

Bloch theorem :- States that the energy eigen states for an electron in a crystal can be written as Bloch waves.

• If a wave is a Bloch wave, its wave function can be described as

$$\psi(r) = \underbrace{e^{i\vec{k} \cdot \vec{r}}}_{\text{Plane wave}} \underbrace{u(r)}_{\text{periodic Bloch cell function}}$$

Eigenstates (ψ)

one dimensional Schrodinger wave eq. for a free e^- which is moving in a constant potential

$$\frac{d^2 \psi}{dx^2} + (2m/\hbar^2)(E - V_0)\psi = 0$$

~~solution of above differential eqn~~

$$\psi = \exp(\pm i k x)$$

K.E. of e^- s $E_{KE} = E - V_0 = \frac{p^2}{2m} = \frac{\hbar^2 k^2}{2m}$ [$p = \hbar k$]

~~To make the wave function~~

Thus, we will consider a periodic solid through Bloch's theorem.

→ Bloch's theorem is just a way to describe the wave function for periodic solids.

Postulates of Bloch's theorem :-

Bloch ① As we have a solid that is periodic at the atomic scale, we get a travelling wave solution ($e^{i\vec{k} \cdot \vec{r}}$) for ψ that is modulated by

the translational symmetry of the lattice ^{free} _{sim}

$$(u_{\vec{k}})$$

$$\rightarrow \psi_{\vec{k}}(\vec{r}) = u_{\vec{k}}(\vec{r}) e^{i\vec{k} \cdot \vec{r}} \quad \text{--- (1)}$$

$$\text{where } u_{\vec{k}}(\vec{r}) = u_{\vec{k}}(\vec{r} + \vec{T}) \quad \text{--- (2)}$$

\rightarrow wavefunction reflects the symmetry of the lattice $u_{\vec{k}}$

Bloch 2) we can describe the wavefunction

$$\psi_{\vec{k}}(\vec{r} + \vec{T}) = C \psi_{\vec{k}}(\vec{r})$$

\Rightarrow wavefunction at $(\vec{r} + \vec{T})$ is very related to the wavefunction at \vec{r} .

where $C =$ modulation term.

Derivation of Bloch-2 using Bloch-1 :-

We are going to use Bloch-1 postulate,

In eqn (1), replace r with $r + T$ because $u_{\vec{k}}$ is

$$\psi_{\vec{k}}(\vec{r} + \vec{T}) = e^{i\vec{k} \cdot (\vec{r} + \vec{T})} u_{\vec{k}}(\vec{r} + \vec{T}) \quad \leftarrow \begin{matrix} \text{Periodic in } T \\ \text{[eqn (2)]} \end{matrix}$$

$$= e^{i\vec{k} \cdot (\vec{r} + \vec{T})} \frac{\psi_{\vec{k}}(\vec{r})}{e^{i\vec{k} \cdot \vec{r}}} \quad \left[\begin{matrix} \text{combining (1)} \\ \text{& (2) expression} \end{matrix} \right]$$

$$\Rightarrow \boxed{\psi_{\vec{k}}(\vec{r} + \vec{T}) = e^{i\vec{k} \cdot \vec{T}} \psi_{\vec{k}}(\vec{r})}$$

Implementation of Bloch: The vanishing potential :-

① If the potential is periodic but so weak that it is vanishing.

\Rightarrow we can estimate our potential as a flat box & our dispersion will basically be parabolic like the

Free electron model.

→ Since we still have periodic conditions,

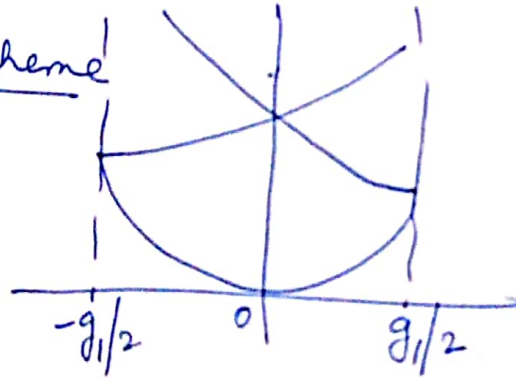
$$\psi_{\vec{k}}(\vec{r}) = \psi_{\vec{k}+\vec{G}}(\vec{r})$$

⇒ we repeat our parabolas at $k+g, k-g, k+2g$
& so on for the dispersion [where k is shift to $+g$ or $-g$ (in 1-D crystal)]

→ Call one parabola centred at $k=0$ (the extended zone scheme)

→ & multiple parabolas (the periodic zone scheme)

→ Reduced zone scheme



1st Brillouin zone
is b/w the dotted
lines.

Physical Interpretation of Bloch's Theorem :-

Wavefunction is periodic at two length scales:

- U is periodic in T (the atomic scale, lattice spacing)
- $e^{i\vec{k}\cdot\vec{r}}$ is periodic in L (sample scale)

