

PH 1027: Maxwell's Equations and Electromagnetic Waves

Instructor: Shubho Roy¹
(Dept. of Physics)

Lecture 5

April 3, 2019

¹Office: C 313 D, Office hrs: Walk in or Email Appointment
Email: sroy@iith.ac.in

Agenda

- ▶ Electrostatics: Coulomb force law, System of units, Nature of electrical charge

Agenda

- ▶ Electrostatics: Coulomb force law, System of units, Nature of electrical charge
- ▶ Electric field strength vector, Gauss law, Conservative nature of Electric field strength: Scalar potential

Agenda

- ▶ Electrostatics: Coulomb force law, System of units, Nature of electrical charge
- ▶ Electric field strength vector, Gauss law, Conservative nature of Electric field strength: Scalar potential
- ▶ Poisson Equation, Laplace Equation, Boundary value problem in Electrostatics: Boundary conditions

Agenda

- ▶ Electrostatics: Coulomb force law, System of units, Nature of electrical charge
- ▶ Electric field strength vector, Gauss law, Conservative nature of Electric field strength: Scalar potential
- ▶ Poisson Equation, Laplace Equation, Boundary value problem in Electrostatics: Boundary conditions
- ▶ Earnshaw, Mean-value and Uniqueness theorems.

Agenda

- ▶ Electrostatics: Coulomb force law, System of units, Nature of electrical charge
- ▶ Electric field strength vector, Gauss law, Conservative nature of Electric field strength: Scalar potential
- ▶ Poisson Equation, Laplace Equation, Boundary value problem in Electrostatics: Boundary conditions
- ▶ Earnshaw, Mean-value and Uniqueness theorems.
- ▶ Electric Dipole

Agenda

- ▶ Electrostatics: Coulomb force law, System of units, Nature of electrical charge
- ▶ Electric field strength vector, Gauss law, Conservative nature of Electric field strength: Scalar potential
- ▶ Poisson Equation, Laplace Equation, Boundary value problem in Electrostatics: Boundary conditions
- ▶ Earnshaw, Mean-value and Uniqueness theorems.
- ▶ Electric Dipole
- ▶ Electric fields in matter (Dielectrics): Induced Polarization, Electric Displacement vector, Linear Dielectrics, Boundary Value problems and boundary conditions

Agenda

- ▶ Electrostatics: Coulomb force law, System of units, Nature of electrical charge
- ▶ Electric field strength vector, Gauss law, Conservative nature of Electric field strength: Scalar potential
- ▶ Poisson Equation, Laplace Equation, Boundary value problem in Electrostatics: Boundary conditions
- ▶ Earnshaw, Mean-value and Uniqueness theorems.
- ▶ Electric Dipole
- ▶ Electric fields in matter (Dielectrics): Induced Polarization, Electric Displacement vector, Linear Dielectrics, Boundary Value problems and boundary conditions
- ▶ Energy in electric fields in vacuum and matter

Agenda

- ▶ Electrostatics: Coulomb force law, System of units, Nature of electrical charge
- ▶ Electric field strength vector, Gauss law, Conservative nature of Electric field strength: Scalar potential
- ▶ Poisson Equation, Laplace Equation, Boundary value problem in Electrostatics: Boundary conditions
- ▶ Earnshaw, Mean-value and Uniqueness theorems.
- ▶ Electric Dipole
- ▶ Electric fields in matter (Dielectrics): Induced Polarization, Electric Displacement vector, Linear Dielectrics, Boundary Value problems and boundary conditions
- ▶ Energy in electric fields in vacuum and matter
- ▶ Conductors

References/Readings

References/Readings

- ▶ Griffiths, D.J., **Introduction to Electrodynamics, Ch. 2, 4**

Electrostatics: Coulomb law

Electrostatics: Coulomb law

- ▶ **Observation 1:** “Electrons and Protons” exert force on each other due to some property “Electric charge”

Electrostatics: Coulomb law

- ▶ **Observation 1:** “Electrons and Protons” exert force on each other due to some property “Electric charge”
- ▶ **Observation 2:** Charges can be two types/signs (determining the direction of force)

Electrostatics: Coulomb law

- ▶ **Observation 1:** “Electrons and Protons” exert force on each other due to some property “Electric charge”
- ▶ **Observation 2:** Charges can be two types/signs (determining the direction of force)
- ▶ **Observation 3:** This force falls off “inverse square” of separation distance and is “central”.

Electrostatics: Coulomb law

- ▶ **Observation 1:** “Electrons and Protons” exert force on each other due to some property “Electric charge”
- ▶ **Observation 2:** Charges can be two types/signs (determining the direction of force)
- ▶ **Observation 3:** This force falls off “inverse square” of separation distance and is “central”.
- ▶ **Combine all three:** Coulomb law

$$\mathbf{F}_1 = \lambda_e \frac{Q_1 Q_2}{|\mathbf{x}_1 - \mathbf{x}_2|^3} (\mathbf{x}_1 - \mathbf{x}_2),$$

$$\mathbf{F}_2 = -\mathbf{F}_1$$

Electrostatics: Coulomb law

- ▶ **Observation 1:** “Electrons and Protons” exert force on each other due to some property “Electric charge”
- ▶ **Observation 2:** Charges can be two types/signs (determining the direction of force)
- ▶ **Observation 3:** This force falls off “inverse square” of separation distance and is “central”.
- ▶ **Combine all three:** Coulomb law

$$\mathbf{F}_1 = \lambda_e \frac{Q_1 Q_2}{|\mathbf{x}_1 - \mathbf{x}_2|^3} (\mathbf{x}_1 - \mathbf{x}_2),$$

$$\mathbf{F}_2 = -\mathbf{F}_1$$

- ▶ λ_e depends on the system of units. In Gaussian units (CGS), its just 1. In MKS, $\lambda = 9 \times 10^9 Nm^2/C^2$.

Electrostatics: Electric Field Strength

Electrostatics: Electric Field Strength

- ▶ **Conservation:** Charges before and after some process remain intact \implies Continuity equation

Electrostatics: Electric Field Strength

- ▶ **Conservation:** Charges before and after some process remain intact \implies Continuity equation
- ▶ **Discreteness:** Charges appear only in multiples of the electronic charge,

$$Q = N e, \quad e = 1.6 \times 10^{-19} C$$

Electrostatics: Electric Field Strength

- ▶ **Conservation:** Charges before and after some process remain intact \implies Continuity equation
- ▶ **Discreteness:** Charges appear only in multiples of the electronic charge,

$$Q = N e, \quad e = 1.6 \times 10^{-19} \text{C}$$

- ▶ **Superposition:** Net force on a charge, Q due to a distribution of other charges, Q_i 's, given by the **vector sum**

$$\mathbf{F}(\mathbf{x}) = \sum_{i=1}^N \lambda_e \frac{Q Q_i}{|\mathbf{x} - \mathbf{x}'_i|^3} (\mathbf{x} - \mathbf{x}'_i)$$

Electrostatics: Electric Field Strength

- ▶ **Conservation:** Charges before and after some process remain intact \implies Continuity equation
- ▶ **Discreteness:** Charges appear only in multiples of the electronic charge,

$$Q = N e, \quad e = 1.6 \times 10^{-19} \text{C}$$

- ▶ **Superposition:** Net force on a charge, Q due to a distribution of other charges, Q_i 's, given by the **vector sum**

$$\mathbf{F}(\mathbf{x}) = \sum_{i=1}^N \lambda_e \frac{Q Q_i}{|\mathbf{x} - \mathbf{x}'_i|^3} (\mathbf{x} - \mathbf{x}'_i)$$

- ▶ **Electric field Strength, $\mathbf{E}(\mathbf{x})$**

$$\mathbf{E}(\mathbf{x}) = \frac{\mathbf{F}(\mathbf{x})}{Q} = \sum_{i=1}^N \lambda_e \frac{Q_i}{|\mathbf{x} - \mathbf{x}'_i|^3} (\mathbf{x} - \mathbf{x}'_i)$$

Electrostatics: Electric Field Strength

- ▶ **Conservation:** Charges before and after some process remain intact \implies Continuity equation
- ▶ **Discreteness:** Charges appear only in multiples of the electronic charge,

$$Q = N e, \quad e = 1.6 \times 10^{-19} \text{C}$$

- ▶ **Superposition:** Net force on a charge, Q due to a distribution of other charges, Q_i 's, given by the **vector sum**

$$\mathbf{F}(\mathbf{x}) = \sum_{i=1}^N \lambda_e \frac{Q Q_i}{|\mathbf{x} - \mathbf{x}'_i|^3} (\mathbf{x} - \mathbf{x}'_i)$$

- ▶ **Electric field Strength, $\mathbf{E}(\mathbf{x})$**

$$\mathbf{E}(\mathbf{x}) = \frac{\mathbf{F}(\mathbf{x})}{Q} = \sum_{i=1}^N \lambda_e \frac{Q_i}{|\mathbf{x} - \mathbf{x}'_i|^3} (\mathbf{x} - \mathbf{x}'_i)$$

- ▶ Region containing charges. Electric field (plot)

Electrostatics: Gauss law

Electrostatics: Gauss law

- ▶ Coulomb Force on a static point charge, q in an electric field,

$$\mathbf{F} = q\mathbf{E}$$

Electrostatics: Gauss law

- ▶ Coulomb Force on a static point charge, q in an electric field,

$$\mathbf{F} = q\mathbf{E}$$

- ▶ **Continuous distribution:**

$$\mathbf{E}(\mathbf{x}) = \lambda_e \iiint \frac{\rho(\mathbf{x}') d^3\mathbf{x}'}{|\mathbf{x} - \mathbf{x}'|^3} (\mathbf{x} - \mathbf{x}')$$

Electrostatics: Gauss law

- ▶ Coulomb Force on a static point charge, q in an electric field,

$$\mathbf{F} = q\mathbf{E}$$

- ▶ **Continuous distribution:**

$$\mathbf{E}(\mathbf{x}) = \lambda_e \iiint \frac{\rho(\mathbf{x}') d^3\mathbf{x}'}{|\mathbf{x} - \mathbf{x}'|^3} (\mathbf{x} - \mathbf{x}')$$

- ▶ **Gauss Law:** Take divergence of both sides,

$$\nabla \cdot \mathbf{E} = 4\pi\lambda_e \rho(\mathbf{x}).$$

Electrostatics: Gauss law

- ▶ Coulomb Force on a static point charge, q in an electric field,

$$\mathbf{F} = q\mathbf{E}$$

- ▶ **Continuous distribution:**

$$\mathbf{E}(\mathbf{x}) = \lambda_e \iiint \frac{\rho(\mathbf{x}') d^3\mathbf{x}'}{|\mathbf{x} - \mathbf{x}'|^3} (\mathbf{x} - \mathbf{x}')$$

- ▶ **Gauss Law:** Take divergence of both sides,

$$\nabla \cdot \mathbf{E} = 4\pi\lambda_e \rho(\mathbf{x}).$$

- ▶ **Permittivity of free space (ϵ_0):** Rename, $4\pi\lambda_e = \frac{1}{\epsilon_0}$,

$$\nabla \cdot \mathbf{E}(\mathbf{x}) = \frac{\rho(\mathbf{x})}{\epsilon_0}.$$

The first Maxwell's equation.

Electrostatics: Scalar Electrostatic potential

Electrostatics: Scalar Electrostatic potential

- ▶ **Electric field is conservative:**

$$\nabla \times \mathbf{E}(\mathbf{x}) = 0.$$

Electrostatics: Scalar Electrostatic potential

- ▶ **Electric field is conservative:**

$$\nabla \times \mathbf{E}(\mathbf{x}) = 0.$$

- ▶ **Use Stokes theorem:**

$$\oint_C \mathbf{E} \cdot d\mathbf{l} = 0.$$

Electrostatics: Scalar Electrostatic potential

- ▶ **Electric field is conservative:**

$$\nabla \times \mathbf{E}(\mathbf{x}) = 0.$$

- ▶ **Use Stokes theorem:**

$$\oint_C \mathbf{E} \cdot d\mathbf{l} = 0.$$

- ▶ **Electrostatic scalar potential: $\Phi(\mathbf{x})$**

$$\mathbf{E}(\mathbf{x}) = -\nabla\Phi(\mathbf{x})$$

Electrostatics: Scalar Electrostatic potential

- ▶ **Electric field is conservative:**

$$\nabla \times \mathbf{E}(\mathbf{x}) = 0.$$

- ▶ **Use Stokes theorem:**

$$\oint_C \mathbf{E} \cdot d\mathbf{l} = 0.$$

- ▶ **Electrostatic scalar potential: $\Phi(\mathbf{x})$**

$$\mathbf{E}(\mathbf{x}) = -\nabla\Phi(\mathbf{x})$$

- ▶ **Solution:**

$$\Phi(\mathbf{x}) = - \int_{\mathbf{x}_0}^{\mathbf{x}} \mathbf{E} \cdot d\mathbf{x}, \quad \Phi(\mathbf{x}_0) = 0.$$

$$\Phi(\mathbf{x}) = \sum_{i=1}^N \frac{1}{4\pi\epsilon_0} \frac{Q_i}{|\mathbf{x} - \mathbf{x}'_i|} = \frac{1}{4\pi\epsilon_0} \iiint \frac{\rho(\mathbf{x}') d^3\mathbf{x}'}{|\mathbf{x} - \mathbf{x}'|}.$$

Electrostatics: Poisson and Laplace Equation

Electrostatics: Poisson and Laplace Equation

- ▶ **Gauss law in term of potential: Poisson Equation**

$$\nabla^2 \Phi(\mathbf{x}) = -\frac{\rho(\mathbf{x})}{\epsilon_0}.$$

Electrostatics: Poisson and Laplace Equation

- ▶ **Gauss law in term of potential: Poisson Equation**

$$\nabla^2 \Phi(\mathbf{x}) = -\frac{\rho(\mathbf{x})}{\epsilon_0}.$$

- ▶ **In regions outside sources: Laplace Equation**

$$\nabla^2 \Phi(\mathbf{x}) = 0.$$

Electrostatics: Poisson and Laplace Equation

- ▶ **Gauss law in term of potential: Poisson Equation**

$$\nabla^2 \Phi(\mathbf{x}) = -\frac{\rho(\mathbf{x})}{\epsilon_0}.$$

- ▶ **In regions outside sources: Laplace Equation**

$$\nabla^2 \Phi(\mathbf{x}) = 0.$$

- ▶ **Fact:** The electrostatic potential cannot have an extremum in a charge-free region.

Electrostatics: Poisson and Laplace Equation

- ▶ **Gauss law in term of potential: Poisson Equation**

$$\nabla^2 \Phi(\mathbf{x}) = -\frac{\rho(\mathbf{x})}{\epsilon_0}.$$

- ▶ **In regions outside sources: Laplace Equation**

$$\nabla^2 \Phi(\mathbf{x}) = 0.$$

- ▶ **Fact:** The electrostatic potential cannot have an extremum in a charge-free region.
- ▶ **Earnshaw's Theorem:** System of point-like charges *cannot* be in stable, static equilibrium under the influence of purely electric forces.

Electrostatics: Mean-Value and Uniqueness Theorem

Electrostatics: Mean-Value and Uniqueness Theorem

- **Mean-Value Theorem:** If S is the surface of a sphere whose interior contains no charge, then the potential at the center is equal to the average potential over the surface S .

$$\Phi(\mathbf{0}) = \langle \Phi(\mathbf{x}) \rangle_S,$$
$$\langle \Phi(\mathbf{x}) \rangle_S = \frac{\oint_S dS \Phi(\mathbf{x})}{4\pi R^2}$$

Electrostatics: Mean-Value and Uniqueness Theorem

- ▶ **Mean-Value Theorem:** If S is the surface of a sphere whose interior contains no charge, then the potential at the center is equal to the average potential over the surface S .

$$\Phi(\mathbf{0}) = \langle \Phi(\mathbf{x}) \rangle_S,$$
$$\langle \Phi(\mathbf{x}) \rangle_S = \frac{\oint_S dS \Phi(\mathbf{x})}{4\pi R^2}$$

- ▶ **Boundary conditions (BC's) for Poisson/Laplace :**

Dirichlet : Specify $\Phi(\mathbf{x})$ on the boundary

Neumann : Specify $\hat{\mathbf{n}} \cdot \mathbf{E}(\mathbf{x}) = \frac{\partial \Phi}{\partial n}$ on the boundary

Electrostatics: Mean-Value and Uniqueness Theorem

- ▶ **Mean-Value Theorem:** If S is the surface of a sphere whose interior contains no charge, then the potential at the center is equal to the average potential over the surface S .

$$\Phi(\mathbf{0}) = \langle \Phi(\mathbf{x}) \rangle_S,$$
$$\langle \Phi(\mathbf{x}) \rangle_S = \frac{\oint_S dS \Phi(\mathbf{x})}{4\pi R^2}$$

- ▶ **Boundary conditions (BC's) for Poisson/Laplace :**

Dirichlet : Specify $\Phi(\mathbf{x})$ on the boundary

Neumann : Specify $\hat{\mathbf{n}} \cdot \mathbf{E}(\mathbf{x}) = \frac{\partial \Phi}{\partial n}$ on the boundary

- ▶ **Uniqueness Theorem:** If $\Phi_1(\mathbf{x})$ and $\Phi_2(\mathbf{x})$ are two soln.s of Poisson Eq. with the *same* charge density and the *same* BC's, then they differ *at most* by an additive constant.

$$\Phi_1(\mathbf{x}) = \Phi_2(\mathbf{x}) + \Phi_0.$$

Field and Potential of an electric dipole

Field and Potential of an electric dipole

- ▶ **Electric dipole:** Two charges, $+Q$ and $-Q$, placed close by at \mathbf{x}_+ and \mathbf{x}_-

Field and Potential of an electric dipole

- ▶ **Electric dipole:** Two charges, $+Q$ and $-Q$, placed close by at \mathbf{x}_+ and \mathbf{x}_-
- ▶ **Ideal/Point dipole:** Take the limit, $|\mathbf{x} - \mathbf{x}_-| \rightarrow 0$ and $Q \rightarrow \infty$ so that their product is finite,

$$Q |\mathbf{x}_+ - \mathbf{x}_-| \rightarrow p$$

- ▶ Define the electric dipole moment,

$$\mathbf{p} = Q (\mathbf{x}_+ - \mathbf{x}_-)$$

Field and Potential of an electric dipole

- ▶ **Electric dipole:** Two charges, $+Q$ and $-Q$, placed close by at \mathbf{x}_+ and \mathbf{x}_-
- ▶ **Ideal/Point dipole:** Take the limit, $|\mathbf{x} - \mathbf{x}_-| \rightarrow 0$ and $Q \rightarrow \infty$ so that their product is finite,

$$Q |\mathbf{x}_+ - \mathbf{x}_-| \rightarrow p$$

- ▶ Define the electric dipole moment,

$$\mathbf{p} = Q (\mathbf{x}_+ - \mathbf{x}_-)$$

- ▶ Scalar potential at a point, \mathbf{x} due to an ideal dipole, located at \mathbf{x}' ,

$$\begin{aligned}\Phi(\mathbf{x}) &= \lim_{Q \rightarrow \infty; \mathbf{x}_+ \rightarrow \mathbf{x}_-} \frac{Q}{4\pi\epsilon_0} \left(\frac{1}{|\mathbf{x} - \mathbf{x}_+|} - \frac{1}{|\mathbf{x} - \mathbf{x}_-|} \right) \\ &= \frac{1}{4\pi\epsilon_0} \frac{\mathbf{p} \cdot (\mathbf{x} - \mathbf{x}')}{|\mathbf{x} - \mathbf{x}'|^3}, \quad \mathbf{x}' = \frac{\mathbf{x}_+ + \mathbf{x}_-}{2} \\ &= -\frac{1}{4\pi\epsilon_0} \mathbf{p}|_{\mathbf{x}'} \cdot \nabla \frac{1}{|\mathbf{x} - \mathbf{x}'|} = \frac{1}{4\pi\epsilon_0} \mathbf{p}|_{\mathbf{x}'} \cdot \nabla' \frac{1}{|\mathbf{x} - \mathbf{x}'|}.\end{aligned}$$

Field and Potential of an electric dipole

- ▶ **Electric dipole:** Two charges, $+Q$ and $-Q$, placed close by at \mathbf{x}_+ and \mathbf{x}_-
- ▶ **Ideal/Point dipole:** Take the limit, $|\mathbf{x} - \mathbf{x}_-| \rightarrow 0$ and $Q \rightarrow \infty$ so that their product is finite,

$$Q |\mathbf{x}_+ - \mathbf{x}_-| \rightarrow p$$

- ▶ Define the electric dipole moment,

$$\mathbf{p} = Q (\mathbf{x}_+ - \mathbf{x}_-)$$

- ▶ Scalar potential at a point, \mathbf{x} due to an ideal dipole, located at \mathbf{x}' ,

$$\begin{aligned}\Phi(\mathbf{x}) &= \lim_{Q \rightarrow \infty; \mathbf{x}_+ \rightarrow \mathbf{x}_-} \frac{Q}{4\pi\epsilon_0} \left(\frac{1}{|\mathbf{x} - \mathbf{x}_+|} - \frac{1}{|\mathbf{x} - \mathbf{x}_-|} \right) \\ &= \frac{1}{4\pi\epsilon_0} \frac{\mathbf{p} \cdot (\mathbf{x} - \mathbf{x}')}{|\mathbf{x} - \mathbf{x}'|^3}, \quad \mathbf{x}' = \frac{\mathbf{x}_+ + \mathbf{x}_-}{2} \\ &= -\frac{1}{4\pi\epsilon_0} \mathbf{p}|_{\mathbf{x}'} \cdot \nabla \frac{1}{|\mathbf{x} - \mathbf{x}'|} = \frac{1}{4\pi\epsilon_0} \mathbf{p}|_{\mathbf{x}'} \cdot \nabla' \frac{1}{|\mathbf{x} - \mathbf{x}'|}.\end{aligned}$$

- ▶ Electric field at \mathbf{x} ,

$$\mathbf{E}(\mathbf{x}) = -\nabla \Phi(\mathbf{x}) = ?$$

Electric fields in matter: Induced Polarization

Electric fields in matter: Induced Polarization

- ▶ **Conductors (free charges) or Dielectrics (bound charges)**

Electric fields in matter: Induced Polarization

- ▶ **Conductors (free charges) or Dielectrics (bound charges)**
- ▶ Effect of Electric fields on Dielectrics: **Induces Polarization (molecules acquire a dipole moment)**

$$\mathbf{P} = \mathbf{P}(\mathbf{E})$$

Electric fields in matter: Induced Polarization

- ▶ **Conductors (free charges) or Dielectrics (bound charges)**
- ▶ Effect of Electric fields on Dielectrics: **Induces Polarization (molecules acquire a dipole moment)**

$$\mathbf{P} = \mathbf{P}(\mathbf{E})$$

- ▶ **Potential** produced by dipoles

$$\Phi_{\mathbf{P}}(\mathbf{x}) = \frac{1}{4\pi\epsilon_0} \int d^3\mathbf{x}' \mathbf{P}(\mathbf{x}') \cdot \nabla' \frac{1}{|\mathbf{x} - \mathbf{x}'|}$$

- ▶ **Induced volume and surface charge:**

$$\rho_{bound} = -\nabla \cdot \mathbf{P}, \sigma_{bound} = \mathbf{P} \cdot \mathbf{n}$$

Electric fields in matter: Electric displacement field

Electric fields in matter: Electric displacement field

- Gauss law gets modified:

$$\begin{aligned}\nabla \cdot \mathbf{E} &= \frac{\rho_{\text{free}} + \rho_{\text{bound}}}{\epsilon_0} \\ \nabla \cdot \underbrace{\mathbf{D}}_{=\epsilon_0 \mathbf{E} + \mathbf{P}} &= \rho_{\text{free}}\end{aligned}$$

Electric fields in matter: Electric displacement field

- ▶ Gauss law gets modified:

$$\begin{aligned}\nabla \cdot \mathbf{E} &= \frac{\rho_{\text{free}} + \rho_{\text{bound}}}{\epsilon_0} \\ \nabla \cdot \underbrace{\mathbf{D}}_{=\epsilon_0 \mathbf{E} + \mathbf{P}} &= \rho_{\text{free}}\end{aligned}$$

- ▶ Linear dielectrics

$$\mathbf{P} = \epsilon_0 \boldsymbol{\chi} \cdot \mathbf{E}, \quad P_i = \epsilon_0 \chi_{ij} E_j,$$

$\boldsymbol{\chi}$ is the **Electric Susceptibility tensor**.

$$\mathbf{D} = \boldsymbol{\epsilon} \cdot \mathbf{E}, \quad \boldsymbol{\epsilon} = \epsilon_0 (\mathcal{I} + \boldsymbol{\chi}).$$

Electric fields in matter: Electric displacement field

- ▶ Gauss law gets modified:

$$\begin{aligned}\nabla \cdot \mathbf{E} &= \frac{\rho_{free} + \rho_{bound}}{\epsilon_0} \\ \nabla \cdot \underbrace{\mathbf{D}}_{=\epsilon_0 \mathbf{E} + \mathbf{P}} &= \rho_{free}\end{aligned}$$

- ▶ Linear dielectrics

$$\mathbf{P} = \epsilon_0 \boldsymbol{\chi} \cdot \mathbf{E}, \quad P_i = \epsilon_0 \chi_{ij} E_j,$$

$\boldsymbol{\chi}$ is the **Electric Susceptibility tensor**.

$$\mathbf{D} = \boldsymbol{\epsilon} \cdot \mathbf{E}, \quad \boldsymbol{\epsilon} = \epsilon_0 (\mathcal{I} + \boldsymbol{\chi}).$$

- ▶ Boundary value problem **inside Dielectrics**:

$$\epsilon_{ij} \partial_i \partial_j \Phi = -\rho_{free}.$$

- ▶ **Linear Isotropic dielectric**: $\chi_{ij} = \chi \delta_{ij}$

$$\nabla^2 \Phi = -\epsilon^{-1} \rho_{free}, \quad \epsilon = \epsilon_0 (1 + \chi)$$

Boundary Conditions at interface

Boundary Conditions at interface

- ▶ **Gauss law over a pillbox:**

$$D_{\perp}^{over} - D_{\perp}^{under} = \sigma_{free}$$

Boundary Conditions at interface

- ▶ **Gauss law over a pillbox:**

$$D_{\perp}^{over} - D_{\perp}^{uder} = \sigma_{free}$$

- ▶ **Stokes theorem over a rectangle**

$$E_{\parallel}^{over} = E_{\parallel}^{uder}$$

Boundary Conditions at interface

- ▶ **Gauss law over a pillbox:**

$$D_{\perp}^{over} - D_{\perp}^{under} = \sigma_{free}$$

- ▶ **Stokes theorem over a rectangle**

$$E_{\parallel}^{over} = E_{\parallel}^{under}$$

- ▶ **In terms of potential:**

$$-\epsilon^{over} \partial_n \Phi^{over} + \epsilon^{under} \partial_n \Phi^{under} = \sigma_f,$$

$$\partial_t \Phi^{over} = \partial_t \Phi^{under}.$$

Electrostatics: Energy in the Electric field

Electrostatics: Energy in the Electric field

- ▶ **Electrostatic self energy**

$$U = \frac{1}{2} \int d^3\mathbf{x} \rho(\mathbf{x}) \Phi(\mathbf{x})$$

Electrostatics: Energy in the Electric field

- ▶ **Electrostatic self energy**

$$U = \frac{1}{2} \int d^3\mathbf{x} \rho(\mathbf{x}) \Phi(\mathbf{x})$$

- ▶ **Energy in the field**

$$U_{EF} = \frac{\epsilon_0}{2} \int d^3\mathbf{x} \mathbf{E} \cdot \mathbf{E},$$
$$u = \frac{\epsilon_0}{2} \mathbf{E}^2$$

Electrostatics: Energy in the Electric field

- ▶ **Electrostatic self energy**

$$U = \frac{1}{2} \int d^3\mathbf{x} \rho(\mathbf{x}) \Phi(\mathbf{x})$$

- ▶ **Energy in the field**

$$U_{EF} = \frac{\epsilon_0}{2} \int d^3\mathbf{x} \mathbf{E} \cdot \mathbf{E},$$
$$u = \frac{\epsilon_0}{2} \mathbf{E}^2$$

- ▶ **Inside Dielectrics:**

$$U_{Dielectric} = \frac{1}{2} \int d^3\mathbf{x} \mathbf{D}(\mathbf{x}) \cdot \mathbf{E}(\mathbf{x}),$$
$$u_{Dielectrics} = \frac{1}{2} \mathbf{D} \cdot \mathbf{E}.$$

Conductors

Conductors

- ▶ **Inside conductors (due infinite number of free charges formed by dissociating molecules)**

$$\mathbf{E} = 0.$$

Conductors

- ▶ Inside conductors (due infinite number of free charges formed by dissociating molecules)

$$\mathbf{E} = 0.$$

- ▶ Entire conductor is an equipotential:

$$\Phi_2 - \Phi_1 = - \int \mathbf{E} \cdot d\mathbf{l} = 0$$

Conductors

- ▶ Inside conductors (due infinite number of free charges formed by dissociating molecules)

$$\mathbf{E} = 0.$$

- ▶ Entire conductor is an equipotential:

$$\Phi_2 - \Phi_1 = - \int \mathbf{E} \cdot d\mathbf{l} = 0$$

- ▶ Electric field is normal at the boundary,

$$\mathbf{E}^{out} = \frac{\sigma}{\epsilon_0} \hat{\mathbf{n}}.$$

- ▶ Capacitance,

$$\Delta Q = \mathbf{C} \Delta V$$

C measures the quantity of charge need to increase the potential by 1V. The better the conductor, the higher this capacity is.

Maxwell's equations for Electric fields in vacuum and matter

Maxwell's equations for Electric fields in vacuum and matter

- In vacuum,

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}, \quad \nabla \times \mathbf{E} = 0.$$

Maxwell's equations for Electric fields in vacuum and matter

- ▶ **In vacuum,**

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}, \quad \nabla \times \mathbf{E} = 0.$$

- ▶ **in terms of scalar potential**

$$\nabla^2 \Phi = -\frac{\rho}{\epsilon_0}.$$

Maxwell's equations for Electric fields in vacuum and matter

- ▶ In vacuum,

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}, \quad \nabla \times \mathbf{E} = 0.$$

- ▶ in terms of scalar potential

$$\nabla^2 \Phi = -\frac{\rho}{\epsilon_0}.$$

- ▶ In linear dielectrics

$$\nabla \cdot \mathbf{D} = \rho_{free}, \quad \nabla \times \mathbf{E} = 0,$$

$$\mathbf{D} = \epsilon \cdot \mathbf{E}.$$

Maxwell's equations for Electric fields in vacuum and matter

- ▶ In vacuum,

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}, \quad \nabla \times \mathbf{E} = 0.$$

- ▶ in terms of scalar potential

$$\nabla^2 \Phi = -\frac{\rho}{\epsilon_0}.$$

- ▶ In linear dielectrics

$$\nabla \cdot \mathbf{D} = \rho_{free}, \quad \nabla \times \mathbf{E} = 0,$$

$$\mathbf{D} = \epsilon \cdot \mathbf{E}.$$

- ▶ in terms of scalar potential (linear isotropic)

$$\nabla^2 \Phi = -\epsilon^{-1} \rho.$$

Magnetic fields: Produced by charges in motion
and affects charges in motion

Magnetic fields: Produced by charges in motion and affects charges in motion

- Consider a point charge, q moving in the Lab frame at a velocity, \mathbf{v} . What fields does it produce?

Magnetic fields: Produced by charges in motion and affects charges in motion

- ▶ Consider a point charge, q moving in the Lab frame at a velocity, \mathbf{v} . What fields does it produce?
- ▶ In the frame attached to the charge aka the REST FRAME: Only observe an electric field,

$$\mathbf{E}'(\mathbf{x}') = \frac{1}{4\pi\epsilon_0} \frac{q}{|\mathbf{x}'|^3} \mathbf{x}'$$

Magnetic fields: Produced by charges in motion and affects charges in motion

- ▶ Consider a point charge, q moving in the Lab frame at a velocity, \mathbf{v} . What fields does it produce?
- ▶ In the frame attached to the charge aka the REST FRAME: Only observe an electric field,

$$\mathbf{E}'(\mathbf{x}') = \frac{1}{4\pi\epsilon_0} \frac{q}{|\mathbf{x}'|^3} \mathbf{x}'$$

- ▶ But in the Lab frame: Observe both Electric *and* Magnetic field!,

$$\mathbf{E}(\mathbf{x}) = \frac{\gamma}{4\pi\epsilon_0} \frac{q}{|\mathbf{x}'|^3} \mathbf{x},$$

$$\mathbf{B}(\mathbf{x}) = \gamma \frac{\mathbf{v}}{c^2} \times \mathbf{E}'(\mathbf{x}'), \quad \gamma = \sqrt{1 - \mathbf{v}^2/c^2}$$

Magnetic fields: Produced by charges in motion and affects charges in motion

- ▶ Consider a point charge, q moving in the Lab frame at a velocity, \mathbf{v} . What fields does it produce?
- ▶ In the frame attached to the charge aka the REST FRAME: Only observe an electric field,

$$\mathbf{E}'(\mathbf{x}') = \frac{1}{4\pi\epsilon_0} \frac{q}{|\mathbf{x}'|^3} \mathbf{x}'$$

- ▶ But in the Lab frame: Observe both Electric *and* Magnetic field!,

$$\mathbf{E}(\mathbf{x}) = \frac{\gamma}{4\pi\epsilon_0} \frac{q}{|\mathbf{x}'|^3} \mathbf{x},$$

$$\mathbf{B}(\mathbf{x}) = \gamma \frac{\mathbf{v}}{c^2} \times \mathbf{E}'(\mathbf{x}'), \quad \gamma = \sqrt{1 - \mathbf{v}^2/c^2}$$

- ▶ How are Lab frame coordinates (t, \mathbf{x}) and Rest frame coordinates (t', \mathbf{x}') related? (they label the same point in space and time)

Magnetic fields: Produced by charges in motion and affects charges in motion

- ▶ Consider a point charge, q moving in the Lab frame at a velocity, \mathbf{v} . What fields does it produce?
- ▶ In the frame attached to the charge aka the REST FRAME: Only observe an electric field,

$$\mathbf{E}'(\mathbf{x}') = \frac{1}{4\pi\epsilon_0} \frac{q}{|\mathbf{x}'|^3} \mathbf{x}'$$

- ▶ But in the Lab frame: Observe both Electric *and* Magnetic field!,

$$\mathbf{E}(\mathbf{x}) = \frac{\gamma}{4\pi\epsilon_0} \frac{q}{|\mathbf{x}'|^3} \mathbf{x},$$

$$\mathbf{B}(\mathbf{x}) = \gamma \frac{\mathbf{v}}{c^2} \times \mathbf{E}'(\mathbf{x}'), \quad \gamma = \sqrt{1 - \mathbf{v}^2/c^2}$$

- ▶ How are Lab frame coordinates (t, \mathbf{x}) and Rest frame coordinates (t', \mathbf{x}') related? (they label the same point in space and time)
- ▶ Answer is provided by Lorentz transformation (special relativity)

$$\mathbf{x}' = \mathbf{x} + \left(\frac{\gamma - 1}{\mathbf{v}^2} \right) (\mathbf{v} \cdot \mathbf{x}) \mathbf{v} - \gamma \mathbf{v} t,$$

$$t' = \gamma t - \gamma \mathbf{v} \cdot \mathbf{x}$$

Magnetic fields: Produced by charges in motion and affects charges in motion

- ▶ Consider a point charge, q moving in the Lab frame at a velocity, \mathbf{v} . What fields does it produce?
- ▶ In the frame attached to the charge aka the REST FRAME: Only observe an electric field,

$$\mathbf{E}'(\mathbf{x}') = \frac{1}{4\pi\epsilon_0} \frac{q}{|\mathbf{x}'|^3} \mathbf{x}'$$

- ▶ But in the Lab frame: Observe both Electric *and* Magnetic field!,

$$\mathbf{E}(\mathbf{x}) = \frac{\gamma}{4\pi\epsilon_0} \frac{q}{|\mathbf{x}'|^3} \mathbf{x},$$

$$\mathbf{B}(\mathbf{x}) = \gamma \frac{\mathbf{v}}{c^2} \times \mathbf{E}'(\mathbf{x}'), \quad \gamma = \sqrt{1 - \mathbf{v}^2/c^2}$$

- ▶ How are Lab frame coordinates (t, \mathbf{x}) and Rest frame coordinates (t', \mathbf{x}') related? (they label the same point in space and time)
- ▶ Answer is provided by Lorentz transformation (special relativity)

$$\mathbf{x}' = \mathbf{x} + \left(\frac{\gamma - 1}{\mathbf{v}^2} \right) (\mathbf{v} \cdot \mathbf{x}) \mathbf{v} - \gamma \mathbf{v} t,$$

$$t' = \gamma t - \gamma \mathbf{v} \cdot \mathbf{x}$$

Magnetic fields: Produced by charges in motion and affects charges in motion

²I am using upper case letters to denote this charge and velocity of this test charge to distinguish it from the source charge which produced the magnetic field

Magnetic fields: Produced by charges in motion and affects charges in motion

- For Low speeds., $|\mathbf{v}| \ll c$, $\gamma \sim 1$, we get back familiar Galilean transformation

$$\begin{aligned}\mathbf{x}' &= \mathbf{x} - \mathbf{v}t, \\ t' &= t.\end{aligned}$$

²I am using upper case letters to denote this charge and velocity of this test charge to distinguish it from the source charge which produced the magnetic field

Magnetic fields: Produced by charges in motion and affects charges in motion

- For Low speeds., $|\mathbf{v}| \ll c$, $\gamma \sim 1$, we get back familiar Galilean transformation

$$\begin{aligned}\mathbf{x}' &= \mathbf{x} - \mathbf{v}t, \\ t' &= t.\end{aligned}$$

- For charges moving with low speeds,

$$\mathbf{E}(\mathbf{x}, t) \approx \frac{q}{4\pi\epsilon_0} \frac{\mathbf{x}}{|\mathbf{x} - \mathbf{v}t|^3},$$

$$\mathbf{B}(\mathbf{x}, t) \approx q \frac{\mathbf{v}}{c^2} \times \frac{\mathbf{x}}{|\mathbf{x} - \mathbf{v}t|^3},$$

²I am using upper case letters to denote this charge and velocity of this test charge to distinguish it from the source charge which produced the magnetic field

Magnetic fields: Produced by charges in motion and affects charges in motion

- ▶ For Low speeds., $|\mathbf{v}| \ll c$, $\gamma \sim 1$, we get back familiar Galilean transformation

$$\begin{aligned}\mathbf{x}' &= \mathbf{x} - \mathbf{v}t, \\ t' &= t.\end{aligned}$$

- ▶ For charges moving with low speeds,

$$\mathbf{E}(\mathbf{x}, t) \approx \frac{q}{4\pi\epsilon_0} \frac{\mathbf{x}}{|\mathbf{x} - \mathbf{v}t|^3},$$

$$\mathbf{B}(\mathbf{x}, t) \approx q \frac{\mathbf{v}}{c^2} \times \frac{\mathbf{x}}{|\mathbf{x} - \mathbf{v}t|^3},$$

- ▶ How do magnetic field affect charges, say a charge of Q moving with velocity \mathbf{V} in a region with magnetic field \mathbf{B} ²?

²I am using upper case letters to denote this charge and velocity of this test charge to distinguish it from the source charge which produced the magnetic field

Magnetic fields: Produced by charges in motion and affects charges in motion

- ▶ For Low speeds., $|\mathbf{v}| \ll c$, $\gamma \sim 1$, we get back familiar Galilean transformation

$$\begin{aligned}\mathbf{x}' &= \mathbf{x} - \mathbf{v}t, \\ t' &= t.\end{aligned}$$

- ▶ For charges moving with low speeds,

$$\mathbf{E}(\mathbf{x}, t) \approx \frac{q}{4\pi\epsilon_0} \frac{\mathbf{x}}{|\mathbf{x} - \mathbf{v}t|^3},$$

$$\mathbf{B}(\mathbf{x}, t) \approx q \frac{\mathbf{v}}{c^2} \times \frac{\mathbf{x}}{|\mathbf{x} - \mathbf{v}t|^3},$$

- ▶ How do magnetic field affect charges, say a charge of Q moving with velocity \mathbf{V} in a region with magnetic field \mathbf{B} ²?
- ▶ Answer provided by Lorentz force law,

$$\mathbf{F} = Q \mathbf{V} \times \mathbf{B}.$$

²I am using upper case letters to denote this charge and velocity of this test charge to distinguish it from the source charge which produced the magnetic field