

Let  $c(x, t)$  = concentration of a chemical in a region  $(x, x + dx)$  at time  $t$ .

In 3 dimensions:  $[c(x, t)] = ML^{-3}$ .

In 1 dimension:  $[c(x, t)] = ML^{-1}$ .

Key:  $c(x, t)$  is measurable

$$\int_a^b c(x, t) dx = \text{total concentration of chemical in } (a, b) \text{ at time } t$$

Check the units

$$\left[ \int_a^b c(x, t) dx \right] = ML^{-1} \times L = M$$

How does a chemical move? **flux**  $J(x, t)$  – amount of substance that passes through  $x$  in the positive direction per unit time.

The rate of change of the amount is

$$\frac{d}{dt} \int_a^b c(x, t) dx = J(a, t) - J(b, t)$$

or

$$\int_a^b \frac{\partial c}{\partial t} dx = - \int_a^b \frac{\partial J}{\partial x} dx \Rightarrow \int_a^b \left( \frac{\partial c}{\partial t} + \frac{\partial J}{\partial x} \right) dx = 0$$

This is the **Conservation of Mass!**

$$\frac{\partial c}{\partial t} + \frac{\partial J}{\partial x} = 0 \tag{1}$$

What is the flux  $J$ ?

How is it related to  $c(x, t)$ ?

## Chemical Diffusion

**Fick's Law** – chemical moves from regions of higher concentration to lower concentration.

$$J \propto -\frac{\partial c}{\partial x}$$

or

$$J = -D \frac{\partial c}{\partial x}$$

where  $D$  is the diffusion constant.

Substitute  $J = -D \frac{\partial c}{\partial x}$  into (1)

$$\frac{\partial c}{\partial t} - D \frac{\partial^2 c}{\partial x^2} = 0 \tag{2}$$

Diffusion equation is a partial differential equation (PDE)

We need auxiliary equation conditions:

$$\left. \begin{aligned} \frac{\partial c}{\partial t} &= D \frac{\partial^2 c}{\partial x^2}, & a < x < b \\ & & t > 0 \end{aligned} \right\}$$

- What was the concentration initially (at the start,  $t = 0$ )
- What happens at the endpoints  $a$  and  $b$ .

### 0.1.1 Aside

Diffusion equations arise in many other settings

#### Heat transfer

Heat energy is measured by temperature  $[\theta]$

Let  $u(x, t)$  = temperature in the bar at  $x$  and time  $t$ .

Same derivation as before:

$$\frac{\partial u}{\partial t} = -\frac{\partial J}{\partial x}$$

Now:  $J(x, t)$  = heat flux

Fourier Law of Heat Conduction:

$$\begin{aligned} J &= -D \frac{\partial u}{\partial x} \\ \frac{\partial u}{\partial t} &= D \frac{\partial^2 u}{\partial x^2} \end{aligned} \tag{3}$$

where  $D$  = thermal diffusion constant

#### Probability

Let  $p(x, t)$  = probability density function of a stochastic process  $X(t)$

$$Pr[a \leq X(t) < b] = \int_a^b p(x, t) dx \tag{4}$$

Example: Random walk (discrete)  $\rightarrow$  Brownian motion (continuous)

**Brownian motion's pdf:**

$$\frac{\partial p}{\partial t} = \frac{1}{2} \frac{\partial^2 p}{\partial x^2}, \quad -\infty < x < \infty, t > 0 \tag{5}$$

where the initial condition ( $X(0) = 0$ )  $\Rightarrow p(x, 0) = \delta$  (The dirac delta function)

$$\delta(x) = \begin{cases} 0 & x \neq 0 \\ \infty & x = 0 \end{cases}$$

such that

$$\int_{-\infty}^{\infty} \delta(x) dx = 1$$

#### Biomedical Application

Drug patch concentration =  $u_0$  (fixed)

$u(x, t)$  = concentration at position  $x$  (depth) and time  $t$  ( $x = 0$  is the surface of the skin)

$$\frac{\partial u}{\partial t} = D \frac{\partial^2 u}{\partial x^2}, \quad 0 < x < \infty, \quad t > 0$$

Initial condition:  $u(x, 0) = 0$ .

Boundary condition:  $u(0, t) = u_0, u \rightarrow 0, x \rightarrow \infty$

## Dimensional Analysis

What is  $[D]$ ?

$$\begin{aligned}\left[\frac{\partial u}{\partial t}\right] &= \left[D \frac{\partial^2 u}{\partial x^2}\right] \\ \left[\frac{\partial u}{\partial t}\right] &= [D] \left[\frac{\partial^2 u}{\partial x^2}\right] \\ ML^{-1}T^{-1} &= [D]ML^{-3} \\ L^2T^{-1} &= [D]\end{aligned}$$

Exact solution of  $u$  is hard!  
Try Laplace transforms

### 0.2.1 Alternate Approach

1. Dimensional reduction to understand the form of the solution
2. Get non-dimensional problem
3. Use dimensionless variables to convert PDE  $\rightarrow$  ODE

$$\begin{aligned}u &= f(x, t, D, u_0) \\ [u] &= [x]^a [t]^b [D]^c [u_0]^d \\ ML^{-1} &= L^a T^b (L^2 T^{-1})^c (ML^{-1})^d \\ M^1 L^{-1} T^0 &= L^a T^b L^{2c} T^{-c} M^d L^{-d} \\ M^1 L^{-1} T^0 &= L^{a+2c-d} T^{b-c} M^d\end{aligned}$$

$$\begin{bmatrix} 1 & 0 & 2 & -1 \\ 0 & 1 & -1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} a \\ b \\ c \\ d \end{bmatrix} = \begin{bmatrix} -1 \\ 0 \\ 1 \end{bmatrix}$$

$$\begin{aligned}a + 2c - d &= -1 & b - c &= 0 & d &= 1 \\ a &= -2c & b &= c & d &= 1\end{aligned}$$

$$\begin{aligned}
\begin{bmatrix} a \\ b \\ c \\ d \end{bmatrix} &= \begin{bmatrix} -2c \\ c \\ c \\ 1 \end{bmatrix} \\
&= \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix} + c \begin{bmatrix} -2 \\ 1 \\ 1 \\ 0 \end{bmatrix} \\
&= u_0^1 + (t^1 D^1 x^{-2})^c \\
&= u_0 + \left( \frac{tD}{x^2} \right)^c \\
u(x, t) &= u_0 F \left( \frac{tD}{x^2} \right)
\end{aligned}$$

Boundary condition makes us multiple the vector within the nullspace by  $-\frac{1}{2}$ .

$$\begin{aligned}
-\frac{1}{2} \begin{bmatrix} -2 \\ 1 \\ 1 \\ 0 \end{bmatrix} &= \begin{bmatrix} 1 \\ -\frac{1}{2} \\ -\frac{1}{2} \\ 0 \end{bmatrix} \\
\begin{bmatrix} a \\ b \\ c \\ d \end{bmatrix} &= \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix} + c \begin{bmatrix} 1 \\ -\frac{1}{2} \\ -\frac{1}{2} \\ 0 \end{bmatrix} \\
&= u_0^1 + \left( t^{-\frac{1}{2}} D^{-\frac{1}{2}} x^1 \right)^c \\
&= u_0 + \left( \frac{x}{\sqrt{tD}} \right)^c \\
u(x, t) &= u_0 F \left( \frac{x}{\sqrt{tD}} \right) \\
\frac{u(x, t)}{u_0} &= F \left( \frac{x}{\sqrt{tD}} \right) \\
\frac{u}{u_0} &= F \left( \frac{x}{\sqrt{tD}} \right)
\end{aligned}$$

is a huge reduction.

Introduce a non-dimensional variable into  $u$ .

$$v(x, t) = \frac{u(x, t)}{u_0},$$

$$[v] = 1, 0 \leq v < 1$$

We can make the similarity variable  $\eta = \frac{x}{\sqrt{tD}}$ .  $[\eta] = 1$

$$u(x, t) = v(\eta)$$

$$u_t = Du_{xx} \Rightarrow v_t = Dv_{xx}$$

The chain rule from Multivariate calculus says that  $v(x, t) = v(\eta(x, t))$

$$\frac{\partial}{\partial t} = \frac{\partial}{\partial \eta} \frac{\partial \eta}{\partial t} = -\frac{1}{2} \frac{x}{\sqrt{D}} \frac{1}{t^{3/2}} \frac{\partial}{\partial \eta}$$

$$\frac{\partial}{\partial x} = \frac{\partial}{\partial \eta} \frac{\partial \eta}{\partial x} = \frac{1}{\sqrt{tD}} \frac{\partial}{\partial \eta}$$

$$\frac{\partial^2}{\partial x^2} = \frac{1}{Dt} \frac{\partial^2}{\partial \eta^2} \leftarrow \text{be able to derive!!}$$

Plug into  $u_t = Du_{xx}$

$$-\frac{1}{2} \frac{x}{\sqrt{D}} \frac{1}{t^{3/2}} v' = D \frac{1}{Dt} v''$$

$$-\frac{1}{2} \frac{x}{\sqrt{D}} \frac{1}{t^{3/2}} v' = \frac{1}{t} v''$$

$$-\frac{1}{2} \frac{x}{\sqrt{D}} \frac{1}{t^{1/2}} v' = v''$$

$$-\frac{1}{2} \frac{x}{\sqrt{Dt}^{1/2}} v' = v''$$

$$-\frac{1}{2} \frac{x}{\sqrt{Dt}} v' = v''$$

$$-\frac{1}{2} \eta v' = v''$$

$$v'' + \frac{1}{2} \eta v' = 0$$

Initial conditional(s) and boundary conditions must be consistent Boundary conditions:

$$\frac{u(0, t)}{u_0} = \frac{u_0}{u_0} \Rightarrow v(0) = 1$$

$$x = 0 \Rightarrow \eta = \frac{x}{\sqrt{Dt}} = 0$$

$$u \rightarrow 0 \text{ as } x \rightarrow \infty \Rightarrow v \rightarrow 0 \text{ if } \eta \rightarrow \infty$$

Initial condition:

$$u(x, 0) = 0$$

$$t = 0 \Rightarrow \eta \rightarrow \infty \Rightarrow v(\infty) = 0$$

Boundary Value ODE Problem

$$v'' + \frac{1}{2} \eta v' = 0, \quad 0 < \eta < \infty$$

$$v(0) = 1, \quad v(\infty) = 0$$

2nd order, linear, not constant coefficient homogeneous.

Let  $g = v'$

$$g' + \frac{1}{2}\eta g = 0$$

$$I = e^{\int \frac{1}{2}\eta d\eta}$$

$$= e^{\frac{\eta^2}{4}}$$

$$e^{\frac{\eta^2}{4}} g' + \frac{1}{2}\eta e^{\frac{\eta^2}{4}} g = 0$$

$$\left( e^{\frac{\eta^2}{4}} g \right)' = 0$$

$$e^{\frac{\eta^2}{4}} g = C$$

$$g = C e^{-\frac{\eta^2}{4}}$$

$$= c_2 e^{-\frac{\eta^2}{4}}$$

$$v = \int g d\eta$$

$$= c_1 + c_2 \int_0^\eta e^{-\frac{\zeta^2}{2}} d\zeta$$

$$v(0) = 1 \Rightarrow c_1 = 1$$

$$v = 1 + c_2 \int_0^\eta e^{-\frac{\zeta^2}{2}} d\zeta$$

$$v(\infty) = 0$$

$$0 = 1 + c_2 \int_0^\infty e^{-\frac{\zeta^2}{2}} d\zeta$$

$$0 = 1 + c_2 [\sqrt{\pi}]$$

$$-1 = c_2 \sqrt{\pi}$$

$$-\frac{1}{\sqrt{\pi}} = c_2$$

$$v(\eta) = 1 - \frac{1}{\sqrt{\pi}} \int_0^\eta e^{-\frac{s^2}{4}} ds$$

Change this back to the original variable:  $\eta = \frac{x}{\sqrt{Dt}}, v = \frac{u}{u_0}$

$$\frac{u(\eta)}{u_0} = 1 - \frac{1}{\sqrt{\pi}} \int_0^\eta e^{-\frac{s^2}{4}} ds$$

$$\frac{u(x, t)}{u_0} = 1 - \frac{1}{\sqrt{\pi}} \int_0^{x/\sqrt{Dt}} e^{-\frac{s^2}{4}} ds$$

$$u(x, t) = u_0 \left[ 1 - \frac{1}{\sqrt{\pi}} \int_0^{x/\sqrt{Dt}} e^{-\frac{s^2}{4}} ds \right]$$

$$= u_0 \left[ 1 - \operatorname{erf} \left( \frac{x}{2\sqrt{Dt}} \right) \right]$$

where erf is the error function

$$\operatorname{erf}(y) = \frac{2}{\sqrt{\pi}} \int_0^y e^{-s^2} ds \quad (6)$$

$$\operatorname{erfc} = 1 - \operatorname{erf}$$

$$\begin{aligned} u(x, t) &= u_0 \left[ 1 - \operatorname{erf} \left( \frac{x}{2\sqrt{Dt}} \right) \right] \\ &= u_0 \operatorname{erfc} \left( \frac{x}{2\sqrt{Dt}} \right) \end{aligned}$$

### Method of Similarity Variables

#### 0.2.2 Aside 2 (Another way of writing erfc)

$$N(x) = \text{cumulative normal distribution} = \int_{-\infty}^x \frac{e^{-\frac{s^2}{2}}}{\sqrt{2\pi}} ds \quad (7)$$

$$N(-\infty) = 0, N(\infty) = 1$$

$$N(x) = \frac{1 + \operatorname{erf} \left( \frac{x}{\sqrt{2}} \right)}{2}$$