for electrical conductivity & Drift anductal which is but extent to an E.F of sirength E n= mass of els on= mass of els e: charge of els According to Newton's second law of molicin the force & acquired by ells is equal to the force excerted by the field on the ell, i egn of motion ma = -eE

so fadt = J-et dt (Integrate [a= gy] deity > V = - et t FC Integration const. During the absence of the E.F., the average velocity of the ed is zero, t=0, LV>=0 > [V=-eEt] aurage relocity for two successive collisions V = feft)dt $\overline{V} = -eE \int t dt = -eE \overline{Z}^2 = -eEZ$ relaxation time of w two Successive collisions = Z 1 = mean free parts of V = - e E 1 2 (3KT) [: 2 mv = 3 KT = $mV = \frac{3kT}{V}$ V = - EEAV _ (D) Of n is no density of Os in conductor then the aurrent density Jesis! - who government density J=-BenV · cov [Fol untales of = -en [-eEAV 1 from O J = netEdV

flowing through a conductor

A in time t If a charge is Cooss-Section area 9 = or AEt $\frac{q}{T} = \sigma A \mathcal{E}$ = $\sigma A \in \Xi$ [current density is directly for tional to applied E.F.] T = I [For writ area of cross-see.] o = ne21v

Different conductivities of diff. materials are due to different no. of feel electrons.

(i) Electrical conductivity

It is defined as the quantity of electricity that flows in unit time per unit area of cross-section of the conductor per unit potential gradient.

According to free electron theory, in a solid the electrons move freely. If E is the applied electric field, then the acceleration of an electron having charge is given by

$$a = \frac{d^2x}{dt^2} = \frac{eE}{m}$$
 ... (5.1)

If λ is the mean free path of electrons, then the relaxation time τ between two successive collisions is given by

$$\tau = \frac{\lambda}{\nu} \qquad (5.2)$$

Integrating eq. 5.1, we get

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$$\frac{dx}{dt} = \frac{e\mathbf{E}}{m}t + \mathbf{C}$$

$$t = 0, \frac{dx}{dt} = 0$$

At $\frac{dx}{dt} = v = \frac{dx}{dt}$

So average velocity between two successive collisions

$$\bar{v} = \frac{1}{\tau} \int_{0}^{eE} t \, dt = \frac{eE}{\tau m} \int_{0}^{t} t \, dt$$

$$= \frac{eE}{\tau m} \frac{\tau^{2}}{2}$$

$$\bar{v} = \frac{eE}{\tau} \frac{2b^{2}}{2}$$

ot

Putting the value of τ from eq. 5.2, we get

Since
$$\overline{v} = \frac{eE\lambda}{2mv}$$

$$\frac{1}{2}mv^2 = \frac{3}{2}kT$$

$$mv = \frac{3kT}{v}$$

$$\overline{v} = \frac{eE\lambda v}{6kT}$$
So
$$\overline{v} = \frac{eE\lambda v}{6kT}$$
If *n* is number density of electrons in the conductor, then the current density *i* is given by
$$i = ne\overline{v}$$

$$ve^2E\lambda v$$
(7 is absolute temperature and *k* is
Boltzmann constant)

$$i = \frac{ne^2 E \lambda v}{6kT}$$
 ... (5.3)

or

or

or

Since

If q charge is flowing through a conductor of cross-section area A in time t, then

$$q = \sigma A E t$$

$$\frac{q}{t} = \sigma AE$$

$$i = \sigma AE$$

$$\sigma = \frac{i}{AE}$$

$$o = \overline{AE}$$

For unit area of cross-section

$$\sigma = \frac{i}{E}$$

$$\sigma = \frac{ne^2\lambda v}{6kTA}$$



This expression shows that different conductivities of different materials are due to different number of free electrons.

(ii) Ohm's law

From eq. (5.4) we have
$$\sigma E = \frac{I}{A}$$
or
$$\sigma E = J$$
or
$$J = \sigma E$$
or This is microscopic form of Ohm's law.

...(5.5)

...(5.4)