

STUDYING THE TWO-PHOTON INTERFERENCE

*“A photon interferes only with
itself” - Dirac*

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AGENDA

- 1 Interference
- 2 Two-photon entangled state
- 3 Temporal two-photon interference
- 4 Gaussian-shape interference
- 5 Conclusions
- 6 Perspectives



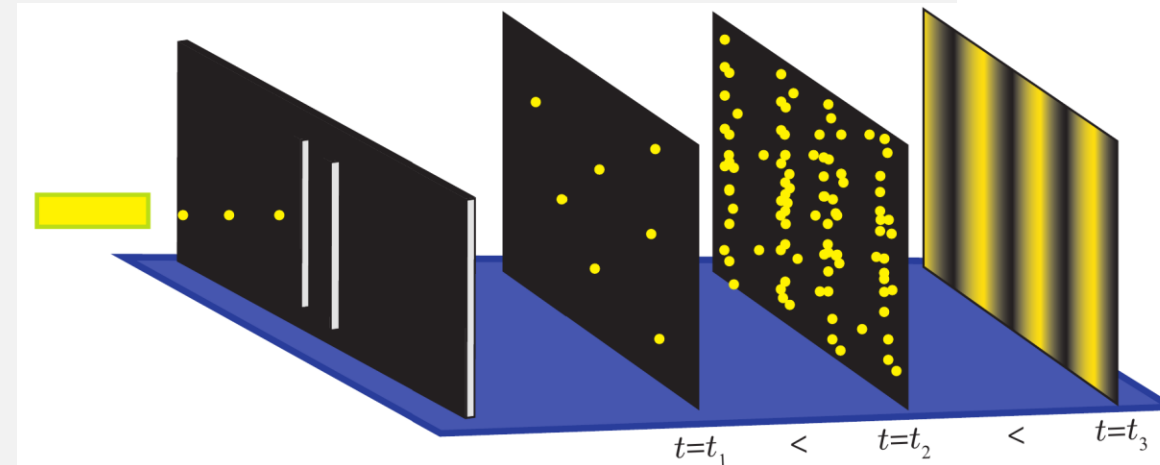
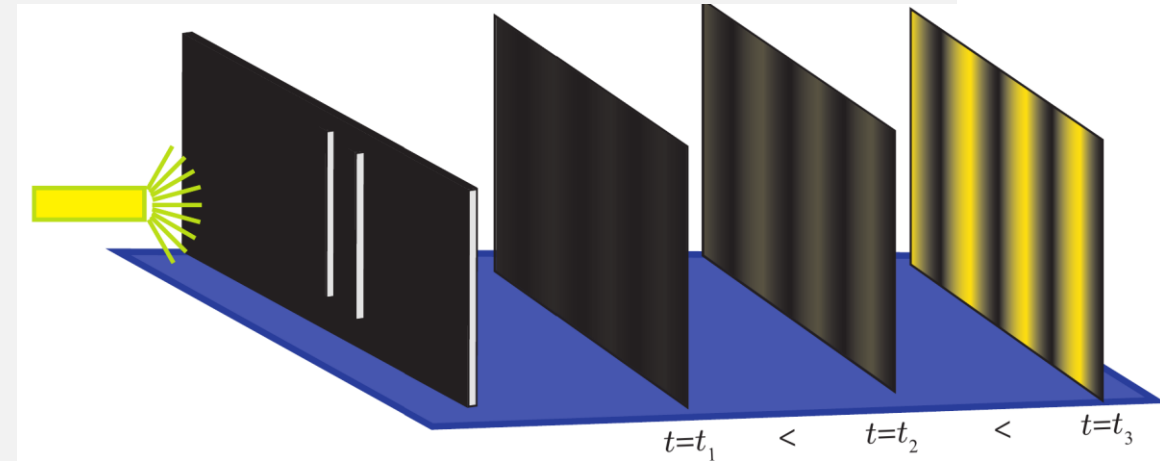
INTERFERENCE

correlations

CORRELATION FUNCTION

Coherence is the degree of order in a random field.

- In the case of random fields, interference occurs to the extent that the fields at two different space-time points are mutually correlated.
- Interference is a periodic variation of the intensity as a function of a parameter. What conditions? (coherent, indistinguishable)
- The nature and the degree of this correlation is described through what is known as a correlation function.



Coherence and Quantum Entanglement lectures.
Anand Kumar Jha

CORRELATION FUNCTION

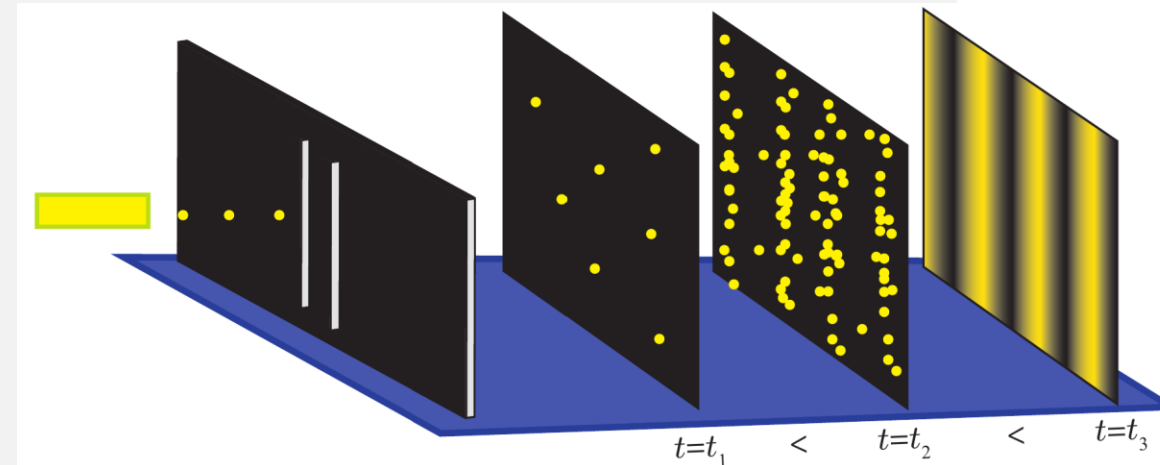
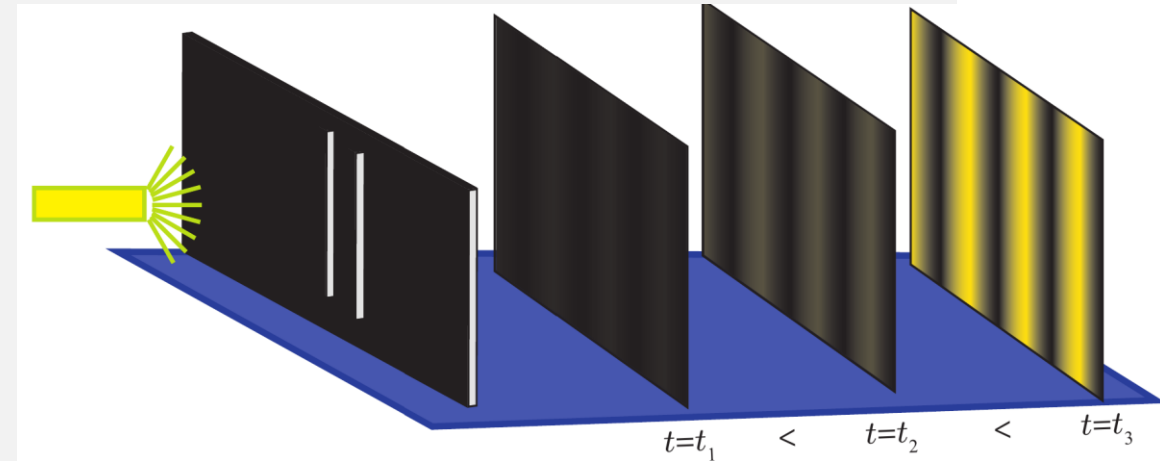
Coherence is the degree of order in a random field.

Two interfering field-amplitudes in the classical theory and the two interfering wave-functions in the quantum theory need to be mutually coherent for the interference to take place.

Physical amplitude \rightarrow Probability amplitude

Types of light interference

- Spatial
- Polarization
- Temporal



Coherence and Quantum Entanglement lectures.
Anand Kumar Jha

QUANTIFYING TEMPORAL CORRELATIONS

Analogy with classical mechanics

Classical

$$V(\mathbf{r}, t) = k_1 V(\mathbf{r}_0, t - t_1) + k_2 V(\mathbf{r}_0, t - t_2).$$

$$I(\mathbf{r}, t) = \langle V^*(\mathbf{r}, t) V(\mathbf{r}, t) \rangle$$

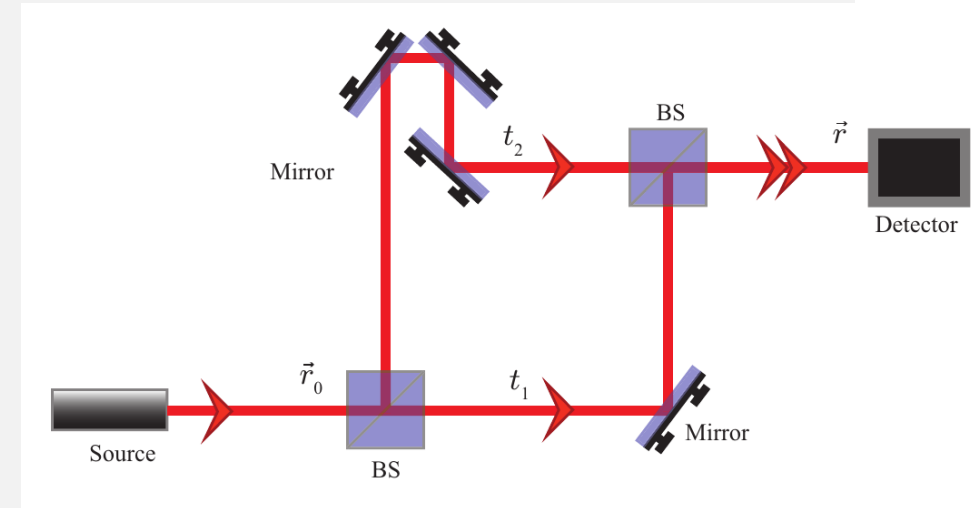
$$= |k_1|^2 I(t - t_1) + |k_2|^2 I(t - t_2) + k_1^* k_2 \Gamma(t - t_1, t - t_2) + \text{c.c.}$$

Quantum

$$\hat{E}^{(+)}(\mathbf{r}, t) = k_1 \hat{E}^{(+)}(\mathbf{r}_0, t - t_1) + k_2 \hat{E}^{(+)}(\mathbf{r}_0, t - t_2)$$

$$P(\mathbf{r}, t) = \langle \psi | \hat{E}^{(-)}(\mathbf{r}, t) \hat{E}^{(+)}(\mathbf{r}, t) | \psi \rangle$$

$$= |k_1|^2 P(\mathbf{r}_0, t - t_1) + |k_2|^2 P(\mathbf{r}_0, t - t_2) + 2|k_1||k_2|G^{(1)}(\mathbf{r}_0, t - t_1; \mathbf{r}_0, t - t_2)$$

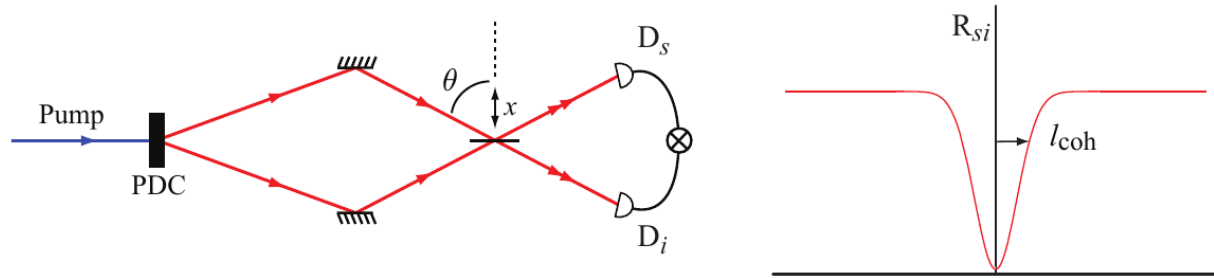


Coherence and Quantum Entanglement lectures.
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Γ
 $G^{(1)}$ \Rightarrow Correlation
function

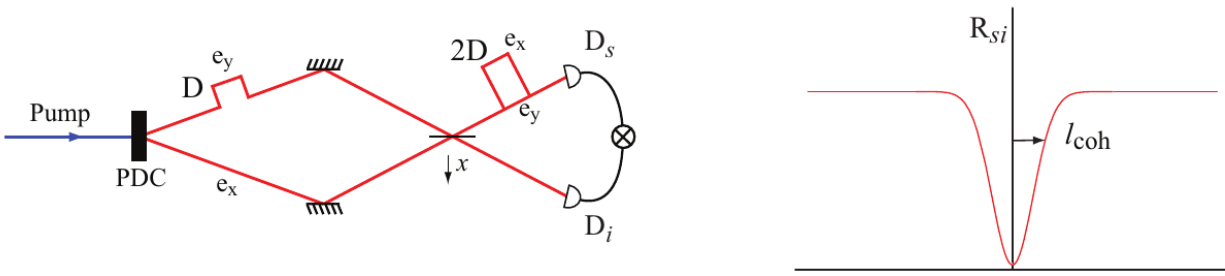
FAMOUS TWO-PHOTON INTERFERENCE

Hong-Ou-Mandel effect



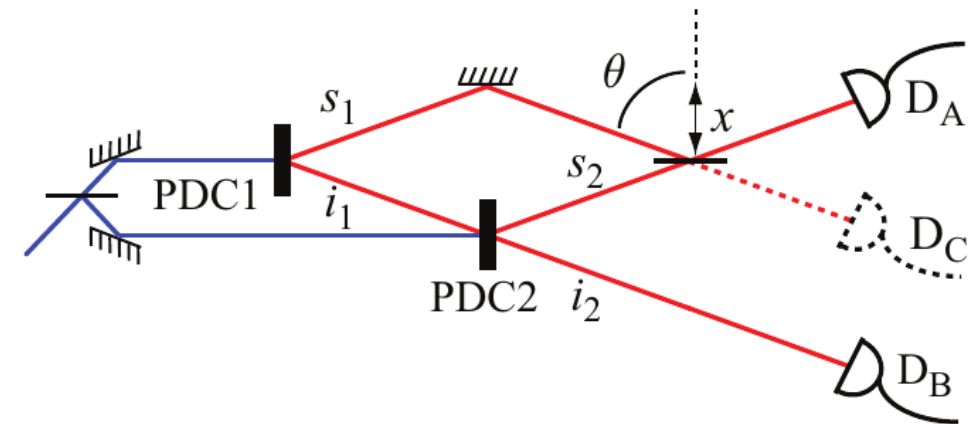
PRL 59, 2044 (1987)

Postponed compensation



PRL 77, 1917 (1996)

Induced coherence



PRL 67, 318 (1991)

IMPORTANCE OF TWO-PHOTON INTERFERENCE

Quantum Spectroscopy and sensing

- Enables super-resolution imaging & Lidar.
LSA 13:163 (2024)
- Used in astronomical measurements, direct imaging of black hole accretion discs. OE 31, 26, 44246 (2023)

Applications in OCT & LiDAR

- Enhances quantum optical coherence tomography.
PRR 5, 023170 (2023)
- Improves depth sensing in Quantum LiDAR.
PRL 131, 033603 (2023)

Quantum Information & Computing

- Key for entanglement & Linear Optical Quantum Computing (LOQC).
JH APL 25, 2 (2004)
- Essential for Boson Sampling
PRA 104, 032204 (2021)

Quantum Communication & Cryptography

- Secures Quantum Key Distribution (QKD).
arXiv 2411.07884 (2024)
- Fundamental for quantum teleportation & Bell tests.
OE 15, 16, 10188 (2007)

ENTANGLED TWO-PHOTON STATE

*Spontaneous
Parametric Down
Conversion*

SPDC

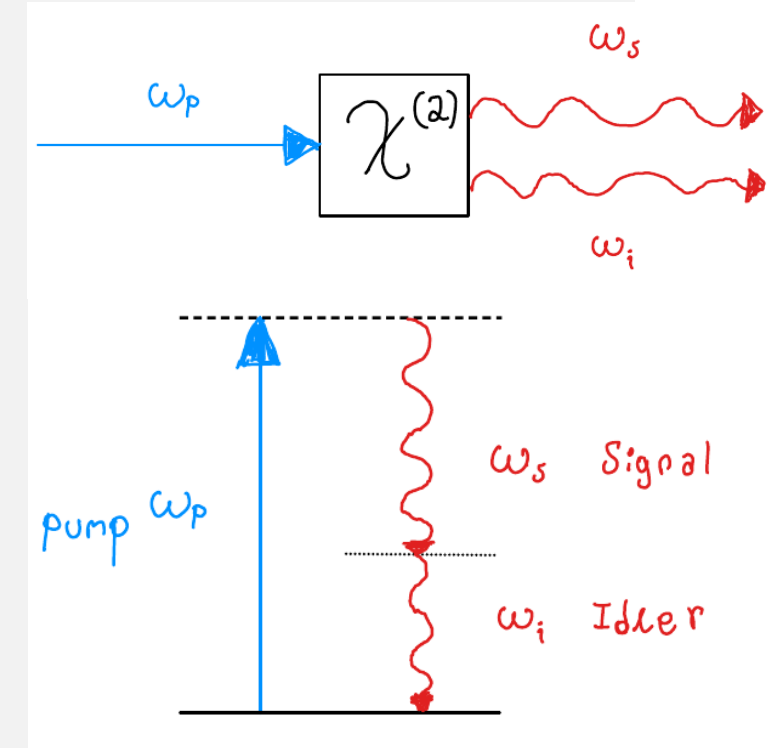
Process able to produce two-photon entangled states

- Inherently random
- Due to second-order effects in a non-linear environment
- Generated photon pairs can be entangled in various degrees of freedom

$$\hat{H}(t) = \epsilon_0 \int_{-L}^0 dz \chi^{(2)} \hat{E}_p^{(+)}(z, t) \hat{E}_s^{(-)}(z, t) \hat{E}_i^{(-)}(z, t) + \text{H.c.}$$

$$\hat{E}_p^{(+)}(z, t) = \int_0^\infty A_p d\omega_p V(\omega_p) e^{i[k_p z(\omega_p)z - \omega_p t]} e^{i(\omega_p \tau_p + \phi_p)},$$

$$\hat{E}_s^{(-)}(z, t) = \int_0^\infty A_s^* d\omega_s \hat{a}_s^\dagger(\omega_s) e^{i[\omega_s t - k_s z(\omega_s)z]}.$$



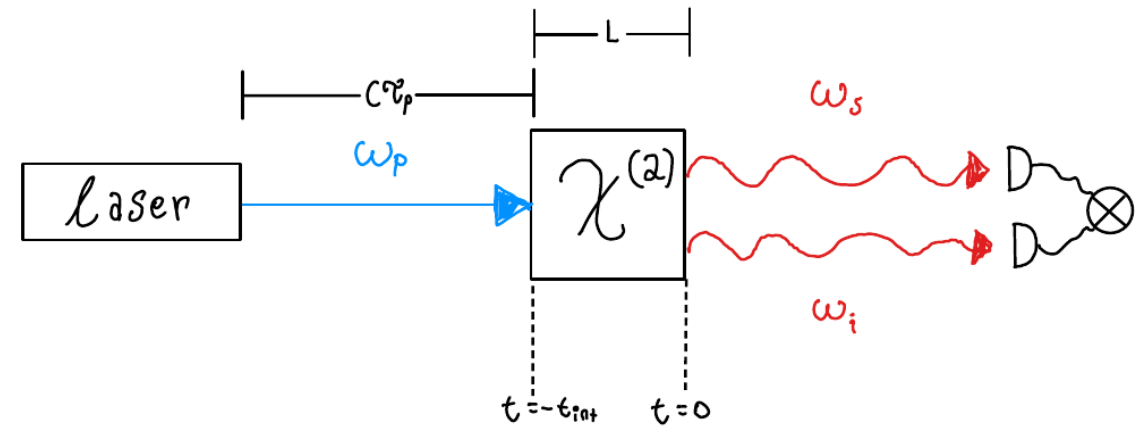
TIME EVOLUTION

Obtaining an entangled state

$$\hat{H}(t) = \epsilon_0 \int_{-L}^0 dz \chi^{(2)} \hat{E}_p^{(+)}(z, t) \hat{E}_s^{(-)}(z, t) \hat{E}_i^{(-)}(z, t) + \text{H.c.}$$

$$|\psi(-t_{\text{int}})\rangle = |\text{vac}\rangle_s |\text{vac}\rangle_i$$

$$|\psi(0)\rangle = \exp \left[\frac{1}{i\hbar} \int_{-t_{\text{int}}}^0 dt \hat{H}(t) \right] |\psi(-t_{\text{int}})\rangle$$



$$|\psi_{t_p}\rangle = A \int \int_0^\infty d\omega_s d\omega_i V(\omega_s + \omega_i) \Phi(\omega_s, \omega_i) e^{i[(\omega_s + \omega_i)\tau_p + \phi_p]} |\omega_s\rangle_s |\omega_i\rangle_i,$$

$$\Phi(\omega_s, \omega_i) = \int_{-L}^0 dz e^{i[k_{pz}(\omega_p) - k_{sz}(\omega_s) - k_{iz}(\omega_i)]z}.$$

ENTANGLED STATES

Phase matching function and entanglement

$$|\psi_{t_p}\rangle = A \int \int_0^\infty d\omega_s d\omega_i V(\omega_s + \omega_i) \Phi(\omega_s, \omega_i) e^{i[(\omega_s + \omega_i)\tau_p + \phi_p]} |\omega_s\rangle_s |\omega_i\rangle_i$$

$$\omega_p = \omega_s + \omega_i \quad \text{and} \quad \omega_d = \frac{\omega_s - \omega_i}{2} \quad \text{such that} \quad d\omega_s d\omega_i \rightarrow d\omega_p d\omega_d.$$

$$|\psi_{t_p}\rangle = A \int \int_0^\infty d\omega_p d\omega_d V(\omega_p) \Phi\left(\frac{\omega_p}{2} + \omega_d, \frac{\omega_p}{2} - \omega_d\right) e^{i(\omega_p \tau_p + \phi_p)} \left|\frac{\omega_p}{2} + \omega_d\right\rangle_s \left|\frac{\omega_p}{2} - \omega_d\right\rangle_i$$

$$V(\omega_p) = V_0 \delta(\omega_0 - \omega_p)$$

$$|\psi_{t_p}\rangle = AV_0 e^{i(\omega_0 \tau_p + \phi_p)} \int_0^\infty d\omega_d \Phi(\omega_d) \left|\frac{\omega_0}{2} + \omega_d\right\rangle_s \left|\frac{\omega_0}{2} - \omega_d\right\rangle_i$$



TEMPORAL TWO-PHOTON INTERFERENCE

*Coincidence
detection*

SUPERPOSITION OF TWO-PHOTON STATES

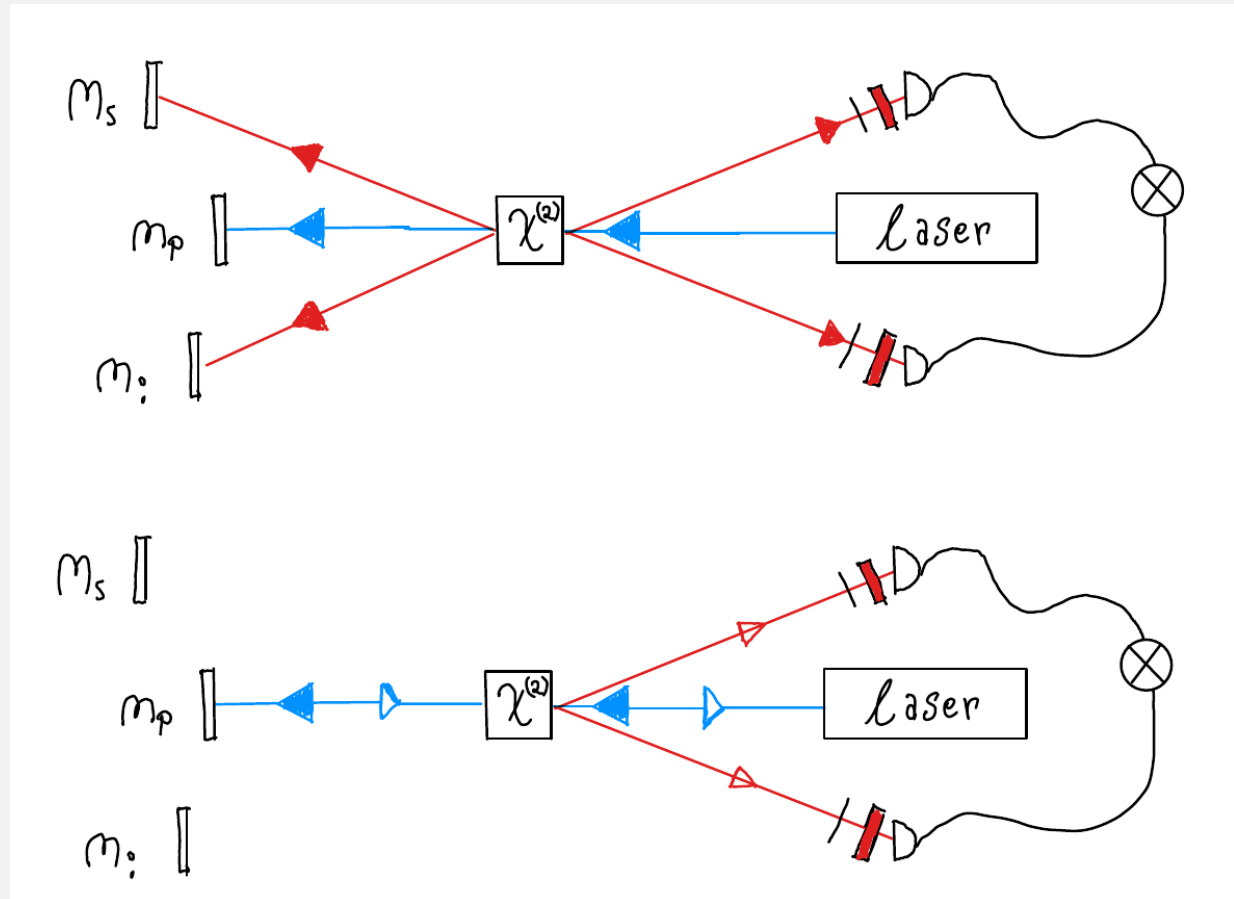
Interference of two possibilities

$$|\psi_{tpa}\rangle = A \int \int_0^\infty d\omega_p d\omega_d V_a(\omega_p) \Phi_a \left(\frac{\omega_p}{2} + \omega_d, \frac{\omega_p}{2} - \omega_d \right) e^{i(\omega_p \tau_p + \phi_p)} \left| \frac{\omega_p}{2} + \omega_d \right\rangle_{sa} \left| \frac{\omega_p}{2} - \omega_d \right\rangle_{ia}$$

$$|\psi\rangle = \sum_{a=1}^2 |\psi_{tpa}\rangle = |\psi_{tp1}\rangle + |\psi_{tp2}\rangle.$$

- Perfect spatial coherence
- Quasi-monochromatic approximation

$$\omega_0 \gg \Delta\omega$$



FIELD AT THE DETECTORS

Field operators

$$R_{si}(t, t + \tau) = \alpha_s \alpha_i \langle \psi | \hat{E}_s^{(-)}(t) \hat{E}_i^{(-)}(t + \tau) \hat{E}_i^{(+)}(t + \tau) \hat{E}_s^{(+)}(t) | \psi \rangle$$

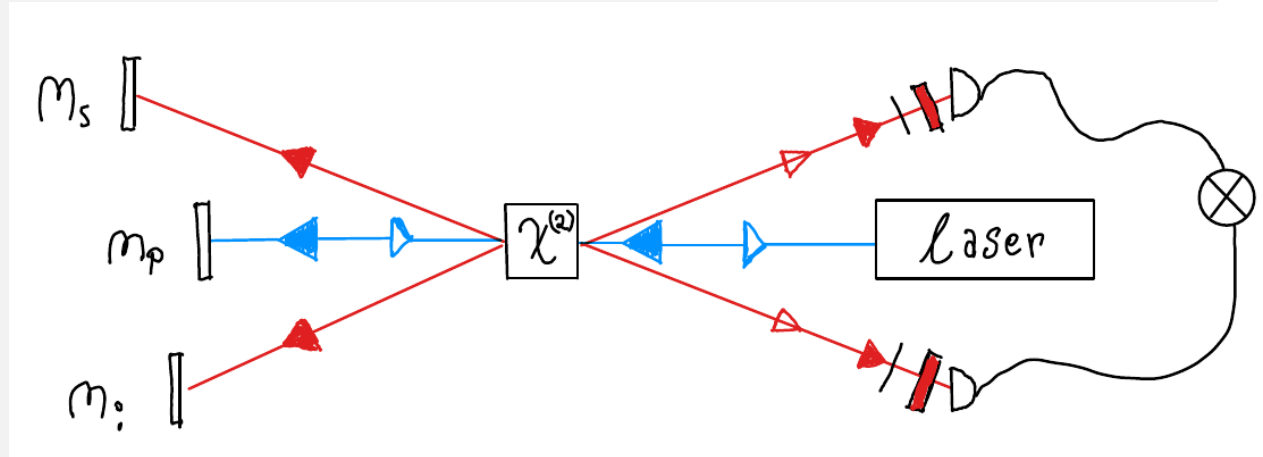
- α_i, α_s Quantum efficiency of detectors
- Amplitude filters f

$$|\psi\rangle = |\psi_{tp1}\rangle + |\psi_{tp2}\rangle$$

$$\hat{E}_s^{(+)}(t) = \hat{E}_{s1}^{(+)}(t - \tau_{s1}) + \hat{E}_{s2}^{(+)}(t - \tau_{s2})$$

$$= c_{s1} e^{i\phi_{s1}} \int_0^\infty d\omega_s f_s(\omega_s - \omega_{s0}) e^{-i\omega_s(t - \tau_{s1})} \hat{a}_{s1}(\omega_s)$$

$$+ c_{s2} e^{i\phi_{s2}} \int_0^\infty d\omega_s f_s(\omega_s - \omega_{s0}) e^{-i\omega_s(t - \tau_{s2})} \hat{a}_{s2}(\omega_s).$$



COINCIDENCES DETECTION

What is the probability to detect a coincidence in a given window of time?

$$R_{si}(t, t + \tau) = \alpha_s \alpha_i \langle \psi | \hat{E}_s^{(-)}(t) \hat{E}_i^{(-)}(t + \tau) \hat{E}_i^{(+)}(t + \tau) \hat{E}_s^{(+)}(t) | \psi \rangle \rightarrow 16 \text{ terms}$$

$$\begin{aligned} R_{si}(t, t + \tau) = & \alpha_s \alpha_i \langle \psi_{tp1} | \hat{E}_{s1}^{(-)} \hat{E}_{i1}^{(-)} \hat{E}_{i1}^{(+)} \hat{E}_{s1}^{(+)} | \psi_{tp1} \rangle + \alpha_s \alpha_i \langle \psi_{tp2} | \hat{E}_{s2}^{(-)} \hat{E}_{i2}^{(-)} \hat{E}_{i2}^{(+)} \hat{E}_{s1}^{(+)} | \psi_{tp2} \rangle \\ & + \alpha_s \alpha_i \langle \psi_{tp1} | \hat{E}_{s1}^{(-)} \hat{E}_{i1}^{(-)} \hat{E}_{i2}^{(+)} \hat{E}_{s2}^{(+)} | \psi_{tp2} \rangle + \text{H.c.} \end{aligned}$$

$$\equiv R_{si}^{11}(t, t + \tau) + R_{si}^{22}(t, t + \tau) + R_{si}^{12}(t, t + \tau) + \text{H.c.}$$

$$\hat{E}_{sa}^{(\pm)} = \hat{E}_{sa}^{(\pm)}(t - \tau_{s2})$$

$$\hat{E}_{ia}^{(\pm)} = \hat{E}_{ia}^{(\pm)}(t + \tau - \tau_{i2})$$

FOURTH-POINT CORRELATION FUNCTION

Finding correlation

$$\hat{E}_{i2}^{(+)}(t + \tau - \tau_{i2}) \hat{E}_{s2}^{(+)}(t - \tau_{s2}) |\psi_{tp2}\rangle = A c_2 e^{i\phi_2} e^{-i\omega_0(t + \frac{\tau}{2} - \tau_2)} e^{i\omega_{d0}(\tau + \tau'_2)} \\ \times v_2 \left(t + \frac{\tau}{2} - \tau_2 \right) g_2^*(\tau + \tau'_2) |\text{vac}\rangle_{s2} |\text{vac}\rangle_{i2}$$

$$\tau_2 = \tau_{p2} + \frac{\tau_{s2} + \tau_{i2}}{2}, \quad \text{and} \quad \tau'_2 = \tau_{s2} - \tau_{i2}$$

$$g_2^*(t) \equiv \int_{-\infty}^{\infty} d\omega'_d \Phi_2(\omega'_d + \omega_{s0}, -\omega'_d + \omega_{i0}) f_s(\omega'_d) f_i(-\omega'_d) e^{-i\omega'_d t} \quad v_2(t) \equiv \int_{-\infty}^{\infty} d\omega'_p V_2(\omega'_p + \omega_0) e^{-i\omega'_p t}$$

$$\omega_p = \omega'_p + \omega_0 \quad \text{and} \quad \omega_d = \omega'_d + \omega_{d0}, \quad \text{where} \quad \omega_{d0} = \frac{\omega_{s0} - \omega_{i0}}{2}$$

$$R_{si}^{12}(t, t+\tau) = K c_1^* c_2 \left\langle e^{-i(\omega_0 \Delta\tau + \omega_{d0} \Delta\tau' + \Delta\phi)} v_1^* \left(t + \frac{\tau}{2} - \tau_1 \right) v_2 \left(t + \frac{\tau}{2} - \tau_2 \right) g_1(\tau + \tau'_1) g_2^*(\tau + \tau'_2) \right\rangle_{t, \tau}$$

FOURTH-POINT CORRELATION FUNCTION

Finding correlation

Detection of coincidences

$$R_{si}^{12}(t, t+\tau) = K c_1^* c_2 \left\langle e^{-i(\omega_0 \Delta\tau + \omega_{d0} \Delta\tau' + \Delta\phi)} v_1^* \left(t + \frac{\tau}{2} - \tau_1\right) v_2 \left(t + \frac{\tau}{2} - \tau_2\right) g_1(\tau + \tau'_1) g_2^*(\tau + \tau'_2) \right\rangle_{t, \tau}$$

$$R_{si}^{12} = K c_1^* c_2 e^{-i(\omega_0 \Delta\tau + \omega_{d0} \Delta\tau' + \Delta\phi)} \sqrt{|v_1|^2 |v_2|^2} \sqrt{|g_1|^2 |g_2|^2} \gamma(\Delta\tau) \gamma(\Delta\tau')$$

$$R_{si} = R_1 + R_2 + 2\sqrt{R_1 R_2} \gamma(\Delta\tau) \gamma'(\Delta\tau') \cos(\omega_0 \Delta\tau + \omega_{d0} \Delta\tau' + \Delta\phi)$$

$$R_1 = K |c_1 g_1 v_1|^2 \quad \text{and} \quad R_2 = K |c_2 g_2 v_2|^2$$

$$R_{si} = R_1 + R_2 + 2\sqrt{R_1 R_2} \gamma(\Delta L) \gamma'(\Delta L') \cos(k_0 \Delta L + k_{d0} \Delta L' + \Delta\phi)$$

FOURTH-POINT CORRELATION FUNCTION

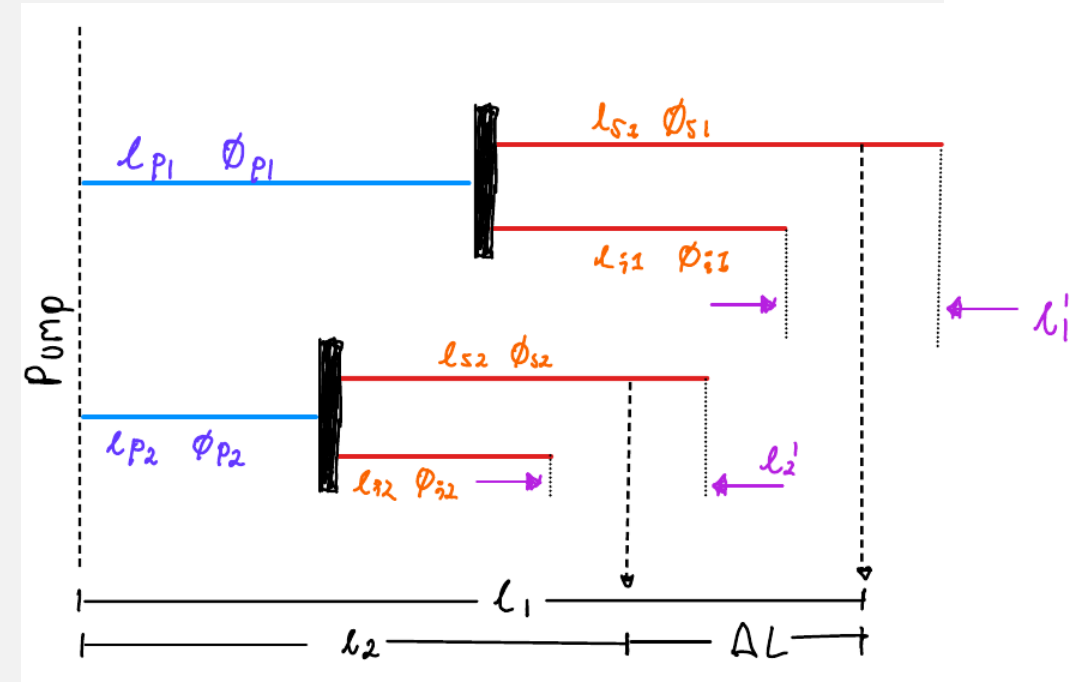
Defining geometrical parameters

$$R_{si} = R_1 + R_2 + 2\sqrt{R_1 R_2} \gamma(\Delta L) \gamma'(\Delta L') \cos(k_0 \Delta L + k_{d0} \Delta L' + \Delta \phi)$$

$$\Delta L \equiv l_1 - l_2 = \left(\frac{l_{s1} + l_{i1}}{2} + l_{p1} \right) - \left(\frac{l_{s2} + l_{i2}}{2} + l_{p2} \right)$$

$$\Delta L' \equiv l'_1 - l'_2 = (l_{s1} - l_{i1}) - (l_{s2} - l_{i2})$$

$$\Delta \phi \equiv \phi_1 - \phi_2 = (\phi_{s1} + \phi_{i1} + \phi_{p1}) - (\phi_{s2} + \phi_{i2} + \phi_{p2})$$



GAUSSIAN-SHAPE CASE

*Shape of interference
pattern*

ASSUMING A GAUSSIAN SPECTRAL DISTRIBUTION

$$R_{si} = R_1 + R_2 + 2\sqrt{R_1 R_2} \gamma(\Delta L) \gamma'(\Delta L') \cos(k_0 \Delta L + k_{d0} \Delta L' + \Delta \phi)$$

$$k_0 = \frac{\omega_0}{c} \quad \text{and} \quad k_{d0} = \frac{\omega_{d0}}{c} \quad \rightarrow \quad \omega_s = \omega_i$$

$$R_{si} = C[1 + \gamma'(\Delta L') \gamma(\Delta L) \cos(k_0 \Delta L + \Delta \phi)].$$

- when the pump is a stationary (continuous-wave) field having a Gaussian spectrum of rms frequency width $\Delta\omega$, the time-averaged degree of correlation of the pump field

$$\gamma(\Delta L) = \exp \left[-\frac{1}{2} \left(\frac{\Delta L}{l_{\text{coh}}^p} \right)^2 \right]. \quad l_{\text{coh}}^p = \frac{c}{\Delta\omega_p}$$

- In situations in which the signal-idler field has a Gaussian spectrum of width $\Delta\omega$,

$$\gamma'(\Delta L') = \exp \left[-\frac{1}{2} \left(\frac{\Delta L'}{l_{\text{coh}}} \right)^2 \right], \quad l_{\text{coh}} = \frac{c}{\Delta\omega}$$

$$\begin{aligned} \Delta L &\ll l_{\text{coh}}^p \\ \Delta L' &\ll l_{\text{coh}} \end{aligned}$$

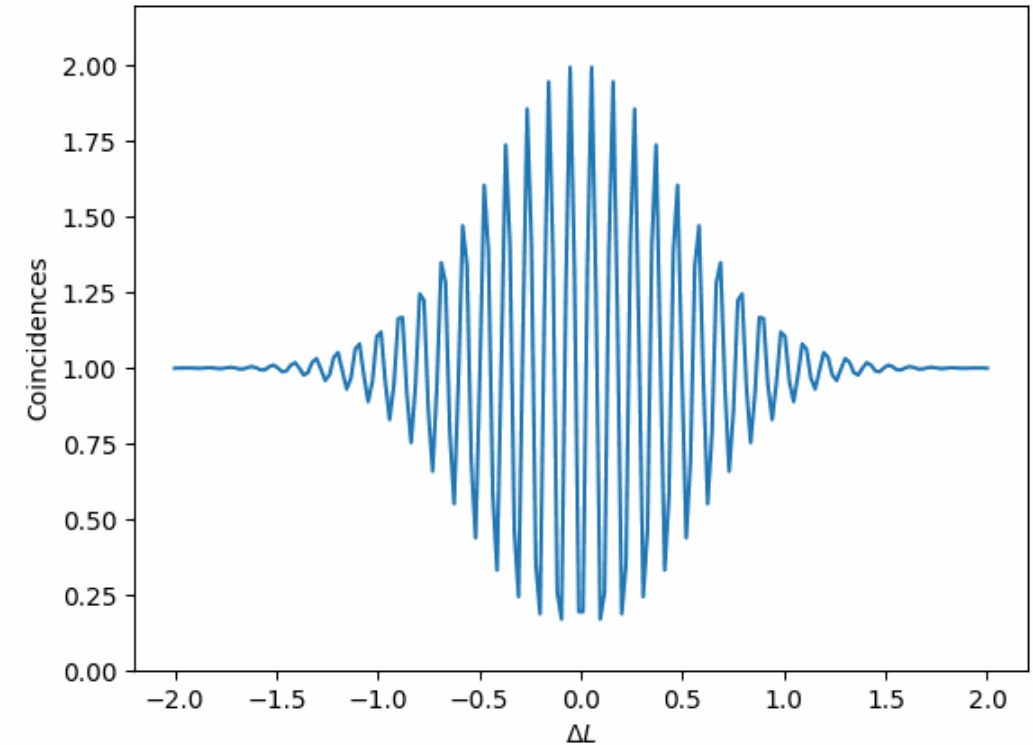
LIMITING CASES

$$\Delta L' = 0 \quad \Delta \phi = 0$$

$$R_{si} = C [1 + \gamma(\Delta L) \cos(k_0 \Delta L)] .$$

$$\gamma(\Delta L) = \exp \left[-\frac{1}{2} \left(\frac{\Delta L}{l_{\text{coh}}^p} \right)^2 \right] .$$

$$\begin{aligned} \Delta L &\ll l_{\text{coh}}^p \\ \Delta L' &\ll l_{\text{coh}} \end{aligned}$$



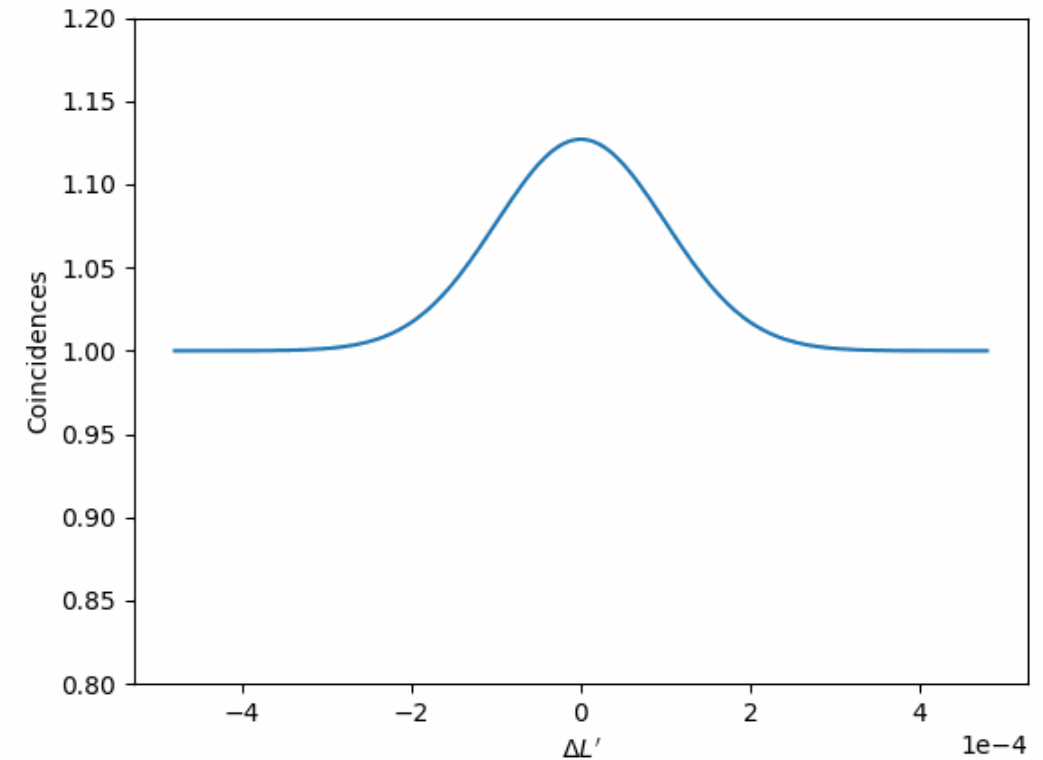
LIMITING CASES

ΔL $\Delta\phi$ Fixed

$$R_{si} = C [1 + K\gamma'(\Delta L')]$$

$$K = \gamma(\Delta L) \cos(k_0\Delta L + \Delta\phi)$$

- The coincidence count rate can show a dip when the two alternatives interfere destructively ($K < 0$)
- Hump when the two alternatives interfere constructively ($K > 0$)



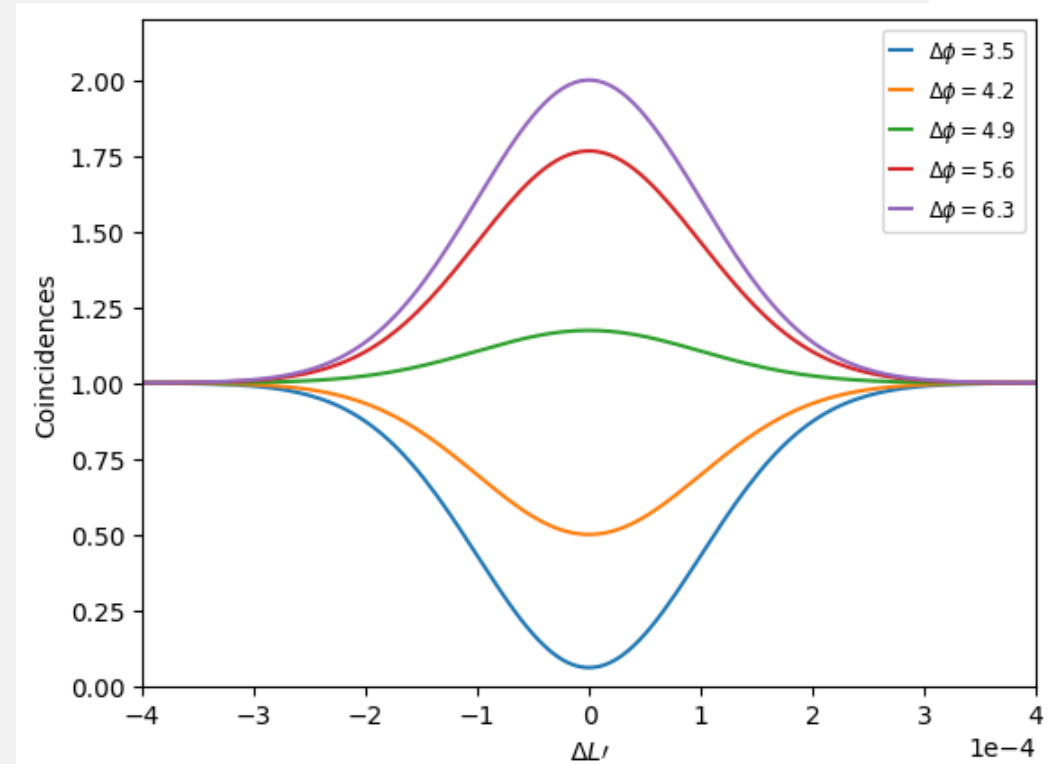
LIMITING CASES

ΔL $\Delta\phi$ Fixed

$$R_{si} = C [1 + K\gamma'(\Delta L')]$$

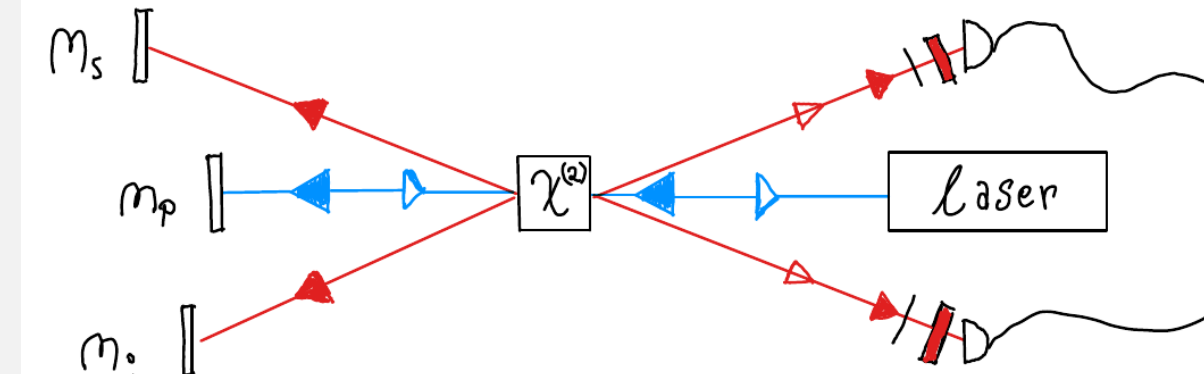
$$K = \gamma(\Delta L) \cos(k_0\Delta L + \Delta\phi)$$

- The coincidence count rate can show a dip when the two alternatives interfere destructively ($K < 0$)
- Hump when the two alternatives interfere constructively ($K > 0$)



CONCLUSIONS

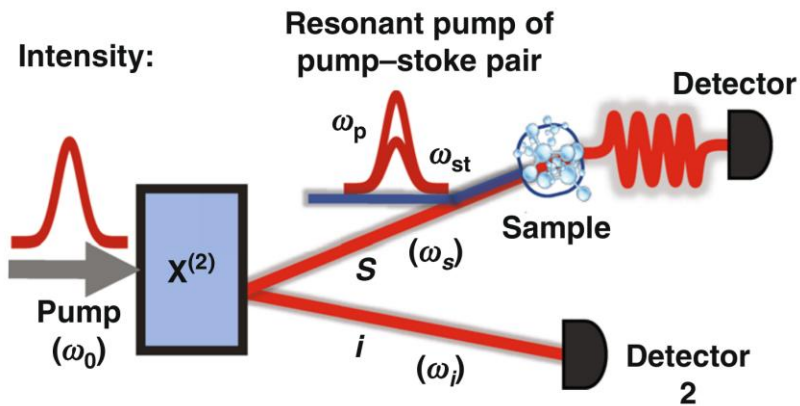
- Interference is a group effect due to interference of individual photon (two-photon).
- photon interference arises from the indistinguishability of quantum states and the superposition of probability amplitudes.
- With gaussian-shape any two-photon interference effect is a superposition of the two limiting cases presented.
- Spontaneous Parametric Down-Conversion (SPDC) is a crucial process for generating biphoton states with tunable spectral and spatial properties.
- This theory works for many different interferometer structure.
- Calculation of quantum correlations is not that simple.



PERSPECTIVES

spectroscopy

Quantum CARS



$$V(t) = \sum_{j=1}^N \sum_b \alpha_{bg}^{(j)} (|b\rangle\langle g|)_j E_s(t) E_s^\dagger(t) e^{i(\omega_b - \omega_g)t} + h.c.$$

Spectral and time resolution at same time

Zhang, Z., Peng, I., Nie, X. *et al.* Entangled photons enabled time-frequency-resolved coherent Raman spectroscopy and applications to electronic coherences at femtosecond scale. *Light Sci Appl* **11**, 274 (2022).

Absorption spectroscopy with quantum light with arbitrary spectral shape

Study two scheme with and without interference

- Advanced optical filtering technology at the telecom frequency can be leveraged to analyse absorption at different frequencies,
- The research includes theoretical analysis in the Heisenberg picture, establishing high-resolution spectroscopy using quantum interference.

Padilla Camargo, A. A. (2023). Absorption Spectroscopy with Quantum Light: Experimental and Theoretical Simulations Based on Arbitrary Spectral Shapes Generated with a Programmable Filter (Master's thesis). Universitat Politècnica de Catalunya.



THANK YOU

EXTRAS

*Spontaneous
Parametric Down
Conversion*

LIMITING CASES

Where the biphoton comes from

Company Name

- Explicar la física
- Mensaje
- Enque trabajo y motivacion
- Hice recapitulacion, que he estudiado
- Perspectivas con interes

SPDC

Process that produces entangled two-photon field.

Electromagnetic field in a non-linear medium

$$P(\mathbf{r}, t) = \epsilon_0 \chi^{(1)} E(\mathbf{r}, t) + \epsilon_0 \chi^{(2)} E^2(\mathbf{r}, t) + \dots$$

$$D(\mathbf{r}, t) = \epsilon_0 E(\mathbf{r}, t) + P(\mathbf{r}, t)$$

$$D(\mathbf{r}, t) \cdot E(\mathbf{r}, t) = [\epsilon_0 E(\mathbf{r}, t) + \epsilon_0 \chi^{(1)} E(\mathbf{r}, t)] \cdot E(\mathbf{r}, t) + \epsilon_0 \chi^{(2)} E^2(\mathbf{r}, t) \cdot E(\mathbf{r}, t) \dots$$

- Nulla a erat eget nunc hendrerit ultrices eu nec nulla. Donec viverra leo aliquet, auctor quam id, convallis orci.
- Sed in molestie est. Cras ornare turpis at ligula posuere, sit amet accumsan neque lobortis.
- Maecenas mattis risus ligula, sed ullamcorper nunc efficitur sed.

Ambiente
Motivacion
Interference

SPDC

Process that produces entangled two-photon field.

Electromagnetic field in a non-linear medium

$$D(\mathbf{r}, t) \cdot E(\mathbf{r}, t) = [\epsilon_0 E(\mathbf{r}, t) + \epsilon_0 \chi^{(1)} E(\mathbf{r}, t)] \cdot E(\mathbf{r}, t) + \epsilon_0 \chi^{(2)} E^2(\mathbf{r}, t) \cdot E(\mathbf{r}, t) \cdots$$

$$H(t) = \int_V P^{(2)}(\mathbf{r}, t) E(\mathbf{r}, t) d^3 r = \epsilon_0 \int_V \chi^{(2)} E_p(\mathbf{r}, t) E_s(\mathbf{r}, t) E_i(\mathbf{r}, t) d^3 r$$

$$\hat{H}(t) = \epsilon_0 \int_V \chi^{(2)} \hat{E}_p(\mathbf{r}, t) \hat{E}_s(\mathbf{r}, t) \hat{E}_i(\mathbf{r}, t) d^3 r \quad \hat{E}_p(\mathbf{r}, t) = \hat{E}_p^{(+)}(\mathbf{r}, t) + \hat{E}_p^{(-)}(\mathbf{r}, t),$$

$$\hat{E}^{(+)}(\mathbf{r}, t) = \int A_{\mathbf{k}} \hat{a}_{\mathbf{k}}(t) e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} d^3 k.$$

$$\hat{H}(t) = \epsilon_0 \int_V d^3 r \chi^{(2)} \hat{E}_p^{(+)}(\mathbf{r}, t) \hat{E}_s^{(-)}(\mathbf{r}, t) \hat{E}_i^{(-)}(\mathbf{r}, t) + \text{H.c.}$$

SPDC

Process that produces entangled two-photon field.

Field In detectors

$$\begin{aligned}
 R_{si}(t, t+\tau) = & \alpha_s \alpha_i \langle \psi_{tp1} | \hat{E}_{s1}^{(-)}(t-\tau_{s1}) \hat{E}_{i1}^{(-)}(t+\tau-\tau_{i1}) \hat{E}_{i1}^{(+)}(t+\tau-\tau_{i1}) \hat{E}_{s1}^{(+)}(t-\tau_{s1}) | \psi_{tp1} \rangle \\
 & + \alpha_s \alpha_i \langle \psi_{tp2} | \hat{E}_{s2}^{(-)}(t-\tau_{s2}) \hat{E}_{i2}^{(-)}(t+\tau-\tau_{i2}) \hat{E}_{i2}^{(+)}(t+\tau-\tau_{i2}) \hat{E}_{s1}^{(+)}(t-\tau_{s1}) | \psi_{tp2} \rangle \\
 & + \alpha_s \alpha_i \langle \psi_{tp1} | \hat{E}_{s1}^{(-)}(t-\tau_{s1}) \hat{E}_{i1}^{(-)}(t+\tau-\tau_{i1}) \hat{E}_{i2}^{(+)}(t+\tau-\tau_{i2}) \hat{E}_{s2}^{(+)}(t-\tau_{s2}) | \psi_{tp2} \rangle + \text{H.c.}
 \end{aligned}$$

COINCIDENCES DETECTION

Process that produces entangled two-photon field.

$$R_{si}(t, t + \tau) = \alpha_s \alpha_i \langle \psi | \hat{E}_s^{(-)}(t) \hat{E}_i^{(-)}(t + \tau) \hat{E}_i^{(+)}(t + \tau) \hat{E}_s^{(+)}(t) | \psi \rangle \rightarrow 16 \text{ terms}$$

$$\begin{aligned} R_{si}(t, t + \tau) &= \alpha_s \alpha_i \langle \psi | [\hat{E}_{s1}^{(-)}(t - \tau_{s1}) + \hat{E}_{s2}^{(-)}(t - \tau_{s2})][\hat{E}_{i1}^{(-)}(t + \tau - \tau_{i1}) + \hat{E}_{i2}^{(-)}(t + \tau - \tau_{i2})] \\ &\quad \times [\hat{E}_{i1}^{(+)}(t + \tau - \tau_{i1}) + \hat{E}_{i2}^{(+)}(t + \tau - \tau_{i2})][\hat{E}_{s1}^{(+)}(t - \tau_{s1}) + \hat{E}_{s2}^{(+)}(t - \tau_{s2})] | \psi \rangle \\ &\equiv R_{si}^{11}(t, t + \tau) + R_{si}^{22}(t, t + \tau) + R_{si}^{12}(t, t + \tau) + \text{H.c.} \end{aligned}$$

SPDC

Process that produces entangled two-photon field.

Detection of coincidences

$$g_2^*(t) \equiv \int_{-\infty}^{\infty} d\omega'_d \Phi_2(\omega'_d + \omega_{s0}, -\omega'_d + \omega_{i0}) f_s(\omega'_d) f_i(-\omega'_d) e^{-i\omega'_d t}$$

$$v_2(t) \equiv \int_{-\infty}^{\infty} d\omega'_p V_2(\omega'_p + \omega_0) e^{-i\omega'_p t}$$

$$\begin{aligned} \hat{E}_{i2}^{(+)}(t + \tau - \tau_{i2}) \hat{E}_{s2}^{(+)}(t - \tau_{s2}) |\psi_{tp2}\rangle &= A c_2 e^{i\phi_2} e^{-i\omega_0(t + \frac{\tau}{2} - \tau_2)} e^{i\omega_{d0}(\tau + \tau'_2)} \\ &\times v_2\left(t + \frac{\tau}{2} - \tau_2\right) g_2^*(\tau + \tau'_2) |\text{vac}\rangle_{s2} |\text{vac}\rangle_{i2} \end{aligned}$$

$$R_{si}^{12}(t, t+\tau) = K c_1^* c_2 \left\langle e^{-i(\omega_0 \Delta\tau + \omega_{d0} \Delta\tau' + \Delta\phi)} v_1^* \left(t + \frac{\tau}{2} - \tau_1\right) v_2 \left(t + \frac{\tau}{2} - \tau_2\right) g_1(\tau + \tau'_1) g_2^*(\tau + \tau'_2) \right\rangle$$