ASTRO C207 Radiative Processes in Astrophysics Lydia Lee

Problem Set 5

1. Good Rovibrations

(1) Assume we end in n = 0.

$$\begin{aligned} k_{\text{center}} &= \frac{v_{\text{center}}}{c} \\ v_{\text{center}} &= ck_{\text{center}} \\ &\approx 3(10^{10}) \cdot 2145 \\ &\approx 6.4(10^{13}) \text{Hz} \\ &\approx 1 \times v_{0,\text{CO}} \end{aligned}$$

where $v_{0,CO} \approx 6.7(10^{13})$ Hz is the natural frequency of CO's vibrational transition.

$$|\Delta n| = 1$$

- (2) Boltzmann statistics for populations in each J state
 - Line intensity $\propto n_{J_{\text{upper}}}$

$$\begin{split} \frac{n_{J+1}}{n_J} &= \frac{g_{J+1}}{g_J} e^{-\frac{E_{J+1}-E_J}{k_BT}} \\ \text{where } E_{J+1} - E_J &= \frac{\hbar^2}{2I} \left[(J+1)(J+2) - J(J+1) \right] \\ &= \frac{\hbar^2}{I} (J+1) \end{split}$$

Sweeping $\frac{n_{J+1}}{n_J}$ with respect to T and choosing $J_{\text{infl}} = 7$ where $\frac{n_{J_{\text{infl}}}}{n_{J_{\text{infl}}-1}} > 1$ and $\frac{n_{J_{\text{infl}}+1}}{n_{J_{\text{infl}}}} < 1$ yields a rough temperature range $T \in [288, 365)\text{K}$

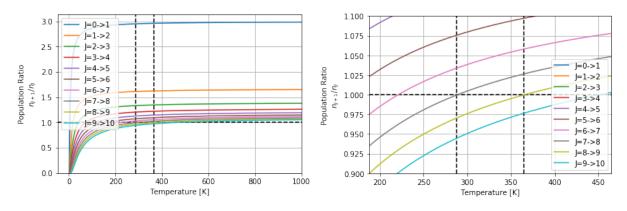


Figure 1: The population ratios vs. temperature for different values of $J \rightarrow J + 1$

Because
$$n_{J=6} \approx n_{J=7}$$

$$T \approx 288$$
K

(3) Looking up the dipole moment $d_{\rm CO} \approx 0.122$ esu.cm and calculating values relative to the Lyman- α

$$A_{\rm CO} = A_{\rm Ly\alpha} \left(\frac{d_{\rm CO}}{d_{\rm H}}\right)^2 \left(\frac{\omega_{\rm CO}}{\omega_{\rm Ly\alpha}}\right)^3$$

with

$$\omega_{\text{CO}} = \frac{\left(\Delta E\right)_{\Delta n = 1}}{\hbar} \tag{1}$$

$$= \omega_{0,CO} \pm \frac{\hbar}{I} J_{\text{upper}} \tag{2}$$

Eqn. ?? assumes a constant moment of inertia. In reality, there will be a change in interatomic distance dependent on n, i.e. the moment of inertia will increase with n.

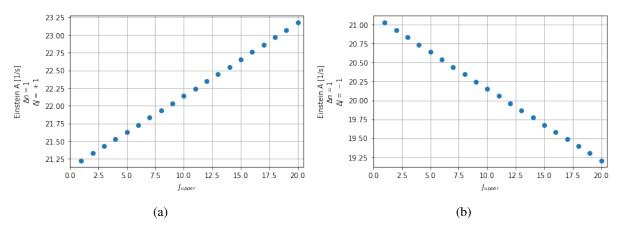


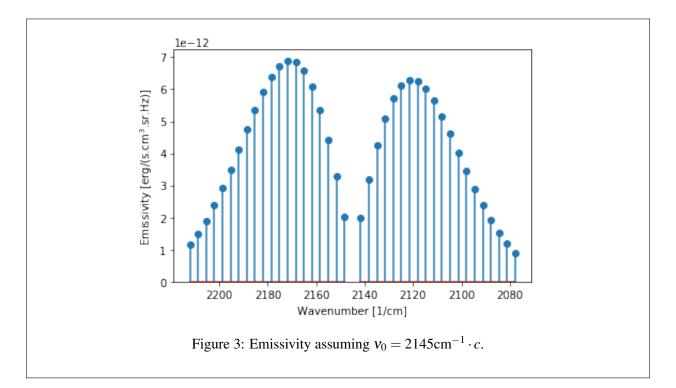
Figure 2: Calculated Einstein A coefficients for the R (??) and P (??) branches

A varies by more than 10% over the range of J—probably enough to warrant inclusion.

(4)

$$j_{m{v}} = rac{h m{v}}{4\pi} n_{J_{
m upper}} A_{21} m{\phi}(m{v})$$

Supposing $\phi(v)$ is constant across frequency,



(5)

Some final discrepancies:

- non-constant line profile function vs. frequency (i.e. Heisenberg uncertainty and Doppler spread) would have each peak smeared out across frequencies
- change in the interatomic distance and consequently the moment of inertia with respect to vibrational energy would cause spreading between the peaks of the P branch vs. the R branch

q_good_rovibrations

October 11, 2021

```
In [1]: ## Lydia Lee
        %matplotlib inline
        import numpy as np
        import matplotlib.pyplot as plt
        from pprint import pprint
In [2]: hbar = 1e-27 # erg.s, Planck's constant
       h_Planck = hbar * 2*np.pi # erg.s, Planck's constant
        c = 3e10 \# cm/s, speed of light
       m_p = 1.6e-24 \# g, proton mass
       a0 = 0.5e-8 \# cm, Bohr radius
       kB = 1.4e-16 # erg/K, Boltzmann constant
       q = 5e-10 # esu, electron charge
       lamda LyA = 121.6e-7 # cm, wavelength of LyA
       nu_LyA = c/lamda_LyA # Hz, frequency of LyA
       A_LyA = 5e8 # 1/s, Einstein A of LyA
       d_CO = .122e-18 # esu.cm, dipole moment of carbon monoxide
       d_H = q*a0 \# esu.cm, dipole moment of hydrogen
       nuO_CO = 4.2e-13/h_Planck # Hz, value given in class
       x0 = 2.4*a0
       m_C = 12*m_p
       m_0 = 16*m_p
       q_mult = 3
       parallel = lambda lst: (sum([1/x for x in lst]))**-1 if 0 not in lst else 0
In [3]: k_center = 2145 # cycles/cm, center wavenumber
       nu_center = k_center*c # Hz
In [4]: def omega0_diatom(m1=m_C, m2=m_0, x=x0, q_mult=q_mult) -> float:
            111
            Inputs:
                m1: Float, g. Mass of one of the atoms in the diatom.
                m2: Float, g. Mass of the other atom in the diatom.
                x: Float, cm. Internuclear distance between the atoms.
                q_mult: Float. The multiplier for the electron charge
```

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used to calculate the coulomb force between the atoms.
    Returns:
        Float. The natural frequency of a diatomic molecule, in rad/s.
   return nu_center * 2*np.pi # pumpkin-eating in the spirit of October
     k\_spring = (q*q\_mult)**2 / x**3 # spring constant
     mu = parallel([m1, m2]) # reduced mass
      return np.sqrt(k_spring/mu)
def E_vib_diatom(n, m1=m_C, m2=m_0, x0=x0, q_mult=q_mult) -> float:
    Inputs:
        n: Integer. Vibrational quantum number.
        m1: Float, g. Mass of one of the atoms in the diatom.
        m2: Float, g. Mass of the other atom in the diatom.
        x0: Float, cm. Atomic separation.
        q_mult: Float. The multiplier for the electron charge
            used to calculate the coulomb force between the atoms.
    Returns:
        Float. The energy of a diatom's nth vibrational state, in ergs.
   omega0 = omega0_diatom(m1, m2, x0, q_mult) # TODO second-order effects don't end u
     k_spring = (q*q_mult)**2 / x0**3 # spring constant
     De = 0.5 * k_spring * x0**2
    return hbar*omega0*(n+0.5) # - (hbar*omega0*(n+0.5))**2/(4*De)
def I_diatom(J, n, m1=m_C, m2=m_0, x0=x0, q_mult=q_mult) -> float:
    Inputs:
        J: Integer. Rotational quantum number.
        n: Integer. Vibrational quantum number.
        m1: Float, g. Mass of one of the atoms in the diatom.
        m2: Float, g. Mass of the other atom in the diatom.
        x0: Float, cm. Atomic separation n=0.
        q_mult: Float. The multiplier for the electron charge
            used to calculate the coulomb force between the atoms.
    Returns:
        Float. The moment of inertia of a diatom in g.cm<sup>2</sup>.
#
     E\_vib = E\_vib\_diatom(n, m1, m2, x0, q\_mult)
#
     k\_spring = (E\_vib/hbar)**2 * parallel([m1, m2])
     k\_spring = (q*q\_mult)**2/x0**3
     De = 1/2 * k_spring * x0**2
#
      a = np.sqrt(k_spring/(2*De))
    x = x0 \# TODO \ adjust \ internuclear \ separation
      x = x0 - 1/a * np.log(1-np.sqrt(E_vib/De))
      x = x0 + np.sqrt(2/k_spring * E_vib)
```

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return parallel([m1, m2]) * x**2
def E_rot_diatom(J, n, m1=m_C, m2=m_0, x0=x0) -> float:
    111
    Inputs:
        J: Integer. Rotational quantum number.
        n: Integer. Vibrational quantum number.
        m1: Float, g. Mass of one of the atoms in the diatom.
        m2: Float, g. Mass of the other atom in the diatom.
        x0: Float, cm. Atomic separation at n=0.
    Returns:
        Float. The energy of a diatom's rotational state with quantum
        numbers J and n in ergs.
    ,,,
    return hbar**2/(2*I_diatom(J,n,m1,m2,x0)) * J * (J+1)
def E_rovib_diatom(J, n, m1=m_C, m2=m_0, x0=x0, q_mult=q_mult) -> float:
    , , ,
    Inputs:
        J: Integer. Rotational quantum number.
        n: Integer. Vibrational quantum number.
        m1: Float, g. Mass of one of the atoms in the diatom.
        m2: Float, g. Mass of the other atom in the diatom.
        x0: Float, cm. Atomic separation at n=0.
        q_mult: Float. The multiplier for the electron charge
            used to calculate the coulomb force between the atoms.
    Returns:
        Float. The energy of the state with vibrational quantum number n
        and rotational quantum number J, in ergs.
    E_vib = E_vib_diatom(n, m1, m2, x0, q_mult)
    E_{rot} = E_{rot\_diatom}(J, n, m1, m2, x0)
    return E_vib + E_rot
def g_degen(J) -> int:
    Inputs:
        J: Integer. Rotational quantum number.
        Integer. Degeneracy of the state associated with rotational quantum
        number J.
    111
    return 2*J + 1
def population_rel(J, n, T, m1=m_C, m2=m_0, x0=x0) -> float:
    111
    Inputs:
    Returns:
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Float. The ratio of populations n_{J}=J_{n_{J}} assuming Boltzmann statistics.
    dE_rot = E_rot_diatom(J+1, n, m1, m2, x0) - E_rot_diatom(J, n, m1, m2, x0)
    return g_degen(J+1)/g_degen(J) * np.exp(-dE_rot/(kB*T))
def omega_rovib(Ji, Jf, delta_n=1, m1=m_C, m2=m_0, x0=x0, q_mult=q_mult) -> float:
    111
    Inputs:
        Ji: Integer. Initial rotational quantum number.
        Jf: Integer. Final rotational quantum number.
        delta_n: Integer. Change in vibrational quantum number.
        m1: Float, g. Mass of one of the atoms in the diatom.
        m2: Float, g. Mass of the other atom in the diatom.
        x0: Float, cm. Atomic separation at n=0.
        q_mult: Float. The multiplier for the electron charge
            used to calculate the coulomb force between the atoms.
    Returns:
        Float. Frequency (in rad/s) of the photon associated with a
        transition\ with\ delta\_n=delta\_n\ and\ delta\_J=Jf-Ji.
    E = E_rovib_diatom(Jf, delta_n, m1, m2, x0, q_mult) - \
        E_rovib_diatom(Ji, 0, m1, m2, x0, q_mult)
    return E/hbar
def A_dip(Ji, Jf, delta_n=1, m1=m_C, m2=m_0, x0=x0, d=d_C0, q_mult=q_mult) -> float:
    Inputs:
        Ji: Integer. Initial rotational quantum number.
        Jf: Integer. Final rotational quantum number.
        delta_n: Integer. Change in vibrational quantum number.
        m1: Float, q. Mass of one of the atoms in the diatom.
        m2: Float, g. Mass of the other atom in the diatom.
        x0: Float, cm. Atomic separation at n=0.
        d: Float, esu*cm. Dipole moment.
        q_mult: Float. The multiplier for the electron charge
            used to calculate the coulomb force between the atoms.
    Returns:
        Float. The Einstein A coefficient of the transition for a dipole.
    omega = omega_rovib(Ji, Jf, delta_n, m1, m2, x0, q_mult)
   nu = omega/(2*np.pi)
    A = A_LyA * (d/d_H)**2 * (nu/nu_LyA)**3
    return A
```

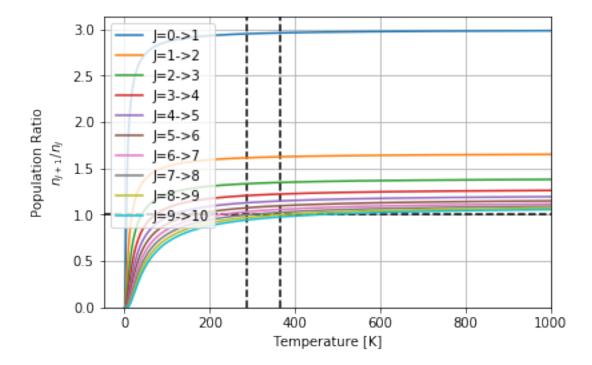
In [5]: k_center = 2145 # cycles/cm, center wavenumber

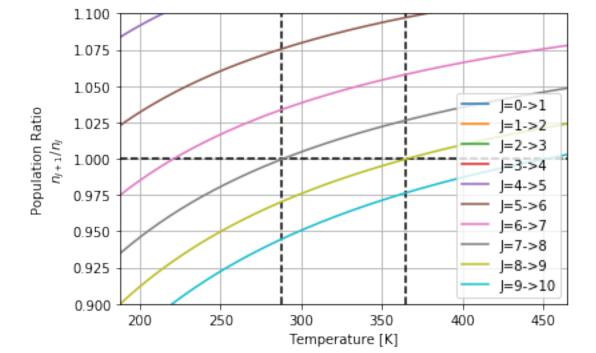
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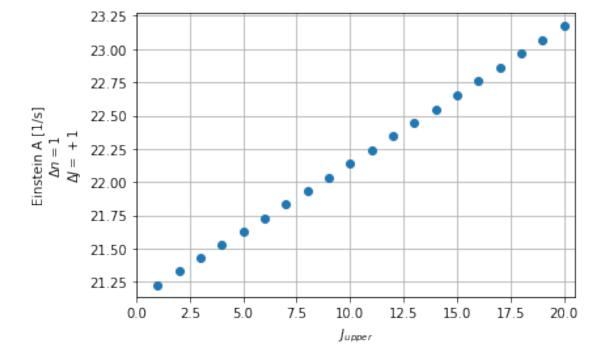
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nu_center = k_center*c # Hz
        \# nu0 = omegaO_diatom(x=x0, q_mult=q_mult)/(2*np.pi)
        Delta_n = nu_center/nu0_CO
        print(f'Delta n (Computed):\t{Delta_n}')
        print("nu_center:\t\t{:.2e} Hz".format(nu_center))
        print(f'n:\t\t\t{int(round(Delta_n))} -> 0')
Delta n (Computed):
                            0.9626737488500153
nu_center:
                           6.44e+13 Hz
                           1 -> 0
n:
   2
In [6]: num_J = 20
        J_vec = np.array(range(num_J))
        T_{\text{vec}} = \text{np.arange}(1, 1000, 1)
        # plt.rcParams['figure.figsize'] = (16, 8)
        # Keeping track of the temperature at which the ratio n_{J+1}/n_{J} exceeds 1
        T_{cross1_J} = [np.inf]*len(J_vec)
        ratio_nJ_dict = dict()
        for i,J in enumerate(J_vec):
            ratio_nJ_dict[J] = population_rel(J, 0, T_vec)
            idx_cross1_J = np.argwhere(np.diff(np.sign(ratio_nJ_dict[J]-1))).flatten()
            assert len(idx_cross1_J) < 2, f'J={J}: ratio does not monotonically increase with
            if len(idx_cross1_J) == 1:
                T_cross1_J[i] = T_vec[idx_cross1_J[0]]
In [7]: J_{infl} = 7 \# value \ of \ J \ where \ n_J/n_{J-1} > 1 \ and \ n_{J+1}/n_{J} < 1
        J_vec_shortened = np.arange(0, 10)
        # Plotting
        for J in J_vec_shortened:
            plt.plot(T_vec, ratio_nJ_dict[J], label=f'J={J}->{J+1}')
        plt.legend()
        plt.grid(True)
        plt.xlabel('Temperature [K]')
        plt.ylabel('Population Ratio\n\$n_{J+1}/n_J\$')
        plt.xlim(*plt.xlim())
        plt.xlim(right=max(T_vec))
        plt.ylim(0)
        # Annotations
        plt.hlines(1, *plt.xlim(), linestyles='dashed')
        T_min = T_cross1_J[J_infl]
```

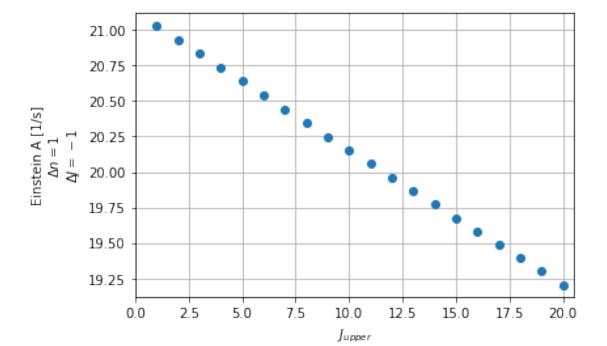
```
T_max = T_cross1_J[J_infl+1]
plt.vlines(T_min, *plt.ylim(), linestyles='dashed')
plt.vlines(T_max, *plt.ylim(), linestyles='dashed')
print(f'Temperature Range: [{T_min}, {T_max}) K')
# For visibility
plt.figure()
for J in J_vec_shortened:
    plt.plot(T_vec, ratio_nJ_dict[J], label=f'J={J}->{J+1}')
plt.legend()
plt.grid(True)
plt.xlabel('Temperature [K]');
plt.ylabel('Population Ratio\n$n_{J+1}/n_J$');
plt.xlim(T_min-100, T_max+100);
plt.ylim(0.9, 1.1);
# Annotations
plt.hlines(1, *plt.xlim(), linestyles='dashed');
plt.vlines(T_min, *plt.ylim(), linestyles='dashed');
plt.vlines(T_max, *plt.ylim(), linestyles='dashed');
```

Temperature Range: [288, 365) K









```
4
In [10]: T_mid = T_min
         idx_T = np.argwhere(np.diff(np.sign(T_vec-T_mid))).flatten()[0]
         nu_vec_incr = omega_rovib(Ji=J_vec, Jf=J_vec+1)/(2*np.pi)
         nu_vec_decr = omega_rovib(Ji=J_vec+1, Jf=J_vec)/(2*np.pi)
         k_vec_incr = nu_vec_incr/c
        k_vec_decr = nu_vec_decr/c
         # Population normalizing with respect to the population @ J=0
         pop_vec = [1] + [0]*len(J_vec)
         for i, J in enumerate(J_vec):
            pop_vec[i+1] = pop_vec[i] * ratio_nJ_dict[J][idx_T]
        pop_upper_vec = pop_vec[1:]
         # Calculating emission coefficient, assuming the line profile function is constant
         j_nu_vec_incr = h_Planck*nu_vec_incr/(4*np.pi) * np.array(A_vec_incr) * pop_upper_vec
         j_nu_vec_decr = h_Planck*nu_vec_decr/(4*np.pi) * np.array(A_vec_decr) * pop_upper_vec
In [11]: plt.stem(k_vec_incr, j_nu_vec_incr)
        plt.stem(k_vec_decr, j_nu_vec_decr)
         xmin, xmax = plt.xlim();
         plt.xlim(xmax, xmin)
```

```
plt.ylim(0)
plt.xlabel('Wavenumber [1/cm]');
plt.ylabel('Emissivity [erg/(s.cm$^3$.sr.Hz)]');
```

