

Coupling of continuous and hybridizable discontinuous Galerkin methods: Application to conjugate heat transfer problem

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Motivation

Continuous Galerkin (CG)

- ✓ Well established
- ✓ Computationally efficient
- ✓ Static condensation
- ✗ Unstable for high convective flows
- ✗ Cumbersome p-adaptivity

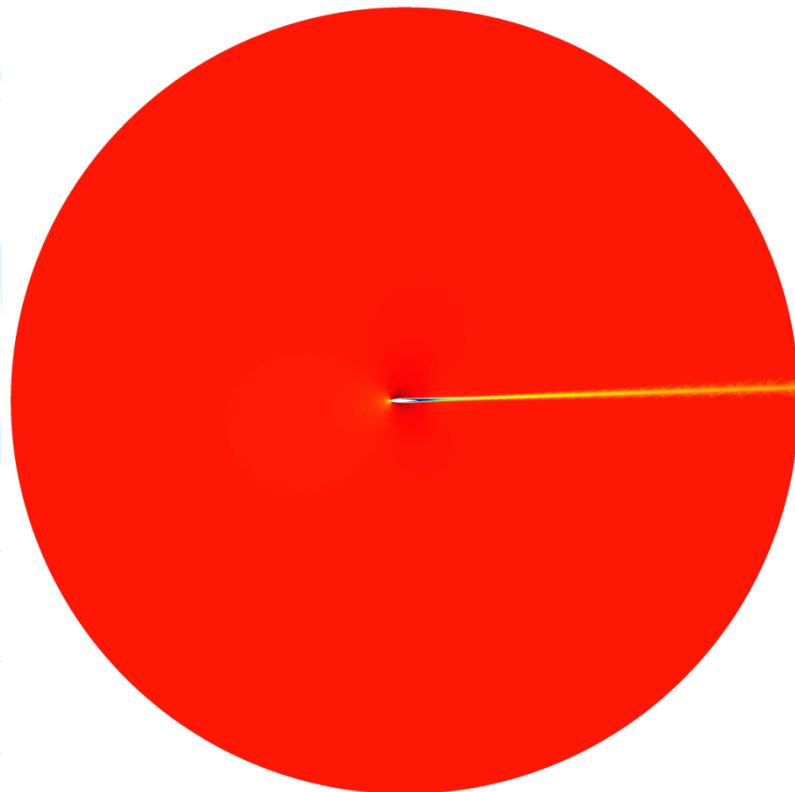
DG (DG, IPM, CDG, LDG..)

- ✗ Duplication of nodes
- ✗ No static condensation
- ✓ Stable for convective flows
- ✓ Suitability for adaptivity, parallel computing, ...

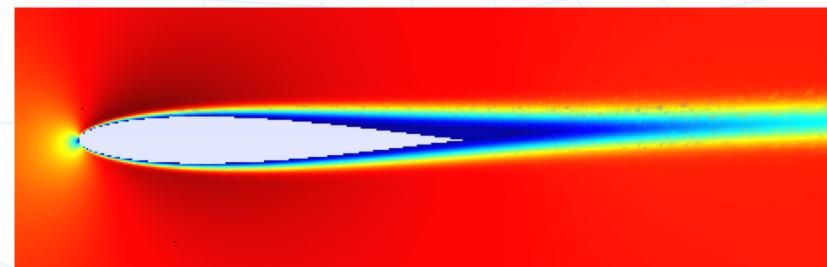
Hybridizable Discontinuous
Galerkin (HDG)

NACA0012 Airfoil

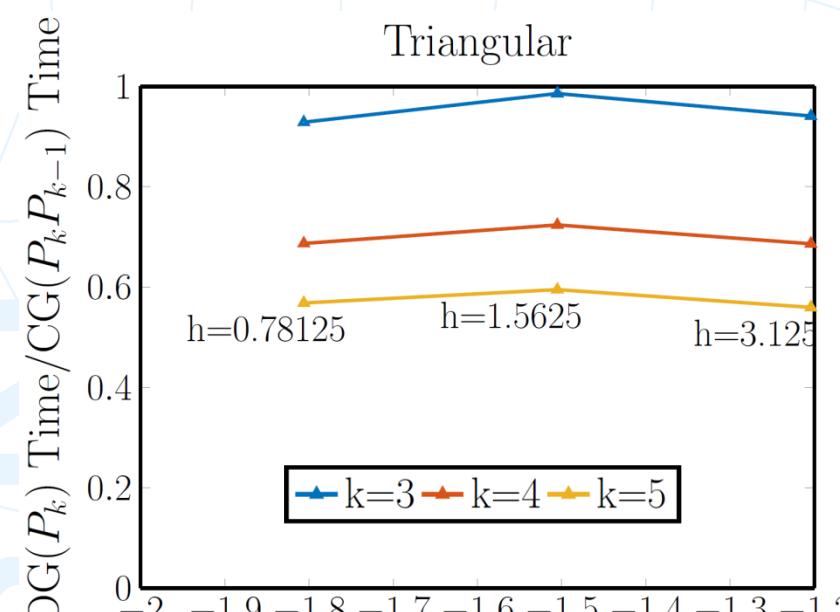
Incompressible Navier-Stokes is solved at $Re = 5000$ and Angle of attack (AoA) = 2° ; Error in lift coefficient is compared



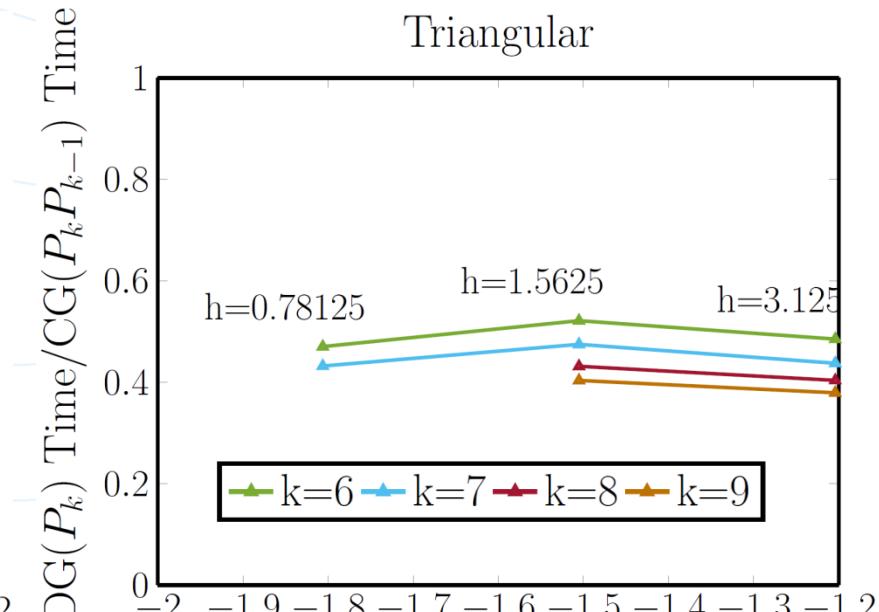
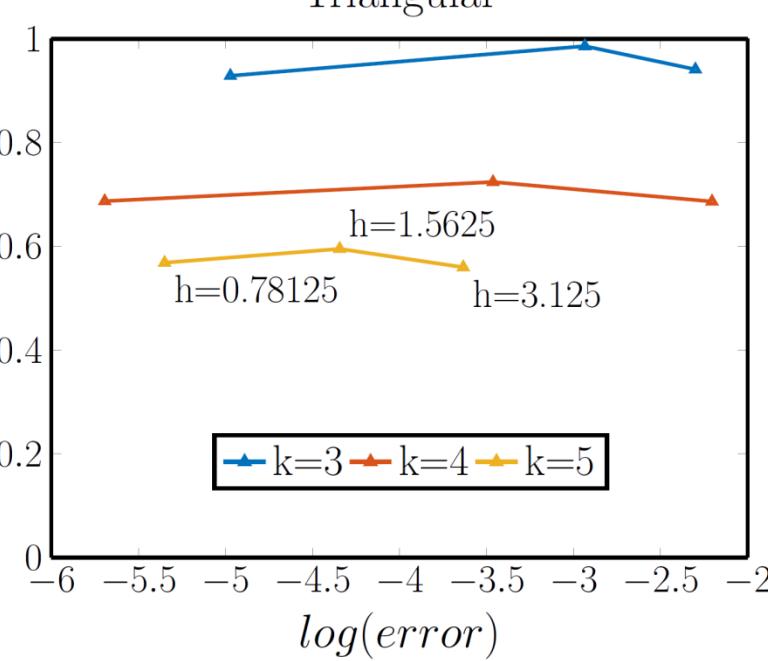
Velocity



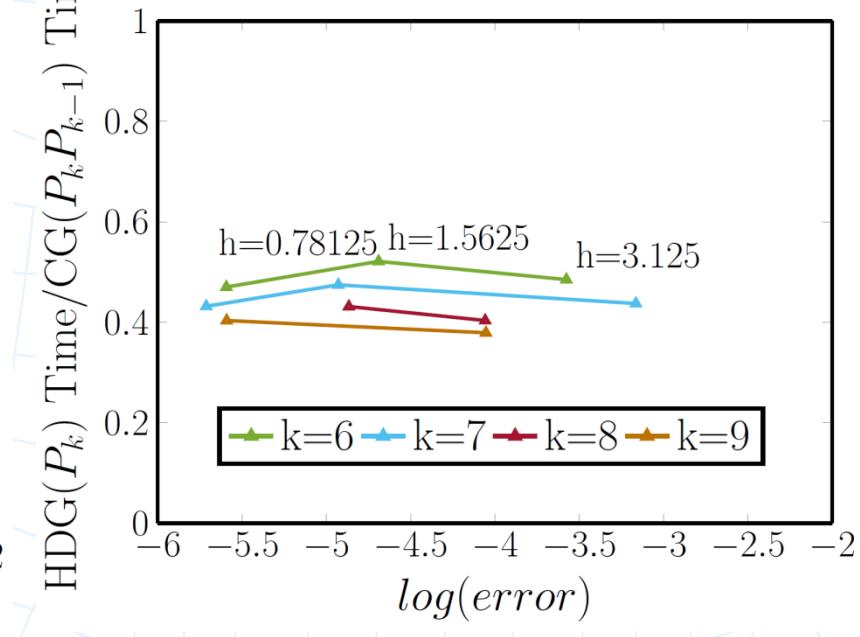
NACA: CPU Time for linear solver



Triangular



Triangular



Motivation

GFRP Problem

- Multi-physics transient problem
- Time varying temperature is prescribed on Γ_b .
- Thermophysical properties of GFRP are temperature dependent
- Bousinesseq approximation is used to model the natural buoyancy of air.

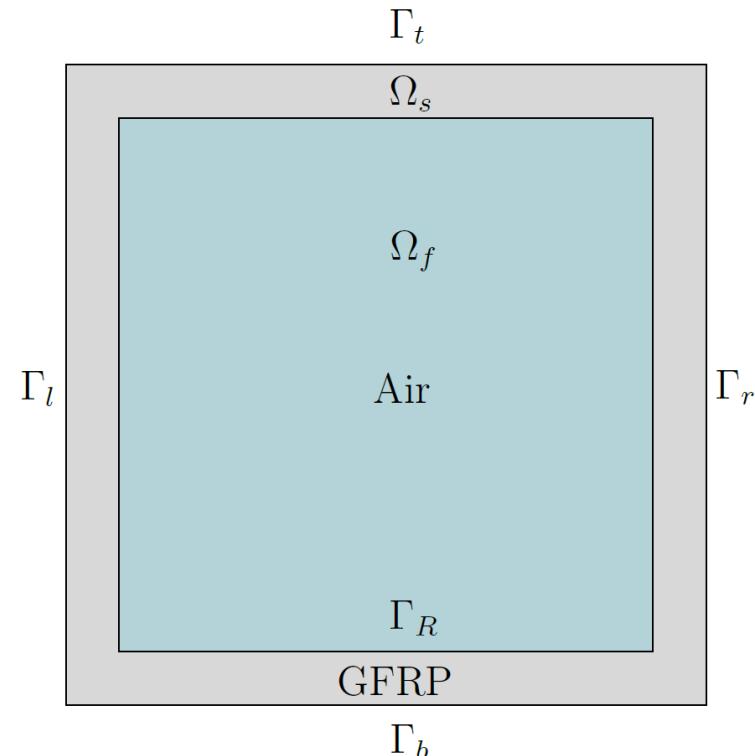
Γ_b, Γ_t - Convective and radiative BC

Γ_l, Γ_r - Adiabatic BC

Γ_R - Internal radiation

1) Coupling of CG and HDG

2) Coupled Navier-Stokes/heat transfer for HDG



Coupling of CG and HDG

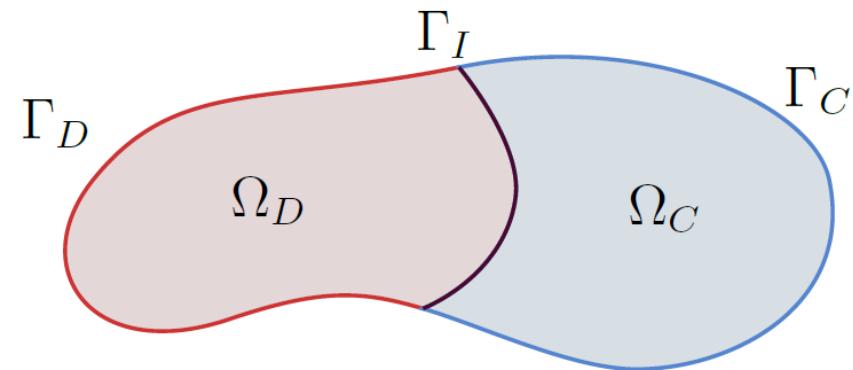
Heat equation

$$-\nabla \cdot (k \nabla u) = f \quad \text{in } \Omega_D$$

$$u = u_D \quad \text{on } \Gamma_D$$

$$-\nabla \cdot (k \nabla u) = f \quad \text{in } \Omega_C$$

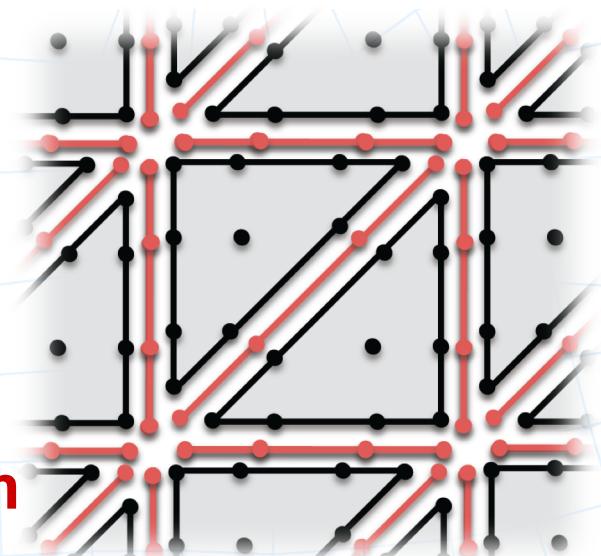
$$u = u_C \quad \text{on } \Gamma_C$$



Suffixes D and C stands for HDG and CG domains. Interface conditions are

$$u|_D - u|_C = 0 \quad \text{on } \Gamma_I$$

$$-k \nabla u \cdot n|_D - k \nabla u \cdot n|_C = 0 \quad \text{on } \Gamma_I$$



Global Problem

$$[\hat{\mathbf{q}} \cdot \mathbf{n}] = 0 \quad \text{in } \Gamma$$

$$\hat{u} = u_D \quad \text{on } \Gamma_D$$

$$\Gamma = \bigcup_{i=1}^{n_{el}} \partial\Omega_{iD}$$

Local Problem

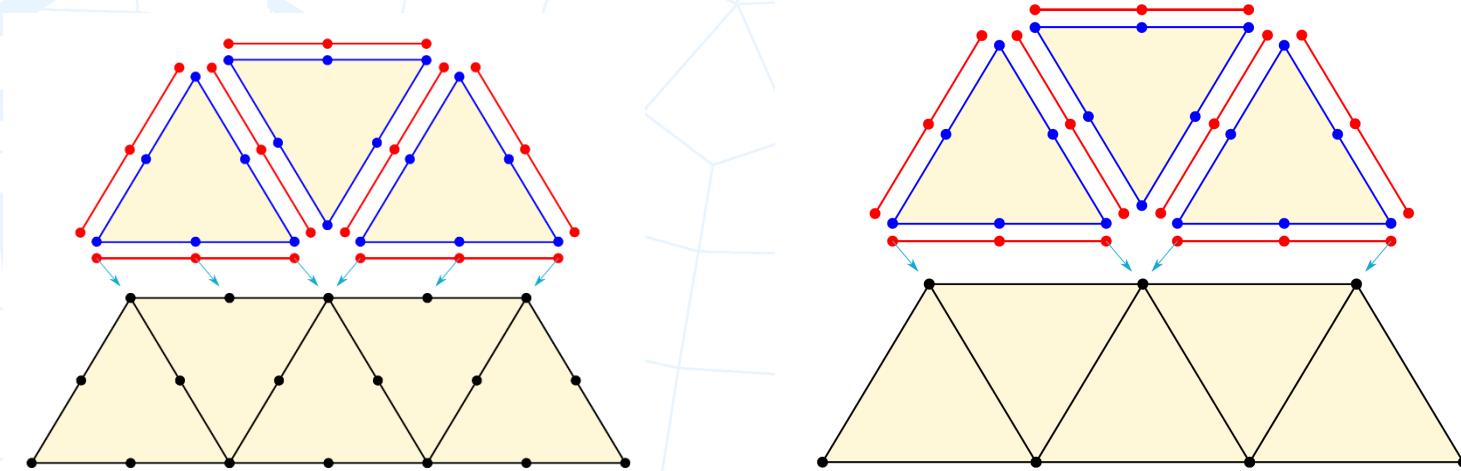
$$\nabla \cdot \mathbf{q} = f \quad \text{in } \Omega_{iD}$$

$$\mathbf{q} + k \nabla u = 0 \quad \text{in } \Omega_{iD}$$

$$u = \hat{u} \quad \text{on } \partial\Omega_{iD}$$

Coupling of CG and HDG

Continuity of u is strongly imposed between u in Ω_C and \hat{u} in Ω_D .



If the degree of approximation is different on both domains, the solution is projected from the domain with **lower degree** to the domain with **higher degree** of approximation

$$\begin{aligned}\mathbf{u}^D &= \mathbf{P} \mathbf{u}^C \\ P_{ij} &= \Psi_j^C(\xi_i^D)\end{aligned}$$

Coupling of CG and HDG

Weak form of CG

$$(\nabla \delta u, \nabla u)_{\Omega_C} - \langle \delta u, k (\nabla u \cdot \mathbf{n}) \rangle_{\Gamma_I} + (\delta u, f)_{\Omega_C} = 0$$

Weak form of HDG (global equations)

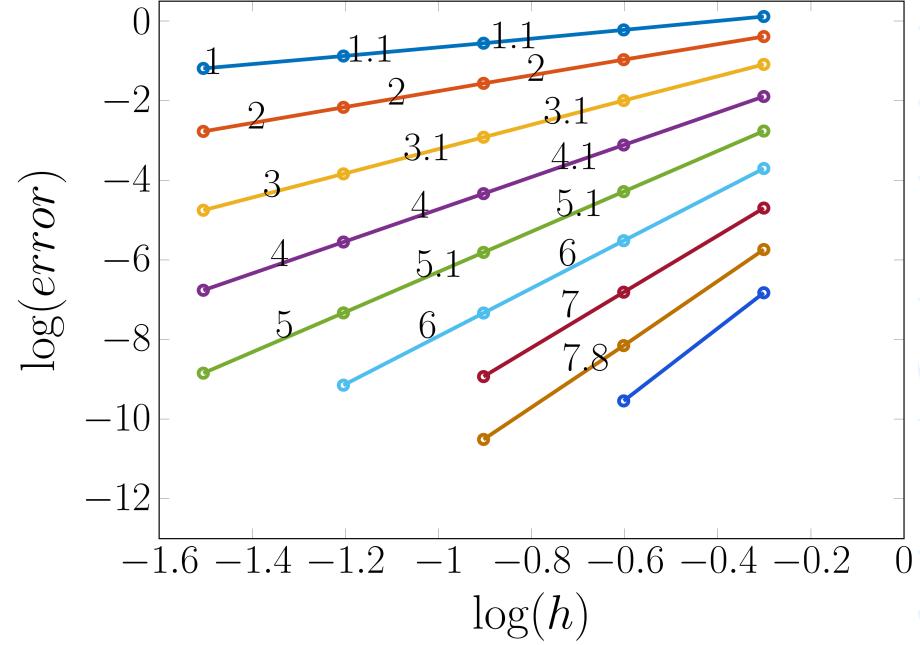
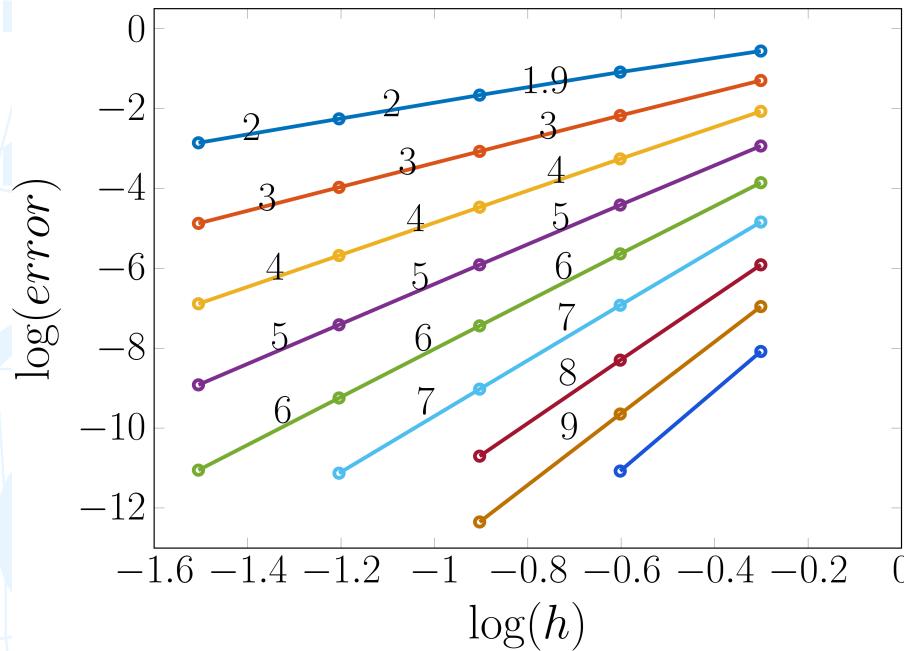
$$\langle \delta \hat{u}, [\![\hat{\mathbf{q}} \cdot \mathbf{n}]\!] \rangle_{\Gamma} = \langle \delta \hat{u}, [\![\hat{\mathbf{q}} \cdot \mathbf{n}]\!] \rangle_{\Gamma_i} + \langle \delta \hat{u}, [\![\hat{\mathbf{q}} \cdot \mathbf{n}]\!] \rangle_{\Gamma_D} + \langle \delta \hat{u}, [\![\hat{\mathbf{q}} \cdot \mathbf{n}]\!] \rangle_{\Gamma_I}$$

$$\langle \delta \hat{u}, [\![\hat{\mathbf{q}} \cdot \mathbf{n}]\!] \rangle_{\Gamma_I} = \langle \delta \hat{u}, \hat{\mathbf{q}} \cdot \mathbf{n} \rangle_{\Gamma_I} |_D + \langle \delta \hat{u}, -k \nabla u \cdot \mathbf{n} \rangle_{\Gamma_I} |_C$$

Weak form on interface

$$\begin{aligned} & \langle \delta \hat{u}, [\![\hat{\mathbf{q}} \cdot \mathbf{n}]\!] \rangle_{\Gamma_i} + \langle \delta \hat{u}, [\![\hat{\mathbf{q}} \cdot \mathbf{n}]\!] \rangle_{\Gamma_D} + \langle \delta \hat{u}, \hat{\mathbf{q}} \cdot \mathbf{n} \rangle_{\Gamma_I} |_D - \langle \delta \hat{u}, k (\nabla u_h \cdot \mathbf{n}) \rangle_{\Gamma_I} |_C \\ & - (\nabla \delta \hat{u}, \nabla v_h)_{\Omega_C} + \langle \delta \hat{u}, k (\nabla u_h \cdot \mathbf{n}) \rangle_{\Gamma_I} |_C - (\delta \hat{u}, f)_{\Omega_C} = 0 \end{aligned}$$

Coupling of CG and HDG



● P1 ● P2 ● P3 ● P4 ● P5
● P6 ● P7 ● P8 ● P9

 u ∇u

Coupled Navier-Stokes/heat transfer with HDG

Governing equations

Incompressible Navier-Stokes and convection diffusion equations.

$$\begin{aligned}\nabla \cdot (\mathbf{u} \otimes \mathbf{u}) - \nabla \cdot (-p\mathbf{I} + \nu \nabla \mathbf{u}) &= \mathbf{f}(\theta) \quad \text{in } \Omega \\ \nabla \cdot \mathbf{u} &= 0 \quad \text{in } \Omega \\ \mathbf{u} \cdot \nabla \theta - \nabla \cdot (k \nabla \theta) &= \mathbf{b} \quad \text{in } \Omega \\ \mathbf{u} &= \mathbf{u}_D \quad \text{on } \partial\Omega \\ \theta &= \theta_D \quad \text{on } \partial\Omega\end{aligned}$$

Bousinesseq approximation is used to model the natural buoyancy flows in closed cavities.

$$\mathbf{f}(\theta) = -\beta \mathbf{g}(\theta - \theta_0)$$

The approximation states that change in density of air due to temperature difference acts as body force on fluid itself

Coupled Navier-Stokes/heat transfer with HDG

Governing equations

Kovasznay flow is considered for Navier-Stokes equations in domain $[0 \ 2]x[-0.5 \ 1.5]$.

$$\mathbf{u} = \begin{bmatrix} 1 - \exp(\lambda x_1) \cos(2\pi x_2) \\ \frac{\lambda}{2\pi} \exp(\lambda x_1) \sin(2\pi x_2) \end{bmatrix}$$
$$p = -\frac{1}{2} \exp(2\lambda x_1) + C,$$

$$\theta = \sin\left(\frac{\pi m x_1}{2a}\right) \sin\left(\frac{\pi n x_2}{2b} + \frac{\pi n}{4b}\right)$$

The body force considered in the analysis is expressed as

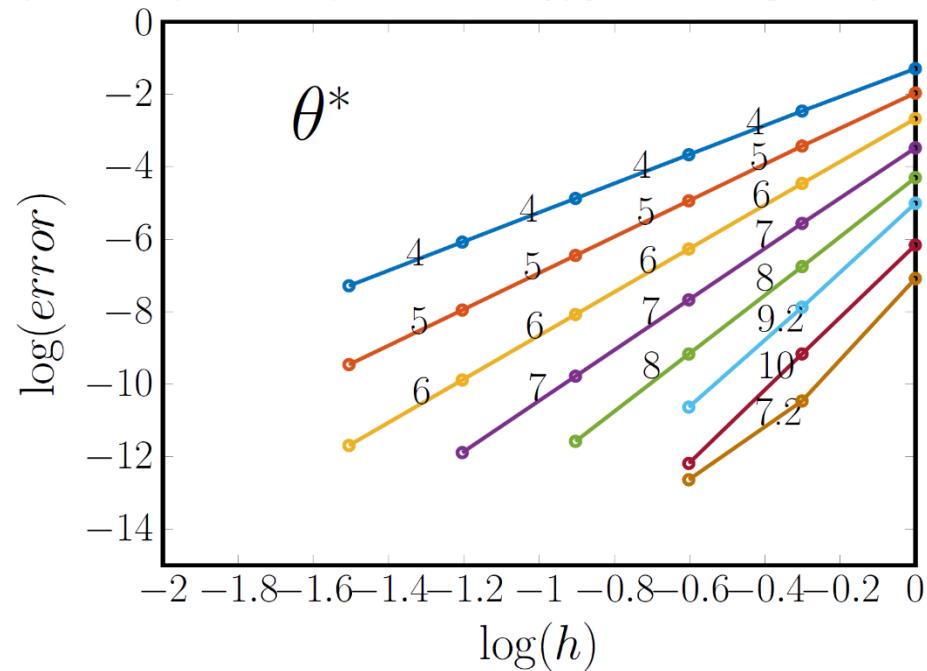
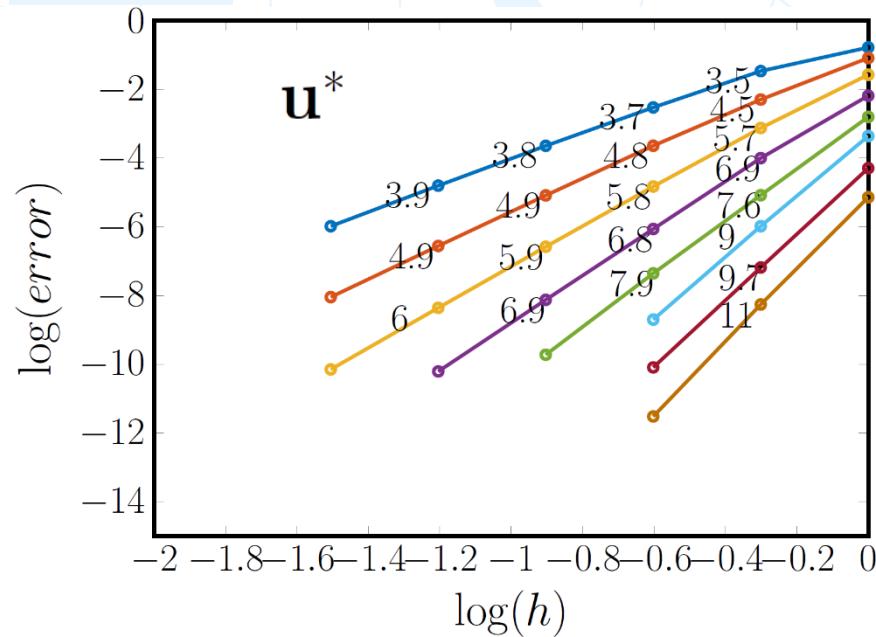
$$\begin{aligned} f(\theta) &= f_n(\theta) - f_a(\theta) \\ f_n(\theta) &= -\beta g \theta \\ f_a(\theta) &= -\beta g \theta_a \end{aligned}$$

$$m = 3, n = 3; a = b = 2; \beta = 1; g = (0, -10)$$

Coupled Navier-Stokes/heat transfer with HDG

Convergence results

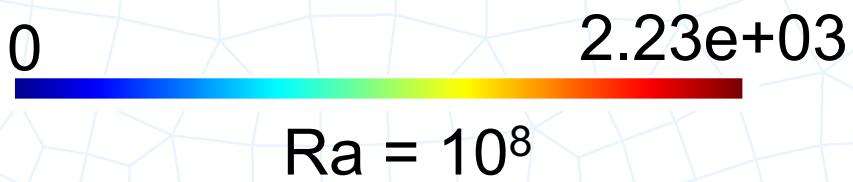
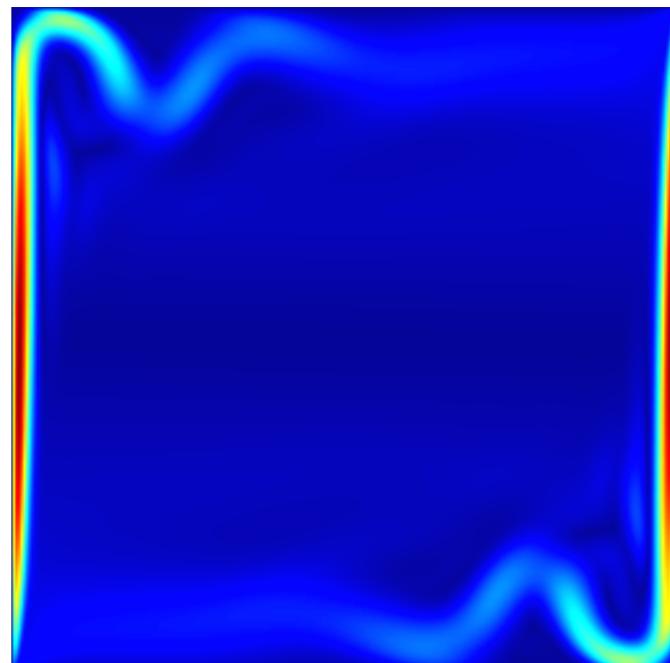
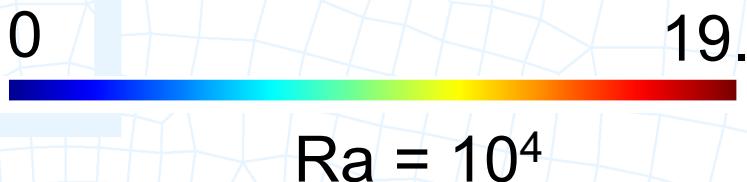
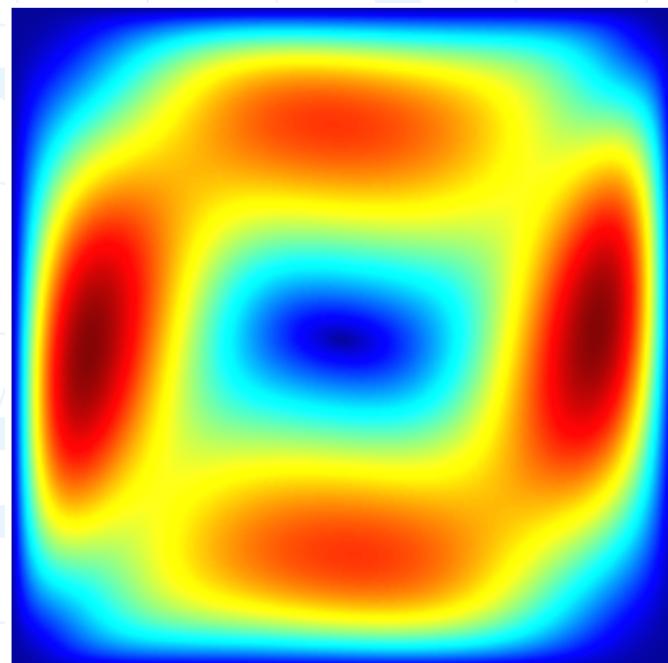
Kovasznay flow is considered for Navier-Stokes equations in domain $[0 \ 2] \times [-0.5 \ 1.5]$.



Coupled Navier-Stokes/heat transfer with HDG

Rayleigh-Bernard flow: Velocity contour plots

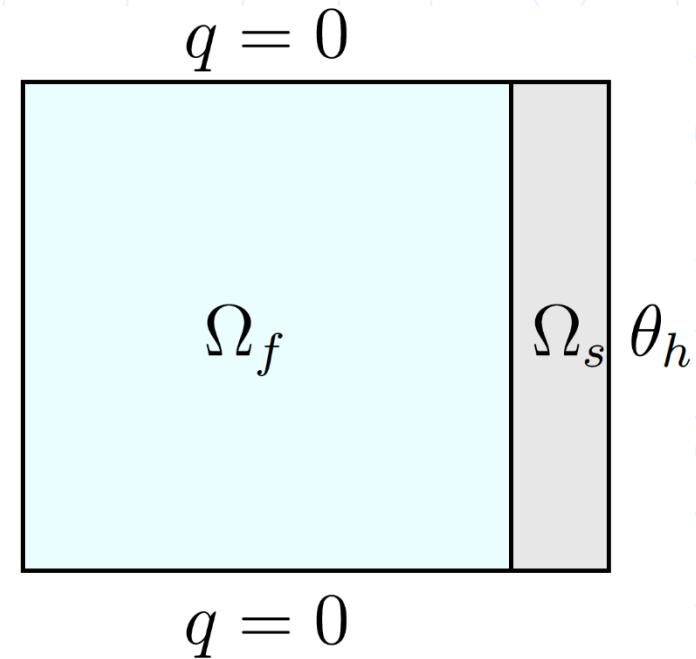
Square cavity of unit length and using Rayleigh number ($g\beta\Delta\theta L^3/v\alpha$). Triangular elements with $h = 0.03125$, $k = 5$.



Coupled HDG and CG method for conjugate HT

Governing equations

- Solid part of the conjugate heat transfer problem is discretized using CG.
- Fluid part is solved using HDG.
- Coupling presented for heat equation is extended to the problem.
- A benchmark results are presented for domain.
- Different Raleigh numbers are considered.



Concluding remarks

- Computational efficiency between HDG and CG for Stokes problem
 - High order-coarser mesh is more efficient than low order-finer mesh in both HDG and CG.
 - HDG is more efficient at order $k \geq 3$ for same level of accuracy, in terms of CPU time of linear solver
- Similar computational efficiency is observed in Navier-Stokes example (NACA).