Master's Thesis (Academic Year 2023)

Link Management For Quantum Network

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#### Abstract of Bachelor's Thesis - Academic Year 20xx

#### Link Management For Quantum Network

I can't write English.

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# Introduction

- 1.1 Background
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## Background

#### 2.1 Quantum State

#### 2.1.1 Quantum Bit

A classical bit has two different states, which are 0 and 1. Instead, those of a quantum bit (or **qubit** in short) are  $|0\rangle$  and  $|1\rangle$ , each of which can be described as a vector. For example

$$|0\rangle = \left[\begin{array}{c} 1\\0 \end{array}\right]$$

$$|1\rangle = \left[ \begin{array}{c} 0\\1 \end{array} \right]$$

The state of a single qubit  $|\psi\rangle$  can be described as follows.

$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle (\alpha, \beta \in \mathbb{C}, |\alpha|^2 + |\beta|^2 = 1)$$

. After the operation called measurement, the quantum state would be collapsed into either 0 or 1. The measurement probability of 0 is  $|\alpha|^2$  and that of 1 is  $|\beta|^2$ . In other words, a single qubit can take both states probabilistically at the same time. For instance, a qubit can be

$$|\psi\rangle = \frac{1}{\sqrt{2}}|0\rangle + \frac{1}{\sqrt{2}}|1\rangle \tag{1}$$

which can be 50% 0 and 50% 1.

#### 2.1.2 Bloch sphere

Because  $|\alpha|^2 + |\beta|^2 = 1$ , the notation of a single qubit state can be represented like this.

$$|\psi\rangle = e^{i\gamma}(\cos\frac{\theta}{2} + e^{i\phi}\sin\frac{\theta}{2})(\gamma, \phi, \theta \in \mathbb{R})$$
 (2.1)

.

Because  $e^{i\gamma}$  is just a global state, it can be ignored and the same state can be rewritten like this.

$$|\psi\rangle = \cos\frac{\theta}{2} + e^{i\phi}\sin\frac{\theta}{2}(\phi, \theta \in \mathbb{C})$$
 (2.2)

Because the equation above has two parameters, any pure single qubit state can be considered as a point on the surface and its geometric representation is called **Bloch sphere**.

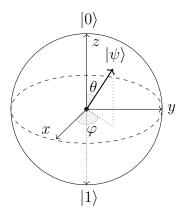


Figure 2.1: Bloch Sphere

#### 2.1.3 Multi-Qubit State

The quantum state for multi-qubits is a **tensor product** of a state vector of each qubit. The general notation of two qubit state is

$$|\psi\rangle = (\alpha|0\rangle + \beta|1\rangle) \otimes (\gamma|0\rangle + \delta|1\rangle) \tag{2.3}$$

$$= \alpha \gamma |00\rangle + \alpha \delta |01\rangle + \beta \gamma |10\rangle + \beta \delta |11\rangle \tag{2.4}$$

$$(\alpha, \beta, \gamma, \delta \in \mathbb{C}, |\alpha|^2 + |\beta|^2 + |\gamma|^2 + |\delta|^2 = 1) \tag{2.5}$$

For example, the state  $|00\rangle$  is equal to

$$\begin{bmatrix} 1 \\ 0 \end{bmatrix} \otimes \begin{bmatrix} 1 \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$0$$
(2.6)

However, some quantum states such as

$$|\psi\rangle = \frac{1}{\sqrt{2}}|00\rangle + \frac{1}{\sqrt{2}}|11\rangle \tag{2.7}$$

cannot be decomposed into quantum state of each qubit. These special quantum states are called **entangled** states.

#### 2.2 Quantum Operations

#### 2.2.1 I gate

I gate is equal to the 2x2 identity matrix, which is

$$I = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \tag{2.8}$$

For example,

$$I|0\rangle = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \end{bmatrix} = |0\rangle$$
 (2.9)

$$I|1\rangle = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} = \begin{bmatrix} 0 \\ 1 \end{bmatrix} = |1\rangle$$
 (2.10)

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#### 2.2.2 X Gate

#### X gate

X gate flips the logical value of a qubit.

$$X = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \tag{2.11}$$

For example,

$$X|0\rangle = \begin{bmatrix} 0 & 1\\ 1 & 0 \end{bmatrix} \begin{bmatrix} 1\\ 0 \end{bmatrix} = \begin{bmatrix} 0\\ 1 \end{bmatrix} = |1\rangle \tag{2.12}$$

$$X|1\rangle = \begin{bmatrix} 0 & 1\\ 1 & 0 \end{bmatrix} \begin{bmatrix} 0\\ 1 \end{bmatrix} = \begin{bmatrix} 1\\ 0 \end{bmatrix} = |0\rangle \tag{2.13}$$

#### 2.2.3 Y gate

Y gate flips the logical value of a qubit and add an imaginary number.

$$Y = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \tag{2.14}$$

For example,

$$Y|0\rangle = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ i \end{bmatrix} = i|1\rangle \tag{2.15}$$

$$Y|1\rangle = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} = \begin{bmatrix} -i \\ 0 \end{bmatrix} = -i|0\rangle \tag{2.16}$$

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#### 2.2.4 Z Gate

Z gate flips the phase of  $|1\rangle$ 

$$Z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \tag{2.17}$$

For example,

$$Z|0\rangle = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \end{bmatrix} = |0\rangle$$
 (2.18)

$$Z|1\rangle = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} = \begin{bmatrix} 0 \\ -1 \end{bmatrix} = -|1\rangle$$
 (2.19)

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#### 2.2.5 H Gate

H gate creates superposition.

$$H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1\\ 1 & -1 \end{bmatrix} \tag{2.20}$$

For example,

$$H|0\rangle = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1\\ 1 & -1 \end{bmatrix} \begin{bmatrix} 1\\ 0 \end{bmatrix} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1\\ 1 \end{bmatrix} = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) \tag{2.21}$$

$$H|1\rangle = \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ -1 \end{bmatrix} = \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle) \tag{2.22}$$

#### 2.2.6 CNOT Gate

A CNOT gate involves two qubits, one is called **controlled qubit** and the other is called **target qubit**. If the controlled qubit is 1, the bit value of the target qubit is flipped.

$$CNOT = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}$$
 (2.23)

For example,

$$CNOT_{0,1}|10\rangle = \begin{bmatrix} 1 & 0 & 0 & 0\\ 0 & 1 & 0 & 0\\ 0 & 0 & 0 & 1\\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} 0\\0\\1\\0 \end{bmatrix} = \begin{bmatrix} 0\\0\\0\\1 \end{bmatrix} = |11\rangle$$
 (2.24)

$$CNOT_{0,1}|11\rangle = \begin{bmatrix} 1 & 0 & 0 & 0\\ 0 & 1 & 0 & 0\\ 0 & 0 & 0 & 1\\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} 0\\0\\0\\1 \end{bmatrix} = \begin{bmatrix} 0\\0\\1\\0 \end{bmatrix} = |10\rangle$$
 (2.25)

#### 2.2.7 Measurement

Quantum measurement can be described by using a group of measurement operators  $\{M_m\}$  (m is the measurement result that is expected to get). If the quantum state before measurement is  $|\psi\rangle$ , the measurement probability of value m is

$$p(m) = \langle \psi | M_m^{\dagger} M_m | \psi \rangle$$

The quantum state after the measurement is

$$\frac{M_m|\psi\rangle}{\sqrt{\langle\psi|M_m^{\dagger}M_m|\psi\rangle}}$$

The measurement operators satisfy the completeness equation

$$\sum_{m} M_{m}^{\dagger} M_{m} = I$$

Also, the sum of the measurement probability of each possible measurement outcome is equal to one.

$$\sum_{m} p(m) = \langle \psi | \sum_{m} M_{m}^{\dagger} M_{m} | \psi \rangle = 1$$

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### 4.1 Problem Definition

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## Conclusion

- 8.1 Conclusion
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# Appendix A Appendix

A.1 hoge

# ${\bf Acknowledgement}$

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