

# Coupled Problems in Nonlinear Solid Mechanics: Non-Fickian Diffusion

## Deal.II Workshop

August 23 - 25 2010  
Heidelberg

Andrew McBride and Paul Steinmann

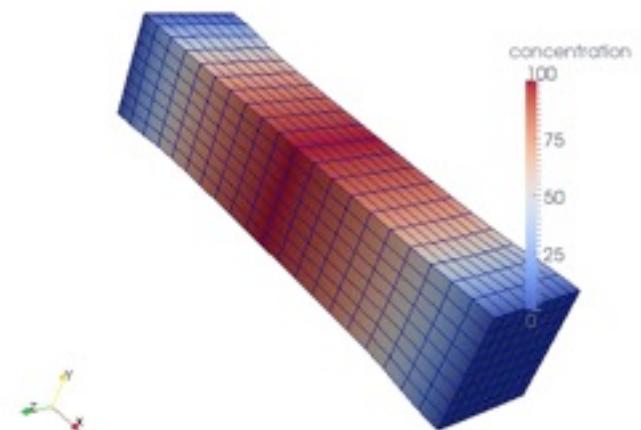
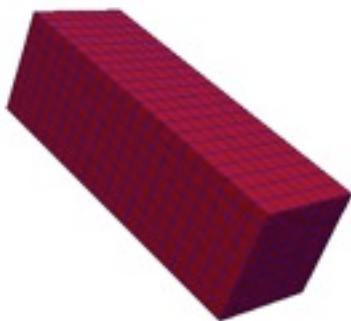
Chair of Applied Mechanics

Friedrich-Alexander-University Erlangen-Nuremberg

Swantje Bargmann

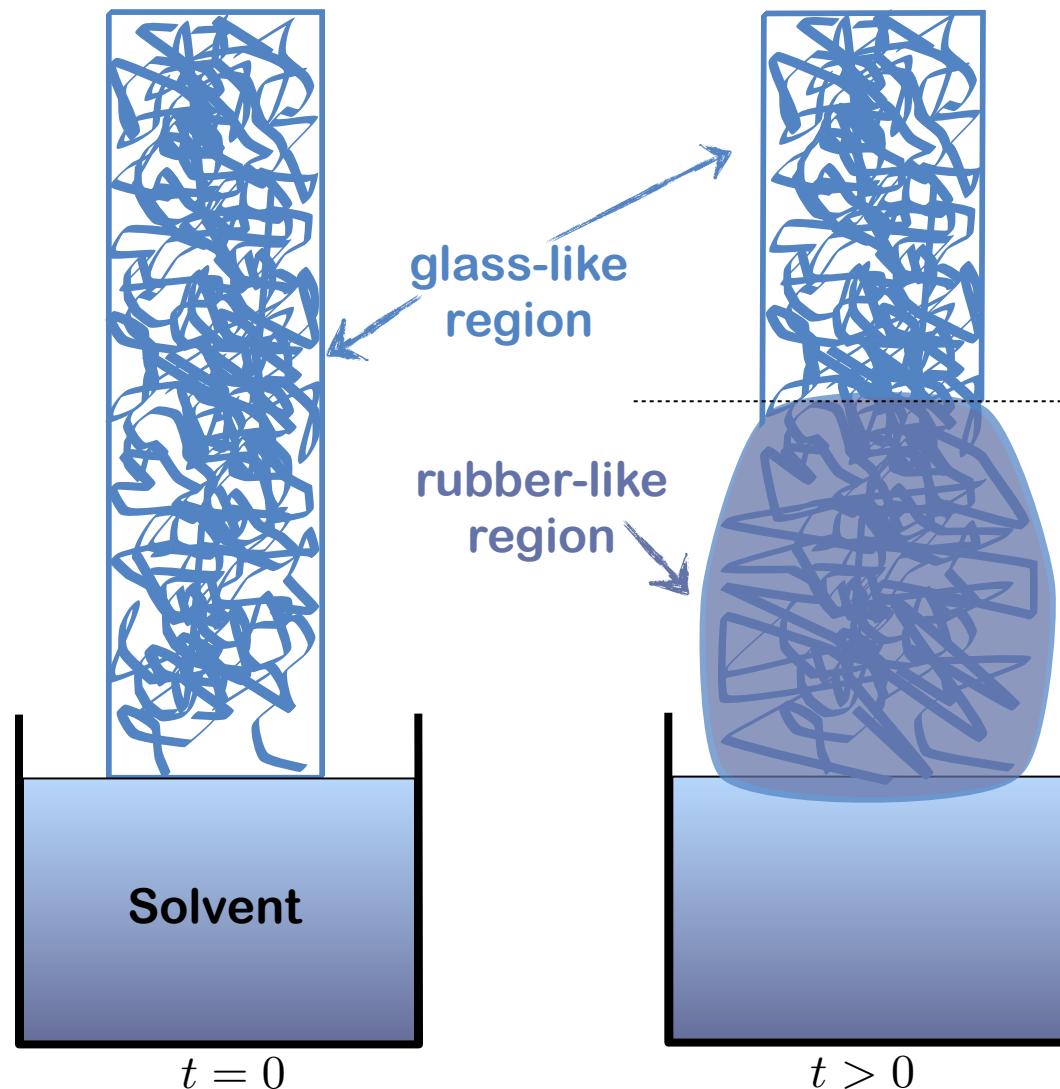
Institute of Mechanics

Technische Universität Dortmund



# Motivation: Case II diffusion

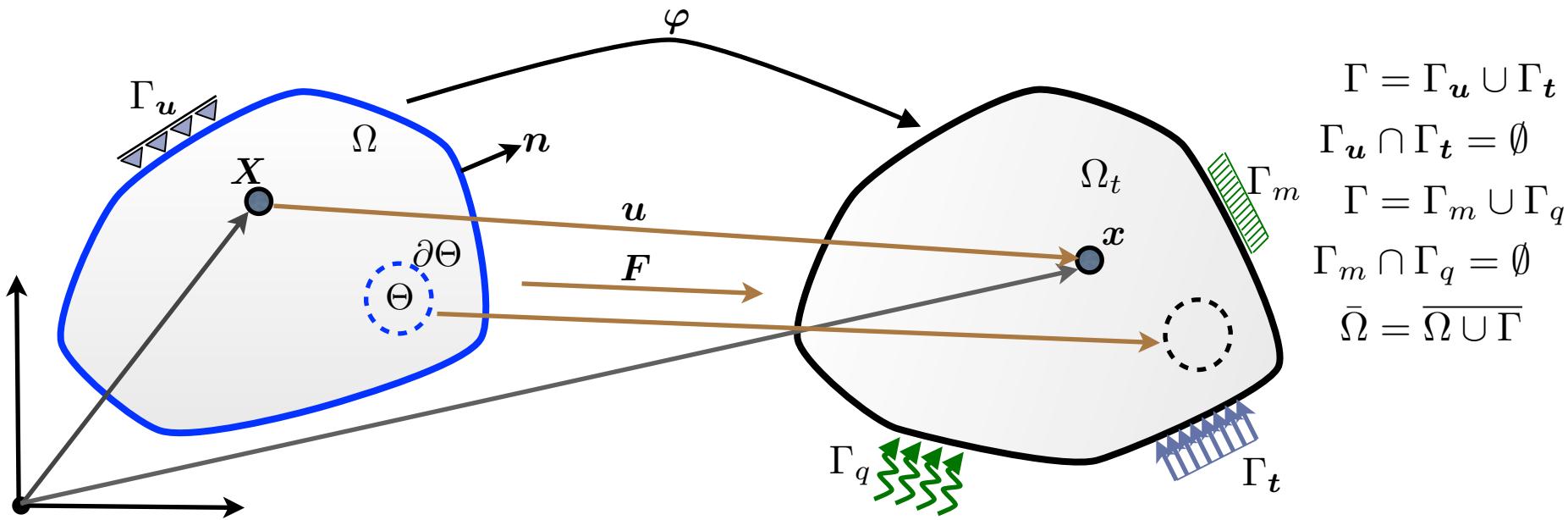
- Case II diffusion (CII) occurs during diffusion of low molecular weight **solvent** in **polymeric solid**
- Solid is originally in a **glass-like** state
- Solvent **wave** progresses through polymer at fixed rate (**constant velocity**)
- Solvent causes polymer change to a **rubber-like** material
- **Highly-coupled non-linear** phenomenon
  - Relaxation time of polymer depends on concentration and **swelling**
  - Solvent exerts a **pressure** on polymer and visa-versa (swelling and concentration dependent)
  - Diffusivity depends on **swelling** and concentration
  - Polymer needs finite amount time to rearrange to accommodate solvent limiting diffusion rate
  - Large swelling in rubber-like region
  - Polymer near-incompressible
- See **DE KEE ET AL (2005)** and **VESELY (2008)** for reviews



- Applications based presentation
  - Case II a prototype for highly non-linear coupled problems
- Describe the governing equations
  - Focus on **internal variable** formulation of viscoelasticity
  - Integration algorithms for internal inelastic variables
- Review models for case II diffusion
  - Strongly coupled diffusion-deformation
- Solution strategies using finite elements
- Consider a reduced model
  - Focus on **spatial adaptivity** with **internal variables** in **deal.II**
- Example problem

Work in progress (does not actually all work yet): comments and suggestions greatly appreciated. Actually, a lot of the questions I ask in the talk have been answered and better strategies proposed. Thanks

# Governing equations



## Kinematics

$$\varepsilon(\mathbf{u}) = \frac{1}{2} \left[ \nabla \mathbf{u} + [\nabla \mathbf{u}]^T \right]$$

$$\mathbf{F} = \frac{\partial \mathbf{x}}{\partial \mathbf{X}}$$

$$j = \det(\mathbf{F})$$

## Equilibrium

$$\operatorname{div} \boldsymbol{\sigma} = \mathbf{0} \quad \text{in } \Omega$$

$$\mathbf{u} = \bar{\mathbf{u}} \quad \text{on } \Gamma_u$$

$$\mathbf{t} := \boldsymbol{\sigma} \cdot \mathbf{n} = \bar{\mathbf{t}} \quad \text{on } \Gamma_t$$

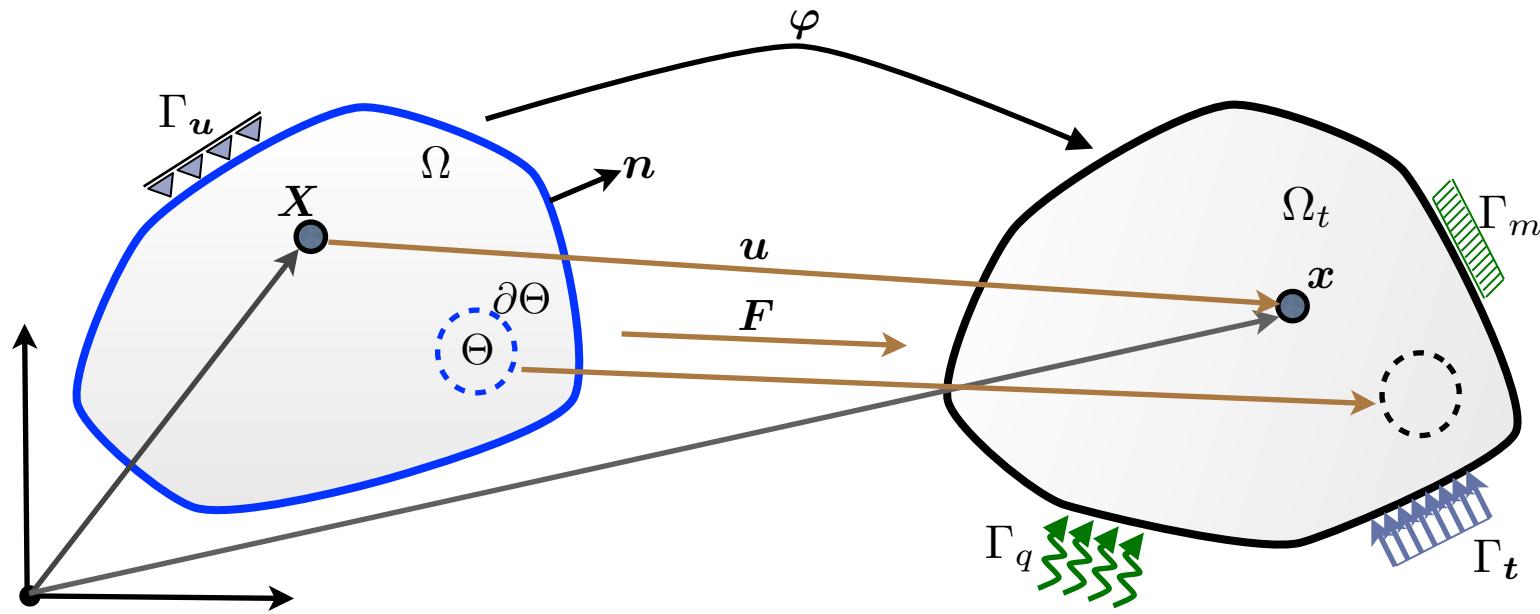
## Conservation of solvent mass

$$\dot{m} = -\operatorname{div} \mathbf{h} \quad \text{in } \Omega \times [0, T]$$

$$m = \bar{m} \quad \text{on } \Gamma_m$$

$$q := -\mathbf{h} \cdot \mathbf{n} = \bar{q} \quad \text{on } \Gamma_q$$

# Constitutive relations



## Free energy

$$\psi = \psi(\varepsilon, \alpha, m)$$

↑  
inelastic (viscous) strain

## Stress

$$\sigma = \frac{\partial \psi}{\partial \varepsilon}$$

## Solvent flux

ideal mixtures:

$$\mathbf{h} = -\mathbf{D}(m) \cdot \nabla m$$

non-ideal mixtures:

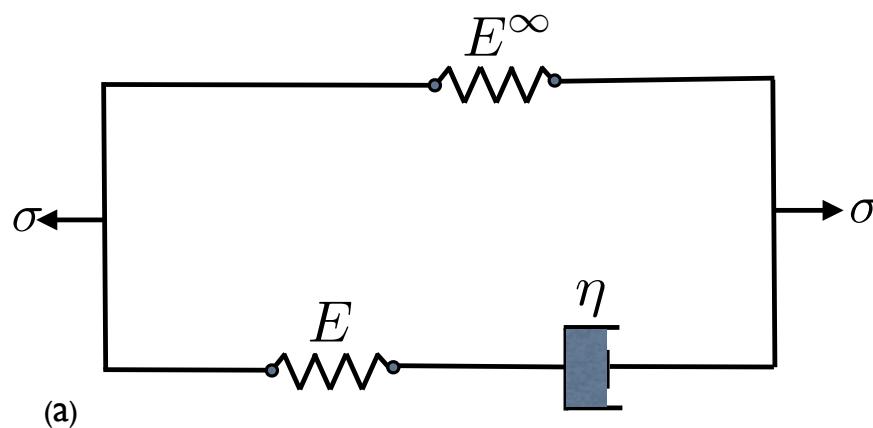
$$\mathbf{h} = -\mathbf{D}(m)m \cdot \nabla \mu(m, j)$$

chemical potential →

# Constitutive relations: viscoelasticity

LM

- Polymeric solid response is well described as a **viscoelastic** solid
- Adopt a model due to SIMO & HUGHES (1998)
  - Inelastic strain treated as an **internal variable**
- Consider the one-dimensional standard linear solid:



inelastic (viscous) strain

$$\sigma = E_\infty \varepsilon + \sigma^v$$
$$\sigma^v = \eta \dot{\alpha} = E [\varepsilon - \alpha]$$
$$\sigma = \underbrace{E_0}_{E_\infty + E} \varepsilon - E\alpha$$

viscosity

$$\tau = \frac{\eta}{E}$$

relaxation time

Evolution of  $\alpha$

$$\left\{ \begin{array}{l} \dot{\alpha} + \frac{\alpha}{\tau} = \frac{\varepsilon}{\tau} \\ \lim_{t \rightarrow -\infty} \alpha(t) = 0 \end{array} \right.$$

# Viscoelasticity

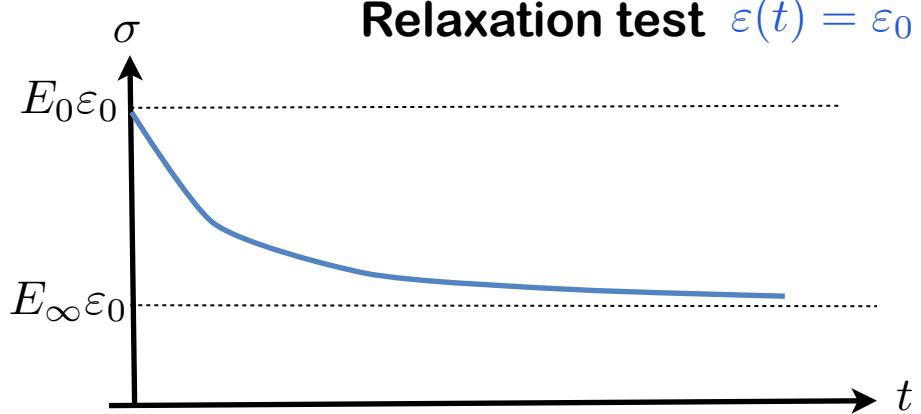
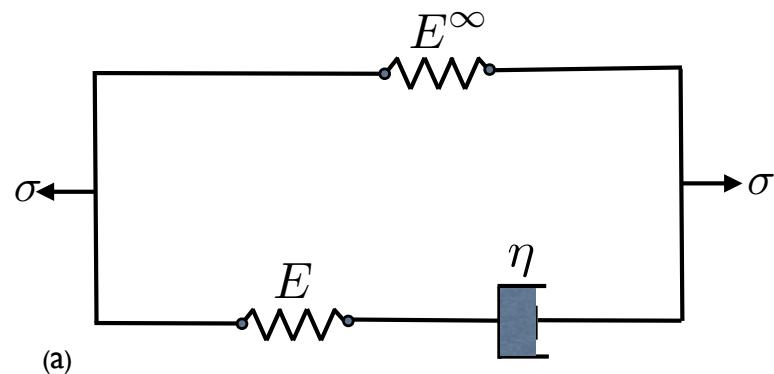
LM

## Convolution representation

$$\sigma(t) = \int_{-\infty}^t G(t-s)\dot{\varepsilon}(s) \, ds$$

$$G(t) = E_\infty + E \exp\left(\frac{-t}{\tau}\right)$$

relaxation function



## Thermodynamics

$$\psi(\varepsilon, \alpha) := \frac{1}{2}E_\infty \varepsilon^2 + \frac{1}{2}E [\varepsilon - \alpha]^2$$

$$\mathcal{D} = \sigma^v \dot{\alpha} = \eta \dot{\alpha}^2 \geq 0 \quad \text{dissipation}$$

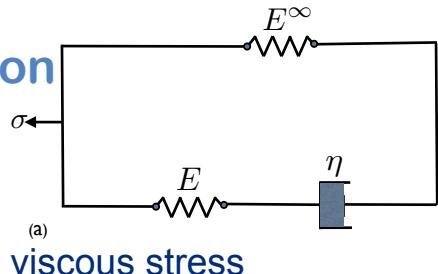
$$\sigma^v = -\frac{\partial \psi(\varepsilon, \alpha)}{\partial \alpha} \quad \text{viscous stress}$$

$$\sigma = \frac{\partial \psi(\varepsilon, \alpha)}{\partial \varepsilon} \quad \text{stress}$$

# Viscoelasticity

LM

## Alternative representation



## Extension to $\mathbb{R}^3$

$$q := E\alpha$$

$$\sigma = E_0\varepsilon - q$$

$$\begin{cases} \gamma := \frac{E}{E_0} \\ \gamma_\infty := \frac{E_\infty}{E_0} \end{cases}$$

$$W^0(\varepsilon) := \frac{1}{2}\varepsilon E_0\varepsilon$$

$$\sigma = \frac{\partial W^0(\varepsilon)}{\partial \varepsilon} - q$$

$$\dot{q} + \frac{q}{\tau} = \frac{\gamma}{\tau} \frac{\partial W^0(\varepsilon)}{\partial \varepsilon}$$

$$\lim_{t \rightarrow -\infty} q(t) = 0$$

$$\varepsilon = e + \frac{1}{3}\Theta\mathbf{1}$$

$$e := \text{dev } \varepsilon \quad \text{and} \quad \Theta := \text{tr } \varepsilon$$

$$W^0(\varepsilon) = \bar{W}^0(e) + U^0(\Theta)$$

$$\sigma^0 := \frac{\partial W^0(\varepsilon)}{\partial \varepsilon} = \text{dev} \left( \frac{\partial \bar{W}^0}{\partial e} \right) + U^{0'}\mathbf{1}$$

$$\sigma(t) = \sigma^0(t) - q$$

$$\dot{q} + \frac{q}{\tau} = \frac{\gamma}{\tau} \text{dev} \left( \frac{\partial W^0(e)}{\partial e} \right)$$

$$\lim_{t \rightarrow -\infty} q(t) = 0$$

$$\sigma(t) = U^{0'}\mathbf{1} + \int_{-\infty}^t g(t-s) \frac{d}{ds} \left[ \text{dev} \left( \frac{\partial \bar{W}^0(e(s))}{\partial e} \right) \right] ds$$

$$g(t) := \gamma_\infty + \gamma_i \exp \left( \frac{-t}{\tau} \right)$$

deviatoric /  
volumetric split

initial stored-  
energy  
function

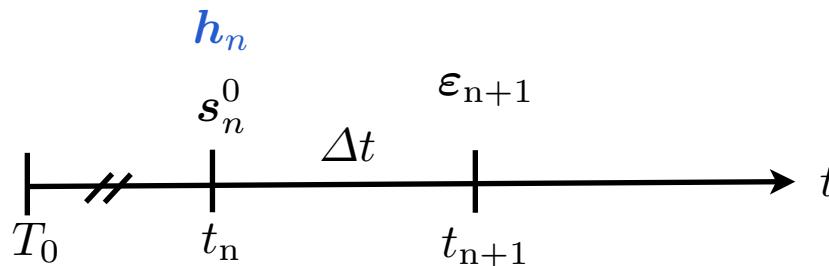
elastic stress

viscous stress  
evolution

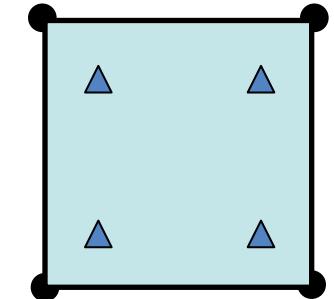
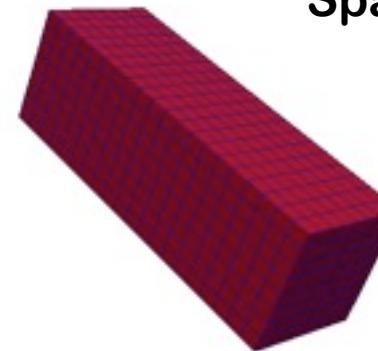
# Integration algorithm for viscoelasticity

LM

## Temporal discretisation



## Spatial discretisation



- Strain driven formulation
- Transform convolution representation for internal variables via two-step recurrence relationship
  - (approach restricted to relaxation functions consisting of linear combinations of functions in time that possess semi-group property)

$$e_{n+1} = \text{dev}(\varepsilon_{n+1})$$

$$s_{n+1}^0 = \text{dev}\left(\frac{\partial \bar{W}^0(e_{n+1})}{\partial e}\right)$$

$$h_{n+1} = \exp\left(\frac{-\Delta t_n}{\tau}\right) h_n + \exp\left(\frac{-\Delta t_n}{2\tau}\right) [s_{n+1}^0 - s_n^0]$$

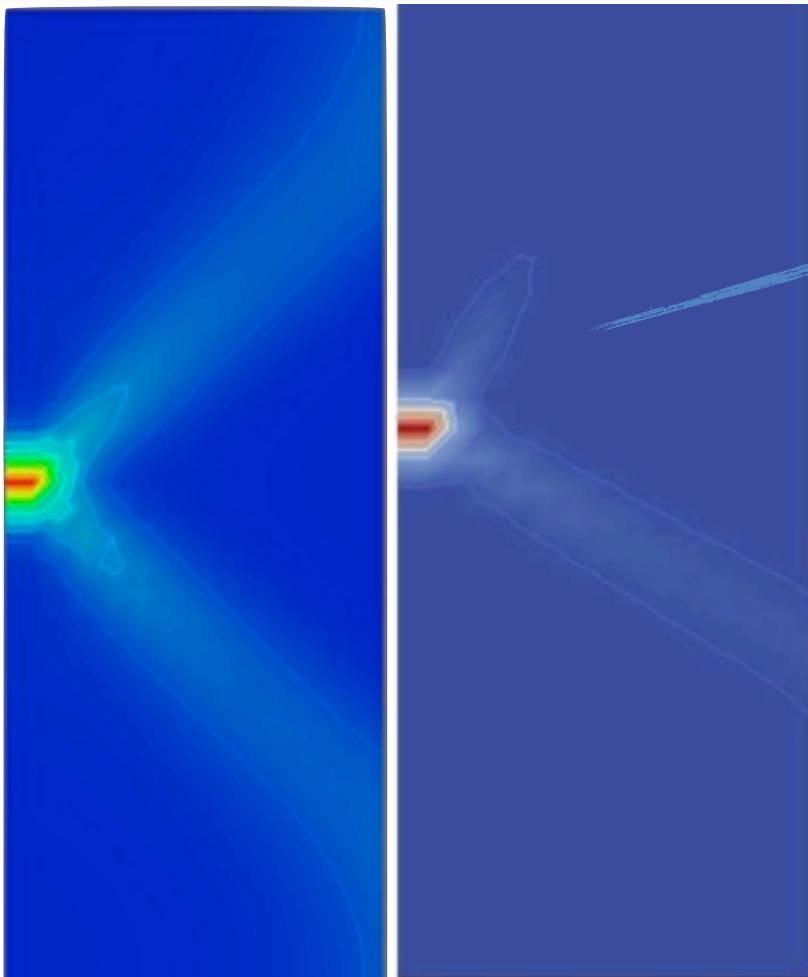
$$\sigma_{n+1} = U^{0'}(\Theta_{n+1})\mathbf{1} + \gamma_\infty s_{n+1}^0 + \gamma h_{n+1}$$



- data at level quadrature point
- no continuity relations
- State not determined from the nodal variables alone

# Internal variable formulations

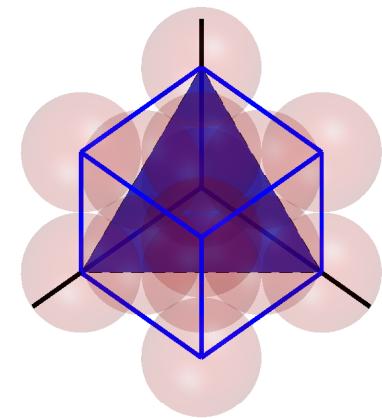
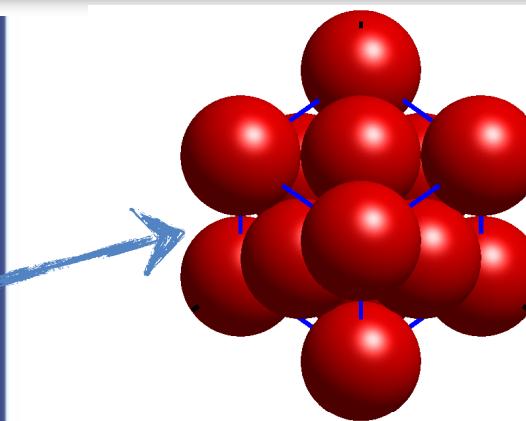
LM



$$v_1 = 0$$

$$v_1 = 15$$

**Algorithms for crystal plasticity**  
McB, Reddy, Richardson



## Internal variables:

- Additional state variables
- Evolve subject to an evolution equation that is possibly dependent on the primary nodal variables
- Possess no continuity requirements
- Generally treated directly in strong form at the level of the quadrature point

## Another example: Plasticity

Rate independent continuum or crystal plasticity

- Plastic strain (multiplier) generally treated as an internal variable
- More complex as the evolution of the plastic strain is subject to KTT constraints
- Damage, etc.

## Hyperbolic diffusion

- Classical diffusion relations are parabolic
  - CII concentration propagates as a wave at a fixed velocity: hyperbolic behaviour
- Following CATTANEO (1948) and VERNOTTE (1958) for Fourier's law of heat conduction:

$$\dot{m} = - \operatorname{div} \mathbf{h}$$

$$\dot{\mathbf{h}} + \beta \mathbf{h} = -\beta D \nabla_x m$$

assuming constant diffusivity

$$\underbrace{\operatorname{div} \dot{\mathbf{h}}}_{\ddot{m}} + \beta \underbrace{\operatorname{div} \mathbf{h}}_{\dot{m}} = \beta D \nabla^2 m$$

$$\dot{m} + \frac{1}{\beta} \ddot{m} = D \nabla^2 m \quad \leftarrow \text{hyperbolic diffusion}$$

- Propagation of heat as thermal wave (second sound) in fluids:

PESHKOV (1944, 1946) and PELLAM (1948) and gases:

ACKERMANN & OVERTON (1966), JACKSON ET AL (1970), NARAYANAMURTI & DYNES (1972)

- Also see framework of AIFANTIS (1980)
- Case II diffusion see KALOSPIROS ET AL (1991) and the GENERIC formulation of EL AFIF & GRMELA (2002)

# Models for Case II diffusion



## Strongly coupled non-Fickian models

- Non-Fickian diffusion coupled to a viscoelastic solid
  - Wu & PEPPAS (1993), GOVINDJEE & SIMO (1993), VIJALAPURA & GOVINDJEE (2003, 2005)
- Models of GOVINDJEE ET AL are far more well developed and advanced

Equilibrium and Solvent mass balance

$$\operatorname{Div}(\mathbf{F} \cdot \mathbf{S}) = \mathbf{0}$$

$$\dot{M} = \operatorname{Div}(\mathbf{D}(M, J, R)) M \mathbf{C}^{-1} \cdot \nabla_{\mathbf{X}} \mu$$

↑  
concentration reacted sites

$$(\mathbf{c} := \mathbf{F}^T \mathbf{F})$$

- Finite deformation setting
- Solid modelled Neo-Hookean viscoelastic (incompressible)
- Highly coupled and non-linear
- Perfect mixing not assumed: extension of Flory-Huggins model to transient regime

PKII stress

$$\mathbf{S} = J p \mathbf{C}^{-1} + \mathbf{S}_{\text{dev}}$$

$$\mathbf{S}_{\text{dev}} = \mathbf{S}_{\text{dev}}^\infty + 2\rho_0 j^{\frac{2}{3}} \operatorname{dev} \mathbf{Q}$$

elastic stress

$$\mathbf{S}_{\text{dev}}^\infty = 2\rho_0 \frac{\partial \psi^E}{\partial \mathbf{C}}$$

pressure  $p = p^\infty(j) + p_s(j, M) + q(j, M)$

$$\dot{q} + \frac{q}{\tau(j, M)} = \gamma_1 \dot{p}^\infty$$

$$\dot{\mathbf{Q}} + \frac{\mathbf{Q}}{\tau(j, M)} = \gamma_2 \left[ \frac{\partial \psi^E}{\partial \mathbf{C}} \right]$$

# Proposed models



Introduce the following features into two reduced models for case II diffusion:

## Non-linear coupled diffusion and deformation

- Key features of the GOVINDJEE ET AL model but restricted to small strains
- Prototype for coupled mixed formulations
- Work in progress

1

## Non-linear coupled hyperbolic diffusion deformation

- As 1 but adopt a hyperbolic diffusion model
- Treat reflection of concentration waves
- Future work

2

### Spatial adaptivity

- Solvent propagates through medium as a wave
- Exploit existing methodologies for hyperbolic problems
- Need to consider the projection of inelastic internal variables

### Space-time finite elements

- Extend spatial adaptivity features of deal.II to temporal adaptivity use finite element in space and time

### Optimal solvers and automatic differentiation routines (Sacado in Trilinos)

### Parallel implementation

### Finite strain formulation

Good problem for MeshWorker and NoX in Trilinos

# Reduced model 1

## Governing DAEs

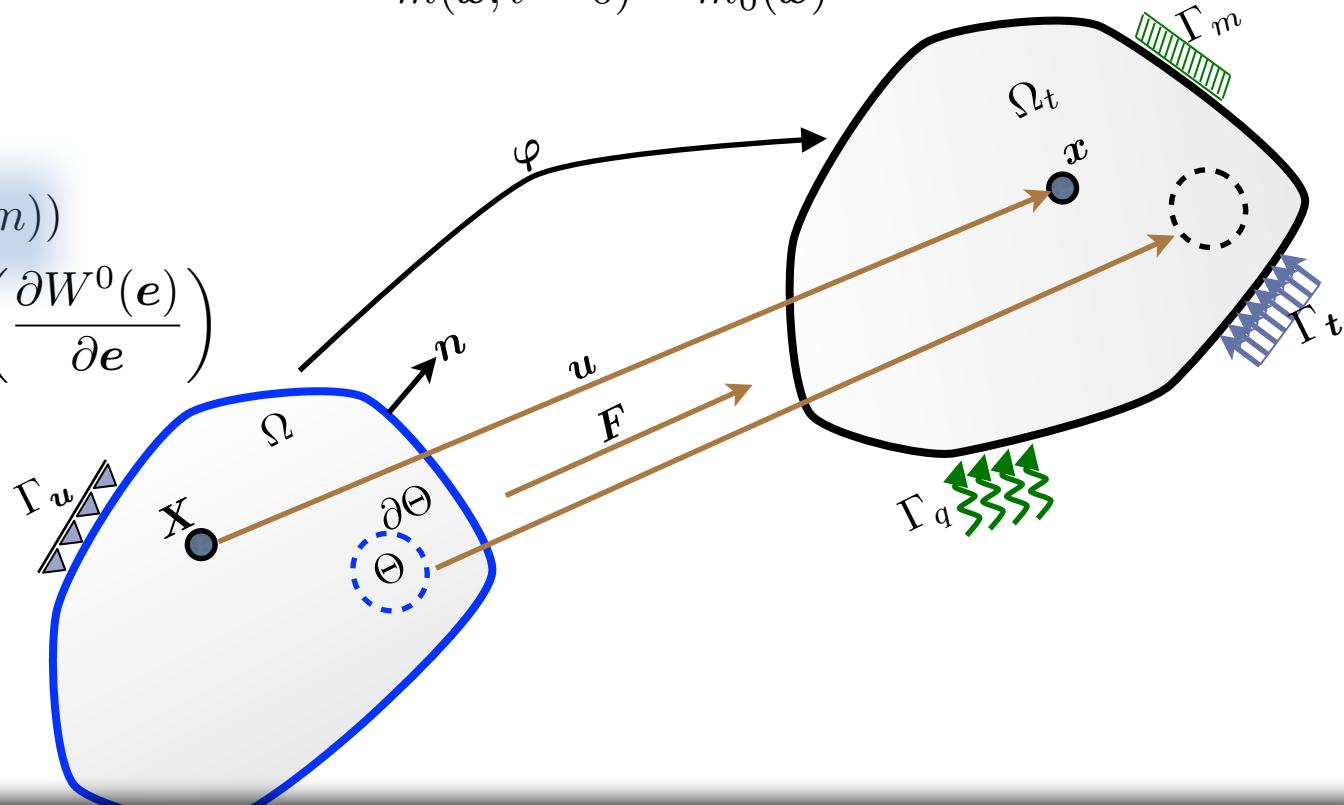
$$\begin{aligned}\operatorname{div} \boldsymbol{\sigma}(\boldsymbol{u}, \alpha, m) &= \mathbf{0} & \text{in } \Omega \times [0, T] \\ \dot{m} &= -\operatorname{div} \boldsymbol{h}(\boldsymbol{u}, m) & \text{in } \Omega \times [0, T]\end{aligned}$$

## Constitutive relations

$$\begin{aligned}\boldsymbol{\sigma} &= \boldsymbol{\sigma}(\boldsymbol{u}, q, m) \\ \boldsymbol{h} &= -D(m)\nabla m \\ (\boldsymbol{h} &= -D(m)m\nabla\mu(J, m)) \\ \dot{\boldsymbol{q}} + \frac{\boldsymbol{q}}{\tau(m)} &= \frac{\gamma}{\tau(m)} \operatorname{dev} \left( \frac{\partial W^0(\boldsymbol{e})}{\partial \boldsymbol{e}} \right)\end{aligned}$$

## Boundary and initial conditions

$$\begin{aligned}\boldsymbol{u} &= \bar{\boldsymbol{u}} && \text{on } \Gamma_u \\ t := \boldsymbol{\sigma} \cdot \boldsymbol{n} &= \bar{t} && \text{on } \Gamma_t \\ m &= \bar{m} && \text{on } \Gamma_m \\ q = -\boldsymbol{h} \cdot \boldsymbol{n} &= \bar{h} && \text{on } \Gamma_q \\ m(\boldsymbol{x}, t = 0) &= m_0(\boldsymbol{x})\end{aligned}$$



## Weak form equilibrium equation

$$\begin{aligned} (\operatorname{div} \boldsymbol{\sigma}, \mathbf{v})_{\Omega} &= 0 \quad \forall \mathbf{v} \in H_0^1(\Omega)^{n_{\text{dim}}} \\ (\boldsymbol{\sigma}, \boldsymbol{\varepsilon}(\mathbf{v}))_{\Omega} &= (\bar{\mathbf{t}}, \mathbf{v})_{\Gamma_t} \end{aligned}$$

## Temporal discretisation

$$\begin{aligned} [0, T] &\approx [0, t^1, \dots, t^n, t^{n+1}, \dots, T] \\ \Delta t &= t^{n+1} - t^n \end{aligned}$$

## Conservation of solvent mass

$$\begin{aligned} \frac{m^{n+1} - m^n}{\Delta t} &= - [\theta \operatorname{div} \mathbf{h}^{n+1} + (1 - \theta) \operatorname{div} \mathbf{h}^n] \quad \theta \in [0, 1] \\ m^{n+1} + \theta \Delta t \operatorname{div} \mathbf{h}^{n+1} &= m^n - (1 - \theta) \Delta t \operatorname{div} \mathbf{h}^n \end{aligned}$$

## Weak form

$$\begin{aligned} (m^{n+1}, v^{n+1})_{\Omega} + \theta \Delta t (\operatorname{div} \mathbf{h}^{n+1}, v^{n+1})_{\Omega} &= (m^n, v^{n+1})_{\Omega} \\ &\quad - (1 - \theta) \Delta t (\operatorname{div} \mathbf{h}^n, v^{n+1})_{\Omega} \quad \forall v \in H_0^1(\Omega) \end{aligned}$$

$$\begin{aligned} (m^{n+1}, v^{n+1})_{\Omega} - \theta \Delta t (\mathbf{h}^{n+1}, \nabla v^{n+1})_{\Omega} &= (m^n, v^{n+1})_{\Omega} + (1 - \theta) \Delta t (\mathbf{h}^n, \nabla v^{n+1})_{\Omega} \\ &\quad + \theta \Delta t (q^{n+1}, v^{n+1})_{\Gamma_q} + (1 - \theta) \Delta t (q^n, v^{n+1})_{\Gamma_q} \end{aligned}$$

- Following Step-23 we discretise in time first and then space: Rothe's method
- Allow for spatial adaptivity

# Reduced model 1

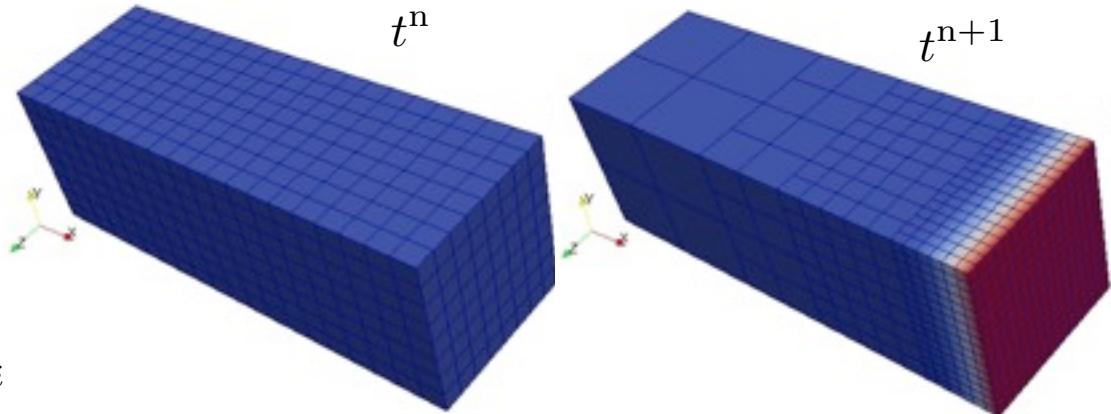
LM

## Displacement

$$\mathbf{u}(\mathbf{x}) \approx \mathbf{u}_h(\mathbf{x}) = \sum_{i=1}^{n_{\text{dof}}} \Phi_i^{\mathbf{u}}(\mathbf{x}) U_i$$

$$\mathbf{v}(\mathbf{x}) \approx \mathbf{v}_h(\mathbf{x}) = \sum_{i=1}^{n_{\text{dof}}} \Phi_i^{\mathbf{u}}(\mathbf{x}) \bar{U}_i$$

$$\boldsymbol{\varepsilon}(\mathbf{v}) \approx \boldsymbol{\varepsilon}(\mathbf{v}_h) = \sum_{i=1}^{n_{\text{dof}}} \nabla^{\text{sym}} (\Phi_i^{\mathbf{u}}(\mathbf{x})) \bar{U}_i$$



## Concentration

$$m^{n+1}(\mathbf{x}) \approx m_h^{n+1}(\mathbf{x}) = \sum_{i=1}^{n_{\text{dof}}} \Phi_i^{m,(n+1)}(\mathbf{x}) M_i^{n+1}$$

$$m^n(\mathbf{x}) \approx m_h^n(\mathbf{x}) = \sum_{i=1}^{n_{\text{dof}}} \Phi_i^{m,(n)}(\mathbf{x}) M_i^n$$

$$\nabla v^{n+1}(\mathbf{x}) \approx \nabla m_h^{n+1}(\mathbf{x}) = \sum_{i=1}^{n_{\text{dof}}} \nabla (\Phi_i^{m,(n+1)}(\mathbf{x})) \bar{M}_i^{n+1}$$

$$(m^n, v^{n+1})_{\Omega} \approx \sum_{i=1}^{n_{\text{dof}}} M_i^n \sum_{j=1}^{n_{\text{dof}}} \left( \Phi_i^{m,(n)}, \Phi_j^{m,(n+1)} \right)_{\Omega} \bar{M}_j^{n+1}$$

- Shape functions are defined on different meshes.
  1. Perform the integration on the finest common mesh (Step 28)
  2. Project solution at  $n$  to  $n+1$  (Steps 31-33)

# Reduced model 1



## Fully-discrete residual equations

$$\mathbf{R}(\mathbf{U}, \mathbf{M}) := [\mathbf{R}^u \ \mathbf{R}^m]^T$$

$$\mathbf{R}^u := (\boldsymbol{\sigma}_h, \boldsymbol{\varepsilon}(\mathbf{v}_h))_\Omega - (\bar{\mathbf{t}}_h, \mathbf{v}_h)_{\Gamma_t}$$

$$\begin{aligned}\mathbf{R}^m := & (m_h^{n+1}, v_h^{n+1})_\Omega - (m_h^n, v_h^{n+1})_\Omega \\ & - \theta \Delta t (\mathbf{h}_h^{n+1}, \nabla v_h^{n+1})_\Omega - [1 - \theta] \Delta t (\mathbf{h}_h^n, \nabla v_h^{n+1})_\Omega \\ & - \theta \Delta t (q_h^{n+1}, v_h^{n+1})_{\Gamma_q} - [1 - \theta] \Delta t (q_h^n, v_h^{n+1})_{\Gamma_q}\end{aligned}$$

$$\mathbf{R}_{i+1} = \mathbf{0}$$

$$\mathbf{R}_i + \frac{\partial \mathbf{R}}{\partial \mathbf{U}} \delta \mathbf{U} + \frac{\partial \mathbf{R}}{\partial \mathbf{M}} \delta \mathbf{M} = \mathbf{0}$$

$$\begin{bmatrix} \frac{\partial \mathbf{R}^u}{\partial \mathbf{U}} & \frac{\partial \mathbf{R}^u}{\partial \mathbf{M}} \\ \frac{\partial \mathbf{R}^m}{\partial \mathbf{U}} & \frac{\partial \mathbf{R}^m}{\partial \mathbf{M}} \end{bmatrix} \begin{bmatrix} \delta \mathbf{U} \\ \delta \mathbf{M} \end{bmatrix} = \begin{bmatrix} -\mathbf{R}_i^u \\ -\mathbf{R}_i^m \end{bmatrix}$$

• Non-symmetric, potentially highly non-linear even for the reduced problem. Full finite deformation problem is horrendously non-linear! (see Govindjee et al.)

- Currently solve using a monolithic Newton scheme to solve
  - ToDo: Investigate the use of split schemes
- Currently using an approximation to tangent
  - ToDo: Investigate the use of automatic differentiation tools in Sacado (Trilinos)

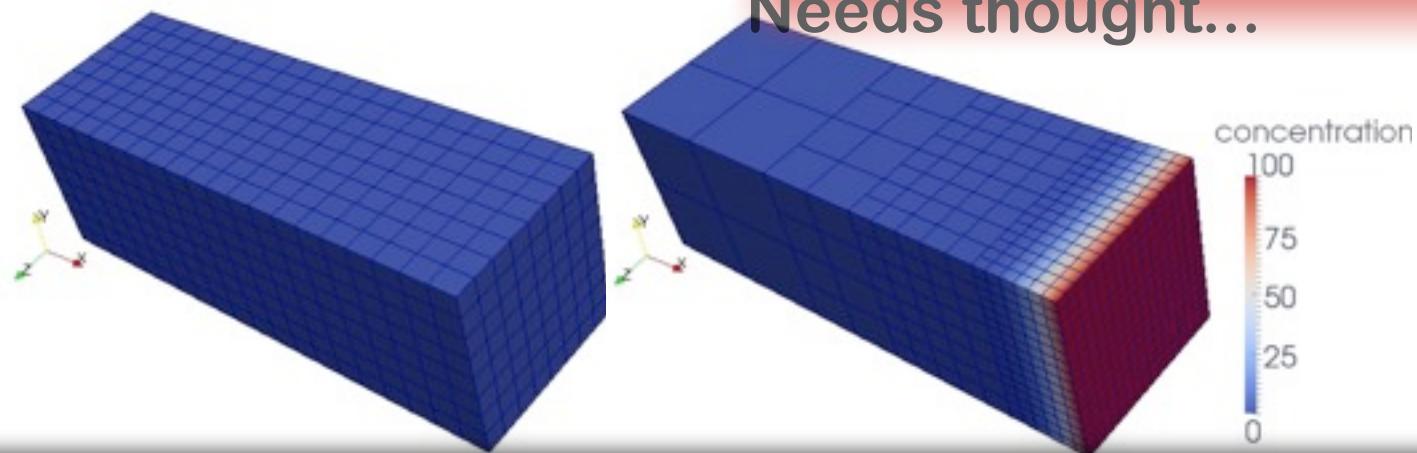
(Or several other suggestions made so far)

# Mesh adaptivity with internal variables



- Spatial adaptivity with nodal unknowns is mature within DEAL.II
  - Hanging nodes with continuity imposed via linear constraints
  - Fully parallelised implementation
- Spatial adaptivity with inelastic internal variables less well developed
  - Projection of quadrature point data between refinement levels
  - Storage of cell related data cumbersome and restricts parallel implementation

Discussed during workshop.  
Needs thought...

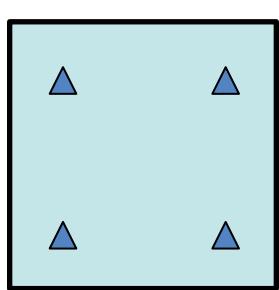


# Mesh adaptivity with internal variables

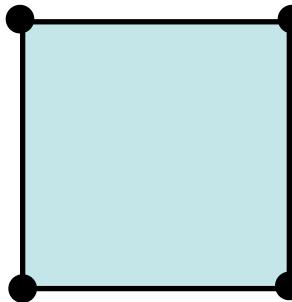
LM

1

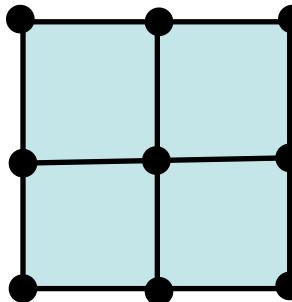
## Refinement



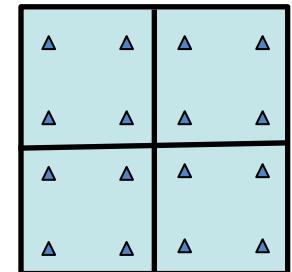
$$P_1^{\text{cG}}$$



$$P_2^{\text{pro}}[cc]$$



$$P_3^{\text{cG}}$$



$$\mathbf{y}' \in \mathbb{R}^{n_{\text{qp}}}$$

$$\mathbf{P}^{\text{ref}}[cc] = P_3^{\text{cG}} * P_2^{\text{pro}}[cc] * P_1^{\text{cG}}$$

$$\mathbf{y}[cc] = \mathbf{P}^{\text{ref}}[cc] * \mathbf{y}'$$

$$\mathbf{y}[cc] \in \mathbb{R}^{n_{\text{qp}}}$$

```
FE_Q<deal_II_dimension> fe_projection_cg(u_degree);

FETools::compute_projection_from_quadrature_points_matrix
(fe_projection_cg,...,P_1_cg);
FETools::compute_interpolation_to_quadrature_points_matrix
(fe_projection_cg,..., P_3_cg);
for (unsigned int cc = 0 ; cc < n_children ; cc++) {
    FullMatrix<double> P_2_i_prolongation =
        fe_projection_cg.get_prolongation_matrix(cc);
    P_refine[cc] = P3_P2_pro_P1;
}
```

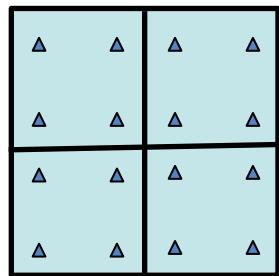
- All projection matrices are computed once and for all on the reference cell
- Method valid for all components of internal variables

# Mesh adaptivity with internal variables

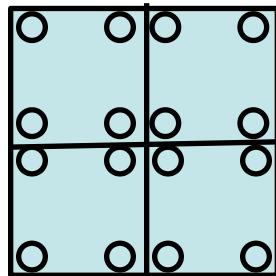


2

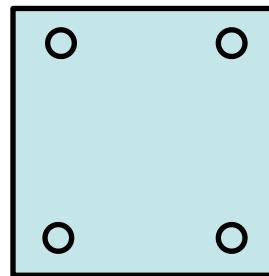
## Coarsening



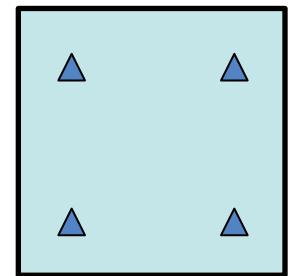
$$\mathbf{P}_1^{\text{dG}}$$



$$\sum_{cc} \mathbf{P}_2^{\text{res}}[cc]$$



$$\mathbf{P}_3^{\text{dG}}$$



$$\mathbf{y}'[cc] \in \mathbb{R}^{n_{\text{qp}}}$$

$$\mathbf{P}^{\text{coarse}}[cc] = \mathbf{P}_3^{\text{dG}} * \mathbf{P}_2^{\text{res}}[cc] * \mathbf{P}_1^{\text{dG}}$$

$$\mathbf{y} = \sum_{cc} \mathbf{P}^{\text{coarse}}[cc] * \mathbf{y}'[cc]$$

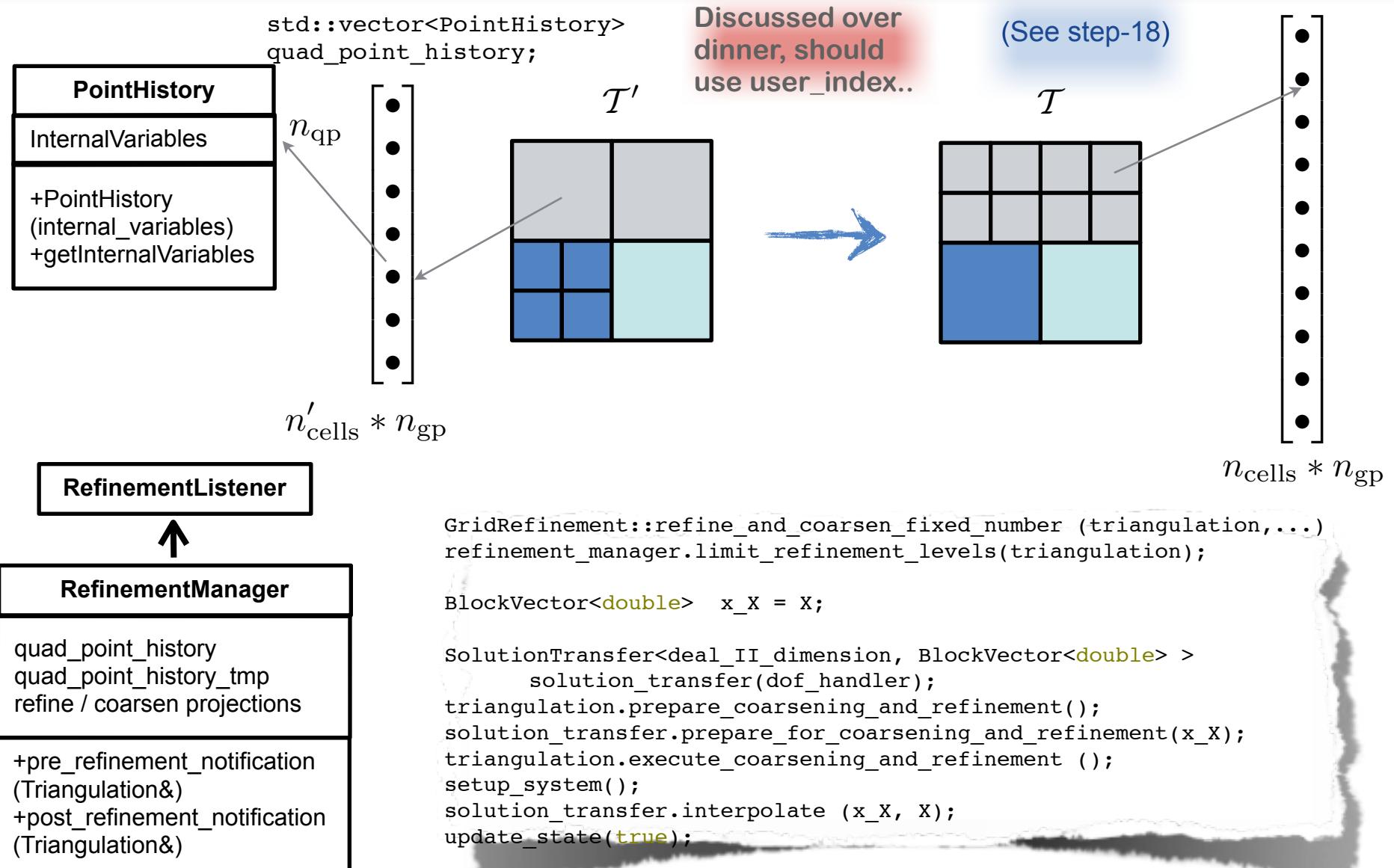
$$\mathbf{y} \in \mathbb{R}^{n_{\text{qp}}}$$

```
FE_DGQ<deal_II_dimension> fe_projection_dg(u_degree);

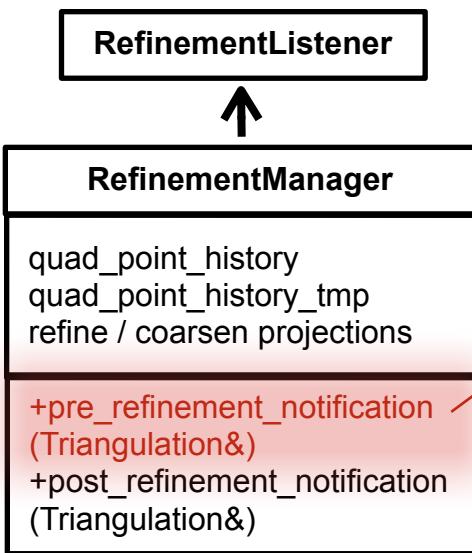
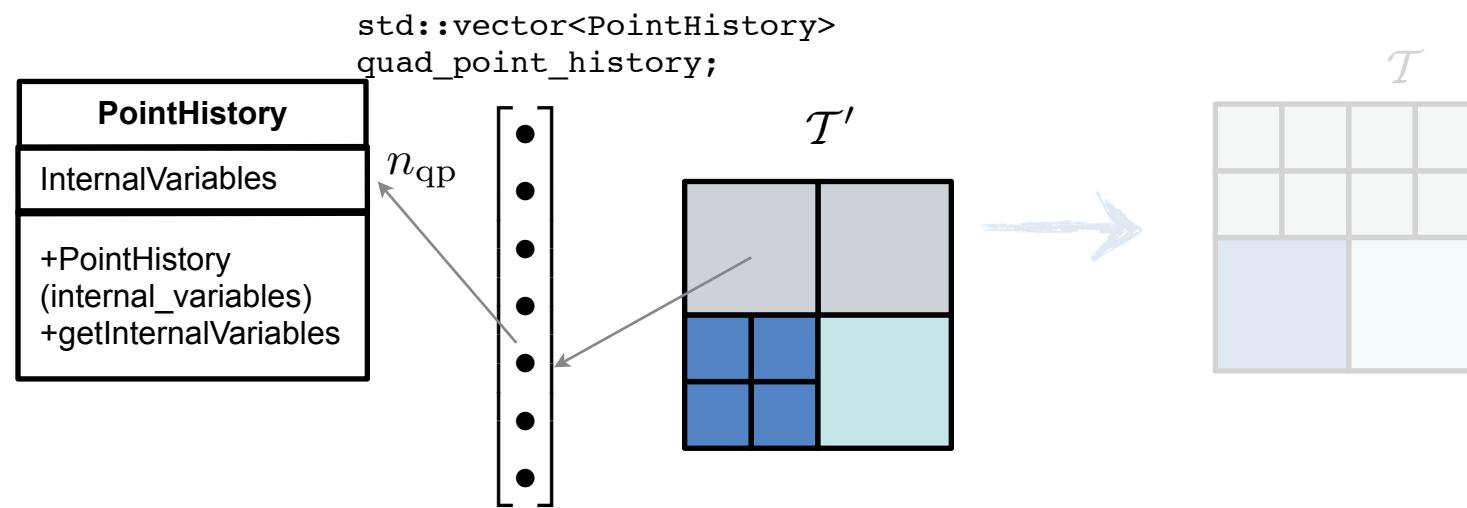
FETools::compute_projection_from_quadrature_points_matrix
(fe_projection_dg,...,P_1_dg);
FETools::compute_interpolation_to_quadrature_points_matrix
(fe_projection_dg,..., P_3_dg);
for (unsigned int cc = 0 ; cc < n_children ; cc++) {
    FullMatrix<double> P_2_i_restriction =
    fe_projection_dg.get_restriction_matrix(cc);
    P_coarsen[cc] = P3_P2_res_P1;
}
```

- All projection matrices are computed once and for all on the reference cell
- Note: the summation
- Method valid for all components of internal variables: i.e. works for scalars, vectors and symmetric second order tensors

# Mesh adaptivity with internal variables

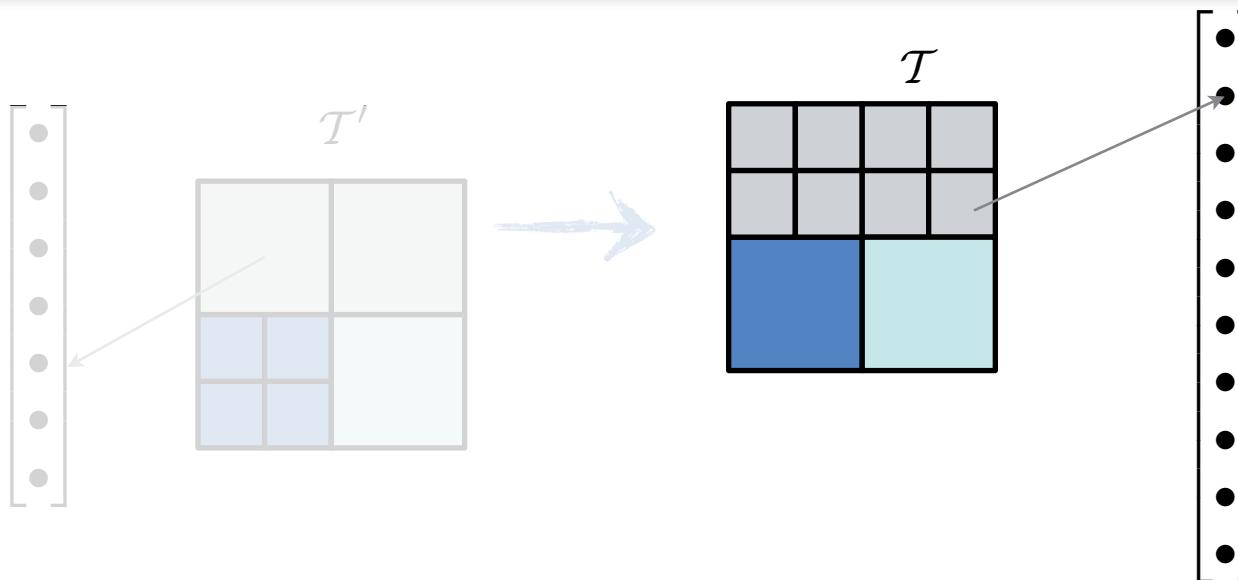


# Mesh adaptivity with internal variables

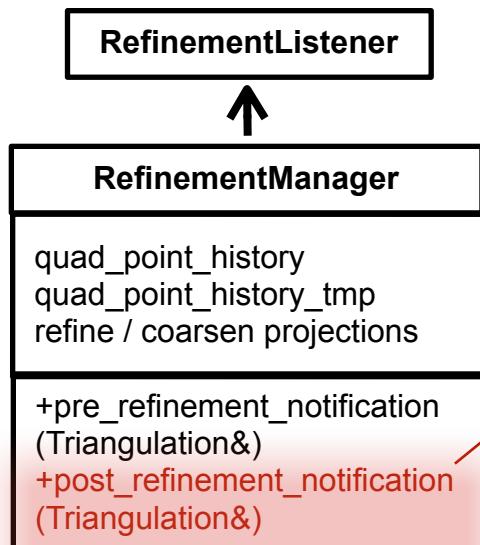


- Copy `quad_point_history` to `quad_point_history_tmp`
- `quad_point_history.clear()`
- Redirect the `cell->user_pointer()` to `quad_point_history_tmp`
- Record the cells to be refined, coarsened and to remain

# Mesh adaptivity with internal variables



Discussed over dinner, should use user\_index..



- Cells to remain:
  - Copy cell data from `quad_point_history_tmp` to `quad_point_history`
  - Redirect `cell->user_pointer`
- Cells to coarsen:
  - project data from children to parents
  - Add parent data to `quad_point_history`
  - Redirect `parent->user_pointer`
- Cells to refine:
  - project data from parents to children
  - Add children's data to `quad_point_history`
  - Redirect `children->user_pointer`

## Issues

- Children do not know their parents :(
  - Complicates the coarsening routine as one can't simply ask those children flagged for coarsening for their parent
- The need to store the cell data in a std::vector causes headaches
  - The data should be associated more tightly with a cell and management abstracted from the user
  - Abstract class CellData that the user can overload
  - Member data of this class include the projections to perform coarsening and refinement
- Simple to handle memory as the CellData is coupled to cell?
- Parallelisation friendly?
- The cell user\_pointer is used by deal.II code itself. Ideally a user should be “me” or “you” and not deal.II

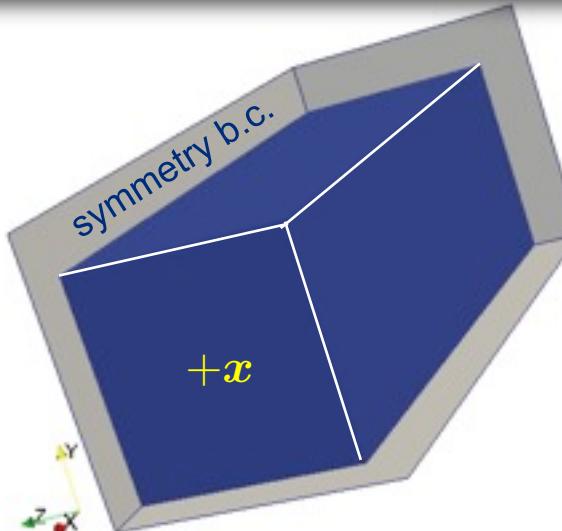
Already discussed during  
workshop

# Example problem

## Constitutive laws

$$\begin{aligned}
 W^0(\boldsymbol{\varepsilon}) &= \bar{W}^0(\boldsymbol{e}) + U^0(\Theta) \\
 &= \mu \boldsymbol{e} : \boldsymbol{e} + \frac{1}{2} \left[ \lambda + \frac{2}{3}\mu \right] \Theta^2 \\
 \nu &= \frac{1}{3} \quad \text{and} \quad \mu = 1000 \\
 D(c) &= \begin{cases} \frac{1}{10} & \text{if } m > 50 \\ \frac{1}{1000} & \text{otherwise} \end{cases} \\
 \tau(m) &= \begin{cases} \frac{1}{10} & \text{if } m > 50 \\ 10000 & \text{otherwise} \end{cases}
 \end{aligned}$$

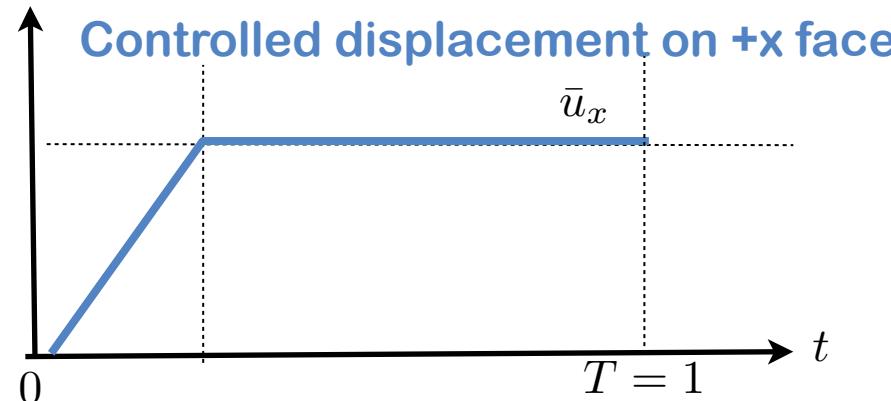
- Coupling is still only one-way: need to account for influence of the solvent on the swelling of the polymer
- Diffusivity and relaxation times are not functions of the deformation
- Polymeric solid still compressible



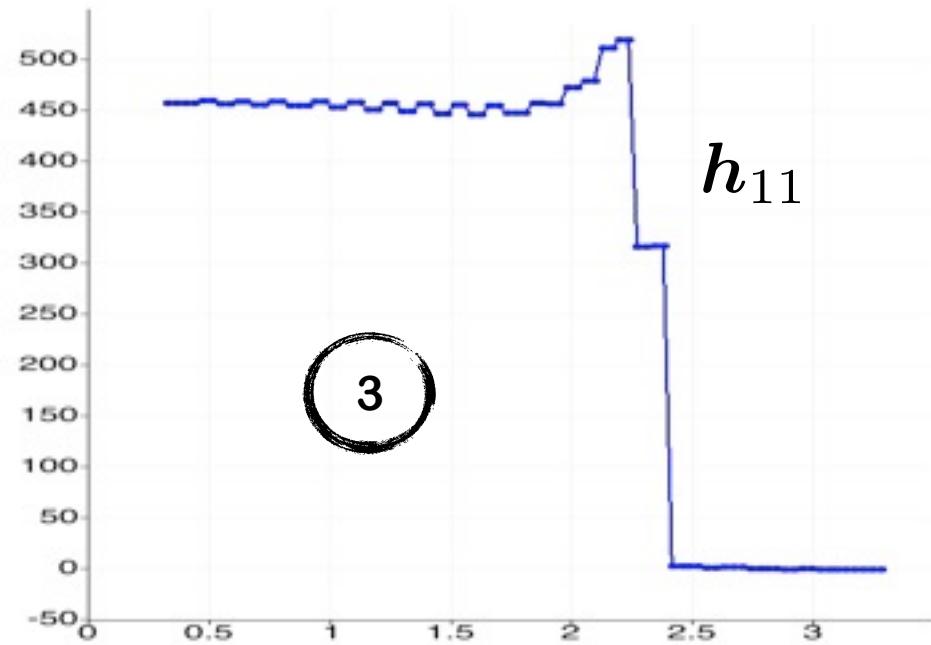
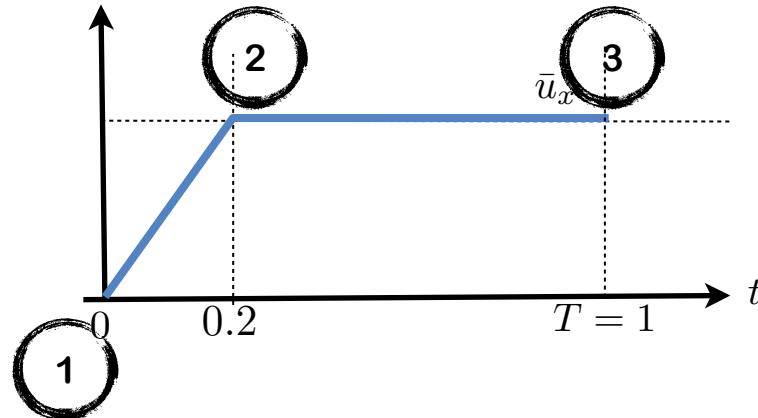
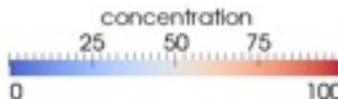
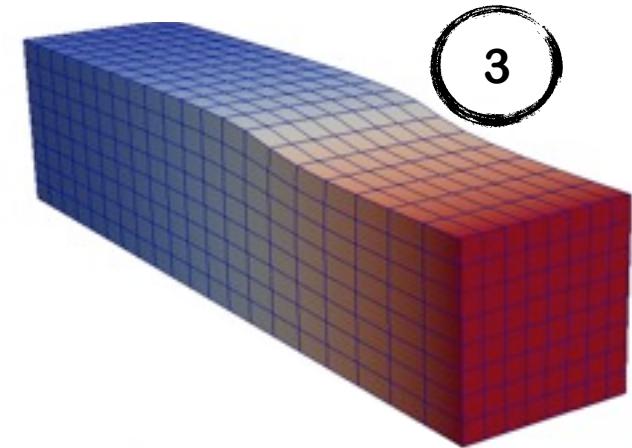
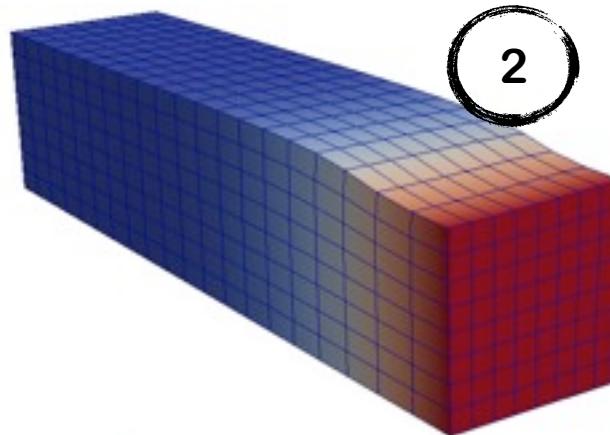
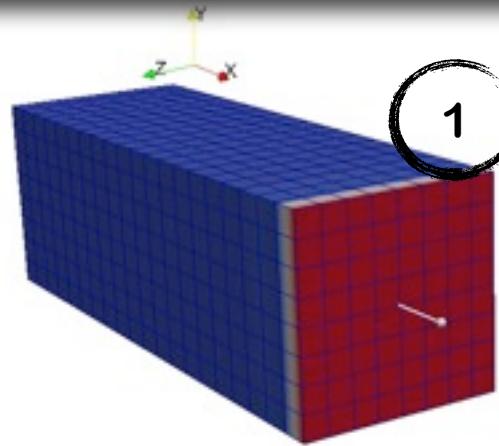
**Initial conditions and concentration shock boundary condition**

$$m(\boldsymbol{x}, t = 0) = 0$$

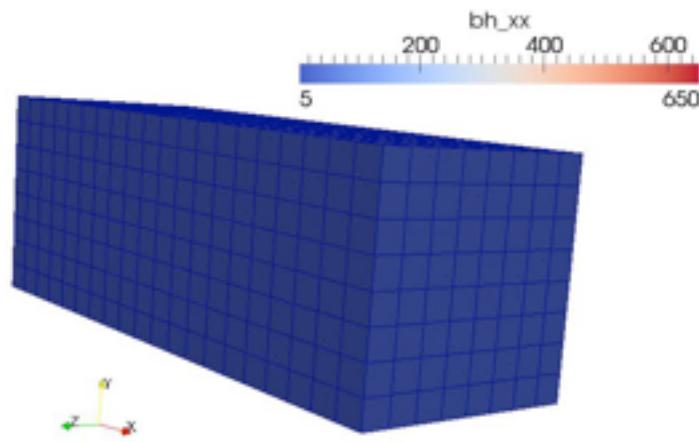
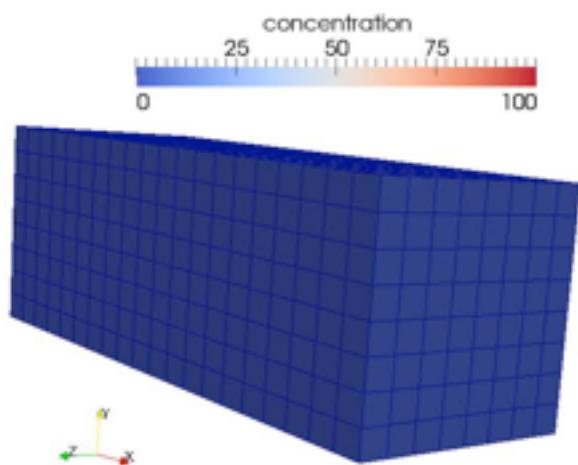
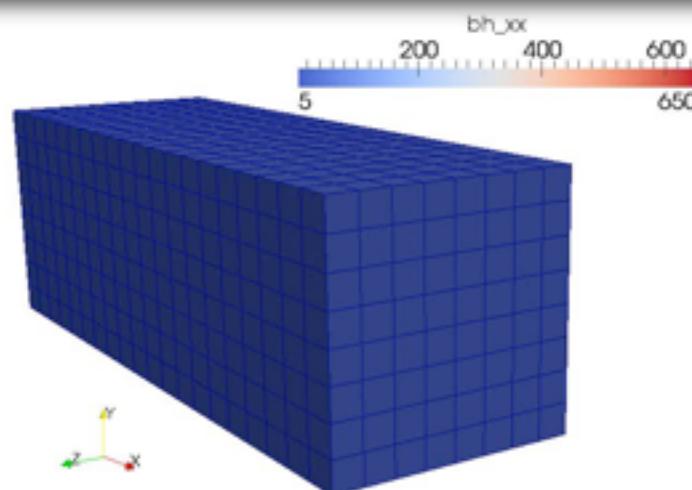
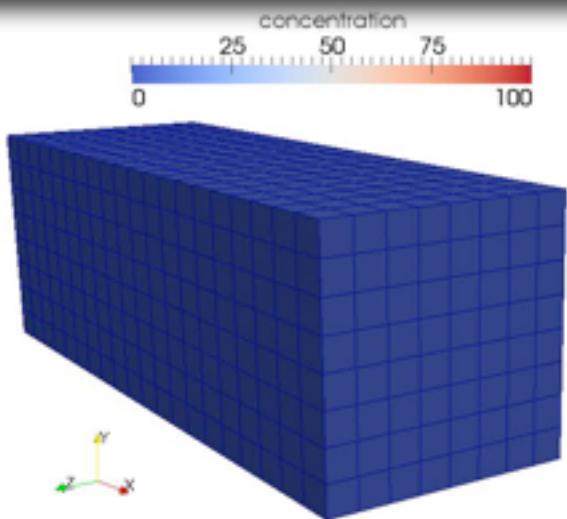
$$m(+\boldsymbol{x}, t = 0) = 100$$



# Example problem



# Example problem

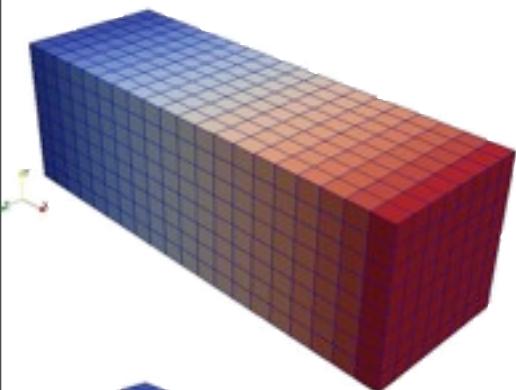


Refinement only...

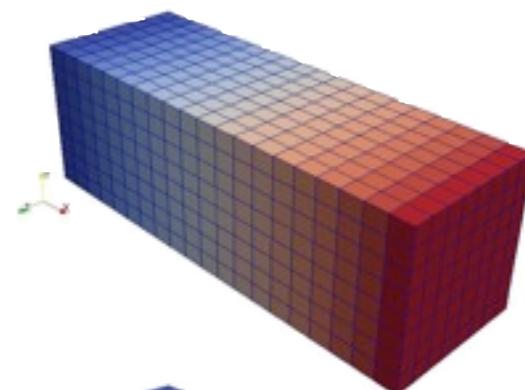
Example problem: the value of test cases. Add the x-coordinate as an internal variable...



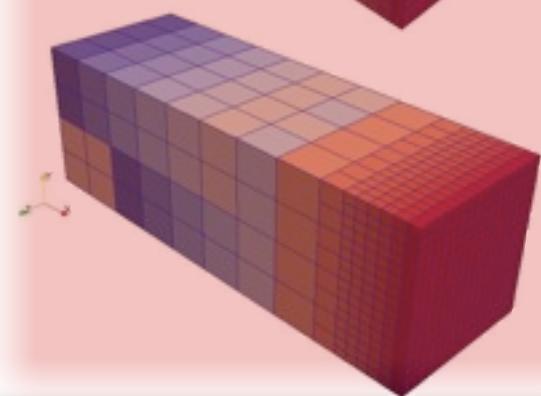
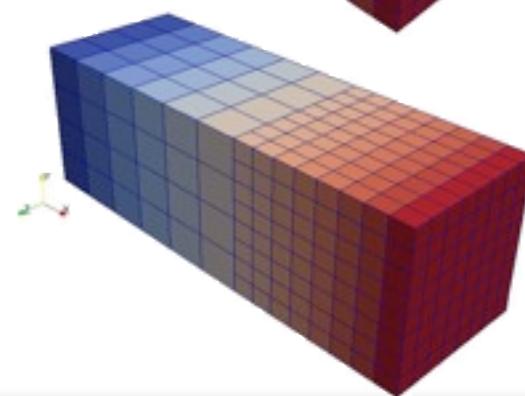
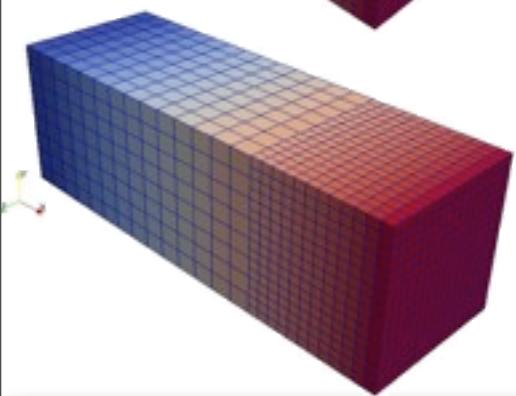
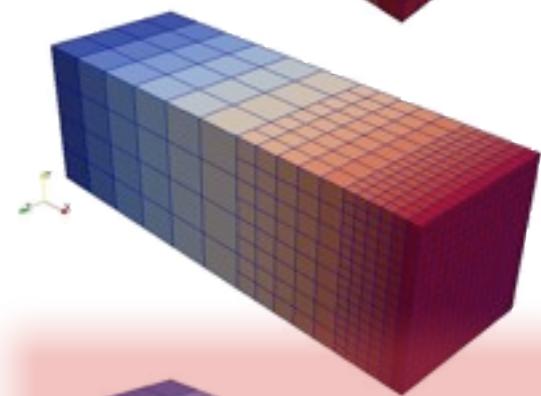
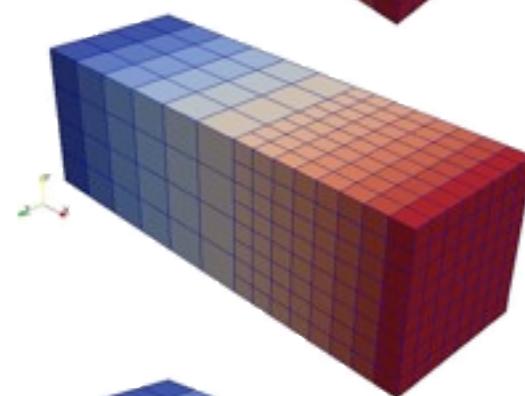
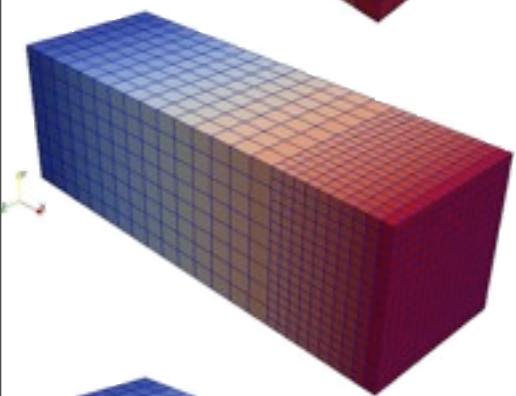
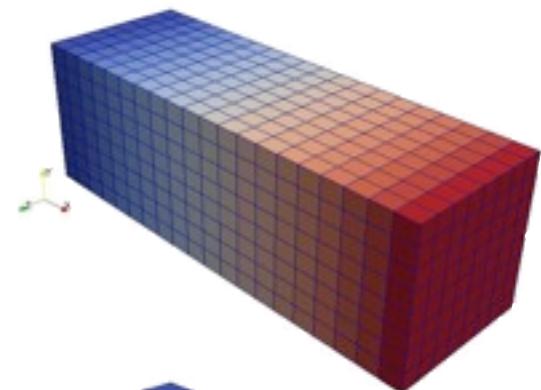
Refine



Coarsen



Refine and coarsen



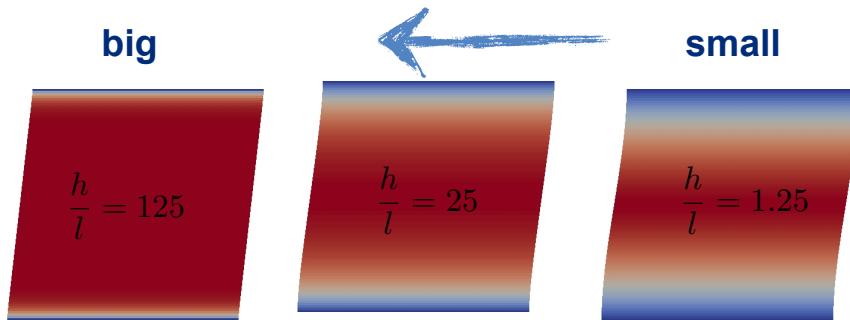
- Spatial adaptivity with internal variables is viable within current structure of deal.II
- Internal variable formulations are widely used in solid mechanics but few (no) solid mechanics codes offer the flexibility and capability of deal.II
- Extend the user base of deal.II if one could facilitate such formulations
  - Strategy to give internal variable data a similar status to nodal variables
  - refinement and projection of internal variables
  - parallel implementation

Discussed in deal.II future ideas

Work in progress: comments and suggestions greatly appreciated

# A word of thanks

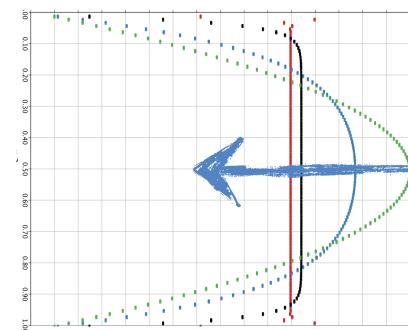
LM



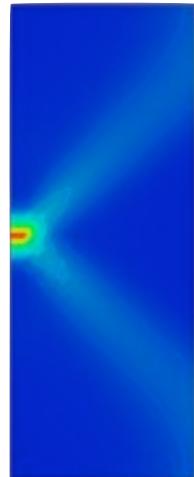
**Non-local crystal plasticity**  
McB, Reddy, Richardson, Gurtin

## Deal.II at the University of Cape Town,

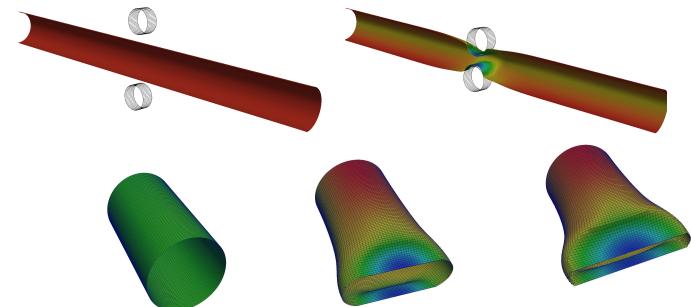
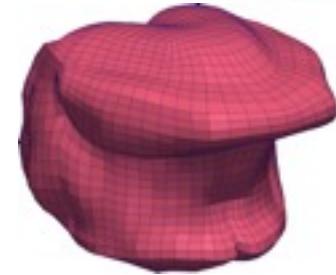
- Geographically distant and (very) small local support community
- Adopted deal.II as in-house code in 2008
- Small group of ~20 students in DEM, FEM, CFD, particle methods
- Completed or in-progress using deal.II:
  - 4 MSc and 4 PhD
- User group that meets regularly
- **Support and code greatly appreciated!  
Made projects possible and is a fantastic learning tool**



**Algorithms for  
crystal plasticity**  
McB, Reddy,  
Richardson



**Shell formulations with  
applications in biomechanics**  
Bartle, Reddy, McB



# Thanks. Comments, suggestions and questions please

The financial support of the German Science Foundation (Deutsche Forschungsgemeinschaft, DFG), grant STE 544/39-1, and the National Research Foundation of South Africa is gratefully acknowledged.

The support of all those involved in the deal.II project is most appreciated! thanks.

Thanks to Michael Rapson and Jean-Paul Pelteret

