# **Machine Learning Theory**

Lecture 12: Chaining

David A. Hirshberg May 24, 2024

Emory University

#### **Notation**

Today, we're using the sample mean inner product and sample mean squared error. To keep notation simple, we're going to write this without any special subscripts.

$$\begin{split} \langle u, v \rangle &= \langle u, v \rangle_{L_2(\mathbf{P_n})} = \frac{1}{n} \sum_{i=1}^n u(X_i) v(X_i) \\ \|v\|^2 &= \|v\|_{L_2(\mathbf{P_n})}^2 = \frac{1}{n} \sum_{i=1}^n v(X_i)^2. \end{split}$$

Keep in mind that for a gaussian vector  $g \sim N(0, I_{n \times n})$  and any function v,

$$\langle g, v \rangle = \frac{1}{n} \sum_{i=1}^{n} g_i \ v(X_i) \sim N\left(0, \ \frac{\|v\|^2}{n}\right).$$

We'll also write  $\mathcal{M}_s$  as a shorthand for what we've called  $\mathcal{M}_s - \mu^*$  before.

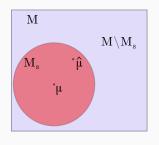
$$\mathcal{M}_s = \{ m - \mu^* : ||m - \mu^*|| \le s \}.$$

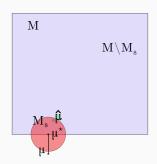
And we'll ignore some constant factors:  $a_n \lesssim b_n$  means  $a_n \leq cb_n$  for some constant c.

# Review

#### What determines our error bounds

It's the gaussian width of neighborhoods of  $\mu^{\star}$  in our model.





$$\|\hat{\mu} - \mu^{\star}\| < s \text{ with high probability if } s^2 \ge 2\sigma c_{\delta} \operatorname{w}(\mathcal{M}_s)$$
$$\mathcal{M}_s = \{m - \mu^{\star} \in \mathcal{M} : \|m - \mu^{\star}\| \le s\}.$$

So what we need is a way to bound the gaussian width of these neighborhoods.

#### Finite Models

- In finite models, bounding width is easy.
- It's the maximum of gaussians with standard deviation  $\leq s/\sqrt{n}$ .

$$E\langle g, m - \mu^* \rangle^2 = \frac{\|m - \mu^*\|^2}{n}.$$

 We can bound this via union bound. It's down to counting the curves in the model.

$$\mathrm{w}(\mathcal{M}_s)\lesssim s\sqrt{\dfrac{\log(K)}{n}}$$
 if  $\mathcal{M}$  contains  $K$  curves  $m$  all with  $\|m-\mu^\star\|_{L_2(\mathrm{Pn})}\leq s.$ 

- $\cdot$  We may be overcounting. This bounds the max of K totally different gaussians.
- That's the case in which it's largest, so if there's correlation we're overcounting.
- · And this definitely won't work for models with infinitely many curves.
- We'll need to take advantage of this correlation to tackle infinite models.

### **Counting Curves in Infinite Models**

Gaussian width is the mean of the maximum of a set of gaussians.

$$w(\mathcal{M}_s) = E \max_{v \in \mathcal{M}_s} \langle g, v \rangle$$
 for  $g \sim N(0, I_{n \times n})$ .

And the difference between many of these gaussians  $\langle g, v \rangle$  will be small.

- So small, sometimes, that we don't need to 'pay probability' to bound them all using the union bound. They needn't contribute to K.
- We can just use the Cauchy-Schwarz inequality to bound differences.

$$|\langle g, u \rangle - \langle g, v \rangle| = |\langle g, u - v \rangle| \le ||g|| ||u - v|| \approx ||u - v||.$$

If the curves u and v are close enough, by bounding  $\langle g, u \rangle$ , we bound  $\langle g, v \rangle$  for free.

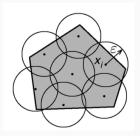
- This means we can take K above to be smaller than the total number of curves.
- · It's enough that some set  $u_1 \dots u_K$  gets close enough to all curves  $v \in \mathcal{M}$ .

This means we have to talk about how many meaningfully different curves we have.

### $\epsilon$ -covers and snapping

We can quantify this using a set's  $\epsilon$ -covering number  $K_{\epsilon}$ . That's the number of balls of size  $\epsilon$  of radius  $\epsilon$  it takes to cover the set. That is, it's the size of the set's smallest  $\epsilon$ -cover.

We call a set  $\mathcal{V}^{\epsilon}$  an  $\epsilon$ -cover for the set  $\mathcal{V}$  if every curve in the set  $\mathcal{V}$  is within a distance  $\epsilon$  of some curve in  $\mathcal{V}^{\epsilon}$ .



We can think of this as the set of curves we get by snapping each curve in  $\mathcal V$  to one of finitely many curves—one that's an approximation with error  $\leq \epsilon$ .

$$\mathcal{V}_{\epsilon} = \{\pi_{\epsilon}(v) : v \in \mathcal{V}\}$$
 where  $\|\pi_{\epsilon}(v) - v\| \le \epsilon$ 

I'll call the function  $\pi_{\epsilon}$  that does this an  $\epsilon$ -snapping map.

That's not standard terminology. As far as I know there isn't a standard name for this.

If we've got an  $\epsilon$ -snapping map, we've got an  $\epsilon$ -cover.

$$\mathcal{V}_{\epsilon} = \{\pi_{\epsilon}(v) : v \in \mathcal{V}\}$$
 where  $\|\pi_{\epsilon}(v) - v\| \leq \epsilon$ 

We can go the other way, too. If we've got an  $\epsilon$ -cover, we can define an  $\epsilon$ -snapping map. How?

If we've got an  $\epsilon$ -snapping map, we've got an  $\epsilon$ -cover.

$$\mathcal{V}_{\epsilon} = \{\pi_{\epsilon}(v) : v \in \mathcal{V}\}$$
 where  $\|\pi_{\epsilon}(v) - v\| \le \epsilon$ 

We can go the other way, too. If we've got an  $\epsilon$ -cover, we can define an  $\epsilon$ -snapping map. How?

We snap to the closest curve in our cover.

$$\pi_{\epsilon}(v) = \underset{v_{\epsilon} \in \mathcal{V}^{\epsilon}}{\operatorname{argmin}} \|v_{\epsilon} - v\|$$

This means snapping maps and covers are more-or-less interchangeable.

#### Terminology.

I'll refer to the *size* of a snapping map as the size of the cover induced by it, i.e., the number of different curves it outputs.

### **Snapping and Gaussian Width**

If we have an  $\epsilon$ -snapping map of size  $K_{\epsilon}$  for a set  $\mathcal{V}$ , then we've got a bound on its gaussian width. We use  $\epsilon$ -closeness together with with our bound for finite sets.

$$\begin{split} \mathbf{w}(\mathcal{V}) &= \mathbf{E} \max_{v \in \mathcal{V}} \left\langle g, \ v \right\rangle \\ &= \mathbf{E} \max_{v \in \mathcal{V}} \{ \left\langle g, \ v - \pi_{\epsilon}(v) \right\rangle + \left\langle g, \pi_{\epsilon}(v) \right\rangle \} \\ &\lesssim \| v - \pi_{\epsilon}(v) \| + \max_{v \in \mathcal{V}} \| \pi(v) \| \sqrt{\frac{\log(K_{\epsilon})}{n}} \end{split}$$

When we're talking about a centered neighborhood  $\mathcal{V} = \mathcal{M}_s - \mu$ , this second term is small because  $\|\pi(v)\|$  is small for every  $v \in \mathcal{V}$ .

$$\|\pi(v)\| \leq \|v-\pi(v)\| + \|v\| \leq \epsilon + s$$
 by the triangle inequality and therefore

$$\mathrm{w}(\mathcal{M}_s - \mu) \lesssim \epsilon + s \sqrt{\dfrac{\log(K_\epsilon)}{n}}$$
 when  $\dfrac{\log(K_\epsilon)}{n} \lesssim 1$ 

Gaussian width doesn't change when we center, so the same bound holds for the neighborhood itself.

# **Dissatisfying Implications**

- · We showed last class that  $\log(K_\epsilon) \approx 1/\epsilon$  for the Lipschitz model.
- If we choose the resolution  $\epsilon$  to minimize our bound, it's roughly  $\sqrt[3]{s^2/n}$ .

$$\mathrm{w}(\mathcal{M}_s) \lesssim \epsilon + s \sqrt{\frac{\log(K_\epsilon)}{n}} \approx \epsilon + \frac{s}{\sqrt{\epsilon n}} \approx s^{2/3} n^{-1/3} \quad \text{ at optimal } \quad \epsilon \approx s^{2/3} n^{-1/3}.$$

• This tells us that our estimator converges at a fourth-root rate.

$$s^2 \ge w(\mathcal{M}_s)$$
 if  $s^2 \gtrsim s^{2/3} n^{-1/3}$  i.e. if  $s \approx n^{-1/4}$ .

- · But we know it converges faster.
- The Lipschitz model is contained in the Sobolev model of order 1.
- And we proved the rate of convergence  $s \approx n^{-1/3}$  for that using Fourier series.

We can do better by looking at covering numbers at multiple resolutions.

$$\mathrm{w}(\mathcal{M}_s) \lesssim \frac{1}{\sqrt{n}} \int_0^s \sqrt{\log(K_\epsilon)} d\epsilon$$

This is called *Dudley's Integral Bound*. Today we'll prove it.

## Dudley's Integral Bound

$$\mathrm{w}(\mathcal{M}_s) \lesssim \frac{1}{\sqrt{n}} \int_0^s \sqrt{\log(K_\epsilon)} d\epsilon$$

- · It's based on an idea called *chaining*.
- The idea is to use approximations  $\pi_0(m), \pi_1(m), \ldots$  at increasing resolutions  $\epsilon_0, \epsilon_1, \ldots$
- We write each function as a sum of differences between finer and finer approximations.

$$m = \pi_0(m) + \sum_{j=0}^{\infty} \pi_{j+1}(m) - \pi_j(m)$$

- · We call these differences links in a chain that goes
  - from the coarsest approximation,  $\pi_0(m)$ , which is the same for all functions.
  - · to the finest approximation,  $m=\pi_\infty(m)$  itself.
- · Before we dig into this too much, let's warm up.

Warm-up

#### Our One-Link Bound

Think about the width bound implied by an  $\epsilon$ -snapping map  $\pi_\epsilon$  for very small  $\epsilon$ .

$$\begin{split} \mathbf{w}(\mathcal{V}) &\leq \mathbf{E} \max_{v \in \mathcal{V}} \langle g, v - \pi_{\epsilon}(v) \rangle + \mathbf{E} \max_{v \in \mathcal{V}} \langle g, \pi(v) \rangle \\ &\leq \mathbf{E} \|g\| \max_{v \in \mathcal{M}_s} \|v - \pi_{\epsilon}(v)\| + \mathbf{E} \max_{v \in \mathcal{V}} \langle g, \pi(v) \rangle \\ &\lesssim \underline{\epsilon} + \mathrm{rad}(\mathcal{V}) \sqrt{\frac{\log(K_{\epsilon})}{n}} \quad \text{where} \quad \mathrm{rad}(\mathcal{V}) = \max_{v \in \mathcal{V}} \|v\| \end{split}$$

- This is what we've been doing. But we have a sense that we're being wasteful.
- When our  $\epsilon$ -cover is fine, it'll contain vectors that are close to one another.
- The corresponding gaussians will be highly correlated, so our  $\sqrt{\log(K)}$  bound on their maximum will be loose. Our second term will be bigger than we want.

We could reduce  $K_\epsilon$  by snapping to coarser approximations—taking  $\epsilon$  to be large. But that makes our first term big.

We can do better by using two approximations—one coarse and one fine.

$$\langle g, \pi_{\epsilon}(v) \rangle = \langle g, \pi_{\epsilon}(v) - \pi_{\epsilon'}(v) \rangle + \langle g, \pi_{\epsilon'}(v) \rangle$$

where  $\pi_{\epsilon'}(v)$  is a snapping map that gives coarser approximations. One with coarser resolution  $\epsilon' \geq \epsilon$  and therefore smaller size  $K'_{\epsilon} \leq K_{\epsilon}$ . We bound the pieces as before.

$$\begin{split} & \mathbf{w}(\mathcal{V}) \lesssim \mathbf{E} \max_{v \in \mathcal{V}} \langle g, v - \pi(v) \rangle + \mathbf{E} \max_{v \in \mathcal{V}} \langle g, \pi(v) - \pi'(v) \rangle + \mathbf{E} \max_{v \in \mathcal{V}} \langle g, \pi'(v) - \mu^{\star} \rangle \\ & \lesssim \max_{v \in \mathcal{V}} \lVert v - \pi(v) \rVert + \underbrace{\max_{v \in \mathcal{V}} \lVert \pi(v) - \pi'(v) \rVert}_{\leq \epsilon + \epsilon'} \sqrt{\frac{\log(K_{\epsilon}K_{\epsilon'})}{n}} + \underbrace{\max_{v \in \mathcal{V}} \lVert \pi'(v) \rVert}_{\leq \operatorname{rad}(\mathcal{V}) + \epsilon'} \sqrt{\frac{\log(K_{\epsilon'})}{n}} \\ & \approx \epsilon + \epsilon' \sqrt{\frac{\log(K_{\epsilon})}{n}} + \operatorname{rad}(\mathcal{V}) \sqrt{\frac{\log(K_{\epsilon'})}{n}}. \end{split}$$

Q: Where do we get this second bound with  $\log(K_{\epsilon}K_{\epsilon'})$ ?

$$\begin{split} & \mathbf{w}(\mathcal{V}) \lesssim \mathbf{E} \max_{v \in \mathcal{V}} \langle g, v - \pi(v) \rangle + \mathbf{E} \max_{v \in \mathcal{V}} \langle g, \pi(v) - \pi'(v) \rangle + \mathbf{E} \max_{v \in \mathcal{V}} \langle g, \pi'(v) - \mu^{\star} \rangle \\ & \lesssim \max_{v \in \mathcal{V}} \lVert v - \pi(v) \rVert + \underbrace{\max_{v \in \mathcal{V}} \lVert \pi(v) - \pi'(v) \rVert}_{\leq \epsilon + \epsilon'} \sqrt{\frac{\log(K_{\epsilon}K_{\epsilon'})}{n}} + \underbrace{\max_{v \in \mathcal{V}} \lVert \pi'(v) \rVert}_{\leq \operatorname{rad}(\mathcal{V}) + \epsilon'} \sqrt{\frac{\log(K_{\epsilon'})}{n}} \\ & \approx \epsilon + \epsilon' \sqrt{\frac{\log(K_{\epsilon})}{n}} + \operatorname{rad}(\mathcal{V}) \sqrt{\frac{\log(K_{\epsilon'})}{n}}. \end{split}$$

Q: Where do we get this second bound with  $\log(K_{\epsilon}K_{\epsilon'})$ ?

- There are  $K_{\epsilon}K_{\epsilon'}$  pairs of the  $K_{\epsilon}$  values of  $\pi$  and the  $K_{\epsilon'}$  values of  $\pi'$ .
- We could probably find a better bound.
- Probably not many more than  $K_{\epsilon}$  occur as  $\pi(v)$  and  $\pi'(v)$  for some point v.
- · But the difference between  $K_{\epsilon}$  and  $K_{\epsilon}K_{\epsilon'}$  doesn't matter here.

$$\begin{split} & \mathbf{w}(\mathcal{V}) \lesssim \mathbf{E} \max_{v \in \mathcal{V}} \langle g, v - \pi(v) \rangle + \mathbf{E} \max_{v \in \mathcal{V}} \langle g, \pi(v) - \pi'(v) \rangle + \mathbf{E} \max_{v \in \mathcal{V}} \langle g, \pi'(v) - \mu^{\star} \rangle \\ & \lesssim \max_{v \in \mathcal{V}} \lVert v - \pi(v) \rVert + \underbrace{\max_{v \in \mathcal{V}} \lVert \pi(v) - \pi'(v) \rVert}_{\leq \epsilon + \epsilon'} \sqrt{\frac{\log(K_{\epsilon}K_{\epsilon'})}{n}} + \underbrace{\max_{v \in \mathcal{V}} \lVert \pi'(v) \rVert}_{\leq \operatorname{rad}(\mathcal{V}) + \epsilon'} \sqrt{\frac{\log(K_{\epsilon'})}{n}} \\ & \approx \epsilon + \epsilon' \sqrt{\frac{\log(K_{\epsilon})}{n}} + \operatorname{rad}(\mathcal{V}) \sqrt{\frac{\log(K_{\epsilon'})}{n}}. \end{split}$$

Q: Why is the last approximation valid?

$$\begin{split} & \mathbf{w}(\mathcal{V}) \lesssim \mathbf{E} \max_{v \in \mathcal{V}} \langle g, v - \pi(v) \rangle + \mathbf{E} \max_{v \in \mathcal{V}} \langle g, \pi(v) - \pi'(v) \rangle + \mathbf{E} \max_{v \in \mathcal{V}} \langle g, \pi'(v) - \mu^{\star} \rangle \\ & \lesssim \max_{v \in \mathcal{V}} \lVert v - \pi(v) \rVert + \underbrace{\max_{v \in \mathcal{V}} \lVert \pi(v) - \pi'(v) \rVert}_{\leq \epsilon + \epsilon'} \sqrt{\frac{\log(K_{\epsilon}K_{\epsilon'})}{n}} + \underbrace{\max_{v \in \mathcal{V}} \lVert \pi'(v) \rVert}_{\leq \operatorname{rad}(\mathcal{V}) + \epsilon'} \sqrt{\frac{\log(K_{\epsilon'})}{n}} \\ & \approx \epsilon + \epsilon' \sqrt{\frac{\log(K_{\epsilon})}{n}} + \operatorname{rad}(\mathcal{V}) \sqrt{\frac{\log(K_{\epsilon'})}{n}}. \end{split}$$

Q: Why is the last approximation valid?

· Triangle inequality.

$$\|\pi(v) - \pi'(v)\| = \|\pi(v) - v + v - \pi'(v)\|$$

$$\leq \|\pi(v) - v\| + \|\pi'(v) - v\| \leq \epsilon + \epsilon' \leq 2\epsilon'.$$

· Log of products is sum of logs.

$$\log(K_{\epsilon}K_{\epsilon}') \le \log(K_{\epsilon}) + \log(K_{\epsilon'}) \le 2\log(K_{\epsilon}).$$

Let's think about the Lipschitz model again.

$$\log(K_{\epsilon}) \approx 1/\epsilon$$
.

Old Bound

$$\mathrm{w}(\mathcal{M}_s) \lesssim \epsilon + s \sqrt{\frac{\log(K_\epsilon)}{n}} \approx s^{2/3} n^{-1/3}$$
 at optimal  $\epsilon \approx s^{2/3} n^{-1/3}$   $\implies s^2 \geq \mathrm{w}(\mathcal{M}_s)$  for  $s^{4/3} \approx n^{-1/3}$  i.e.  $s \approx n^{-1/4}$ .

New Bound

$$\begin{split} \mathrm{w}(\mathcal{M}_s) &\lesssim \epsilon + \epsilon' \sqrt{\frac{\log(K_\epsilon)}{n}} + s \sqrt{\frac{\log(K_\epsilon')}{n}} \approx s^{4/7} n^{-3/7} \quad \text{ at optimal } \quad \epsilon \approx s^{4/7} n^{-3/7} \\ &\qquad \qquad \epsilon' \approx n^{1/2} \epsilon^{3/2} \\ &\Longrightarrow \quad s^2 \geq \mathrm{w}(\mathcal{M}_s) \quad \text{for} \quad s^{10/7} \approx n^{-3/7} \quad \text{i.e.} \quad s \approx n^{-3/10}. \end{split}$$

This isn't the  $s \approx n^{-1/3}$  bound we got using Fourier series, but it's closer. Let's see what happens when we use a longer chain of approximations.

Let's think about the Lipschitz model again.

$$\log(K_{\epsilon}) \approx 1/\epsilon$$
.

Old Bound

$$\mathrm{w}(\mathcal{M}_s) \lesssim \epsilon + s \sqrt{\frac{\log(K_\epsilon)}{n}} \approx s^{2/3} n^{-1/3}$$
 at optimal  $\epsilon \approx s^{2/3} n^{-1/3}$   $\implies s^2 \geq \mathrm{w}(\mathcal{M}_s)$  for  $s^{4/3} \approx n^{-1/3}$  i.e.  $s \approx n^{-1/4}$ .

New Bound

$$\begin{split} \mathrm{w}(\mathcal{M}_s) &\lesssim \epsilon + \epsilon' \sqrt{\frac{\log(K_\epsilon)}{n}} + s \sqrt{\frac{\log(K_\epsilon')}{n}} \approx s^{4/7} n^{-3/7} \quad \text{at optimal} \quad \epsilon \approx s^{4/7} n^{-3/7} \\ &\qquad \qquad \epsilon' \approx n^{1/2} \epsilon^{3/2} \\ &\Longrightarrow \quad s^2 \geq \mathrm{w}(\mathcal{M}_s) \quad \text{for} \quad s^{10/7} \approx n^{-3/7} \quad \text{i.e.} \quad s \approx n^{-3/10}. \end{split}$$

No magic here. We optimize as usual.

- 1. Set the derivative with respect to  $\epsilon'$  to zero and solve for  $\epsilon'$  in terms of  $\epsilon$ .
- 2. Set the derivative with respect to  $\epsilon$  to zero and solve for  $\epsilon$ .



Suppose we want to bound the gaussian width of a set  $\mathcal{V}$ .

$$w(V) = E \max_{v \in V} \langle g, v \rangle.$$

- And we have, for each  $v \in \mathcal{V}$ , increasingly fine approximations  $\pi_0(v) \dots \pi_M(v)$ .
- · These are the closest vectors to v in  $\epsilon$ -covers for increasingly small  $\epsilon_0 \ldots \epsilon_M$ .
- Then we write each  $v \in \mathcal{V}$  as the sum over links in a chain from  $\pi_0(v)$  to  $\pi_M(v)$ .
- Plus a final link from the finest approximation,  $\pi_M(v)$ , to v itself.

$$v = v - \pi_M(v) + \sum_{j=1}^M \pi_j(v) - \pi_{j-1}(v).$$
 a link in the chain

- $\cdot$  We can expand v this way when we write our gaussian width.
- And we can bound it by maximizing each term separately.
- Just like we did in our warm-up, but with more terms.

The thing we're bounding.

$$w(\mathcal{V}) = E \max_{v \in \mathcal{V}} \langle g, v \rangle.$$

The decomposition we're working with.

$$v = v - \pi_M(v) + \sum_{j=1}^M \pi_j(v) - \pi_{j-1}(v).$$
 a link in the chain

The bound we get.

$$\begin{aligned} \mathbf{w}(\mathcal{V}) &= \mathbf{E}\left[\max_{v \in \mathcal{V}} \langle g, v - \pi_M(v) \rangle + \sum_{j=1}^{M} \langle g, \ \pi_j(v) - \pi_{j-1}(v) \rangle\right] \\ &\leq \mathbf{E}\max_{v \in \mathcal{V}} \langle g, v - \pi_M(v) \rangle + \sum_{j=1}^{M} \mathbf{E}\max_{v \in \mathcal{V}} \langle g, \ \pi_j(v) - \pi_{j-1}(v) \rangle \\ &\lesssim \epsilon_M + \sum_{j=1}^{M} \epsilon_{j-1} \sqrt{\frac{\log\left(K_{\epsilon_j}\right)}{n}}. \end{aligned}$$

Now all we've got to do is choose  $\epsilon_0 \dots \epsilon_M$ 

#### Considerations

$$w(\mathcal{V}) \leq E \max_{v \in \mathcal{V}} \langle g, v - \pi_M(v) \rangle + \sum_{j=1}^M E \max_{v \in \mathcal{V}} \langle g, \pi_j(v) - \pi_{j-1}(v) \rangle$$
$$\lesssim \epsilon_M + \sum_{j=1}^M \epsilon_{j-1} \sqrt{\frac{\log(K_{\epsilon_j})}{n}}.$$

- We want K to be small.
  - That is, we want there to be few distinct values of each link  $\pi_j(v) \pi_{j-1}(v)$  for  $v \in \mathcal{V}$ .
  - The more values, the more gaussians  $\langle g, \pi_j(v) \pi_{j-1}(v) \rangle$  we have to deal with in our union bound.
- We want  $\epsilon$  to be small.
  - That is, we want all the links to be short in the sense that their variance  $\|\pi_j(v)-\pi_{j-1}\|^2/n$  is small.
  - The longer the links, the bigger the individual gaussians we need to bound.

We can't get both at any one resolution.

- The finer our resolution  $\epsilon_j$ , the more vectors we need in our cover.
- To balance these considerations, we use a lot of short links and a few large ones.
- · Since  $\epsilon_j$  and  $\sqrt{\log(K_{\epsilon_{j-1}})}$  are multiplied, this can make the product small.

#### Considerations

$$w(\mathcal{V}) \leq E \max_{v \in \mathcal{V}} \langle g, v - \pi_M(v) \rangle + \sum_{j=1}^M E \max_{v \in \mathcal{V}} \langle g, \pi_j(v) - \pi_{j-1}(v) \rangle$$
$$\lesssim \epsilon_M + \sum_{j=1}^M \epsilon_{j-1} \sqrt{\frac{\log(K_{\epsilon_j})}{n}}.$$

- · We want K to be small.
  - That is, we want there to be few distinct values of each link  $\pi_j(v) \pi_{j-1}(v)$  for  $v \in \mathcal{V}$ .
  - The more values, the more gaussians  $\langle g, \pi_j(v) \pi_{j-1}(v) \rangle$  we have to deal with in our union bound.
- We want  $\epsilon$  to be small.
  - That is, we want all the links to be short in the sense that their variance  $\|\pi_i(v) \pi_{i-1}\|^2/n$  is small.
  - · The longer the links, the bigger the individual gaussians we need to bound.

A sensible choice: halve  $\epsilon$  each time.  $\epsilon_j = 1/2^j$ .

Assuming all elements of  $\mathcal V$  are  $\epsilon=1$ -close, i.e.  $\epsilon_0=1$  is big enough that  $K_1=1$ .

$$\|\pi_j(v) - \pi_{j-1}(v)\| \le \|\pi_j(v) - v\| + \|v - \pi_{j-1}(v)\|$$
  
 
$$\le \epsilon_j + \epsilon_{j-1} = 1/2^j + 2/2^j = 3/2^j.$$

#### Considerations

$$w(\mathcal{V}) \leq E \max_{v \in \mathcal{V}} \langle g, v - \pi_M(v) \rangle + \sum_{j=1}^M E \max_{v \in \mathcal{V}} \langle g, \pi_j(v) - \pi_{j-1}(v) \rangle$$
$$\lesssim \epsilon_M + \sum_{j=1}^M \epsilon_{j-1} \sqrt{\frac{\log(K_{\epsilon_j})}{n}}.$$

- · We want K to be small.
  - That is, we want there to be few distinct values of each link  $\pi_j(v) \pi_{j-1}(v)$  for  $v \in \mathcal{V}$ .
  - The more values, the more gaussians  $\langle g, \pi_j(v) \pi_{j-1}(v) \rangle$  we have to deal with in our union bound.
- We want  $\epsilon$  to be small.
  - That is, we want all the links to be short in the sense that their variance  $\|\pi_i(v) \pi_{i-1}\|^2/n$  is small.
  - The longer the links, the bigger the individual gaussians we need to bound.

Plugging these in yields a bound in terms of cover sizes  $K_{\epsilon_j}$ 

$$\mathrm{w}(\mathcal{V}) \lesssim 2^{-M} + \sum_{j=1}^{M} \frac{3}{2^{j}} \sqrt{\frac{\log(K_{1/2^{j}})}{n}}.$$

# Chaining all the way

$$\mathrm{w}(\mathcal{V}) \lesssim 2^{-M} + \sum_{j=1}^{M} \frac{3}{2^j} \sqrt{\frac{\log\left(K_{1/2^j}\right)}{n}}.$$

- If  ${\mathcal V}$  is small enough in the right sense, the terms of the sum get small quickly.
- $\cdot$  And if terms get small quickly enough, the sum doesn't really depend much on M.
- This happens if  $\mathcal V$  has  $\epsilon$ -covers of size  $K_\epsilon \lesssim 2^{1/\epsilon^\alpha}$  for  $\alpha < 2$ .

$$\sum_{j=1}^M \frac{1}{2^j} \sqrt{\log \Bigl(K_{1/2^j}\Bigr)} \lesssim \sum_{j=1}^M \frac{1}{2^j} \sqrt{2^{\alpha j}} = \sum_{j=1}^M 2^{(\alpha/2-1)j} \leq \frac{2^{\alpha/2-1}}{1-2^{\alpha/2-1}}.$$

This means we can drop the special term for our final link from  $\pi_M(v) \to v$ .

- $\cdot$  If it doesn't matter how big M is, we can have this link be arbitrarily short.
- That is, we can use the limit of this bound as  $M \to \infty$ .

Often people approximate this sum by an integral

$$\begin{split} \mathbf{w}(\mathcal{V}) &\lesssim \frac{1}{\sqrt{n}} \sum_{j=1}^{M} \frac{1}{2^{j}} \sqrt{\log\left(K_{1/2^{j}}\right)} & \stackrel{(a)}{=} & \frac{1}{\sqrt{n}} \sum_{j=1}^{M} \int_{1/2^{j+1}}^{1/2^{j}} 2\sqrt{\log\left(K_{1/2^{j}}\right)} \\ &\stackrel{(b)}{\leq} & \frac{1}{\sqrt{n}} \sum_{j=1}^{M} \int_{1/2^{j+1}}^{1/2^{j}} 2\sqrt{\log(K_{\epsilon})} d\epsilon & = & \frac{2}{\sqrt{n}} \int_{1/2^{M+1}}^{1} \sqrt{\log(K_{\epsilon})} d\epsilon \\ &\stackrel{(c)}{\leq} & \frac{2}{\sqrt{n}} \int_{0}^{1} \sqrt{\log(K_{\epsilon})} d\epsilon \end{split}$$

(a) We're integrating a constant.

$$\int_{1/2^{j+1}}^{1/2^j} 2c = \left(\frac{1}{2^j} - \frac{1}{2^{j+1}}\right) 2c = \frac{1}{2^j} \left(1 - \frac{1}{2}\right) 2c$$

(b) Smaller  $\epsilon$ , bigger  $\epsilon$ -cover.

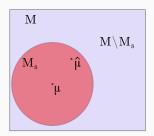
$$K_{\epsilon} \ge K_{1/2^j}$$
 for  $\epsilon \le 1/2^j$ .

(c) Bigger range, bigger integral — our integrand is non-negative.

### Neighborhoods

$$w(\mathcal{V}) \lesssim \frac{12}{\sqrt{n}} \int_0^1 \sqrt{\log(K_{\epsilon})} d\epsilon$$

- If all  $v \in \mathcal{V}$  are small, we don't have to integrate all the way to one.
- If we can cover  $\mathcal{V}$  with one ball of radius s, we're integrating zero for  $\epsilon \geq s$ .
- · For example, for our centered neighborhood  $\mathcal{V}=\mathcal{M}_s$  or its boundary.



$$\mathrm{w}(\mathcal{V}) \lesssim \frac{1}{\sqrt{n}} \int_0^s \sqrt{\log(K_\epsilon)} d\epsilon \quad \text{ for } \quad s := \max_{v \in \mathcal{V}} \lVert v \rVert.$$

### The Lipschitz Regression Case

$$\log(K_{\epsilon}) \lesssim 1/\epsilon \quad \text{ for } \quad \mathcal{M} = \big\{ f : \rho_{Lip}(f) \leq 1, \ |f| \leq 1 \big\}.$$

Integrating, we can bound the width of a neighborhood

$$w(\mathcal{M}_s) \lesssim \frac{1}{\sqrt{n}} \int_0^s \sqrt{\log(K_\epsilon)} d\epsilon \lesssim \frac{1}{\sqrt{n}} \int_0^s \sqrt{\frac{1}{\epsilon}} d\epsilon = \frac{1}{\sqrt{n}} 2\sqrt{\epsilon} \mid_0^s = 2\sqrt{\frac{s}{n}}.$$

And solve for the radius s for a least squares estimator

$$s^2 \gtrsim \mathrm{w}(\mathcal{M}_s)$$
 for  $s^{-3/2} \approx n^{-1/2}$  i.e.  $s \approx n^{-1/3}$ .

This agrees with what we see based on Fourier series.

## Chaining and Gaussian Width in General

- This isn't just another bound it's pretty tight.
- This bound with  $K_{\epsilon}$  the size of the smallest  $\epsilon$ -cover can barely be improved.
- It's off by at most a factor of log(n). Proving it isn't so hard.
- See Vershynin [2018, Chapter 8.1.2] if you're interested.

$$\frac{1}{\sqrt{n}\log(n)}\int_0^1 \sqrt{\log(K_\epsilon)}d\epsilon \lesssim \mathrm{w}(\mathcal{V}) \lesssim \frac{1}{\sqrt{n}}\int_0^1 \sqrt{\log(K_\epsilon)}d\epsilon$$

- In fact, if we're a bit more careful about how we choose  $\pi_k(v)$ , chaining gives us a bound that's off by no more than a constant factor.
- · This fancier chaining is pretty straightforward conceptually.
- We just do the bound thinking of  $\pi_k(v)$  as an arbitrary function taking on  $2^{2^k}$  distinct values, then minimize the chaining bound over all the  $\pi_k$ .
- · It's easy to prove this is no worse than what we've talked about.
- But proving it's tight up to constants is a feat. See Talagrand [2014].

### **Chaining and Fourier Series**

Chaining is, in a sense, approximating our analysis using Fourier series.

- Using Fourier series, we were able to decompose the functions in Sobolev models into combinations of orthogonal functions.
- · There were infinitely many such functions, but only a few were allowed to be big.

$$\{m = \sum_j m_j \phi_j : \sum_j m_j^2 \lambda_j \le B\} \implies \|m_j \phi_j\|_{L_2} = m_j \le B / \sqrt{\lambda_j}.$$

The links in our chains play the role of the Fourier basis functions  $\phi_i$ .

- These links,  $\phi_{j,v}(x) = \{\pi_j(v) \pi_{j-1}(v)\}(x)$ , are approximately orthogonal.
  - for different resolutions j
  - $\cdot$  for the same resolution and different v- unless they're the same curve.
  - i.e. unless  $\pi_j(v) = \phi_j(v')$  and  $\pi_{j-1}(v) = \pi_{j-1}(v')$ , so  $\phi_{j,v} = \phi_{j,v'}$ .
- $\boldsymbol{\cdot}$  And as a result, the corresponding gaussians are approximately uncorrelated.

$$E\langle g, u\rangle \langle g, v\rangle = \frac{1}{n^2} \sum_{ij} u_i v_j E g_i g_j = \frac{1}{n^2} \sum_{i=1}^n u_i v_j = \frac{\langle u, v\rangle}{n}.$$

and therefore

$$\operatorname{Cov}\left\{\left\langle g, \, \pi_{j}(u) - \pi_{j-1}(u)\right\rangle, \, \left\langle g, \, \pi_{j'}(v) - \pi_{j'-1}(v)\right\rangle\right\}$$

$$= \frac{\left\langle \pi_{j}(v) - \pi_{j-1}(v), \, \pi_{j'}(u) - \pi_{i-1}(u)\right\rangle}{r}.$$

#### The Point

- The point is to make sure that when we use the union bound, we're not being wasteful and bounding more-or-less the same thing twice.
- Decomposing the curves in our model into sums of approximately orthogonal functions helps us keep track of what we're bounding more accurately.
- It helps us not overcount when we're bounding gaussian width.

Let's look into how orthogonal our links are.

# $\pi_k(v)$ for our Lipschitz cover

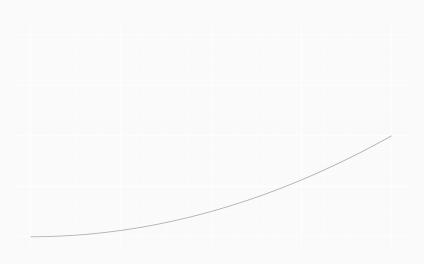
- 1. Draw an  $\epsilon_k \times \epsilon_k$  grid.
- 2. Snap v(x) to it at each x on the grid.
- 3. Piecewise-linear between grid points.

Use the small squares for  $\pi_{i+1}$ , two for  $\pi_i$ , and four for  $\pi_{i-1}$ .



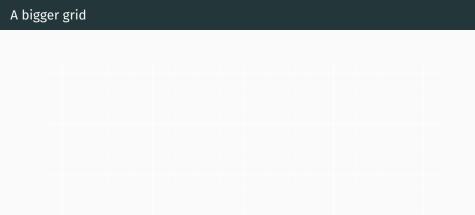
Check the inner product between links  $\ell_j(v)=\pi_j(v)-\pi_{j-1}(v)$ . Do it both for different j and different curves.

# A bigger grid



## You can try it for more curves and more resolutions

- Do it by hand on the blank grid on the next slide.
- $\boldsymbol{\cdot}$  Or code it up in R so you can try more stuff.



#### References

Michel Talagrand. *Upper and lower bounds for stochastic processes*, volume 60. Springer, 2014.

Roman Vershynin. *High-dimensional probability: An introduction with applications in data science*, volume 47. Cambridge university press, 2018.