



UNIVERSITÀ
DEGLI STUDI
DI PADOVA

Road to Fluxonium implementation

Quantum information and Computing

Cavazza Marco, Sardo Infirri Giosuè

April 13, 2024

1. A bit of theory
2. Implementation and results
3. Fluxonium
4. Conclusions
5. Bibliography

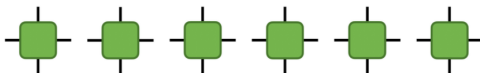
MPS:

$$|\psi\rangle = \sum_{\{\sigma_j\}, \{a_j\}} A_{a_1}^{[1]\sigma_1} A_{a_1 a_2}^{[2]\sigma_2} \dots A_{a_{N-2} a_{N-1}}^{[N-1]\sigma_{N-1}} A_{a_{N-1}}^{[N]\sigma_N} |\sigma_1 \sigma_2 \dots \sigma_N\rangle$$



MPO:

$$H = \sum_{\sigma'_1 \dots \sigma'_N} \sum_{\sigma_1 \dots \sigma_N} \sum_{b_1 \dots b_{N-1}} H_{b_1}^{[\sigma'_1 \sigma_1]} H_{b_1 b_2}^{[\sigma'_2 \sigma_2]} \dots H_{b_{N-1}}^{[\sigma'_N \sigma_N]} |\sigma'_1 \sigma'_2 \dots \sigma'_N\rangle \langle \sigma_1 \sigma_2 \dots \sigma_N|$$



How to build the MPO (1)

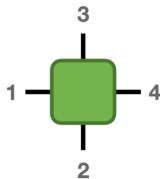


Ising Hamiltonian for a chain of spins:

$$\hat{H}_{\text{Ising}}^{\text{tot}} = -J \sum_{\langle i,j \rangle} \hat{S}_i^z \hat{S}_j^z - H \sum_i \hat{S}_i^x$$

Single tensor in the MPO:

$$H_{b_j b_{j+1}}^{[\sigma'_i \sigma_i]} = \begin{pmatrix} \hat{\mathbb{I}} & \hat{S}_z & -H \hat{S}_x \\ \hat{\sigma} & \hat{\sigma} & -J \hat{S}_z \\ \hat{\sigma} & \hat{\sigma} & \hat{\mathbb{I}} \end{pmatrix}$$



- 1 and 4: bond dim. (χ).
- 2 and 3: physical dim. (d).
- rank 4: ($\chi \times \chi \times d \times d$).

How to build the MPO (2)



In the case of Open Boundary Condition (OBC) the first and the last tensor need to be contracted as shown below:

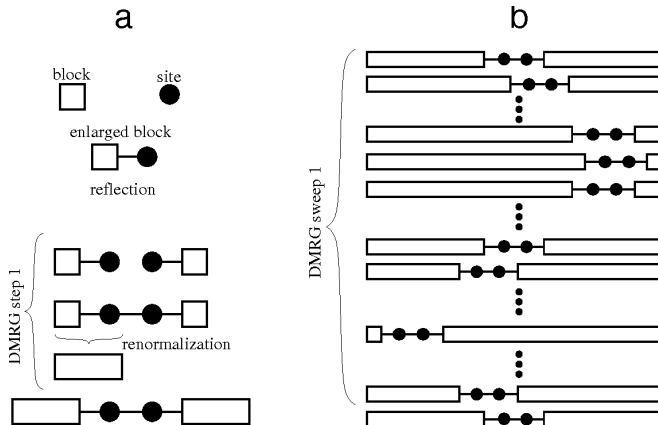
$$(1 \quad 0 \quad 0) \begin{pmatrix} \hat{\mathbb{I}} & \hat{S}_z & -H\hat{S}_x \\ \hat{O} & \hat{O} & -J\hat{S}_z \\ \hat{O} & \hat{O} & \hat{\mathbb{I}} \end{pmatrix} \begin{pmatrix} \hat{\mathbb{I}} & \hat{S}_z & -H\hat{S}_x \\ \hat{O} & \hat{O} & -J\hat{S}_z \\ \hat{O} & \hat{O} & \hat{\mathbb{I}} \end{pmatrix} \dots \begin{pmatrix} \hat{\mathbb{I}} & \hat{S}_z & -H\hat{S}_x \\ \hat{O} & \hat{O} & -J\hat{S}_z \\ \hat{O} & \hat{O} & \hat{\mathbb{I}} \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

The product of the matrices should yield the exact Hamiltonian of the system

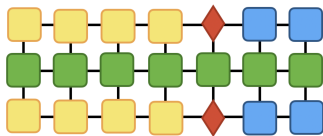
Infinite vs Finite DMRG



Infinite(a) vs Finite(b) DMRG

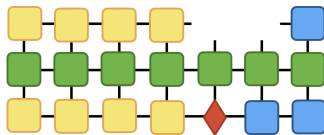


1) The initial network is shown below, where the initial MPS is randomly generated.

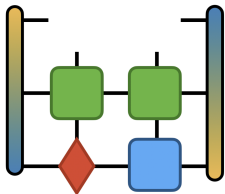


The initial MPS must be in the canonical form with the orthogonality center (red diamond) in the i -th position.

2) Select the i -th and $(i+1)$ -th sites, one of which must be the center of orthogonality. Remove those sites with a partial derivative, so that the network can be variationally optimized.

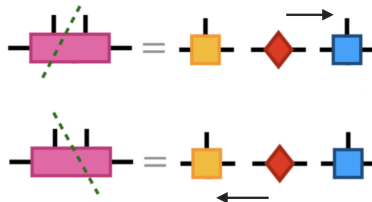


3) Contract the tensors. The tensor on the left represents the left environment. Same thing for the right.

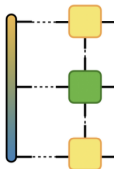


4) Apply Lanczos to update the tensors (discussed later).

5) Apply the SVD to the output of the Lanczos algorithm. And move the orthogonality center on the right/left (depending on the sweep).



6) Substitute the optimized tensors in the previous MPS and update the left/right environment (if you are moving to the right/left).



7) Repeat 2,3,4,5 and 6 scanning back and forth across the lattice until a set number of iterations (called sweeps) or until some relevant quantities reach convergence.

Lanczos algorithm (1)



Start with a random vector: $|\psi_1\rangle \in \mathbb{C}^n$.

For $i = 1$:

$$\alpha_1 = \langle \psi_1 | \mathcal{H} | \psi_1 \rangle \quad \beta_1 = 0$$

$$\phi_1 = \mathcal{H}|\psi_1\rangle - \alpha_1|\psi_1\rangle$$

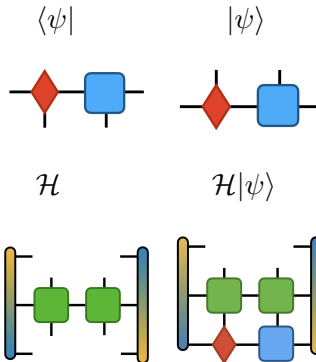
For $i = 2, \dots, m$:

$$\beta_i^2 = \langle \phi_{i-1} | \phi_{i-1} \rangle$$

$$|\psi_i\rangle = |\phi_{i-1}\rangle / \beta_i$$

$$\alpha_i = \langle \psi_i | \mathcal{H} | \psi_i \rangle$$

$$\phi_i = \mathcal{H}|\psi_i\rangle - \alpha_i|\psi_i\rangle - \beta_i|\psi_{i-1}\rangle$$



At the end of the procedure, the Hamiltonian in the Krylov subspace takes the tridiagonal form (in our case, only 2 iterations are needed):

$$\mathcal{M} = \begin{pmatrix} \alpha_1 & \beta_2 & 0 & \cdots \\ \beta_2 & \alpha_2 & \beta_3 & \cdots \\ 0 & \beta_3 & \alpha_3 & \ddots \\ \vdots & \vdots & \ddots & \ddots \end{pmatrix}$$

$$\mathcal{M} = U \mathbf{D} U^\dagger$$

$$|\psi_{GS}^{optimized}\rangle = (U_{[1,1]})^* |\psi_1\rangle + (U_{[2,1]})^* |\psi_2\rangle + \cdots + (U_{[m,1]})^* |\psi_m\rangle$$

Ising model Hamiltonian

$$\hat{H}_{\text{Ising}}^{\text{n.n.}} = -J \underbrace{\sum_i \hat{S}_i^z \hat{S}_{i+1}^z}_{\text{nearest neighbor}} - H \underbrace{\sum_i \hat{S}_i^x}_{\text{local}}$$

The Ising interaction is **short-range** due to nearest neighbor (n.n.)

Ising model Hamiltonian

$$\hat{H}_{\text{Ising}}^{\text{n.n.}} = -J \underbrace{\sum_i \hat{S}_i^z \hat{S}_{i+1}^z}_{\text{nearest neighbor}} - H \underbrace{\sum_i \hat{S}_i^x}_{\text{local}}$$

The Ising interaction is **short-range** due to nearest neighbor (n.n.)

Ising model MPO

$$\text{MPO} = \begin{pmatrix} \hat{\mathbb{I}} & \hat{S}_z & -H\hat{S}_x \\ \hat{\mathbb{O}} & \hat{\mathbb{O}} & -J\hat{S}_z \\ \hat{\mathbb{O}} & \hat{\mathbb{O}} & \hat{\mathbb{I}} \end{pmatrix}$$

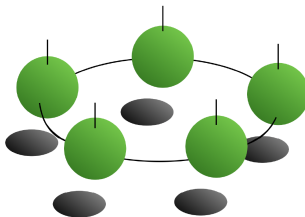
MPO bond dimensions are small ($\text{dim} = 3$) due to the Hamiltonian form

Periodic Boundary Conditions (PBC):

- MPS non-physical dimension is constant

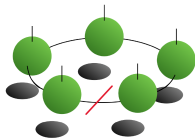


- MPS closed in a loop



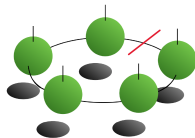
Operative way:

1. Cut the MPS in a specific position



2. Optimize the MPS using the normal Finite DMRG

3. Cut the MPS in another (different) position



4. Repeat cut and optimize until energy converges

To address the interaction between the initial and final links in the chain it is necessary to modify the MPO



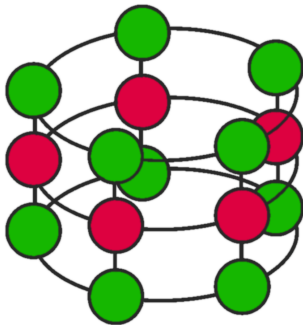
Let us consider the MPO for 3 sites:

$$\begin{array}{c} \text{pink} \\ \text{green} \\ \text{green} \end{array} = \left(\begin{array}{c} ? \\ ? \\ ? \end{array} \right) \begin{pmatrix} \hat{I} & \hat{S}_z & -H\hat{S}_x \\ \hat{O} & \hat{O} & -J\hat{S}_z \\ \hat{O} & \hat{O} & \hat{I} \end{pmatrix} \begin{pmatrix} \hat{I} & \hat{S}_z & -H\hat{S}_x \\ \hat{O} & \hat{O} & -J\hat{S}_z \\ \hat{O} & \hat{O} & \hat{I} \end{pmatrix}$$

The trace of the full contracted MPO must represent the correct Hamiltonian of the system (chain of 3 spins):

$$\begin{aligned}
 & \text{Diagram: A box containing three MPO tensors (pink, green, green) connected in a chain, with their top and bottom indices contracted in a loop.} \\
 & = -J\hat{S}_z^1\hat{S}_z^2 - J\hat{S}_z^2\hat{S}_z^3 - J\hat{S}_z^3\hat{S}_z^1 - H\hat{S}_x^1 - H\hat{S}_x^2 - H\hat{S}_x^3 \\
 & \left(\begin{array}{ccc} \hat{\mathbb{I}} & \hat{\mathbb{I}} & \hat{\mathbb{I}} \\ -J\hat{S}_z & \hat{\mathbb{I}} & \hat{\mathbb{I}} \\ \hat{\mathbb{I}} & \hat{S}_z & -H\hat{S}_x \end{array} \right) \left(\begin{array}{ccc} \hat{\mathbb{I}} & \hat{S}_z & -H\hat{S}_x \\ \hat{\mathbb{I}} & \hat{\mathbb{I}} & -J\hat{S}_z \\ \hat{\mathbb{I}} & \hat{\mathbb{I}} & \hat{\mathbb{I}} \end{array} \right) \left(\begin{array}{ccc} \hat{\mathbb{I}} & \hat{S}_z & -H\hat{S}_x \\ \hat{\mathbb{I}} & \hat{\mathbb{I}} & -J\hat{S}_z \\ \hat{\mathbb{I}} & \hat{\mathbb{I}} & \hat{\mathbb{I}} \end{array} \right) \\
 & \left(\begin{array}{ccc} \hat{\mathbb{I}} & \hat{\mathbb{I}} & \hat{\mathbb{I}} \\ -J\hat{S}_z^1 & -J\hat{S}_z^1\hat{S}_z^3 & J\hat{S}_z^1(H\hat{S}_x^3 + H\hat{S}_x^2 + J\hat{S}_z^2\hat{S}_z^3) \\ \hat{\mathbb{I}} & -J\hat{S}_z^3 & -H\hat{S}_x^3 - H\hat{S}_x^2 - H\hat{S}_x^1 - J\hat{S}_z^1\hat{S}_z^2 - J\hat{S}_z^2\hat{S}_z^3 \end{array} \right)
 \end{aligned}$$

The **ANSATZ** is correct.



$$\langle \psi_{tot} | \mathcal{H} | \psi_{tot} \rangle$$

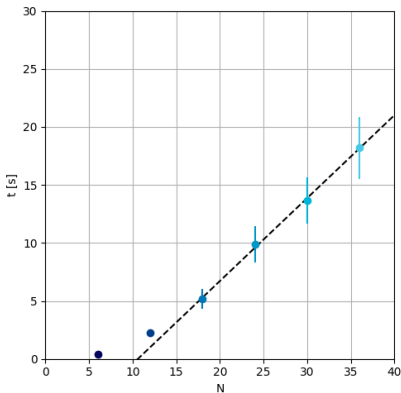
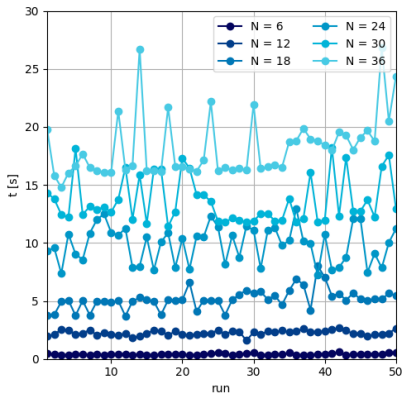
To verify the convergence of energy:

- Combine the MPS in one cycle (green)
- Normalize the MPS in this form
- Merge the modified MPO into one cycle (red)
- Contract everything

DMRG Results (Time)



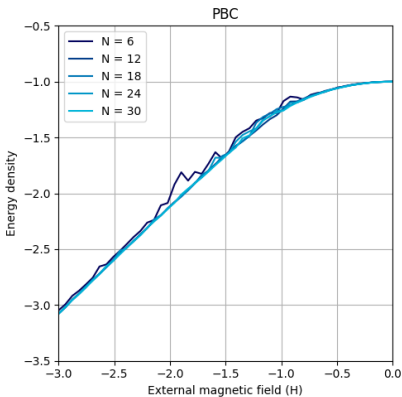
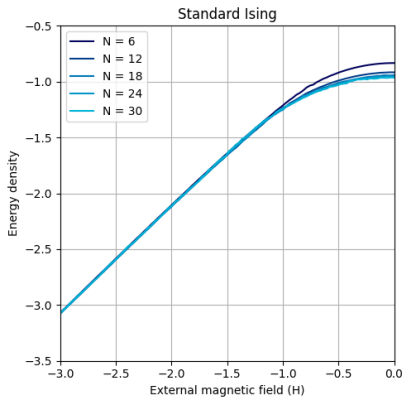
Time taken for the algorithm to run: linear behavior



DMRG Results (Correctness)



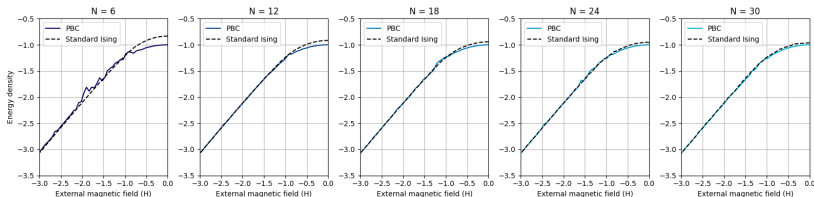
Ground state energy for different values of the external field: go beyond the writable hamiltonian limit



DMRG Results (Correctness)



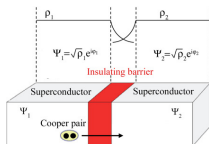
Comparison between the ground state with and without the PBC



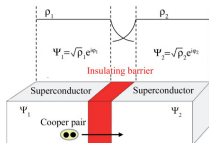
Attention: the results suffer of “over-optimization” if the convergence limit $\epsilon = E_{new} - E_{old}$ is too small in the sweep optimization (impossible to eliminate for small N)

Attention: bond dimension $\chi = 1$ is sufficient to retrieve the ground state energy results

Josephson Junction (JJ):

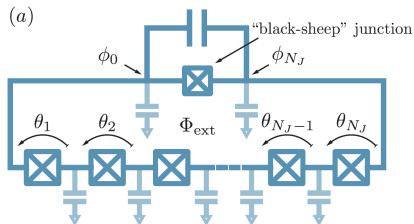


Josephson Junction (JJ):



Fluxonium components:

- N Josephson Junctions
- Black-sheep Josephson Junction
- Shunted capacitances



Strengths:

- big anharmonicity
- resistance to charge noise
- better control

Drowbacks:

- hard to compute theoretically
- hard to simulate

Strengths:

- big anharmonicity
- resistance to charge noise
- better control

Drowbacks:

- hard to compute theoretically
- hard to simulate

Fluxonium Hamiltonian:

$$H = \sum_{i=1}^{N_j} H_{0,i} + \sum_j h g_{ij} \hat{n}_i \hat{n}_j - E_{J_B} \cos \left(\sum_{i=1}^{N_j} \theta_i + \phi_{\text{ext}} \right)$$

$$H_{0,i} = 4E_{C_i} \hat{n}_i^2 + E_{J_i} \cos(\theta_i)$$

Canonical quantization:

$$\begin{cases} \Phi \rightarrow \hat{\Phi} \\ Q \rightarrow \hat{Q} \end{cases} \rightarrow \begin{cases} \hat{\phi} = 2\pi\hat{\Phi}/\Phi_0 \\ \hat{n} = \hat{Q}/2e \end{cases}$$

Classical energy:

$$\text{Capacitance} \rightarrow C\dot{\Phi}^2/2$$

$$\text{Inductance} \rightarrow \dot{\Phi}^2/2L$$

Canonical quantization:

$$\begin{cases} \Phi \rightarrow \hat{\Phi} \\ Q \rightarrow \hat{Q} \end{cases} \rightarrow \begin{cases} \hat{\phi} = 2\pi\hat{\Phi}/\Phi_0 \\ \hat{n} = \hat{Q}/2e \end{cases}$$

Classical energy:

$$\text{Capacitance} \rightarrow C\dot{\Phi}^2/2$$

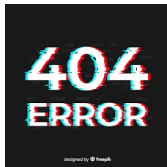
$$\text{Inductance} \rightarrow \dot{\Phi}^2/2L$$

Fluxonium Lagrangian:

$$\begin{aligned} L(\phi, \dot{\phi}) = & \frac{C_{JB}}{2} (\dot{\phi}_N - \dot{\phi}_0)^2 + \sum_{i=1}^{N_j} C_{ij} (\dot{\phi}_i - \dot{\phi}_{i-1})^2 + \sum_{i=0}^{N_j} \frac{C_{0i}}{2} \dot{\phi}_i^2 \\ & + \sum_{i=1}^{N_j} E_{J_i} \cos\left(\frac{\phi_i - \phi_{i-1}}{\phi_0}\right) + E_{J_B} \cos\left(\frac{\phi_N - \phi_0 + \phi_{\text{ext}}}{\phi_0}\right) \end{aligned}$$

Fluxonium Hamiltonian:

$$H = \sum_{i=1}^{N_j} H_{0,i} + \underbrace{\sum_j hg_{ij} \hat{n}_i \hat{n}_j}_{\text{All-to-all}} - E_{J_B} \cos \left(\sum_{i=1}^{N_j} \theta_i + \phi_{\text{ext}} \right)$$



We were not able to implement the all-to-all connection in an optimized way

Topics of our project:

- How to write MPS and orthogonality center handling
- How to write MPO with and without PBC
- Finite DMRG algorithm using Lanczos in both the canonical way and the tensor network implementation
- Circuit quantization and fluxonium description

Topics of our project:

- How to write MPS and orthogonality center handling
- How to write MPO with and without PBC
- Finite DMRG algorithm using Lanczos in both the canonical way and the tensor network implementation
- Circuit quantization and fluxonium description

Further developments:

- How to shrink the MPO operator for the all-to-all connection
- Implement Block-Lanczos in DMRG algorithm

T. E. Baker, S. Desrosiers, M. Tremblay, and M. P. Thompson, "Méthodes de calcul avec réseaux de tenseurs en physique", *Canadian Journal of Physics* **99**, 207–221 (2021).

T. E. Baker and M. P. Thompson, *Build your own tensor network library: dmrjulia i. basic library for the density matrix renormalization group*, 2021.

T. E. Baker, A. Foley, and D. Sénéchal, *Direct solution of multiple excitations in a matrix product state with block lanczos*, 2023.

A. Di Paolo, T. E. Baker, A. Foley, D. Sénéchal, and A. Blais, "Efficient modeling of superconducting quantum circuits with tensor networks", *npj Quantum Information* **7**, 10.1038/s41534-020-00352-4 (2021).

M. Weyrauch and M. Rakov, "Efficient mps algorithm for periodic boundary conditions and applications section.", *Ukrainian Journal of Physics* **58**, 657–665 (2013).

U. Schollwöck, "The density-matrix renormalization group in the age of matrix product states", *Annals of Physics* **326**, 96–192 (2011).

U. Schollwöck, "The density-matrix renormalization group", *Reviews of Modern Physics* **77**, 259–315 (2005).

C. Hubig, I. P. McCulloch, and U. Schollwöck, "Generic construction of efficient matrix product operators", *Physical Review B* **95**, 10.1103/physrevb.95.035129 (2017).

S. Montangero, *Introduction to tensor network methods: numerical simulations of low-dimensional many-body quantum systems*, 1st ed., Physics and Astronomy (R0) (Springer Cham, Dec. 2018), pp. XV, 172.

E. F. D'Azevedo, W. R. Elwasif, N. D. Patel, and G. Alvarez, *Targeting multiple states in the density matrix renormalization group with the singular value decomposition*, 2019.

R. Orús, "A practical introduction to tensor networks: matrix product states and projected entangled pair states", *Annals of Physics* **349**, 117–158 (2014).

J. Hauschild and F. Pollmann, "Efficient numerical simulations with tensor networks: tensor network python (tenpy)", *SciPost Physics Lecture Notes*, 10.21468/scipostphyslectnotes.5 (2018).

G. Catarina and B. Murta, "Density-matrix renormalization group: a pedagogical introduction", *The European Physical Journal B* **96**, 10.1140/epjb/s10051-023-00575-2 (2023).