

**Example 1 Pendulum**

$\phi(t)$  = angular displacement.

$\phi(0) = \theta$  = initial angular displacement.

Let  $t_p$  = period of the pendulum (time for 1 complete cycle)

**What factors determine  $t_p$  ( $[t_p] = T$ )?**

1.  $g$

$$\bullet [g] = LT^{-2}$$

2.  $m$  – mass

$$\bullet [m] = M$$

3.  $r$  – length

$$\bullet [r] = L$$

4.  $\theta$  = initial displacement

$$\bullet [\theta] = 1$$

5. air resistance (drag)

6. friction in hinge

Assume  $t_p = f(m, g, r, \theta)$  Dimension reduction:

$$\begin{aligned} [t_p] &= [m^a g^b r^c \theta^d] \\ [t_p] &= [m]^a [g]^b [r]^c [\theta]^d \\ M^0 L^0 T^1 &= M^a (LT^2)^b L^c (M^0 L^0 T^0)^d \\ &= M^a L^b T^{2b} L^c M^{0d} L^{0d} T^{0d} \\ &= M^a L^b T^{2b} L^c (1) \\ &= M^a L^{b+c} T^{2b} \end{aligned}$$

$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 1 & 0 \\ 0 & -2 & 0 & 0 \end{bmatrix} \begin{bmatrix} a \\ b \\ c \\ d \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}$$

$$\begin{aligned}
\left[ \begin{array}{cccc|c} 1 & 0 & 0 & 0 & 0 \\ 0 & 1 & 1 & 0 & 0 \\ 0 & -2 & 0 & 0 & 1 \end{array} \right] &= (r_2 \leftarrow 2r_2 + r_3) \left[ \begin{array}{cccc|c} 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 2 & 0 & 1 \\ 0 & -2 & 0 & 0 & 1 \end{array} \right] \\
&= (r_2 \leftrightarrow r_3) \left[ \begin{array}{cccc|c} 1 & 0 & 0 & 0 & 0 \\ 0 & -2 & 0 & 0 & 1 \\ 0 & 0 & 2 & 0 & 1 \end{array} \right] \\
&= \left( r_2 \leftarrow -\frac{r_2}{2} \right) \left[ \begin{array}{cccc|c} 1 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & -\frac{1}{2} \\ 0 & 0 & 2 & 0 & 1 \end{array} \right] \\
&= \left( r_3 \leftarrow \frac{r_3}{2} \right) \left[ \begin{array}{cccc|c} 1 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & -\frac{1}{2} \\ 0 & 0 & 1 & 0 & \frac{1}{2} \end{array} \right]
\end{aligned}$$

$$a = 0 \quad b = -\frac{1}{2} \quad c = \frac{1}{2} \quad d = \text{arbitrary constant}$$

$$\begin{aligned}
\begin{bmatrix} a \\ b \\ c \\ d \end{bmatrix} &= \begin{bmatrix} 0 \\ -\frac{1}{2} \\ \frac{1}{2} \\ d \end{bmatrix} \\
&= \begin{bmatrix} 0 \\ -\frac{1}{2} \\ \frac{1}{2} \\ 0 \end{bmatrix} + d \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix} \\
&= g^{-\frac{1}{2}} r^{\frac{1}{2}} + \theta^1 \\
&= \sqrt{g^{-1}r} + \theta \\
&= \sqrt{\frac{r}{g}} + \theta
\end{aligned}$$

Note the dimensions:

$$\begin{aligned}
\left[ \sqrt{\frac{r}{g}} \right] &= \sqrt{\frac{[r]}{[g]}} \\
&= \sqrt{\frac{L}{LT^{-2}}} \\
&= \sqrt{\frac{1}{T^{-2}}} \\
&= \sqrt{T^2} \\
&= T \\
[\theta] &= 1 \\
\text{nullspace} &= \left\{ \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix} \right\} \Rightarrow \text{non-dimensional quantity!}
\end{aligned}$$

$$t_p = \alpha \sqrt{\frac{r}{g}} \theta^d$$

where  $\alpha$  is a non-dimensional constant ( $[\alpha] = 1$ )

$$t_p = \sqrt{\frac{r}{g}} [\alpha, d_1 \theta^d + d_2 \theta^d + d_3 \theta^d + \dots]$$

where all the terms within the brackets can be defined as  $h(\theta)$  (and the  $\alpha$ 's are hidden in  $h(\theta)$  as well). So:

$$t_p = \sqrt{\frac{r}{g}} h(\theta)$$

where  $h(\theta)$  can be found using data. The completely dimensionless form would be:

$$\frac{t_p}{\sqrt{\frac{r}{g}}} = h(\theta)$$

## Example 2 Pendulum ODE

Derivation of ODE describing motion.

Newton's 2nd Law - rotational form

$\tau$  = torque (force)

$$\bullet [\tau] = \frac{ML}{T^2}$$

$I$  = moment of inertia (like mass; measures how difficult it is to change the rotational force)

$\alpha$  = angular acceleration (a)

$$\bullet \frac{1}{T^2}$$

$$\alpha = \frac{\tau}{I} \text{ or } \tau = I\alpha$$

What is the dimension for  $I$ ?

$$\begin{aligned} [\tau] &= [I\alpha] \\ &= [I][\alpha] \\ \frac{ML^2}{T^2} &= [I] \frac{1}{T^2} \\ ML^2 &= [I] \end{aligned}$$

In physics,  $I$  = mass  $\times$  square of the distance from axis =  $mr^2$

$$\begin{aligned}
\tau &= I \frac{d^2\psi}{dt^2} \\
&= \vec{r} \times \vec{f} \\
&= -mgr \sin(\psi) \\
I \frac{d^2\psi}{dt^2} &= -mgr \sin(\psi) \\
mr^2 \frac{d^2\psi}{dt^2} &= -mgr \sin(\psi) \\
r \frac{d^2\psi}{dt^2} &= -g \sin(\psi)
\end{aligned}$$

$$r \frac{d^2\psi}{dt^2} + g \sin(\psi) = 0$$

$$\frac{d^2\psi}{dt^2} + \frac{g}{r} \sin(\psi) = 0$$

Assuming that  $\theta$  is small,  $\sin(\psi) \approx \psi$ .

$$\frac{d^2\psi}{dt^2} + \frac{g}{r}\psi = 0, \quad \psi(0) = \theta, \quad \psi'(0) = 0$$

Since this is a 2nd order, homogenous function, we can assume  $\psi = e^{\lambda t}$

$$\begin{aligned}
0 &= \lambda^2 e^{\lambda t} + \frac{g}{r} e^{\lambda t} \\
&= e^{\lambda t} \left( \lambda^2 + \frac{g}{r} \right) \\
&= \lambda^2 + \frac{g}{r} \\
-\lambda^2 &= \frac{g}{r} \\
\lambda^2 &= -\frac{g}{r} \\
\lambda &= \pm \sqrt{-\frac{g}{r}} \\
&= \pm \sqrt{\frac{g}{r}} \times \sqrt{-1} \\
&= \pm \sqrt{\frac{g}{r}} i
\end{aligned}$$

The fundamental set is  $\left\{ e^{i\sqrt{\frac{g}{r}}t}, e^{-i\sqrt{\frac{g}{r}}t} \right\}$

Their real form (from ODE course) is

$$\psi = c_1 \cos\left(\sqrt{\frac{g}{r}}t\right) + c_2 \sin\left(\sqrt{\frac{g}{r}}t\right) \quad (1)$$

$c_1$  and  $c_2$  can be found using the initial conditions  $\psi(0) = \theta$  and  $\psi'(0) = 0$ .

$$\psi(0) = c_1 \cos \left( \sqrt{\frac{g}{r}}(0) \right) + c_2 \sin \left( \sqrt{\frac{g}{r}}(0) \right)$$

$$\theta = c_1 \cos(0) + c_2 \sin(0)$$

$$\theta = c_1(1) + c_2(0)$$

$$\theta = c_1$$

$$c_1 = \theta$$

$$\psi'(t) = -\sqrt{\frac{g}{r}}c_1 \sin \left( \sqrt{\frac{g}{r}}t \right) + \sqrt{\frac{g}{r}}c_2 \cos \left( \sqrt{\frac{g}{r}}t \right)$$

$$\psi'(0) = -\sqrt{\frac{g}{r}}c_1 \sin \left( \sqrt{\frac{g}{r}}(0) \right) + \sqrt{\frac{g}{r}}c_2 \cos \left( \sqrt{\frac{g}{r}}(0) \right)$$

$$0 = -\sqrt{\frac{g}{r}}c_1 \sin(0) + \sqrt{\frac{g}{r}}c_2 \cos(0)$$

$$0 = -\sqrt{\frac{g}{r}}c_1(0) + \sqrt{\frac{g}{r}}c_2(1)$$

$$0 = \sqrt{\frac{g}{r}}c_2$$

$$0 = c_2$$

$$c_2 = 0$$

$$\psi = c_1 \cos \left( \sqrt{\frac{g}{r}}t \right) + c_2 \sin \left( \sqrt{\frac{g}{r}}t \right)$$

$$\psi = \theta \cos \left( \sqrt{\frac{g}{r}}t \right) + 0 \sin \left( \sqrt{\frac{g}{r}}t \right)$$

$$\psi = \theta \cos \left( \sqrt{\frac{g}{r}}t \right)$$

### Buckingham's Pi Theorem

Let  $q$  = physical quantity (measurable).  $q$  depends on  $n$  parameters and variables  $p_1, \dots, p_n$

$$q = f(p_1, p_2, \dots, p_n)$$

Assume fundamental dimensions  $h, T, M$ .

$$[q] = L^{l_0} T^{t_0} M^{m_0}$$

$$[p_i] = L^{l_i} T^{t_i} M^{m_i}$$

$$[q] = [p_1]^{a_1} [p_2]^{a_2} \dots [p_n]^{a_n}$$

Objective: find  $a_j$  where  $j = 1, 2, \dots, n$

$$L^{l_0} T^{t_0} M^{m_0} = (L^{l_1} T^{t_1} M^{m_1})^{a_1} \times (L^{l_2} T^{t_2} M^{m_2})^{a_2} \times \cdots \times (L^{l_n} T^{t_n} M^{m_n})^{a_n}$$

Equate exponents!

$$L : l_0 = l_1 a_1 + l_2 a_2 + \cdots + l_n a_n$$

$$T : t_0 = t_1 a_1 + t_2 a_2 + \cdots + t_n a_n$$

$$M : m_0 = m_1 a_1 + m_2 a_2 + \cdots + m_n a_n$$

$$A = \begin{bmatrix} l_1 & l_2 & \cdots & l_n \\ t_1 & t_2 & \cdots & t_n \\ m_1 & m_2 & \cdots & m_n \end{bmatrix}_{3 \times n} \quad \vec{x} = \begin{bmatrix} a_1 \\ a_2 \\ \vdots \\ a_n \end{bmatrix}_{n \times 1} \quad \vec{b} = \begin{bmatrix} l_0 \\ t_0 \\ m_0 \end{bmatrix}_{3 \times 1}$$

In general,  $n > 3$  – undetermined  $\Rightarrow \infty$  set of solutions.

$A\vec{x} = \vec{b}$  has a solution of form

$$\vec{x} = \vec{x}_p + \gamma_1 \vec{x}_1 + \gamma_2 \vec{x}_2 + \cdots + \gamma_k \vec{x}_k$$

where  $\gamma$  is arbitrary,  $\vec{x}_p$  is the particular solution, and  $\{\vec{x}_j\}$  is the null space of  $A$ .

- $\vec{x}_j$  correspond to exponents that yield dimensionless quantities

**defines:**  $\Pi_j, j = 1, \dots, k$

- $\vec{x}_p$  exponents that yield a quantity with  $[q]$ , call it  $Q$

So

$$q = QF(\Pi_1, \Pi_1, \dots, \Pi_k)$$

$$\pi_0 = \frac{q}{Q} = QF(\Pi_1, \Pi_1, \dots, \Pi_k)$$

where the second equation is all in terms of dimensionless quantities

### Example 3 Drag on a sphere

Drag on a sphere within a fluid. The sphere has radius  $R$  moving at velocity  $v$  with a resistant drag force  $D_F$ . What factors play a role?

$R$  = radius  $[L]$

$v$  = velocity  $[LT^{-1}]$

$\rho$  = density  $[ML^{-3}]$

$\mu$  = dynamic viscosity  $[ML^{-1}T^{-1}]$

$$D_F = f(R, v, \rho, \mu)$$

$$[D_F] = [R]^a [v]^b [\rho]^c [\mu]^d$$

$$MLT^{-2} = L^a (LT^{-1})^b (ML^{-3})^c (ML^{-1}T^{-1})^d$$

$$MLT^{-2} = L^a \times L^b T^{-b} \times M^c L^{-3c} \times M^d L^{-d} T^{-d}$$

$$MLT^{-2} = M^{c+d} L^{a+b-3c-d} T^{-b-d}$$

$$\begin{bmatrix} 1 & 1 & -3 & -1 \\ 0 & -1 & 0 & -1 \\ 0 & 0 & 1 & 1 \end{bmatrix} \begin{bmatrix} a \\ b \\ c \\ d \end{bmatrix} = \begin{bmatrix} 1 \\ -2 \\ 1 \end{bmatrix}$$

$$\begin{aligned} \left[ \begin{array}{cccc|c} 1 & 1 & -3 & -1 & 1 \\ 0 & -1 & 0 & -1 & -2 \\ 0 & 0 & 1 & 1 & 1 \end{array} \right] &= (r_2 \leftarrow -r_2) \left[ \begin{array}{cccc|c} 1 & 1 & -3 & -1 & 1 \\ 0 & 1 & 0 & 1 & 2 \\ 0 & 0 & 1 & 1 & 1 \end{array} \right] \\ &= (r_1 \leftarrow r_1 - r_2) \left[ \begin{array}{cccc|c} 1 & 0 & -3 & -2 & -1 \\ 0 & 1 & 0 & 1 & 2 \\ 0 & 0 & 1 & 1 & 1 \end{array} \right] \\ &= (r_1 \leftarrow r_1 + 3r_3) \left[ \begin{array}{cccc|c} 1 & 0 & 0 & 1 & 2 \\ 0 & 1 & 0 & 1 & 2 \\ 0 & 0 & 1 & 1 & 1 \end{array} \right] \end{aligned}$$

$$\begin{aligned} a + d &= 2 & b + d &= 2 & c + d &= 1 \\ a &= 2 - d & b &= -d + 2 & c &= -d + 1 \end{aligned}$$

$$\begin{aligned} \begin{bmatrix} a \\ b \\ c \\ d \end{bmatrix} &= \begin{bmatrix} 2 - d \\ 2 - d \\ 1 - d \\ d \end{bmatrix} \\ &= \begin{bmatrix} 2 \\ 2 \\ 1 \\ 0 \end{bmatrix} + d \begin{bmatrix} -1 \\ -1 \\ -1 \\ 1 \end{bmatrix} \\ &= R^2 v^2 \rho^1 + R^{-1} v^{-1} \rho^{-1} \mu^1 \\ &= R^2 v^2 \rho + \frac{\mu}{Rv\rho} \\ &= \begin{bmatrix} 2 \\ 2 \\ 1 \\ 0 \end{bmatrix} + \textcolor{red}{-d} \begin{bmatrix} 1 \\ 1 \\ 1 \\ -1 \end{bmatrix} \\ &= R^2 v^2 \rho^1 + R^1 v^1 \rho^1 \mu^{-1} \\ &= R^2 v^2 \rho + \frac{Rv\rho}{\mu} \\ D_F &= R^2 v^2 \rho F \left( \frac{Rv\rho}{\mu} \right) \end{aligned}$$

where  $\frac{Rv\rho}{\mu}$  represents the Reynold's number  $Re$

### Reynold's number

$$Re = \frac{Rv\rho}{\mu} = \frac{\text{inertial}}{\text{viscous}} \quad (2)$$

When  $Re$  is small, there will be laminar (smooth) flow, when  $Re$  is large, the flow will be turbulent (having vortices).

### Example 4 T-Rex Top Speed

How fast can a Tyrannosaurus Rex (T-Rex) walk or run?

Develop a model!

Find  $v = \text{velocity}$   $[v] = \left[\frac{L}{T}\right]$

What variables are important?

$m$  mass,  $[M]$

$g$  gravity,  $[LT^{-2}]$

$h$  hip length,  $[L]$

$s$  stride length – distance between 2 steps,  $[s] = L$

$$v = f(s, h, g, m)$$

$$[v] = [s^a h^b g^c m^d]$$

$$LT^{-1} = L^a L^b (LT^{-2})^c M^d$$

$$LT^{-1} = L^{a+b} L^c T^{-2c} M^d$$

$$LT^{-1} M^0 = L^{a+b+c} T^{-2c} M^d$$

Equate exponents:

$$L: \quad 1 = a + b + c$$

$$T: \quad -1 = -2c$$

$$M: \quad 0 = d \text{ (independent of mass)}$$

$$\begin{bmatrix} 1 & 1 & 1 \\ 0 & 0 & -2 \end{bmatrix} \begin{bmatrix} a \\ c \end{bmatrix} = \begin{bmatrix} 1 \\ -1 \end{bmatrix}$$

$$\left[ \begin{array}{ccc|c} 1 & 1 & 1 & 1 \\ 0 & 0 & -2 & -1 \end{array} \right] = \left( r_2 \leftarrow -\frac{r_2}{2} \right) \left[ \begin{array}{ccc|c} 1 & 1 & 1 & 1 \\ 0 & 0 & 1 & \frac{1}{2} \end{array} \right]$$

$$a + b + c = 1 \quad b = \text{free } c = \frac{1}{2}$$

$$a + b + \frac{1}{2} = 1 \quad b = \text{free } c = \frac{1}{2}$$

$$a = -b + \frac{1}{2} \quad b = \text{free } c = \frac{1}{2}$$



$$\begin{aligned}
\begin{bmatrix} a \\ b \\ c \end{bmatrix} &= \begin{bmatrix} \frac{1}{2} - b \\ b \\ \frac{1}{2} \end{bmatrix} \\
&= \begin{bmatrix} \frac{1}{2} \\ 0 \\ \frac{1}{2} \end{bmatrix} + b \begin{bmatrix} -1 \\ 1 \\ 0 \end{bmatrix} \\
&= s^{\frac{1}{2}} g^{\frac{1}{2}} + s^{-1} h^1 \\
&= \sqrt{sg} + \frac{h}{s} \\
b \begin{bmatrix} -1 \\ 1 \\ 0 \end{bmatrix} &\Rightarrow (-1) \begin{bmatrix} -1 \\ 1 \\ 0 \end{bmatrix} \Rightarrow \begin{bmatrix} 1 \\ -1 \\ 0 \end{bmatrix} \Rightarrow s^1 h^{-1} \Rightarrow \frac{s}{h} \\
v &= \sqrt{gh} f\left(\frac{s}{h}\right)
\end{aligned}$$

### Froude number

As found in the previous example:

$$\frac{v}{\sqrt{gh}} = \frac{\text{inertial force}}{\text{gravitational force}}$$