

EE225E Project Proposal: Metabolism mapping with prior information

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0.1 Dynamic metabolic MRI using hyperpolarized carbon-13 pyruvate

Hyperpolarized carbon-13 magnetic resonance imaging (MRI) has enabled the real-time observation of perfusion and metabolism in preclinical and clinical studies [1, 2, 3, 4, 5, 6]. This technology is made possible by techniques for dynamic nuclear polarization (DNP) that have led to signal-to-noise ratio (SNR) increases of four to five orders of magnitude compared with endogenous signal in dissolved ^{13}C -labelled molecules [7, 8]. Injected $[1-^{13}\text{C}]$ pyruvate is frequently used as a substrate in metabolism experiments and its rate of conversion to $[1-^{13}\text{C}]$ lactate has been shown to distinguish between healthy and diseased tissues in animal [2], and recently human [4], studies.

In contrast with conventional MRI, hyperpolarized experiments are inherently dynamic as images must be acquired as the injected substrate spreads through the body and is metabolized. This necessitates dynamical system modelling and estimation for quantifying metabolic reaction rates.

0.2 Mathematical model

We consider a two-dimensional system of ordinary differential equations

$$\frac{dx}{dt}(t) = \begin{bmatrix} -k_{PL} - R_{1P} & 0 \\ k_{PL} & -R_{1L} \end{bmatrix} x(t) + \begin{bmatrix} k_{TRANS} \\ 0 \end{bmatrix} u(t) \quad (1)$$

that models the magnetization dynamics in a tissue with an arterial input function $u(t)$ and uni-directional conversion from the substrate (pyruvate) to a metabolic product (lactate), which has been commonly applied for hyperpolarized ^{13}C pyruvate experiments. The state $x_1(t)$ denotes the longitudinal magnetization of pyruvate contained in a particular voxel in the tissue and $x_2(t)$ the longitudinal magnetization of lactate in the same voxel. The rate of metabolism of pyruvate to lactate is denoted k_{PL} , the perfusion rate from the arterial input to the tissue is denoted k_{TRANS} , and R_{1P} and R_{1L} are lumped parameters that account for T_1 decay in the magnetization along with other effects, such as metabolism of pyruvate into products other than lactate as well as flow of magnetization out of the slice. The input to the system $u(t)$ is an unmeasured arterial input function (AIF) resulting from the injection of hyperpolarized $[1-^{13}\text{C}]$ pyruvate and is assumed to be of gamma-variate shape

$$u(t) = A_0(t - t_0)^\gamma e^{-(t-t_0)/\beta}.$$

We acquire data at N time points separated by intervals of length T_R . Each time t an acquisition is made, we must choose a flip angle $\alpha_{k,t}$ for each compound k to be measured. If the magnetization of the k -th compound before the acquisition is x_k , then this choice of flip angle allows us to measure a signal of magnitude $\sin(\alpha_{k,t})x_k$, after which $\cos(\alpha_{k,t})x_k$ magnetization remains for future acquisitions. This causes discrete jumps, or resets, in the system state, leading to a hybrid dynamical system [9, 10]. Since we are only interested in the system's state at acquisition times, we can avoid technicalities associated with hybrid system modelling by discretizing the system in time and considering a discrete-time dynamical system that

simultaneously captures the evolution of (1) between acquisitions and the discrete jumps induced by the acquisitions. We define the transition matrices A_d and B_d

$$A_d = \exp \left(T_R \begin{bmatrix} -k_{PL} - R_{1P} & 0 \\ k_{PL} & -R_{1L} \end{bmatrix} \right)$$

$$B_d = \begin{bmatrix} -k_{PL} - R_{1P} & 0 \\ k_{PL} & -R_{1L} \end{bmatrix}^{-1} (A_d - I) \begin{bmatrix} k_{TRANS} \\ 0 \end{bmatrix}$$

that correspond to the discretization of (1) assuming a zero-order hold on the input between each acquisition [11].

We will construct metabolite maps using magnitude image data. This allows us to avoid modelling sources of phase in the image and keeps the model dimension small, requiring only a single state for each metabolite. Accordingly, we model the measurements as independent Rician-distributed random variables [12], which have probability density

$$p_{x,\sigma}(y) = \frac{y}{\sigma^2} \exp \left(-\frac{y^2 + x^2}{2\sigma^2} \right) I_0 \left(\frac{yx}{\sigma^2} \right)$$

where I_ν denotes the modified Bessel function of the first kind of order ν . All together, we have the discrete-time model for the observed data from a single voxel:

$$x_0 = 0$$

$$x_{t+1} = A_d(\theta) \begin{bmatrix} \cos \alpha_{1,t} & 0 \\ 0 & \cos \alpha_{2,t} \end{bmatrix} x_t + B_d(\theta) u_t(\theta) \quad (2)$$

$$\tilde{x}_{k,t} = \sin(\alpha_{k,t}) x_{k,t} \quad k = 1, 2$$

$$Y_{k,t} \sim \text{Rice}(\tilde{x}_{k,t}, \sigma_k) \quad k = 1, 2.$$

0.3 Low SNR degrades parameter map quality

Even with dynamic nuclear polarization, carbon-13 MRI suffers from low SNR relative to conventional MRI. This makes it difficult to acquire high-resolution images, as large voxel sizes are required to achieve sufficient SNR. In Figure 1 we show the effect of increased resolution on the quality of parameter maps computed by an independent maximum-likelihood fit to each voxel, assuming that SNR of the trajectories of (2) is proportional to voxel volume. We see that when voxels are treated independently (*i.e* no spatial structure is assumed in the parameter maps) increased resolution can be detrimental to parameter map quality.

0.4 A Douglas-Rachford splitting algorithm for spatial regularization

In order to overcome this limitation, we propose to investigate incorporating prior information about parameter maps by including regularization in the mapping. This will allow spatial structure in the data to be exploited, which was ignored in the previous analysis. We hope that this will help to solve the problem of low SNR in high-resolution carbon-13 images.

Let y_i denote a dynamic dataset generated by the model (2) corresponding to some voxel i from the set of all voxels \mathcal{V} . We wish to solve the optimization problem

$$\text{minimize} \quad \sum_{i \in \mathcal{V}} -\ell(\theta_i | y_i) + \lambda r(\theta)$$

where $\ell(\theta_i | y_i)$ denotes the log likelihood function corresponding to the model (2) and r is a regularizer that corresponds to prior information we have about the parameter map. For example, for the parameter map in Figure 1 we might use a total-variation regularization term in order to exploit the fact that we know that the map is piecewise-constant.

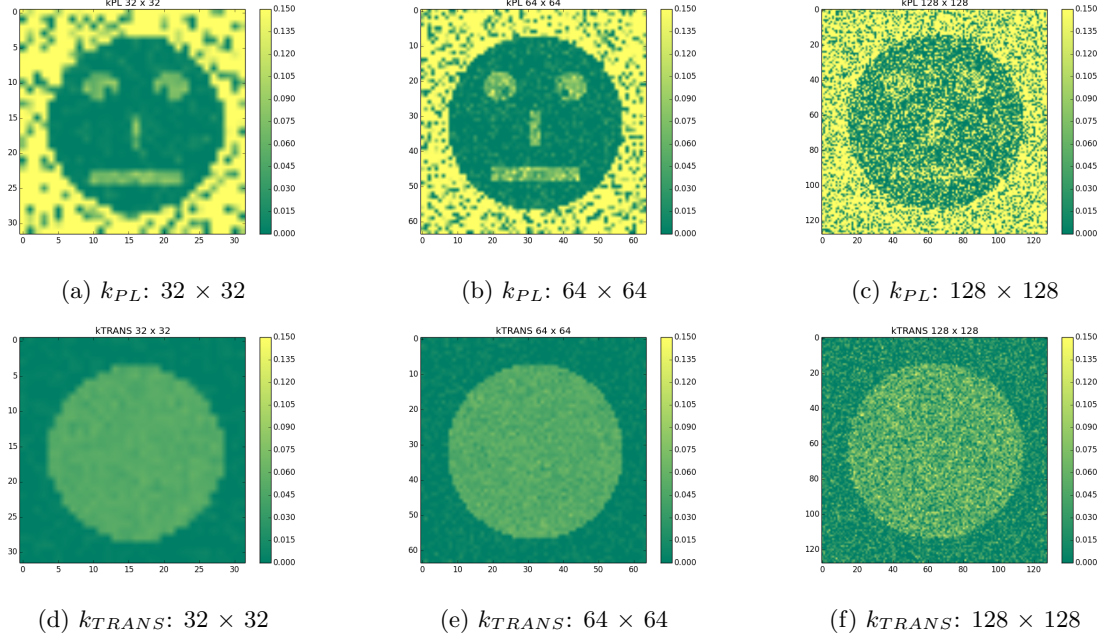


Figure 1: Parameter maps of varying resolution fit to simulated data sets. The first row shows parameter maps for the metabolic rate parameter k_{PL} ; the second row shows parameter maps for the perfusion parameter k_{TRANS} .

To solve 0.4 we introduce $z = \theta$ and solve

$$\begin{aligned} & \text{minimize} && \sum_{i \in \mathcal{V}} -\ell(\theta_i | y_i) + \lambda r(z) \\ & \text{subject to} && \theta - z = 0 \end{aligned} \tag{3}$$

This can be solved using the Douglas-Rachford iteration (a special case of ADMM)

$$\begin{aligned} \theta^{k+1} &= \operatorname{argmin}_{\theta} \sum_{i \in \mathcal{V}} -\ell(\theta_i | y_i) + \frac{\rho}{2} \|\theta - z^k + u^k\|_2^2 \\ z^{k+1} &= \operatorname{argmin}_z \lambda r(z) + \frac{\rho}{2} \|\theta^{k+1} - z + u^k\|_2^2 \\ u^{k+1} &= u^k + \theta^{k+1} - z^{k+1} \end{aligned}$$

Note that both the θ update is additively separable. Introducing the proximity operator

$$\operatorname{prox}_f(x) = \operatorname{argmin}_u f(u) + \frac{1}{2} \|u - x\|_2^2$$

we can re-write this iteration as

$$\begin{aligned} \theta_i^{k+1} &= \operatorname{prox}_{-\frac{1}{\rho} \ell(\cdot | y_i)}(z_i^k - u_i^k) \quad i \in \mathcal{V} \\ z^{k+1} &= \operatorname{prox}_{\frac{\lambda}{\rho} r}(\theta^{k+1} + u^k) \\ u^{k+1} &= u^k + \theta^{k+1} - z^{k+1}. \end{aligned}$$

0.5 Project goals:

- Implement this algorithm and investigate convergence properties (since the negative log likelihood in this case is non-convex).
- Explore whether similar things have been done before (*e.g.* in dynamic contrast enhanced MRI).
- Investigate applying to more realistic simulated data sets (*e.g.* subsampled k-space).

References

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