

Sine Gordon model Renormalization

$$L[\Psi] = \frac{1}{2V} (\partial_t \Psi)^2 - \frac{V}{2} (\partial_x \Psi)^2 + \widehat{g} \cos \beta \Psi$$

integrations by part

$$S[\Psi] = \iint dt dx \underbrace{\left[\frac{V}{2} \Psi \partial_x^2 \Psi - \frac{1}{2V} \Psi (\partial_t \Psi)^2 + \widehat{g} \cos \beta \Psi \right]}_{S_0[\Psi]} + S_I[\Psi]$$

split into fast and slow modes:

$$\Psi(x) = \Psi^S(x) + \delta(\Psi(x)) (\text{slow})$$

$$(\partial_t \Psi^S(x) + \delta \Psi(x)) (\partial_x^2 \Psi^S(x) + \partial_x^2 \delta \Psi(x))$$

$$\begin{aligned} & \int dx \delta \Psi \nabla_x^2 \Psi^S = \int dx \Psi^S \nabla_x^2 \delta \Psi \xrightarrow{\text{Fourier transform}} \\ &= \int dx \int_{\text{bulk}} \frac{d\vec{q}}{(2\pi)^2} \int_{\text{shell}} \frac{d\vec{q}'}{(2\pi)^2} \Psi(q) \left(-q^2 - \frac{\omega^2}{q^2} \right) \Psi(q') \\ & e^{i(\vec{q} + \vec{q}') \cdot \vec{x}} \\ &= \int dx \int_{\text{bulk}} \frac{d\vec{q}}{(2\pi)^2} \Psi(q) \left(-q^2 - \frac{\omega^2}{q^2} \right) \Theta(|q| - \frac{1}{s}) \end{aligned}$$

Definition:

$$\frac{1}{2} \nabla_{x'}^2 G_0(x, x') = \delta(x - x')$$

$$\begin{aligned} & \frac{1}{2} \nabla_x^2 \int \frac{d\mathbf{q}}{(2\pi)^2} G_0(\mathbf{q}) e^{i\mathbf{q}(x-x')} = \int \frac{d\mathbf{q}}{(2\pi)^2} e^{i\mathbf{q}(x-x')} \\ &= \int \frac{d\mathbf{q}}{(2\pi)^2} G_0(\mathbf{q}) \cdot \frac{1}{2} \left(-\mathbf{q}^2 - \frac{\omega^2}{\mathbf{q}^2} \right) e^{i\mathbf{q}(x-x')} \\ & G_0(\mathbf{q}) = -\frac{2\gamma^2}{\mathbf{q}^2 \gamma^2 + \omega^2} \end{aligned}$$

Dyson Eq:

$$G^{-1}(x, x') = G_0^{-1}(x, x') - \sum (x_i x'^i) \quad \downarrow \text{reverse}$$

$$\begin{aligned} & \frac{1}{2} \nabla_x^2 G(x - x') = - \int d\mathbf{q}' \sum (x_i x'^i) G(x', x') \\ &= \frac{1}{(2\pi)^2} \delta(x - x') \\ \Rightarrow & \int \frac{d\mathbf{q}^2}{(2\pi)^4} \frac{d\mathbf{q}'^2}{(2\pi)^2} G(\mathbf{q}, \mathbf{q}') \frac{1}{2} \left(-\mathbf{q}^2 - \frac{\omega^2}{\mathbf{q}^2} \right) e^{i\mathbf{q}x + i\mathbf{q}'x'} \quad \downarrow \end{aligned}$$

$$-\int d\vec{q}'' \int \frac{d\vec{q}}{(2\pi)^2} \frac{d\vec{k}}{(2\pi)^2} \frac{d\vec{k}'}{(2\pi)^2} \frac{d\vec{q}'}{(2\pi)^2} \sum(\vec{q}, \vec{k}) G(\vec{k}', \vec{q}')$$

$$e^{i\vec{q}\vec{x} + i(\vec{k} + \vec{k}')\vec{x}'' + i\vec{q}'\vec{x}'} = \int \frac{d\vec{q}^2}{(2\pi)^4} e^{i\vec{q}(\vec{x} - \vec{x}')} \\ \text{integrate out}$$

$$\int \frac{d\vec{q}}{(2\pi)^2} \frac{d\vec{q}'}{(2\pi)^2} \left[G(\vec{q}, \vec{q}') G_0(\vec{q}') - \int \frac{d\vec{q}''}{(2\pi)^2} \sum(\vec{q}, \vec{q}'') \right]$$

$$G(-\vec{q}'', \vec{q}') \cdot e^{i\vec{q}\vec{x} + i\vec{q}'\vec{x}'} = \int \frac{d\vec{q}^2}{(2\pi)^2} \frac{d\vec{q}''}{(2\pi)^2} \delta(\vec{q} + \vec{q}'')$$

$$e^{i\vec{q}\vec{x} + i\vec{q}'\vec{x}'}$$

$$\Rightarrow G(\vec{q}, \vec{q}') = G_0(\vec{q}) \delta(\vec{q} + \vec{q}') + G_0(\vec{q})$$

$$\int \frac{d\vec{q}''}{(2\pi)^2} \sum(\vec{q}, \vec{q}'') G(-\vec{q}'', \vec{q}')$$

• perturbative theory in powers of g

$$G^{(0)}(\vec{q}, \vec{q}') = G_0(\vec{q}) \delta(\vec{q} + \vec{q}')$$

up to first order (zero-th order)

$$G^{(1)}(\vec{q}, \vec{q}') = G_0(\vec{q}) \delta(\vec{q} + \vec{q}') - G_0(\vec{q}) \int \frac{d\vec{q}''}{(2\pi)^2} b^S(\vec{q}, \vec{q}'', \vec{q}')$$

$$G(-\vec{q}'', \vec{q}'')$$

$$G^{(1)}(q, q') = G_0(q) \delta(q + q') - \frac{1}{(2\pi)^2} G_0(q) b^S(q, q') G_0(q')$$

Second order

$$G^{(2)}(q, q') = G_0(q) \delta(q + q') - G_0(q) \int \frac{d^2 q''}{(2\pi)^2} b^S(q, q'')$$

$$\left[G_0(-q''), \delta(-q'' + q') - \frac{1}{(2\pi)^2} G_0(q'') b^S(-q'', q') G_0(-q') \right]$$

$$\Rightarrow G^{(2)}(q, q') = G_0(q) \delta(q + q') - \frac{1}{(2\pi)^2} G_0(q) b^S(q, q')$$

$$G_0(-q') + \frac{1}{(2\pi)^4} G_0(q) \int \frac{d^2 q''}{(2\pi)^2} b^S(q, q'') G_0(-q'')$$

$$b^S(-q'', q') G_0(-q'')$$

with

$$S[\Psi^S, \delta\Psi] = S_0[\delta\Psi] + \int d\mathbf{x} a^S(\mathbf{x}) \delta\Psi(\mathbf{x}) +$$

$$\int d\mathbf{x} d\mathbf{x}' \delta\Psi(\mathbf{x}) b^S(\mathbf{x}, \mathbf{x}') \delta\Psi(\mathbf{x}')$$

$$\rightarrow \int d\mathbf{x} d\mathbf{x}' \delta\Psi(\mathbf{x}) \left[\delta(\mathbf{x} - \mathbf{x}') \frac{1}{2} \nabla_x^2 + b^S(\mathbf{x}, \mathbf{x}') \right] \delta\Psi(\mathbf{x}')$$

$$+ \int d\mathbf{x} a^S(\mathbf{x}) \delta\Psi(\mathbf{x})$$

$$-\int d\vec{q}'' \int \frac{d\vec{q}}{(2\pi)^2} \frac{d\vec{k}}{(2\pi)^2} \frac{d\vec{k}'}{(2\pi)^2} \frac{d\vec{q}'}{(2\pi)^2} \sum(\vec{q}, \vec{k}) G(\vec{k}', \vec{q}')$$

$$e^{i\vec{q}\vec{x} + i(\vec{k} + \vec{k}')\vec{x}'' + i\vec{q}'\vec{x}'} = \int \frac{d\vec{q}}{(2\pi)^4} e^{i\vec{q}(\vec{x} - \vec{x}')}$$

integrate out

$$\int \frac{d\vec{q}}{(2\pi)^4} \frac{d\vec{q}'}{(2\pi)^4} \left[G(\vec{q}, \vec{q}') G_0(\vec{q}) - \int \frac{d\vec{q}''}{(2\pi)^2} \sum(\vec{q}, \vec{q}'') \right]$$

$$G(-\vec{q}'', \vec{q}') e^{i\vec{q}\vec{x} + i\vec{q}'\vec{x}'} = \int \frac{d\vec{q}''}{(2\pi)^2} \frac{d\vec{q}'}{(2\pi)^2} \delta(\vec{q} + \vec{q}')$$

$$e^{i\vec{q}\vec{x} + i\vec{q}'\vec{x}'}$$

$$\Rightarrow G(\vec{q}, \vec{q}') = G_0(\vec{q}) \delta(\vec{q} + \vec{q}') + G_0(\vec{q}')$$

$$\int \frac{d\vec{q}'}{(2\pi)^4} \sum(\vec{q}, \vec{q}'') G(-\vec{q}'', \vec{q}')$$

- Perturbative theory in powers of g

$$G^{(0)}(\vec{q}, \vec{q}') = G_0(\vec{q}) \delta(\vec{q} + \vec{q}')$$

up to first order

$$G^{(1)}(\vec{q}, \vec{q}') = G_0(\vec{q}) \delta(\vec{q} + \vec{q}') - G_0(\vec{q}) \int \frac{d\vec{q}''}{(2\pi)^2} b^S(\vec{q}, \vec{q}'') G(-\vec{q}'', \vec{q}')$$

$$G^{(1)}(q, q') = G_0(q) \delta(q + q') - \frac{1}{(2\pi)^2} G_0(q) b^S(q, q') G_0(-q')$$

Second order

$$\begin{aligned} G^{(2)}(q, q') &= G_0(q) \delta(q + q') - G_0(q) \int \frac{dq''}{(2\pi)^2} b^S(q, q'') \\ &\quad [G_0(-q'') \delta(-q'' + q') - \frac{1}{(2\pi)^2} G_0(q'') b^S(-q'', q') G_0(-q')] \\ \Rightarrow G^{(2)}(q, q') &= G_0(q) \delta(q + q') - \frac{1}{(2\pi)^2} G_0(q) b^S(q, q') \\ &\quad G_0(-q') + \frac{1}{(2\pi)^4} G_0(q) \int \frac{dq''}{(2\pi)^2} b^S(q, q'') G_0(-q') \\ &\quad b^S(-q'', q') G_0(-q') \end{aligned}$$

with

$$\begin{aligned} S[\Psi^S, \delta\Psi] &= S_0[\delta\Psi] + \int dx a^S(x) \delta\Psi(x) + \\ &\quad \int dx dx' \delta\Psi(x) b^S(x, x') \delta\Psi(x') \\ \rightarrow \int dx dx' \delta\Psi(x) &\left[\delta(x-x') \frac{1}{2} \nabla_x^2 + b^S(x, x') \right] \delta\Psi(x') \\ &+ \int dx a^S(x) \delta\Psi(x) \end{aligned}$$

$$\rightarrow \delta S[\Psi^S, \delta \Psi] = \int dx dx' \delta \Psi(x) G^{-1}(x, x') \delta \Psi(x')$$

$$+ \int dx \alpha^S(x) \Psi(x)$$

$$\delta \Psi(x) = \widetilde{\Psi}(x) + r(x)$$

$$\delta S[\Psi^S, \delta \Psi] = \int dx dx' \{ \widetilde{\Psi}(x) G^+(x, x') \widetilde{\Psi}(x')$$

$$+ \widetilde{\Psi}(x) [2G^+(x, x') r(x') + \alpha^S(x') \delta(x - x')]$$

$$+ r(x) [G^-(x, x') \Psi(x) + \alpha^S(x) \delta(x - x')]$$

Now if

$$\int dx' G(\bar{x}, x') \Psi(x') = -\frac{1}{2} \alpha^S(x)$$

$$\delta S[\Psi^S, \delta \Psi] = \int dx dx' [\widetilde{\Psi}(x) G^+(\bar{x}, x') \widetilde{\Psi}(x') - \frac{1}{4} \alpha^S(x) G^+(x, x') \alpha^S(x')]$$

$$\Rightarrow \int dx \underbrace{\int dx' G(x'', x) G(\bar{x}, x')}_{\delta(x'' - x)} \Psi(x') +$$

$$= r(x) = -\frac{1}{2} \int dx' G(x, x') \alpha^S(x')$$

$$\Psi(\alpha) = \Psi^s(\alpha) + r(\alpha) + \tilde{\Psi}(\alpha) = \tilde{\Psi}(\alpha) + \Psi^s(\alpha)$$

$$H_0[\Pi, \Psi] = \int d\alpha \frac{\partial}{2} [\Pi^2 + (\partial_\alpha \Psi)^2] \rightarrow H_0[\Pi^s, \Psi]$$

$$+ H_0[\tilde{\Pi}, \tilde{\Psi}]$$

unperturbed Green function

$$\begin{aligned} G_0(\alpha, \alpha') &= \langle 0 | \Psi(\alpha) \Psi(\alpha') | 0 \rangle^\Psi \\ &= \langle 0^\Psi | (\tilde{\Psi}(\alpha) + \Psi^s(\alpha)) (\tilde{\Psi}(\alpha') + \Psi^s(\alpha')) | 0 \rangle^\Psi \\ \Rightarrow G_0(\alpha, \alpha') &= 2 \langle 0 | \tilde{\Psi}(\alpha) \tilde{\Psi}(\alpha') | 0 \rangle^\Psi \\ &= 2 \langle 0 | \Psi(\alpha) \Psi(\alpha') | 0 \rangle^\tilde{\Psi} \end{aligned}$$

the effective slow mode residual contribution to
the action

$$S_{\text{eff}}[\Psi^s] = \langle 0 | S[S[\Psi^s, \Phi]] | 0 \rangle^\Psi$$

$$\begin{aligned} S_{\text{eff}}[\Psi^s] &= \int d\alpha d\alpha' \left[\frac{1}{2} G_0(\alpha, \alpha') G^{-1}(\alpha', \alpha') \right. \\ &\quad \left. - \frac{1}{4} \alpha^s(\alpha) G(\alpha, \alpha') \alpha^s(\alpha') \right] \end{aligned}$$

$$S_{\text{eff}}[\Psi^s] = \int d\alpha d\alpha' \left[\frac{1}{2} G_0(\alpha, \alpha') [G_0(\alpha, \alpha') + b^s(\alpha, \alpha')] \right]$$

$$- \frac{1}{4} a^s(\alpha) G(\alpha, \alpha') a^s(\alpha')$$

$$= \langle 0 | S_0[\Psi] | 0 \rangle^\Psi + \int d\alpha d\alpha' \left[\frac{1}{2} G_0(\alpha, \alpha') b^s(\alpha, \alpha') \right]$$

$$- \frac{1}{4} a^s(\alpha) G(\alpha, \alpha') a^s(\alpha') \leftarrow S_{\text{eff}}[\Psi]$$

Let's compute first term:

$$F_1[\Psi^s] = \frac{1}{2} \int d\alpha d\alpha' \int_{\text{Shell}} \frac{d\vec{q}}{(2\pi)^2} \frac{d\vec{q}'}{(2\pi)^2} G_0(\vec{q})$$

$$b^s(q', q'') e^{i\vec{q}(\alpha - \alpha') + i\vec{q}'\alpha'' + i\vec{q}''\alpha'}$$

\downarrow integrate out α, α'

$$= \frac{1}{2} \int_{\text{Shell}} \frac{d\vec{q}}{(2\pi)^2} G_0(\vec{q}) \Theta' b^s(-\vec{q}, \vec{q})$$

$$b^s(\alpha, \alpha') = -\frac{\beta^2}{2} \vec{q} \cdot \cos \beta \Psi^s \delta(\alpha - \alpha')$$

$$= -\frac{\beta^2}{2} d_I[\Psi^s] \Psi(\alpha - \alpha')$$

$$b^s(p, q') = \int dx dx' b^s_{\cancel{q} \cancel{x}}, e^{iqx + iq'x'} g(x - x')$$

$$-\frac{p^2}{2} g \cos p \Phi$$

$$= b_s(q + q')$$

- $F_I[\Psi^s] = \frac{q}{2} b^s(q=0) \int_{\text{shape}} \frac{dq}{(2\pi)} G_0(q)$
- $= -b^s(q=0) \int_{\text{shape}} \frac{dq}{(2\pi)^2} \frac{1}{q^2 + \omega^2/v^2}$
- $= -\frac{1}{2\pi} b^s(q=0) \int_{NS} \lambda \frac{|q|}{q^2} dq$
- $= -\frac{1}{2\pi} b^s(q=0) \log s$
- $\rightarrow \frac{p^2}{4\pi} \log s \int dx f_I[\Psi^s] = -\frac{g dl}{4\pi} p^2 \int dx \cos p \Phi$

$$F_2[\Psi^s] = - \int dx dx' \frac{1}{4} G_0(x, x') a^s(x) a^s(x')$$

=

$$S_R[\psi] = \int d\lambda \left[\frac{1}{2\beta^2} \left(1 + \frac{3\beta^4 g^2 d\lambda}{4\pi \lambda^3} \right) \psi \partial_\lambda \psi + g s^{-2} \left(1 - \frac{\beta^2 d\lambda}{4\pi} \right) \cos \psi \right]$$

$$S[\psi] = \int d\lambda \left[\frac{1}{2\beta^2} \psi \nabla_\lambda^2 \psi + g \cos \psi \right]$$

$$\Rightarrow \begin{cases} \beta_R^{-2} = \beta^{-2} \left(1 + \frac{3\beta^4 g^2 d\lambda}{4\pi \lambda^3} \right) \\ g_R = g s^{-2} \left(1 - \frac{\beta^2 d\lambda}{4\pi} \right) \end{cases}$$

$$\Rightarrow \frac{d\beta^{-2}}{d\lambda} = \frac{3\beta^2 g^2}{4\pi \lambda^3} = \frac{-2}{\beta^{-3}} \frac{d\beta}{d\lambda}$$

$$\boxed{\frac{d\beta^2}{d\lambda} = -\frac{3\beta^6 g^2}{4\pi \lambda^3}}$$

$$g_R = g \left(1 + 2dg \right) \left(1 - \frac{\beta^2 d\lambda}{4\pi} \right) = g + g \left(2 - \frac{\beta^2}{4\pi} \right) d\lambda$$

$$\boxed{\frac{dg}{d\lambda} = 2g \left(1 - \frac{\beta^2}{8\pi} \right)}$$

$$\text{Let: } K = \frac{\beta^2}{8\pi}, \quad \pi u = 4\sqrt{\frac{3\pi}{\lambda^3}} g$$

$$\frac{du}{dl} = 2u(1-k^2)$$

$$\frac{dk}{dl} = -\pi u^3 k^3$$

$$K = 1 + \gamma \quad \left\{ \begin{array}{l} \frac{du}{dl} = -2\pi u \gamma \\ \frac{dv}{dl} = -\pi u^2 (1+2v) \end{array} \right.$$

$$\Rightarrow 2u \frac{du}{dl} = \frac{du^2}{dl} = -4u^2 v$$

$$2v \frac{dv}{dl} = -2\pi^2 v - \frac{4u^2 v^2}{\downarrow} = \frac{dv^2}{dl}$$

second order

$$\Rightarrow \frac{d}{dl}(u^2 - 2v^2) = \frac{d}{dl}(u^2 - 2(K-1)^2) = 0$$

