



Minimal pole method for spectral functions

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J. Chem. Phys. 162, 214111 (2025), L. Zhang et al.

Introduction – Spectral function

Spectral function (measured by ARPES):

$$A(\omega) = \sum_l A_l^{(\text{real})} \delta(\omega - \xi_l^{(\text{real})}) \text{ with } A_l \text{ (fermion) or } \text{sgn}(\omega)A_l \text{ (boson) being positive}$$

Finite system: a finite number of delta peaks

Infinite system: infinite delta peaks with infinitesimal weights \rightarrow broadened peaks

Connect with imaginary axis:

$$G^{\text{Mat}}(i\omega_n) = \int_{-\infty}^{+\infty} d\omega \frac{A(\omega)}{i\omega_n - \omega} \text{ and } G(\tau) = \frac{1}{\beta} \sum_{n=-\infty}^{+\infty} e^{-i\omega_n \tau} G(i\omega_n)$$

Connect with real axis:

$$\text{Im}[G^{\text{Ret}}(\omega)] = -\pi A(\omega), \text{ Re}[G^{\text{Ret}}(\omega)] = -\mathcal{H}[\text{Im}[G^{\text{Ret}}(\omega)]] \text{ and } G^{\text{Ret}}(t) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} d\omega e^{-i\omega t} G^{\text{Ret}}(\omega)$$

$$\text{BCF: } C(t) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} d\omega e^{-i\omega t} A(\omega) \left[\coth\left(\frac{\beta\omega}{2}\right) + 1 \right]$$

Introduction – Complex pole representation

Ansatz:

$$A(z) = \sum_{l=1}^M \left(\frac{A_l^{(\text{dn})}}{z - \xi_l^{(\text{dn})}} + \frac{A_l^{(\text{up})}}{z - \xi_l^{(\text{up})}} \right) \text{ with } \xi_l^{(\text{up})} = (\xi_l^{(\text{dn})})^* \text{ and } A_l^{(\text{up})} = (A_l^{(\text{dn})})^*$$

With $z \in \mathbb{C}$ and $A(z = \omega) \approx A_{\text{exact}}(\omega)$

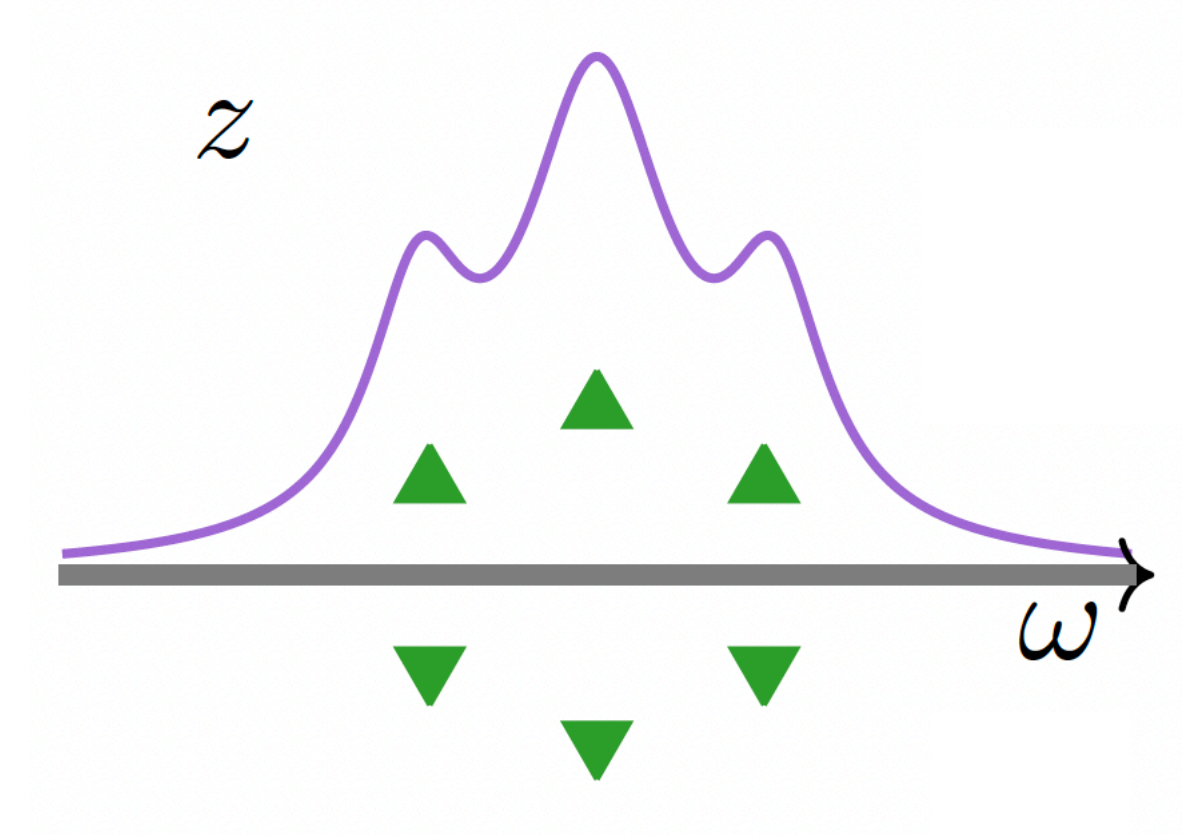
Applications:

$$1. \text{ Imaginary axis: } G(z) = -2\pi i \sum_{l=1}^M \frac{A_l^{(\text{dn})}}{z - \xi_l^{(\text{dn})}}, \quad \text{Im} z > 0$$

Hilbert transform; recover GF in the upper-half plane, e.g., $G(z = i\omega_n) = G^{\text{Mat}}(i\omega_n)$

2. Real-time simulations, e.g., HEOM:

$$A(\omega) = J(\omega) \left[\coth \left(\frac{\beta\omega}{2} \right) + 1 \right] \rightarrow C(t) = \sum_{l=1}^M \eta_l e^{-\gamma_l t} \text{ with } \eta_l = -iA_l^{(\text{dn})} \text{ and } \gamma_l = i\xi_l^{(\text{dn})}$$



Introduction — Prony-like methods

Input: $\{f(t_k), t_k\}$ on a uniform grid

$$\text{Ansatz: } f(t) \approx \sum_{i=1}^M R_i e^{s_i t} \xrightarrow{\text{discretize}} f(t_k) \approx \sum_{i=1}^M R_i z_i^k$$

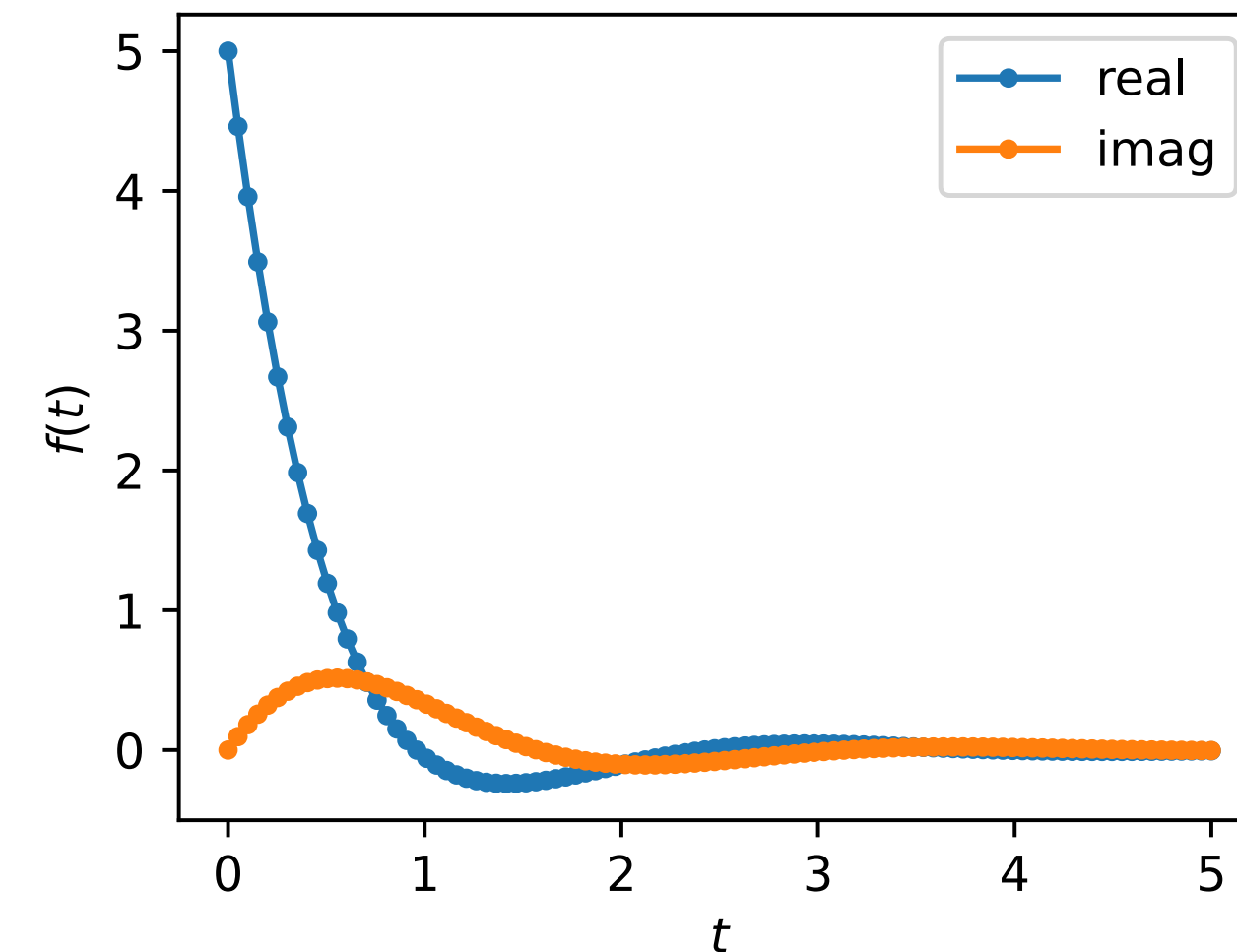
Goal: estimate $M \in \mathbb{Z}$, $R_i \in \mathbb{C}$ and $z_i \in \mathbb{C}$

Solution: ESPRIT, Matrix Pencil, Prony approximation...

IEEE Trans. Acoust. Speech,
Signal Process. 37, 984 (1989)

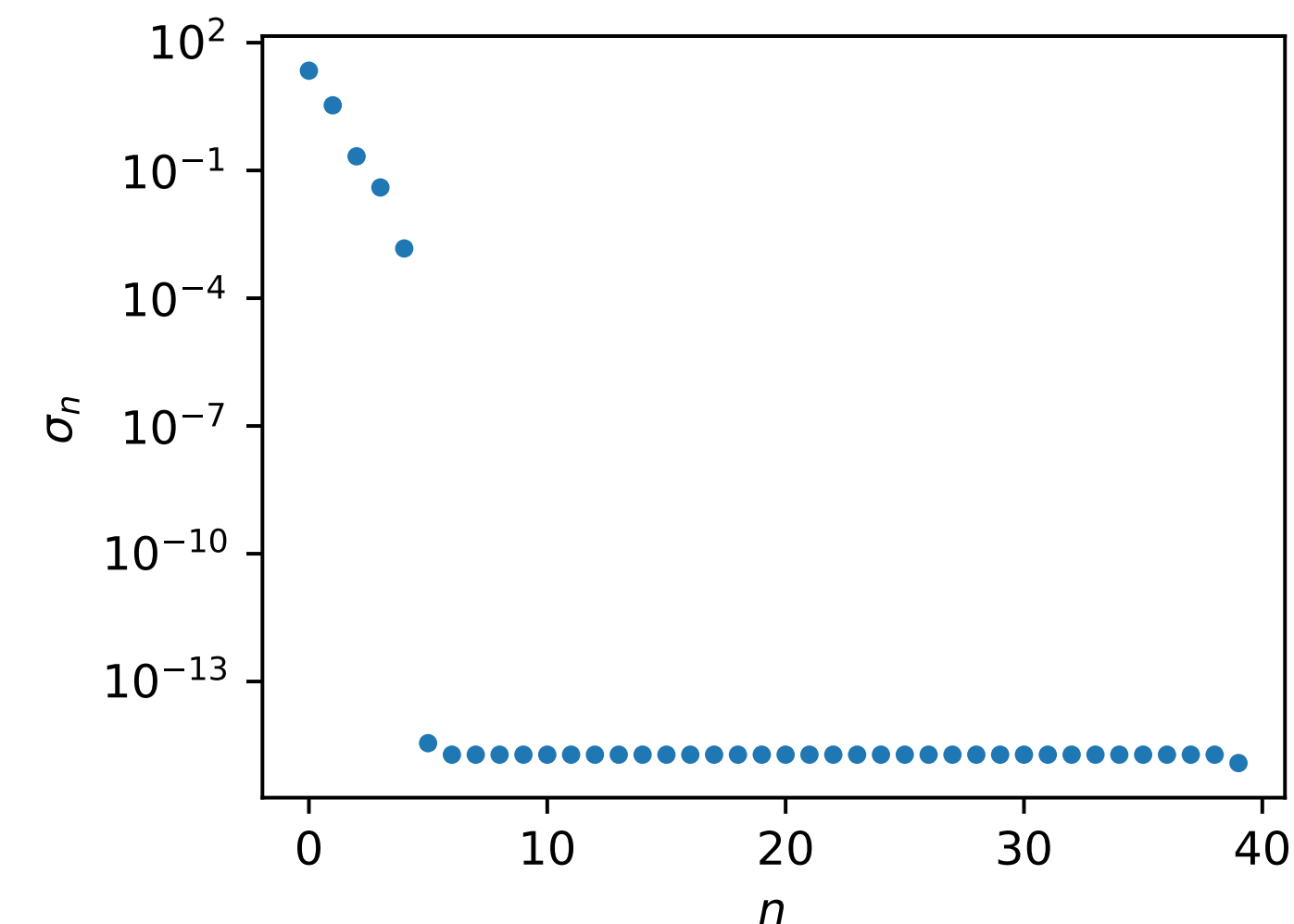
IEEE Trans. Acoustics, Speech,
Signal Process. 38, 814 (1990)

Appl. Comput. Harmonic Anal. 19, 17 (2005)



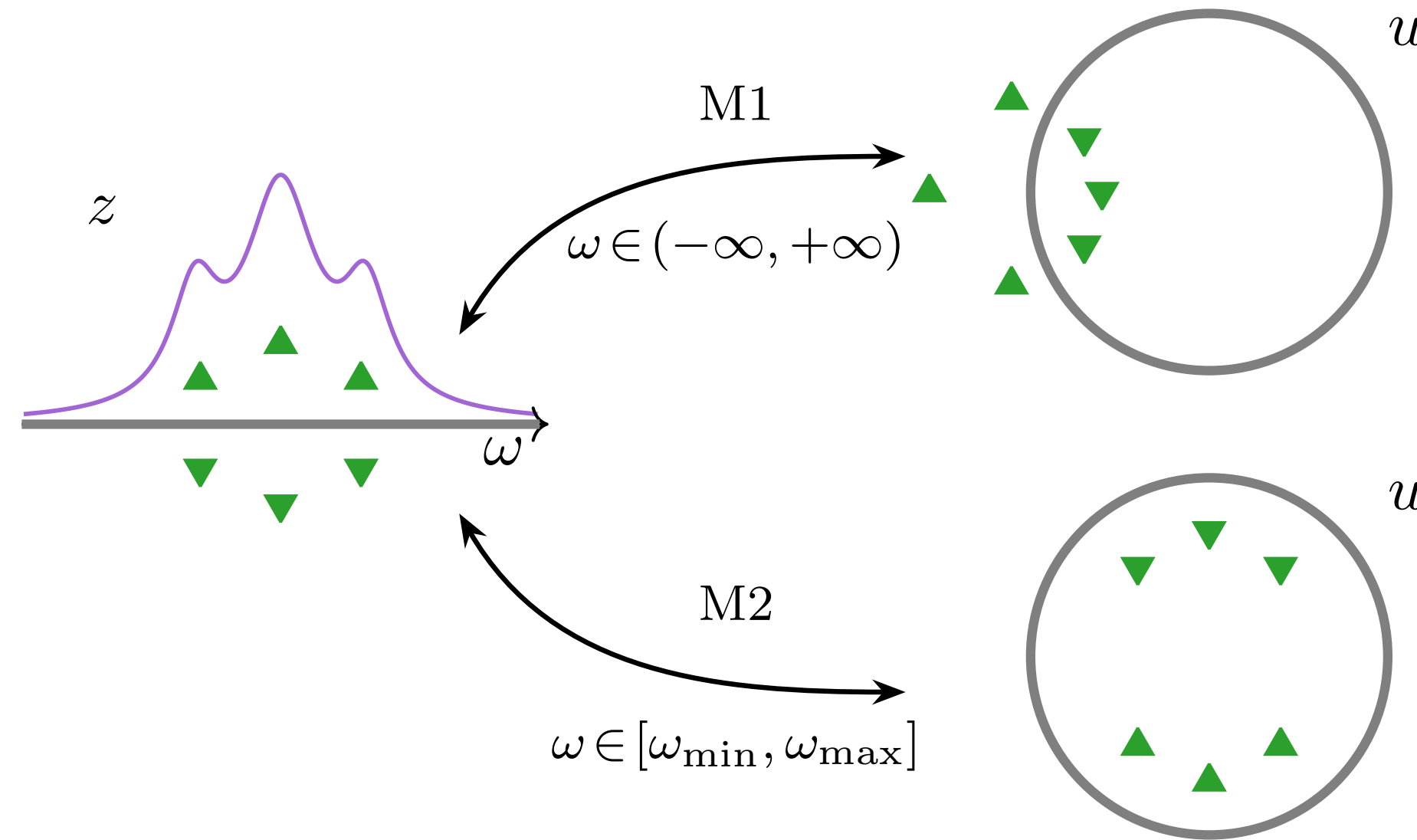
ESPRIT:

1. Construct $(N - L) \times (L + 1)$ Hankel matrix $H_{ij} = f(t_{i+j})$, with $\frac{N}{3} \leq L \leq \frac{N}{2}$
2. Perform SVD: $H = U \Sigma W$
3. Obtain M from $\sigma_{M+1} < \varepsilon$
4. Obtain z_i from eigenvalues of $(W(1 : M, 1 : L)^T)^+ W(1 : M, 2 : L + 1)^T$
5. Obtain R_i from least-square fits



Robust to noise; use a minimal number of complex exponentials to fit the data to the given precision ε

Minimal Complex Pole Method (MPM)



Mapping 1 (M1):

$$\begin{cases} u = f(z) = \frac{z+i\omega_p}{z-i\omega_p} \\ z = f^{-1}(u) = i\omega_p \frac{u+1}{u-1} \end{cases}$$

Mapping 2 (M2):

$$\begin{cases} u = f(z) = z_s + \sqrt{z_s^2 - 1} \text{ with } z_s = \frac{z-\omega_m}{\Delta\omega_h} \\ z = f^{-1}(u) = \frac{\Delta\omega_h}{2} \left(u + \frac{1}{u}\right) + \omega_m \end{cases}$$

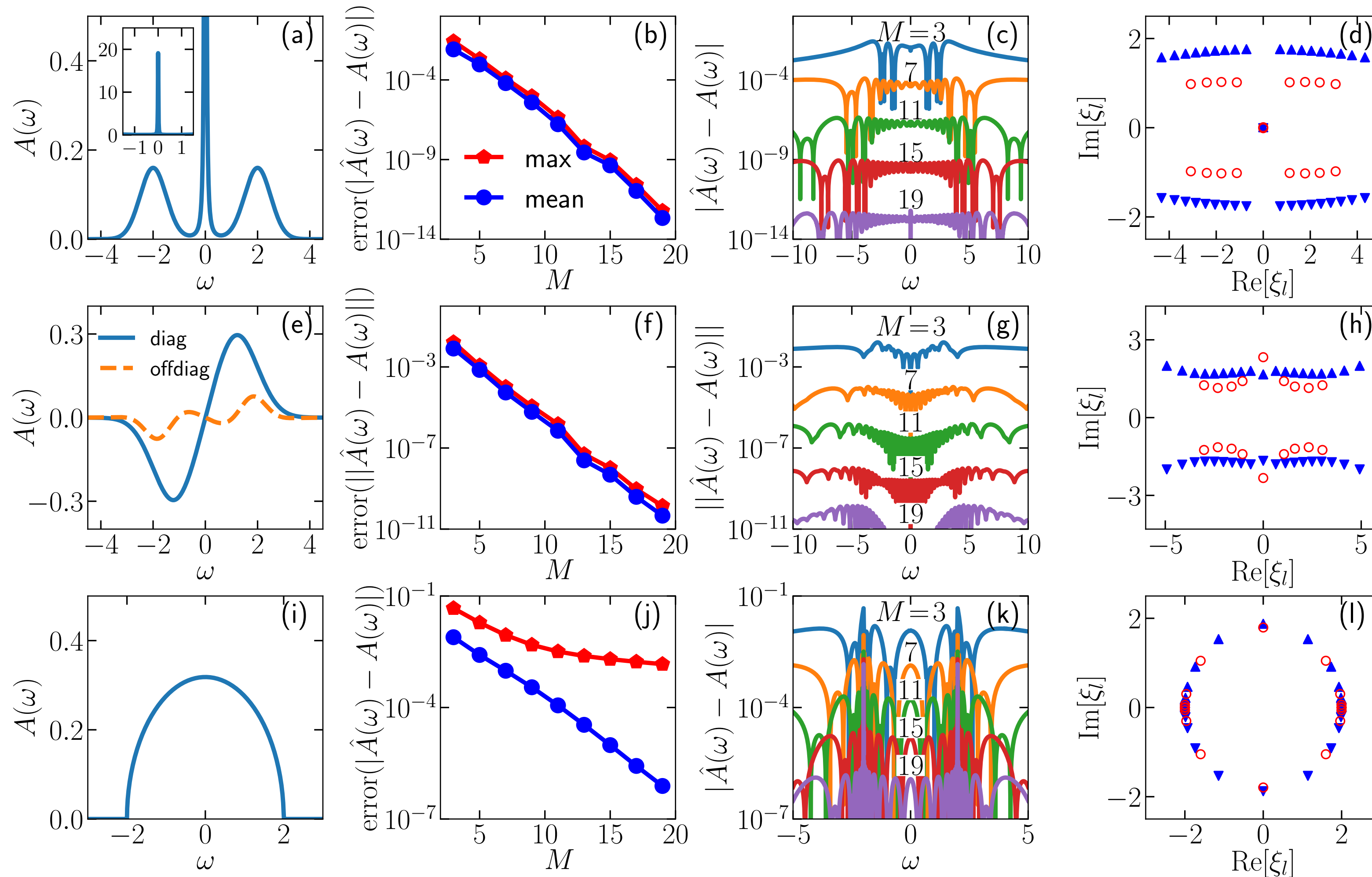
Assume: analytic expression of $A(\omega)$ is known

Step 1 $A(z) = \sum_{l=1}^M \left(\frac{A_l^{(\text{dn})}}{z - \xi_l^{(\text{dn})}} + \frac{A_l^{(\text{up})}}{z - \xi_l^{(\text{up})}} \right) \longleftrightarrow A'(u) = \sum_{l=1}^M \left(\frac{A_l'^{(\text{dn})}}{u - \xi_l'^{(\text{dn})}} + \frac{A_l'^{(\text{up})}}{u - \xi_l'^{(\text{up})}} \right) + \text{analytic part}$ Origin: Numerical analytic continuation L. Zhang et al., Phys. Rev. B 110, 035154 (2024)
L. Zhang et al., Phys. Rev. B 110, 235131 (2024)

Step 2 Transformed moments $h_k := \frac{1}{2\pi i} \int_{\partial \bar{D}} du A'(u) u^k, \quad k \geq 0, \quad h_k = \sum_l A_l' \xi_l'^k$ Imperfect input: utilize AAA, trapezoidal rule...

Step 3 ESPRIT to solve A_l' and ξ_l' and transform back to obtain A_l and ξ_l Multiple orbitals: matrix-valued ESPRIT L. Zhang et al., Phys. Rev. B 110, 235131 (2024)

Results – spectral functions

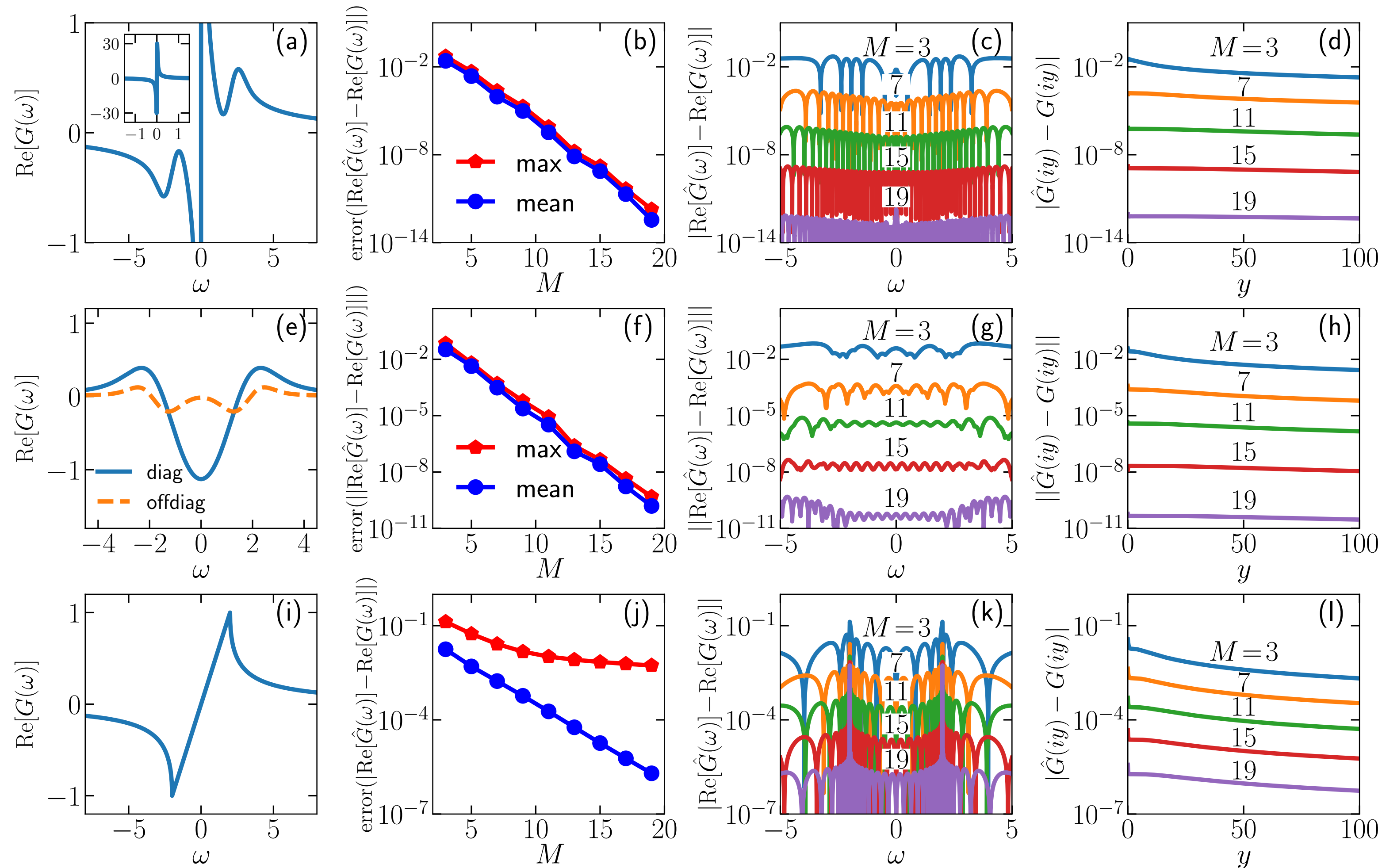


- (a) Kondo-like
- (b) Bosonic multi-orbital
- (c) Semicircular

MPM (M1) was used

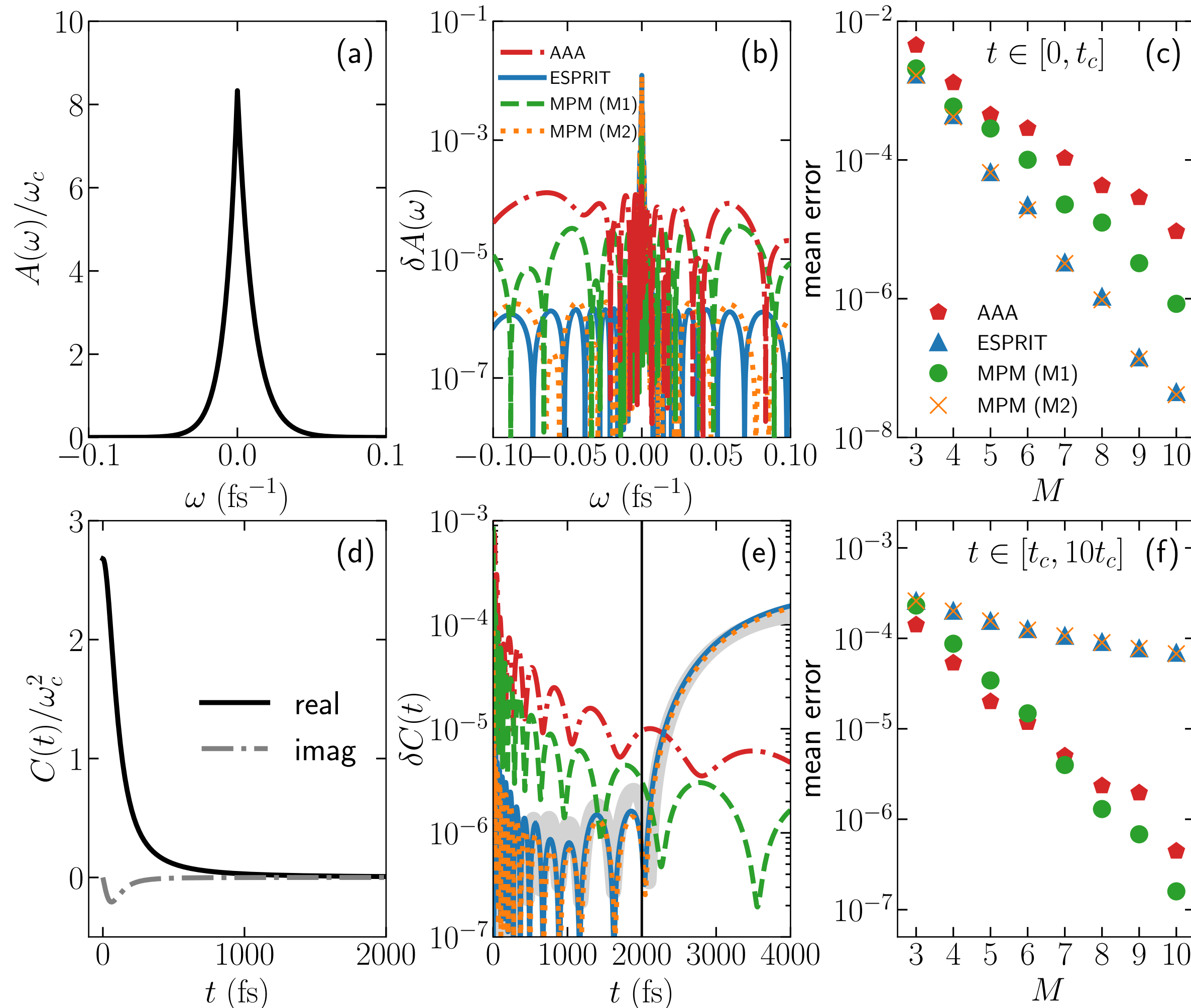
- Converges exponentially away from singularities
- Convergence slows down near the singularities

Results — recovery of GF's



Results — BCF (Ohmic bath)

$T = 300K$ ($\beta\omega_c \approx 0.240$)



$$J(\omega > 0) = \omega e^{-\omega/\omega_c} \text{ with } J(\omega < 0) = -J(-\omega)$$

$$A(\omega) = J(\omega) \left[\coth \left(\frac{\beta\omega}{2} \right) + 1 \right]$$

AAA: logarithmic discretized grid on $\omega \in [-\omega_{\max}, \omega_{\max}]$

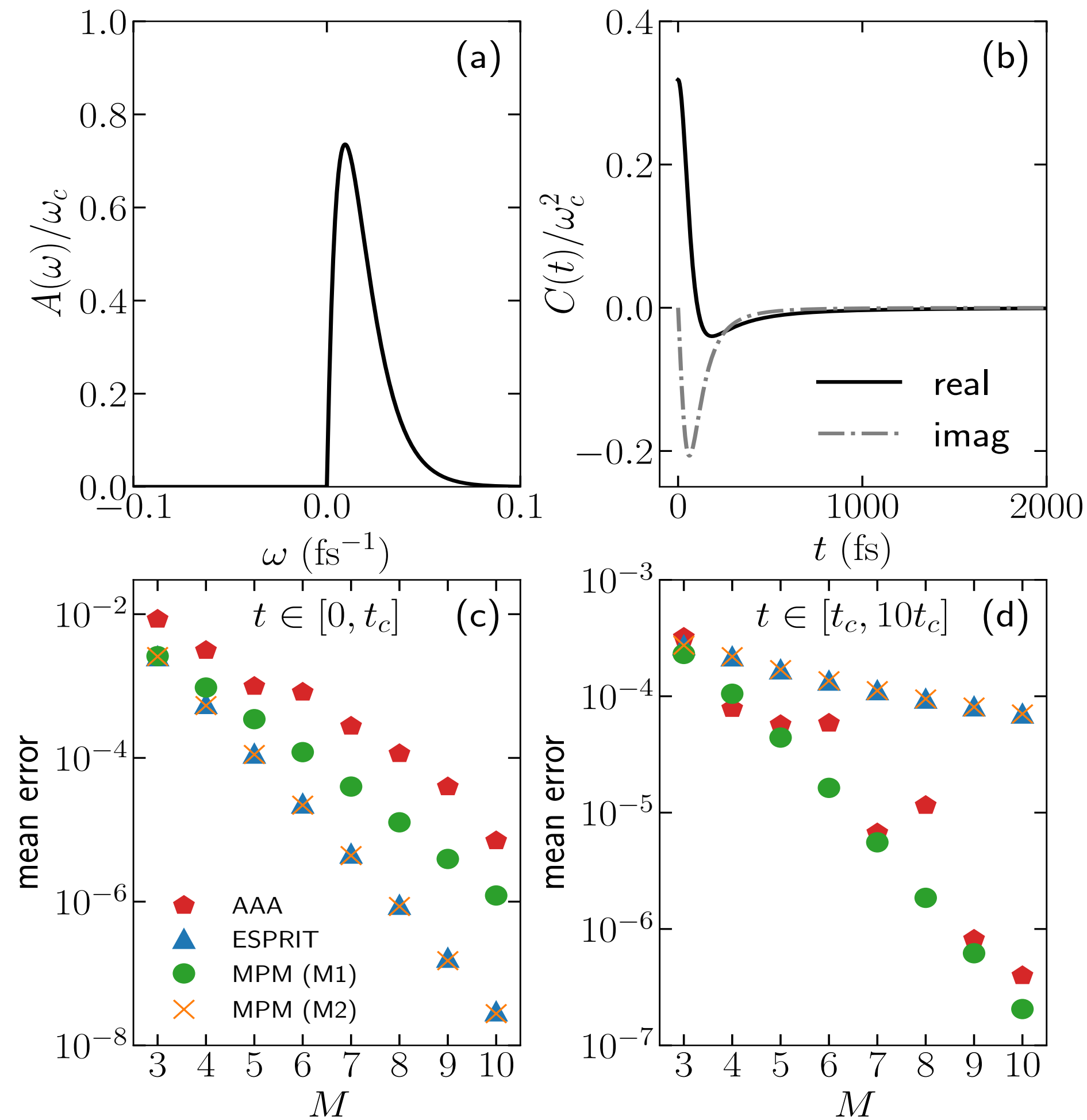
ESPRIT: uniform grid on $t \in [0, t_c]$

Gray line: $|\hat{C}^{(L_1)}(t) - \hat{C}^{(L_2)}(t)| / |\hat{C}^{(L_1)}(t)|$ for MPM (M2)

- MPM (M2) has similar performance compared to ESPRIT
- MPM (M1) has better control over the long tail
- MPM converges faster than AAA

Results — BCF (Ohmic bath)

$$T = 0.001K (\beta\omega_c \approx 7.19 \times 10^4)$$



$$J(\omega > 0) = \omega e^{-\omega/\omega_c} \text{ with } J(\omega < 0) = -J(-\omega)$$

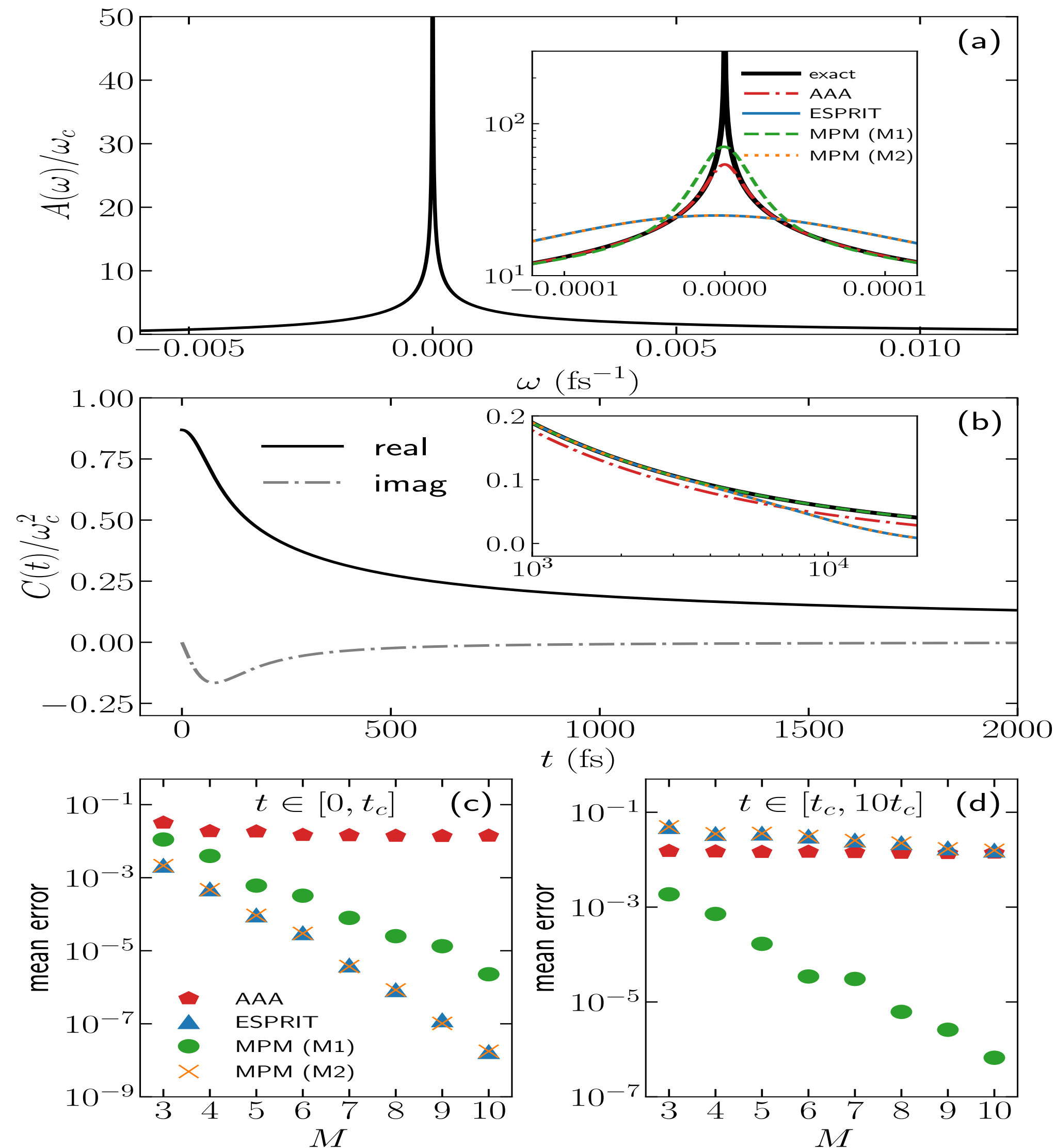
$$A(\omega) = J(\omega) \left[\coth \left(\frac{\beta\omega}{2} \right) + 1 \right]$$

AAA: logarithmic discretized grid on $\omega \in [-\omega_{\max}, \omega_{\max}]$

ESPRIT: uniform grid on $t \in [0, t_c]$

Similar performance as for $T = 300K$

Results — BCF (Sub-Ohmic bath)



$$J(\omega > 0) = (\omega_c \omega)^{0.5} e^{-\omega/\omega_c} \text{ with } J(\omega < 0) = -J(-\omega)$$

$$A(\omega) = J(\omega) \left[\coth \left(\frac{\beta \omega}{2} \right) + 1 \right]$$

AAA: logarithmic discretized grid on $\omega \in [-\omega_{\max}, \omega_{\max}]$

ESPRIT: uniform grid on $t \in [0, t_c]$

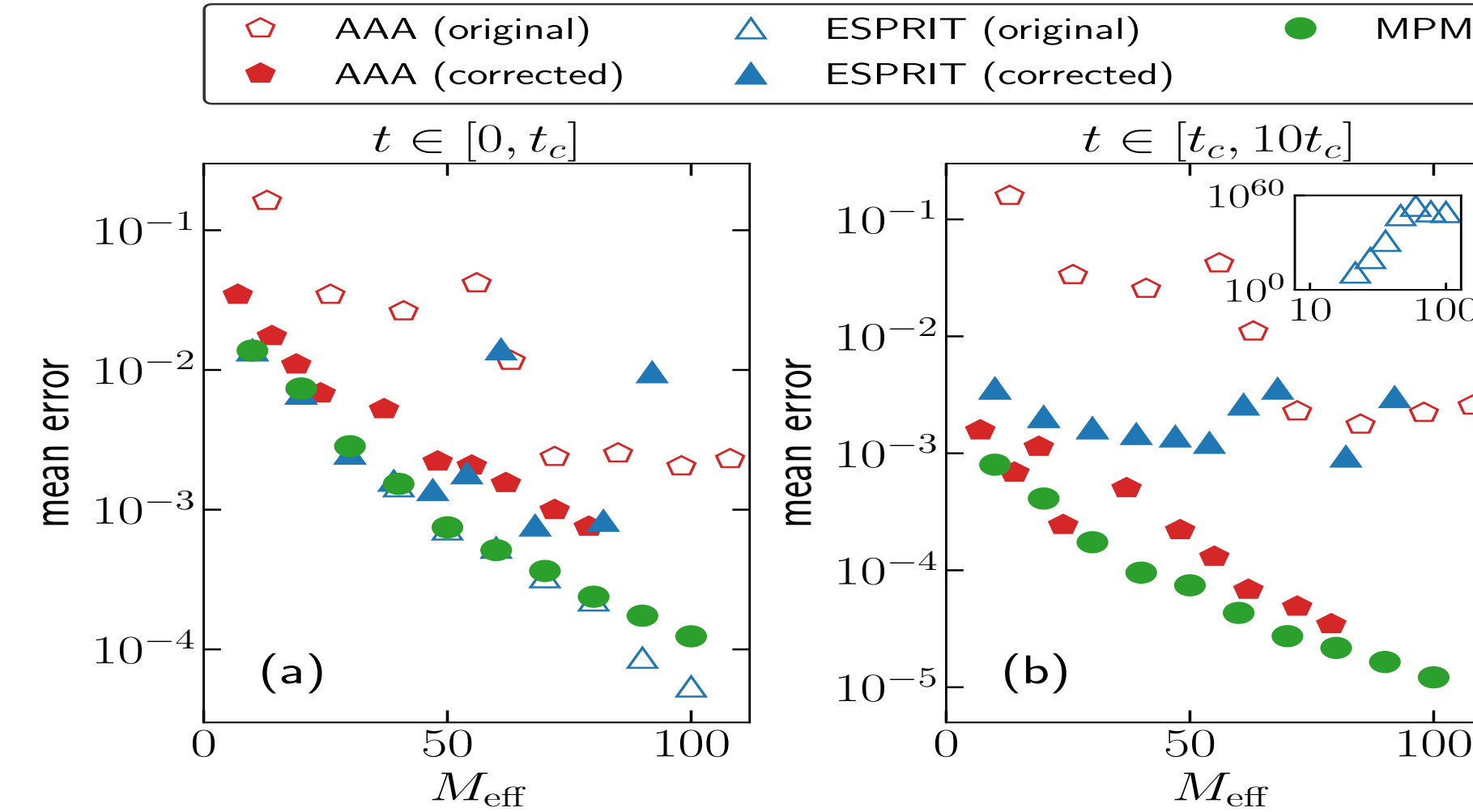
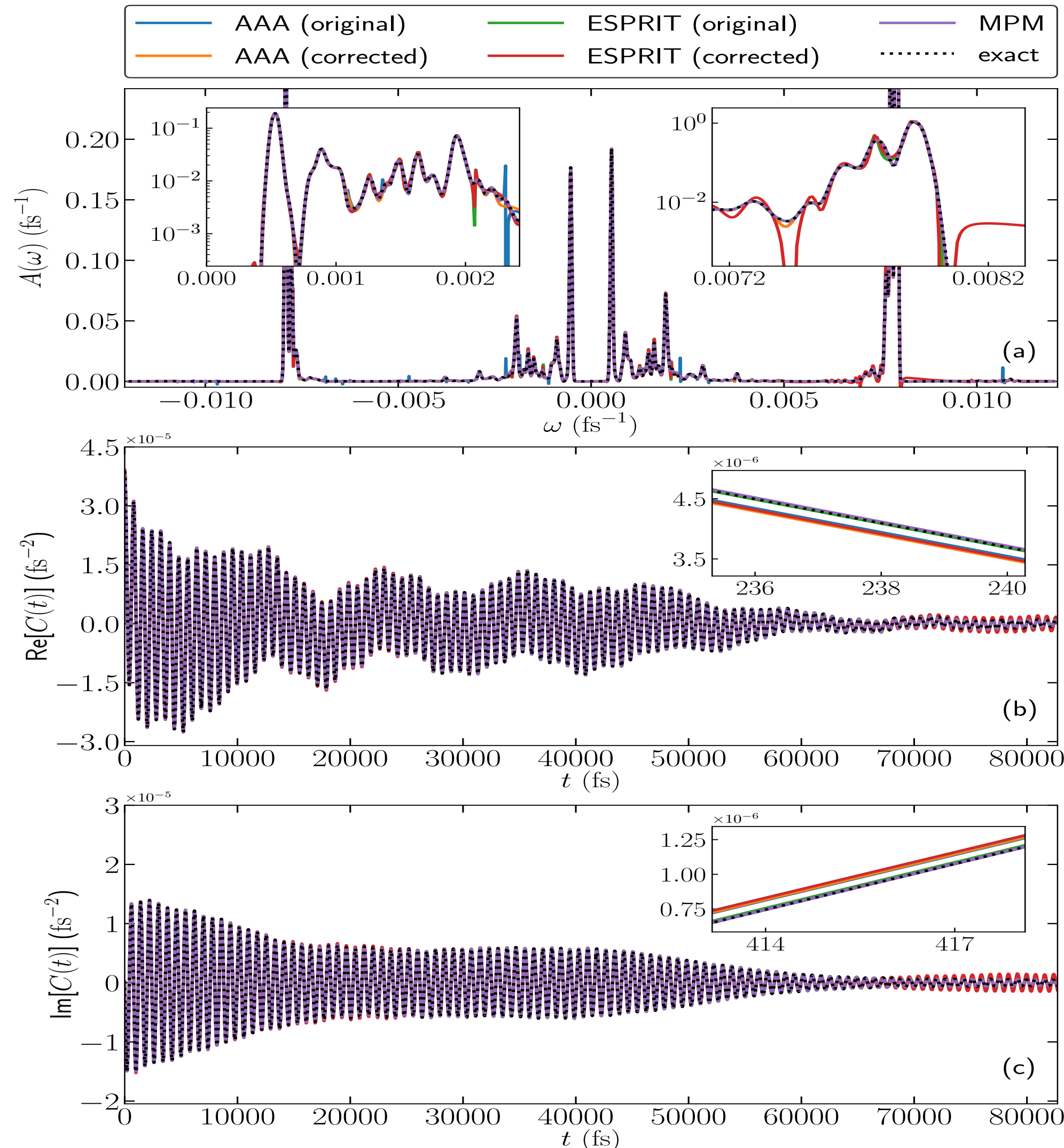
- AAA completely fails
- MPM (M2) has similar performance compared to ESPRIT
- MPM (M1) is the only one that has control over the long tail

Reason: AAA approximates spectral shapes,
MPM approximates spectral moments

Results – BCF (Structured spectral function)

data from exciton–phonon couplings in quantum dots

Nano Lett. 21, 8741–8748 (2021)
npj Comput. Mater. 9, 145 (2023)



Input: $\{J(\omega_i), \omega_i\}$ on a uniform grid

Preprocessing: perform AAA with $\varepsilon = 10^{-8}$, leading to 704 poles,
served as the exact solution

Issues: 1. ESPRIT recovers divergent exponentials

2. AAA recovers some real poles

Results: MPM remains stable and has better performance

Code

Thanks!



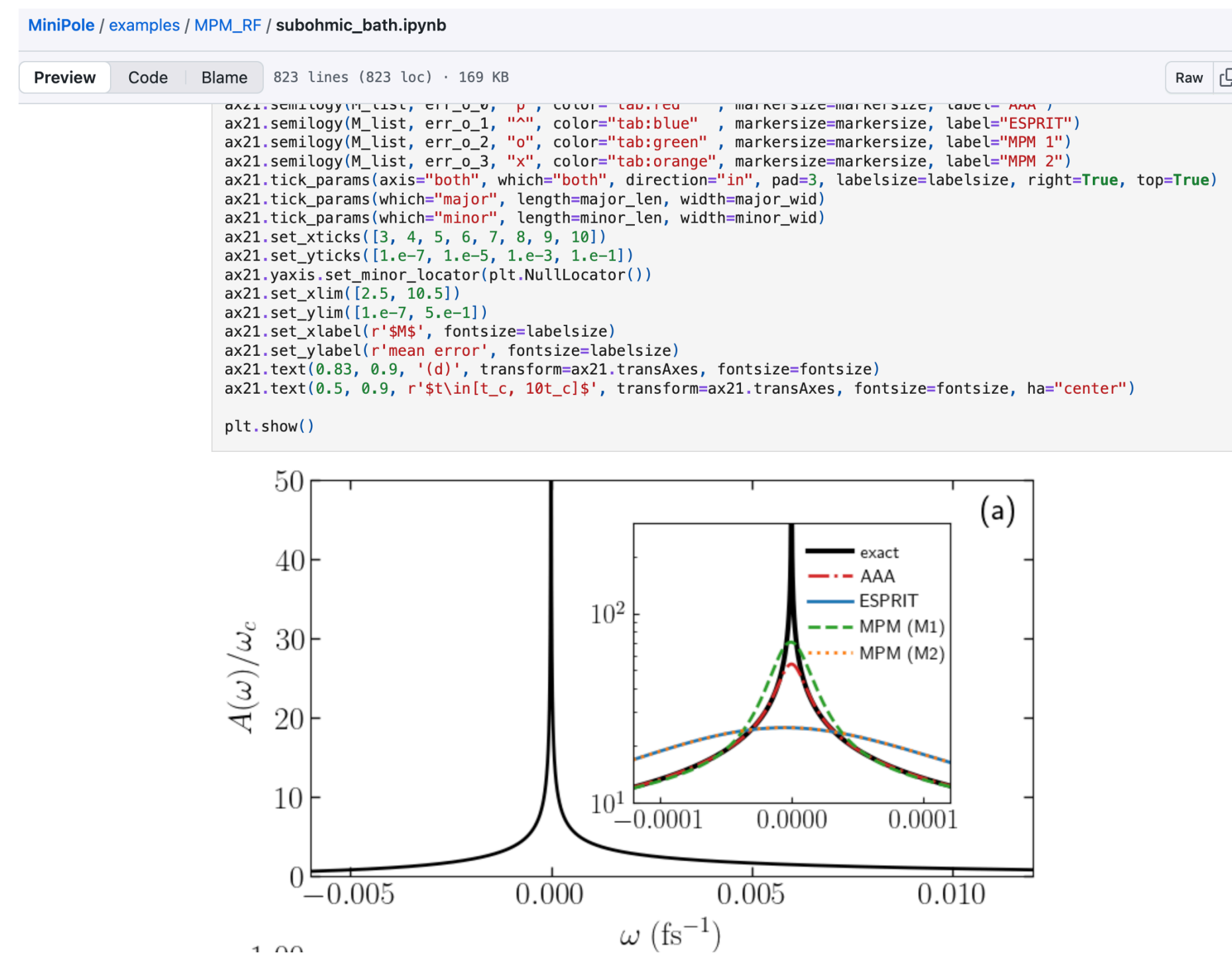
Green Software Package

<https://github.com/Green-Phys/MiniPole>
(also available: pypi, zenodo)



pip install mini_pole

MiniPoleRf: standard real-frequency MPM
MiniPoleRfDPR: real-frequency MPM combined with AAA



If get stuck, feel free to ask me for help!