3. PRINCIPLES OF QM

Axiomatic principles

State vector axiom: State vector at t is ket $\psi(t)$, or $|\psi\rangle$.

Probability axiom: Given a system in state $|\psi\rangle$, a measurement will find it in state $|\phi\rangle$ with probability amplitude $\langle\phi|\psi\rangle$.

Hermitian operator axiom: Physical observable is represented by a linear and Hermitian operator.

Measurement axiom: Measurement of a physical observable results in eigenvalue of observable. Observable $\widehat{A}$, we have $\widehat{A}|a\rangle = a|a\rangle$, where a is eigenvalue and $|a\rangle$ is eigenvector. Measurement of the physical quantity represented by $\widehat{A}$ collapses the state $|\psi\rangle$ before measurement into an eigenstate $|a\rangle$ of $\widehat{A}$.

Time evolution axiom: $i\hbar\frac{\partial}{\partial t}|\psi(t)\rangle = \widehat{H}|\psi(t)\rangle$, w/o consider x or p.

Vector space

State vector is neither in position nor momentum space. Basis vectors: $|0\rangle, |1\rangle, |n\rangle$

Linearity : Because the SE is linear, given two states $|\psi_1(t)\rangle$ and $|\psi_2(t)\rangle$, $|\psi(t)\rangle = c_1|\psi_1(t)\rangle + c_2|\psi_2(t)\rangle$ is also a sol. (c's are complex).

Dual vector space $c|\psi\rangle$ is mapped to $c^*\langle\psi|$. Given a vector, $|\psi\rangle = \begin{bmatrix}\vdots\\ \alpha\\ \vdots\end{bmatrix}$, the dual vector is $\langle\psi| = [\cdots \quad \alpha^* \quad \cdots]$.

Dual basis vectors are $\langle 0| = \begin{bmatrix}1 & 0 & \cdots\end{bmatrix}, \cdots, \langle n| \begin{bmatrix}0 & \cdots & 1\end{bmatrix}$.

Inner product : $\langle\phi|\psi\rangle = c$, where c is complex.

$\langle\phi|\psi\rangle = \langle\psi|\phi\rangle^* \to \langle\psi|\psi\rangle$ is real, positive, and finite for a normalizable ket vector. Can choose $\langle\psi|\psi\rangle = 1$. $\langle\psi_m|\psi_n\rangle = \delta_{mn}$

Operators

A matrix operator $\widehat{A}$ acting on a state vector $|\psi\rangle$ transforms it into another state vector $|\phi\rangle$, $\widehat{A}|\psi\rangle = |\phi\rangle$. It is linear.

Hermitian conjugate (Hermitian adjoint) operator in the dual space

Hermitian adjoint operator $\widehat{A}^\dagger$ acts on the dual vector $\langle\psi|$ from the right as $\langle\psi|\widehat{A}^\dagger\rangle$, where $\widehat{A}^\dagger = (\widehat{A})^{T*}$.

$$(\widehat{A}|\psi\rangle)^\dagger = |\psi\rangle^\dagger\widehat{A}^\dagger = \langle\psi|\widehat{A}^\dagger \qquad \langle\psi| = |\psi\rangle^\dagger \qquad \langle\psi|^\dagger = |\psi\rangle$$

$$(\widehat{A}\widehat{B})^\dagger = (\widehat{A}\widehat{B})^{T*} = (\widehat{B}^T\widehat{A}^T)^* = \widehat{B}^{T*}\widehat{A}^{T*} = \widehat{B}^\dagger\widehat{A}^\dagger, \quad (c\widehat{A})^\dagger = c^*\widehat{A}^\dagger$$

Outer product operators : $|\psi\rangle\langle\phi| \qquad [|\psi\rangle\langle\phi|]\chi\rangle = |\psi\rangle\langle\phi|\chi\rangle$

Matrix elements of operators $\langle\phi|\widehat{A}|\psi\rangle$ (complex num)

Hermitian equiv to complex conj $\langle\phi|\widehat{A}|\psi\rangle^\dagger = \langle\psi|\widehat{A}^\dagger|\phi\rangle = \langle\phi|\widehat{A}|\psi\rangle^*$

Hermitian operators : $\widehat{A}^\dagger = \widehat{A}$, so given $\widehat{A}|\phi\rangle$ in the vector space, we have $\langle\psi|\widehat{A}^\dagger = \langle\phi|\widehat{A}$ in the dual vector space.

Matrix elements of a Hermitian operator

$\langle\phi|\widehat{A}|\psi\rangle^\dagger = \langle\phi|\widehat{A}|\psi\rangle^* = \langle\psi|\widehat{A}^\dagger|\phi\rangle = \langle\psi|\widehat{A}|\phi\rangle$

Hermitian operator, real expectation vals: $\langle\psi|\widehat{A}|\phi\rangle^* = \langle\psi|\widehat{A}|\phi\rangle \equiv \langle\widehat{A}\rangle$

Same result whether $\widehat{A}$ acts to right or left: $\langle\phi|\widehat{A}|\psi\rangle = \langle\phi|\widehat{A}^\dagger|\psi\rangle$

Eigenvals and eigenvecs of Hermitian operators : $\widehat{A}|a_n\rangle = a_n|a_n\rangle$

Normalized eigvecs $\langle a_m|a_n\rangle = \delta_{mn}$. Gram-Schmidt, degenerate evec.

Completeness of eigenvector of a Hermitian operator Set $|a_n\rangle$ is complete if $\sum_n|\langle a_n|\psi\rangle|^2 = 1$. $\sum_n|a_n\rangle\langle a_n| = 1$ (identity operator)

Continuous spectra of a Hermitian operator Hermitian operator $\widehat{A}$, $\widehat{A}|a\rangle = a|a\rangle$, where a is continuous.

$\int da'\langle a'|\widehat{A}|a\rangle = a\int da'\langle a'|a\rangle = \int da' a'\langle a'|a\rangle \to \langle a'|a\rangle = \delta(a'-a)$

Continuous condition: $\int da|a\rangle\langle a| = 1$

Gram-Schmidt orthogonalization procedure Eigval (like energy level) is n-fold degenerate: n states w same eigval.

Orthogonal eigenstates $\to$ no degeneracy.

1. Normalize each state and define $\alpha_i = \frac{\alpha_i}{\sqrt{\langle a_i|a_i\rangle}}$. 2. $|\alpha_1'\rangle = |\alpha_1\rangle$.

3. $|\alpha_2'\rangle = \frac{|\alpha_2\rangle - |\alpha_1\rangle\langle\alpha_1|\alpha_2\rangle}{\sqrt{\langle\alpha_2|\alpha_2\rangle - \langle\alpha_1|\alpha_2\rangle\langle\alpha_2|\alpha_1\rangle}} = \frac{|\alpha_2\rangle - |\alpha_1\rangle\langle\alpha_1|\alpha_2\rangle}{\sqrt{1-\langle\alpha_1|\alpha_2\rangle\langle\alpha_2|\alpha_1\rangle}}$

4. Subtract components of $|\alpha_3\rangle$ along $|\alpha_1\rangle$ and $|\alpha_2\rangle$, $|\alpha_3\rangle - |\alpha_1\rangle\langle\alpha_1|\alpha_3\rangle - |\alpha_2\rangle\langle\alpha_2|\alpha_3\rangle$, normalize and promote to $|\alpha_3'\rangle$. ...

Position and momentum representation

$\widehat{\vec{r}}|\vec{r}\rangle = \vec{r}|\vec{r}\rangle \quad \langle\vec{r'}|\vec{r}\rangle = \delta^3(\vec{r'}-\vec{r}), \int d^3\vec{r}|\vec{r}\rangle\langle\vec{r}| = 1, \langle\vec{r'}|\widehat{\vec{r}}|\vec{r}\rangle = \vec{r}\delta^3(\vec{r'}-\vec{r})$

$\widehat{\vec{p}}|\vec{p}\rangle = \vec{p}|\vec{p}\rangle \quad \langle\vec{p'}|\vec{p}\rangle = \delta^3(\vec{p'}-\vec{p}), \int d^3\vec{p}|\vec{p}\rangle\langle\vec{p}| = 1$

State vector $|\psi(t)\rangle$ in position space (scalar): $\langle\vec{r}|\psi(x,t)\rangle \equiv \psi(\vec{r},t)$

$\langle\psi|\widehat{\vec{p}}|\psi\rangle = \frac{d}{dt}\langle\psi|\widehat{\vec{r}}|\psi\rangle m$

Representation of momentum operator in position space: $\widehat{\vec{p}} = -i\hbar\vec{\nabla}$.

$\langle x|\widehat{p}|x'\rangle = -i\hbar\frac{\partial}{\partial x}\delta(x-x') = -i\hbar\frac{\partial}{\partial x}\langle x|x'\rangle$.

$\widehat{p} = -i\hbar\frac{\partial}{\partial x}$ is Hermitian, $\frac{\partial}{\partial x}$ is not.

$\langle x|\widehat{p}|p\rangle = p\langle x|p\rangle = -i\hbar\frac{\partial}{\partial x}\langle x|p\rangle$. The solution is $\langle x|p\rangle = \frac{1}{\sqrt{2\pi\hbar}}e^{\frac{i}{\hbar}px}$.

In 3D, $\langle\vec{r}|\vec{p}\rangle = \frac{1}{(2\pi\hbar)^{3/2}}e^{\frac{i}{\hbar}\vec{p}\vec{r}}$.

We can write the normalized wavefunction of definite position in momentum space, $\langle p|x\rangle = \langle x|p\rangle^*$. So, $\langle p|x\rangle = \frac{1}{\sqrt{2\pi\hbar}}e^{-\frac{i}{\hbar}px}$ (particle moving to the left, or with momentum $-p$, in the momentum space). $[x,p] = i\hbar$

Operators and wavefunction in position representation Position and momentum operators in pos space: $\widehat{\vec{r}} = \vec{r}$, $\widehat{\vec{p}} = -i\hbar\vec{\nabla}$.

$\widehat{\vec{r}}$ is Hermitian and $\langle\phi|\widehat{\vec{r}}^\dagger|\psi\rangle = \langle\phi|\widehat{\vec{r}}|\psi\rangle$.

$\widehat{O}(\widehat{\vec{r}},\widehat{\vec{p}}) = \widehat{O}(\vec{r},-i\hbar\vec{\nabla})$

The expectation val of the observable should be indep of representation. In state $\psi(t)$, $\langle\widehat{O}\rangle = \langle\psi(t)|\widehat{O}|\psi(t)\rangle$.

Insert $\int d^2\vec{r}|\vec{r}\rangle\langle\vec{r}| = 1$ to get $\langle\widehat{O}\rangle = \int d^2\vec{r}\langle\psi(t)|\vec{r}\rangle\langle\vec{r}|\widehat{O}|\psi(t)\rangle$

$\psi(\vec{r},t) = \langle\vec{r}|\psi(t)\rangle, \qquad \psi(\vec{r},t)^* = \langle\vec{r}|\psi(t)\rangle^* = \langle\psi(t)|\vec{r}\rangle$,

$\langle\vec{r}|\widehat{O}|\psi(t)\rangle = \widehat{O}(\vec{r},-i\hbar\vec{\nabla})\psi(\vec{r},t), \langle\widehat{O}\rangle = \int d^3\vec{r}\psi(\vec{r},t)^*\vec{O}(\vec{r},-i\hbar\vec{\nabla})\psi(\vec{r},t)$

Operators and wavefunction in momentum representation $\widehat{\vec{r}} = i\hbar\vec{\nabla}_{\vec{p}}$, or in 1D, $\widehat{x} = i\hbar\frac{\partial}{\partial p}$, $\widehat{\vec{p}} = \vec{p}$, where $\vec{p}^* = \vec{p}$.

$\widehat{O}(\widehat{\vec{r}},\widehat{\vec{p}}) = \widehat{O}(i\hbar\vec{\nabla}_{\vec{p}},\vec{p})$

$\langle\widehat{O}\rangle = \langle\psi(t)|\widehat{O}|\psi(t)\rangle \to \langle\widehat{O}\rangle = \int d^2\vec{p}\langle\psi(t)|\vec{p}\rangle\langle\vec{p}|\widehat{O}|\psi(t)\rangle$.

$\psi(\vec{p},t) = \langle\vec{p}|\psi(t)\rangle, \qquad \psi(\vec{p},t)^* = \langle\vec{p}\psi(t)\rangle^* = \langle\psi(t)|\vec{p}\rangle$

$\langle\vec{p}|\widehat{O}|\psi(t)\rangle = \widehat{O}(i\hbar\vec{\nabla}_{\vec{p}},\vec{p}), \langle\widehat{O}\rangle = \int d^3\vec{p}\psi(\vec{p},t)^*\widehat{O}(i\hbar\vec{\nabla}_{\vec{p}},\vec{p})\psi(\vec{p},t)$.

$i\hbar\frac{\partial}{\partial t}|\psi(t)\rangle = \widehat{H}|\psi(t)\rangle$, where $\widehat{H} = \frac{\widehat{\vec{p}}^2}{2m} + V(\widehat{\vec{r}},t)$ becomes

$i\hbar\frac{\partial\psi(\vec{r},t)}{\partial t} = -\frac{\hbar^2}{2m}\vec{\nabla}^2\psi(\vec{r},t) + V(\vec{r},t)\psi(\vec{r},t)$

Commuting operators

If $[\widehat{A},\widehat{B}] = 0$ and the states are nondegenerate, $|\psi\rangle$ is a simultaneous eigenstate of $\widehat{A}$ and $\widehat{B}$.

$|\psi\rangle = |ab\rangle$, and $\widehat{A}|ab\rangle = a|ab\rangle$, $\widehat{B}|ab\rangle = b|ab\rangle$

Non-commuting operators and the general uncertainty principle

$(\Delta A)^2(\Delta B)^2 \geq (\frac{1}{2i}\langle[\widehat{A},\widehat{B}]\rangle)^2$

Cannot construct simulatneous eigenstates (which correspond to definite eigenvalues) of non-commuting observables.

Time evolution of expectation value of an operator and Ehrenfest's theorem

Ehrenfest's theorem: how observable $\widehat{O}$'s expectation value in state $|\psi(t)\rangle$ evolves in time, $\frac{d}{dt}\langle\widehat{O}\rangle = \langle\frac{\partial\widehat{O}}{\partial t}\rangle + \frac{i}{\hbar}\langle[\widehat{H},\widehat{O}]\rangle$. If operator has no explicit time dep, $\frac{d}{dt}\langle\widehat{O}\rangle = \frac{1}{i\hbar}\langle[\widehat{O},\widehat{H}]\rangle$.

For $\widehat{O} = \widehat{\vec{p}}$ and a Hamiltonian that is TI, $\frac{d}{dt}\langle\widehat{\vec{p}}\rangle = -\langle\vec{\nabla}V(\widehat{\vec{r}})\rangle$, which is just Newton's Second Law! $\to$ QM contains all of classical mech.

The simple harmonic oscillator

$\widehat{H} = \frac{\widehat{p}^2}{2m} + \frac{1}{2}m\omega^2\widehat{x}^2$

Raising and lowering operators Lowering op: $\widehat{a} = \sqrt{\frac{m\omega}{2\hbar}}(\widehat{x} + \frac{i}{m\omega}\widehat{p})$, Raising op: $\widehat{a}^\dagger = \sqrt{\frac{m\omega}{2\hbar}}(\widehat{x} - \frac{i}{m\omega}\widehat{p})$.

$[\widehat{a},\widehat{a}^\dagger] = 1 \qquad \widehat{x} = \sqrt{\frac{\hbar}{2m\omega}}(\widehat{a}^\dagger + \widehat{a}), \widehat{p} = i\sqrt{\frac{m\omega\hbar}{2}}(\widehat{a}^\dagger - \widehat{a})$

$\widehat{H} = (\widehat{N} + \frac{1}{2})\hbar\omega$, where $\widehat{N} = \widehat{a}^\dagger\widehat{a}$. Now $\widehat{N}$ is Hermitian, and $\widehat{N}|n\rangle = n|n\rangle$.

$[\widehat{N},\widehat{a}] = -\widehat{a}$, $[\widehat{N},\widehat{a}^\dagger] = \widehat{a}^\dagger$

$\widehat{N}(\widehat{a}|n\rangle) = (n-1)(\widehat{a}|n\rangle)$, $\widehat{N}(\widehat{a}^\dagger|n\rangle) = (n+1)(\widehat{a}^\dagger|n\rangle)$

Normalized number state vectors Energy levels are not degenerate, so $|n-1\rangle = c_n\widehat{a}|n\rangle \to c_n = \frac{1}{\sqrt{n}} \to \widehat{a}|n\rangle = \sqrt{n}|n-1\rangle$.

$|n+1\rangle = d_n\widehat{a}^\dagger|n\rangle \to d_n = \frac{1}{\sqrt{n+1}} \to \widehat{a}^\dagger|n\rangle = \sqrt{n+1}|n+1\rangle$

Ground state: $|0\rangle$, excited state: $|n\rangle = \frac{(\widehat{a}^\dagger)^n}{\sqrt{n!}}|0\rangle$, $n = 0, 1, 2, \ldots$

$\langle n'|\widehat{x}|n\rangle = \sqrt{\frac{\hbar}{2m\omega}}\langle n'|(\widehat{a}^\dagger + \widehat{a})|n\rangle = \sqrt{\frac{\hbar}{2m\omega}}(\sqrt{n+1}\delta_{n',n+1} + \sqrt{n}\delta_{n',n-1})$

$\langle n'|\widehat{p}|n\rangle = i\sqrt{\frac{m\omega\hbar}{2}}\langle n'|(\widehat{a}^\dagger - \widehat{a})|n\rangle = i\sqrt{\frac{m\omega\hbar}{2}}(\sqrt{n+1}\delta_{n',n+1} - \sqrt{n}\delta_{n',n-1})$

Wavefunctions in position representation $E_n = (n+\frac{1}{2})\hbar\omega, n = 0, 1, 2, \ldots$

The stationary wavefunctions of definite energy: $\psi_n(x) = \langle x|n\rangle$

$\langle x'|\widehat{a}^\dagger|x''\rangle = \delta(x'-x'')\frac{1}{\sqrt{2}\sigma}(x'' - \sigma^2\frac{\partial}{\partial x''})$, where $\sigma \equiv \sqrt{\frac{\hbar}{m\omega}}$

$\xi = \frac{x}{\sigma}, \qquad \langle x|n\rangle = \frac{1}{\sqrt{\sqrt{\pi}n!2^n\sigma}}(\xi - \frac{\partial}{\partial\xi})^n e^{-\frac{1}{2}\xi^2}$

$\langle x|0\rangle = (\frac{m\omega}{\pi\hbar})^{1/4}e^{-\frac{m\omega}{2\hbar}x^2}, \qquad \langle x|1\rangle = \sqrt{2}(\frac{m^3\omega^3}{\pi\hbar^3})^{1/4}xe^{-\frac{m\omega}{2\hbar}x^2}$

Classical simple harmonic oscillator Hamiltonian of a simple harmonic is $H = \frac{p^2}{2m} + \frac{1}{2}m\omega^2x^2. \qquad \dot{x} = \frac{\partial H}{\partial p} = \frac{p}{m}, \qquad \dot{p} = -\frac{\partial H}{\partial x} = -m\omega^2x$

Define $\sqrt{\hbar\omega}\alpha = \sqrt{\frac{m\omega^2}{2}}x + \frac{i}{\sqrt{2m}}p$, so $x = \sqrt{\frac{2\hbar}{m\omega}}\alpha_R$ and $p = \sqrt{2m\hbar\omega}\alpha_I$

Rewrite Hamiltonian, $H = \hbar\omega|\alpha|^2, \qquad \dot{\alpha} = -i\omega\alpha$. The sol is $\alpha = \alpha_0e^{-i\omega t}$.

The quantum simple harmonic oscillator and coherent state Coherent state, superpos of stat states $|n\rangle$: $|\alpha\rangle = e^{-\frac{1}{2}|\alpha|^2}\sum_{n=0}^\infty\frac{\alpha^n}{\sqrt{n!}}|n\rangle$

$P(n) = |\langle n|\alpha\rangle|^2 = |\alpha_n|^2 = \frac{\langle n\rangle^n e^{-\langle n\rangle}}{n!}$, where $\langle n\rangle = \langle\alpha|a^\dagger a|\alpha\rangle = |\alpha|^2$.

Linear superpos. of all quantum nums which represent the class oscill the most. Has shape of Gaussian of min uncertainty satisfying $\Delta x\Delta p \geq \frac{\hbar}{2}$ regardless of value of energy. Oscillates like a class oscill, w only diff being that the particle's loc is not represented by a point (or a delta func) but by a Gaussian func.

4. 3D SYSTEMS

Three-dimensional infinite square well

$-\frac{\hbar^2}{2m}(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2})\psi(x,y,z) = E\psi(x,y,z)$ for $0 \leq x \leq l_x, \ldots$ while $\psi(x,y,z) = 0$ outside.

Separation of vars: $\psi(x,y,z) = \psi_1(x)\psi_2(y)\psi_3(z)$

$\to$ SE becomes $-\frac{\hbar^2}{2m}\frac{d^2}{dx^2}\psi_1(x) = E_1\psi_1(x), \ldots$, where $E = E_1 + E_2 + E_3$.

$\psi_{n_xn_yn_z}(x,y,z) = \sqrt{\frac{8}{l_xl_yl_z}}\sin\left(\frac{n_x\pi}{l_x}x\right)\sin\left(\frac{n_y\pi}{l_y}z\right)\sin\left(\frac{n_z\pi}{l_z}z\right)$

$E_{n_xn_yn_z} = \frac{\hbar^2\pi^2}{2m}(\frac{n_x^2}{l_x^2} + \frac{n_y^2}{l_y^2} + \frac{n_z^2}{l_z^2})$, with $n_x, n_y, n_z = 1, 2, \ldots$.

Wave vector: $\vec{k} = (k_x, k_y, k_z) = (\frac{n_x\pi}{l_x}, \frac{n_y\pi}{l_y}, \frac{n_z\pi}{l_z})$

The Schrödinger equation in spherical coordinates

$i\hbar\frac{\partial\psi(\vec{r},t)}{\partial t} = -\frac{\hbar^2}{2m}\vec{\nabla}^2\psi(\vec{r},t) + V(\vec{r})\psi(\vec{r},t)$, where $\vec{r} = (r,\theta,\phi)$,

$\psi(\vec{r},t) = \psi(r,\theta,\phi,t)$ and $\vec{\nabla}^2 = \frac{1}{r^2}\frac{\partial}{\partial r} + \frac{1}{r^2\sin\theta}\frac{\partial}{\partial\theta} + \frac{1}{r^2\sin^2\theta}\frac{\partial^2}{\partial\phi^2}$.

For a TI and central potential, potential depends only on r, $V(\vec{r}) = V(r)$.

$\frac{1}{R(r)}[\frac{d}{dr} - \frac{2mr^2}{\hbar^2}(V(r)-E)] = -\frac{1}{Y(\theta,\phi)}[\frac{1}{\sin\theta}\frac{d}{d\theta} + \frac{1}{\sin^2\theta}\frac{d^2Y(\theta,\phi)}{d\phi^2}]$

Each side must be constant and equal (let it be $l(l+1)$).

$\frac{1}{\sin\theta}\frac{d}{d\theta} + \frac{1}{\sin^2\theta}\frac{d^2Y(\theta,\phi)}{d\phi^2} = -l(l+1)Y(\theta,\phi)$

$\frac{d}{dr} - \frac{2mr^2}{\hbar^2}(V(r)-E) = l(l+1)R(r)$

Orbital angular momentum

$\widehat{L}_x = \widehat{y}\widehat{p}_z - \widehat{z}\widehat{p}_y, \widehat{L}_y = \widehat{z}\widehat{p}_x - \widehat{x}\widehat{p}_z, \widehat{L}_z = \widehat{x}\widehat{p}_y - \widehat{y}\widehat{p}_x$

$[\widehat{L}_i,\widehat{L}_j] = i\hbar\epsilon_{ijk}\widehat{L}_k$, with $i = 1, 2, 3$ representing the x, y, and z components, and $\epsilon_{123} = \epsilon_{231} = \epsilon_{312} = 1$, which is -1 for odd perms of indices, and vanishes when repeated.

$\widehat{\vec{L}}^2 = \widehat{L}_x^2 + \widehat{L}_y^2 + \widehat{L}_z^2, [\widehat{L}^2,\widehat{L}_i] = 0$

In pos rep, $\widehat{\vec{L}} = \widehat{\vec{r}}\times\widehat{\vec{p}} = -i\hbar\vec{r}\times\vec{\nabla}$. In sph coords,

$\widehat{\vec{L}} = -i\hbar r\widehat{r}\times(\frac{\partial}{\partial r}\widehat{r} + \frac{1}{r}\frac{\partial}{\partial\theta}\widehat{\theta} + \frac{1}{r\sin\theta}\frac{\partial}{\partial\phi}\widehat{\phi} = -i\hbar(\widehat{\phi}\frac{\partial}{\partial\theta} - \widehat{\theta}\frac{1}{\sin\theta}\frac{\partial}{\partial\phi})$

$\widehat{r} = \sin\theta\cos\psi\widehat{x} + \sin\theta\sin\phi\widehat{y} + \cos\theta\widehat{z}$

$\widehat{\theta} = \cos\theta\cos\phi\widehat{x} + \cos\theta\sin\phi\widehat{y} - \sin\theta\widehat{z} \qquad \widehat{\phi} = -\sin\phi\widehat{x} - \cos\phi\widehat{y}$

$\widehat{L}_x = i\hbar(\sin\theta\frac{\partial}{\partial\theta} + \cot\theta\cos\phi\frac{\partial}{\partial\phi}) \qquad \widehat{L}_y = i\hbar(-\cos\phi\frac{\partial}{\partial\theta} + \cot\theta\sin\phi\frac{\partial}{\partial\phi})$

$\widehat{L}_z = -i\hbar\frac{\partial}{\partial\phi} \qquad \widehat{\vec{L}}^2 = -\hbar^2[\frac{1}{\sin\theta}\frac{\partial}{\partial\theta} + \frac{1}{\sin^2\theta}\frac{\partial^2}{\partial\phi^2}]$

$\widehat{\vec{L}}^2 Y(\theta,\phi) = l(l+1)\hbar^2Y(\theta,\phi)$

$-\frac{\hbar^2}{2m}\frac{1}{r^2}\frac{d}{dr} - V_{\text{eff}}(r)R(r) = ER(r)$, $V_{\text{eff}}(r) = V(r) + \frac{l(l+1)\hbar^2}{2mr^2}$, centrifugal

Spherical harmonics Find sols to angular eqn. Sep vars $Y(\theta,\phi)=\Theta(\theta)\Phi(\phi)$.

$\frac{1}{\Theta}[\sin\theta\frac{\mathrm{d}}{\mathrm{d}\theta}+l(l+1)\sin^2\theta=-\frac{1}{\Theta}\frac{\mathrm{d}^2\Phi}{\mathrm{d}\phi^2}=\text{constant}=m^2$

$\Phi(\phi)=e^{im\phi}$, periodic in ϕ w period 2π gives constraint $m=0,\pm1,\pm2,\cdots$

$\Theta(\theta)$ can be written in terms of $x\equiv\cos\theta$ as

$(1-x^2)\frac{\mathrm{d}^2P(x)}{\mathrm{d}x^2}-2x\frac{\mathrm{d}P(x)}{\mathrm{d}x}+(l(l+1)-\frac{m^2}{1-x^2})P(x)=0$

Associated Legendre functions: $P_l^{m_l}(x)=(1-x^2)^{|m_l|/2}(\frac{\mathrm{d}}{\mathrm{d}x})^{|m_l|}P_l(x)$, where $P_l(x)$ is the l^{th} Legendre polynomial given by the Rodrigues formula $P_l(x)=\frac{1}{2^l l!}(\frac{\mathrm{d}}{\mathrm{d}x})^l(x^2-1)^l$, with l taking values $l=0,1,2,...$ and for each l, m_l takes $2l+1$ values $m_l=-l,-l+1,...,l-1,l$.

Spherical harmonics, normalized angular wave functions:

$Y_l^m(\theta,\phi)=\epsilon\sqrt{\frac{(2l+1)}{4\pi}\frac{(l-|m|)!}{(l+|m|)!}}e^{im\phi}P_l^m(\cos\theta)$, where $\epsilon=(-1)^m$ for $m\geq0$ and $\epsilon=1$ for $m\leq0$.

$\widehat{\vec{L}}^2Y_l^{m_l}=l(l+1)\hbar^2Y_l^{m_l}, \qquad \widehat{\vec{L}}_zY_l^{m_l}=m\hbar Y_l^{m_l}$

The Legendre polynomials are normalized s.t. they satisfy the ortho relation

$\int_{-1}^1P_{l'}P_l(x)dx=\int_0^\pi P_{l'}(\theta)P_l(\theta)\sin\theta d\theta=\frac{2}{2l+1}\delta_{l'l}$

$P_0^0(\theta)=1, P_1^1(\theta)=\sin\theta, P_1^1(\theta)=\cos\theta$, with $P_l^{-m_l}(x)=P_l^{m_l}(x)$

$\int_{-1}^1P_{l'}^{m'}(x)P_l^m(x)dx=\int_0^\pi P_{l'}^{m'}(\theta)P_l^m(\theta)\sin\theta d\theta=\frac{(l+m)!}{(2l+1)(l-m)!}\delta_{l'l}\delta_{m',m}$

Satisfy the orthogonality relation

$\int_0^{2\pi}d\phi\int_0^\pi d\theta\sin\theta Y_{l'}^{m_l'*}(\theta,\phi)Y_l^{m_l}(\theta,\phi)=\delta_{l'l}\delta_{m_l'm_l}$

$\widehat{\vec{L}}^2|lm_l\rangle=l(l+1)\hbar^2|lm_l\rangle$, $\widehat{\vec{L}}_z|lm_l\rangle=m\hbar|lm_l\rangle$

$\widehat{\vec{L}}_+=L_x+iL_y$, $\widehat{\vec{L}}_-L_x-iL_y$, $L_x=\frac{1}{2}(L_-+L_+)$, $\langle L_x^2\rangle=\frac{1}{2}\langle L^2-L_z^2\rangle$

$L_\pm|lm\rangle=\hbar\sqrt{(l\mp m)(l\pm m+1)}|l,m\pm1\rangle$

Spherical harmonics are the wavefunctions in pos rep, $Y_l^{m_l}(\theta,\phi)=\langle\vec{r}|lm_l\rangle$

Parity of the spherical harmonics

$\widehat{P}\psi(x,y,z)=\psi(-x,-y,-z), \qquad \widehat{P}\psi(r,\theta,\phi)=\psi(r,\pi-\theta,\phi+\theta)$

For the Legendre polynomials, $\widehat{P}P_l^{m_l}(\theta)=(-1)^{l-|m_l|}P_l^{m_l}(\theta)$

$\rightarrow$ even for $l+|m_l|$ even and odd for $l+|m_l|$ odd.

Azimuthal part of the wavefunction, $\widehat{P}e^{im_l\phi}=e^{im_l(\phi+\pi)}=(-1)^{m_l}e^{im_l\phi}$.

The spherical harmonics are products of two, and $\widehat{P}Y_l^{m_l}(\theta,\phi)=$
$Y_l^{m_l}(\pi-\theta,\phi+\pi)=(-1)^{l-|m_l|+m_l}Y_l^{m_l}(\theta,\phi)=(-1)^lY_l^{m_l}(\theta,\phi)$

The hydrogen atom

Coulomb's law, $\widehat{V}=-\frac{e^2}{4\pi\epsilon_0}\frac{1}{r}$

Let $u(r)\equiv rR(r)$, Radial eq: $-\frac{\hbar^2}{2m}\frac{\mathrm{d}^2u}{\mathrm{d}r^2}+[-\frac{e^2}{4\pi\epsilon_0}\frac{1}{r}+\frac{\hbar^2}{2m}\frac{l(l+1)}{r^2}]u=Eu$

The radial wave function

$\kappa\equiv\frac{\sqrt{-2mE}}{\hbar}$. Divide by E, $\frac{1}{\kappa^2}\frac{\mathrm{d}^2u}{\mathrm{d}r^2}=[1-\frac{me^2}{2\pi\epsilon_0\hbar^2\kappa}\frac{1}{(\kappa r)}+\frac{l(l+1)}{(\kappa r)^2}]u$

Introduce $\rho\equiv\kappa r$, $\rho_0\equiv\frac{me^2}{2\pi\epsilon\hbar^2\kappa}$, $\frac{\mathrm{d}^2u}{\mathrm{d}\rho^2}=[1-\frac{\rho_0}{\rho}+\frac{l(l+1)}{\rho^2}]u$

As $\rho\rightarrow\infty$, the constant term in the brackets dominates, so $\frac{\mathrm{d}^2u}{\mathrm{d}\rho^2}=u$.

General sol is $u(\rho)=Ae^{-\rho}+Be^{\rho}$, but $B=0\rightarrow u(\rho)=Ae^{-\rho}$ for large ρ.

As $\rho\rightarrow0$, centriugal term dominates, $\frac{\mathrm{d}^2u}{\mathrm{d}\rho^2}=\frac{l(l+1)}{\rho^2}u$

The general sol is $u(\rho)=C\rho^{l+1}+D\rho^{-l}$, but ρ^{-l} blows up as $\rho\rightarrow0$, so $D=0$. Thus, $u(\rho)\approx Cp^{l+1}$ for small ρ.

Peel off the asymptotic behavior, let $u(\rho)=\rho^{l+1}e^{-\rho}v(\rho)$

$\frac{\mathrm{d}u}{\mathrm{d}\rho}=\rho^le^{-\rho}[(l+1-\rho)v+\rho\frac{\mathrm{d}v}{\mathrm{d}\rho}]$

$\frac{\mathrm{d}^2u}{\mathrm{d}\rho^2}=\rho^le^{-\rho}\{[-2l-2+\rho+\frac{l(l+1)}{\rho}]v+2(l+1-\rho)\frac{\mathrm{d}v}{\mathrm{d}\rho}+\rho\frac{\mathrm{d}^2v}{\mathrm{d}\rho^2}\}$

Radial eq in terms of $v(\rho)$, $\rho\frac{\mathrm{d}^2v}{\mathrm{d}\rho^2}+2(l+1-\rho)\frac{\mathrm{d}v}{\mathrm{d}\rho}+[\rho_0-2(l+1)]v=0$

Assume $v(p)$ can be expressed as a power series in ρ: $v(\rho)=\sum_{j=0}^\infty c_j\rho^j$.

$\frac{\mathrm{d}v}{\mathrm{d}\rho}=\sum_{j=0}^\infty jc_j\rho^{j-1}=\sum_{j=0}^\infty(j+1)c_{j+1}\rho^j$,

$\frac{\mathrm{d}^2v}{\mathrm{d}\rho^2}=\sum_{j=0}^\infty j(j+1)c_{j+1}\rho^{j-1}$

$j(j+1)c_{j+1}+2(l+1)(j+1)c_{j+1}-2jc_j+[\rho_0-2(l+1)]c_j=0$

$c_{j+1}=\frac{2(j+l+1)-\rho_0}{(j+1)(j+2l+2)}c_j$

For large j (corresponding to large ρ), $c_{j+1}=\frac{2j}{j(j+1)}c_j=\frac{2}{j+1}c_j$

If this were exact, $c_j=\frac{2^j}{j!}c_0$, $v(\rho)=c_0\sum_{j=0}^\infty\frac{2^j}{j!}\rho^j=c_0e^{2\rho}$, and hence $u(\rho)=c_0\rho^{l+1}e^{\rho}$, which blows up at large ρ

$\exists c_{j_{\max}+1}=0$, so $2(j_{max}+l+1)-\rho_0=0$.

Define principle quantum number, $n\equiv j_{\max}+l+1$, so $\rho_0=2n$

$E=-\frac{\hbar^2\kappa^2}{2m}=-\frac{me^3}{8\pi^2\epsilon_0^2\hbar^2\rho_0^2}$

Bohr formula: $E_n=-[\frac{m}{2\hbar^2}(\frac{e^2}{4\pi\epsilon})^2]\frac{1}{n^2}=\frac{E_1}{n^2}=\frac{-13.6\text{ eV}}{n^2}$, $n=1,2,3,...$

$\kappa=(\frac{me^2}{4\pi\epsilon_0\hbar^2})\frac{1}{n}=\frac{1}{an}$, Bohr radius: $a\equiv\frac{4\pi\epsilon_0\hbar^2}{me^2}=0.529\times10^{-10}\mathrm{m}$

$\psi_{nlm}(r,\theta,\phi)=R_{nl}(r)Y_l^m(\theta,\phi)$, $\psi_{100}(r,\theta,\phi)=\frac{1}{\sqrt{\pi a^3}}e^{-r/a}$

For arbitrary n, $l=0,1,...,n-1$, so $d(n)=2\sum_{l=0}^{n-1}(2l+1)=2n^2$

$v(\rho)=L_{n-l-1}^{2l+1}(2\rho)$, where $L_{q-p}^p(x)\equiv(-1)^p(\frac{\mathrm{d}}{\mathrm{d}x})^pL_q(x)$ is an associated Laguerre polynomial. $L_q(x)\equiv e^x(\frac{\mathrm{d}}{\mathrm{d}x})^q(e^{-x}x^q)$ is the qth Lag. poly.

Normalized hydrogen wavefunctions:

$$\psi_{nlm}=\sqrt{(\frac{2}{na})^3\frac{(n-l-1)!}{2n[(n+1)!]^3}}e^{-r/na}(\frac{2r}{na})^l[L_{n-l-1}^{2l+1}(2r/na)Y_l^m(\theta,\phi)$$

Wavefunctions are mutually orthogonal.

$\int\psi*_{n'l'm_l'}\psi_{nlm_l}r^2\sin\theta drd\theta d\phi=\delta_{n'n}\delta_{l'l}\delta_{m_l'm_l}$

Spectrum Transitions: $E_\gamma=E_i-E_f=-13.6eV(\frac{1}{n_i^2}-\frac{1}{n_f^2})$

Planck formula, $E_\gamma=h\nu$, wavefunction is $\lambda=c/\nu$.

Rydberg: $\frac{1}{\lambda}=R(\frac{1}{n_f^2}-\frac{1}{n_i^2})$, $R\equiv\frac{m}{4\pi c\hbar^3}(\frac{e^2}{4\pi\epsilon_0})^2=1.097\times10^7\ \mathrm{m}^{-1}$

General angular momentum

$\widehat{\vec{J}}=(\widehat{J}_x,\widehat{J}_y,\widehat{J}_z)=(\widehat{J}_1,\widehat{J}_2,\widehat{J}_3) \qquad \widehat{\vec{J}}^2=\widehat{\vec{J}}_x^2+\widehat{\vec{J}}_y^2+\widehat{\vec{J}}_z^2$

$[\widehat{J}_i,\widehat{J}_j]=i\hbar\epsilon_{ijk}\widehat{J}_k$, $[\widehat{\vec{J}}^2,J_i]=0$

Take commuting set to be $\widehat{\vec{J}}^2$ and $\widehat{J}_z$. Trade $\widehat{J}_x$ and $\widehat{J}_y$ for $\widehat{J}_\pm=\widehat{J}_x\pm i\widehat{J}_y$

Commutation relations: $[\widehat{J}_+,\widehat{J}_-]=2\hbar\widehat{J}_z$, $[\widehat{J}_z,\widehat{J}_\pm]=\pm\hbar\widehat{J}_\pm$, $[\widehat{\vec{J}}^2,\widehat{J}_\pm]=0$

$\widehat{\vec{J}}^2$ and $\widehat{J}_z$ commute $\rightarrow$ we can simulaneously diagonalize them. Let the simultaneous eigenstate be $|ab\rangle$ s.t. $\widehat{\vec{J}}^2|ab\rangle=a|ab\rangle$, $\widehat{\vec{J}}_z|ab\rangle=b|ab\rangle$

$\widehat{\vec{J}}^2(\widehat{J}_\pm|ab\rangle)=a(\widehat{J}_\pm|ab\rangle \qquad \widehat{J}_z(\widehat{J}_\pm|ab\rangle)=(b\pm\hbar)(\widehat{J}_\pm|ab\rangle)$

$\widehat{J}_+$ raises and $\widehat{J}_-$ lowers the eigenvalue b of $\widehat{J}_z$. Assuming $|ab\rangle$ is normalized, $\widehat{J}_\pm|ab\rangle=c_\pm|ab\pm\hbar\rangle$, where $c_\pm$ are normalization constants.

$\widehat{J}_\pm\widehat{J}_\mp=\widehat{\vec{J}}^2-\widehat{J}_z^2\pm\hbar\widehat{J}_z$

$0=\langle ab_{\max}|\widehat{J}_-\widehat{J}_+|ab_{\max}\rangle=a-b_{\max}^2-\hbar b_{\max}$, $0=a-b_{\min}^2+\hbar b_{\min}$

$b_{\max}=\frac{-\hbar+\sqrt{\hbar^2+4a}}{2}$, $b_{\min}=\frac{\hbar-\sqrt{\hbar^2+4a^2}}{2}$, $b_{\max}=-b_{\min}=j\hbar$, $j=0,\frac{1}{2},1,\dots$

$j\equiv\frac{n}{2}$, then $a=b_{\max}^2+\hbar b_{\max}=j^2\hbar^2+\hbar^2j=j(j+1)\hbar^2$

$\widehat{J}_\pm|jm_j\rangle=\hbar\sqrt{(j\mp m_j)(j\pm m_j+1)}|jm_j\pm1\rangle$

$\langle j'm_j'|\widehat{J}_\pm|jm_j\rangle=\hbar\sqrt{(j\mp m_j)(j\pm m_j+1)}\langle j'm_j'|jm_j\pm1\rangle=$
$\hbar\sqrt{(j\mp m_j)(j\pm m_j+1)}\delta_{j'j}\delta_{m_j'm_j\pm1}$

Spin

Classical orbital and spinning motion Infinitesimal classical angular momentum corresponding to an infinite linear momentum $d\vec{p}=dm\vec{v}$ at position $\vec{r}$ from the axis of rotation is $d\vec{L}=\vec{r}\times d\vec{p}$

The total angular momentum is $\vec{L}=\int\vec{r}\times d\vec{p}=\int\vec{r}\times dm\vec{v}$

Point particle of mass m at radius r spinning w constant angular velocity ω about the z-axis, $\vec{L}=I\omega\widehat{z}=m\omega r^2\widehat{z}$

Considering a particle of mass m and charge q rotating with angular velocity ω at radius r about the z-axis, the angular momentum $\vec{L}$ and the momentum dipole momentum $\vec{\mu}$ are given by $\vec{L}=m\omega r^2\widehat{z}$, $\vec{\mu}=\frac{q}{2}\omega r^2\widehat{z}$, where we used $\mu=I\pi r^2$ with current $I=\frac{q}{2\pi/\omega}=\frac{q\omega}{2\pi}$. Thus, $\vec{\mu}=\frac{q}{2m}\vec{L}$

Spin Electron: $j=\frac{1}{2}$, $m_j=\pm\frac{1}{2}$. Spin-$\frac{1}{2}$: $s=\frac{1}{2}$, use $\widehat{\vec{J}}\rightarrow\widehat{\vec{S}}$.

Basis vectors are $|\frac{1}{2},\frac{1}{2}\rangle\equiv|1\rangle=\begin{bmatrix}1\\0\end{bmatrix}$, $|\frac{1}{2},-\frac{1}{2}\rangle\equiv|2\rangle=\begin{bmatrix}0\\1\end{bmatrix}$

$\widehat{S}_z$ and $\widehat{\vec{S}}^2$ are diagonal, since simultaneously diagonalized. Matrix elements:

$\langle s'm_s'|\widehat{\vec{S}}^2|sm_s\rangle=s(s+1)\hbar^2\delta_{s's}\delta_{m_s'm_s}$,

$\langle s'm_s'|\widehat{S}_z|sm_s\rangle=m_s\hbar\delta_{s's}\delta_{m_s'm_s}$

$\widehat{\vec{S}}^2=\frac{3}{4}\hbar^2\begin{bmatrix}1&0\\0&1\end{bmatrix}$, $\widehat{S}_z=\frac{\hbar}{2}\begin{bmatrix}1&0\\0&-1\end{bmatrix}$. $\widehat{S}_+=\hbar\begin{bmatrix}0&1\\0&0\end{bmatrix}$, $\widehat{S}_-=\frac{\hbar}{2}\begin{bmatrix}1&0\\0&-1\end{bmatrix}$,

$\widehat{S}_x=\frac{1}{2}(\widehat{S}_++\widehat{S}_-)$, $\widehat{S}_y=\frac{1}{2i}(\widehat{S}_+-\widehat{S}_-)$, $\widehat{S}_x=\frac{\hbar}{2}\begin{bmatrix}0&1\\1&0\end{bmatrix}$, $\widehat{S}_y=\frac{\hbar}{2}\begin{bmatrix}0&-i\\i&0\end{bmatrix}$

Spin angular momentum: $\vec{S}=\frac{\hbar}{2}\vec{\sigma}$. Pauli m: $\sigma_x=\begin{bmatrix}0&1\\1&0\end{bmatrix}$, $\sigma_y=\begin{bmatrix}0&-i\\i&0\end{bmatrix}$, $\sigma_z=\begin{bmatrix}1&0\\0&-1\end{bmatrix}$

$[\widehat{S}_i,\widehat{S}_j]=i\hbar\epsilon_{ijk}\widehat{S}_k$ and $[\sigma_i,\sigma_j]=2i\epsilon_{ijk}\sigma_k$

A general state of a spin-half system is given by a spinor,

$|\chi\rangle=\alpha|\frac{1}{2},\frac{1}{2}\rangle+\beta|\frac{1}{2},\frac{1}{2}\rangle=\begin{bmatrix}\alpha\\\beta\end{bmatrix}$, where α and β are complex constants.

$a_x=\sin\theta\cos\phi, a_y=\sin\theta\sin\phi, a_z=\cos\theta$

Magnetic moment of the electron $\vec{\mu}=g\frac{q}{2m}\vec{S}$, gyromagentic factor (distribution of mass != charge). For the electron, $q=-e$, and $\vec{\mu}=-g\frac{e}{2m}\vec{S}$

$\widehat{\vec{\mu}}=-g\frac{e}{2m}\widehat{\vec{S}}=-\frac{g}{2}\frac{e\hbar}{2m}\vec{\sigma}=-\frac{g}{2}\mu_B\vec{\sigma}$, where $\mu_B=\frac{e\hbar}{2m}$ is Bohr magneton.

Electron in a magnetic field Intrinsic spin angular momentum $\rightarrow$ intrinsic magnetic moment. Energy from spin & external mag field: $\widehat{H}=\widehat{V}=-\widehat{\vec{\mu}}\cdot\vec{B}$

For a magnetic field along the z-axis, $\vec{B}=B\widehat{z}$, and

$\widehat{H}=-\widehat{\mu}_zB=-(-\frac{g}{2}\frac{e}{m}\vec{S})\dot{B}\widehat{z}=\frac{g}{2}\frac{eB}{m}S_z=\omega_sS_z=\frac{g}{2}\frac{eB\hbar}{2m}\sigma_z$, where $\omega_s=\frac{g}{2}\frac{eB}{m}=\frac{g}{2}\omega_c$ is the spin precession (or Larmor) frequency and $w_c=\frac{eB}{m}$ is cyclotron frequency. $g\approx2$ but $g\neq2\rightarrow\omega_s\neq\omega_c$.

Rewrite Hamiltonian as $\widehat{H}=\omega_sS_z$. In the bases in which $\widehat{\vec{S}}$ and $\widehat{S}_z$ are diagonalized, the eigenstates are given by

$\widehat{H}|\frac{1}{2},\frac{1}{2}\rangle=\omega_s\widehat{S}_z|\frac{1}{2},\frac{1}{2}\rangle=\frac{1}{2}\hbar\omega_s|\frac{1}{2},\frac{1}{2}\rangle$,

$\widehat{H}|\frac{1}{2},-\frac{1}{2}\rangle=\omega_s\widehat{S}_z|\frac{1}{2},-\frac{1}{2}\rangle=-\frac{1}{2}\hbar\omega_s|\frac{1}{2},-\frac{1}{2}\rangle$

Interaction of electron spin w external magentic field $\rightarrow$ energies $\pm\frac{1}{2}\hbar\omega_s$.

Spin-up $|\frac{1}{2},\frac{1}{2}\rangle$ & spin-down state $|\frac{1}{2},-\frac{1}{2}\rangle$, with a gap of $\hbar\omega_s$ btwn them.

Consider $\vec{B}=B_x\widehat{e}_x+B_y\widehat{e}_y+B_z\widehat{e}_z$.

$\widehat{H}=(\frac{g}{2}\frac{e}{m}\vec{S})\cdot\vec{B}=\frac{g}{2}\frac{e\hbar}{2m}\begin{bmatrix}B_z&B_x-iB_y\\B_x+iB_y&-B_z\end{bmatrix}$

Eigenvals of matrix $\begin{vmatrix}B_z-\lambda&B_x-iB_y\\B_x+iB_y&-B_z-\lambda\end{vmatrix}=0$, which gives $\lambda=\pm B$, where $B=|\vec{B}|$. Therefore, eigenvals of $\widehat{H}$ are $\pm\frac{g}{2}\frac{e\hbar B}{2m}=\pm\frac{1}{2}\hbar\omega_s$.

The Stern-Gerlach experiment

Force on electron w spin-up: $\vec{F}_1=-\vec{\nabla}V_1=\frac{1}{2}\hbar\vec{\nabla}\omega_s=\frac{g}{2}\frac{e\hbar}{2m}\frac{\partial B(z)}{\partial z}$

Force on electron w spin-down: $\vec{F}_2=-\vec{\nabla}V_2=-\frac{1}{2}\hbar\vec{\nabla}\omega_s=-\frac{g}{2}\frac{e\hbar}{2m}\frac{\partial B(z)}{\partial z}$

Electrons deflected up/down depending on whether spin-up/spin-down along $\vec{B}$.

Spin precession $|\chi(0)\rangle=\begin{bmatrix}a\\b\end{bmatrix}$, $|a|^2+|b|^2=1$ and $a=\cos\frac{\alpha}{2}$, $b=\sin\frac{\alpha}{2}$

$|\chi(0)\rangle=\cos\frac{\alpha}{2}|\frac{1}{2}\frac{1}{2}\rangle+\sin\frac{\alpha}{2}|\frac{1}{2}-\frac{1}{2}\rangle=\begin{bmatrix}\cos\frac{\alpha}{2}\\\sin\frac{\alpha}{2}\end{bmatrix}$, $|\chi(t)\rangle=\begin{bmatrix}e^{-\frac{i}{2}\omega_st}\cos\frac{\alpha}{2}\\e^{\frac{i}{2}\omega_st}\sin\frac{\alpha}{2}\end{bmatrix}$

$\langle\widehat{S}_z\rangle=|e^{-\frac{i}{2}\omega_st}\cos\frac{\alpha}{2}|^2\frac{\hbar}{2}-|e^{-\frac{i}{2}\omega_st}\sin\frac{\alpha}{2}{}^2\frac{\hbar}{2}=(\cos^2\frac{\alpha}{2}-\sin^2\frac{\alpha}{2})\frac{\hbar}{2}$

$\langle\widehat{S}_x\rangle=\frac{\hbar}{2}\sin\alpha\cos\omega_st, \quad \langle\widehat{S}_y\rangle\frac{\hbar}{2}\sin\alpha\sin\omega_st$, $\langle\widehat{S}_z\rangle=\frac{\hbar}{2}\cos\alpha$

Angle $\alpha\rangle\pi-\alpha$ for spin-down. Spin-up, $\widehat{S}_z$ eigenval is $\frac{\hbar}{2}$, $|\widehat{\vec{S}}^2|$ is $\frac{\sqrt{3}\hbar}{2}$.

Space quantization: angular momentum along any fixed direction take only discrete $(2j+1)$ values.

Addition of angular momentum

$\widehat{\vec{J}}_1$, $|j_1,m_{j1}\rangle$. $\widehat{\vec{J}}_2$, $|j_2,m_{j2}\rangle$. $\widehat{\vec{J}}=\widehat{\vec{J}}_1+\widehat{\vec{J}}_2$. $\widehat{\vec{J}}^2$ & $\widehat{\vec{J}}_z$: sim diag set. $|j,m_j\rangle$

Triplet and singlet states of a system of two spin-halves

$|j_1,m_{j1}\rangle\otimes|j_2,m_{j2}\rangle$

The triplet states ($j=1$ multiplet): $|1,1\rangle=|\frac{1}{2},\frac{1}{2}\rangle\otimes|\frac{1}{2},\frac{1}{2}\rangle$,

$|1,0\rangle=\frac{1}{\sqrt{2}}(|\frac{1}{2},\frac{1}{2}\rangle\otimes|\frac{1}{2},-\frac{1}{2}\rangle+|\frac{1}{2},-\frac{1}{2}\rangle\otimes|\frac{1}{2},\frac{1}{2}\rangle)$,

$|1,-1\rangle=|\frac{1}{2},-\frac{1}{2}\rangle\otimes|\frac{1}{2},-\frac{1}{2}\rangle$

Singlet state ($j=0$): $|0,0\rangle=\frac{1}{\sqrt{2}}(|\frac{1}{2},\frac{1}{2}\rangle\otimes|\frac{1}{2},-\frac{1}{2}\rangle-|\frac{1}{2},-\frac{1}{2}\rangle\otimes|\frac{1}{2},\frac{1}{2}\rangle)$

$s=1,0$ out of s_1 and s_2 as $\frac{1}{2}\otimes\frac{1}{2}=1\oplus0$

$\widehat{\vec{J}}^2=\widehat{\vec{J}}_1^2\otimes1+1\otimes\widehat{\vec{J}}_2^2+2\widehat{\vec{J}}_{1z}\otimes\widehat{\vec{J}}_{2z}+\widehat{\vec{J}}_{1+}\otimes\widehat{\vec{J}}_{2-}+\widehat{\vec{J}}_{1-}\otimes\widehat{\vec{J}}_{2+}$

Spin angular momentum, interchan. use $\widehat{\vec{S}}$ for $\widehat{\vec{J}}$, and s and m_s for j and m_j.

Addition of general angular momentum

$|j_1+j_2,j_1+j_2\rangle=|j_1,m_{j1}\rangle\otimes|j_2,m_{j2}\rangle$

$j=j_1\otimes j_2=j_1+j_2\oplus j_1+j_2-1\oplus j_1+j_2-2\oplus\cdots\oplus|j_1-j_2|$

Clebsch-Gordon coefficients

Complete states: $\sum_{m_{j1},m_{j2}}|j_1,m_{j1};j_2,m_{j2}\rangle\langle j_1,m_{j1};j_2,m_{j2}|=1$

$|j,m_j\rangle=\sum_{m_j=m_{j1}+m_{j2}}\langle j_1,m_{j1};j_2,m_{j2}|j,m_j\rangle|j_1,m_{j1};j_2,m_{j2}\rangle$

where $\langle j_1,m_{j1};j_2,m_{j2};j,m_j\rangle$ are Clebsch-Gordon coefficients.