Solutions - Chapter 6

Kevin S. Huang

Problem 6.1

a) Base case:

$$[\hat{x}, \hat{p}_x] = i\hbar = i\hbar(1)\hat{x}^{1-1}$$

Assume:

$$[\hat{x}^k, \hat{p}_x] = i\hbar k\hat{x}^{k-1}$$

Induction:

$$\begin{split} [\hat{x}^{k+1}, \hat{p}_x] &= \hat{x}^{k+1} \hat{p}_x - \hat{p}_x \hat{x}^{k+1} = \hat{x} (\hat{x}^k \hat{p}_x) - (\hat{p}_x \hat{x}) \hat{x}^k \\ &= \hat{x} ([\hat{x}^k, \hat{p}_x] + \hat{p}_x \hat{x}^k) + ([\hat{x}, \hat{p}_x] - \hat{x} \hat{p}_x) \hat{x}^k = \hat{x} [\hat{x}^k, \hat{p}_x] + [\hat{x}, \hat{p}_x] \hat{x}^k \\ &= \hat{x} i \hbar k \hat{x}^{k-1} + i \hbar \hat{x}^k = i \hbar (k+1) \hat{x}^k \end{split}$$

$$[\hat{x}^n, \hat{p}_x] = i\hbar n\hat{x}^{n-1}$$

b)
$$F(x) = F(a) + F'(a)(x - a) + \frac{F''(a)}{2!}(x - a)^2 + \dots = \sum_{0}^{\infty} \frac{F^{(n)}(a)}{n!}(x - a)^n$$

$$F'(x) = F'(a) + F''(a)(x - a) + \frac{F'''(a)}{2!}(x - a)^2 + \dots = \sum_{1}^{\infty} \frac{F^{(n)}(a)}{(n - 1)!}(x - a)^{(n - 1)}$$

$$\left[\sum_{0}^{\infty} \frac{F^{(n)}(0)}{n!} \hat{x}^n, \hat{p}_x\right] = \sum_{0}^{\infty} \frac{F^{(n)}(0)}{n!} [\hat{x}^n, \hat{p}_x] = i\hbar \sum_{1}^{\infty} \frac{F^{(n)}(0)}{(n - 1)!} \hat{x}^{n - 1} = i\hbar \frac{\partial F}{\partial x}(\hat{x})$$

$$[F(\hat{x}), \hat{p}_x] = i\hbar \frac{\partial F}{\partial x}(\hat{x})$$

c)
$$\hat{H} = \frac{\hat{p}_x^2}{2m} + V(\hat{x})$$

$$[\hat{H}, \hat{p}_x] = [V(\hat{x}), \hat{p}_x] = i\hbar \frac{dV}{dx}(\hat{x})$$

$$\frac{d\langle p_x \rangle}{dt} = \frac{i}{\hbar} \langle \psi | [\hat{H}, \hat{p}_x] | \psi \rangle = \frac{i}{\hbar} \langle \psi | i\hbar \frac{dV}{dx}(\hat{x}) | \psi \rangle = \left\langle -\frac{dV}{dx} \right\rangle$$

$$\langle p|x\rangle = \frac{1}{\sqrt{2\pi\hbar}} e^{-ipx/\hbar}$$

$$i\hbar \frac{\partial}{\partial p} \langle p|x\rangle = x \langle p|x\rangle$$

$$\langle x|\hat{x}|\psi\rangle = \int dx'x' \langle x|x'\rangle \langle x'|\psi\rangle = x \langle x|\psi\rangle$$

$$\langle p|\hat{x}|\psi\rangle = \int dx \langle p|x\rangle \langle x|\hat{x}|\psi\rangle = \int dx x \langle p|x\rangle \langle x|\psi\rangle = \int dx i\hbar \frac{\partial}{\partial p} \langle p|x\rangle \langle x|\psi\rangle$$

$$\langle p|\hat{x}|\psi\rangle = i\hbar \frac{\partial}{\partial p} \langle p|\psi\rangle$$

$$\langle \varphi|\hat{x}|\psi\rangle = \int dp \langle \varphi|p\rangle \langle p|\hat{x}\psi\rangle = \int dp \langle p|\varphi\rangle^* i\hbar \frac{\partial}{\partial p} \langle p|\varphi\rangle$$

$$\hat{x} \xrightarrow[p-basis]{} i\hbar \frac{\partial}{\partial p}$$

$$|\psi\rangle = \int dx |x\rangle \langle x|\psi\rangle$$

$$|\psi'\rangle = \hat{T}(\delta x) |\psi\rangle = \int dx \, \hat{T}(\delta x) |x\rangle \langle x|\psi\rangle = \int dx |x + \delta x\rangle \langle x|\psi\rangle$$

$$\psi'(x) = \langle x|\psi'\rangle = \int dx' \langle x|x' + \delta x\rangle \langle x'|\psi\rangle = \psi(x - \delta x)$$

$$\langle x\rangle = \langle \psi|x|\psi\rangle = \int dx \, \psi^*(x)x\psi(x)$$

$$\langle x\rangle' = \langle \psi'|x|\psi'\rangle = \int dx \, \psi'^*(x)x\psi'(x) = \int dx \, \psi^*(x - \delta x)x\psi(x - \delta x)$$

$$= \int dx \, \left[\psi^*(x) - \frac{\partial \psi^*}{\partial x} \delta x\right] x \, \left[\psi(x) - \frac{\partial \psi}{\partial x} \delta x\right]$$

Keep first-order terms and integrate by parts (boundary term goes to zero due to normalization constraint):

$$\langle x \rangle' = \int dx \, \psi^*(x) x \psi(x) - \delta x \int dx \, x \left(\psi \frac{\partial \psi^*}{\partial x} + \frac{\partial \psi}{\partial x} \psi^* \right)$$
$$= \langle x \rangle - \delta x \int dx \, x \frac{\partial}{\partial x} (\psi^* \psi) = \langle x \rangle + \delta x \int dx \, |\psi|^2$$

$$\langle x \rangle' = \langle x \rangle + \delta x$$

$$\langle p_x \rangle = \int dx \, \psi^*(x) \frac{\hbar}{i} \frac{\partial}{\partial x} \psi(x)$$

$$\langle p_x \rangle' = \int dx \, \left[\psi^*(x) - \frac{\partial \psi^*}{\partial x} \delta x \right] \frac{\hbar}{i} \frac{\partial}{\partial x} \left[\psi(x) - \frac{\partial \psi}{\partial x} \delta x \right]$$

Keep first-order terms:

$$= \int dx \, \psi^*(x) \frac{\hbar}{i} \frac{\partial}{\partial x} \psi(x) - \frac{(\delta x)\hbar}{i} \int dx \, \left[\frac{\partial \psi^*}{\partial x} \frac{\partial \psi}{\partial x} + \psi^*(x) \frac{\partial^2 \psi}{\partial x^2} \right]$$

Integrate by parts:

$$= \langle p_x \rangle - \frac{(\delta x)\hbar}{i} \int dx \, \frac{\partial}{\partial x} \left[\psi^*(x) \frac{\partial \psi}{\partial x} \right]$$
$$\langle p_x \rangle' = \langle p_x \rangle$$

Problem 6.4

a) Free particle:

$$\hat{H} = \frac{\hat{p}_x^2}{2m}$$

Gaussian wave-packet:

$$\langle x|\psi\rangle = \frac{1}{\sqrt{\sqrt{\pi}a}}e^{-x^2/2a^2}$$

$$\langle p|\psi\rangle = \sqrt{\frac{a}{\hbar\sqrt{\pi}}}e^{-p^2a^2/2\hbar^2}$$

$$|\psi(t)\rangle = e^{-i\hat{H}t/\hbar}\int dp\ |p\rangle\ \langle p|\psi\rangle = \int dp\ e^{-ip^2t/2m\hbar}\ |p\rangle\ \langle p|\psi\rangle$$

$$\psi(x,t) = \int dp\ e^{-ip^2t/2m\hbar}\ \langle x|p\rangle\ \langle p|\psi\rangle = \int dp\ e^{-ip^2t/2m\hbar}\frac{1}{\sqrt{2\pi\hbar}}e^{ipx/\hbar}\sqrt{\frac{a}{\hbar\sqrt{\pi}}}e^{-p^2a^2/2\hbar^2}$$

$$I(a,b) = \int_{-\infty}^{+\infty} dx\ e^{-ax^2+bx} = e^{b^2/4a}\sqrt{\frac{\pi}{a}}$$

$$\psi(x,t) = \frac{1}{\sqrt{2\pi\hbar}}\sqrt{\frac{a}{\hbar\sqrt{\pi}}}\int dp\ e^{-ip^2t/2m\hbar}e^{ipx/\hbar}e^{-p^2a^2/2\hbar^2}$$

$$A = \frac{it}{2m\hbar} + \frac{a^2}{2\hbar^2} = \frac{i\hbar t + ma^2}{2m\hbar^2}$$

$$B = \frac{ix}{\hbar}$$

$$\frac{B^2}{4A} = \frac{-x^2}{\hbar^2} \frac{m\hbar^2}{2(i\hbar t + ma^2)} = -\frac{x^2}{2a^2(1 + i\hbar t/ma^2)}$$

$$\psi(x,t) = \frac{1}{\sqrt{2\pi\hbar}} \sqrt{\frac{a}{\hbar\sqrt{\pi}}} \sqrt{\frac{2\pi m\hbar^2}{i\hbar t + ma^2}} e^{-x^2/2a^2[1 + (i\hbar t/ma^2)]}$$

$$\psi(x,t) = \frac{1}{\sqrt{\sqrt{\pi}[a + (i\hbar t/ma)]}} e^{-x^2/2a^2[1 + (i\hbar t/ma^2)]}$$

$$|\psi(x,t)|^2 = \frac{1}{\sqrt{\pi[a^2 + (\hbar t/ma)^2]}} e^{-x^2/a^2[1 + (i\hbar t/ma^2)^2]}$$

$$\Delta x = \frac{a}{\sqrt{2}} \sqrt{1 + \left(\frac{\hbar t}{ma^2}\right)^2}$$
b)
$$\langle p|\psi(t)\rangle = \int dp' \, e^{-ip'^2t/2m\hbar} \, \langle p|p'\rangle \, \langle p'|\psi\rangle = e^{-ip^2t/2m\hbar} \sqrt{\frac{a}{\hbar\sqrt{\pi}}} e^{-p^2a^2/2\hbar^2}$$

$$\psi(p,t) = \sqrt{\frac{a}{\hbar\sqrt{\pi}}} e^{-p^2(i\hbar t + ma^2)/2m\hbar^2}$$

$$|\psi(p,t)|^2 = \frac{a}{\hbar\sqrt{\pi}} e^{-p^2a^2/\hbar^2}$$
Odd function:
$$\langle p\rangle = \int dp \, |\psi(p,t)|^2 p = \frac{a}{\hbar\sqrt{\pi}} \int dp \, e^{-p^2a^2/\hbar^2} p = 0$$

$$\int dx \, e^{-x^2/a^2} x^2 = \frac{a^3\sqrt{\pi}}{2}$$

$$\langle p^2\rangle = \int dp \, |\psi(p,t)|^2 p^2 = \frac{a}{\hbar\sqrt{\pi}} \int dp \, e^{-p^2a^2/\hbar^2} p^2 = \frac{a}{\hbar\sqrt{\pi}} \left(\frac{\hbar}{a}\right)^3 \frac{\sqrt{\pi}}{2} = \frac{\hbar^2}{2a^2}$$

$$\Delta p = \sqrt{\langle p^2 \rangle - \langle p \rangle^2} = \frac{\hbar}{a\sqrt{2}}$$

$$\langle p|\psi\rangle = \left\{ \begin{array}{ll} 0 & p < -P/2 \\ N & -P/2 < P/2 \\ 0 & p > P/2 \end{array} \right.$$

a) Normalization:

$$\int dp \, |\langle p|\psi\rangle|^2 = 1$$

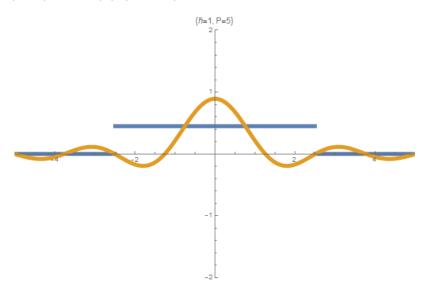
$$N^2 P = 1$$

$$N = \frac{1}{\sqrt{P}}$$

b)
$$\psi(x) = \int dp \, \langle x|p \rangle \, \psi(p) = \int_{-P/2}^{P/2} dp \, \frac{1}{\sqrt{2\pi\hbar}} e^{ipx/\hbar} N = \frac{N}{\sqrt{2\pi\hbar}} \frac{\hbar}{ix} e^{ipx/\hbar} \Big|_{-P/2}^{P/2}$$

$$\psi(x) = \frac{1}{ix} \sqrt{\frac{\hbar}{2\pi P}} (e^{iPx/2\hbar} - e^{-iPx/2\hbar}) = \frac{1}{x} \sqrt{\frac{2\hbar}{\pi P}} \sin\left(\frac{Px}{2\hbar}\right)$$

c) Plot of $\psi(p)$ (blue) and $\psi(x)$ (orange) for $\hbar=1, P=5$



Estimation:

$$\Delta p \approx \frac{P}{4}$$

$$\frac{P\Delta x}{2\hbar} = \pi$$

$$\Delta x \approx \frac{2\pi\hbar}{P}$$

Independent of P:

$$\Delta x \Delta p \approx \frac{\pi \hbar}{2}$$

Problem 6.6

a) Real wave function:

$$\langle x|\psi\rangle = \langle \psi|x\rangle$$

$$\langle p \rangle = \int dx \ \langle \psi | x \rangle \frac{\hbar}{i} \frac{\partial}{\partial x} \langle x | \psi \rangle = \frac{\hbar}{i} \int dx \ \psi(x) \frac{\partial \psi}{\partial x}$$
$$\int dx \ \psi(x) \frac{\partial \psi}{\partial x} = \int dx \ \frac{\partial \psi}{\partial x} \psi(x)$$

Normalization constraint: $\psi \to 0, x \to \pm \infty$

b)
$$\psi(x)' = e^{ip_0x/\hbar}\psi(x)$$

$$\langle x\rangle' = \int dx \, \psi'^* x \psi' = \int dx \, e^{-ip_0x/\hbar}\psi^*(x) x e^{ip_0x/\hbar}\psi(x) = \langle x\rangle$$

$$\langle p\rangle' = \int dx \, \psi'^* \frac{\hbar}{i} \frac{\partial}{\partial x} \psi' = \frac{\hbar}{i} \int dx \, e^{-p_0x/\hbar}\psi^*(x) \frac{\partial}{\partial x} [e^{ip_0x/\hbar}\psi(x)]$$

$$= \frac{\hbar}{i} \int dx \, e^{-p_0x/\hbar}\psi^*(x) \left[\frac{ip_0}{\hbar} e^{ip_0x/\hbar}\psi(x) + e^{ip_0x/\hbar} \frac{\partial}{\partial x}\psi(x) \right]$$

$$= p_0 \int dx \, \psi^*(x)\psi(x) + \int dx \, \psi^*(x) \frac{\hbar}{i} \frac{\partial}{\partial x}\psi(x)$$

$$\langle p\rangle' = \langle p\rangle + p_0$$

Problem 6.7

a)

$$\langle \varphi | \hat{x} | \psi \rangle = \int dx \ \langle \varphi | x \rangle \, x \, \langle x | \psi \rangle = \int dx \ \langle x | \varphi \rangle^* \, x \, \langle \psi | x \rangle^* = \left(\int dx \ \langle \psi | x \rangle \, x \, \langle x | \varphi \rangle \right)^* = \langle \psi | \hat{x} | \varphi \rangle^*$$
b)
$$-[\hat{x}, \hat{p}_x] = \hat{p}_x \hat{x} - \hat{x} \hat{p}_x = -i\hbar$$

$$[\hat{x}, \hat{p}_x]^{\dagger} = (\hat{x} \hat{p}_x)^{\dagger} - (\hat{p}_x \hat{x})^{\dagger} = \hat{p}_x^{\dagger} \hat{x}^{\dagger} - \hat{x}^{\dagger} \hat{p}_x^{\dagger} = \hat{p}_x \hat{x}^{\dagger} - \hat{x}^{\dagger} \hat{p}_x = -i\hbar$$

$$\hat{x}^{\dagger} = \hat{x}$$

Problem 6.8 - SKIPPED

$$\frac{d^2\psi}{dx^2} + \frac{2m}{\hbar^2} (E - a|x|)\psi = 0$$

$$E = \epsilon (\hbar^2 a^2/m)^{1/3}$$

$$x = z(\hbar^2/ma)^{1/3}$$
a)
$$[E] = \frac{ML^2}{T^2}, [x] = L, [\hbar] = \frac{ML^2}{T}, [a] = \frac{ML}{T^2}, [m] = M$$

$$\frac{ML^2}{T^2} = [\epsilon] \left(\frac{M^2L^4}{T^2} \frac{M^2L^2}{T^4} \frac{1}{M} \right)^{1/3} = [\epsilon] \frac{ML^2}{T^2}$$

$$L = [z] \left(\frac{M^2L^4}{T^2} \frac{1}{M} \frac{T^2}{ML} \right)^{1/3} = [z]L$$

$$[\epsilon] = [z] = [\]$$
b)
$$\frac{d\psi}{dx} = \frac{d\psi}{dz} \frac{dz}{dx}$$

$$\frac{dz}{dx} = (\hbar^2/ma)^{-1/3}$$

$$\frac{d\psi}{dx} = \frac{d\psi}{dz} (\hbar^2/ma)^{-1/3}$$

$$\frac{d^2\psi}{dx^2} = \frac{d}{dz} \left(\frac{d\psi}{dx} \right) \frac{dz}{dx} = \frac{d^2\psi}{dz^2} (\hbar^2/ma)^{-2/3}$$

$$\frac{d^2\psi}{dz^2} (\hbar^2/ma)^{-2/3} + \frac{2m}{\hbar^2} [\epsilon(\hbar^2a^2/m)^{1/3} - a|z|(\hbar^2/ma)^{1/3}]\psi = 0$$

$$\frac{d^2\psi}{dz^2} + \frac{2m}{\hbar^2} \left(\frac{\hbar^2a^2}{m} \right)^{1/3} \left(\frac{\hbar^2}{ma} \right)^{2/3} (\epsilon - |z|)\psi = 0$$

$$\frac{d^2\psi}{dz^2} + 2(\epsilon - |z|)\psi = 0$$

c) - SKIPPED

Problem 6.10

$$V(x) = \begin{cases} 0 & 0 < x < L \\ \infty & \text{elsewhere} \end{cases}$$

Outside the potential well:

$$\psi(x) = 0$$

Inside the potential well:

$$-\frac{\hbar^2}{2m}\frac{d^2\psi}{dx^2} = E\psi(x)$$

$$\frac{d^2\psi}{dx^2} = -\omega^2\psi(x), \quad \omega = \sqrt{\frac{2mE}{\hbar^2}}$$

$$\psi(x) = A\sin\omega x + B\cos\omega x$$

Boundary conditions:

$$\psi(0) = B = 0$$

$$\psi(L) = A \sin \omega L = 0$$

Energy eigenvalues:

$$\omega L = n\pi, \quad n = 1, 2, 3, ...$$

$$\sqrt{\frac{2mE_n}{\hbar^2}} = \frac{n\pi}{L}$$

$$E_n = \frac{\hbar^2 \pi^2 n^2}{2mL^2}, \quad n = 1, 2, 3, ...$$

Normalization:

$$A^{2} \int_{0}^{L} dx \sin^{2} \frac{n\pi x}{L} = 1$$
$$A = \sqrt{\frac{2}{L}}$$

Energy eigenfunctions:

$$\psi_n(x) = \begin{cases} \sqrt{\frac{2}{L}} \sin \frac{n\pi x}{L} & 0 < x < L \\ 0 & \text{elsewhere} \end{cases} \quad n = 1, 2, 3, \dots$$

Problem 6.11

Ground state:

e.
$$\psi(x) = \begin{cases} \sqrt{\frac{2}{L}} \sin \frac{\pi x}{L} & 0 < x < L \\ 0 & \text{elsewhere} \end{cases}$$

$$\langle x \rangle = \int dx \, \psi^*(x) x \psi(x) = \frac{2}{L} \int_0^L dx \, x \sin^2 \frac{\pi x}{L} = \frac{L}{2}$$

$$\langle x^2 \rangle = \int dx \, \psi^*(x) x^2 \psi(x) = \frac{2}{L} \int_0^L dx \, x^2 \sin^2 \frac{\pi x}{L} = \frac{1}{6} \left(2 - \frac{3}{\pi^2} \right) L^2$$

$$\Delta x = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = \sqrt{\frac{1}{6} \left(2 - \frac{3}{\pi^2} \right) L^2 - \frac{L^2}{4}} = L \sqrt{\frac{1}{12} - \frac{1}{2\pi^2}}$$

$$\langle p \rangle = \int dx \, \psi^*(x) \frac{\hbar}{i} \frac{d}{dx} \psi(x) = \frac{2\hbar \pi}{i L^2} \int_0^L dx \, \sin \frac{\pi x}{L} \cos \frac{\pi x}{L} = 0$$

$$\langle p^2 \rangle = \int dx \, \psi^*(x) \left(\frac{\hbar}{i} \right)^2 \frac{d^2}{dx^2} \psi(x) = \frac{2\hbar^2 \pi^2}{L^3} \int_0^L dx \, \sin^2 \frac{\pi x}{L} = \frac{\pi^2 \hbar^2}{L^2}$$

$$\Delta p = \sqrt{\langle p^2 \rangle - \langle p \rangle^2} = \frac{\pi \hbar}{L}$$

$$\Delta x \Delta p = \pi \hbar \sqrt{\frac{1}{12} - \frac{1}{2\pi^2}} = \hbar \sqrt{\frac{\pi^2 - 6}{12}} > \frac{\hbar}{2}$$

$$V(x) = \begin{cases} 0 & 0 < x < L \\ \infty & \text{elsewhere} \end{cases}$$

$$\psi(x) = \begin{cases} \left(\frac{1+i}{2}\right)\sqrt{\frac{2}{L}}\sin\frac{\pi x}{L} + \frac{1}{\sqrt{2}}\sqrt{\frac{2}{L}}\sin\frac{2\pi x}{L} & 0 < x < L \\ 0 & \text{elsewhere} \end{cases}$$
a)
$$\psi(x) = \left(\frac{1+i}{2}\right)\psi_1(x) + \frac{1}{\sqrt{2}}\psi_2(x)$$

$$\psi(x,t) = \left(\frac{1+i}{2}\right)e^{-iE_1t/\hbar}\psi_1(x) + \frac{1}{\sqrt{2}}e^{-iE_2t/\hbar}\psi_2(x)$$
b)
$$\langle E \rangle = \left|\frac{1+i}{2}\right|^2 E_1 + \left(\frac{1}{\sqrt{2}}\right)^2 E_2 = \frac{E_1 + E_2}{2}$$

$$\langle E \rangle = \langle \psi | \hat{H} | \psi \rangle = \left[\left(\frac{1-i}{2}\right)\psi_1^*(x) + \frac{1}{\sqrt{2}}\psi_2^*(x)\right] \left[\left(\frac{1+i}{2}\right)E_1\psi_1(x) + \frac{1}{\sqrt{2}}E_2\psi_2(x)\right] = \frac{E_1 + E_2}{2}$$
c)
$$|\psi(t)\rangle = \left(\frac{1+i}{2}\right)e^{-iE_1t/\hbar}|E_1\rangle + \frac{1}{\sqrt{2}}e^{-iE_2t/\hbar}|E_2\rangle$$

$$|\langle E_1|\psi\rangle|^2 = \frac{1}{2}$$

d) $\langle x \rangle$ is time-dependent due to the relative phase difference between the two energy eigenstates. Time enters through the cross terms between the wave functions.

Problem 6.13

a)

$$V(x) = \begin{cases} 0 & 0 < x < L \\ \infty & \text{elsewhere} \end{cases}$$
$$|\psi(0)\rangle = |E_1^L\rangle$$

t=0: Wall pulled back to x=2L instantaneously

$$\langle E_1^{2L} | E_1^L \rangle = \int_0^L dx \, \psi_1^{2L*}(x) \psi_1^L(x) = \int_0^L dx \, \sqrt{\frac{1}{L}} \sin\left(\frac{\pi x}{2L}\right) \sqrt{\frac{2}{L}} \sin\left(\frac{\pi x}{L}\right) = \frac{4\sqrt{2}}{3\pi}$$

$$|\langle E_1^{2L}|E_1^L\rangle|^2 = \frac{32}{9\pi^2} = 0.36$$
 b)
$$|\psi(0)\rangle = \sum_n |E_n^{2L}\rangle \langle E_n^{2L}|E_1^L\rangle$$

$$|\psi(t)\rangle = e^{-i\hat{H}t/\hbar} \sum_n |E_n^{2L}\rangle \langle E_n^{2L}|E_1^L\rangle = \sum_n e^{-iE_n^{2L}t/\hbar} |E_n^{2L}\rangle \langle E_n^{2L}|E_1^L\rangle$$

The system is no longer is an energy eigenstate. Calculate amplitudes $\langle E_n^{2L}|E_1^L\rangle$ and let computer carry out the sum numerically.

Problem 6.14

$$V(x) = \begin{cases} 0 & 0 < x < L \\ \infty & \text{elsewhere} \end{cases}$$

$$\psi(x) = \begin{cases} Nx(x-L) & 0 < x < L \\ 0 & \text{elsewhere} \end{cases}$$

a) Normalization:

$$\int dx \, |\psi(x)|^2 = N^2 \int_0^L dx \, x^2 (x - L)^2 = \frac{N^2 L^5}{30} = 1$$

$$N = \sqrt{\frac{30}{L^5}}$$

$$\langle E_1 | \psi \rangle = \int dx \ \langle E_1 | x \rangle \langle x | \psi \rangle = \int_0^L dx \sqrt{\frac{2}{L}} \sin\left(\frac{\pi x}{L}\right) Nx(x - L) = \sqrt{\frac{60}{L^6}} \left(-\frac{4L^3}{\pi^3}\right) = -\frac{8\sqrt{15}}{\pi^3}$$
$$|\langle E_1 | \psi \rangle|^2 = \frac{960}{\pi^6} = 0.998555$$

c)
$$\langle E \rangle = \sum_{n} |\langle E_n | \psi \rangle|^2 E_n = \left(\frac{960}{\pi^6}\right) \frac{\hbar^2 \pi^2}{2mL^2} + \left(\frac{320}{243\pi^6}\right) \frac{9\hbar^2 \pi^2}{2mL^2} + \dots \approx \frac{13120\hbar^2}{27\pi^4 mL^2}$$

Problem 6.15 - SKIPPED

$$\frac{2m}{\hbar^2}V(x) = -\frac{\lambda}{b}\delta(x)$$

$$\left(\frac{d\psi}{dx}\right)_{x+\epsilon} - \left(\frac{d\psi}{dx}\right)_{x-\epsilon} = \int_{x-\epsilon}^{x+\epsilon} dx \, \frac{d}{dx} \frac{d\psi}{dx} = \int_{x-\epsilon}^{x+\epsilon} dx \, \frac{2m}{\hbar^2} (V - E)\psi$$

$$x = 0$$
:

$$\begin{split} \left(\frac{d\psi}{dx}\right)_{0^{+}} - \left(\frac{d\psi}{dx}\right)_{0^{-}} &= \lim_{\epsilon \to 0} \int_{0-\epsilon}^{0+\epsilon} dx \, \frac{d}{dx} \frac{d\psi}{dx} = \lim_{\epsilon \to 0} \int_{0-\epsilon}^{0+\epsilon} dx \, \frac{2m}{\hbar^{2}} \left[-\frac{\lambda\hbar^{2}}{2bm} \delta(x) - E \right] \psi \\ &= \lim_{\epsilon \to 0} \int_{0-\epsilon}^{0+\epsilon} dx \, \left[-\frac{\lambda}{b} \delta(x) - \frac{2mE}{\hbar^{2}} \right] \psi = -\frac{\lambda}{b} \psi(0) \end{split}$$

b)

$$k = \sqrt{\frac{2mE}{\hbar^2}}$$

E < 0:

$$\frac{d^2\psi}{dx^2} = -k^2\psi$$

 $x < 0 \ (B = 0)$:

$$\psi = Ae^{ikx}$$

$$\psi' = ikAe^{ikx}$$

 $x > 0 \ (C = 0)$:

$$\varphi = De^{-ikx}$$
$$\varphi' = -ikDe^{-ikx}$$

Boundary conditions: $\psi(0) = \varphi(0)$

$$A = D$$

$$\varphi'(0) - \psi'(0) = -\frac{\lambda}{h}\varphi(0) = -\frac{\lambda}{h}\psi(0)$$

$$-ikD - ikA = -2ikA = -\frac{\lambda}{h}A$$

$$2ik = \frac{\lambda}{b}$$

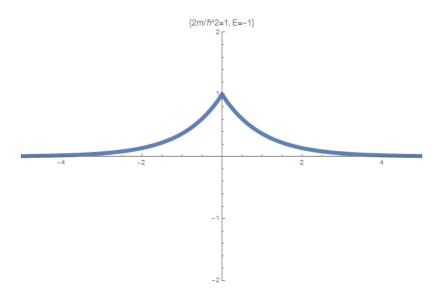
$$2\sqrt{\frac{2m|E|}{\hbar^2}} = \frac{\lambda}{b}$$

$$E_e = -\frac{\hbar^2 \lambda^2}{8mb^2}$$

Normalization:

$$A = \left(\frac{2m|E|}{\hbar^2}\right)^{1/4}$$

Plot of $\psi(x)$ for $\frac{2m}{\hbar^2} = 1, E = -1$



$$\frac{2m}{\hbar^2}V(x) = -\frac{\lambda}{b}\delta(x)$$

$$k = \sqrt{\frac{2mE}{\hbar^2}}$$

E > 0:

$$\frac{d^2\psi}{dx^2} = -k^2\psi$$

x < 0:

$$\psi = Ae^{ikx} + Be^{-ikx}$$
$$\psi' = ikAe^{ikx} - ikBe^{-ikx}$$

x > 0:

$$\varphi = Ce^{ikx} + De^{-ikx}$$
$$\varphi' = ikCe^{ikx} - ikDe^{-ikx}$$

Boundary conditions: $\psi(0) = \varphi(0)$

$$A + B = C + D$$

$$\varphi'(0) - \psi'(0) = -\frac{\lambda}{b}\varphi(0) = -\frac{\lambda}{b}\psi(0)$$

$$(ikC - ikD) - (ikA - ikB) = -\frac{\lambda}{b}(A + B)$$

Scattering:

$$D = 0$$
$$A + B = C$$

$$ikC - ikA + ikB = -\frac{\lambda}{b}(A+B)$$

$$ikC - ikA + ik(C-A) = -\frac{\lambda}{b}C$$

$$\left(2ik + \frac{\lambda}{b}\right)C = 2ikA$$

$$T = \frac{|C|^2}{|A|^2} = \left(\frac{2ik}{2ik + \frac{\lambda}{b}}\right)\left(\frac{-2ik}{-2ik + \frac{\lambda}{b}}\right) = \frac{4k^2}{4k^2 + (\lambda/b)^2}$$

$$ik(A+B) - ikA + ikB = -\frac{\lambda}{b}(A+B)$$

$$\left(2ik + \frac{\lambda}{b}\right)B = -\frac{\lambda}{b}A$$

$$R = \frac{|B|^2}{|A|^2} = \left(\frac{-\lambda/b}{2ik + \lambda/b}\right)\left(\frac{-\lambda/b}{-2ik + \lambda/b}\right) = \frac{(\lambda/b)^2}{4k^2 + (\lambda/b)^2}$$

$$T + R = 1$$

Probability current:

$$j = \frac{\hbar}{2mi} \left(\psi^* \frac{\partial \psi}{\partial x} - \psi \frac{\partial \psi^*}{\partial x} \right)$$

x < 0:

$$\psi^* = Ae^{-ikx} + Be^{ikx}$$

$$j = \frac{\hbar}{2mi} \left[\left(Ae^{-ikx} + Be^{ikx} \right) (ikAe^{ikx} - ikBe^{-ikx}) - (Ae^{ikx} + Be^{-ikx}) (-ikAe^{-ikx} + ikBe^{ikx}) \right]$$

$$j = \frac{\hbar}{2mi} (2ik|A|^2 - 2ik|B|^2) = \frac{\hbar k}{m} (|A|^2 - |B|^2)$$

$$j_{inc} = \frac{\hbar k}{m} |A|^2$$

$$j_{ref} = \frac{\hbar k}{m} |B|^2$$

$$R = \frac{j_{ref}}{j_{inc}} = \frac{|B|^2}{|A|^2}$$

$$x > 0 \ (D = 0):$$

$$\psi^* = Ce^{-ikx}$$

$$j = \frac{\hbar}{2mi} [Ce^{-ikx}(ikCe^{ikx}) - Ce^{ikx}(-ikCe^{-ikx})] = \frac{\hbar k}{m} (|C|^2)$$
$$j_{trans} = \frac{\hbar k}{m} |C|^2$$
$$T = \frac{j_{trans}}{j_{inc}} = \frac{|C|^2}{|A|^2}$$

 $E > V_0$

$$k\sqrt{\frac{2mE}{\hbar^2}}$$

$$Q = \sqrt{\frac{2m(E - V_0)}{\hbar^2}}$$

x < 0:

$$\psi(x) = Ae^{ikx} + Be^{-ikx}$$

x > 0:

$$\psi(x) = Ce^{iQx} + De^{-iQx}$$

Boundary conditions:

$$A + B = C + D$$
$$ik(A - B) = iQ(C - D)$$

Particles are incident from the left:

$$R = \frac{(k-Q)^2}{(k+Q)^2}$$
$$T = \frac{4kQ}{(k+Q)^2}$$

Particles are incident from the right:

$$B = C + D$$
$$-ikB = iQ(C - D)$$
$$-ik(C + D) = iQ(C - D)$$
$$(-k + Q)D = (k + Q)C$$
$$R = \frac{|C|^2}{|D|^2} = \frac{(k - Q)^2}{(k + Q)^2}$$

$$-ikB = iQ(B - 2D)$$

$$(k + Q)B = (2Q)D$$

$$T = \frac{k}{Q} \frac{|B|^2}{|D|^2} = \frac{k}{Q} \frac{4Q^2}{(k + Q)^2} = \frac{4kQ}{(k + Q)^2}$$

Square potential barrier:

$$V = \begin{cases} 0 & x < 0 \\ V_0 & 0 < x < a \\ 0 & x > a \end{cases}$$

$$k = \sqrt{\frac{2mE}{\hbar^2}}$$

$$q = \sqrt{\frac{2m(V_0 - E)}{\hbar^2}}$$

$$\psi = \begin{cases} Ae^{ikx} + Be^{-ikx} & x < 0 \\ Fe^{qx} + Ge^{-qx} & 0 < x < a \\ Ce^{ikx} & x > a \end{cases}$$

Boundary conditions:

$$A + B = F + G$$

$$ik(A - B) = q(F - G)$$

$$Fe^{qa} + Ge^{-qa} = Ce^{ika}$$

$$q(Fe^{qa} - Ge^{-qa}) = ikCe^{ika}$$

$$F = \frac{1}{2}\left(1 + \frac{ik}{q}\right)A + \frac{1}{2}\left(1 - \frac{ik}{q}\right)B$$

$$G = \frac{1}{2}\left(1 - \frac{ik}{q}\right)A + \frac{1}{2}\left(1 + \frac{ik}{q}\right)B$$

$$2Fe^{qa} = \left(1 + \frac{ik}{q}\right)Ce^{ika}$$

$$2Ge^{-qa} = \left(1 - \frac{ik}{q}\right)Ce^{ika}$$

$$\left(1 + \frac{ik}{q}\right)Ae^{qa} + \left(1 - \frac{ik}{q}\right)Be^{qa} = \left(1 + \frac{ik}{q}\right)Ce^{ika}$$

$$\left(1 - \frac{ik}{q}\right)Ae^{-qa} + \left(1 + \frac{ik}{q}\right)Be^{-qa} = \left(1 - \frac{ik}{q}\right)Ce^{ika}$$

$$\begin{split} \left(1 - \frac{ik}{q}\right) A e^{-qa} + \left(\frac{1 + ik/q}{1 - ik/q}\right) \left[\left(1 + \frac{ik}{q}\right) C e^{ika} - \left(1 + \frac{ik}{q}\right) A e^{qa}\right] e^{-2qa} &= \left(1 - \frac{ik}{q}\right) C e^{ika} \\ \left(1 - \frac{ik}{q}\right)^2 A e^{-qa} + \left[\left(1 + \frac{ik}{q}\right)^2 C e^{(ik-2q)a} - \left(1 + \frac{ik}{q}\right)^2 A e^{-qa}\right] &= \left(1 - \frac{ik}{q}\right)^2 C e^{ika} \\ \left(\frac{4ik}{q}\right) A e^{-qa} &= \left[\left(1 + \frac{ik}{q}\right)^2 e^{-2qa} - \left(1 - \frac{ik}{q}\right)^2\right] C e^{ika} \\ \frac{C}{A} &= \frac{\left(\frac{4ik}{q}\right) e^{-(q+ik)a}}{\left[\left(1 + \frac{ik}{q}\right)^2 e^{-2qa} - \left(1 - \frac{ik}{q}\right)^2\right]} \\ T &= \frac{j_{x>a}}{j_{inc}} &= \frac{|C|^2}{|A|^2} &= \frac{\left(\frac{4ik}{q}\right) e^{-(q+ik)a}}{\left[\left(1 + \frac{ik}{q}\right)^2 e^{-2qa} - \left(1 - \frac{ik}{q}\right)^2\right]} \left[\left(1 - \frac{ik}{q}\right)^2 e^{-2qa} - \left(1 + \frac{ik}{q}\right)^2\right] \\ &= \frac{\left(4k/q\right)^2 e^{-2qa}}{\left(1 + \frac{k^2}{q^2}\right)^2 e^{-4qa} - \left[\left(1 + \frac{ik}{q}\right)^4 + \left(1 - \frac{ik}{q}\right)^4\right] e^{-2qa} + \left(1 + \frac{k^2}{q^2}\right)^2} \\ &= \frac{\left(4k/q\right)^2}{\left(1 + \frac{k^2}{q^2}\right)^2 \left(e^{2qa} + e^{-2qa}\right) - 2\left(1 - \frac{6k^2}{q^2} + \frac{k^4}{q^4}\right)} \\ &= \frac{(4k/q)^2}{\left(1 + \frac{k^2}{q^2}\right)^2 \left(4\sinh^2 qa + 2\right) - 2\left(1 - \frac{6k^2}{q^2} + \frac{k^4}{q^4}\right)} \\ &= \frac{\left(4k/q\right)^2}{4\left(1 + \frac{k^2}{q^2}\right)^2 \sinh^2(qa) + \left(16k^2/q^2\right)} = \frac{1}{\left(\frac{q}{2k} + \frac{k}{2q}\right)^2 \sinh^2(qa) + 1} \\ T &= \frac{1}{1 + \left(\frac{k^2+q^2}{q^2}\right)^2 \sinh^2(qa)} \end{split}$$

$$V = \begin{cases} 0 & x < 0 \\ -V_0 & 0 < x < a \\ 0 & x > a \end{cases}$$

$$q' \to i\sqrt{\frac{2m(E+V_0)}{\hbar^2}}$$

$$T = \frac{1}{1 + \left(\frac{k^2 + q'^2}{2kq'}\right)^2 \sinh^2(q'a)}$$

$$\sinh(ix) = i\sin(x)$$

$$\sinh^2(ix) = -\sin^2(x)$$

$$\left(\frac{k^2 + q'^2}{2kq'}\right)^2 = \left[\frac{\frac{2mE}{\hbar^2} - \frac{2m(E+V_0)}{\hbar^2}}{2i\sqrt{\frac{2mE}{\hbar^2}}\sqrt{\frac{2m(E+V_0)}{\hbar^2}}}\right]^2 = -\frac{(2m/\hbar^2)^2(V_0)^2}{4(2m/\hbar^2)^2 E(E+V_0)}$$

$$\left(\frac{k^2 + q'^2}{2kq'}\right)^2 \sinh^2(q'a) = \frac{V_0^2}{4E(E+V_0)} \sin^2\sqrt{\frac{2m(E+V_0)}{\hbar^2}}a$$

$$T = \left[1 + \frac{\sin^2\sqrt{\frac{2m}{\hbar^2}(E+V_0)}a}{4\frac{E}{V_0}\frac{(E+V_0)}{V_0}}\right]^{-1}$$

b) - SKIPPED

Problem 6.21

Approximation:

$$T \simeq \exp\left(-2\int dx \sqrt{\frac{2m}{\hbar^2}}[V(x) - E]\right)$$

$$V(x) = -e|E|x + (E_f + W)$$

$$L = \frac{W}{e|E|}$$

$$-2\int dx \sqrt{\frac{2m}{\hbar^2}}[V(x) - E] = -2\sqrt{\frac{2m}{\hbar^2}} \int_0^L dx \sqrt{-e|E|x + W} = -2\sqrt{\frac{2m}{\hbar^2}} \frac{2}{3e|E|} W^{3/2}$$

$$T \simeq \exp\left(\frac{-4\sqrt{2m}W^{3/2}}{3e|E|\hbar}\right)$$