SC 618: Flows, derivatives and brackets

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Outline

- 1 Linear and nonlinear systems
- 2 The flow of a vector field
- 3 Lie groups
- 4 Push-forward and pull-back
- **5** The Lie derivative and the Jacobi-Lie bracket
- 6 Lie algebras

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- 1 Linear and nonlinear systems
- 2 The flow of a vector field

- 4 Push-forward and pull-back
- **6** The Lie derivative and the Jacobi-Lie bracket
- 6 Lie algebras

3 Lie groups

Linear systems - preliminaries

A linear system with an input

$$\dot{x} = Ax + Bu \quad x(t) \in \mathbb{R}^n$$

x(t) lives in \mathbb{R}^n , a vector space.

A linear autonomous system

$$\dot{x} = Ax, \quad x(t) \in \mathbb{R}^n$$

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Lie algebras

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A linear autonomous system

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x(t) lives in \mathbb{R}^n , a vector space.

The right-hand side of the differential equation is termed a vector field. For the linear system, it is a linear vector field. Linearity

$$A(\alpha_1 x_1 + \alpha_2 x_2) = \alpha_1 A x_1 + \alpha_2 A x_2$$

Solution and flow

Solution to the set of differential equations

$$x(t) = e^{At}x_0, \quad x(0) = x_0, \quad e^{At} \stackrel{\triangle}{=} I + A + A^2/2! + \dots$$

The term $e^{At}x_0$ is termed the flow associated with the linear vector field Ax.

Nonlinear systems - preliminaries

A nonlinear system with an input

$$\dot{x} = f(x) + g(x)u \quad x(t) \in M$$

 $f(\cdot), g(\cdot)$ are smooth functions, x(t) lives in M, a smooth manifold.

A nonlinear autonomous system

$$\dot{x} = f(x), \quad x(t) \in \mathbb{R}^n$$

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Nonlinear systems - preliminaries

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Nonlinear systems - preliminaries

A nonlinear system with an input

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 $f(\cdot), g(\cdot)$ are smooth functions, x(t) lives in M, a smooth manifold.

A nonlinear autonomous system

$$\dot{x} = f(x), \quad x(t) \in \mathbb{R}^n$$

x(t) lives in M, a smooth manifold.

The right-hand side of the differential equation is a nonlinear vector field. Linearity does not hold.

$$f(\alpha_1 x_1 + \alpha_2 x_2) \neq \alpha_1 f(x_1) + \alpha_2 f(x_2)$$

Solution and flow

Solution to the set of differential equations

$$x(t) = \Phi(t, x_0) \quad x(0) = x_0$$

The term $\Phi(t, x_0)$ is termed the flow associated with the nonlinear vector field f(x).

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Flow of a vector field

Flow of X(x)

The flow of the vector field X(x), denoted by $\Phi(t, x_0)$, is a mapping from $(-a, a) \times U \to \mathbb{R}^n$ (where $a(>0) \in \mathbb{R}$ and U is an open region in the state-space) and satisfies the differential equation

$$\frac{d\Phi(t, x_0)}{dt} = X(\Phi(t, x_0)) \quad \forall t \in (-a, a), x(0) = x_0 \in U.$$

over the interval (-a, a) and with initial conditions starting in the region U.

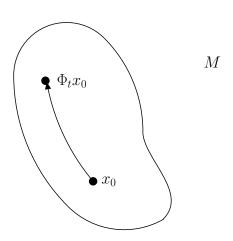


Figure: Flow of a vector field

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Properties of flows

The group structure

Denote

$$\Phi_t(x_0) \stackrel{\triangle}{=} \Phi(t, x_0)$$

The set of transformations $\{\Phi_t\}: U \to \mathbb{R}^n$ satisfies the following properties.

- $\Phi_{t+s}x_0 = \Phi_t \circ \Phi_s x_0 \quad \forall t, s, t+s \in (-a,a)$ (the group binary operation.)
- $\Phi_0 x_0 = x_0$ (the group identity.)
- For a fixed $t \in (-a, a)$ we have $\Phi_t \Phi_{-t} x_0 = x_0 \Rightarrow [\Phi_t]^{-1} = \Phi_{-t}$ (existence of an inverse.)

The flow of a linear system

The group property

Remark

The three properties mentioned above impart a group structure to the set $\{\Phi_t\}$. This set is called a one-parameter (time) group of diffeomorphisms $(\Phi_t$ and its inverse are smooth mappings).

The flow of a linear system

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Linear flow

Remark

For a linear system described by

$$\dot{x} = Ax \quad A \in \mathbb{R}^{n \times n}$$

the flow $\Phi_t x_0 = e^{At} x_0$ where $\{e^{At} : t \in (-\infty, \infty)\}$ constitutes the one-parameter group of diffeomorphisms.

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A group

Definition

A group is a set \mathcal{G} with a binary operation + such that

- For any $x,y\in\mathcal{G},\,x+y\in\mathcal{G}$ (Closure) and (x+y)+z=x+(y+z) (Associativity)
- There exists a unique $i \in \mathcal{G}$ such that x+i=i+x=x for every $x \in \mathcal{G}$ (Existence of the identity element)
- For every $x \in \mathcal{G}$ there exists a unique $y \in \mathcal{G}$ such that x + y = i. (Existence of inverse)

A Lie group

Introduction

Definition

A smooth manifold M together with a group structure is called a Lie group G if the group operation + is smooth.

$$(g,h) \to g + h \quad (\forall g, h \in G)$$
 is smooth

- The identity element of the Lie group is usually denoted by e.
- Left translation of a group

$$L_g: G \to G \qquad h \to g + h$$

• Right translation of a group

$$R_g: G \to G \qquad h \to h + g$$

Examples of Lie groups

- \mathbb{R} or multiple copies of \mathbb{R} (as \mathbb{R}^n) with the bianry operation being the usual component-wise addition +.
- The unit circle S^1 with elements denoted as $\theta \in [0, 2\pi)$ and the binary operation being the usual addition. Similarly, multiple copies of S^1 (as $S^1 \times \ldots \times S^1$).
- The set of $n \times n$ invertible matrices with real entries with the binary operation being matrix multiplication. This group is called $GL(n, \mathbb{R})$.
- The set of $n \times n$ real-orthogonal matrices O(n), a subset of GL(n,R). The set of $n \times n$ rotation matrices SO(n), a subset of O(n,R).

Rigid body motion

Definition

Rigid body motion is characterized by two properties

- The distance between any two points remains invariant
- The orientation of the body is preserved. (A right-handed coordinate system remains right-handed)

SO(3) and SE(3)

Introduction

- Two groups which are of particular interest to us in mechanics and control are SO(3) the special orthogonal group that represents rotations and SE(3) the special Euclidean group that represents rigid body motions. These are Lie groups.
- Elements of SO(3) are represented as 3×3 real matrices and satisfy

$$R^T R = I$$

with det(R) = 1.

• An element of SE(3) is of the form (p, R) where $p \in \mathbb{R}^3$ and $R \in SO(3)$.

Frames of reference or coordinate frames

In describing rigid body motions we always fix two frames of reference.
 One is called the body frame that remains fixed to the body and the other is the inertial frame that remains fixed in inertial space.

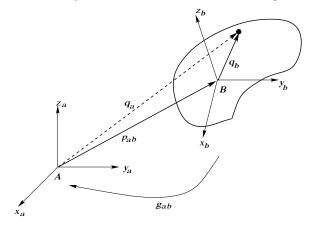


Figure: Rigid body motion

Rigid body motions and groups

 Suppose q_a and q_b are coordinates of a point q relative to frames A and B, respectively.

$$q_a = p_{ab} + R_{ab}q_b$$

Here p_{ab} represents the position of the origin of the frame B with respect to frame A in frame A coordinates and R_{ab} is the orientation of frame B with respect to frame A.

• Appending a "1" to the coordinates of a point (to render the group operation as the usual matrix multiplication)

$$\bar{q}_a = \begin{pmatrix} q_a \\ 1 \end{pmatrix} = \begin{pmatrix} R_{ab} & p_{ab} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} q_b \\ 1 \end{pmatrix} = \bar{g}_{ab}\bar{q}_b$$

Rotation in a plane

Claim

The rotation group SO(2) can be identified with S^1 (the unit circle).

Proof:

$$S^{1} = \{ x \in \mathbb{R}^{2} : ||x|| = 1 \}$$

Parametrize the elements of S^1 in terms of $\theta \in [0, 2\pi]$. For each $\theta \in [0, 2\pi]$, the counter-clockwise rotation of the vectors $\{(1,0),(0,1)\}$ in \mathbb{R}^2 (these form a basis) by the angle θ

$$(1,0) \rightarrow (\cos \theta \quad \sin \theta) \qquad (0,1) \rightarrow (-\sin \theta \quad \cos \theta)$$

is given by the matrix

$$R_{\theta} = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix}$$

which is an element of SO(2).

Proof (contd.)

Conversely, take an element of SO(2) of the form

$$R = \left[\begin{array}{cc} a_1 & a_2 \\ a_3 & a_4 \end{array} \right]$$

Then from the properties of an element of SO(2), we have

$$a_1a_4 - a_2a_3 = 1; a_1^2 + a_3^2 = 1; a_2^2 + a_4^2 = 1; a_1a_2 + a_3a_4 = 0$$

It is possible to find a $\theta \in [0, 2\pi]$ such that that R can be represented in the form R_{θ} .

Lie algebras

Euler's theorem

Theorem

(Euler's theorem)

Every $A \in SO(3)$ is a rotation through an angle $\theta \in S^1$ about an axis $\omega \in \mathbb{R}^3$.

Proof: Since 1 is an eigen value of A, we have Aw = w where $w \in \mathbb{R}^3$ is an eigen vector. Choose two vectors e_1 and e_2 that are orthogonal to each other as well as w. So

$$\langle w, e_1 \rangle = 0, \langle w, e_2 \rangle = 0, \langle e_1, e_2 \rangle = 0$$

The two vectors $\{e_1, e_2\}$ lie in the plane perpendicular to w and it follows that $\{w, e_1, e_2\}$ form a basis for \mathbb{R}^3 . Since A is orthogonal, the matrix of A in this basis is of the form

$$\left[\begin{array}{ccc} 1 & 0 & 0 \\ 0 & a_1 & a_3 \\ 0 & a_2 & a_4 \end{array}\right].$$

Proof (contd.)

(why?) Now

$$\left[\begin{array}{cc} a_1 & a_3 \\ a_2 & a_4 \end{array}\right]$$

is an element of SO(2) and hence there exists a $\theta \in [0, 2\pi]$ such that

$$\begin{bmatrix} a_1 & a_3 \\ a_2 & a_4 \end{bmatrix} = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix}$$

It follows that A is a rotation about w through the angle θ .

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The "star" map (f_*) associated with a smooth function f

Consider a smooth map $f: X \to Y$. At each $p \in X$ we define a linear transformation as follows

$$f_{*p}:T_p(X)\to T_{f(p)}(Y)$$

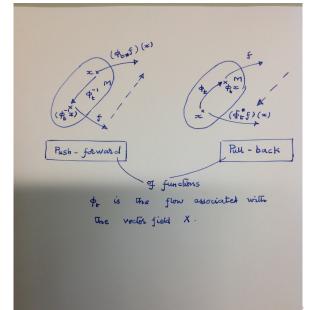
called the derivative of f at p, which is intended to serve as a "linear approximation to f near p,". Visualize this as follows.

- Choose a parametrized curve $c(\cdot): (-\epsilon, \epsilon) \to X$ with c(0) = p and $\frac{dc}{dt}|_{t=0} = v_p$.
- Construct the curve $f \circ c$. Then define

$$f_{*p}(v_p) \stackrel{\triangle}{=} T_p f \cdot v_p = \frac{d}{dt}|_{t=0} (f \circ c)(t)$$

• The rank of f at p is the rank of the Jacobian matrix at x(p) and this is independent of the choice of coordinates x.

Push-forward and pull-back of a function



Push-forward and pull-back of a function

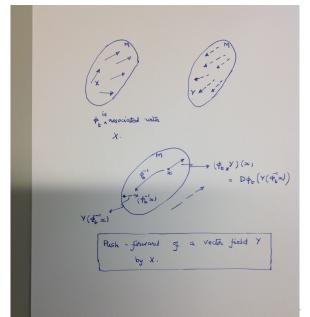
• Suppose X is a vector field on M and $f: M \to \mathbb{R}$ is a smooth function. Then the *push-forward* of the function f on M by the flow of X is the function $\Phi_{t*}f$ defined by

$$(\Phi_{t*}f)(x) \stackrel{\triangle}{=} f \circ \Phi_t^{-1}(x) \quad \forall x \in M$$

• Suppose X is a vector field on M and $f: M \to \mathbb{R}$ is a smooth function. Then the *pull-back* of the function f on M by the flow of X is the function $\Phi_t^* f$ defined by

$$(\Phi_t^* f)(x) \stackrel{\triangle}{=} f \circ \Phi_t(x) \quad \forall x \in M$$

Push-forward and pull-back of a vector field



Push-forward of a vector field by another vector field

Push-forward

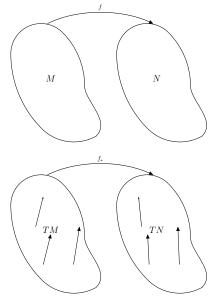
Suppose Φ_t: M → M is the flow associated with a vector field X, then
the push-forward of a vector field Y on M by f is the vector field
(Φ_{t*}Y) on M defined by

$$(\Phi_{t*}Y)(x) = T_{(\Phi_{t}^{-1}x)}[Y(\Phi_{t}^{-1}x)] \quad \forall x \in M$$

• In coordinates

$$(\Phi_{t*}Y)(x) = (D\Phi_t)(Y(\Phi_t^{-1}x))$$

Push forward of vector fields under a diffeomorphism \boldsymbol{f}



Push-forward of vector fields under a diffeomorphism f

Push-forward

Introduction

• Suppose $f: M \to N$ is a diffeomorphism, then the *push-forward* of a vector field X on M by f is the vector field f_*X on N defined by

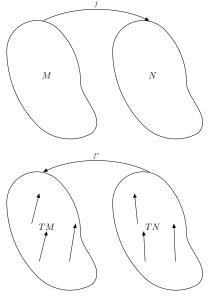
$$(f_*X)(f(x)) = T_x f(X(x)) \quad \forall x \in M$$

In coordinates

$$y = f(x)$$
 $(f_*X)(y) = Df(x).X(x) = \frac{dy}{dx} \cdot X(x)$

Lie algebras

Pull back of a vector field under a diffeomorphism f



Vector Field f^*Y

Pull-back of vector fields under a diffeomorphism f

The pull-back

Introduction

• Suppose $f: M \to N$ is a diffeomorphism, then the *pull-back* of a vector field Y on N by f is the vector field f^*Y on M defined by

$$f^*Y = (f^{-1})_*Y = Tf^{-1} \circ Y \circ f$$

• In coordinates

$$y = f(x)$$
 $(f_*X)(y) = Df(x).X(x) = \frac{dy}{dx} \cdot X(x)$

Lie algebras

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Operations on vector fields

The Gradient

Consider a smooth function $g(\cdot):U\to\mathbb{R}$. The gradient of such a function, denoted by ∇g , is defined as

$$\nabla g(x) = \left[\begin{array}{ccc} \frac{\partial g}{\partial x_1} & \cdots & \frac{\partial g}{\partial x_n} \end{array} \right]$$

alternate notation: grad(g).

Lie algebras

The Lie derivative of a function

The Lie derivative

The Lie derivative of a function f along X is

$$(\mathcal{L}_X f)(x) = \frac{d}{dt}|_{t=0}(\Phi_t^* f)(x) = \frac{d}{dt}|_{t=0} f \circ \Phi_t(x)$$

In coordinates we have the familiar

$$(\mathcal{L}_X f)(x) = \begin{bmatrix} \frac{\partial f}{\partial x_1} & \cdots & \frac{\partial f}{\partial x_n} \end{bmatrix} X(x)$$

The Lie derivative of a function

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In coordinates we have the familiar

$$(\mathcal{L}_X f)(x) = \begin{bmatrix} \frac{\partial f}{\partial x_1} & \cdots & \frac{\partial f}{\partial x_n} \end{bmatrix} X(x)$$

Alternate notation

$$(Xf)(x) = \frac{d}{dt}|_{t=0}f \circ \Phi_t(x) = \lim_{t \to 0} \frac{f(\Phi_t(x)) - f(x)}{t}$$

High school physics

The cross product

- Vector space \mathbb{R}^3 and the cross-product operation \times .
 - $(\alpha_1 a_1 + \alpha_2 a_2) \times b = \alpha_1 (a_1 \times b_1) + \alpha_2 (a_2 \times b_2)$ linearity. (holds in the second argument as well.)
 - $a \times b = -b \times a$ skew-commutative.
 - $a \times (b \times c) + c \times (a \times b) + b \times (c \times a) = 0$ the Jacobi-Lie identity.

Comment: the cross-product of two linearly independent vectors in \mathbb{R}^3 yields a vector in a new *direction*.

Flow

Introduction

The cross product

- Vector space \mathbb{R}^3 and the cross-product operation \times .
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Comment: the cross-product of two linearly independent vectors in \mathbb{R}^3 yields a vector in a new direction.

An alternate notation

$$a \times b \leftrightarrow \hat{a}b$$
 $\hat{a} = \begin{bmatrix} 0 & -a_3 & a_2 \\ a_3 & 0 & -a_1 \\ -a_2 & a_1 & 0 \end{bmatrix}$

The Lie derivative of a vector field

The pull back of a vector field

The Lie derivative of Y along X is

$$\mathcal{L}_X Y \stackrel{\triangle}{=} \frac{d}{dt}|_{t=0} \Phi_t^* Y$$

where Φ is the flow of X.

Lie algebras

The Lie derivative of a vector field

The pull back of a vector field

The Lie derivative of Y along X is

$$\mathcal{L}_X Y \stackrel{\triangle}{=} \frac{d}{dt}|_{t=0} \Phi_t^* Y$$

Push-forward and pull-back

where Φ is the flow of X.

Explicitly

$$(\mathcal{L}_X Y)(x) = \frac{d}{dt}|_{t=0} (D\Phi_t(x))^{-1} \cdot Y(\Phi_t(x))$$

The Lie bracket

In coordinates we have the familiar expression

$$\frac{d}{dt}|_{t=0}(D\Phi_t(x))^{-1} \cdot Y(\Phi_t(x)) = \frac{\partial Y}{\partial x}X(x) - \frac{\partial X}{\partial x}Y(x) = [X,Y](x)$$

The Jacobi-Lie bracket

The Jacobi-Lie bracket of two vector fields is an operation between two vector fields that yields another vector field. For two vector fields X and Y, both defined from U to \mathbb{R}^n , it is defined as

$$[X,Y] = (\mathcal{L}_X Y) = (DY) \cdot X - (DX) \cdot Y$$

and satisfies the following properties (for any three vector fields X, Y, Z)

- $[\alpha X + \beta Y, Z] = \alpha [X, Z] + \beta [Y, Z]$ linearity in the first argument (also hold for the second argument)
- [X,Y] = -[Y,X] skew-commutative.
- [X, [Y, Z]] + [Z, [X, Y]] + [Y, [X, Z]] = 0 the Jacobi-Lie identity

More properties

The Jacobi-Lie bracket

Let X generate the flow $\{\Phi_t\}$ and Y generate the flow $\{\Psi_t\}$. Then [X,Y]=0 if and only if $\Phi_t\circ\Psi_s=\Psi_s\circ\Phi_t$ for all $s,t\in\mathbb{R}$.

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The Lie algebra - $\mathfrak{so}(3)$

3×3 skew-symmetric matrices

Recall

$$\omega \times x \leftrightarrow \hat{\omega}x \qquad \hat{\omega} = \begin{bmatrix} 0 & -\omega_3 & \omega_2 \\ \omega_3 & 0 & -\omega_1 \\ -\omega_2 & \omega_1 & 0 \end{bmatrix}$$

The eigen values of $\hat{\omega}$ are $0, \pm \|\omega\|$ i (Hint: The trace of a matrix is the sum of its eigen values.)

Claim

Exponential of a skew-symmetric matrix is a rotation matrix

To show $e^{\hat{\omega}} \in SO(3)$

$$(e^{\hat{\omega}})(e^{\hat{\omega}})^T = (e^{(\hat{\omega} - \hat{\omega})}) = I \Rightarrow det(e^{\hat{\omega}}) = \pm 1$$

Now from $\omega = 0$, $e^{\hat{\omega}} = I$ and $\det(I) = 1$. The determinant is a continuous function of the elements of the matrix

$$\Rightarrow det(e^{\hat{\omega}}) = 1$$

Properties of the Lie algebra - $\mathfrak{so}(3)$

It is a vector space of dimension 3.

The tangent space of the identity of SO(3) i. e. $T_e SO(3) = \mathfrak{so}(3)$.

The bracket operation $[\cdot,\cdot]:\mathfrak{so}(3)\times\mathfrak{so}(3)\to\mathfrak{so}(3)$ satisfies

- $[\alpha \hat{x} + \beta \hat{y}, \hat{z}] = \alpha[\hat{x}, \hat{z}] + \beta[\hat{y}, \hat{z}]$ linearity in the first argument (also hold for the second argument)
- $[\hat{x}, \hat{z}] = -[\hat{z}, \hat{x}]$ skew-commutative.
- $[\hat{x},[\hat{y},\hat{z}]]+[\hat{z},[\hat{x},\hat{y}]]+[\hat{y},[\hat{z},\hat{x}]]=0$ the Jacobi-Lie identity

Notice that the cross product relates as

$$[\hat{x},\hat{z}] = \widehat{x \times z}$$