ADVANCED QUANTUM MECHANICS EXERCISE SERIES 8

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Exercise 19

Let $n \in \mathbb{N}$ be a natural number, $A \in \mathbb{R}^{n \times n}$ with $A^{\mathrm{T}} = A$ a symmetric real matrix and $\langle \cdot, \cdot \rangle$ be the dot product on \mathbb{R}^n . Using the spectral theorem, we know that there exists an orthonormal basis $\mathcal{B} := \{b_i \in \mathbb{R}^n \mid i \in \mathbb{N}, i \leq n\}$ such that b is an eigenvector of A for all $b \in \mathcal{B}$. It is now convenient to define the following orthogonal matrix $S \in \mathbb{R}^{n \times n}$.

$$S := (b_1, \dots, b_n)$$

Because S is an orthogonal matrix, we can conclude the following propositions.

$$\det S = \pm 1, \quad S^{\mathrm{T}} S = S S^{\mathrm{T}} = I$$

Additionally, we define $\lambda_i \in \mathbb{R}$ with $Ab_i = \lambda_i b_i$ as the eigenvalue of the i-basis-vector for all $i \in \mathbb{N}$ with $i \leq n$. Of course, these eigenvalues do not have to be distinct. Using $\langle b_i, b_j \rangle = \delta_{ij}$ for all $i, j \in \mathbb{N}$ with $i, j \leq n$ we get the following result.

$$\tilde{A} := S^{T} A S = (b_1, \dots, b_n)^{T} A (b_1, \dots, b_n)$$

$$= (b_1, \dots, b_n)^{T} (A b_1, \dots, A b_n)$$

$$= (b_1, \dots, b_n)^{T} (\lambda_1 b_1, \dots, \lambda_n b_n)$$

$$= \operatorname{diag} (\lambda_1, \dots, \lambda_n)$$

 \tilde{A} is a diagonal matrix. Hence, the determinants of A and \tilde{A} are given by the product of all eigenvalues of A.

$$\det A = \det \tilde{A} = \prod_{k=1}^{n} \lambda_k$$

We apply this result to the the functions $f, \tilde{f} : \mathbb{R}^n \to \mathbb{C}$ with the following definitions for all $x \in \mathbb{R}^n$.

$$f(x) := \exp\left(\frac{\mathrm{i}}{2\hbar} \langle x, Ax \rangle\right)$$

$$\tilde{f}(x) \coloneqq \exp\left(\frac{\mathrm{i}}{2\hbar} \left\langle x, \tilde{A}x \right\rangle\right)$$

By using the bijective function $\varphi \colon \mathbb{R}^n \to \mathbb{R}^n$ with $\varphi(x) := S^{\mathrm{T}}x$ for all $x \in \mathbb{R}^n$, we get the following.

$$f(x) = \exp\left(\frac{\mathrm{i}}{2\hbar} \left\langle SS^{\mathrm{T}} x, ASS^{\mathrm{T}} x \right\rangle\right)$$
$$= \exp\left[\frac{\mathrm{i}}{2\hbar} \left\langle \left(S^{\mathrm{T}} x\right), S^{\mathrm{T}} AS \left(S^{\mathrm{T}} x\right) \right\rangle\right]$$

$$= \exp\left(\frac{\mathrm{i}}{2\hbar} \left\langle \varphi(x), \tilde{A}\varphi(x) \right\rangle\right)$$
$$= \tilde{f} \circ \varphi(x)$$

Now we have to rearrange the explicit form of $\tilde{f}(x)$ for all $x \in \mathbb{R}^n$ by using the definition of \tilde{A} .

$$\tilde{f}(x) = \exp\left(\frac{\mathrm{i}}{2\hbar} \langle x, \operatorname{diag}(\lambda_1, \dots, \lambda_n) x \rangle\right)$$

$$= \exp\left(\frac{\mathrm{i}}{2\hbar} \sum_{k=1}^n \lambda_k x_k^2\right)$$

$$= \prod_{k=1}^n \exp\left(\frac{\mathrm{i}}{2\hbar} \lambda_k x_k^2\right)$$

Since \mathbb{R}^n is an unbounded set and $|f|=\operatorname{id}$ we have $\int_{\mathbb{R}^n}|f|\;\mathrm{d}\lambda^n=\infty$ with λ^n as the n-dimensional Lebesguemeasure on $(\mathbb{R}^n,\mathcal{B}(\mathbb{R}^n))$. Therefore the function f is not Lebesgue-integrable. We have to rely on the improper Riemann-integral. Using $\det \mathrm{D}\varphi(x)=\det S^{\mathrm{T}}=\pm 1$, the transformation theorem and the theorem of Fubini, we are able to cut down the actual calculation to a one-dimensional improper Riemann-integral over complex numbers.

$$\int_{\mathbb{R}^n} f(x) \, \mathrm{d}^n x = \int_{\mathbb{R}^n} \tilde{f}(\varphi(x)) \, \mathrm{d}^n x$$

$$= \int_{\mathbb{R}^n} \tilde{f}(\varphi(x)) \, |\det \mathrm{D}\varphi(x)| \, \mathrm{d}^n x$$
(transformation)
$$= \int_{\varphi(\mathbb{R}^n)} \tilde{f}(y) \, \mathrm{d}^n y$$

$$= \int_{\mathbb{R}^n} \prod_{k=1}^n \exp\left(\frac{\mathrm{i}}{2\hbar} \lambda_k y_k^2\right) \, \mathrm{d}^n y$$
(Fubini theorem)
$$= \prod_{k=1}^n \int_{-\infty}^{\infty} \exp\left(\frac{\mathrm{i}}{2\hbar} \lambda_k x^2\right) \, \mathrm{d} x$$

$$= \prod_{k=1}^n 2 \int_0^{\infty} \exp\left(\frac{-\lambda_k x^2}{2\mathrm{i}\hbar}\right) \, \mathrm{d} x$$

Because the improper Riemann-integral is not well-defined on the complete real axis, some problems may arise when using the transformation theorem on \mathbb{R}^n . We could circumvent these problems by observing the actual limits of the integral.

To evaluate the one-dimensional integrals one has to use Cauchy's theorem for holomorphic functions. For $k\in\mathbb{N}$, $k\leq n$ we define the following curves in \mathbb{C} .

$$\gamma_k \colon \mathbb{R}^+ \to \mathbb{C}, \quad \gamma_k(x) \coloneqq \sqrt{\frac{\lambda_k}{2i\hbar}} x$$

1

Using these curves together with the transformation theorem and the definition of line integrals, we achieve the following result.

$$\int_0^\infty \exp\left(\frac{-\lambda_k x^2}{2i\hbar}\right) dx$$

$$= \int_0^\infty \exp\left(-\gamma_k^2(x)\right) dx$$

$$= \int_0^\infty \exp\left(-\gamma_k^2(x)\right) \frac{\gamma_k'(x)}{\gamma_k'(x)} dx$$

$$= \sqrt{\frac{2i\hbar}{\lambda_k}} \int_{\gamma_k(\mathbb{R}^+)} e^{-z^2} d\gamma_k(z)$$

This equation holds even if $\lambda_k = 0$. In this case this integral is evaluated to ∞ .

Applying Cauchy's theorem on the sector-shaped region given by the direction of γ'_k and the x-axis with radius $R \in \mathbb{R}^+$, we obtain the following with $\tilde{R} := |\gamma'_k| R$.

$$\int_{0}^{R} e^{-t^{2}} dt - \int_{\gamma_{k}((0,\tilde{R}))} e^{-z^{2}} d\gamma_{k}(z) + \underbrace{\int_{0}^{\arg \gamma'_{k}} e^{-R^{2}e^{2i\vartheta}} d\vartheta}_{R \to \infty} = 0$$

As R goes to infinity, the equation becomes the following with the well known half Gaussian integral.

$$\int_{\gamma_k(\mathbb{R}^+)} e^{-z^2} \, \mathrm{d}\gamma_k(z) = \int_0^\infty e^{-t^2} \, \mathrm{d}t = \sqrt{\frac{\pi}{4}}$$

Putting everything together, we are now able to find the solution to the integral.

$$\int_{\mathbb{R}^n} f(x) \, \mathrm{d}^n x = \prod_{k=1}^n 2 \int_0^\infty \exp\left(\frac{-\lambda_k x^2}{2\mathrm{i}\hbar}\right) \, \mathrm{d}x$$

$$= \prod_{k=1}^n 2\sqrt{\frac{2\mathrm{i}\hbar}{\lambda_k}} \int_{\gamma_k(\mathbb{R}^+)} e^{-z^2} \, \mathrm{d}\gamma_k(z)$$

$$= \prod_{k=1}^n \sqrt{\frac{8\mathrm{i}\hbar}{\lambda_k}} \int_0^\infty e^{-t^2} \, \mathrm{d}t$$

$$= \prod_{k=1}^n \sqrt{\frac{2\mathrm{i}\pi\hbar}{\lambda_k}} = \frac{(2\mathrm{i}\pi\hbar)^{\frac{n}{2}}}{\sqrt{\prod_{k=1}^n \lambda_k}}$$

$$= \frac{(2\mathrm{i}\pi\hbar)^{\frac{n}{2}}}{\sqrt{\det A}}$$