

# AST3220 - Project 1

(Dated: March 30, 2023)

This abstract is abstract.

## 1. PROBLEM 1

We can rewrite the number density  $n_i$  of species  $i$  in terms of the relative number density  $Y_i$  as:

$$Y_i = \frac{n_i}{n_b} \Rightarrow n_i = n_b Y_i \quad (1.1)$$

$$= \frac{n_{b0}}{a^3} Y_i \quad (1.2)$$

Where  $n_b(t)$  is the baryon number density,  $n_{b0}$  is the baryon number density today, and  $a(t)$  is the scale factor. Using the product rule for differentiation, we can then write

$$\frac{dn_i}{dt} = n_{b0} \frac{d}{dt} (Y_i a^{-3}) \quad (1.3)$$

$$= n_{b0} \left( \frac{1}{a^3} \frac{dY_i}{dt} - 3Y_i \frac{\dot{a}}{a^4} \right) \quad (1.4)$$

$$= n_b \frac{dY_i}{dt} - 3Y_i n_b H \quad (1.5)$$

$$= n_b \frac{dY_i}{dt} - 3n_i H \quad (1.6)$$

Next we want to switch from  $t$  to  $\ln T$  as our time variable, where  $T$  is the temperature. Using  $T = T_0 a^{-1}$  we get

$$\ln T = \ln T_0 - \ln a(t) \quad (1.7)$$

Then, using the chain rule of differentiation, we can rewrite

$$\frac{dY_i}{dt} = \frac{d(\ln T)}{dt} \frac{dY_i}{d(\ln T)} \quad (1.8)$$

$$= -\frac{\dot{a}}{a} \frac{dY_i}{d(\ln T)} \quad (1.9)$$

$$= -H \frac{dY_i}{d(\ln T)} \quad (1.10)$$

Inserting to equation (1.6) we get

$$\frac{dn_i}{dt} = -n_b H \frac{dY_i}{d(\ln T)} - 3n_i H \quad (1.11)$$

The equations for the evolution of the number densities of protons  $p$  and neutrons  $n$  are given as

$$\frac{dn_n}{dt} + 3Hn_n = n_p \Gamma_{p \rightarrow n} - n_n \Gamma_{n \rightarrow p} \quad (1.12)$$

$$\frac{dn_p}{dt} + 3Hn_p = n_n \Gamma_{n \rightarrow p} - n_p \Gamma_{p \rightarrow n} \quad (1.13)$$

$$= - \left( \frac{dn_n}{dt} + 3Hn_n \right) \quad (1.14)$$

And by inserting Eq. (1.1) and Eq. (1.6) we finally find the evolution of the relative number densities:

## 2. PROBLEM 2

The relation  $T_\nu = (4/11)^{1/3} T$  can be derived from the conservation of entropy, which tells us that

$$g_{*s}(aT)^3 = \text{const.} \quad (2.1)$$

At the time where the universe had a temperature  $k_B T$  0.511 MeV, electrons and positrons were relativistic and the process

$$e^+ + e^- \rightarrow \gamma + \gamma \quad (2.2)$$

occurred in both directions. However, as the temperature universe falls below the rest mass of the electron and positron  $k_B T$  0.511, the average energy of a photon collision is too small for an electrons-positron pair to be created. Since electrons and positrons will still annihilate through the process

$$e^+ + e^- \rightarrow \gamma + \gamma \quad (2.3)$$

most of the positrons and electrons will then disappear. Assuming this happened immediately, and that the universe is in thermal equilibrium ( $T_i = T$ ), the effective number of degrees of freedom before and after can be written

$$g_{*s}^{\text{before}} = g_\nu + \frac{7}{8}(g_{e^-} + g_{e^+}) \quad (2.4)$$

$$= 2 + \frac{7}{8}4 \quad (2.5)$$

$$= \frac{11}{2} \quad (2.6)$$

$$g_{*s}^{\text{after}} = g_\nu \quad (2.7)$$

$$= 2 \quad (2.8)$$

If we also assume the scale factor  $a$  is the same before and after, the conservation of entropy gives us

$$\frac{11}{2}(aT)_{\text{before}}^3 = 2(aT)_{\text{after}}^3 \quad (2.9)$$

$$\Rightarrow T_{\text{after}} = \left(\frac{11}{4}\right)^{1/3} T_{\text{before}} \quad (2.10)$$

Since neutrinos are decoupled, we then have

$$T_{\nu,after} = T_{\nu,before} = T_{before} = \left(\frac{4}{11}\right)^{1/3} T_{after} \quad (2.11)$$

Finally giving us

$$T_{\nu} = \left(\frac{4}{11}\right)^{1/3} T \quad (2.12)$$

## ACKNOWLEDGMENTS

I would like thank myself for writing this beautiful document.

## REFERENCES

- Reference 1

- Reference 2

## Appendix A: Name of appendix

This will be the body of the appendix.

## Appendix B: This is another appendix

Tada.