Electronic Devices Lecture 8 23-08-2018

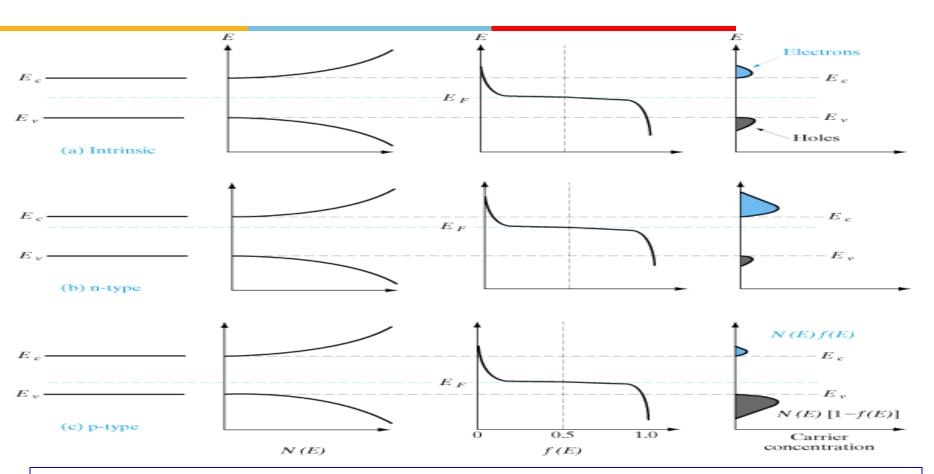
Density of states function

Table 3.1 | Effective density of states function and effective mass values

	N_e (cm ⁻³)	N_v (cm ⁻³)	m_n^*/m_0	m_p^*/m_0
Silicon	2.8×10^{19}	1.04×10^{19}	1.08	0.56
Gallium arsenide	4.7×10^{17}	7.0×10^{18}	0.067	0.48
Germanium	1.04×10^{19}	6.0×10^{18}	0.55	0.37

$$N_C = 2 \left(\frac{2\pi m_n^* kT}{h^2} \right)^{3/2} \qquad N_V = 2 \left(\frac{2\pi m_p^* kT}{h^2} \right)^{3/2}$$

Schematic Band Diagram



Schematic band diagram, density of states, Fermi-Dirac distribution, and the carrier concentrations for (a) intrinsic, (b) n-type, and (c) p-type semiconductors at thermal equilibrium.

2018-08-31

3

• Recall these two equations :

$$egin{aligned} & oldsymbol{n}_0 = oldsymbol{N}_C oldsymbol{f}(oldsymbol{E}_C) = oldsymbol{N}_C oldsymbol{e}^{-(E_C - E_F)/kT} \ & oldsymbol{p}_0 = oldsymbol{N}_V [1 - oldsymbol{f}(oldsymbol{E}_V)] = oldsymbol{N}_V oldsymbol{e}^{-(E_F - E_V)/kT} \end{aligned}$$

Intrinsic electron and hole concentrations where almost $E_F = E_i$:

$$\begin{aligned} & p_{i} = N_{V}e^{-(E_{i}-E_{V})/kT} & n_{i} = N_{C}e^{-(E_{C}-E_{i})/kT} \\ & n_{0}p_{0} = (N_{C}e^{-(E_{C}-E_{F})/kT})(N_{V}e^{-(E_{F}-E_{V})/kT}) = \\ & N_{C}N_{V}e^{-(E_{C}-E_{V})/kT} = N_{C}N_{V}e^{-E_{g}/kT} \\ & n_{i}p_{i} = (N_{C}e^{-(E_{C}-E_{i})/kT})(N_{V}e^{-(E_{i}-E_{V})/kT}) = \\ & N_{C}N_{V}e^{-(E_{C}-E_{V})/kT} = N_{C}N_{V}e^{-E_{g}/kT} \end{aligned}$$

2018-08-31

achieve

innovate

lead

Thus, intrinsic electron and hole concentrations are equal

since the carriers lead

are created in pairs : $n_i = p_i$:

Thus, intrinsic concentration:

Also,

Note: n_i of Si at RT = 1.5 x 10¹⁰ cm⁻³ & $E_C - E_i = E_g/2$ if $N_C = N_V$

$$\boldsymbol{n}_0 \boldsymbol{p}_0 = \boldsymbol{n}_i^2$$

$$n_i = \sqrt{N_C N_V} e^{-E_g/2kT}$$

Two convenient expressions:

$$n_0 = N_C e^{-(E_C - E_F)/kT} = N_C e^{-(E_C - E_i)/kT} e^{(E_F - E_i)/kT} = n_i e^{(E_F - E_i)/kT}$$

Thus,

$$p_0 = N_V e^{-(E_F - E_V)/kT} = N_V e^{-(E_i - E_V)/kT} e^{(E_i - E_F)/kT} = n_i e^{(E_i - E_F)/kT}$$

$$\boldsymbol{n}_0 = \boldsymbol{n}_i \boldsymbol{e}^{(\boldsymbol{E}_F - \boldsymbol{E}_i)/kT}$$

$$\boldsymbol{p}_0 = \boldsymbol{n}_i \boldsymbol{e}^{(\boldsymbol{E}_i - \boldsymbol{E}_F)/kT}$$

Find the equilibrium electron and hole concentrations and the location of the Fermi level (with respect to the intrinsic Fermi level E_i) in silicon at 300K if the silicon contains 8×10^{16} cm⁻³ Arsenic (As) and 2×10^{16} cm⁻³ boron (B) atoms.

$$n = 6 \times 10^{16} \text{ cm}^{-3}$$

$$n_0 = n_i e^{(E_F - E_i)/kT}$$

$$p = \frac{n_i^2}{n} = 3.5 \times 10^3 \text{ cm}^{-3}$$

$$p_0 = n_i e^{(E_i - E_F)/kT}$$

$$E_f - E_i = kT \ln(n/n_i) = 0.393eV$$

$$E_c - E_f = kT \ln(N_c/n) = 0.0258\ln(2.8 \times 10^{19}/6 \times 10_{16})$$

$$= 0.159eV$$

Temperature Dependence of Carrier Concentrations



$$n_i = \sqrt{N_C N_V} e^{-E_g/2kT}$$

$$N_{C} = 2 \left(\frac{2\pi m_{n}^{*} kT}{h^{2}} \right)^{3/2}$$

$$N_{V} = 2 \left(\frac{2\pi m_{p}^{*} kT}{h^{2}} \right)^{3/2} \qquad n_{i}(T) = 2 \left(\frac{2\pi kT}{h^{2}} \right)^{3/2} (m_{n}^{*} m_{p}^{*})^{3/4} e^{-E_{g}/2kT}$$

Take n_i to calculate n_0 and p_0 in the following equations:

$$\boldsymbol{n}_0 = \boldsymbol{n}_i \boldsymbol{e}^{(\boldsymbol{E}_F - \boldsymbol{E}_i)/kT}$$
 $\boldsymbol{p}_0 = \boldsymbol{n}_i \boldsymbol{e}^{(\boldsymbol{E}_i - \boldsymbol{E}_F)/kT}$

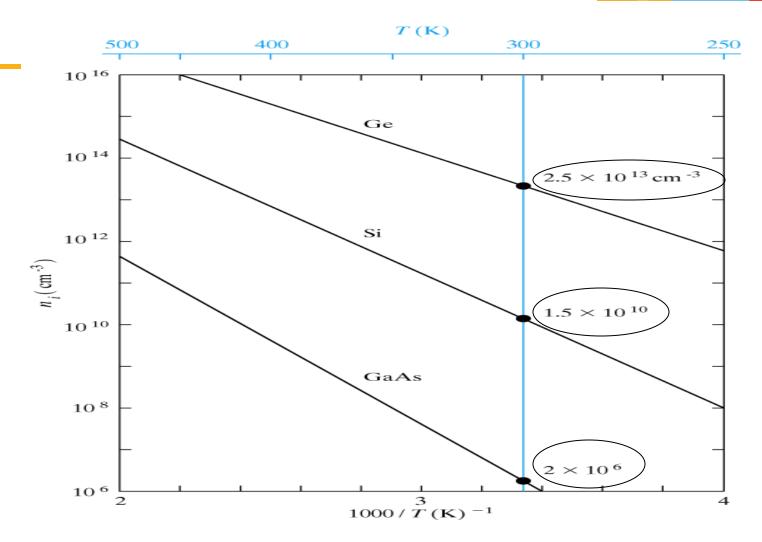


Figure 3—17
Intrinsic carrier concentration for Ge, Si, and GaAs as a function of inverse temperature. The room temperature values are marked for reference.

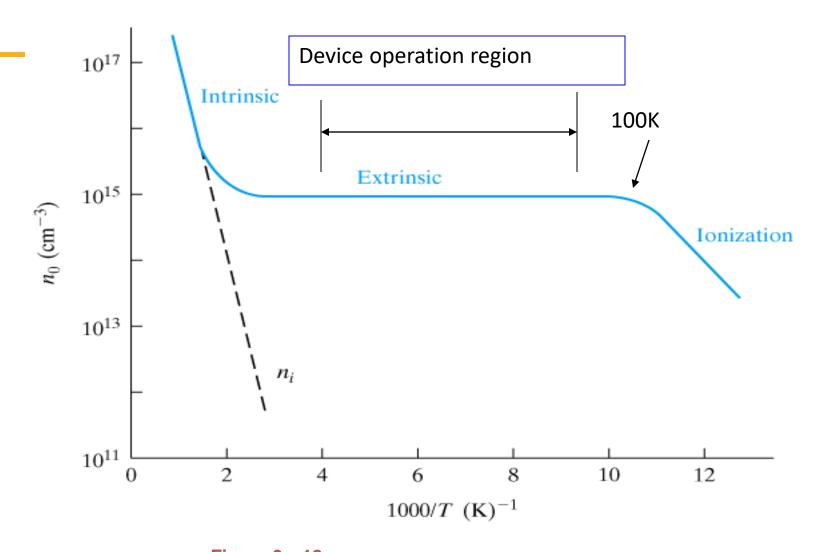
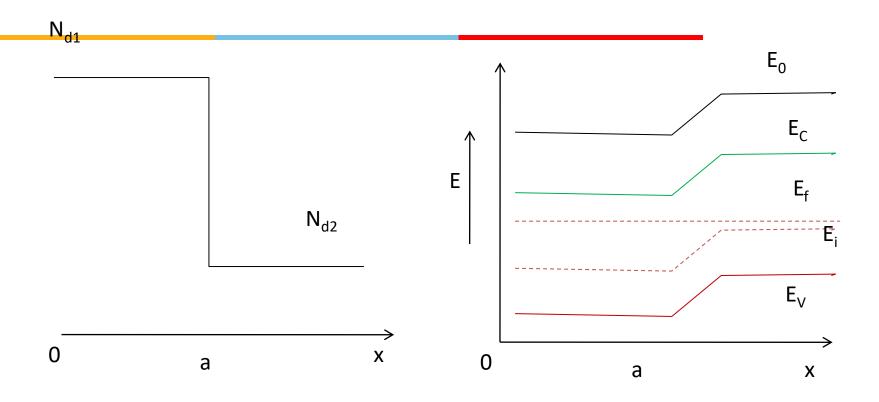


Figure 3—18
Carrier concentration vs. inverse temperature for Si doped with 10¹⁵ donors/cm³.

A silicon crystal is known to contain 10-4 atomic percent of arsenic (As) as an impurity. It then receives a uniform doping 3×1016 cm-3 phosphorous (P) atoms and a subsequent uniform doping of 1018cm-3 boron (B) atoms. A thermal annealing treatment then completely activates all impurities.

- (a) What is the conductivity type of this silicon sample?
- (b) What is the density of the majority carriers?

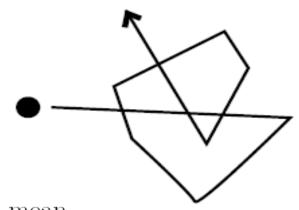


Electronic Devices Lecture 9 28-08-2018

Thermal Motion

In thermal equilibrium, carriers are not sitting still:

- undergo collisions with vibrating Si atoms (*Brownian motion*)
- electrostatically interact with charged dopants and with each other



Characteristic time constant of thermal motion - mean free time between collisions:

$$\tau_c \equiv collision \ time \ [s]$$

In between collisions, carriers acquire high velocity:

$$v_{th} \equiv thermal\ velocity\ [cm/s]$$

$$\lambda \equiv mean free path [cm]$$

$$\lambda = v_{th}\tau_c$$

Put numbers for Si at 300 K:

$$\tau_c \simeq 10^{-14} \sim 10^{-13} \ s$$

$$v_{th} \simeq 10^7 \ cm/s$$

$$\Rightarrow \lambda \simeq 1 \sim 10 \ nm$$

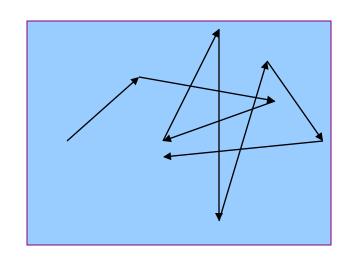
For reference, state-of-the-art MOSFETs today:

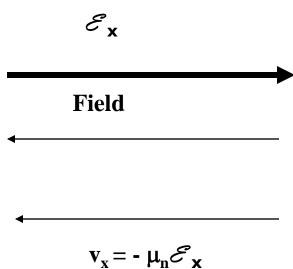
$$L_q \simeq 50 \ nm$$

 \Rightarrow carriers undergo many collisions in modern devices

To calculate current flow in presence of electric and magnetic field we have to take into account the collisions of the charge carriers with the lattice and impurities.

Ease with which electrons and holes can flow through the crystal determines their mobility within the solid.





Net force experienced by each electron $-q\mathcal{E}_{x}$

$$-nq\mathbf{E}_{x} = \frac{dp_{x}}{dt}\bigg|_{field}$$

Let N(t) is the number of electrons that have not undergone collision by time t.

The rate of decrease in N(t) at any time t is

$$-\frac{dN(t)}{dt} = \frac{1}{t}N(t)$$

$$N(t) = N_o e^{-t/t}$$

CONDUCTIVITY AND MOBILITY

The probability that any electron has a collision in time dt is - dt/\bar{t}

The differential change is p_x due to collisions is

$$dp_{x} = -p_{x} \frac{dt}{t}$$

$$\frac{dp_{x}}{dt} \Big|_{collision} = -\frac{p_{x}}{t}$$

$$-\frac{p_{x}}{t} - nq\mathcal{E}_{x} = 0$$

$$\langle p_{x} \rangle = \frac{p_{x}}{n} = -q^{T} \mathcal{E}_{x}$$

$$\langle v_{x} \rangle = \frac{\langle p_{x} \rangle}{m_{n}^{*}} = -\frac{q^{T}}{m_{n}^{*}} \mathcal{E}_{x}$$

-The average acceleration and deceration must be zero for steady state

-Average momentum

-Average velocity m_n^* - conductivity effective mass for electrons

CONDUCTIVITY AND MOBILITY



for electrons:
$$v_{dn} = -\mu_n E$$

for holes:
$$v_{dp} = \mu_p E$$

Drift current

Net velocity of charged particles \Rightarrow electric current:

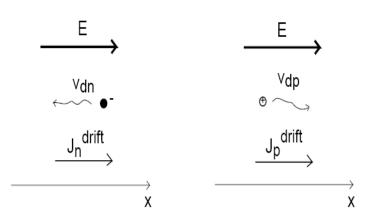
Drift current density $\propto carrier drift velocity$

 $\propto carrier\ concentration$

 $\propto carrier\ charge$

$$J_n^{drift} = -qnv_{dn} = qn\mu_n E$$

$$J_p^{drift} = qpv_{dp} = qp\mu_p E$$



achieve



CONDUCTIVITY AND MOBILITY

Total drift current:

$$J^{drift} = J_n^{drift} + J_p^{drift} = q(n\mu_n + p\mu_p)E$$

Has the shape of *Ohm's Law*:
$$J = \sigma E = \frac{E'}{\rho}$$

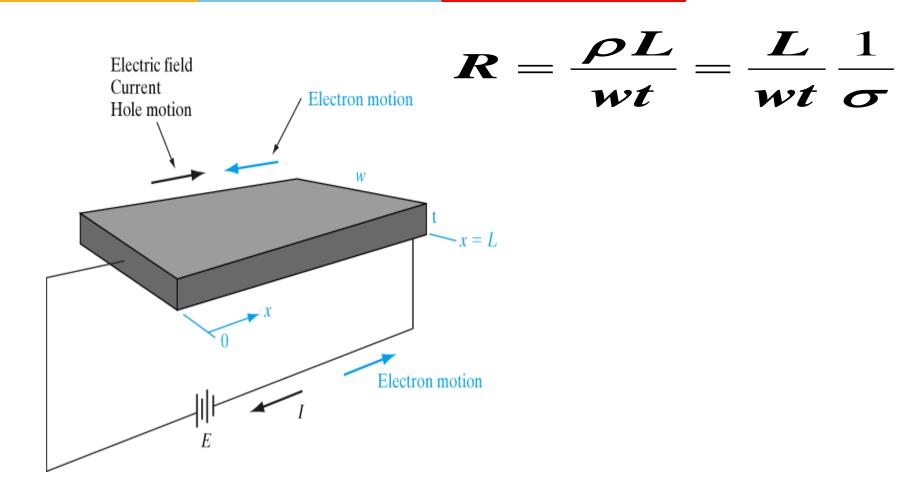
$$\sigma \equiv conductivity \left[\Omega^{-1} \cdot cm^{-1}\right]$$

$$\rho \equiv resistiviy [\Omega \cdot cm]$$

$$\rho = \frac{1}{\sigma} = \frac{1}{q(n\mu_n + p\mu_p)}$$

Drift of electrons and holes in a semiconductor bar





Impurity & Lattice Scatterings

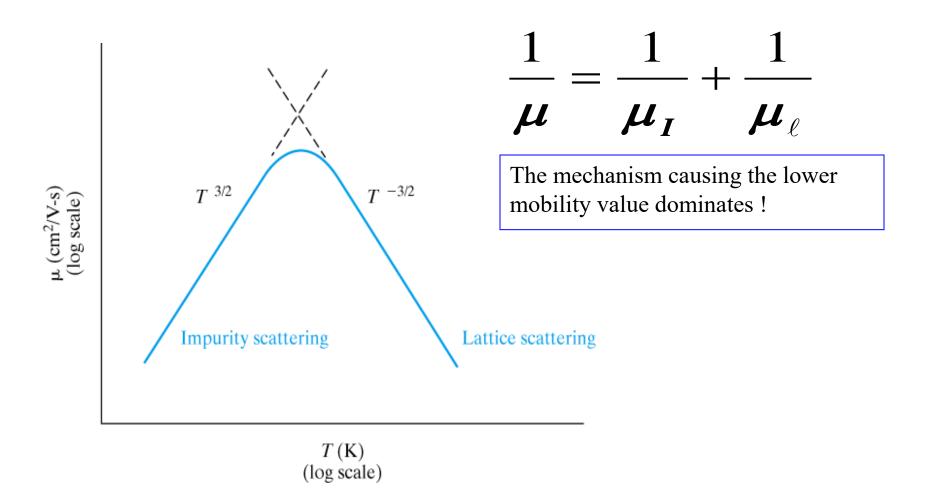
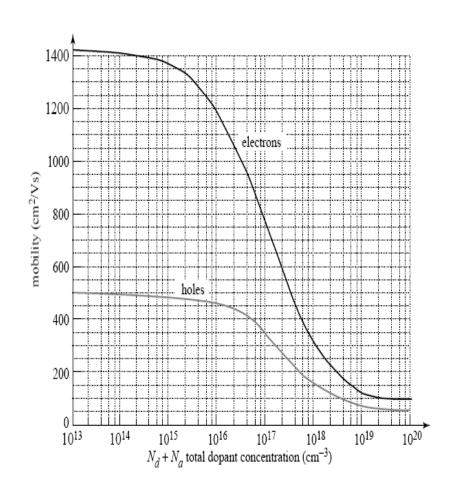


Figure 3—21

Approximate temperature dependence of mobility with both lattice and impurity scattering.

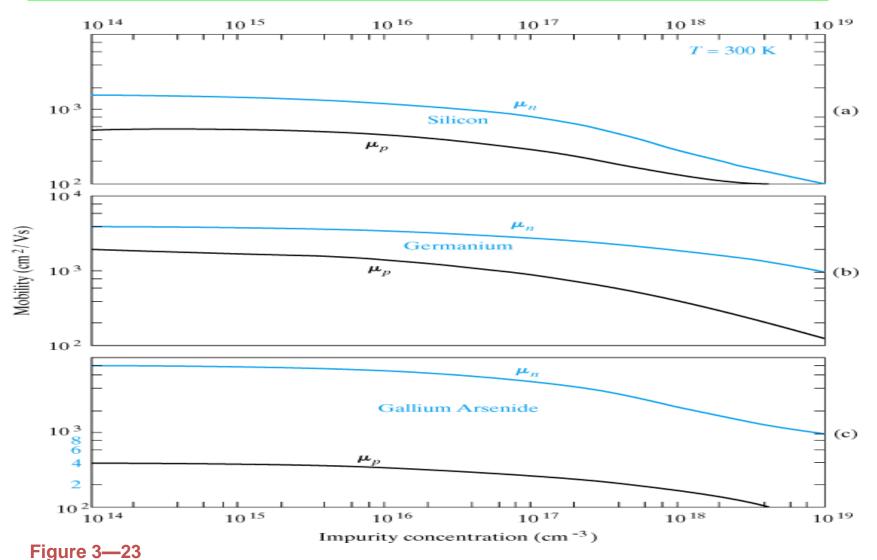
 $\mu - \alpha - T^{-3/2}$ lattice scattering

 μ α $T^{3/2}$ impurity scattering



Mobility depends on doping. For Si at 300K:

Mobility vs. doped impurity concentration

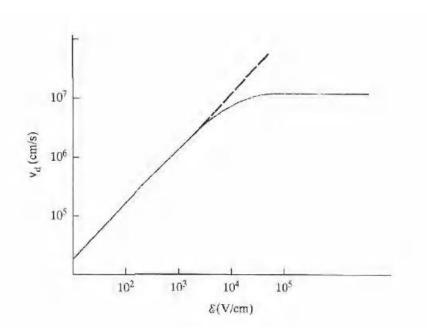


Variation of mobility with total doping impurity concentration ($N_a + N_d$) for Ge, Si, GaAs at 300 K.

2018-08-31

Saturation of electron drift velocity at high electric field

Velocity reaches a saturation value of the mean thermal velocity.



Field required to saturate velocity:

$$\mathcal{E}_{sat} = \frac{v_{sat}}{\mu}$$

Since μ depends on doping, \mathcal{E}_{sat} depends on doping too.

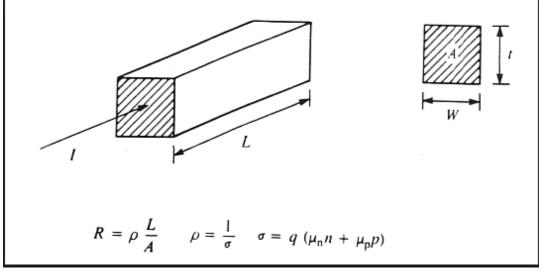
Electrical conductivity

- $\sigma_i = e n_i (\mu_n + \mu_p)$ is the intrinsic conductivity of a semiconductor material
- For extrinsic semiconductors assuming complete ionization, N_d or $N_a\gg n_i$ Hence the conductivity reduces to

 $\sigma_n \approx e N_d \mu_n$ and $\sigma_p \approx e N_a \mu_p$ thus the conductivity is purely dependent on the majority carrier concentration

Electrical resistance

- $J = \sigma E = \sigma \frac{V}{L}$ where L is the length of the SC material
- $I = JA = \frac{\sigma VA}{L}$ and the elctrical resistance is given by
- $R = \frac{V}{I} = \frac{1}{\sigma} \frac{L}{A} = \frac{\rho L}{A}$



Electrical resistivity

- $\rho = \frac{1}{\sigma} = \frac{1}{e(n\mu_n + p\mu_p)}$ is the electrical resistivity
- Resistance, conductance, resistivity and conductivity, depend only on the majority carrier concentration and not on the minority carrier concentration