.

The three nuisance parameters, including all the proper correlations among the jet
 bins, are defined according to Table 4.7, where the f_n constants represent the exclusive
 theoretical n jet bin fractions, i.e. $f_0 = \sigma_0/\sigma_{\geq 0}$, $f_1 = \sigma_1/\sigma_{\geq 0}$, $f_2 = \sigma_2/\sigma_{\geq 0}$.

The nuisance parameters reported in table 4.7 have then been calculated for each p_T^H
 bin embedding the b-jet veto efficiency and using the following formulas:

$$QCDscale_{ggH} = \frac{\Delta_{\geq 0}^0 \cdot f_0 \cdot \epsilon_0 + \Delta_{\geq 1}^0 \cdot f_1 \cdot \epsilon_1}{\Delta_{\geq 0}^0 \cdot f_0 \cdot \epsilon_0 + \Delta_{\geq 1}^0 \cdot f_1 \cdot \epsilon_0} \quad , \quad (4.25)$$

$$QCDscale_{ggH1in} = \frac{\Delta_{\geq 1}^1 \cdot f_1 \cdot \epsilon_1 + \Delta_{\geq 2}^1 \cdot f_2 \cdot \epsilon_2}{\Delta_{\geq 1}^1 \cdot f_1 \cdot \epsilon_1 + \Delta_{\geq 2}^1 \cdot f_2 \cdot \epsilon_1} \quad , \quad (4.26)$$

$$QCDscale_{ggH2in} = 1 \quad , \quad (4.27)$$

Table 4.7.: Numerical calculation for the systematic uncertainties of jet binning.

Nuisance parameter	0-jet bin	1-jet bin	2-jet bin
QCDscale_ggH	$\Delta_{\geq 0}^0 = (\kappa_{\geq 0})^{\frac{1}{f_0}}$		
QCDscale_ggH1in	$\Delta_{\geq 1}^0 = (\kappa_{\geq 1})^{-\frac{f_1+f_2}{f_0}}$	$\Delta_{\geq 1}^1 = (\kappa_{\geq 1})^{\frac{f_1+f_2}{f_1}}$	
QCDscale_ggH2in		$\Delta_{\geq 2}^1 = (\kappa_{\geq 2})^{-\frac{f_2}{f_1}}$	$\Delta_{\geq 2}^2 = (\kappa_{\geq 2})$

where ε_0 , ε_1 and ε_2 are the selection efficiencies for the three jet categories. These nuisance parameters are expected to be equal to one in case the efficiency is independent on the number of jets, i.e if $\varepsilon_0 = \varepsilon_1 = \varepsilon_2$.
The numerical values obtained following this procedure are reported in Table 4.8 for each p_T^H bin.

Table 4.8.: Values of the jet binning nuisance parameters for different p_T^H bins.

Nuisance parameter	p_T^H bin [GeV]					
	[0-15]	[15-45]	[45-85]	[85-125]	[125-165]	[165- ∞]
QCDscale_ggH	0.998	0.993	0.989	1.000	1.000	1.000
QCDscale_ggH1in	0.997	0.993	0.984	0.975	0.946	0.974

4.5.4. Statistics uncertainty of the simulated samples

Due to the large range of weights used to correct the simulated distributions in order to match those in data, the effective size of the MC samples are sometimes smaller than the actual number of events in the sample. The statistical uncertainties of the event yields estimated from MC samples are included as nuisance parameters in the fit and have a small impact on the final result.

1068 **4.5.5. Treatment of systematic uncertainties in the shape
1069 analysis**

1070 One can distinguish between normalization uncertainties, where a systematic effect is
1071 changing the normalization of a given process assuming the shape is not affected, and shape
1072 uncertainties where the actual change in the shape of the distribution is taken into account.
1073 The normalization uncertainties enter the shape analysis as a constant normalization factor,
1074 whereas for shape uncertainties the nominal and the $+1\sigma$ and -1σ shapes enter the analysis
1075 in form of three histograms with the same normalization.

1076 For the W+jets background, the shape differences for different jet p_T thresholds in the
1077 di-jet control sample are considered separately for electron and muon fakes, while the other
1078 sources of systematics are taken as normalization uncertainties as in the cut-based analysis.

1079 Effects from experimental uncertainties are studied by applying a scaling and smearing
1080 of certain variables of the physics objects, followed by a subsequent recalculation of all the
1081 correlated variables. This is done for simulation, to account for possible systematic mis-
1082 measurements of the data. All experimental sources from Section 4.5.2 but luminosity are
1083 treated both as normalization and shape uncertainties. For background with a data-driven
1084 normalization estimation, only the shape uncertainty is considered.

1085 To account for statistical uncertainties, for each distribution going into the shape analysis,
1086 the $+1\sigma$ and -1σ shapes were obtained by adding/subtracting the statistical error in each
1087 bin and renormalizing it to the nominal distribution. In addition to this procedure a
1088 constant normalization uncertainty due to the finite statistics of the MC sample used to
1089 extract the shape is assigned.

1090 **4.6. Signal extraction**

1091 According to the “blinding” policy of the CMS Collaboration, the strategy of the analysis
1092 has been scrutinized and approved by a selected committee of internal reviewers before
1093 looking at the data in the signal region. This approach prevents the analysts from being
1094 biased by the data in the developing phase of the analysis. Below are shown the results
1095 after having looked at the data.

1096 **4.6.1. Fitting procedure**

1097 The signal, including ggH, VBF, and VH production mechanisms, is extracted in each bin
1098 of p_T^H by performing a binned maximum likelihood fit simultaneously in all p_T^H bins to a
1099 two-dimensional template for signals and backgrounds in the $m_{\ell\ell}$ – m_T plane. The variables
1100 used for the two-dimensional template are chosen for their power to discriminate signal

and background contributions. This is shown in Fig. 4.18, where the two-dimensional MC distributions are shown for the signal and background processes in the 0-jets category.

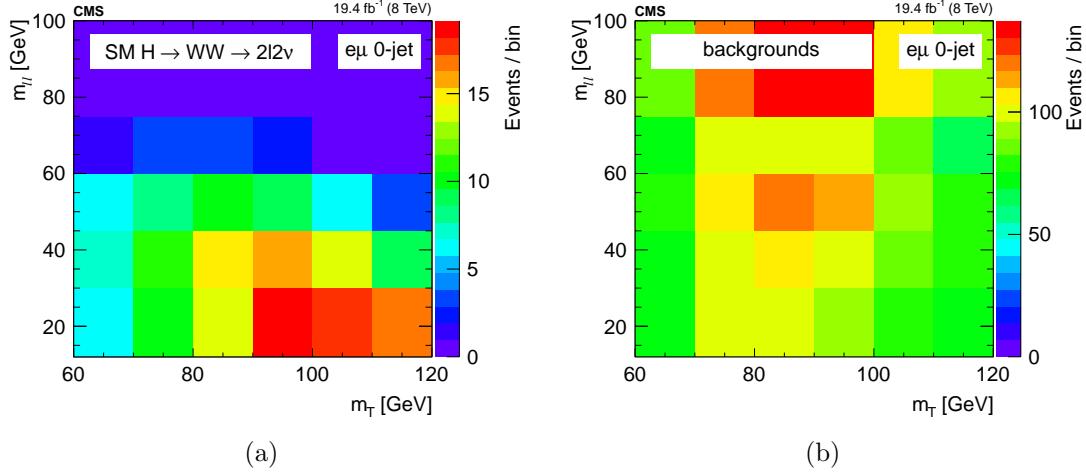


Figure 4.18.: Two-dimensional $m_{\ell\ell}$ – m_T distribution for signal (a) and background (b) processes in the 0-jets category.

Six different signal strength parameters are extracted from the fit, one for each p_T^H bin. The relative contributions of the different Higgs production mechanisms in the signal template are taken to be the same as in the SM. The systematic uncertainty sources are considered as nuisance parameters in the fit.

The binning of the $m_{\ell\ell}$ and m_T templates is chosen to be:

- $m_{\ell\ell}$: [12, 30, 45, 60, 75, 100, 125, 150, 175, 200]
- m_T : [60, 70, 80, 90, 100, 110, 120, 140, 160, 180, 200, 220, 240, 280]

To avoid a dependence of the results on the variables used for the template fit, $m_{\ell\ell}$ and m_T need to be uncorrelated with respect to p_T^H . This has been verified and the correlation between the discriminating variables and p_T^H is shown in Fig. 4.19 and Fig. 4.20 for ggH and VBF production modes respectively.

The relative contribution for different production mechanisms in the input signal template is taken to be the same as the SM. The signal strength μ in each bin, i.e. the ratio between the measured cross section and the SM one, $\mu = \sigma/\sigma_{SM}$, is allowed to float between -10 and +10, thus allowing negative values. This is mainly intended to allow the error bars to float below 0.

Because of detector resolution effects, some of the reconstructed $H \rightarrow WW$ signal events might originate from outside the fiducial phase space. These out-of-fiducial signal events cannot be precisely handled by the unfolding procedure and must be subtracted from the measured spectrum. The p_T^H distribution of the out-of-fiducial signal events is taken from

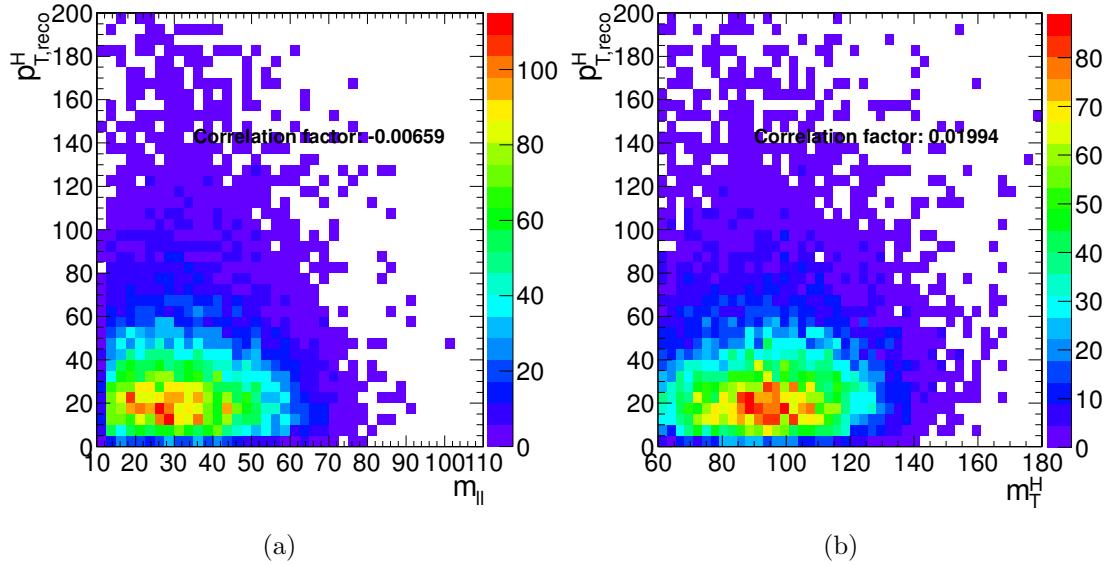


Figure 4.19.: Correlation between p_T^H and $m_{\ell\ell}$ (a) and between p_T^H and m_T (b) after the full selection for the ggH production mode.

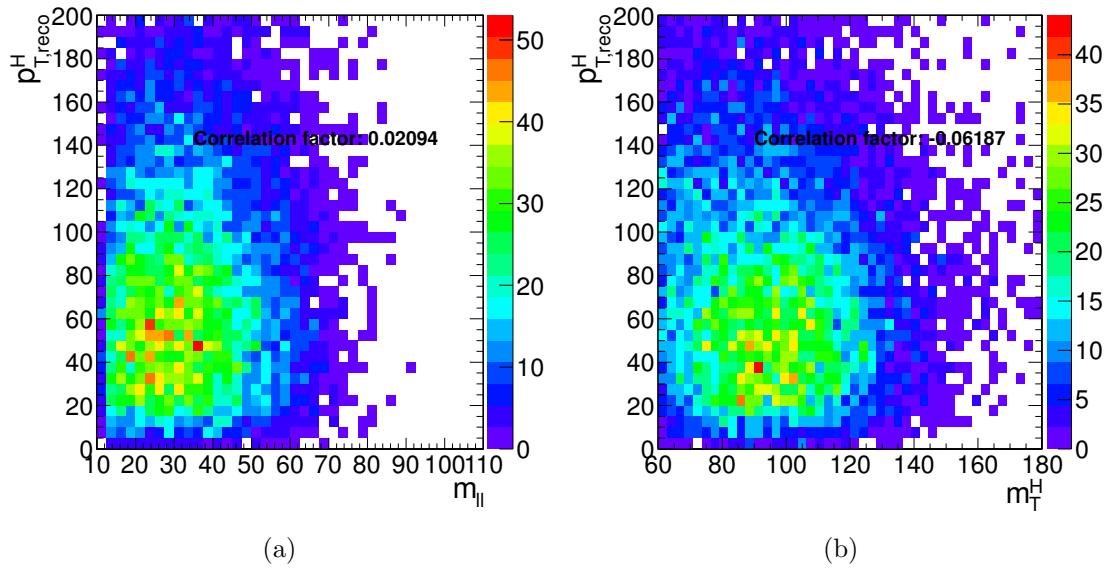


Figure 4.20.: Correlation between p_T^H and $m_{\ell\ell}$ (a) and between p_T^H and m_T (b) after the full selection for the VBF production mode.

¹¹²³ simulation, and each bin is multiplied by the corresponding measured signal strength before performing the subtraction.

¹¹²⁴

1125 At the end, the number of events in each bin i of the measured spectrum is:

$$N_i = \mu_i(s_i - f_i) , \quad (4.28)$$

1126 where s_i and f_i are respectively the number of signal and fake events expected from
 1127 simulation and μ_i is the measured signal strength.

1128 The fit makes use of the binned maximum likelihood approach. The likelihood function,
 1129 \mathcal{L} , restricted to the p_T^H bin j , can be written as: **CHECK!!**

$$\mathcal{L}(data|\mu_j, \theta) = \prod_{i=0}^{N_{\text{bins}}} \frac{(\mu_j s_i(\theta) + b_i(\theta))^{n_i}}{n_i!} e^{-\mu_j s_i(\theta) - b_i(\theta)} \cdot p(\tilde{\theta}|\theta) , \quad (4.29)$$

1130 where *data* corresponds to the experimental observation and μ_j is the signal strength in
 1131 the bin j , i.e. the parameter of interest of the fit, which multiplies the signal yield. The
 1132 index i runs over the bins of the $m_{ll}-m_T$ two-dimensional histogram corresponding the p_T^H
 1133 bin j , s_i and b_i are the expected number of signal and background events respectively in bin
 1134 i , and n_i is the total number of observed events in bin i . The set of parameters θ represents
 1135 the full suite of nuisance parameters used to incorporate the systematic uncertainties. Each
 1136 nuisance parameter is constrained in the fit including the prior distributions functions $p(\tilde{\theta}|\theta)$
 1137 in the likelihood, where $\tilde{\theta}$ is the set of default values for the θ parameters. For the major part
 1138 of the nuisance parameters a log-normal prior distribution is used, with a standard deviation
 1139 corresponding to the given systematic uncertainty. For some nuisance parameters, as the
 1140 ones related to the statistical uncertainty coming from the background measurement in data
 1141 control regions, a Gamma distribution is instead recommended. A log-uniform distribution
 1142 is used for the uncertainty related to the normalization of background contributions that
 1143 are left unconstrained in the fit, such as the WW background process. Finally, some of
 1144 the experimental uncertainties, related to the shape of signal and background processes,
 1145 are modelled by means of additional histograms as explained in Sec. 4.5.5. The nuisance
 1146 parameters correlations across different p_T^H bins are taken into account. Moreover the
 1147 nuisance parameters can also be correlated (or anti-correlated) between signal and different
 1148 background processes. As an example, the uncertainty related to the integrated luminosity
 1149 measurement is fully correlated for all the signal and background processes.

1150 Before running the fit on the data, the same procedure has been applied to the so called
 1151 *Asimov data set*¹, which provides a simple method to estimate the signal sensitivity before
 1152 looking at the data [69].

¹In a parallel reality imagined by the science fiction writer I. Asimov, politics was run in a peculiar way: instead of mobilizing millions of people to cast their vote to deliberate on something, an algorithm was used to select an individual “average” person, and then this person was asked to take the decision on that matter.

₁₁₅₃ **4.6.2. Signal and background yields**

₁₁₅₄ A comparison of data and background prediction is shown in Fig. 4.21, where the $m_{\ell\ell}$
₁₁₅₅ distribution is shown for the six p_T^H bins. Distributions correspond to the m_T window of
₁₁₅₆ [60, 110] GeV, in order to emphasize the signal contribution [31]. The m_T distributions are
₁₁₅₇ shown in Fig. 4.22 and correspond to the $m_{\ell\ell}$ window of [12, 75] GeV.

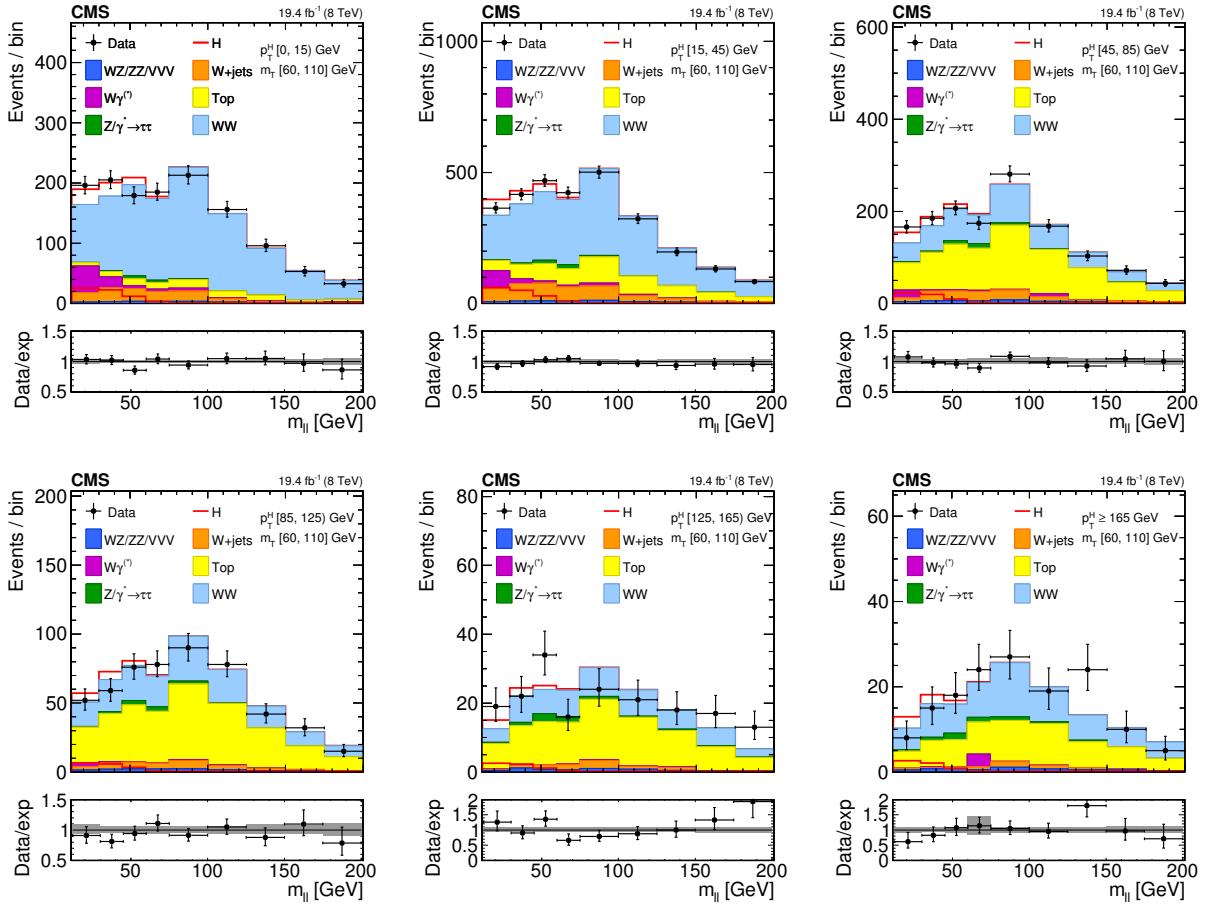


Figure 4.21.: Distributions of the $m_{\ell\ell}$ variable in each of the six p_T^H bins. Background normalizations correspond to the values obtained from the fit. Signal normalization is fixed to the SM expectation. The distributions are shown in an m_T window of [60,110] GeV in order to emphasize the Higgs boson (H) signal. The signal contribution is shown both stacked on top of the background and superimposed to it. Ratios of the expected and observed event yields in individual bins are shown in the panels below the plots. The uncertainty band shown in the ratio plot corresponds to the envelope of systematic uncertainties after performing the fit to the data.

₁₁₅₈ The signal and background yields after the analysis selection are reported in Table 4.9.

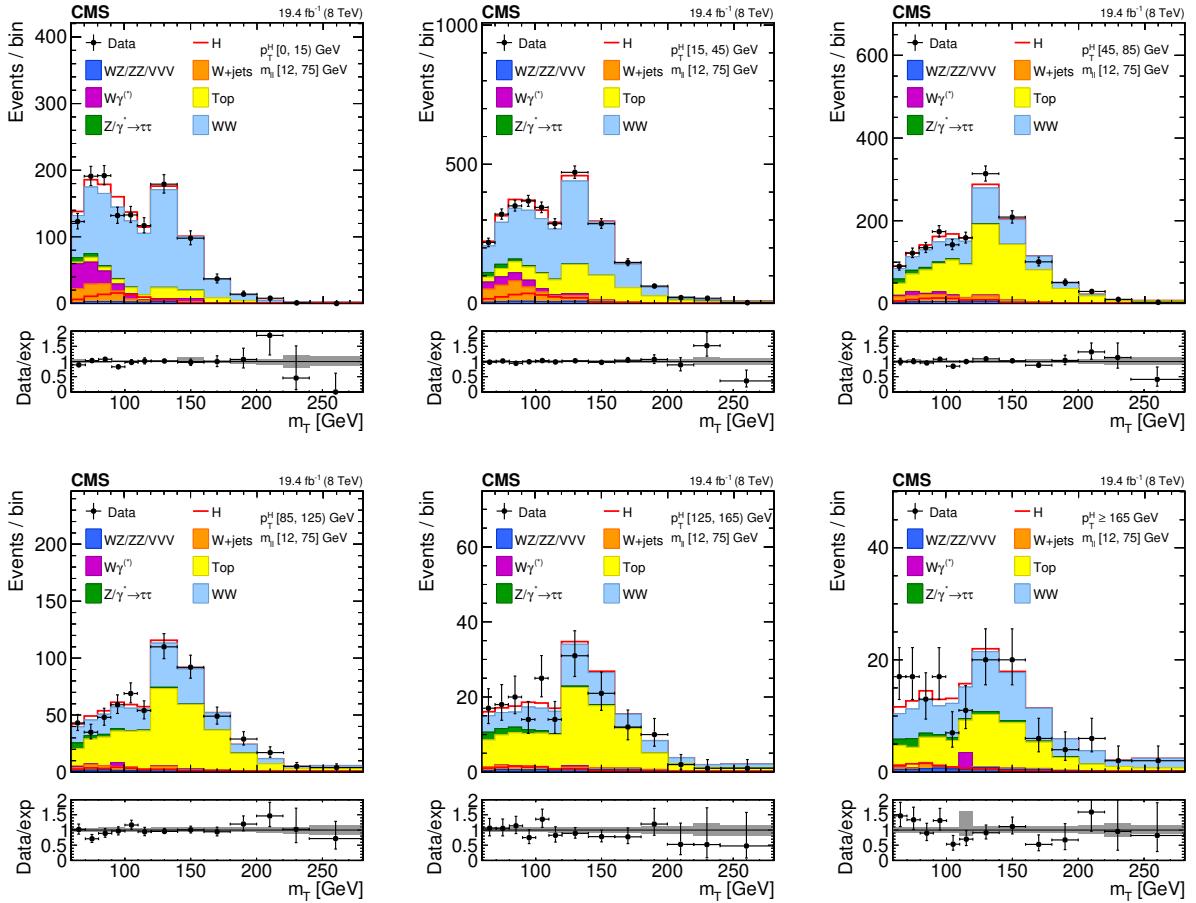


Figure 4.22.: Distributions of the m_T variable in each of the six p_T^H bins. Background normalizations correspond to the values obtained from the fit. Signal normalization is fixed to the SM expectation. The distributions are shown in an $m_{\ell\ell}$ window of $[12, 75]$ GeV in order to emphasize the Higgs boson (H) signal. The signal contribution is shown both stacked on top of the background and superimposed to it. Ratios of the expected and observed event yields in individual bins are shown in the panels below the plots. The uncertainty band shown in the ratio plot corresponds to the envelope of systematic uncertainties after performing the fit to the data.

Table 4.9.: Signal prediction, background estimates and observed number of events in data are shown in each p_T^H bin for the signal after applying the analysis selection requirements. The total uncertainty on the number of events is reported. For signal processes, the yield related to the ggH are shown, separated with respect to the contribution of the other production mechanisms (XH=VBF+VH). The WW process includes both quark and gluon induced contribution, while the Top process takes into account both $t\bar{t}$ and tW .

p_T^H [GeV]	0-15	15-45	45-85	85-125	125-165	165- ∞
ggH	73 ± 3	175 ± 5	59 ± 3	15 ± 2	5.1 ± 1.5	4.9 ± 1.4
XH=VBF+VH	4 ± 2	15 ± 4	16 ± 4	8 ± 2	3.8 ± 1.1	3.0 ± 0.8
Out-of-fiducial	9.2 ± 0.5	19.9 ± 0.7	11.4 ± 0.6	4.4 ± 0.3	1.6 ± 0.2	2.4 ± 0.2
Data	2182	5305	3042	1263	431	343
Total background	2124 ± 128	5170 ± 321	2947 ± 293	1266 ± 175	420 ± 80	336 ± 74
WW	1616 ± 107	3172 ± 249	865 ± 217	421 ± 120	125 ± 60	161 ± 54
Top	184 ± 38	1199 ± 165	1741 ± 192	735 ± 125	243 ± 51	139 ± 49
W+jets	134 ± 5	455 ± 10	174 ± 6	48 ± 4	14 ± 3	9 ± 3
WZ+ZZ+VVV	34 ± 4	107 ± 10	71 ± 7	29 ± 5	14 ± 3	13 ± 4
$Z/\gamma^* \rightarrow \tau^+\tau^-$	23 ± 3	67 ± 5	47 ± 4	22 ± 3	12 ± 2	10 ± 2
$W\gamma^{(*)}$	132 ± 49	170 ± 58	48 ± 30	12 ± 9	3 ± 3	5 ± 10

1159 The spectrum shown in Fig. 4.23 is obtained after having performed the fit and after
 1160 the subtraction of the out-of-fiducial signal events, but before undergoing the unfolding pro-
 1161 cedure. The theoretical distribution after the detector simulation and event reconstruction
 1162 is also shown for comparison.

1163 In order to assess the robustness of the fit, several toy MC samples have been produced,
 1164 with a statistical accuracy corresponding to the one expected in data. The distribution
 1165 of the signal strengths extracted in each bin using the toy MC samples and the their pull
 1166 distributions are shown in Fig. 4.24.

1167 4.7. Unfolding

1168 To facilitate comparisons with theoretical predictions or other experimental results, the
 1169 signal extracted performing the fit has to be corrected for detector resolution and efficiency
 1170 effects and for the efficiency of the selection defined in the analysis. An unfolding procedure
 1171 is used relying on the RooUnfold package [70], which provides the tools to run various
 1172 unfolding algorithms.

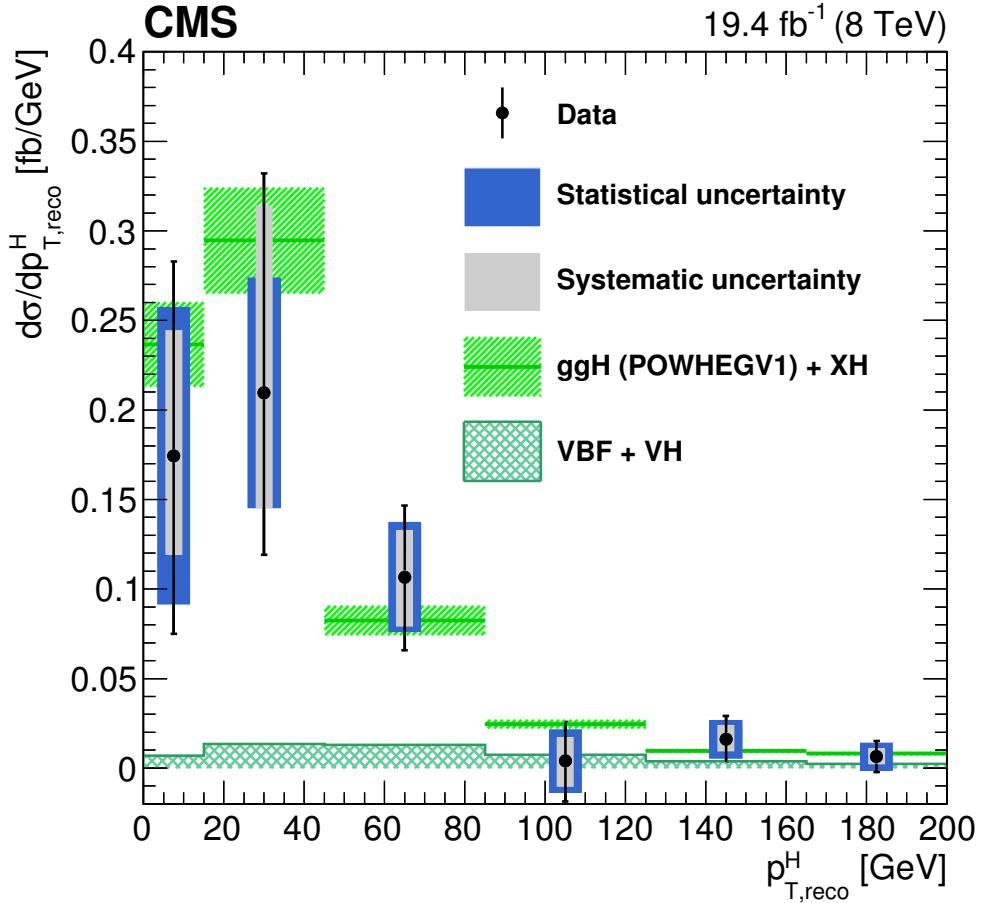


Figure 4.23.: Differential Higgs boson production cross section as a function of the reconstructed p_{T}^H , before applying the unfolding procedure. Data values after the background subtraction are shown together with the statistical and the systematic uncertainties, determined propagating the sources of uncertainty through the fit procedure. The line and dashed area represent the SM theoretical estimates in which the acceptance of the dominant ggH contribution is modelled by POWHEG V1. The sub-dominant component of the signal is denoted as XH=VBF+VH, and is shown with the cross filled area separately.

1173 The basic principle behind the unfolding procedure in this analysis is to use MC signal
 1174 samples to make the “true” distribution of the variable of interest, which is obtained using
 1175 simulated events before particle interaction with the detector, and the same distribution
 1176 obtained using events reconstructed after the full GEANT4 simulation of the CMS detector
 1177 and event reconstruction. These two distributions are used to calculate the detector response
 1178 matrix M :

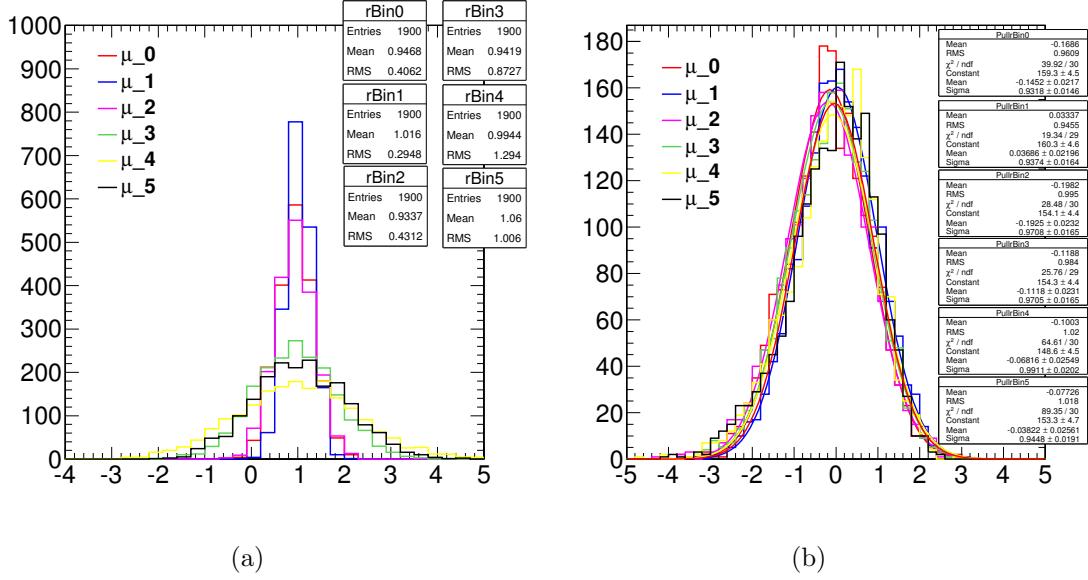


Figure 4.24.: Signal strength distribution as extracted from the fit of toy MC samples (a). Distribution of the pull of the signal strength parameters (b).

$$R_i^{\text{MC}} = \sum_{j=1}^n M_{ij} T_j^{\text{MC}} \quad , \quad (4.30)$$

where R^{MC} and T^{MC} are two n -dimensional vectors representing the distribution before and after event processing through CMS simulation and reconstruction. The dimension n of the two vectors corresponds to the number of bins in the distributions, equal to six in this analysis. The response matrix M includes all the effects related to the detector and analysis selection that affect the R^{MC} distribution. The goal of the unfolding procedure is to obtain the T^{truth} distribution starting from the measured R^{observed} distribution by inverting the matrix M . To avoid the large variance and strong negative correlation between the neighbouring bins [32], the unfolding procedure in this analysis relies on the singular value decomposition [71] method based on the Tikhonov regularization function. Since the response matrix is in general limited by the statistical uncertainties of simulated samples and given the finite data statistical accuracy, a simple inversion could lead to large fluctuations between bins in the unfolded result. In particular, if the off-diagonal elements of the response matrix are sizeable, the unfolded distribution has large variance and strong negative correlations between the neighbouring bins [32]. Several unfolding methods with regularization are available in literature, such as a method based on the Bayes' theorem, which overcome the unfolding instability using an iterative procedure [72]. One possible solution is the utilization of regularization methods. Such methods introduce

1196 a regularization function that controls the smoothness of the distribution and depends
 1197 generally on one regularization parameter, which can be controlled to achieve the desired
 1198 degree of smoothness. The choice of the regularization parameter is particularly critical,
 1199 and it should represent an optimal trade-off between taming the fluctuations in the unfolded
 1200 result, and biasing the unfolded distribution towards the one used to build the response
 1201 matrix. The main feature of this method is the use of the singular value decomposition of
 1202 the response matrix, including an additional term to suppress the oscillatory component
 1203 of the solution, i.e. the regularization term, which represents some *a priori* knowledge
 1204 of the final solution. The regularization parameter is chosen to obtain results that are
 1205 robust against numerical instabilities and statistical fluctuations, following the prescription
 1206 described in Ref. [71]. **Maybe I should add an appendix describing the SVD method in**
 1207 **details**

1208 The response matrix is built as a two-dimensional histogram, with the generator-level
 1209 p_T^H on the y axis and the same variable after the reconstruction on the x axis, using the
 1210 same binning for both distributions. The resulting detector response matrix, including all
 1211 signal sources and normalized by row, is shown in Fig. 4.25(a). The value of the diagonal
 1212 bins corresponds to the stability S . The same matrix, normalized by column, is shown
 1213 in Fig. 4.25(b). In this case the diagonal bins correspond to the purity P . The S and P
 1214 parameters, defined in Sec. 4.3, provide an estimate of the p_T^H resolution and migration
 1215 effects. The main source of bin migrations effects in the response matrix is the limited
 1216 resolution in the measurement of E_T^{miss} .

1217 The resulting detector response matrix, which includes the effects of all signal sources and
 1218 is represented by normalizing each row to unity is shown in Fig. 4.25(a). This representation
 1219 shows the stability S in the diagonal bins, where S is defined as the ratio of the number of
 1220 events generated and reconstructed in a given bin, and the number of events generated in
 1221 that bin. In addition, a deconvolution matrix is constructed by normalizing each column
 1222 to unity and is shown in Fig. 4.25(b). This latter representation shows the purity P in
 1223 the diagonal bins, where P is defined as the ratio of the number of events generated and
 1224 reconstructed in a given bin, and the number of events reconstructed in that bin. The S
 1225 and P parameters provide an estimate of the p_T^H resolution and of migration effects. The
 1226 response matrix built including all signal sources is shown in Fig. 4.25. In order to point
 1227 out either the purity or the stability in diagonal bins, each column or row of the matrix was
 1228 respectively normalized to unity. The matrix obtained in the first case is what is actually
 1229 called detector response matrix, while in the other case the matrix is usually referred to as
 1230 detector deconvolution matrix.

1231 Several closure tests are performed in order to validate the unfolding procedure. To
 1232 estimate the uncertainty in the unfolding procedure due to the particular model adopted for
 1233 building the response matrix, two independent gluon fusion samples are used, corresponding
 1234 to two different generators: POWHEG V1 and JHUGEN generators, both interfaced to
 1235 PYTHIA 6.4. The JHUGEN generator sample is used to build the response matrix while the
 1236 POWHEG V1 sample is used for the measured and the MC distributions at generator level.

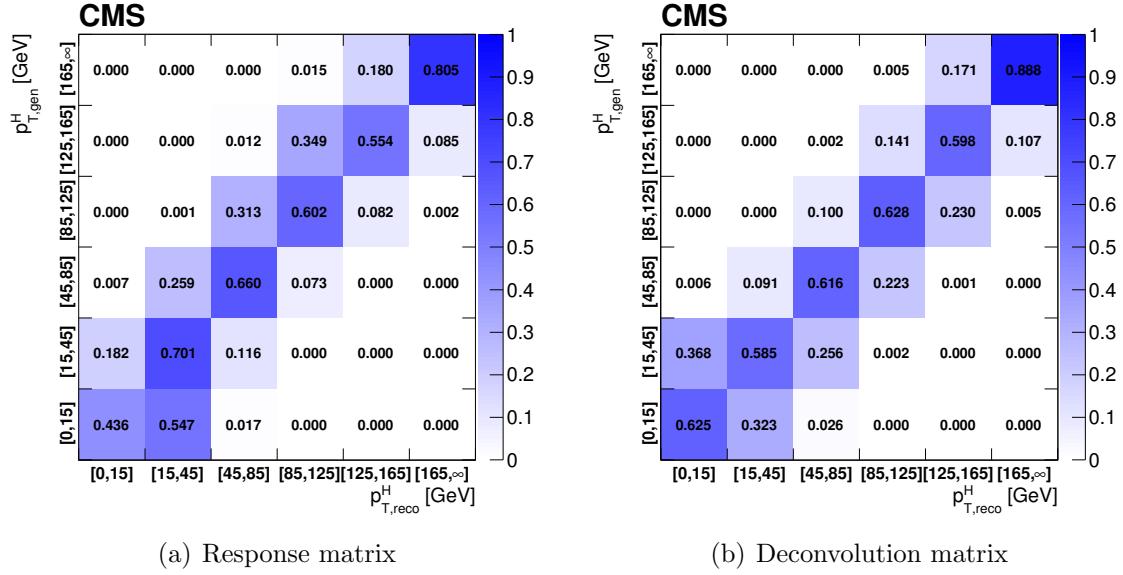


Figure 4.25.: Response matrix (a) and deconvolution matrix (b) including all signal processes. The matrices are normalized either by row (a) or by column (b) in order to show the purity or stability respectively in diagonal bins.

1237 The result of this test shows good agreement between the unfolded and the distribution
1238 from MC simulation.

1239 In order to further prove the choice of the regularization parameter, a large number of
1240 simulated pseudo-experiments has been generated to verify that the coverage of the unfolded
1241 uncertainties obtained with this procedure is as expected. From each pseudo-experiment
1242 the reconstructed p_T^H spectrum is obtained and then unfolded using the procedure described
1243 above, including only the statistical uncertainties. The coverage is calculated for each p_T^H
1244 bin, counting the number of pseudo-experiments for which the statistical uncertainty covers
1245 the true value. The confidence intervals are calculated using the Clopper-Pearson approach,
1246 and the results are shown in Table 4.10 for different values of the regularization parameter:
1247 starting from $k_{\text{reg}} = 2$ (stronger regularization) up to $k_{\text{reg}} = 5$ (weaker regularization). The
1248 criterion for choosing the best k_{reg} value is to increase the regularization as much as possible
1249 without introducing a bias, i.e. until a 68% coverage is fulfilled. This criterion leads to the
1250 same result as the prescription described in Ref. [71], strengthening the choice of $k_{\text{reg}} = 3$.

1251 4.7.1. Treatment of systematic uncertainties

1252 An important aspect of this analysis is the treatment of the systematic uncertainties and
1253 the error propagation through the unfolding procedure. The sources of uncertainty are

Table 4.10.: Coverage interval for each bin and for different values of the regularization parameter, obtained using pseudo-experiments.

p_T^H bin [GeV]	Coverage			
	$k_{\text{reg}} = 2$	$k_{\text{reg}} = 3$	$k_{\text{reg}} = 4$	$k_{\text{reg}} = 5$
0–15	$0.654^{+0.015}_{-0.016}$	$0.704^{+0.015}_{-0.015}$	$0.727^{+0.014}_{-0.015}$	$0.755^{+0.014}_{-0.014}$
15–45	$0.701^{+0.015}_{-0.015}$	$0.665^{+0.015}_{-0.016}$	$0.683^{+0.015}_{-0.015}$	$0.733^{+0.014}_{-0.015}$
45–85	$0.717^{+0.014}_{-0.015}$	$0.706^{+0.015}_{-0.015}$	$0.709^{+0.015}_{-0.015}$	$0.716^{+0.014}_{-0.015}$
85–125	$0.634^{+0.016}_{-0.016}$	$0.681^{+0.015}_{-0.015}$	$0.714^{+0.015}_{-0.015}$	$0.739^{+0.014}_{-0.015}$
125–165	$0.599^{+0.015}_{-0.016}$	$0.650^{+0.015}_{-0.016}$	$0.700^{+0.015}_{-0.015}$	$0.751^{+0.014}_{-0.014}$
165– ∞	$0.632^{+0.016}_{-0.016}$	$0.674^{+0.015}_{-0.015}$	$0.701^{+0.015}_{-0.015}$	$0.722^{+0.014}_{-0.015}$

divided into three categories, depending on whether the uncertainty affects only the signal yield (type A), both the signal yield and the response matrix (type B), or only the response matrix (type C). These three classes propagate differently through the unfolding procedure.

Type A uncertainties are extracted directly from the fit in the form of a covariance matrix, which is passed to the unfolding tool as the covariance matrix of the measured distribution. The nuisance parameters belonging to this category are the background shape and normalization uncertainties. To extract the effect of type A uncertainties a dedicated fit is performed, fixing to constant all the nuisance parameters in the model, but type A nuisance parameters. The correlation matrix among the six signal strengths corresponding to the six p_T^H bins, including all type A uncertainties, is shown in Fig. 4.26. The correlation $\text{cor}(i,j)$ of bins i and j is defined as:

$$\text{cor}(i,j) = \frac{\text{cov}(i,j)}{s_i s_j} \quad , \quad (4.31)$$

where $\text{cov}(i,j)$ is the covariance of bins i and j , and (s_i, s_j) are the standard deviations of bins i and j , respectively.

The nuisance parameters falling in the type B class are:

- the b veto scale factor. It affects the signal and background templates by varying the number of events with jets that enter the selection. It also affects the response matrix because the reconstructed spectrum is harder or softer depending on the number of jets, which in turn depends on the veto.

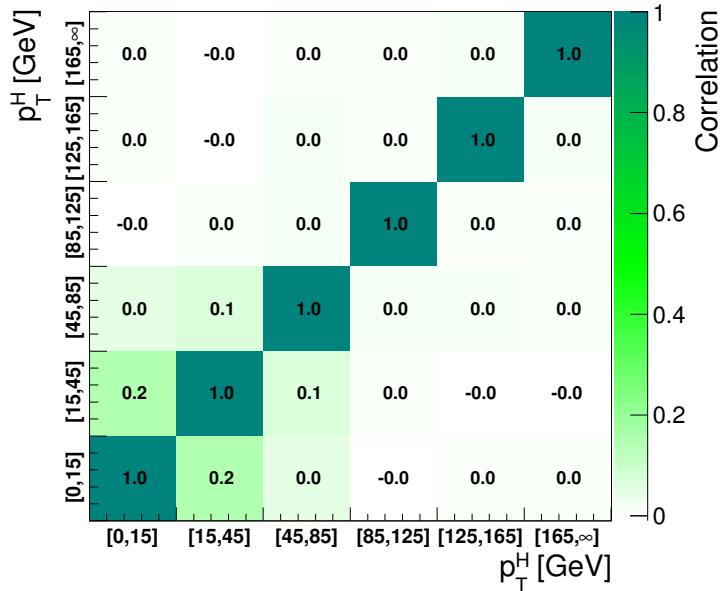


Figure 4.26.: Correlations among the signal strengths corresponding to the six p_T^H bins including all type A uncertainties.

- the lepton efficiency scale factor. It affects the signal and background template shape and normalization. It affects the response matrix by varying the reconstructed spectrum;
- the E_T^{miss} scale and resolution, which have an effect similar to the above;
- lepton scale and resolution. The effect is similar to the above;
- jet energy scale. It affects the signal and background template shape and normalization. It also affects the response matrix because, by varying the fraction of events with jets, the b veto can reject more or fewer events, thus making the reconstructed spectrum harder or softer.

The effect of each type B uncertainty is evaluated separately, since each one changes the response matrix in a different way. In order to evaluate their effect on the signal strengths parameters, two additional fits are performed, each time fixing the nuisance parameter value to ± 1 standard deviation with respect to its nominal value. The results of the fits are then compared to the results of the full fit obtained by floating all the nuisance parameters, thus determining the relative uncertainty on the signal strengths due to each nuisance parameter, as shown in Tab. 4.11. Using these uncertainties, the measured spectra for each type B source are built. The effects are propagated through the unfolding by building the corresponding variations of the response matrix and unfolding the measured spectra with the appropriate matrix.

Table 4.11.: Effect of all the Type B uncertainties on the signal strengths of each bin. In the table are reported the signal strength variations corresponding to an up or down scaling of each nuisance.

Type B uncertainty	Effect on signal strength ($+1\sigma/-1\sigma$ [%])					
	[0–15]	[15–45]	[45–85]	[85–125]	[125–165]	[165– ∞]
b veto	-10.1/-8.8	7.3/12.2	-6.3/3.1	-14.4/-4.8	-5.4/14.5	-7.9/17.8
lepton efficiency	-14.7/-3.9	4.5/15.1	-5.7/2.5	-13.2/-5.3	-0.2/7.6	-0.1/6.8
E_T^{miss} resolution	-12.5/0.0	15.4/-0.0	-12.8/-0.0	8.7/0.0	-20.9/-0.0	10.5/0.0
E_T^{miss} scale	-14.4/-6.8	-0.0/17.7	-6.1/-7.1	9.6/-20.9	2.3/32.4	2.5/2.6
lepton resolution	-12.5/-0.0	11.2/0.0	-2.4/0.0	-13.4/-0.0	9.9/0.0	-4.6/-0.0
electron momentum scale	-2.7/-13.1	15.9/9.9	10.8/-16.8	16.2/-33.1	30.9/-14.4	12.6/-10.9
muon momentum scale	-7.0/-10.7	11.8/8.9	1.1/-8.7	-0.7/-14.4	14.5/-4.6	8.0/-1.6
jet energy scale	-10.9/-10.1	9.0/9.0	-3.0/-2.9	-10.3/-8.9	0.3/3.4	5.2/3.1

1291 Type C uncertainties are related to the underlying assumption on the Higgs boson
 1292 production mechanism used to extract the fiducial cross sections. These are evaluated
 1293 using alternative response matrices that are obtained by varying the relative fraction of
 1294 the VBF and ggH components within the experimental uncertainty, as given by the CMS
 1295 combined measurement [73]. Three different response matrices are built, corresponding to
 1296 the nominal, scaled up, and scaled down VBF/ggH ratio. The nominal matrix assumes
 1297 the SM VBF/ggH ratio, while up- and down-scaled matrices are constructed by varying
 1298 the SM signal strengths within the experimental constraints for VBF and ggH in such a
 1299 way as to obtain the maximal variation of the VBF/ggH ratio allowed by the experimental
 1300 constraints. These three matrices are used to unfold the reconstructed spectrum with the
 1301 nominal VBF/ggH fraction, and obtain an uncertainty on the unfolded spectrum.

1302 4.8. Results

1303 In order to unfold the spectrum, the procedure described in section 4.7 has been pursued.
 1304 The statistical plus type A systematic uncertainties are propagated by the unfolding
 1305 procedure into the final spectrum, taking into account the signal strengths covariance
 1306 matrix. The type B systematic uncertainty has been propagated using the following
 1307 procedure: for each p_T^H bin, we compute the upper bound of the systematic band computing
 1308 the square sum of all the signal strength variations that deviate in the up direction with
 1309 respect to the bin central value, whether or not this variation corresponds to the up or
 1310 down shift of the systematic uncertainty. The same is done for the lower bound of the
 1311 systematic band. If both the up and down shifts of a given nuisance parameter lead to a
 1312 same direction variation of the signal strength, only the larger variation is considered.

¹³¹³ The unfolded p_T^H spectrum is shown in Fig. 4.27. Statistical, systematic, and theoretical uncertainties are shown as separate error bands in the plot. The unfolded spectrum is

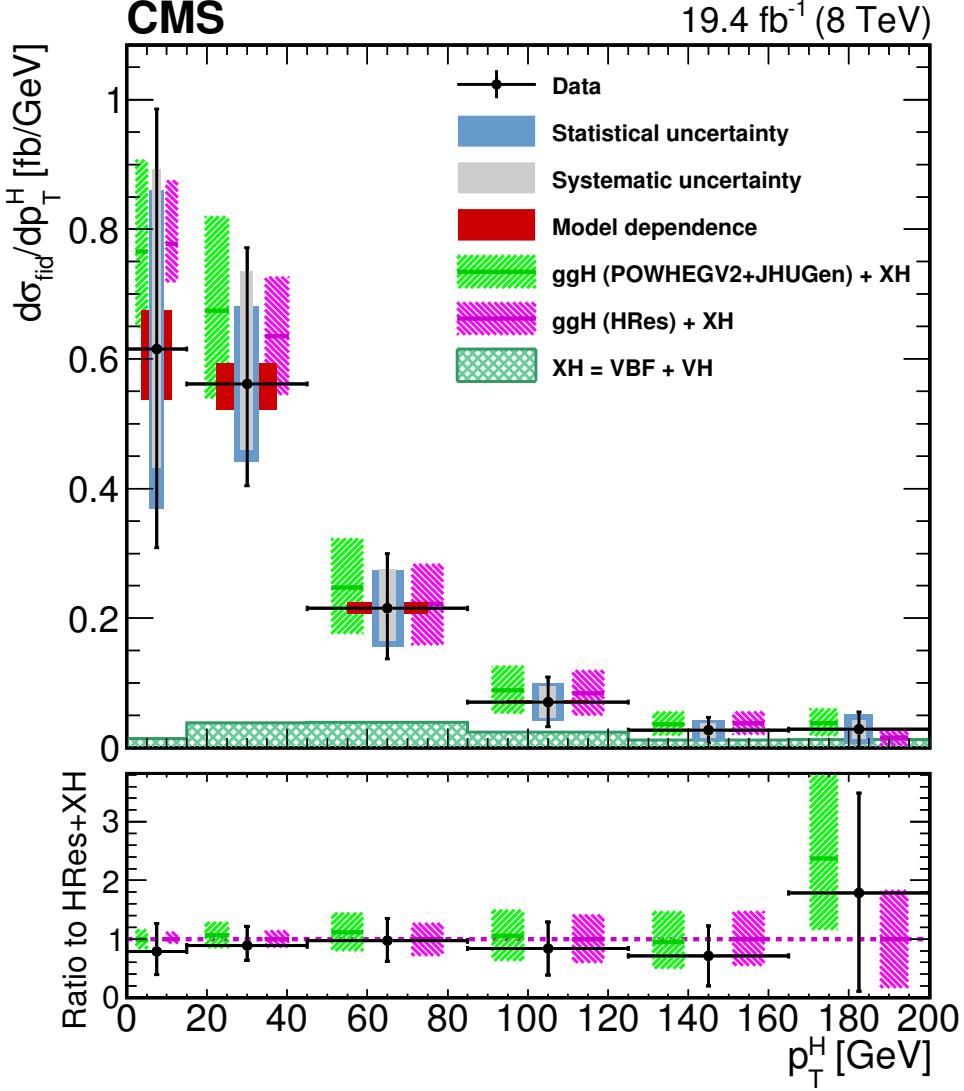


Figure 4.27.: Higgs boson production cross section as a function of p_T^H , after applying the unfolding procedure. Data points are shown, together with statistical and systematic uncertainties. The vertical bars on the data points correspond to the sum in quadrature of the statistical and systematic uncertainties. The model dependence uncertainty is also shown. The pink (and back-sashed filling) and green (and slashed filling) lines and areas represent the SM theoretical estimates in which the acceptance of the dominant ggH contribution is modelled by HRES and POWHEG V2, respectively. The subdominant component of the signal is denoted as $\text{XH}=\text{VBF}+\text{VH}$ and it is shown with the cross filled area separately. The bottom panel shows the ratio of data and POWHEG V2 theoretical estimate to the HRES theoretical prediction.

¹³¹⁴

¹³¹⁵ compared with the SM-based theoretical predictions where the ggH contribution is modelled

Table 4.12.: Differential cross section in each p_T^H bin, together with the total uncertainty and the separate components of the various sources of uncertainty.

p_T^H [GeV]	$d\sigma/dp_T^H$ [fb/GeV]	Total uncertainty [fb/GeV]	Statistical uncertainty [fb/GeV]	Type A uncertainty [fb/GeV]	Type B uncertainty [fb/GeV]	Type C uncertainty [fb/GeV]
0-15	0.615	+0.370/-0.307	± 0.246	± 0.179	+0.211/-0.038	+0.0782/-0.0608
15-45	0.561	+0.210/-0.157	± 0.120	± 0.093	+0.146/-0.041	+0.0395/-0.0327
45-85	0.215	+0.084/-0.078	± 0.059	± 0.037	+0.047/-0.034	+0.0089/-0.0084
85-125	0.071	+0.038/-0.038	± 0.029	± 0.017	+0.018/-0.017	+0.0018/-0.0022
125-165	0.027	+0.020/-0.019	± 0.016	± 0.009	+0.007/-0.007	+0.0003/-0.0006
165- ∞	0.028	+0.027/-0.027	± 0.023	± 0.012	+0.008/-0.007	+0.0002/-0.0006

1316 using the HRES and POWHEG V2 programs. The comparison shows good agreement
 1317 between data and theoretical predictions within the uncertainties. The measured values
 1318 for the differential cross section in each bin of p_T^H are reported together with the total
 1319 uncertainty in Table 4.12.

1320 Figure 4.28 shows the correlation matrix for the six bins of the differential spectrum.
 1321 The correlation of bins is defined as in Eq. (4.31).

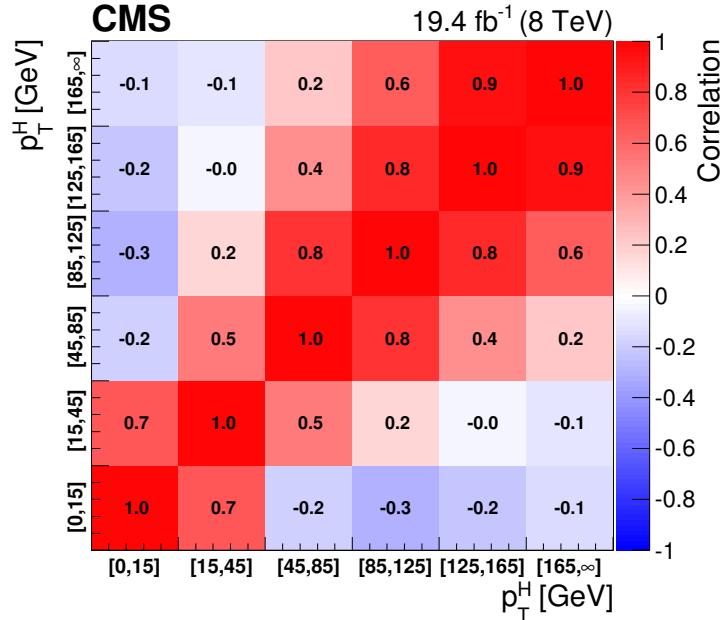


Figure 4.28.: Correlation matrix among the p_T^H bins of the differential spectrum.

1322 To measure the inclusive cross section in the fiducial phase space, the differential
 1323 measured spectrum is integrated over p_T^H . In order to compute the contributions of the bin
 1324 uncertainties of the differential spectrum to the inclusive uncertainty, error propagation
 1325 is performed taking into account the covariance matrix of the six signal strengths. For
 1326 the extrapolation of this result to the fiducial phase space, the unfolding procedure is not
 1327 needed, and the inclusive measurement has only to be corrected for the fiducial phase
 1328 space selection efficiency ϵ_{fid} . Dividing the measured number of events by the integrated
 1329 luminosity and correcting for the overall selection efficiency, which is estimated in simulation
 1330 to be $\epsilon_{fid} = 36.2\%$, the inclusive fiducial $\sigma \times \mathcal{B}$, σ_{fid} , is computed to be:

$$\sigma_{fid} = 39 \pm 8 \text{ (stat)} \pm 9 \text{ (syst)} \text{ fb} \quad , \quad (4.32)$$

1331 in agreement within the uncertainties with the theoretical estimate of 48 ± 8 fb, computed
 1332 integrating the spectrum obtained with the POWHEG V2 program for the ggH process and
 1333 including the XH contribution.

Chapter 5.

Search for the SM Higgs boson in the $H \rightarrow WW$ channel with the first 13 TeV LHC data

5.1. Introduction

In this chapter, the first search for the SM Higgs boson decaying to a W boson pair at 13 TeV is presented, using a total integrated luminosity of 2.3 fb^{-1} , collected during the 2015 proton proton data taking period of the LHC.

Final states in which the two W bosons decay leptonically are studied. Therefore, events with a pair of oppositely-charged leptons, exactly one electron and one muon, a substantial amount of missing transverse energy, E_T^{miss} , due to the presence of neutrinos in the final state, and either zero or one jet are selected. This signature is common to other processes, which enter the analysis as backgrounds. The main background comes from WW production, irreducible background that shares the same final states and can only be separated by the use of certain kinematic properties. Another important background is W+jets, where a jet can mimick a leptonic signature. Background coming from top quark events, i.e. $t\bar{t}$ and single top production, is also important, followed by other processes such as Drell-Yan, WZ, and other EWK production. The analysis strategy follows the one used during Run 1 in the same channel, described in Chapter 4, with a few different aspects that are described in the next sections.

With respect to 8 TeV, the ggH production cross section at 13 TeV is expected to increase of a factor of 2, thus raising the number of expected signal events. In addition, the cross section for the background processes is increasing as well. The WW production cross section increases of a factor of 1.8 and the $t\bar{t}$ cross section of a factor of 3.5, due to the enhancement of the gluon PDFs at higher center of mass energies.

1358 5.2. Data and simulated samples

1359 Data recorded in proton proton collisions at 13 TeV during 2015 was used in the analysis,
 1360 with a total integrated luminosity of 2.3 fb^{-1} . Single and double lepton triggers are used
 1361 similarly to the same analysis at 8 TeV. The HLT paths and descriptions of the triggers used
 1362 in this analysis are described in Tables 5.1 and 5.2 for electrons and muons respectively.

Table 5.1.: HLT paths related to Electrons

HLT Path	Description
HLT_Ele23_WP Loose_Gsf_v*	Single Electron trigger. Best trigger to be used for 2015 data. In HWW, we are using “Trigger safe” Id. Turn on is at around Ele $p_T = 30 \text{ GeV}$
HLT_Ele17_Ele12_CaloIdL_TrackIdL_IsoVL_DZ_v*	Double Electron Trigger. Best trigger to cover the turn on region from single electron trigger. “DZ” filter is also present. Its efficiency is also calculated separately.
HLT_Ele12_CaloIdL_TrackIdL_IsoVL_v*	This electron leg of HLT_Mu17_TrkIsoVVL_Ele12_CaloIdL_TrackIdL_IsoVL_v* same as Ele12 leg of double electron trigger.
HLT_Ele17_CaloIdL_TrackIdL_IsoVL_v*	This electron leg of HLT_Mu8_TrkIsoVVL_Ele17_CaloIdL_TrackIdL_IsoVL_v* same as Ele17 leg of double electron trigger.

1363 The trigger efficiencies are measured in data and applied on simulated events as described
 1364 in Sec. 4.2.1.

1365 Concerning the simulated samples, several different Monte Carlo (MC) generators were
 1366 used. In the simulation, ‘lepton’ includes also τ . Higgs signal samples have been simulated
 1367 in all channels with POWHEG v2 [34, 47, 74], designed to describe the full NLO properties
 1368 of these processes. In particular, for Higgs produced via gluon fusion [36], and vector-boson-
 1369 fusion (VBF) [37], the decay of the Higgs boson into two W boson and subsequently into
 1370 leptons was done using JHUGEN v5.2.5 [75]. For associated production with a vector boson
 1371 (W^+H , W^-H , ZH) [76], including gluon fusion produced ZH ($ggZH$), the Higgs decay was
 1372 done via PYTHIA 8.1 [56]. Alternative signal samples were produced with AMC@NLO [39],
 1373 or with POWHEG v2 but decayed via PYTHIA 8.1 for gluon fusion and VBF assuming a
 1374 Higgs boson mass of 125 GeV. In the following, the mass of the SM Higgs boson is assumed
 1375 to be 125 GeV.

1376 The WW production, irreducible background for the analysis, was simulated in different
 1377 ways. POWHEG v2 [77] was used for $q\bar{q}$ produced WW in different decays. The cross
 1378 section used for normalizing WW processes produced via $q\bar{q}$ was computed at next-to-next-
 1379 to-leading order (NNLO) [78]. In order to control the top quark background processes, the
 1380 analysis is performed with events that have no more than one high- p_T jet. The veto on

Table 5.2.: Muon trigger's elements description

HLT path	
HLT_IsoMu18_v*	single muon trigger
HLT_IsoTrMu20_v*	single muon trigger with tracker isolation
HLT_Mu17_TrkIsoVVL	leg for the HLT_Mu17_TrkIsoVVL_Mu8_TrkIsoVVL_DZ_v*, HLT_Mu17_TrkIsoVVL_TkMu8_TrkIsoVVL_DZ_v* and HLT_Mu17_TrkIsoVVL_Ele12_CaloIdL_TrackIdL_IsoVL_v* double lepton triggers
HLT_Mu8_TrkIsoVVL	leg for the HLT_Mu17_TrkIsoVVL_Mu8_TrkIsoVVL_DZ_v* and HLT_Mu8_TrkIsoVVL_Ele17_CaloIdL_TrackIdL_IsoVL_v* double lepton triggers
HLT_TkMu8_TrkIsoVVL	leg for the HLT_Mu17_TrkIsoVVL_TkMu8_TrkIsoVVL_DZ_v* double muon trigger
$DZ_{\mu\mu}$	efficiency of DZ cut in the HLT_Mu17_TrkIsoVVL_Mu8_TrkIsoVVL_DZ_v* and HLT_Mu17_TrkIsoVVL_TkMu8_TrkIsoVVL_DZ_v* double muon triggers, it is around 95%

1381 high- p_T jets enhances the importance of logarithms of the jet p_T , spoiling the convergence
 1382 of fixed-order calculations of the $q\bar{q} \rightarrow WW$ process and requiring the use of dedicated
 1383 resummation techniques for an accurate prediction of differential distributions [79, 80]. The
 1384 p_T of the jets produced in association with the WW system is strongly correlated with its
 1385 transverse momentum, p_T^{WW} , especially in the case where only one jet is produced. The
 1386 simulated $q\bar{q} \rightarrow WW$ events are reweighted to reproduce the p_T^{WW} distribution from the
 1387 p_T -resummed calculation.

1388 Gluon fusion produced WW was generated, with and without Higgs diagrams, using
 1389 MCFM v7.0 [81]. A $t\bar{t}$ sample dilepton sample was also generated using POWHEG v2. The
 1390 WW and $t\bar{t}$ samples produced specifically for this analysis are presented in Table 5.3. Other
 1391 background samples are used, a list of the most relevant ones is presented in Table 5.4.

1392 All processes are generated using the NNPDF2.3 [82, 83] parton distribution functions
 1393 (PDF) for NLO generators, while the LO version of the same PDF is used for LO generators.
 1394 All the event generators are interfaced to PYTHIA 8.1 [56] for the showering of partons
 1395 and hadronization, as well as including a simulation of the underlying event (UE) and
 1396 multiple interaction (MPI) based on the CUET8PM1 tune [84]. To estimate the systematic
 1397 uncertainties related to the choice of UE and MPI tune, the signal processes and the WW
 1398 events are also generated with two alternative tunes which are representative of the errors
 1399 on the tuning parameters. The showering and hadronization systematic uncertainty is

Table 5.3.: Simulated samples for $t\bar{t}$ and WW production. The $gg \rightarrow WW \rightarrow 2\ell 2\nu$ (H diagr.) sample includes both ggH production, the ggWW component and the interference.

Process	$\sigma \times \mathcal{B}$ [pb]
$t\bar{t} \rightarrow WW b\bar{b} \rightarrow 2\ell 2\nu b\bar{b}$	87.31
$q\bar{q} \rightarrow WW \rightarrow 2\ell 2\nu$	12.178
$gg \rightarrow WW \rightarrow 2\ell 2\nu$	0.5905
$gg \rightarrow WW \rightarrow 2\ell 2\nu$ (H diagr.)	0.9544

Table 5.4.: Simulated samples for other backgrounds used in the analysis.

Process	$\sigma \times \mathcal{B}$ [pb]
Single top	71.7
Drell-Yan ($10 \text{ GeV} < m_{\ell\ell} < 50 \text{ GeV}$)	20471.0
Drell-Yan ($m_{\ell\ell} > 50 \text{ GeV}$)	6025.26
$WZ \rightarrow 2\ell 2q$	5.5950
$ZZ \rightarrow 2\ell 2q$	3.2210
WWZ	0.1651
WZZ	0.05565
ZZZ	0.01398

1400 estimated by interfacing the same MC samples with the HERWIG ++ 2.7 parton shower [85,
1401 86]. For all processes, the detector response is simulated using a detailed description of the
1402 CMS detector, based on the GEANT4 package [48].

1403 The simulated samples are generated with distributions for the number of pileup
1404 interactions that are meant to roughly cover, though not exactly match, the conditions
1405 expected for the different data-taking periods. In order to factorize these effects, the number
1406 of true pileup interactions from the simulation truth is reweighted to match the data. The
1407 re-weighting is propagated automatically to both the in-time pile up and the out-of-time
1408 one. In Fig. 5.1, the effect of this reweighting on a sample enriched in Drell-Yan events is
1409 shown. Before the reweighting the simulation is presented in the open red histogram; after
1410 the reweighting, it is represented by the solid green histogram that matched well the data.
1411 In order to select this sample, events with two leptons with $p_T > 20 \text{ GeV}$, opposite sign,
1412 and same flavour, are selected only if $|m_{\ell\ell} - m_Z| < 15 \text{ GeV}$.

1413 The average number of pileup is approximately 11.5.

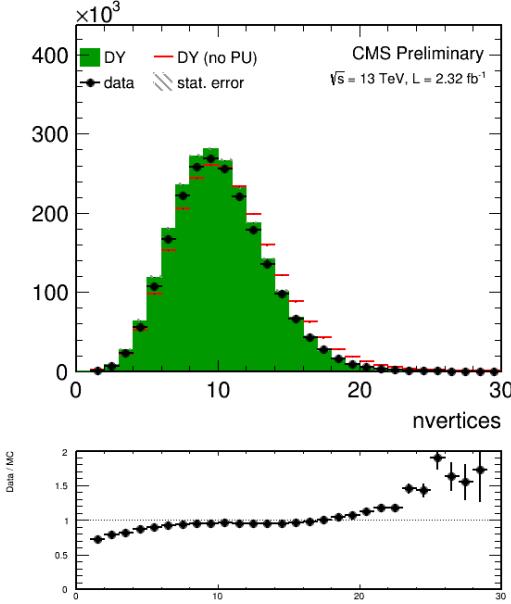


Figure 5.1.: Distributions of the number of vertices in a Drell-Yan enriched sample in data, together with the simulation before (red) and after (solid green) the pileup reweighting.

Different sources and calculations are used to obtain the cross sections for the different processes at 13 TeV. For Higgs signal, the cross sections used are the ones reported by the LHC Higgs Cross Section Working Group [87], computed at NNLO and NNLL QCD and NLO EW for gluon fusion, and at NNLO QCD and NLO EW for the rest of the production modes. The branching fractions are the ones reported in Ref. [30].

The cross section used for normalizing $q\bar{q}$ produced WW processes was computed at next-to-next-to-leading order (NNLO) [78]. The leading-order (LO) cross section for ggWW is obtained directly from MCFM. For gluon fusion, the difference between LO and NNLO cross sections is significantly big. A scale factor of 1.4 is theoretically calculated [88]. For the LO simulation of the interference between $gg \rightarrow WW$ and gluon fusion produced $H \rightarrow WW$ a k-factor of 1.87 is applied. This k-factor is obtained as the average between LO to NNLO ggH scale factor and LO to NLO ggWW scale factor.

The cross sections of the different single top processes are estimated by the LHC Top Working group [89] at NLO. The $t\bar{t}$ cross section is also provided by the LHC Top Working group [90], and it is computed at NNLO, with NNLL soft gluon resummation.

Drell-Yan (DY) production of Z/γ^* is generated using AMC@NLO [39]. Other multi-boson processes, such as WZ, ZZ, and VVV ($V=W/Z$), are generated with AMC@NLO and normalized to the cross section obtained at NLO in generation.

¹⁴³² All processes are generated using the NNPDF2.3 [82, 83] parton distribution functions
¹⁴³³ (PDF) for NLO generators, while the LO version of the same PDF is used for LO generators.
¹⁴³⁴ All the event generators are interfaced to PYTHIA 8.1 for the showering of partons and
¹⁴³⁵ hadronization, as well as including a simulation of the underlying event (UE) and multiple
¹⁴³⁶ interaction (MPI) based on the CUET8PM1 tune [84].

¹⁴³⁷ 5.3. Analysis strategy

¹⁴³⁸ 5.3.1. Event reconstruction

¹⁴³⁹ Regarding the electrons, muons, jets and E_T^{miss} definition and reconstruction, the standard
¹⁴⁴⁰ CMS recommendations described in Chapter 2 are used. The specific selections used in
¹⁴⁴¹ this analysis are briefly summarised below.

¹⁴⁴² Muons are identified according to the CMS recommendations for the medium working
¹⁴⁴³ point, with the addition of some extra cuts, as defined by the following selections:

- ¹⁴⁴⁴ • identified by the standard medium muon selection described in Sec. 2.4; **Not yet**
¹⁴⁴⁵ **defined :**
- ¹⁴⁴⁶ • $p_T > 10 \text{ GeV}$;
- ¹⁴⁴⁷ • $|\eta| < 2.4$;
- ¹⁴⁴⁸ • $|d_{xy}| < 0.01 \text{ cm}$ for $p_T < 20 \text{ GeV}$ and $|d_{xy}| < 0.02 \text{ cm}$ for $p_T > 20 \text{ GeV}$, d_{xy} being the
¹⁴⁴⁹ transverse impact parameter with respect to the primary vertex;
- ¹⁴⁵⁰ • $|d_z| < 0.1 \text{ cm}$, where d_z is the longitudinal distance of the muon track in the tracker
¹⁴⁵¹ extrapolated along the beam direction.

¹⁴⁵² For the muon isolation, the CMS recommended particle flow isolation based on the
¹⁴⁵³ tight working point is used, corresponding to a requirement on the isolation variable of
¹⁴⁵⁴ $ISO_{\text{tight}} < 0.15$. In addition a tracker relative isolation is also applied.

¹⁴⁵⁵ For the electron identification, the tight working point is used. In addition some
¹⁴⁵⁶ additional cuts to make the selection “trigger-safe” are included. This is done because the
¹⁴⁵⁷ electron triggers already include some identification and isolation requirements that are
¹⁴⁵⁸ based on the raw detector information, while the offline selections make use of particle flow
¹⁴⁵⁹ requirements. The “trigger-safe” selections are defined to make the the offline identification
¹⁴⁶⁰ and isolation requirements tighter with respect to the online triggers.

¹⁴⁶¹ The simulated events are corrected for the lepton trigger, identification and isolation
¹⁴⁶² efficiencies measured in data using the same techniques described in Sec. 4.3.1.

1463 Jets are defined clustering the particle flow objects using the anti- k_t algorithm with a
1464 distance parameter of 0.4. The CHS pileup mitigation technique is used. The L1, L2, L3
1465 and L2L3 jet energy correction described in Sec. 2.4 are applied. The reject jets coming from
1466 calorimeter or readout electronics noise, the loose working point for PF jet identification is
1467 used.

1468 The b-tagging algorithm for this analysis is chosen comparing the performances of
1469 different algorithms using simulations for signal and background contributions in the phase
1470 space defined by the analysis kinematic requirements. More precisely, two MC samples are
1471 used, one corresponding to the H \rightarrow WW \rightarrow 2 ℓ 2 ν signal produced via the ggH production
1472 mode and another corresponding to the t \bar{t} process. In fact, the first sample is enriched
1473 in light jets, i.e. originating by the hadronization of light quarks like u,d,c and s quarks,
1474 while the second sample is enriched in b jets, coming from the top quark decay. The b-veto
1475 efficiency, ϵ_{bveto} , is computed separately for the two samples and for the various b tagging
1476 algorithms. To compare the b tagging performance ϵ_{bveto} is computed for different working
1477 points, i.e. different selections on the specific b tagging discriminator, and the results are
1478 reported in the form of a ROC curve. The ROC curves corresponding to events with 0, 1
1479 and ≥ 2 jets are shown in Fig. 5.2. Events considered for this study are the ones passing
1480 the WW baseline selection.

1481 The ROC curves show that the cMVAv2 algorithm has the best performance for the
1482 analysis phase space among the algorithms taken into account. For both the CSVv2 and
1483 cMVAv2 algorithms, three working points are defined corresponding to the mistag rates¹ of
1484 10% for the loose, 1% for the medium and 0.1% for the tight working point. The distribution
1485 of the cMVAv2 discriminator associated to the leading jet both for the ggH and the t \bar{t} MC
1486 sample is shown in figure 5.3.

1487 In order to determine the best working point for this analysis a preliminary significance
1488 assessment is performed, using a complete analysis procedure in which only statistical
1489 effects are taken into account (no systematics are included). The significance assessment
1490 was performed using a two dimensional discriminating variable consisting of the dilepton
1491 invariant mass versus the transverse mass. The assessment was performed with the following
1492 leptonic selection:

- 1493 • two leptons, an electron and a muon with opposite charge, with leading lepton p_T
1494 greater than 20 GeV and sub-leading lepton p_T greater than 13 GeV;
- 1495 • no other lepton (electron or muon) with p_T greater than 10 GeV;
- 1496 • $m_{\ell\ell}$ greater than 12 GeV;
- 1497 • PF type 1 corrected MET greater than 20 GeV;
- 1498 • $p_T^{\ell\ell}$ greater than 30 GeV.

1497¹The mistag rate is defined as the probability for a light jet to be identified as a b-jet by the b tagging
algorithms.

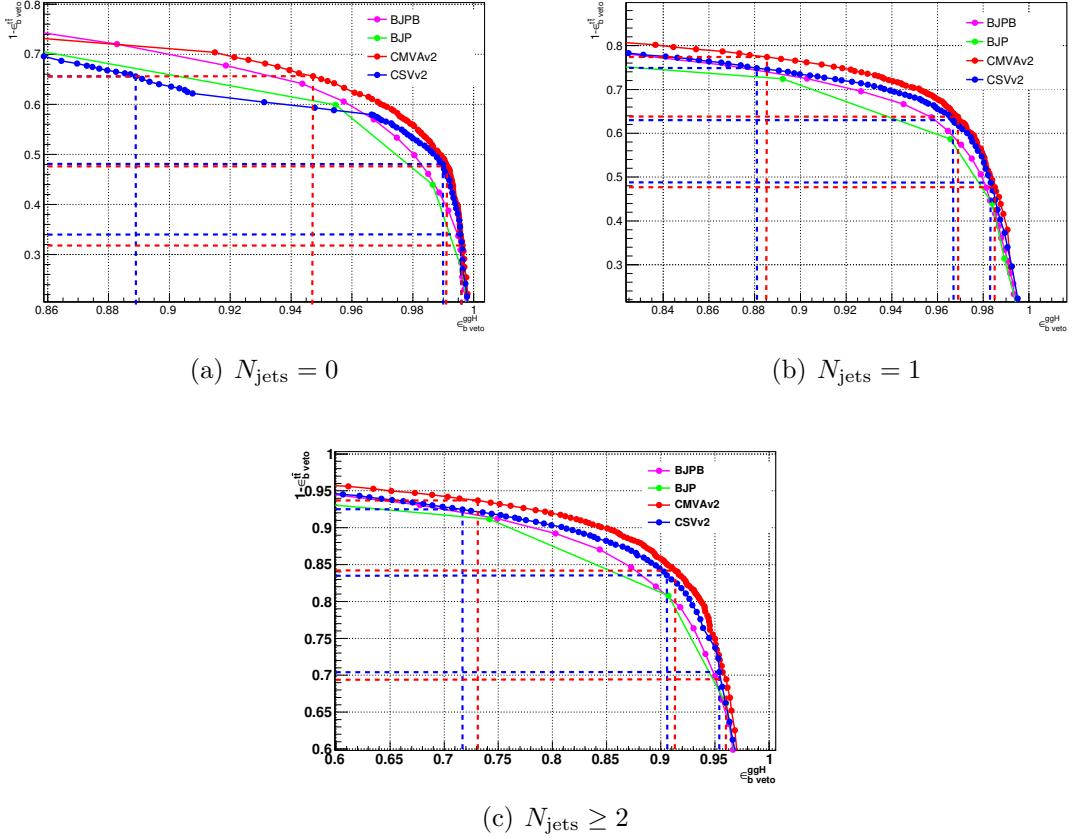


Figure 5.2.: ROC curve for the b veto efficiency on signal and background events. The blue and red lines point out the signal efficiency and the background rejection corresponding to the three working points considered for the CSVv2 and the cMVAv2 algorithms respectively.

1499 In addition to this global selection, two categories were identified:

- 1500 • 0 jets: no jets above 30 GeV, jets between 20 GeV and 30 GeV are b-vetoed with the
1501 cMVAv2 WP under study;
- 1502 • 1 jet: exactly 1 jet above 30 GeV, no b-tagged jets above 30 GeV according to the
1503 cMVAv2 WP under study.

1504 The two categories were eventually combined together and the significance assessment was
1505 repeated for the three working points. With these selection we find the significance values
1506 listed in Table 5.5 for the three working points.

1507 The working point providing the best significance in the combined 0 + 1 jets category is
1508 found to be the loose one.

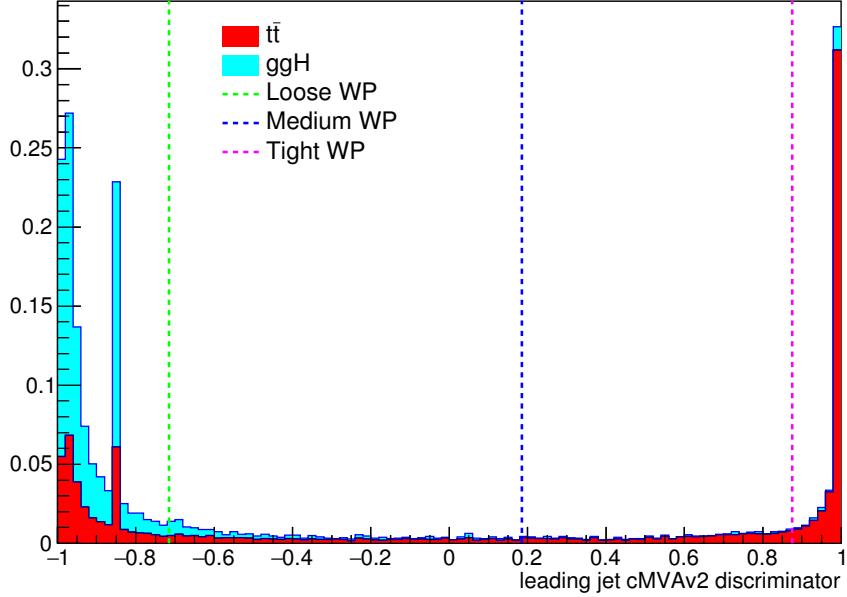


Figure 5.3.: cMVAv2 discriminator associated to the leading jet (with $p_T > 30$ GeV) both for the ggH and the $t\bar{t}$ processes. The two processes are normalized to unity and stacked. The vertical dashed lines show the discriminator value corresponding to the three working points.

Table 5.5.: Significance corresponding to the three working points and for different jet categories using a shape analysis.

Jet category	Loose WP (-0.715)	Medium WP (0.185)	Tight WP (0.875)
0 jets	2.022	2.043	2.036
1 jet	1.439	1.404	1.305
0 + 1 jets	2.481	2.479	2.420

1509 To correct for a possible different b tagging efficiency in data and simulation, the
 1510 simulated events are reweighted using scale factors computed in bins of the jet η and p_T .
 1511 These scale factors and the corresponding uncertainties are centrally calculated for each
 1512 working point, in such a way to be employable by all the CMS analyses. The prescription
 1513 to reweight the simulated events is the following. First of all one has to compute the b
 1514 tagging efficiency using the MC samples, $\varepsilon_{MC}(p_T, \eta, f)$, for the chosen working point in bins
 1515 of jet p_T and η . The efficiency has to be computed for different flavours f of the jets, b,

1516 c and light (u,d,s), using the jet matching information² which is available in all the MC
1517 samples. An MC-based event weight is then calculated computing the probability P_{MC} of a
1518 given b tagging configuration to occur, e.g.:

$$P_{\text{MC}} = \prod_{i \in b\text{-tagged-jets}} \varepsilon_{\text{MC}_i} \prod_{j \in \text{non-}b\text{-tagged-jets}} (1 - \varepsilon_{\text{MC}_j}) \quad (5.1)$$

1519 Afterwards, a similar probability is computed using data:

$$P_{\text{DATA}} = \prod_{i \in b\text{-tagged-jets}} SF_i \varepsilon_{\text{MC}_i} \prod_{j \in \text{non-}b\text{-tagged-jets}} (1 - SF_j \varepsilon_{\text{MC}_j}) , \quad (5.2)$$

1520 where SF_i is the provided scale factor value for the relevant jet flavour, p_T and η . Products
1521 in Eqs. 5.1 and 5.2 run over all jets. The event weight is finally given by the ration
1522 $P_{\text{DATA}}/P_{\text{MC}}$.

1523 The b tagging efficiencies to be fed into Eq. 5.1 and Eq. 5.2 are derived using $t\bar{t}$ simulated
1524 events and applying basic leptonic selections. These efficiencies are shown in Fig. 5.4 for
1525 light (a), c-jets (b) and b-jets (c), in bins of η and p_T . The uncertainties associated to the
1526 efficiencies are representative of the statistics of the simulated $t\bar{t}$ sample, and are computed
1527 according to a binomial distribution.

1528 The effect of the event reweighting is to correct the shape of the b tagging discriminator
1529 in simulation, moving events from the b tag region (discriminator greater than > -0.715)
1530 to the b veto region (discriminator < -0.715) and viceversa. A data/simulation comparison
1531 of the b tagging discriminator for the leading and subleading jets is performed to check the
1532 agreement after the application of the event weights. In order to evaluate the data/simulation
1533 agreement for b-jets, the data and simulation are compared in a top enriched control region,
1534 defined by the following requirements:

- 1535 • two leptons, an electron and a muon with opposite charge, with leading lepton p_T
1536 greater than 20 GeV and sub-leading lepton p_T greater than 15 GeV;
- 1537 • no other lepton (electron or muon) with p_T greater than 10 GeV;
- 1538 • lepton invariant mass greater than 50 GeV;
- 1539 • at least two jets with p_T greater than 30 GeV;
- 1540 • at least one of the two leading jets with cMVAv2 btagging score greater than -0.715
1541 (i.e. the loose working point).

1542 In order to evaluate the agreement for light jets, a second control region is defined, populated
1543 by Z+light jet events, defined as follows:

²There are a couple of techniques developed by the CMS Collaboration to assess the flavour of a reconstructed jet in simulation. The technique used here makes use of the flavour of the hadrons clustered into a jet.

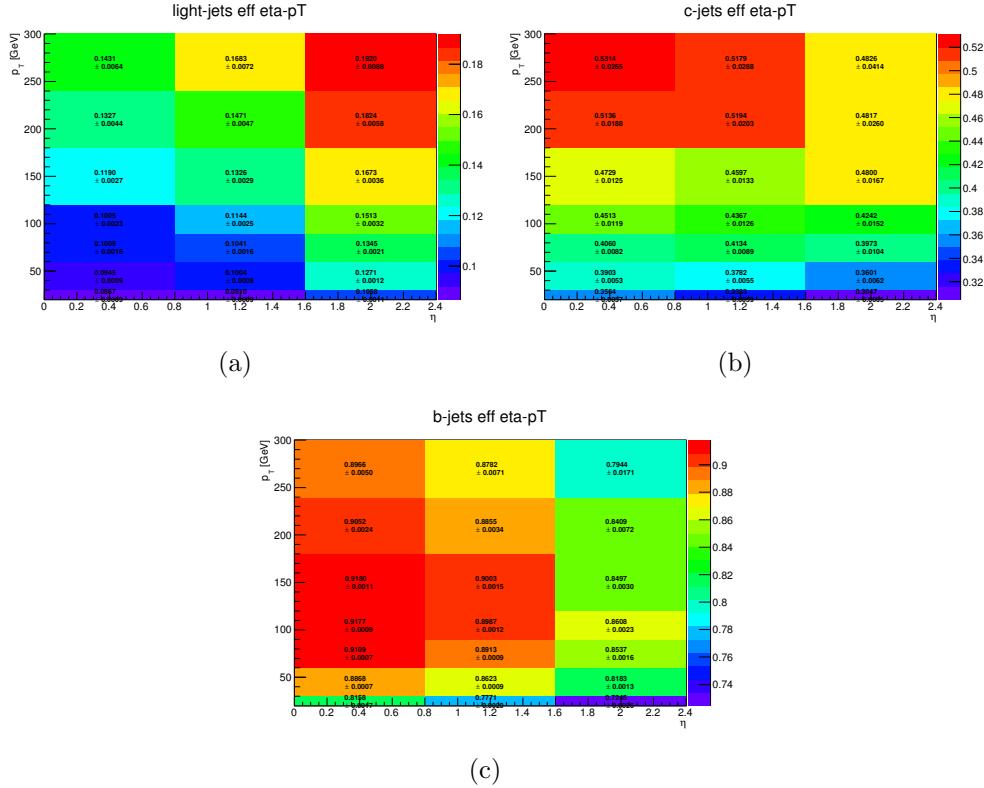


Figure 5.4.: B tagging efficiencies for light jets (a), c-jets (b) and b jets (c), as a function of η and p_T .

- 1544 • two leptons, two electrons or two muons with opposite charge, with leading lepton p_T
1545 greater than 20 GeV and sub-leading lepton p_T greater than 15 GeV.
 - 1546 • no other lepton (electron or muon) with p_T greater than 10 GeV.
 - 1547 • lepton invariant mass greater between 80 GeV and 110 GeV.
 - 1548 • at least two jets with p_T greater than 30 GeV.
 - 1549 • at least one jet above 30 GeV.
 - 1550 • no jets above 20 GeV with a TCHE score above 2.1.
- 1551 Although a Z+jets sample is dominated by light flavor jets, a b-veto on an alternative
1552 algorithm (TCHE) is applied to reduce the contamination from b-jets, especially above the
1553 cMVAv2 cut. This helps mitigating possible data/simulation discrepancies in the modeling
1554 of the heavy/light flavour ratio. The comparison between data and simulation after the
1555 event reweighting is shown in Figs. 5.5 and 5.6 for the b-jets and light jets enriched control
1556 regions, respectively.

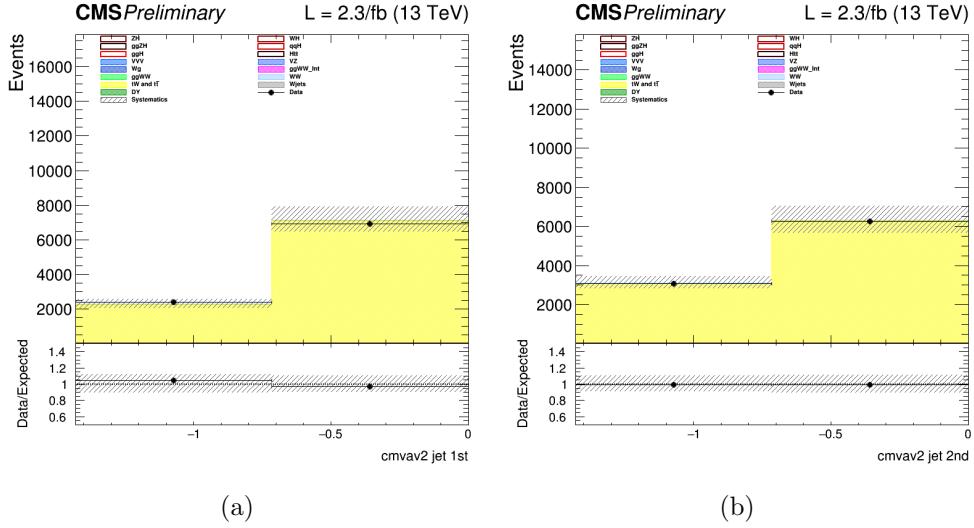


Figure 5.5.: B tagging cMVAv2 discriminator for the leading (a) and the subleading (b) jet in the b-jets enriched control region.

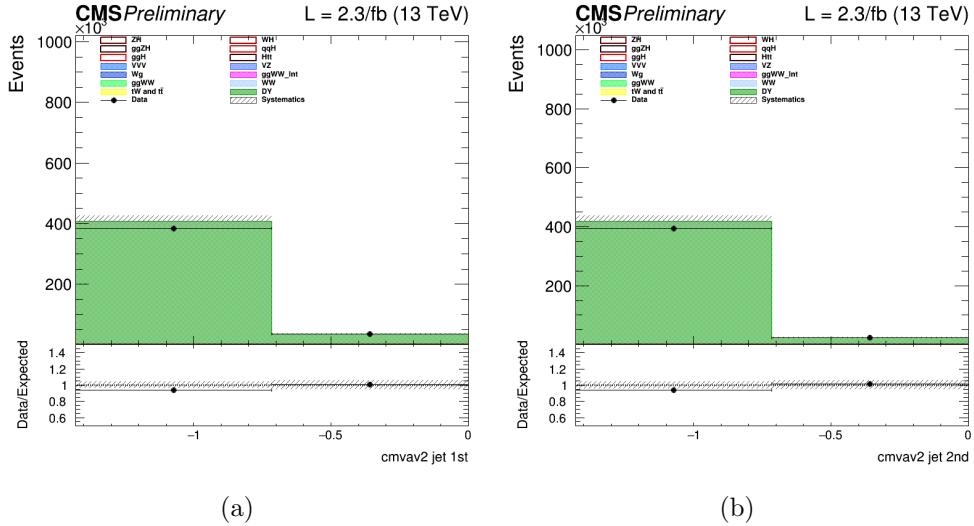


Figure 5.6.: B tagging cMVAv2 discriminator for the leading (a) and the subleading (b) jet in the light jets enriched control region.

1557 5.3.2. Event selection and background rejection

1558 Since the ggH production mechanism, which is the main production mode for a Higgs
 1559 mass of around 125 GeV, is characterized by the emission of few jets arising from initial
 1560 or final state radiation, this analysis is limited to events with no jets or one jet. Due to
 1561 the large DY background in di-electrons and di-muons events, only the $e\mu$ final state is

studied in this early Run 2 data analysis, including the indirect contribution from τ leptons decaying to electron or muons. Exactly one electron and one muon are required to be reconstructed in the event with opposite charges and a minimum p_T of 10 (13) GeV for the muon (electron). One of the two leptons should also have a p_T greater than 20 GeV and both leptons are required to be well identified and isolated to reject fake leptons and leptons coming from QCD sources. To suppress background processes with three or more leptons in the final state, such as ZZ, WZ, Z γ , W γ , or tri-boson production, no additional identified and isolated lepton with $p_T > 10$ GeV should be reconstructed. The low $m_{\ell\ell}$ region dominated by QCD production of leptons is not considered in the analysis and $m_{\ell\ell}$ is requested to be higher than 12 GeV. To suppress the background arising from DY events decaying to a τ lepton pair which subsequently decays to an $e\mu$ final state and suppress processes without genuine E_T^{miss} , a minimal E_T^{miss} of 20 GeV is required. The DY background is further reduced by requesting $p_T^{\ell\ell} > 30$ GeV. Finally the contribution from leptonic decays of single top and t \bar{t} production is reduced by requesting that no jets with $p_T > 20$ GeV are identified by the b tagging algorithm as originating from a b quark in the event.

The requirements described above define the WW baseline selection. After those requirements the data sample is dominated by events arising from the non-resonant WW production and t \bar{t} production. To further reduce the effect of these backgrounds on the signal sensitivity, the events are categorized depending on the jet multiplicity, counting jets with $p_T > 30$ GeV. Events with zero associated jets mainly arise from the WW production, while WW and t \bar{t} productions have a similar contribution in the category with one jet. Higher jet multiplicity categories, which are sensitive to other Higgs production mechanisms, such as VBF, are not included in this analysis, given the very low expected yield for other production modes with the analysed integrated luminosity.

Distributions of some variables of interest for the 0 and 1 jet categories separately, but merging the $e\mu$ and μe final states together, are shown in Figs. 5.7, 5.8 and 5.9 after applying the WW baseline selections, with the addition of a cut on $m_{\ell\ell}$ to remove the Higgs signal contribution ($m_{\ell\ell} > 80$ GeV), and a cut on m_T to be orthogonal to the Z $\gamma^* \rightarrow \tau\tau$ background control region ($m_T > 60$ GeV).

The W+jets background, where one jet can be misidentified as a lepton, is a sub-dominant background in the phase space defined by the analysis kinematic requirements. The 0 ad 1 jets categories are further split according to the lepton flavour to $e\mu$ and μe , where the first lepton refers to the leading one. In this way an improvement of about 10% in terms of the signal significance can be achieved, exploiting the different W+jets background contribution in the two categories. Indeed the probability for a jet to be misidentified as an electron or a muon is not the same.

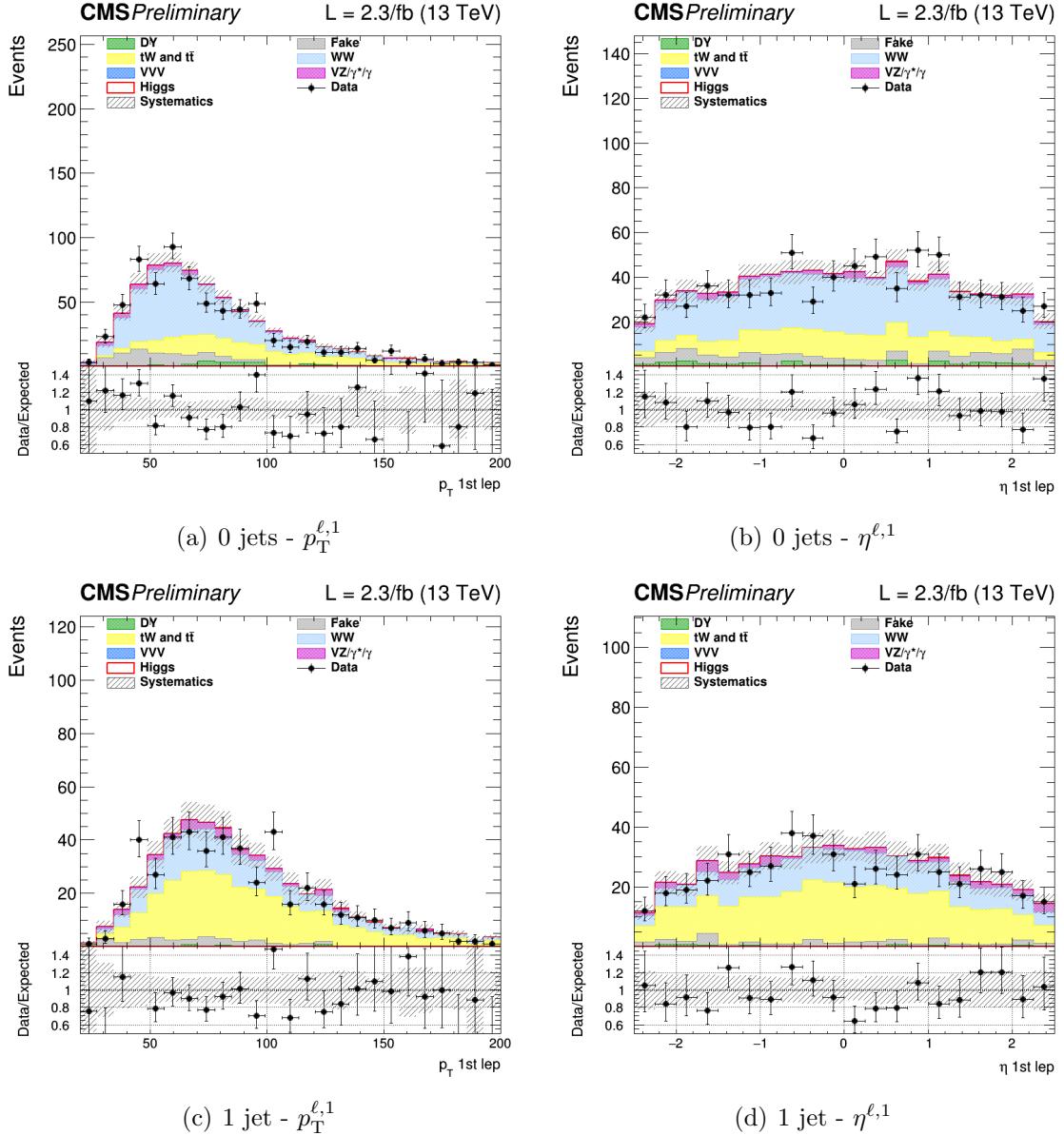


Figure 5.7.: Distributions of p_T (left) and η (right) of the leading lepton for events with 0 jet (upper row) and 1 jet (lower row), for the main backgrounds (stacked histograms), and for a SM Higgs boson signal with $m_H = 125$ GeV (superimposed and stacked red histogram) at the WW selection level. The last bin of the histograms includes overflows. The simulation of the WW background is normalized to data.

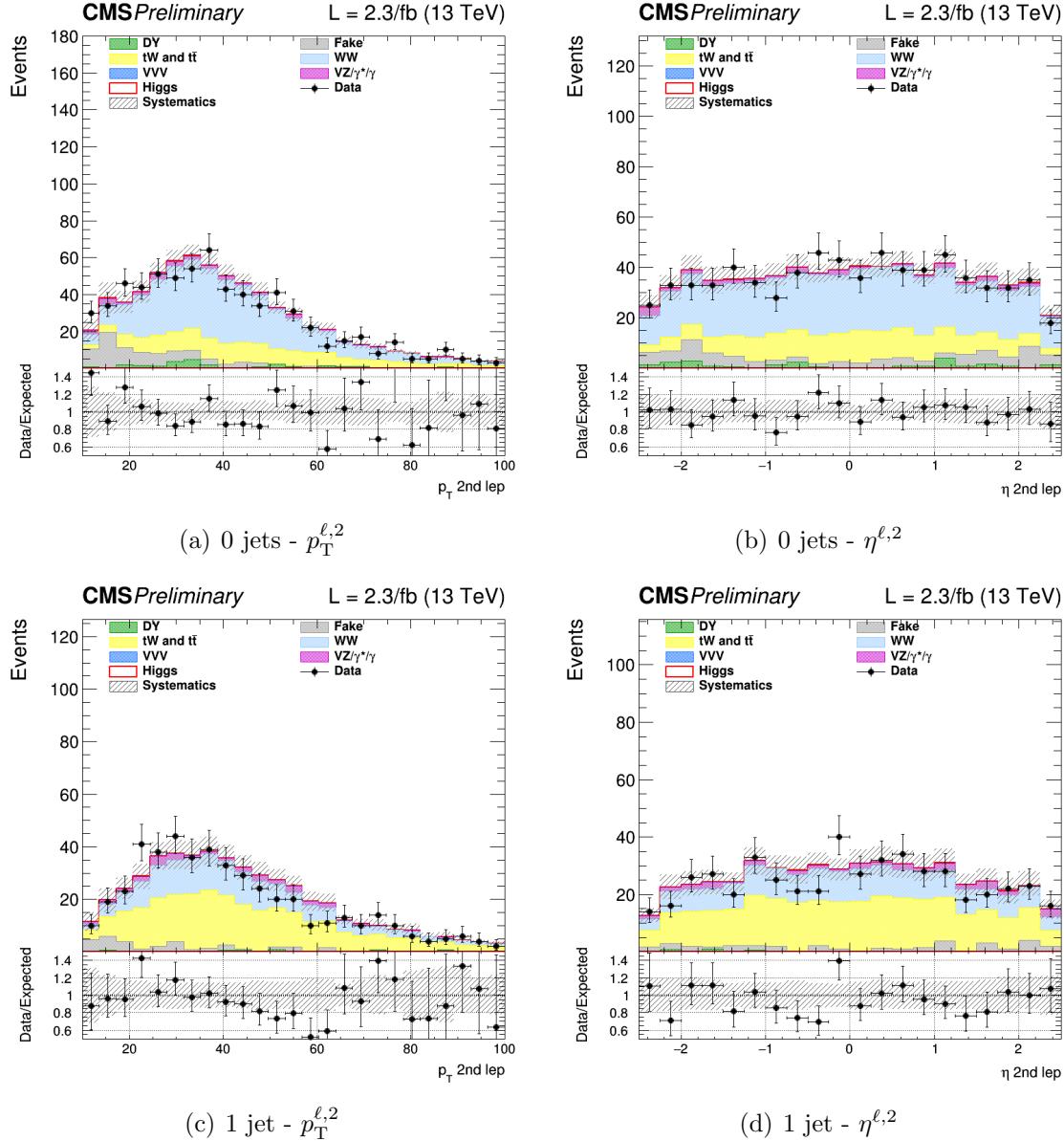


Figure 5.8.: Distributions of p_T (left) and η (right) of the subleading lepton for events with 0 jets (upper row) and 1 jet (lower row), for the main backgrounds (stacked histograms), and for a SM Higgs boson signal with $m_H = 125$ GeV (superimposed and stacked red histogram) at the WW selection level. The last bin of the histograms includes overflows. The simulation of the WW background is normalized to data.

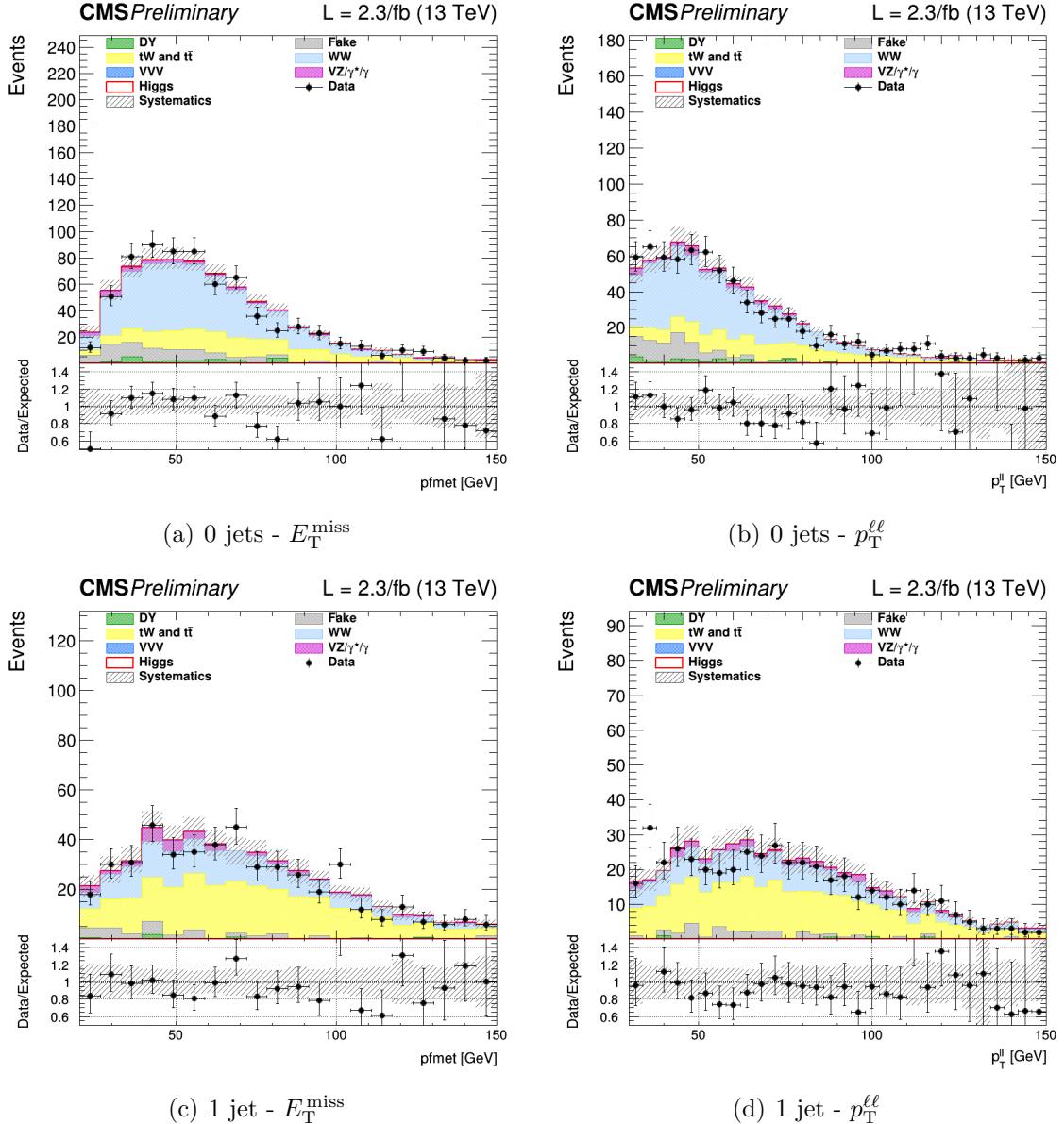


Figure 5.9.: Distributions of E_T^{miss} (left) and $p_T^{\ell\ell}$ (right) for events with 0 jets (upper row) and 1 jet (lower row), for the main backgrounds (stacked histograms), and for a SM Higgs boson signal with $m_H = 125$ GeV (superimposed and stacked red histogram) at the WW selection level. The last bin of the histograms includes overflows. The simulation of the WW background is normalized to data.

¹⁵⁹⁹ **5.3.3. Signal extraction**

¹⁶⁰⁰ To extract the Higgs boson signal contribution in the four previously mentioned categories,
¹⁶⁰¹ a similar approach to the one used in the Run 1 analysis [31] is pursued. The analysis is
¹⁶⁰² based on two-dimensional templates of $m_{\ell\ell}$ versus m_T to discriminate signal and background
¹⁶⁰³ contributions. The $m_{\ell\ell}$ template is defined using 5 bins from $m_{\ell\ell} = 10 \text{ GeV}$ up to $m_{\ell\ell} =$
¹⁶⁰⁴ 110 GeV , while for the m_T template 7 bins are defined in the range $60 \text{ GeV} < m_T < 200 \text{ GeV}$.
¹⁶⁰⁵ The phase space with $m_T < 60 \text{ GeV}$ is used as an orthogonal control region to extract the
¹⁶⁰⁶ normalization of the DY background. A binned maximum likelihood fit to the signal and
¹⁶⁰⁷ background two-dimensional templates is performed to extract the signal strength in the
¹⁶⁰⁸ four categories.

¹⁶⁰⁹ Distributions of the $m_{\ell\ell}$ and m_T variables after the WW level selection are shown in
¹⁶¹⁰ Fig. 5.10 for the 0 and 1 jet categories separately, but merging the $e\mu$ and μe final states
¹⁶¹¹ together.

¹⁶¹² The statistical methodology used to interpret the data and to combine the results from
¹⁶¹³ the independent 0-jet and 1-jet categories in the $e\mu$ and μe final states has been developed
¹⁶¹⁴ by the ATLAS and CMS collaborations in the context of the LHC Higgs Combination
¹⁶¹⁵ Group [73, 91]. The number of events in each category and in each bin of two-dimensional
¹⁶¹⁶ template is modelled as a Poisson random variable, with a mean value given by the sum
¹⁶¹⁷ of the contributions from all the processes under consideration. Systematic uncertainties
¹⁶¹⁸ are represented by individual nuisance parameters with log-normal distributions. The
¹⁶¹⁹ uncertainties affect the overall normalization of the signal and backgrounds as well as the
¹⁶²⁰ shape of the predictions across the distribution of the observables. Correlation between
¹⁶²¹ systematic uncertainties in different categories are taken into account.

¹⁶²² **5.4. Background estimation**

¹⁶²³ The main background processes affecting the analysis signature, non-resonant WW pro-
¹⁶²⁴ duction and top quark processes, are estimated using data. Backgrounds arising from an
¹⁶²⁵ experimental misidentification of the objects, such as W+jets (also called “Fake”), are
¹⁶²⁶ estimated using data as well. The other minor backgrounds are generally estimated directly
¹⁶²⁷ from simulation as described in the following subsections.

¹⁶²⁸ **5.4.1. WW background**

¹⁶²⁹ The quark-induced WW background is simulated with NLO accuracy in perturbative
¹⁶³⁰ QCD, and the transverse momentum of the diboson system is reweighted to match the
¹⁶³¹ NNLO+NNLL accuracy from theoretical calculations [79, 80]. However, given the large

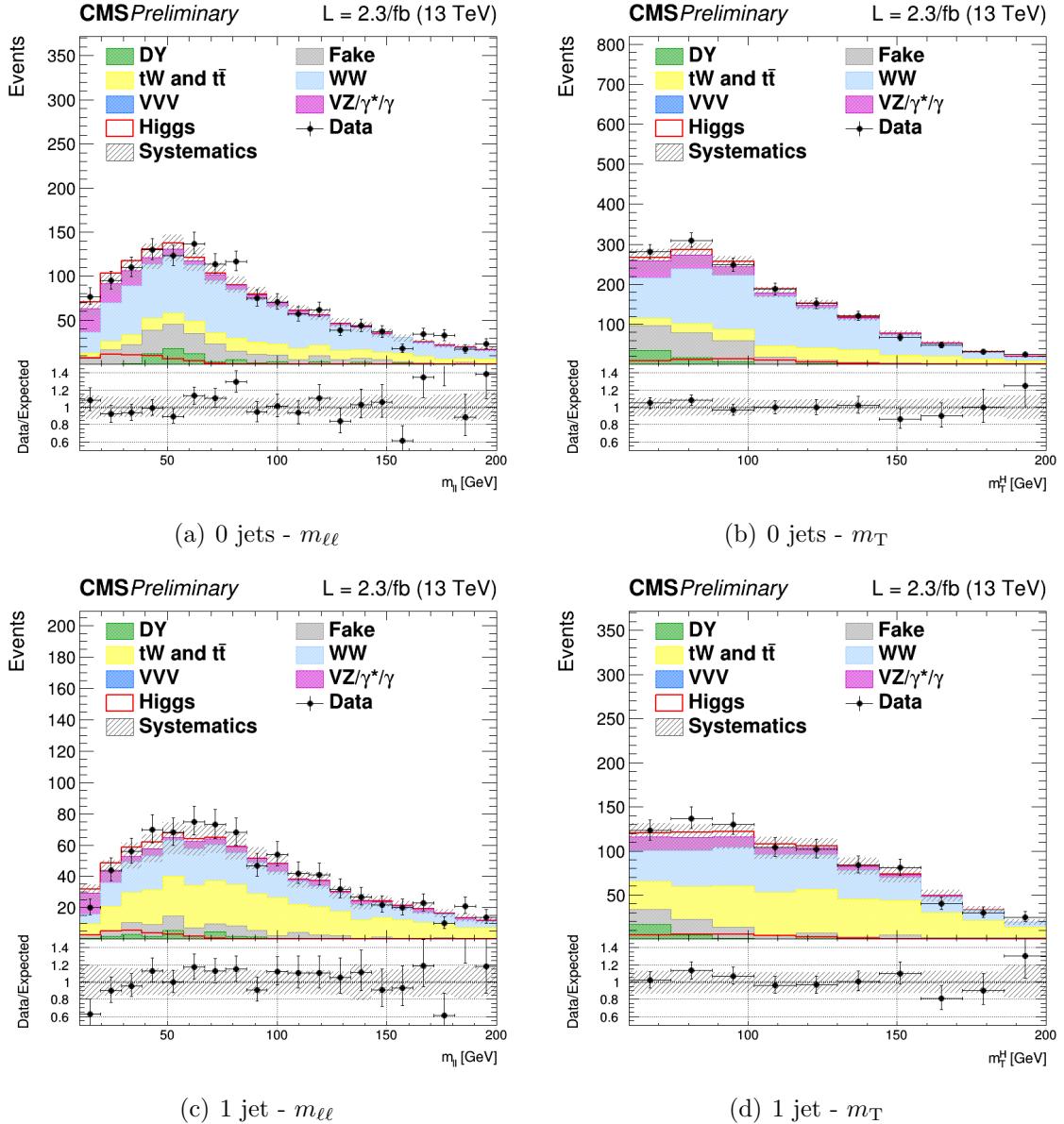


Figure 5.10.: Distributions of $m_{\ell\ell}$ (left) and m_T (right) for events with 0 jets (upper row) and 1 jet (lower row), for the main backgrounds (stacked histograms), and for a SM Higgs boson signal with $m_H = 125$ GeV (superimposed and stacked red histogram) at the WW selection level. The last bin of the histograms includes overflows. The simulation of the WW background is normalized to data.

1632 uncertainties on the jet multiplicity distribution associated to this process, the normalization
 1633 of this background is measured from data separately for the 0 and 1 jet categories. The
 1634 normalization k-factors are extracted directly from the fit together with the signal strengths,
 1635 leaving the WW normalization free to float separately in the two jet multiplicity categories.
 1636 An orthogonal control region for the WW background normalization estimation is not
 1637 needed in this case, owing to the different $m_{\ell\ell}$ - m_T shape for signal and background.

1638 The gluon-induced WW production is sub-dominant with respect to the quark-induced
 1639 production, and its shape and normalization is fully taken from simulation, scaling the
 1640 cross section to the theoretical prediction with NLO accuracy [88].

1641 5.4.2. Top quark background

1642 As explained in Sec. 5.3, the production of top quark pairs represents one of the dominant
 1643 backgrounds in this analysis given its large cross section and a similar final state compared
 1644 to the signal. A b-jet veto, based on the *cMVAv2* b tagging algorithm, is used to suppress
 1645 this background and a reweighting procedure is applied on top of the simulated events to
 1646 correct for different b tagging efficiency in data and simulation.

1647 The top quark background normalization is measured using data, defining a b-jets
 1648 enriched control region by inverting the b-jet veto. More precisely, the b-jets enriched
 1649 control region for the 0-jet category is defined with the same WW baseline selection but
 1650 requiring at least one jet with $20 < p_T < 30 \text{ GeV}$ to be identified as a b jet and no other
 1651 jets with $p_T > 30 \text{ GeV}$. For the 1-jet category, the b-jets enriched region is defined requiring
 1652 exactly one jet with $p_T > 30 \text{ GeV}$ identified as a b-jet. To reduce other backgrounds in
 1653 these two regions, the dilepton mass has to be greater than 50 GeV . Distributions of the
 1654 $m_{\ell\ell}$ and m_T variables in the b-jets enriched control regions after applying the data driven
 1655 estimation are shown in Figure 6.5, for the 0 and 1 jet categories separately.

1656 The top quark background normalization is constrained during the fit procedure sepa-
 1657 rately in the two jet categories, by means of the control regions defined above, which are
 1658 treated in the fit as two additional categories.

1659 5.4.3. Jet-induced (or Fake) background

1660 One of the primary source belonging to this category arises from the misidentification of
 1661 leptons in W+jets processes in the 0 jet category. Also, semileptonic $t\bar{t}$ decays contribute
 1662 especially for higher jet multiplicities. Multijet production and hadronic $t\bar{t}$ decays are also
 1663 taken into account, but have a much smaller contribution.

1664 This background is fully estimated using data, with the technique described in Sec. 4.4.3.
 1665 To check the agreement of the background estimated in this way with data, a control sample

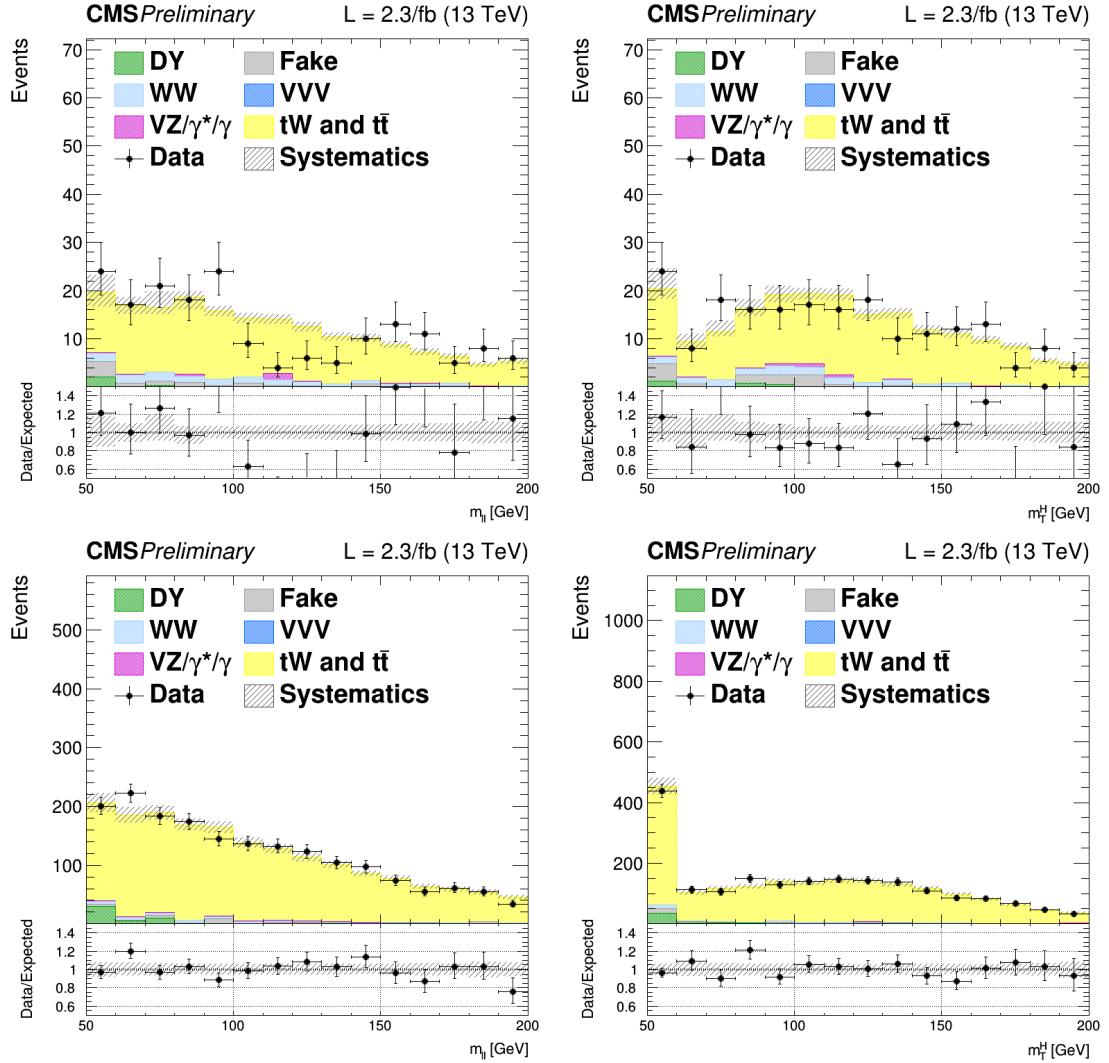


Figure 5.11.: Distributions of $m_{\ell\ell}$ (left) and m_T (right) for events with 0 jet (top) and 1 jet (bottom) in top enriched phase space. Scale factors estimated from data are applied. The first (last) bin includes underflows (overflows).

enriched in jet-induced events is defined. The events in the control sample are selected applying the WW baseline requirements but requesting an $e\mu$ pair with same charge, which significantly suppresses the WW and $t\bar{t}$ processes. The $m_{\ell\ell}$ distributions in this control region for the 0 and 1 jet categories are shown in Fig. 5.12. From the crosscheck in this control region, a global normalization factor of 0.8 is derived and applied to the jet-induced background.

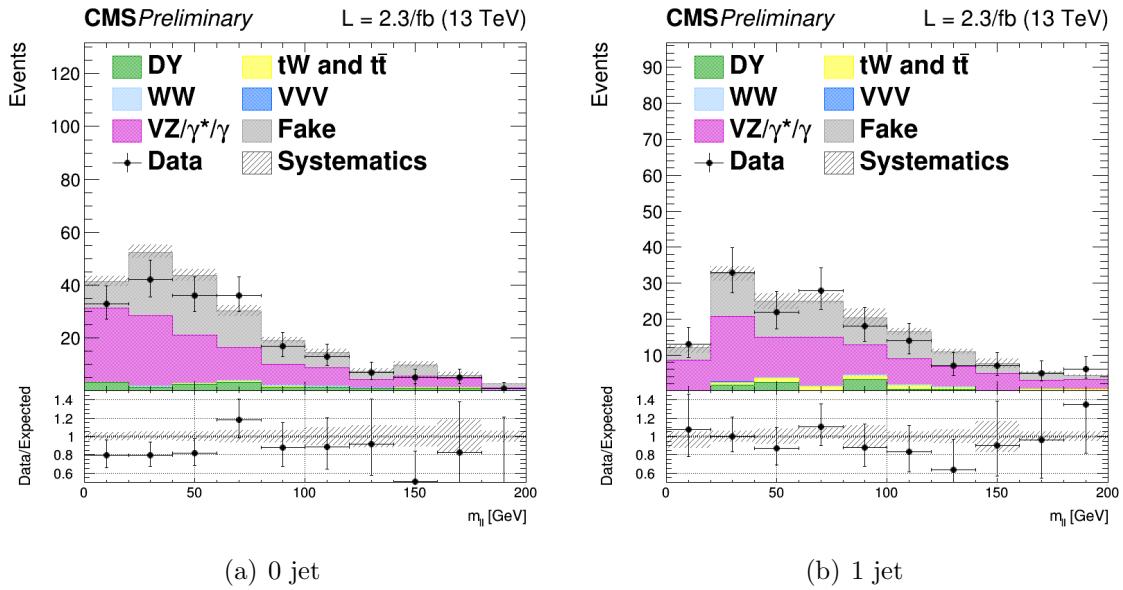


Figure 5.12.: Control plots for $m_{\ell\ell}$ in a fakes enriched phase space for events with 0 and 1 jet with $p_T > 30$ GeV, in $e\mu$ final state. Fake contribution has been scaled by 0.8 to match data.

5.4.4. DY background

This background contributes to the analysis phase space because of the Z/γ^* decays to a pair of τ leptons, which consequently decays to an $e\mu$ pair. This background process is predominant in the low m_T region, which is used as an orthogonal control region to determine the background normalization in the 0 and 1 jet categories separately. In particular this control region is defined by selecting events with $m_T < 60$ GeV and $30 \text{ GeV} < m_{\ell\ell} < 80$ GeV. The $m_{\ell\ell}$ distributions in these control regions for the 0 and 1 jet categories are shown in Fig. 5.13.

As for the top quark background, the normalization of this background in the 0 and 1 jet categories, is constrained directly in the fit by means of the control regions, which are treated as two additional categories.

The kinematics of this background is taken from simulation, after reweighting the Z boson p_T spectrum to match the observed distribution measured in data. In fact, this

variable is not well reproduced by the MC generator used for simulating this process, especially in the bulk of the distribution, the discrepancy being ascribed to the missing contribution from resummed calculations.

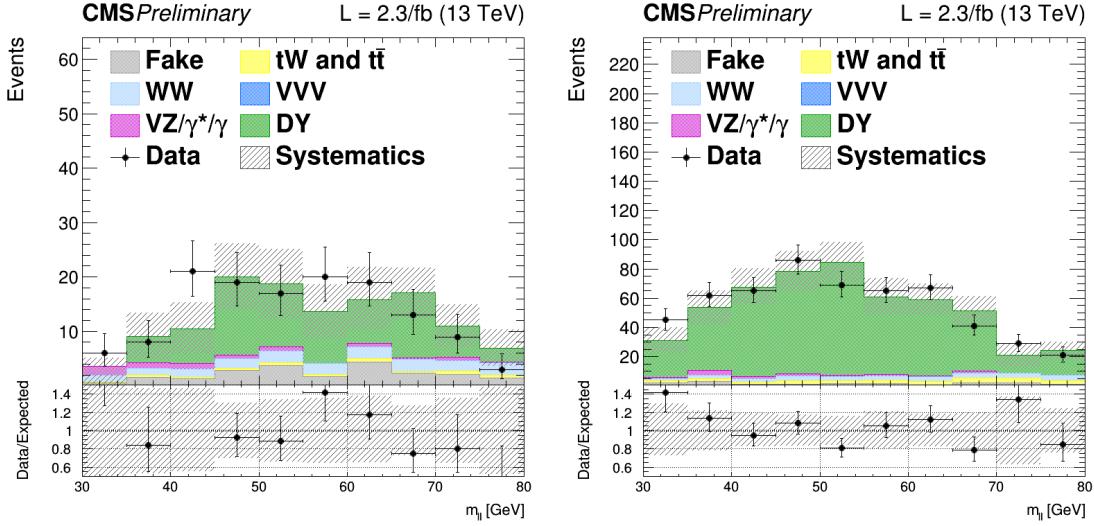


Figure 5.13.: Distributions of $m_{\ell\ell}$ for events with 0 jet (left) and 1 jet (right) in the $DY \rightarrow \tau\tau$ enriched control region. Scale factors estimated from data are applied.

1688 5.4.5. Other backgrounds

1689 The $W\gamma^*$ and the WZ electroweak processes can be gathered in the same physical process,
 1690 although the final state kinematics is rather different. In particular, the invariant mass of
 1691 the leptons arising from the γ^* decays is generally below 4 GeV, while the leptons from the
 1692 Z boson decay are characterized by a larger invariant mass. Another background which can
 1693 be experimentally identical to those is the $W\gamma$ production, where a real photon is produced
 1694 in association with a W boson and consequently undergoes a photon conversion to leptons
 1695 due to the interaction with the material constituting the first layers of the silicon tracker.

1696 All these backgrounds may contribute to the signal phase space whenever one of the three
 1697 leptons escape from the detector acceptance or is not identified. The shape and cross section
 1698 of these backgrounds are taken from simulation. The only exception is the normalization of
 1699 the $W\gamma^*$ background, being this process dominant in the low $m_{\ell\ell}$ region, which is scaled
 1700 to data defining a proper control region. The control region is defined selecting events
 1701 with three isolated muons, with $p_T > 10,5$ and 3 GeV for the first three leading muons
 1702 respectively. The selection is further defined by $E_T^{\text{miss}} < 25$ GeV and E_T^{miss} projected to the
 1703 leading muon < 45 GeV. The pair of muons with the smallest invariant mass is taken as
 1704 coming from the γ^* decay. The k-factor measured in data for this background to be applied
 1705 in the simulation is 1.98 ± 0.54 .

1706 All remaining backgrounds from di-boson and tri-boson production, which are of mi-
1707 nor importance in the analysis phase space, are normalized according to their expected
1708 theoretical cross sections.

1709 5.5. Systematic uncertainties

1710 The systematic uncertainties affecting this measurement can be divided into three categories:
1711 the uncertainties on the background estimation, experimental uncertainties and theoretical
1712 uncertainties.

1713 The first category includes the uncertainties related to the background normalization
1714 and shape. For the non-resonant WW production the shape is taken from simulation.
1715 The input normalization to the fit is set to the expected value from simulation, and an
1716 unconstrained nuisance parameter with a flat distribution is associated to this number.
1717 This is done separately for the two jet categories.

1718 The top quark background shape is taken from simulation after correcting for the b
1719 tagging scale factors. An uncertainty due to these scale factors is included and affects
1720 both the normalization and the shape of the top quark background. The uncertainties on
1721 the normalization are treated similarly to the WW background case, but constraining the
1722 corresponding nuisances by means of the two control regions orthogonal to the signal phase
1723 space. A similar procedure is used for the DY background.

1724 Effects due to experimental uncertainties are studied by applying a scaling and smearing
1725 of certain variables related to the physics objects, e.g. the p_T of the leptons, followed by
1726 a subsequent recalculation of all the correlated variables. This is done for simulation, to
1727 account for possible systematic mismodeling.

1728 All experimental sources, except luminosity, are treated both as normalization and shape
1729 uncertainties, and are correlated among the signal and background processes and all the
1730 categories. The following experimental uncertainties are considered:

- 1731 • the uncertainty determined by the CMS online luminosity monitoring, 2.7% for the
1732 first data collected at $\sqrt{s} = 13 \text{ TeV}$;
- 1733 • the acceptance uncertainty associated with the combination of single and double lepton
1734 triggers, which is 2%;
- 1735 • the lepton reconstruction and identification efficiencies uncertainties, that are in the
1736 range 0.5-5% for electrons and 1-7% for muons depending on p_T and η ;
- 1737 • the muon momentum and electron energy scale and resolution uncertainties, that
1738 amount to 0.01-0.5% for electrons and 0.5-1.5% for muons depending on p_T and η ;

- 1739 • the jet energy scale uncertainties, that vary between 1-11% depending on the p_T and
1740 η of the jet;
- 1741 • the E_T^{miss} resolution uncertainty, that is taken into account by propagating the corre-
1742 sponding uncertainties on the leptons and jets;
- 1743 • the scale factors correcting the b tagging efficiency and mistagging rate, that are varied
1744 within their uncertainties. This systematic uncertainty is anticorrelated between the
1745 top control regions and the other ones.

1746 The uncertainties in the signal and background production rates due to theoretical
1747 uncertainties include several components, which are assumed to be independent: the PDFs
1748 and α_s , the underlying event and parton shower model, and the effect of missing higher-order
1749 corrections via variations of the renormalization and factorization scales.

1750 The effects of the variation of PDFs, α_s and renormalization/factorization QCD scales,
1751 mainly affect the signal processes, being the most important backgrounds estimated using
1752 data driven techniques. However, the uncertainties on minor backgrounds that are estimated
1753 from simulation are taken into account. These uncertainties are split in the uncertainties
1754 on the cross section, which are computed by the LHC cross section working group [92],
1755 and on the selection efficiency [93]. The PDFs and α_s signal cross section normalization
1756 uncertainties are $^{+7.4\%}_{-7.9\%}$ and $^{+7.1\%}_{-6.0\%}$ for ggH and $\pm 0.7\%$ and $\pm 3.2\%$ for VBF Higgs production
1757 mechanism. The PDFs and α_s acceptance uncertainties are less than 1% for gluon- and
1758 quark-induced processes. The effect of the QCD scales variation on the selection efficiency is
1759 around 1-3% depending on the specific process. To estimate these uncertainties, the events
1760 are reweighted according to different QCD scales or different PDF sets and the selection
1761 efficiency is recomputed each time. For the QCD scale uncertainty the maximum variation
1762 with respect to the nominal value is taken as the uncertainty. For the case of PDF and α_s
1763 uncertainties, the distribution of the selection efficiency is built taking into account all the
1764 replicas in the NNPDF3.0 set and the uncertainty is estimated as the standard deviation of
1765 that distribution.

1766 In addition, the categorization of events based on jet multiplicity introduces additional un-
1767 certainties on the ggH production mode related to missing higher order corrections. These un-
1768 certainties are evaluated following the prescription described in Sec. [subsec:stewart-tackman
1769] and correspond to 5.6% for the 0-jet and 13% for the 1-jet bin categories.

1770 The underlying event uncertainty is estimated by comparing two different PYTHIA
1771 8 tunes, while parton shower modelling uncertainty is estimated by comparing samples
1772 interfaced with the PYTHIA 8 and HERWIG++ parton shower programs. The effect on the
1773 ggH (VBF) signal expected yield is about 5% (5%) for the PYTHIA 8 tune variation and
1774 about 7% (10%) for the parton shower description.

1775 Other specific theoretical uncertainties are associated to some backgrounds. An uncer-
1776 tainty on the ratio of the $t\bar{t}$ and tW cross sections is included. Indeed, these two processes

1777 are characterized by a different number of b-jets in the final state (2 b-jets for $t\bar{t}$ and 1
1778 for tW) and the b-veto acts differently for the two. A variation of the relative ratio of the
1779 cross sections can thus cause a migration of events from the 0 to the 1 jet categories and
1780 viceversa. The corresponding uncertainty is of 8%, according to the theoretical cross section
1781 calculations [89, 90].

1782 The $gg \rightarrow WW$ background LO cross section predicted by the MCFM generator is scaled
1783 to the NLO calculation, applying a k-factor of 1.4 with an uncertainty of 15% [88]. The
1784 interference term between the $gg \rightarrow WW$ and the ggH signal is also included and simulated
1785 with LO accuracy using MCFM. The k-factor to scale the interference term is 1.87, given by
1786 the geometrical average of the LO to NNLO $gg \rightarrow H \rightarrow WW$ scale factor (2.5) and the LO
1787 to NLO $gg \rightarrow WW$ scale factor (1.4). The uncertainty on this value is estimated as the
1788 maximum variation with respect to the two scale factors mentioned above, and is found
1789 to be of 25%. Anyway, with the current amount of integrated luminosity, the interference
1790 contribution is found to be negligible.

1791 For what the $qq \rightarrow WW$ background shape is concerned, an uncertainty related to
1792 the diboson p_T reweighting is evaluated varying the renormalization, factorization and
1793 resummation QCD scales.

1794 Finally, the uncertainties due to the limited statistical accuracy of the MC simulations
1795 are also taken into account, including an independent uncertainty for each bin of the
1796 two-dimensional distribution, and for each category. The uncertainty for a certain bin
1797 and process is given by the standard deviation of the Poisson distribution with mean
1798 corresponding to the number of MC events in that bin.

1799 5.6. Results

1800 The expected and observed signal significance are shown in Table 5.6 for all the categories
1801 separately. Also, the observed signal strengths and the corresponding uncertainties are
1802 shown. The best fit signal strength obtained combining all the categories together is found
1803 to be $0.3^{+0.5}_{-0.5}$, corresponding to an observed significance of 0.7σ , to be compared with the
1804 expected significance of 2.0σ for a Higgs boson mass of 125 GeV.

1805 Maybe I should add the nuisances impact plots...

Table 5.6.: Observed and expected significance and signal strength the SM Higgs boson with a mass of 125 GeV for the 0-jet and 1-jet, μe and $e\mu$, categories.

Category	Expected significance	Observed significance	σ/σ_{SM}
0-jet μe	1.1	1.3	$1.13^{+0.9}_{-0.9}$
0-jet $e\mu$	1.3	0.4	$0.33^{+0.7}_{-0.7}$
1-jet μe	0.8	0	$-0.11^{+0.5}_{-1.7}$
1-jet $e\mu$	0.9	0	$-0.54^{+1.4}_{-1.4}$
0-jet	1.6	1.3	$0.71^{+0.6}_{-0.5}$
1-jet	1.2	0	$-0.56^{+1.0}_{-1.0}$
Combination	2.0	0.7	$0.33^{+0.5}_{-0.5}$

Chapter 6.

1806 **Search for high mass resonances 1807 decaying to a W boson pair with first 1808 13 TeV LHC data**

1809 **6.1. Introduction**

1810 In this chapter, a search for a high mass spin-0 particle (from now on denoted as X)
1811 in the $X \rightarrow WW \rightarrow \ell\nu\ell'\nu'$ decay channel is presented, where ℓ and ℓ' refer to an different
1812 flavour lepton pair, i.e. $e\mu$. The search is based upon proton-proton collision data samples
1813 corresponding to an integrated luminosity of up to 2.3 fb^{-1} at $\sqrt{s} = 13 \text{ TeV}$, recorded by
1814 the CMS experiment at the LHC during 2015. This analysis represents a general extension
1815 of the SM Higgs boson search presented in 5 and is performed in a range of heavy scalar
1816 masses from $M_X = 200 \text{ GeV}$ up to 1 TeV , extending the range studied in a similar analysis
1817 performed using Run 1 LHC data [94], which provided upper limits on the production cross
1818 section of new scalar resonances up to 600 GeV .

1819 Despite the discovery of a particle consistent with the SM Higgs boson in 2012, there is a
1820 possibility that this particle is only a part of a larger Higgs sector, and hence only partially
1821 responsible of the EW symmetry breaking. This can be achieved in different theoretical
1822 models that extends the SM, such as the two-Higgs-doublet models [95–97], or models in
1823 which the SM Higgs boson mixes with a heavy EW singlet, which predict the existence
1824 of an additional resonance at high mass, with couplings similar to those of the SM Higgs
1825 boson, as most recently described in [98, 99].

1826 This analysis reports a generic search for a scalar particle with different resonance decay
1827 widths hypothesis, produced via the ggH and VBF production mechanisms. The results
1828 can then be interpreted in terms of different theoretical models. This analysis is heavily
1829 based on the SM Higgs search described in 5, in terms of physics objects, selections and
1830 background estimation. The differences and similarities are discussed in this chapter.

1831 6.2. Data and simulated samples

1832 The data sets, triggers, pile up reweighting, lepton identification and isolation used in this
1833 analysis are the same as the SM Higgs search and are described in Sec. 5.2.

1834 Also, the same MC simulations are used for the background processes, the only exception
1835 being the DY background, for which the MG5_AMC@NLO generator is used with LO
1836 QCD accuracy, matching together events with up to four jets in addition to the vector
1837 boson with the MLM [100] matching scheme. Given that this analysis aims to probe regions
1838 of phase space where the DY contribution is very small, like in the high transverse mass
1839 region, the usage of a simulation of the inclusive DY process may lead to large uncertainties
1840 due to the limited simulation statistics in the sample. To partially overcome this issue,
1841 different DY samples are generated in restricted portions of the phase space defined by
1842 the H_T variable, i.e. the scalar sum of all the partons p_T in the event. For $H_T < 100$ GeV
1843 the inclusive simulation is used, while different samples are used for higher values of H_T .
1844 The samples are merged using the parton level information, and it has been verified that a
1845 smooth transition between different H_T regions is achieved, as shown in Fig. 6.1. The DY
1846 LO cross section obtained from the simulation is scaled using the LO to NNLO k-factor of
1847 1.23.

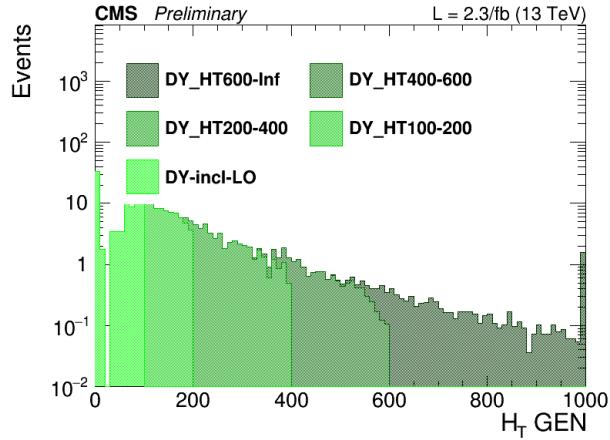


Figure 6.1.: Generator level H_T distribution for the merged DY sample.

1848 In order to perform the resonance search in a large part of the mass spectrum, several
1849 signal samples for the gluon-gluon fusion and the vector boson fusion mechanisms have been
1850 generated corresponding to different Higgs boson masses in the range between 200 GeV
1851 and 1 TeV. The signal width for each mass point corresponds to the one expected for a
1852 SM Higgs boson at that mass. The samples are produced with a mass step of 50 GeV from
1853 250 to 800 GeV and of 100 GeV from 800 to 1000 GeV. A finer stepping is used between
1854 200 and 250 GeV. All the signal samples are generated with the POWHEG V2 generator,
1855 interfaced with the JHUGEN v6.2.8 generator, which handles the decay of the Higgs boson
1856 to $W^+W^- \rightarrow 2\ell 2\nu$.

1857 The interference effects among $gg \rightarrow X \rightarrow WW$, $gg \rightarrow WW$ and $gg \rightarrow H \rightarrow WW$ are evaluated
1858 using the MCFM and JHUGEN generators, as implemented in the MELA (Matrix Element
1859 Likelihood Approach) framework [75]. Details about the interference effects are given in
1860 Sec. 6.3.

1861 6.3. Analysis strategy

1862 The analysis strategy for the first results on the high mass search in the $W^+W^- \rightarrow 2\ell 2\nu$
1863 decay channel closely follows the strategy presented in the 13 TeV SM Higgs search in the
1864 $H \rightarrow W^+W^- \rightarrow 2\ell 2\nu$ channel regarding the 0 and 1 jet categories. In addition a dedicated
1865 category to the VBF production mechanism is added, given that this production mode is
1866 particularly important in the high mass region. Indeed, assuming a SM Higgs boson, the
1867 ratio of cross sections $\sigma_{VBF}/\sigma_{ggH}$ ¹ increases with the Higgs boson mass, making the VBF
1868 production mechanism more and more important as the mass of the resonance approaches
1869 to high values.

1870 This analysis is affected essentially by the same background processes as the SM Higgs
1871 boson search, with the difference that in this case the SM Higgs boson processes, including
1872 all production modes, are treated as backgrounds.

1873 In addition to requiring the events to pass the single or double lepton triggers, exactly
1874 one electron and one muon are required to be reconstructed in the event with opposite
1875 charges and a minimum p_T of 20 GeV for both the muon and electron. Both leptons are
1876 required to be well identified and isolated to reject fake leptons and leptons coming from
1877 decays in flight. To suppress background processes with three or more leptons in the final
1878 state, such as diboson or triboson production, events with any additional identified and
1879 isolated lepton with $p_T > 10$ GeV are rejected. To suppress the contribution of the SM
1880 production of the Higgs boson at 125 GeV, $m_{\ell\ell}$ is requested to be higher than 50 GeV. The
1881 other event requirements are identical to the 125 GeV Higgs boson search and are described
1882 in Sec. 5.3.2.

1883 In addition to the 0 and 1 jet categories, a specific category sensitive to the VBF
1884 production mode is defined exploiting the characteristic signature of this process, where
1885 two energetic jets are emitted in the forward region of the detector and with large $\Delta\eta$
1886 gap. Events belonging to the VBF-enriched category are selected by requiring at least
1887 two jets with $p_T > 30$ GeV, an invariant mass $m_{jj} > 500$ GeV and a gap in pseudorapidity
1888 $\Delta\eta_{jj} > 3.5$.

1889 In addition to the transverse mass variable m_T , which is used in the analysis selection
1890 to define the DY background control region, an additional variable is defined, that from

1The ggH notation is used for the gluon-gluon fusion production mode, even in the cases where a non-SM Higgs boson is created in the process.

now on will be labelled as “improved transverse mass” m_T^i . This variable is defined as the invariant mass of the four momentum resulting from the sum of the two leptons four momenta ($p_{\ell\ell}, \vec{p}_{\ell\ell}$) and four momentum $\mathbf{E}_T^{\text{miss}} = (E_T^{\text{miss}}, \vec{p}_T^{\text{miss}})$, i.e.:

$$m_T^i = \sqrt{(p_{\ell\ell} + E_T^{\text{miss}})^2 - (\vec{p}_{\ell\ell} + \vec{p}_T^{\text{miss}})^2} . \quad (6.1)$$

This variable allows having a better sensitivity to different resonance mass hypothesis as shown in Fig. 6.2, where the shape of the m_T^i variable is shown for different SM Higgs mass hypothesis and it is compared to the standard m_T variable. The usage of this variable also provide a good discriminating power between signal and background, which depends on the particular signal mass hypothesis.

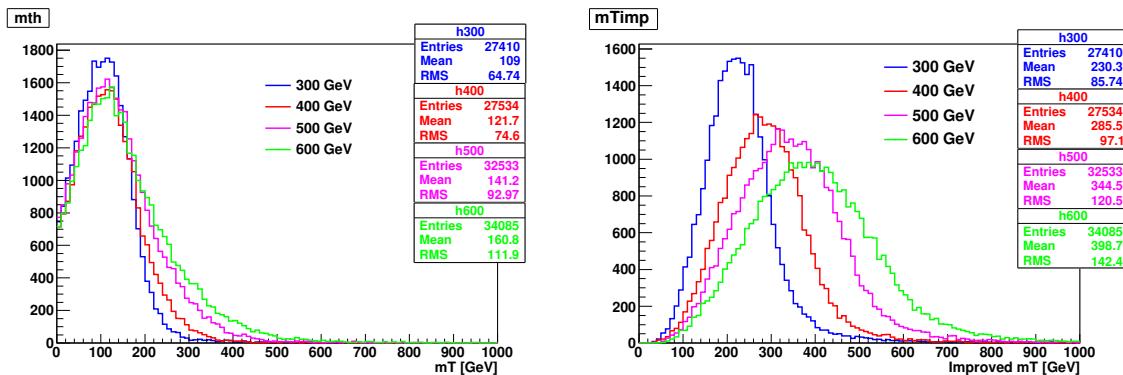


Figure 6.2.: Distribution of the m_T and m_T^i variables at generator level for different resonance mass hypothesis.

The signal extraction is based on a binned maximum likelihood fit using the m_T^i distribution for signal and background contributions as templates. The m_T^i template is defined using the following bin boundaries:

- **0/1 jet:** $[100, 150, 200, 250, 300, 350, 400, 450, 500, 600, 700, 1000]$,
- **VBF:** $[100, 150, 200, 250, 300, 350, 400, 500, 700, 1000]$,

where the first number represents the lower edge of the first bin while the other numbers represent the upper edges. The last bin is an overflow bin.

In order to test different resonance decay widths hypotheses, the signal samples, which are generated with a decay width corresponding to the expected value for a SM Higgs boson at that mass (Γ_{SM}), are reweighted to obtain the desired width value (Γ'). In particular the following values are used: $\Gamma' = \Gamma_{\text{SM}}$, $\Gamma' = 0.49 \times \Gamma_{\text{SM}}$, $\Gamma' = 0.25 \times \Gamma_{\text{SM}}$ and $\Gamma' = 0.09 \times \Gamma_{\text{SM}}$. The reweighting is performed at generator level by computing the ratio of two relativistic

1911 Breit Wigner distributions with different decay widths, $f(E, \Gamma', M_X)/f(E, \Gamma_{\text{SM}}, M_X)$, where:

$$f(E) \propto \frac{1}{(E^2 - M^2)^2 + M^2 \Gamma^2} . \quad (6.2)$$

1912 Here, $f(E, \Gamma_{\text{SM}}, M_X)$ represents the distribution used for the simulation of the signal at
1913 a mass M_X , and $f(E, \Gamma', M_X)$ the distribution with the new decay width. Each event
1914 is multiplied by this ratio (which depends on the energy E of the event) to obtain the
1915 reweighted distribution.

1916 When a resonance with a non negligible width is considered, it is important to take
1917 into account the interference effects both with the $\text{gg} \rightarrow \text{WW}$ background and the SM Higgs
1918 boson off-shell tail. A study of the interference effects for a resonance X produced through
1919 the gluon fusion mechanism is performed within the MCFM+JHUGEN framework, and
1920 including NNLO corrections for cross section using HNNLO program [101] based on
1921 MCFM. The matrix element package MELA supports all of these processes and allows fast
1922 MC re-weighting and optimal discriminant calculation. The basic idea of this approach is to
1923 compute the matrix elements of the processes under study with the MCFM and JHUGEN
1924 generators, including the interference terms, and using these matrix elements to compute
1925 an event weight used to reweight the simulated samples. Using this approach the simulated
1926 events can be reweighted according to different scenarios, for instance including some or all
1927 the interference terms, allowing a detailed study of the interference contribution. The effect
1928 of the various interference terms for the M_X variable at generator level is shown in Fig. 6.3,
1929 after having applied the WW baseline selections. As can be observed the contribution of
1930 the interference of the scalar resonance with the $\text{gg} \rightarrow \text{WW}$ background and with SM Higgs
1931 boson have opposite sign and partially cancel out. This cancellation effect is different for
1932 different resonance masses and depends on the event selection. In particular the interference
1933 term with the SM Higgs off-shell tail is positive for values below M_X while it turns negative
1934 above M_X . The contribution of the interference with the $\text{gg} \rightarrow \text{WW}$ background is instead
1935 characterized by an opposite sign lineshape, thus leading to a partial cancellation when
1936 considering the total interference.

1937 The effect of the resulting interference contribution including all the different terms is
1938 shown in Fig. 6.4 for the m_T^i signal templates, in the three categories separately and for
1939 different M_X hypotheses.

1940 The interference contribution is thus not negligible, especially for large values of M_X ,
1941 and is included in the analysis as part of the signal contribution. More specifically, during
1942 the fit procedure the signal yield is scaled by the signal strength parameter μ (which is the
1943 parameter of interest of the fit), while the interference yield is scaled by $\sqrt{\mu}$.

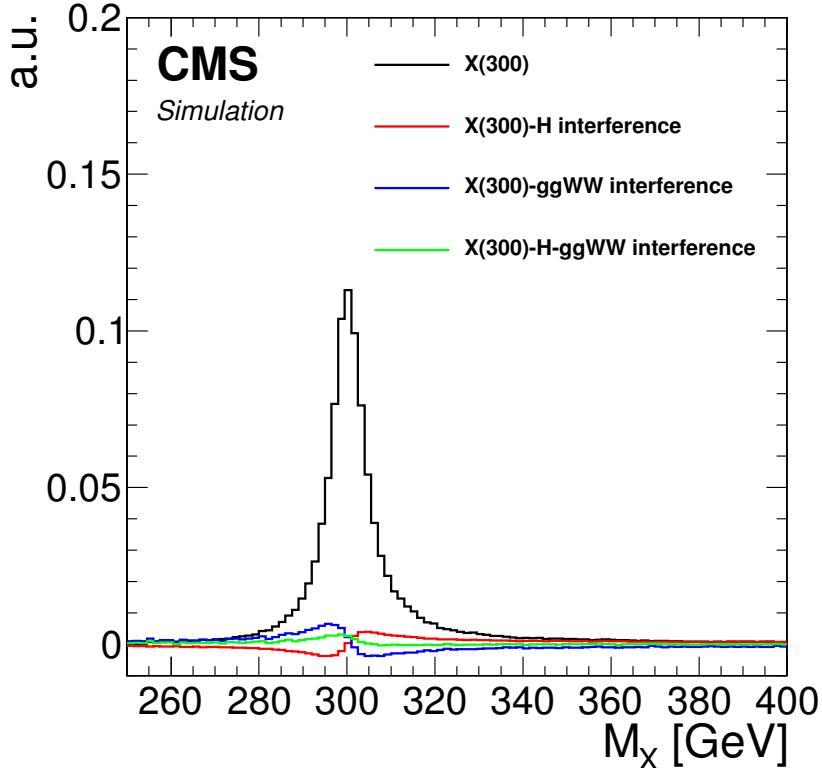


Figure 6.3.: Distribution of the M_X variable for a resonance mass of 300 GeV, showing the various interference terms after the WW baseline selections.

1944 6.4. Background estimation

1945 The background processes affecting the analysis phase space are the same as the ones
 1946 contributing to the SM Higgs search described in Sec. 5.4. The techniques used for the
 1947 background estimation are the same as well.

1948 The most relevant difference is the addition of the 2 jets category. The WW, top
 1949 quark and DY background normalizations are estimated in this category using data driven
 1950 techniques, similarly to the other jet bins.

1951 Given the slightly different WW baseline selection with respect to the SM Higgs search,
 1952 also the control regions for the top quark and DY backgrounds estimation change, while the
 1953 WW background normalization is estimated from data in the three signal regions separately,
 1954 owing to the different m_T^j shapes for signal and background.

1955 For the estimation of the top quark background, three control regions enriched in b-jets
 1956 are defined by selecting events that pass the WW baseline selections and applying a b
 1957 tagging requirement which depends on the jet category as follows:

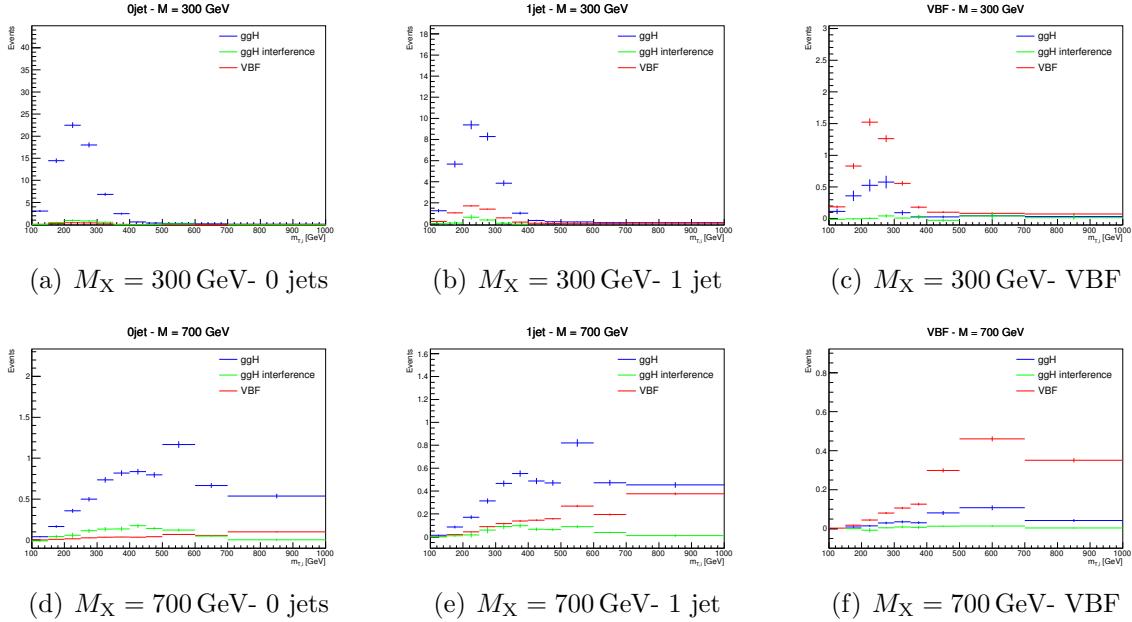


Figure 6.4.: Distributions of the m_T^i variable for $M_X = 300$ and 700 GeV , showing the signal (both the ggH and VBF processes) and the interference contributions in the three jet categories.

• 0 jets category: at least one b-tagged jet with $20 < p_T < 30 \text{ GeV}$ is required;

• 1 jet category: exactly one b-tagged jet with p_T above 30 GeV is required;

• 2 jets category: at least one b-tagged jet with p_T above 30 GeV is required.

Distributions of the m_T^i variable in the 0 jets, 1 jet and 2 jets top quark enriched control regions after applying the data driven estimation are shown in Fig. 6.5.

The jet induced background, also labelled as “non-prompt” background, so as to highlight that these events do not contain prompt leptons, is estimated using the same fake rate method described in 4.4.3. A crosscheck is performed selecting events passing the WW baseline selection but with an $e\mu$ pair with same charge. The m_T^i distributions for this phase space are shown in Fig. 6.6 for the three jet categories separately, showing agreement between data and simulation within the uncertainties.

Due to the cuts on the leptons p_T and on $m_{\ell\ell}$ in the WW baseline requirements, the contribution of the DY background decaying to a pair of τ leptons is very small in the signal regions, especially in the VBF phase space. The normalization of this background is estimated from a control region in data, defined in the same way as explained in 5.4.4, for the 0 and 1 jet categories. In the VBF category, given the very small number of expected events, the normalization of this background is taken from simulation.

Other minor background processes are estimated as described in 5.4.5.

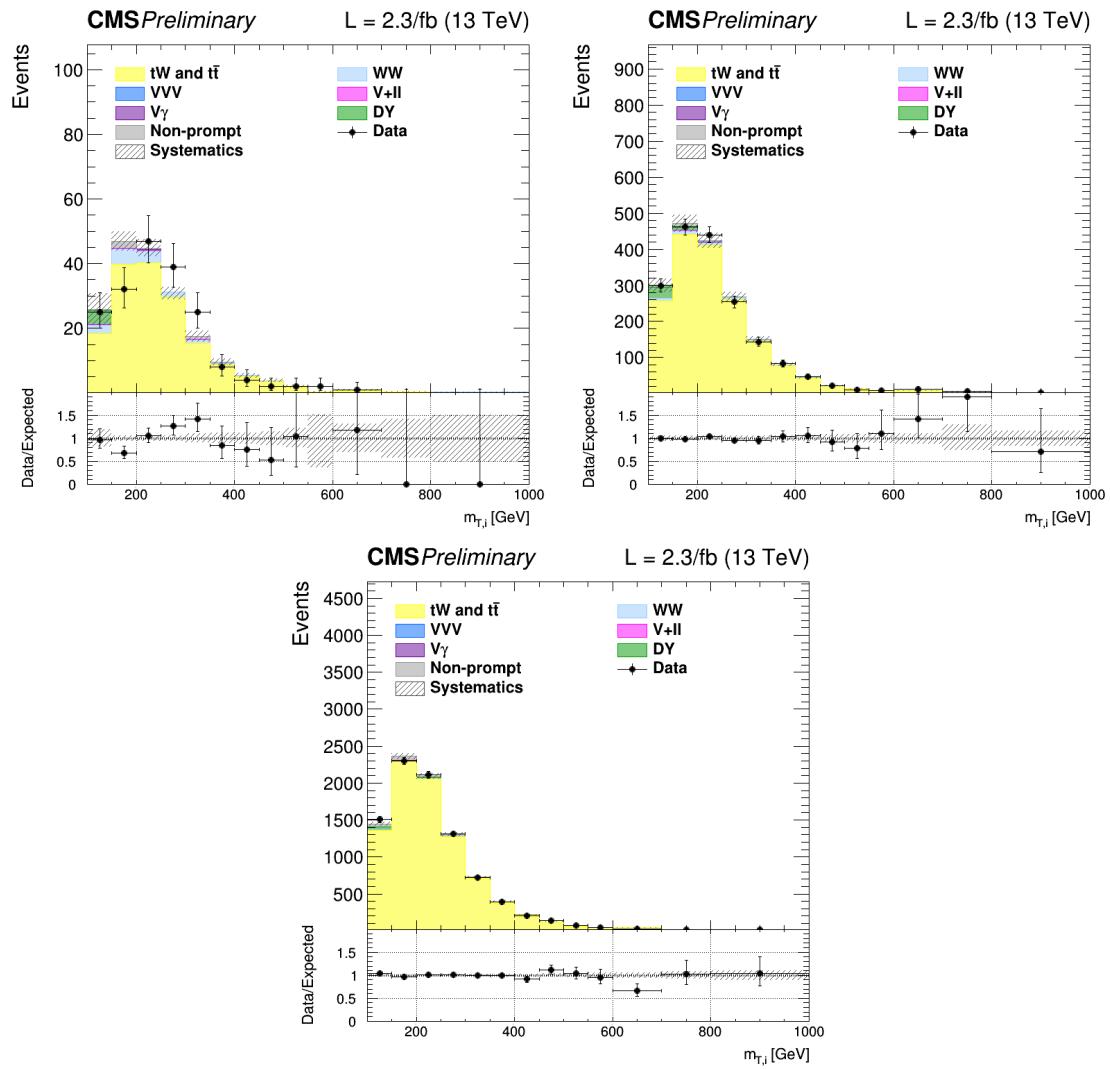


Figure 6.5.: Distributions of $m_{T,j}^j$ for events with 0 jet (top left), 1 jet (top right) and 2 jets (bottom) in top enriched control region. Scale factors estimated from data are not applied in the plots.

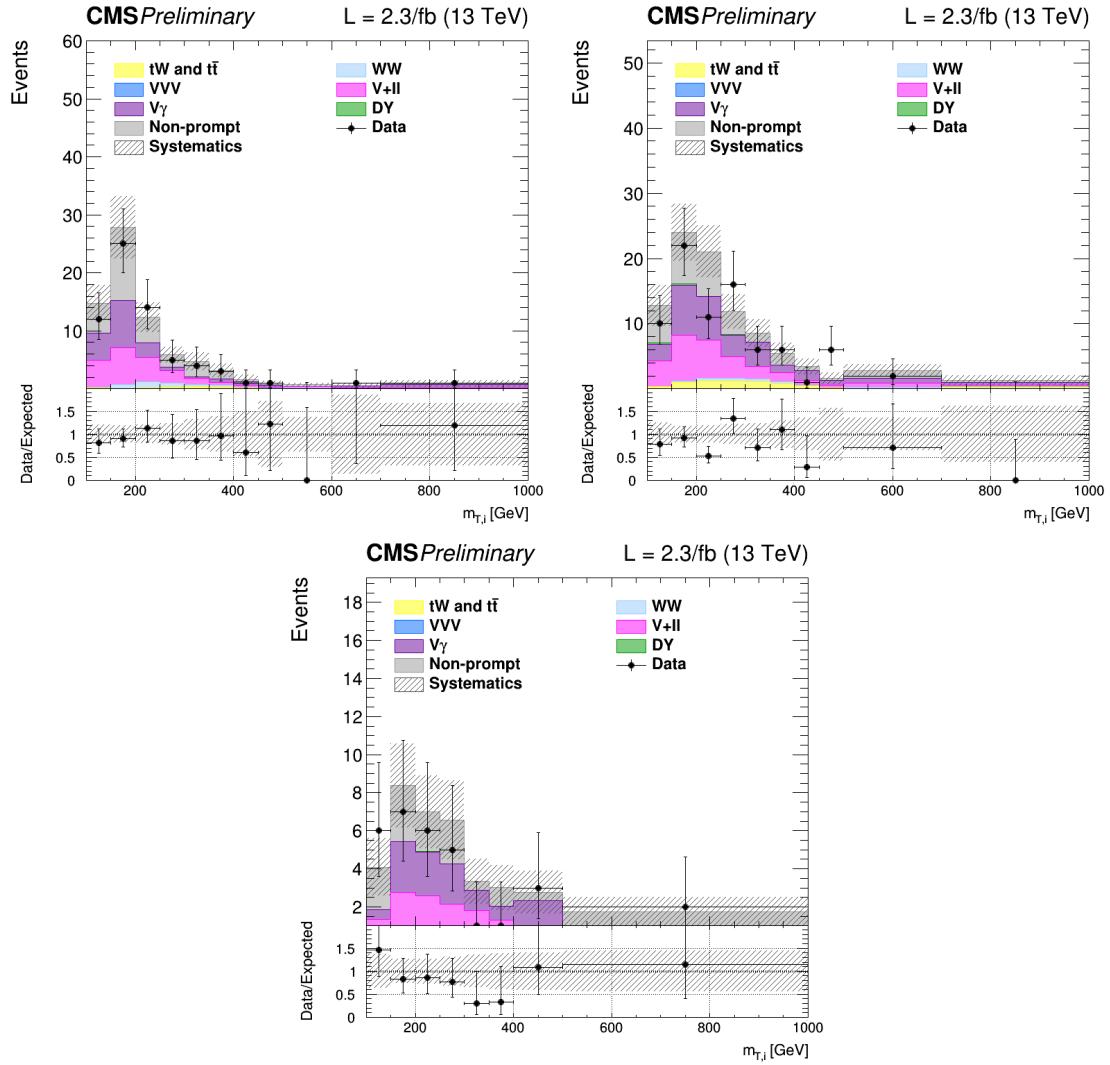


Figure 6.6.: Distributions of $m_{T,i}^i$ for events with 0 jet (top left), 1 jet (top right) and 2 jets (bottom) in the same-charge dilepton control region. The last bin of the histograms includes overflows.

1976 6.5. Systematic uncertainties

1977 The systematic uncertainties affecting this analysis are the same discussed in Sec. 5.5. The
1978 differences with respect to the SM Higgs boson search are described below.

1979 The PDF and α_s uncertainties on the signal cross sections are taken from the compu-
1980 tations performed by the LHC cross section working group [92], and are included for all
1981 the mass points. The value of these uncertainties depends on the resonance mass and vary
1982 from 3 and 5% for ggH and from 2 and 3% for VBF production modes. The PDFs and α_s
1983 uncertainties on the signal selection are evaluated for every resonance mass and are found
1984 to be less than 1% for both ggH and VBF.

1985 The theoretical uncertainties in the signal yields due to the jet categorization are
1986 evaluated for all the ggH signals following the prescription described in Sec. 4.5.3.

1987 An additional uncertainty on the modelling of the top pair background is derived from
1988 the observed discrepancy between data and POWHEG V2 plus PYTHIA 8.1 simulation
1989 on the top quark p_T spectrum [102], which is particularly important in the tail of the m_T^i
1990 distribution. Another uncertainty affecting the m_T^i tail for the top quark background is the
1991 parton shower uncertainty. This is evaluated comparing the generator level m_T^i distributions
1992 corresponding to two different simulations of the $t\bar{t}$ process: one obtained using PYTHIA
1993 8.1 for the showering and hadronization of the simulated events, and the other using
1994 HERWIG++. The difference between the two is used to extract a shape uncertainty, which
1995 is less than 1% for low m_T^i values and reaches about 6% in the m_T^i tail.

1996 6.6. Signal extraction and limit setting

1997 The signal yield, including both the ggH and VBF production modes, is extracted performing
1998 a combined fit of the three categories to the m_T^i simulation templates for backgrounds and
1999 signal, and is repeated for each resonance mass hypothesis. Moreover, fixed the mass of
2000 the resonance, the fit is performed again for the various hypotheses of the resonance decay
2001 width. A single signal strength μ is extracted from the fit, which multiplies both the ggH
2002 and VBF contributions. In other words it is assumed that the ratio of the two production
2003 mechanism stays the same as the one predicted by the SM².

2004 The background yields expected from simulation corresponding to the three jet categories
2005 and after the analysis event selection are shown in Table 6.1. The signal yields corresponding
2006 to a selection of mass points and assuming $\Gamma' = \Gamma_{\text{SM}}$ are shown in Table 6.2.

²This is an approximation which limits the amount of models that can be tested with the provided results.
A future development of this analysis will also include the cases for which different ggH and VBF relative contributions are expected.

Table 6.1.: Expected yields estimated from simulation (except for the non-prompt contribution which is estimated using data) for each background process in the three analysis categories, after the analysis event selection. The uncertainties are shown for the processes estimated from simulation.

Background process	0 jets	1 jet	VBF
qq \rightarrow WW	501.93 \pm 0.00 (0%)	198.72 \pm 0.00 (0%)	4.54 \pm 0.00 (0%)
gg \rightarrow WW	37.28 \pm 5.77 (15%)	19.63 \pm 3.04 (15%)	1.05 \pm 0.16 (15%)
Top quark	188.75 \pm 0.00 (0%)	330.05 \pm 0.00 (0%)	25.06 \pm 0.00 (0%)
DY	33.24 \pm 0.00 (0%)	12.99 \pm 0.00 (0%)	0.28 \pm 0.00 (0%)
Non-prompt	64.21 \pm 19.26 (30%)	31.69 \pm 9.51 (30%)	2.10 \pm 0.63 (30%)
V γ	26.62 \pm 0.72 (3%)	14.18 \pm 0.38 (3%)	0.64 \pm 0.02 (3%)
V γ^*	4.44 \pm 1.12 (25%)	3.39 \pm 0.85 (25%)	0.14 \pm 0.04 (25%)
VZ	13.51 \pm 0.76 (6%)	11.67 \pm 0.66 (6%)	0.28 \pm 0.02 (6%)
VVV	0.01 \pm 0.00 (3%)	0.02 \pm 0.00 (3%)	0.00 \pm 0.00 (3%)
SM H \rightarrow WW	6.04 \pm 0.40 (7%)	3.10 \pm 0.11 (5%)	0.34 \pm 0.02 (7%)
SM H \rightarrow $\tau\tau$	0.50 \pm 0.05 (9%)	0.43 \pm 0.04 (9%)	0.04 \pm 0.00 (9%)
Total background	876.5	625.9	34.5

2007 The strategy for computing the exclusion limits is based on the modified frequentist
 2008 approach, also referred to as CL_s, as described in [91]. The first step is to construct the
 2009 likelihood function $\mathcal{L}(\text{data}|\mu, \theta)$:

$$\mathcal{L}(\text{data}|\mu, \theta) = \text{Poisson}(\text{data}|\mu \cdot s(\theta) + b(\theta)) \cdot p(\tilde{\theta}|\theta) , \quad (6.3)$$

2010 where data represents the experimental observation, s and b are the expected signal and
 2011 background yields respectively and θ is the full set of nuisance parameters constrained by
 2012 the prior distribution functions $p(\tilde{\theta}|\theta)$. The default values for the nuisance parameters are
 2013 labelled as $\tilde{\theta}$.

2014 For a binned shape analysis, $\text{Poisson}(\text{data}|\mu \cdot s + b)$ is the product of the Poisson
 2015 probabilities to observe n_i events in bin i:

$$\prod_i \frac{(\mu \cdot s_i + b_i)^{n_i}}{n_i!} e^{-\mu \cdot s_i - b_i} . \quad (6.4)$$

2016 In order to test the compatibility of the data with the signal plus background (or the
 2017 background only) hypothesis, the test statistic \tilde{q}_μ is constructed based on the profile

Table 6.2.: Expected signal yields for the ggH and VBF production modes estimated from simulation after the analysis event selection for different mass hypothesis assuming $\Gamma' = \Gamma_{\text{SM}}$ in the three analysis categories. The errors correspond to the theoretical uncertainties in the signal estimation.

Mass (GeV)	0 jets	1 jet	VBF
ggH signal yields			
200	90.21 ± 6.67 (7%)	37.47 ± 1.81 (5%)	1.25 ± 0.26 (21%)
400	66.35 ± 4.90 (7%)	32.65 ± 1.57 (5%)	2.04 ± 0.42 (21%)
600	13.86 ± 1.05 (8%)	8.56 ± 0.44 (5%)	0.68 ± 0.14 (21%)
800	3.20 ± 0.25 (8%)	2.32 ± 0.13 (6%)	0.22 ± 0.05 (21%)
1000	0.88 ± 0.07 (8%)	0.70 ± 0.04 (6%)	0.07 ± 0.02 (21%)
VBF signal yields			
200	1.54 ± 0.06 (4%)	6.18 ± 0.25 (4%)	5.05 ± 0.20 (4%)
400	0.91 ± 0.04 (4%)	3.42 ± 0.14 (4%)	3.19 ± 0.13 (4%)
600	0.50 ± 0.02 (4%)	1.95 ± 0.08 (4%)	1.88 ± 0.08 (4%)
800	0.33 ± 0.01 (4%)	1.21 ± 0.05 (4%)	1.16 ± 0.05 (4%)
1000	0.22 ± 0.01 (4%)	0.79 ± 0.03 (4%)	0.69 ± 0.03 (4%)

2018 likelihood ratio:

$$\tilde{q}_\mu = -2 \ln \frac{\mathcal{L}(\text{data}|\mu, \hat{\theta}_\mu)}{\mathcal{L}(\text{data}|\hat{\mu}, \hat{\theta})} \quad \text{with } 0 \leq \hat{\mu} \leq \mu \quad , \quad (6.5)$$

2019 where $\hat{\theta}_\mu$ refers to the conditional maximum likelihood estimators of θ , given the signal
2020 strength μ . The parameter estimators $\hat{\mu}$ and $\hat{\theta}$ correspond to the global maximum of the
2021 likelihood. The $0 \leq \hat{\mu}$ constraint is imposed to have a positive signal yield, e.g. background
2022 underfluctuations are forbidden, while $\hat{\mu} \leq \mu$ is imposed to have a one-sided confidence
2023 interval. The observed test statistic for the signal strength μ under test is referred to as
2024 $\tilde{q}_\mu^{\text{obs}}$. The values of the nuisance parameters obtained maximising the likelihood function are
2025 labelled as $\hat{\theta}_0^{\text{obs}}$ and $\hat{\theta}_\mu^{\text{obs}}$ for the background only and signal plus background hypotheses,
2026 respectively. The pdf of the test statistic is constructed by generating toy MC pseudo-data
2027 for both the background only and signal plus background hypotheses, i.e. $f(\tilde{q}_\mu|\mu, \hat{\theta}_\mu^{\text{obs}})$ and
2028 $f(\tilde{q}_\mu|0, \hat{\theta}_0^{\text{obs}})$. These distributions can be used to define two p-values corresponding to the

2029 two hypotheses, p_μ and p_b :

$$p_\mu = P(\tilde{q}_\mu \geq \tilde{q}_\mu^{\text{obs}} | \text{signal + background}) = \int_{\tilde{q}_\mu^{\text{obs}}}^{\infty} f(\tilde{q}_\mu | \mu, \hat{\theta}_\mu^{\text{obs}}) d\tilde{q}_\mu , \quad (6.6)$$

2030

$$1 - p_b = (\tilde{q}_\mu \geq \tilde{q}_\mu^{\text{obs}} | \text{background only}) = \int_{\tilde{q}_0^{\text{obs}}}^{\infty} f(\tilde{q}_\mu | 0, \hat{\theta}_0^{\text{obs}}) d\tilde{q}_\mu . \quad (6.7)$$

2031 According to these definitions, p_μ and p_b can be identified with CL_{s+b} and $1 - \text{CL}_b$, 2032 respectively. The $\text{CL}_s(\mu)$ is calculated using the following ratio:

$$\text{CL}_s(\mu) = \frac{\text{CL}_{s+b}}{\text{CL}_b} = \frac{p_\mu}{1 - p_b} . \quad (6.8)$$

2033 If, for a given signal strength μ , $\text{CL}_s \leq \alpha$, then the hypothesis is excluded with a $(1 - \alpha)$ 2034 confidence level (CL). For instance, if one wants to quote the upper limit on μ with a 95% 2035 CL, the signal strength has to be adjusted until $\text{CL}_s = 0.05$.

2036 The expected median upper limit, as well as the $\pm 1\sigma$ (68% CL) and $\pm 2\sigma$ (95% CL) 2037 bands, are determined generating a large amount of pseudo-data in the background only 2038 hypothesis and calculating CL_s and the 95% CL upper limit for each of them, as if they 2039 were real data. Then the cumulative distribution of the 95% CL upper limits is built and 2040 the median expected value is identified as the value at which the cumulative distribution 2041 crosses the 50% quantile. The $\pm 1\sigma$ ($\pm 2\sigma$) band is defined by the values at which the 2042 cumulative distribution crosses the 16% (2.5%) and 84% (97.5%) quantiles.

2043 In order to assess the sensitivity of the analysis, the expected upper exclusion limits at 2044 95% CL on the signal strength are shown in Fig. 6.7 for the three jet categories separately. 2045 For a given mass of the resonance, the limits are derived assuming a signal decay width 2046 $\Gamma' = \Gamma_{\text{SM}}$ and a cross section equal to the one expected from a SM Higgs boson at that 2047 mass. The other decay width hypothesis have also been tested, showing a very similar 2048 expected exclusion limit, suggesting that this analysis is not strongly sensitive to variations 2049 of the resonance decay width.

2050 The 0 jets category is the most sensitive especially in the low mass region, while for 2051 very large masses of the resonance the 1 jet and VBF categories start being important. 2052 This is explained mainly by the fact that the VBF contribution increases, with respect to 2053 ggH, as the mass increases. The expected exclusion limit on the signal strength after the 2054 combination of the three categories is shown in Fig. 6.8. Comparing the limits in the single 2055 categories with the combination of the three it is evident how the higher jet multiplicity 2056 categories help in improving the results for large values of M_X .

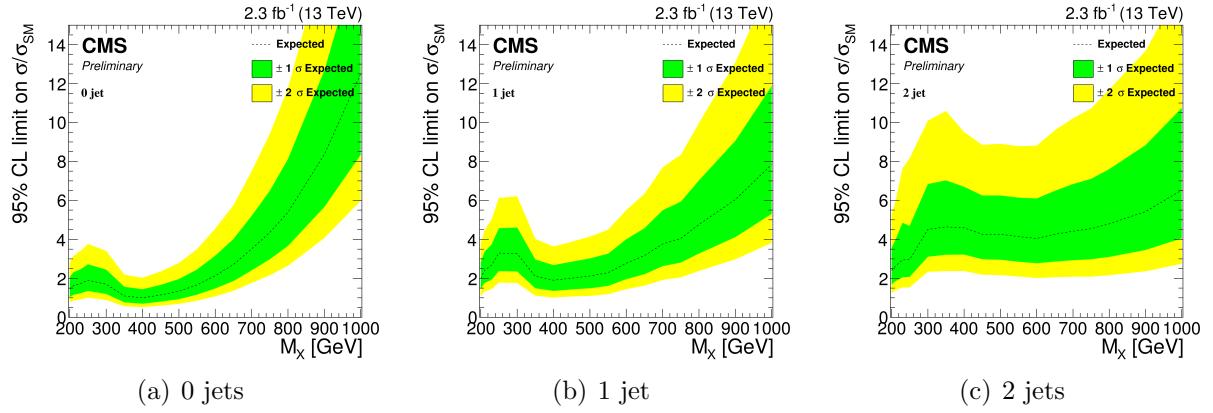


Figure 6.7.: Expected exclusion upper limits at 95% CL on the signal strength in the three categories, as a function of the resonance mass. The dashed line corresponds to median upper limit, while the green and yellow regions represent the $\pm 1\sigma$ and $\pm 2\sigma$ uncertainty bands, respectively. Limits are derived assuming a SM Higgs boson cross section and decay width for each mass point.

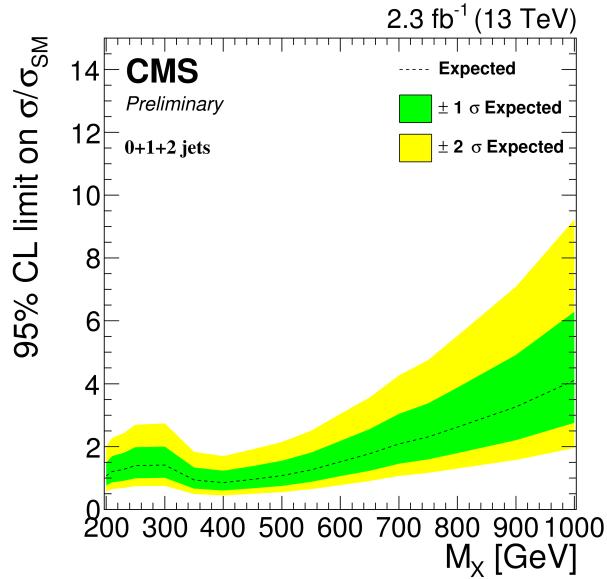


Figure 6.8.: Expected exclusion upper limit at 95% CL on the signal strength for the combination of the three categories, as a function of the resonance mass. The dashed line corresponds to median upper limit, while the green and yellow regions represent the $\pm 1\sigma$ and $\pm 2\sigma$ uncertainty bands, respectively. The limit is derived assuming a SM Higgs boson cross section and decay width for each mass point.

2057 6.7. Results

2058 The m_T^i distributions for the signal region after the full analysis selection are shown
 2059 in Fig. 6.9 for the three jet categories. Two different signal hypotheses corresponding
 2060 to $M_X = 400 \text{ GeV}$ and $M_X = 800 \text{ GeV}$ are shown superimposed on the background for
 2061 comparison.

2062 For every mass point from 200 GeV up to 1 TeV the observed p-value and the 95% CL
 2063 upper exclusion limit are calculated for five hypothesis of the signal width. The observed
 2064 p-value as a function of the resonance mass for the combination of the three jet categories
 2065 is shown in Table 6.3.

Table 6.3.: Observed p-value and corresponding significance (set to 0 in case of underfluctuations of the observed number of events) for the combination of the three jet categories for different resonance masses. Different values of the signal width are shown.

Mass [GeV]	$\Gamma = 0.09 \times \Gamma_{SM}$ p-value (signif.)	$\Gamma = 0.25 \times \Gamma_{SM}$ p-value (signif.)	$\Gamma = 0.49 \times \Gamma_{SM}$ p-value (signif.)	$\Gamma = \Gamma_{SM}$ p-value (signif.)
200	0.50 (0)	0.50 (0)	0.50 (0)	0.56 (0)
210	0.58 (0)	0.45 (0.1)	0.35 (0.4)	0.24 (0.7)
230	0.21 (0.8)	0.22 (0.8)	0.23 (0.7)	0.26 (0.6)
250	0.29 (0.5)	0.20 (0.8)	0.15 (1.0)	0.12 (1.2)
300	0.014 (2.2)	0.015 (2.2)	0.016 (2.1)	0.018 (2.1)
350	0.16 (1.0)	0.17 (1.0)	0.18 (0.9)	0.23 (0.7)
400	0.50 (0)	0.49 (0)	0.49 (0)	0.57 (0)
450	0.51 (0)	0.50 (0)	0.50 (0)	0.52 (0)
500	0.50 (0)	0.51 (0)	0.50 (0)	0.52 (0)
550	0.50 (0)	0.51 (0)	0.51 (0)	0.51 (0)
600	0.50 (0)	0.50 (0)	0.51 (0)	0.51 (0)
650	0.50 (0)	0.50 (0)	0.54 (0)	0.50 (0)
700	0.50 (0)	0.50 (0)	0.50 (0)	0.50 (0)
750	0.50 (0)	0.54 (0)	0.50 (0)	0.40 (0.3)
800	0.50 (0)	0.55 (0)	0.39 (0.3)	0.29 (0.6)
900	0.29 (0.6)	0.27 (0.6)	0.24 (0.7)	0.22 (0.8)
1000	0.18 (0.9)	0.18 (0.9)	0.18 (0.9)	0.18 (0.9)

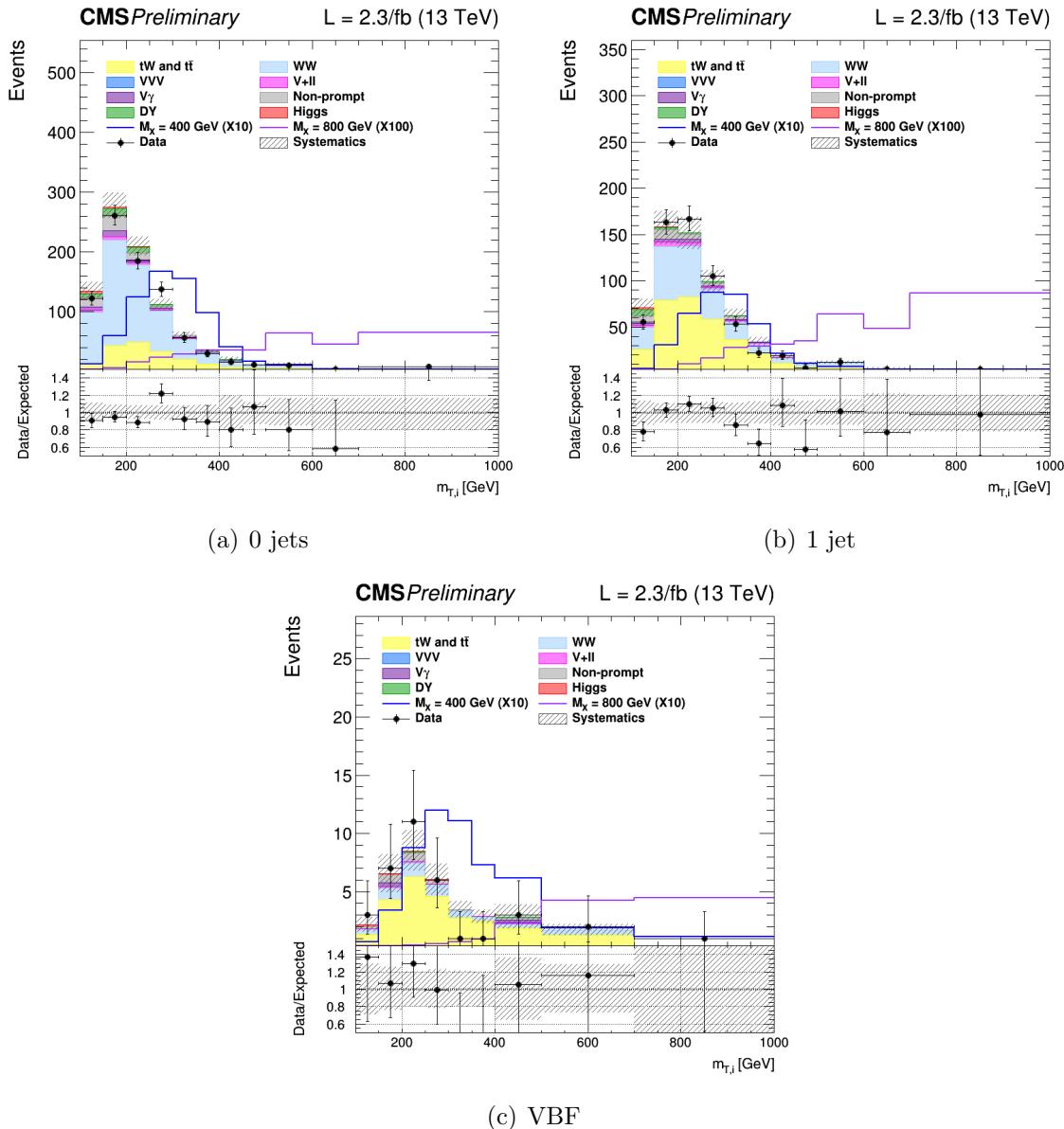


Figure 6.9.: Distributions of m_T^i in the signal region for the 0 jets, 1 jet and VBF categories. Background normalisations correspond to the pre-fit value. Signal contributions for two mass hypotheses, $M_X = 400$ GeV and $M_X = 800$ GeV, are shown superimposed on the background and scaled to facilitate the comparison.

In order to be independent on the particular model assumed for the signal cross section, the results are interpreted as exclusion limits on $\sigma \times \mathcal{B}$, where σ stands for the sum of the ggH and VBF cross sections, and \mathcal{B} represents the $X \rightarrow WW \rightarrow 2\ell 2\nu$ branching ratio including all lepton flavours. The expected and observed upper exclusion limits on $\sigma \times \mathcal{B}$ for $\Gamma' = \Gamma_{\text{SM}}$ are shown in Fig. 6.10.

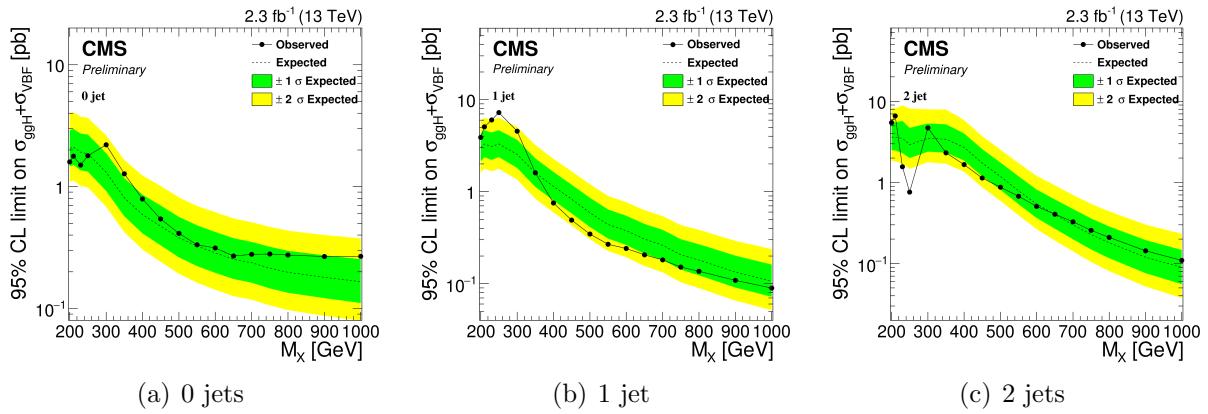


Figure 6.10.: Expected and observed exclusion upper limits at 95% CL on $\sigma \times \mathcal{B}$ in the three categories, as a function of the resonance mass. The dashed line corresponds to median upper limit, while the green and yellow regions represent the $\pm 1\sigma$ and $\pm 2\sigma$ uncertainty bands, respectively. The dotted line represents the observed limit. Limits are derived assuming $\Gamma' = \Gamma_{\text{SM}}$ for each mass point.

2071 A mild excess is observed in the 0 jets category and, more evident, in the 1 jet category
 2072 around 250-300 GeV. A deficit is instead in the VBF category around 250 GeV, which is
 2073 mainly due to an underfluctuation of the background. This effect can be understood looking
 2074 at the VBF shape in Fig. 6.9, where two adjacent data points, corresponding to the fifth
 2075 and sixth bins, clearly underfluctuate with respect to the background prediction, causing
 2076 the dip in the observed limit.

The exclusion limit resulting from the combination of the three categories is shown in Fig. 6.11, for the five Γ' hypotheses discussed before. From the combined exclusion limits no significant evidence of a deviation from the SM prediction is observed. The presence of a scalar resonance with $\sigma \times \mathcal{B}$ higher than the values reported in Fig. 6.11 is thus excluded with a 95% CL for masses ranging from 200 GeV up to 1 TeV.

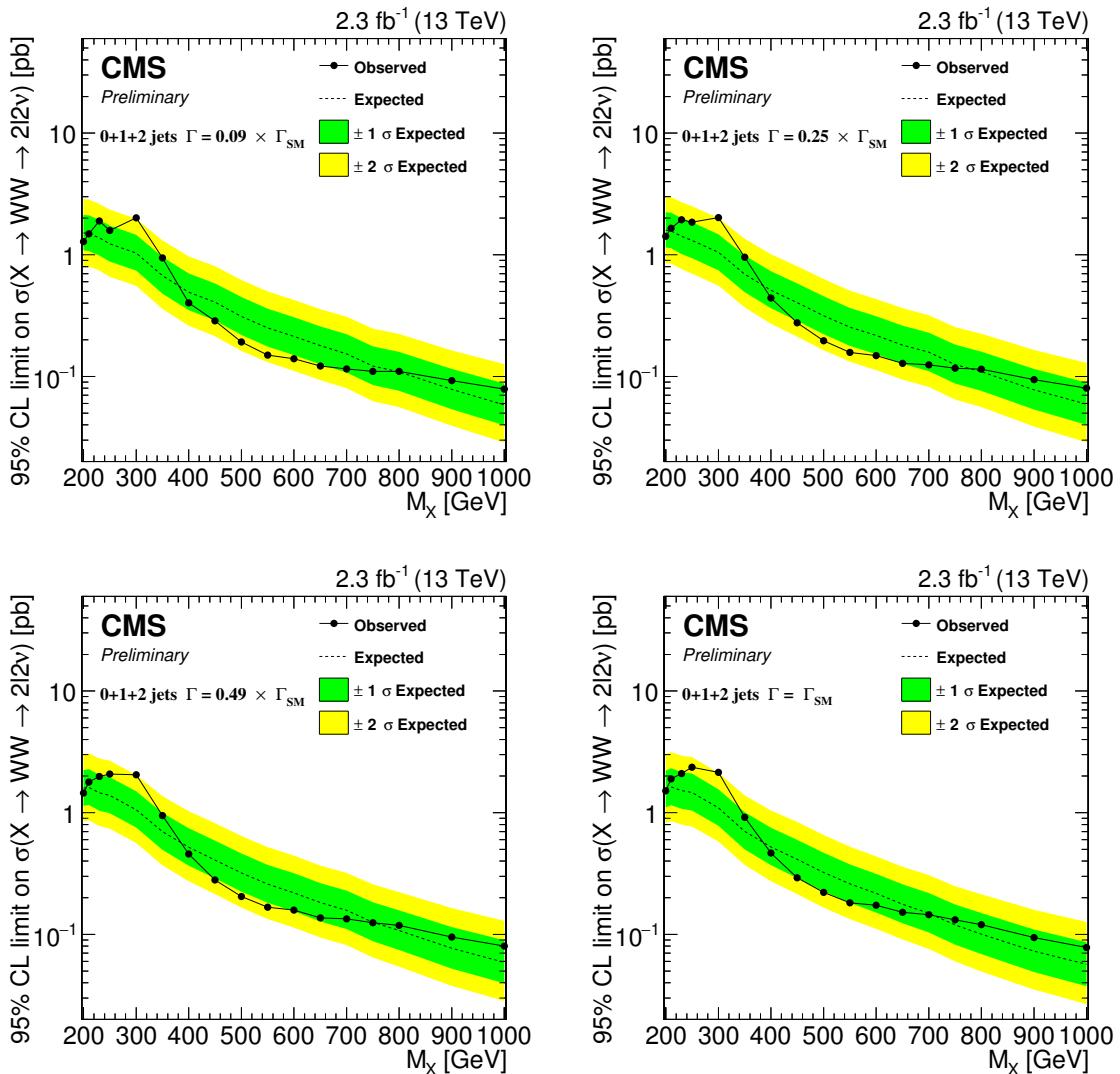


Figure 6.11.: Expected and observed exclusion limits at 95% CL on $\sigma \times \mathcal{B}$ for the combination of the three jet categories as a function of the resonance mass. The black dotted line corresponds to the observed value while the yellow and green bands represent the $\pm 1\sigma$ and $\pm 2\sigma$ uncertainties respectively. Limits are shown for four hypotheses of the resonance decay width.

Chapter 7.

Conclusions

2082

Appendix A.

²⁰⁸³ Fiducial region definition and ²⁰⁸⁴ optimization

²⁰⁸⁵ The fiducial region must be chosen in such a way to be as close as possible to the selections
²⁰⁸⁶ applied in the analysis, in order to reduce the model dependence in the extrapolation step.
²⁰⁸⁷ That means that for optimizing the fiducial volume definition, the efficiency has to be
²⁰⁸⁸ maximized. Another parameter entering the game is the number of fake events, in other
²⁰⁸⁹ words the number of reconstructed events which do not belong to the fiducial phase space.
²⁰⁹⁰ This parameter should instead be as small as possible. Even if we have to observe the trend
²⁰⁹¹ of these two quantities as a function of p_T^H , we can maximize the ratio between the overall
²⁰⁹² efficiency and the overall fake rate as a proxy for establishing the “goodness” of the fiducial
²⁰⁹³ region.

²⁰⁹⁴ Several different fiducial region definitions were tested and the results show that:

- ²⁰⁹⁵ • **of cut:** The fiducial region definition must include only the opposite flavor combination
²⁰⁹⁶ including one electron and one muon. If we include also the combinations involving
²⁰⁹⁷ τ 's the efficiency falls down.
- ²⁰⁹⁸ • **Lepton cut:** Since the resolution on lepton transverse momentum is good, there is no
²⁰⁹⁹ need to loosen the cuts related these variables, i.e. we can use the same cuts defined
²¹⁰⁰ in the analysis selection ($p_T^{\ell,1} > 20 \text{ GeV}$, $p_T^{\ell,2} > 10 \text{ GeV}$).
- ²¹⁰¹ • **Di-lepton p_T cut:** As stated in the previous point, there is no need to loosen this
²¹⁰² cut, so we kept the same value as the analysis selection, i.e. $p_T^{\ell\ell} > 30 \text{ GeV}$.
- ²¹⁰³ • **Di-lepton mass cut:** $m_{\ell\ell} > 12 \text{ GeV}$ as discussed before.
- ²¹⁰⁴ • **neutrino pair p_T cut:** Since the resolution on the measurement of the missing
²¹⁰⁵ transverse energy is poor, the neutrino pair cut should not be included in the definition
²¹⁰⁶ of the fiducial region, because it would increase the fake rate without increasing the
²¹⁰⁷ efficiency, thus resulting in a lower ratio between overall efficiency and fake rate.
- ²¹⁰⁸ • **m_T cut:** Also the m_T cut that we have in the analysis selection, i.e. $m_T > 60 \text{ GeV}$,
²¹⁰⁹ should be loosened or removed because it involves neutrinos and then increase the

fake rate. We decided eventually to keep this cut, loosening it to 50 GeV, because in addition to increase the number of fake events, it increases the efficiency as well.

The fake rate and the efficiency as a function of p_T^H after the optimization discussed before are shown in figure A.1. To obtain these plots the fiducial region was modified adding in sequence the various cuts and computing the efficiency and the fake rate each time. In that way we can asses the composition of those distributions.

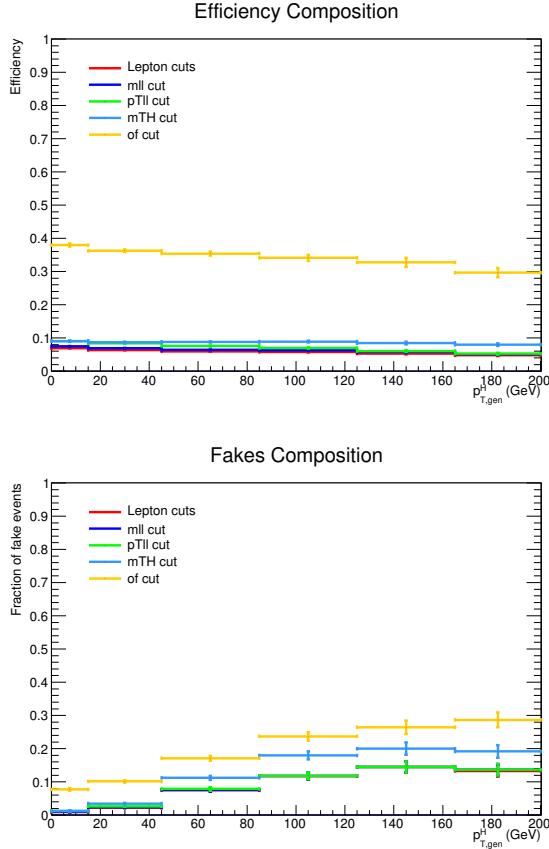


Figure A.1.: Efficiency and fake rate as a function on Higgs transverse momentum. The plots correspond to the optimized fiducial region definition and show the effect of adding each of the mentioned cuts in sequence.

The efficiency and fraction of fake events have been measured also as a function of the E_T^{miss} and m_T cuts in the fiducial region. Since these two variables are correlated, the results are reported as two-dimensional histograms. In Fig. A.2 are reported the efficiency and fraction of fake events for these two variables.

The criterion adopted to define the fiducial region is a tradeoff between having a large efficiency and a small fraction of fake events. Especially when looking at the low resolution variables, such as E_T^{miss} and m_T , a suitable figure of merit has to be chosen for the estimation of the best cuts. Several different figures of merit have been checked, such as ϵ/f , $\epsilon - f$

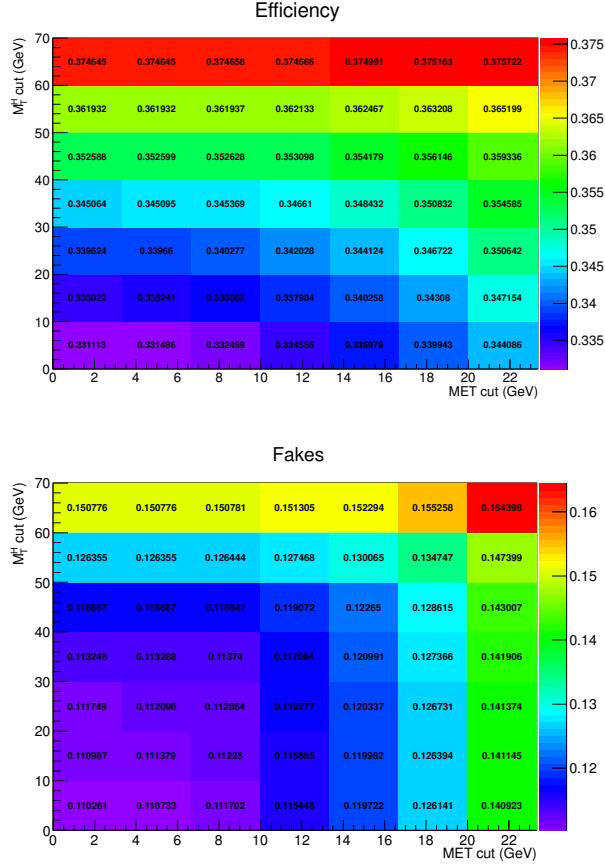


Figure A.2.: Efficiency and fake rate as a function of E_T^{miss} and m_T cuts in the fiducial region.

and $(1 - f)/\epsilon$. The results for these three different figures of merit are shown in Fig. A.3 as a function of the E_T^{miss} and m_T cuts in the fiducial region.

Following the same criterion, similar plots as above have been obtained for an alternative model, given by varying up the ggH/VBF ratio within the experimental uncertainties. The results, shown in Fig. A.4 and Fig. A.5, show a similar trend with respect to the model with nominal ggH/VBF ratio.

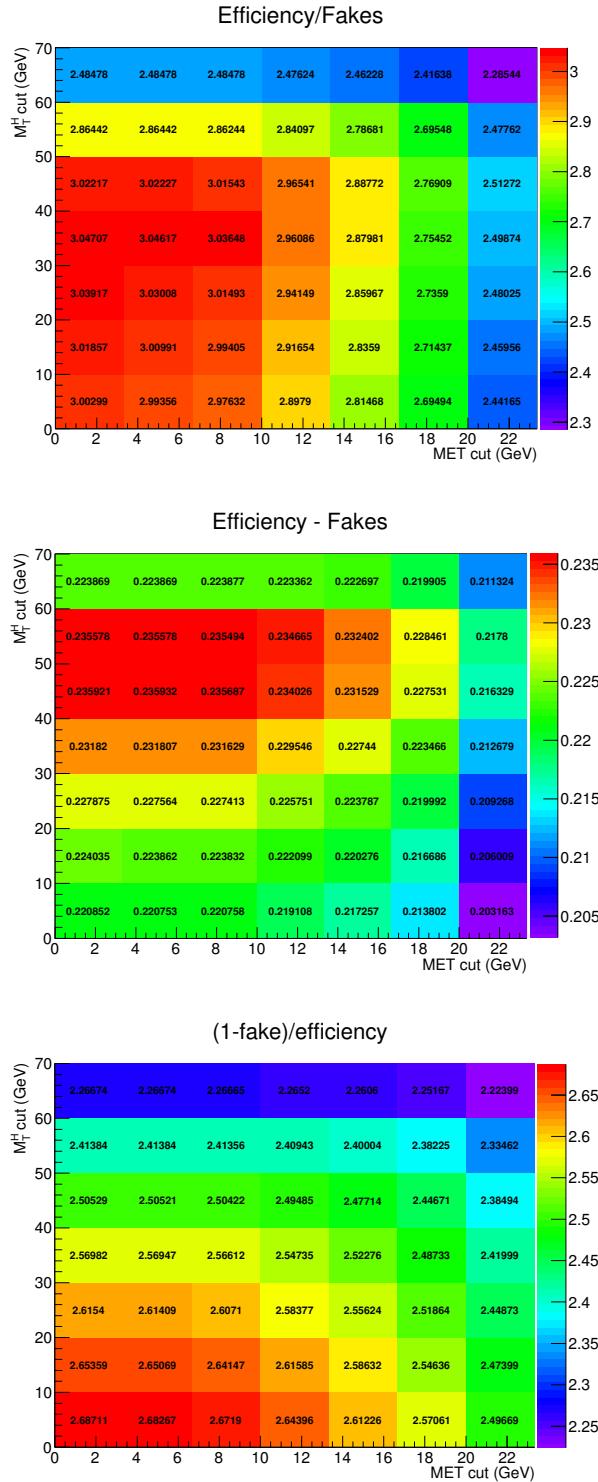


Figure A.3.: Different figures of merit as a function of E_T^{miss} and m_T cuts in the fiducial region.

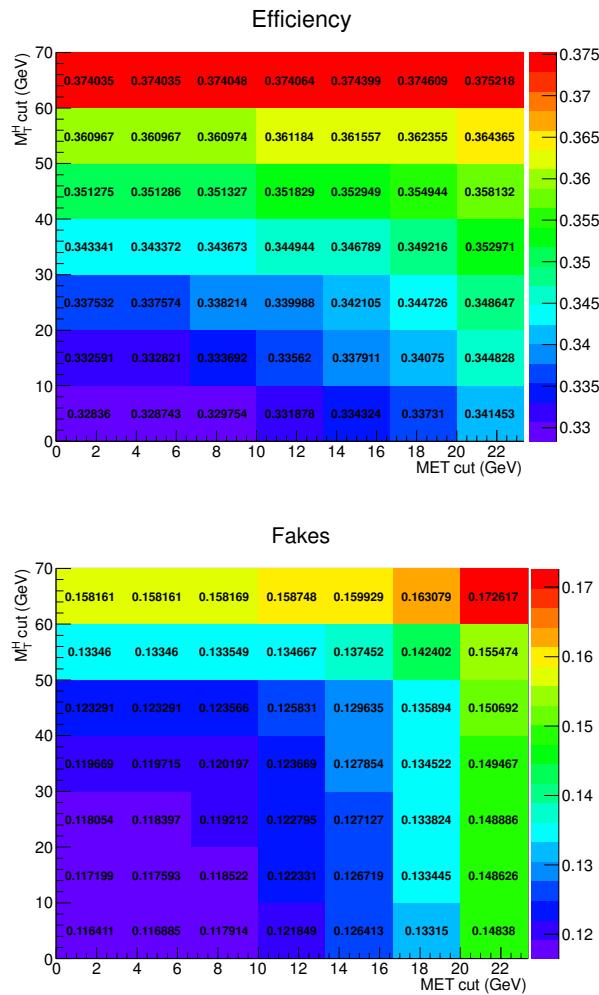


Figure A.4.: Efficiency and fake rate as a function of E_T^{miss} and m_T cuts in the fiducial region, for the alternative model with an up variation of the ggH/VBF ratio.

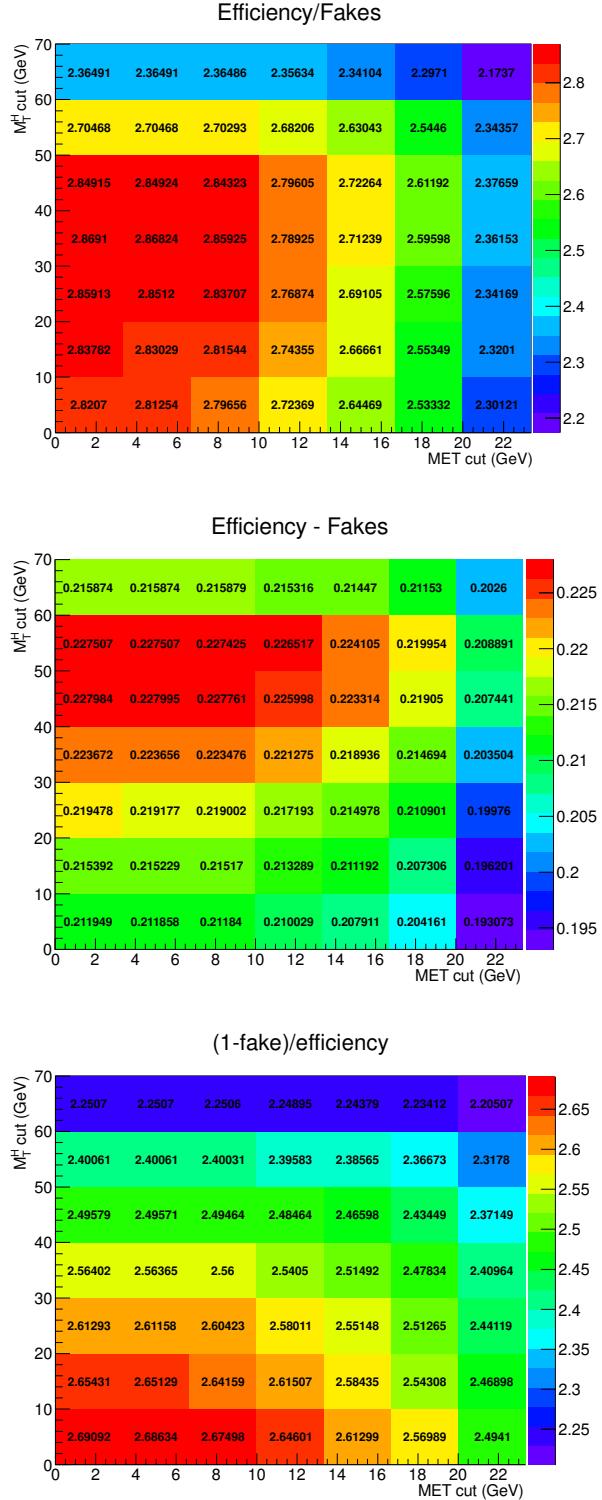


Figure A.5.: Different figures of merit as a function of E_T^{miss} and m_T cuts in the fiducial region, for the alternative model with an up variation of the ggH/VBF ratio.

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