# EP 1027: Maxwell's Equations and Electromagnetic Waves

Instructor: Shubho Roy<sup>1</sup> (Dept. of Physics)

Lecture 9 & 10

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► Wave equation basics: d'Alembert's solution, Bernoulli's solution, superposition

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EM waves in vacuum

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- ► Laws of reflection and refraction for waves

► Fresnel Equations

# References/Readings

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Griffiths, D.J., Introduction to Electrodynamics, Ch.9

Generic Wave equation (in 1D),

$$\left(\frac{\partial^2}{\partial x^2} - \frac{1}{c^2}\frac{\partial^2}{\partial t^2}\right)f(x,t) = 0,$$

c is the wave-velocity

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$$f(x,t) = g(x-ct) + h(x+ct).$$

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Bernoulli's solution: Sinusoidal

$$f(x, t) = A\cos(kx - \omega t + \Delta_1) + B\cos(kx + \omega t + \Delta_2), \omega = kc.$$



Complex form,

$$\tilde{f}(x,t) = \tilde{A} \exp(i kx - i\omega t), \ \tilde{A} = Ae^{i\Delta}$$

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▶ 3D Wave Equations for **E**, **B** 

$$\begin{split} \left(\nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2}\right) \mathbf{E} &= 0, \\ \left(\nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2}\right) \mathbf{B} &= 0. \end{split}$$

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c is the speed of light in vacuum, light is an EM wave!



Wave front is planar & perpendicular to direction of propagation

$$\mathbf{E} = \tilde{\mathbf{E}}_0 e^{i(\mathbf{k} \cdot \mathbf{x} - \omega t)}, \mathbf{B} = \tilde{\mathbf{B}}_0 e^{i(\mathbf{k} \cdot \mathbf{x} - \omega t)},$$

**k** is the wave vector, normal to wave fronts.

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▶ Gauss law  $(\nabla \cdot \mathbf{E} = \nabla \cdot \mathbf{B} = 0)$  implies,

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► Faraday law  $(\nabla \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = 0)$  implies,

$$\mathbf{B} = \hat{\mathbf{k}} \times \frac{\mathbf{E}}{\epsilon}$$
.



▶ Using,  $\mathbf{B} = \hat{\mathbf{k}} \times \frac{\mathbf{E}}{c}$ , can show,

$$rac{1}{2\mu_0}\mathbf{B}^2 = rac{\epsilon_0}{2}\mathbf{E}^2, u_\mathbf{E} = u_\mathbf{B},$$
  $u = \epsilon_0\mathbf{E}^2.$ 

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 Energy flux density (per unit cross-sectional area, per unit time), also called intensity

$$\mathbf{S} = \frac{\mathbf{E}^2}{\mu_0 c} \hat{\mathbf{k}} = \epsilon_0 c \; \mathbf{E}^2 \hat{\mathbf{k}}.$$

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ightharpoonup Momentum density,  $\pi_{\it EM}$ 

$$m{\pi}_{\mathit{EM}} = rac{1}{c^2} \mathbf{S}, m{\pi}_{\mathit{EM}} = rac{u}{c}.$$

<sup>&</sup>lt;sup>2</sup>Note that  $\langle \mathbf{E}^2 \rangle = \frac{1}{2} \mathbf{E}_0^2$ , since  $\langle \sin^2 \theta \rangle = \frac{1}{2}$ .

 Absorption of EM waves: Momentum transferred to the absorber,

$$P=rac{ extit{Normal Force}}{ extit{Area}}=rac{ extit{Force}}{ extit{Cross-sectional area}}=rac{ extit{Momentum Transferred}}{ extit{time} imex extit{Cross-sectional area}}$$
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▶ Perfect reflector, twice as much momentum transfer

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► Microscopic level, **E** and **B** push/pull on charges in the wall.

#### EM waves in matter

Maxwell Equations in absence of sources in matter,

$$\mathbf{\nabla} \cdot \mathbf{D} = 0, \qquad \mathbf{\nabla} \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = 0,$$

$$\boldsymbol{\nabla}\cdot\boldsymbol{\mathbf{B}}=0,\qquad \boldsymbol{\nabla}\times\boldsymbol{\mathbf{H}}-\frac{\partial\boldsymbol{\mathbf{D}}}{\partial t}=0$$

Maxwell Equations in absence of sources in matter,

$$\mathbf{\nabla \cdot D} = 0,$$
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Linear homogenous media: Constitutive relations,

$$\mathbf{D}=\epsilon\mathbf{E},\quad \mathbf{B}=\mu\mathbf{H}.$$

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Maxwell's equations for linear homogenous media

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Speed of EM waves,

$$v = \frac{1}{\sqrt{\mu\epsilon}} = \frac{c}{n}$$

n is called the refractive index.



Energy density in media,

$$u = \frac{1}{2}\mathbf{D} \cdot \mathbf{H} + \frac{1}{2}\mathbf{B} \cdot \mathbf{H}$$
$$= \frac{\epsilon}{2}\mathbf{E}^2 + \frac{1}{2\mu}\mathbf{B}^2.$$

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Energy flux rate: Poynting vector,

$$\mathbf{S} = \frac{1}{\mu} \left( \mathbf{E} \times \mathbf{B} \right)$$

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▶ Boundary Conditions at the interface of media:

$$\epsilon^1 E_{\perp}^1 = \epsilon^2 E_{\perp}^2, \qquad \mathbf{E}_{\parallel}^1 = \mathbf{E}_{\parallel}^2,$$

Monochromatic waves incident on the interface from left (fig. 9.14 of Griffiths' text)

$$\mathbf{E}_{I} = \tilde{\mathbf{E}}_{0I} e^{i(\mathbf{k}_{I} \cdot \mathbf{x} - \omega t)}, \ \mathbf{B}_{I} = \hat{\mathbf{k}}_{I} \times \frac{\mathbf{E}_{I}}{v_{1}}.$$

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Reflected wave

$$\mathbf{E}_R = \tilde{\mathbf{E}}_{0R} e^{i(\mathbf{k}_R \cdot \mathbf{x} - \omega t)}, \ \mathbf{B}_R = \hat{\mathbf{k}}_R \times \frac{\mathbf{E}_R}{v_1}.$$

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Transmitted wave

$$\mathbf{E}_T = \tilde{\mathbf{E}}_{0T} e^{i(\mathbf{k}_T \cdot \mathbf{x} - \omega t)}, \ \mathbf{B}_T = \hat{\mathbf{k}}_T \times \frac{\mathbf{E}_T}{v_2}.$$

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► Transmitted wave

$$\mathbf{E}_T = \tilde{\mathbf{E}}_{0T} e^{i(\mathbf{k}_T \cdot \mathbf{x} - \omega t)}, \ \mathbf{B}_T = \hat{\mathbf{k}}_T \times \frac{\mathbf{E}_T}{V_2}.$$

Frequency same,

$$\omega = |\mathbf{k}|v,$$

$$\implies |\mathbf{k}_I| = |\mathbf{k}_R| = \frac{\omega}{v_1}, \qquad |\mathbf{k}_T| = \frac{\omega}{v_2} = \frac{v_1}{v_2} |\mathbf{k}_I| = \frac{n_2}{n_1} |\mathbf{k}_I|.$$

We take the interface of the media to be the xy-plane i.e. z=0. Then,

$$|\mathbf{k} \cdot \mathbf{x}|_{z=0} = (k_x x + k_y y + k_z z)_{z=0} = k_x x + k_y y = \mathbf{k}^{\parallel} \cdot \mathbf{x}$$

where we have introduced  $\mathbf{k}^{\parallel}$  as the part of the wave-vector parallel to the interface

$$\mathbf{k}^{\parallel} = \mathbf{k} - k_z \hat{\mathbf{z}} = k_x \hat{\mathbf{x}} + k_y \hat{\mathbf{y}}.$$

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Boundary conditions for Maxwell's equations at the interface i.e. z = 0 generically look like

$$(\ldots) e^{i(\mathbf{k}_{I} \cdot \mathbf{x} - \omega t)} + (\ldots) e^{i(\mathbf{k}_{R} \cdot \mathbf{x} - \omega t)} \Big|_{z=0} = (\ldots) e^{i(\mathbf{k}_{T} \cdot \mathbf{x} - \omega t)} \Big|_{z=0} \forall x, y, t,$$

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or equivalently

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➤ Since this is a complex equation the phases on both sides must be equal,

$$\Rightarrow \mathbf{k}_{I}^{\parallel} \cdot \mathbf{x} = \mathbf{k}_{R}^{\parallel} \cdot \mathbf{x} = \mathbf{k}_{T}^{\parallel} \cdot \mathbf{x}, \quad \forall \mathbf{x}, \forall \mathbf{y} \mapsto \mathbf{x} \mapsto \mathbf{x}$$

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From last slide

$$\Rightarrow \mathbf{k}_I^{\parallel} = \mathbf{k}_R^{\parallel} = \mathbf{k}_T^{\parallel}$$

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First law: All three wave vectors lie in a plane (plane of incidence containing **k**<sub>I</sub> and **2**).

$$|\mathbf{k}_I|\sin\theta_I = |\mathbf{k}_R|\sin\theta_R = |\mathbf{k}_T|\sin\theta_T.$$

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First law: All three wave vectors lie in a plane (plane of incidence containing  $\mathbf{k}_l$  and  $\hat{\mathbf{z}}$ ).

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Second law: Angle of reflection = Angle of incidence,

$$\theta_I = \theta_R$$
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Second law: Angle of reflection = Angle of incidence,

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► Third law: Snell's law of Refraction

$$\frac{\sin \theta_I}{\sin \theta_T} = \frac{|\mathbf{k}_T|}{|\mathbf{k}_I|} = \frac{v_1}{v_2} = \frac{n_2}{n_1}$$

## Satisfying boundary conditions

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Boundary conditions relates amplitudes

$$\begin{split} \epsilon^1 \left( \tilde{E}_{0I}^3 + \tilde{E}_{0R}^3 \right) &= \epsilon^2 \tilde{E}_{0T}^3, \\ \tilde{B}_{0I}^3 + \tilde{B}_{0R}^3 &= \tilde{B}_{0T}^3, \\ \tilde{E}_{0I}^{1,2} + \tilde{E}_{0R}^{1,2} &= \tilde{E}_{0T}^{1,2}, \\ \frac{1}{\mu_1} \left( \tilde{B}_{0I}^{1,2} + \tilde{B}_{0R}^{1,2} \right) &= \frac{1}{\mu_2} \tilde{B}_{0T}^{1,2}. \end{split}$$

## Satisfying boundary conditions

Boundary conditions relates amplitudes

$$\begin{split} \epsilon^1 \left( \tilde{E}_{0I}^3 + \tilde{E}_{0R}^3 \right) &= \epsilon^2 \tilde{E}_{0T}^3, \\ \tilde{B}_{0I}^3 + \tilde{B}_{0R}^3 &= \tilde{B}_{0T}^3, \\ \tilde{E}_{0I}^{1,2} + \tilde{E}_{0R}^{1,2} &= \tilde{E}_{0T}^{1,2}, \\ \frac{1}{\mu_1} \left( \tilde{B}_{0I}^{1,2} + \tilde{B}_{0R}^{1,2} \right) &= \frac{1}{\mu_2} \tilde{B}_{0T}^{1,2}. \end{split}$$

► Case I: "In plane" polarization,

$$\begin{split} \epsilon^{1} \left( - \left| \tilde{\mathbf{E}}_{0I} \right| \sin \theta_{I} + \left| \tilde{\mathbf{E}}_{0R} \right| \sin \theta_{R} \right) &= -\epsilon^{2} \left| \tilde{\mathbf{E}}_{0T} \right| \sin \theta_{T} \\ \left| \tilde{\mathbf{E}}_{0I} \right| \cos \theta_{I} + \left| \tilde{\mathbf{E}}_{0R} \right| \cos \theta_{R} &= \left| \tilde{\mathbf{E}}_{0T} \right| \cos \theta_{T} \\ &\frac{1}{\mu_{1} \nu_{1}} \left( \left| \tilde{\mathbf{E}}_{0I} \right| - \left| \tilde{\mathbf{E}}_{0R} \right| \right) &= \frac{1}{\mu_{2} \nu_{2}} \left| \tilde{\mathbf{E}}_{0T} \right| \end{split}$$

► We have,

$$\begin{split} \left| \tilde{\mathbf{E}}_{0I} \right| - \left| \tilde{\mathbf{E}}_{0R} \right| &= \beta \left| \tilde{\mathbf{E}}_{0T} \right|, \beta = \frac{\mu_1 v_1}{\mu_2 v_2} = \frac{\mu_1 n_2}{\mu_2 n_1} \\ \left| \tilde{\mathbf{E}}_{0I} \right| + \left| \tilde{\mathbf{E}}_{0R} \right| &= \alpha \left| \tilde{\mathbf{E}}_{0T} \right|, \alpha = \frac{\cos \theta_T}{\cos \theta_I} \end{split}$$

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Fresnel's Equations for "in plane" component of polarization

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Brewster Angle: Full transmission

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Next class: EM waves incident on conductors



1. Point out where rationalized units are not used

$$(a)\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{|\mathbf{x}|^3} \mathbf{x}, \quad (b)\mathbf{B} = \frac{\mu_0}{4\pi} \oint \frac{ld\mathbf{x}' \times (\mathbf{x} - \mathbf{x}')}{|\mathbf{x} - \mathbf{x}'|^3}, \quad (c)\mathbf{\nabla} \cdot \mathbf{E} = 4\pi\rho, \quad (d)\mathbf{\nabla} \times \mathbf{B} = \frac{4\pi}{6} \mathbf{j}$$

2. Direction of Eof point charge at origin in spherical polar coordinates is along

$$(a)\hat{\mathbf{r}}, \qquad (b)\hat{\boldsymbol{\theta}}, \qquad (c)\hat{\boldsymbol{\phi}}, \qquad (d)\hat{\boldsymbol{\rho}}$$

Direction of B of infinitely long current carrying wire along z-axis in cylindrical coordinates is along

$$(a)\hat{\boldsymbol{\rho}}, \qquad (b)\hat{\boldsymbol{\phi}}, \qquad (c)\hat{\boldsymbol{z}}, \qquad (d) \quad \rho\hat{\boldsymbol{\rho}} + z\hat{\boldsymbol{z}}$$

4. Magnetic field B(x) due to a point charge, q at x' moving with vel, v

(a)absent, (b)approx 
$$\propto q \frac{\mathbf{v}}{c} \times \frac{\mathbf{x}}{|\mathbf{x} - \mathbf{v}t|^3}$$
, (c)approx  $\propto q \frac{\mathbf{v}}{c} \times \frac{\mathbf{x} - \mathbf{v}t}{|\mathbf{x} - \mathbf{v}t|^3}$ , (d)approx  $\propto q \frac{\mathbf{v}}{c} \times \mathbf{E}$ 

5. Inside a conductor, in general

$$(a)\mathbf{E} = 0$$
,  $(b)\rho = 0$ ,  $(c)\mathbf{i} = \sigma\mathbf{E}$   $(d)\mathbf{E} = \sigma\mathbf{i}$ 



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$$\mathbf{E}_{\parallel}$$
, (b) $\mathbf{H}_{\parallel}$  (c) $D_{\perp}$  (d) $B_{\perp}$ 

7. If **E**, **B** and  $\phi$ , **A** contain same amount of information, why are they different in number of components?

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