

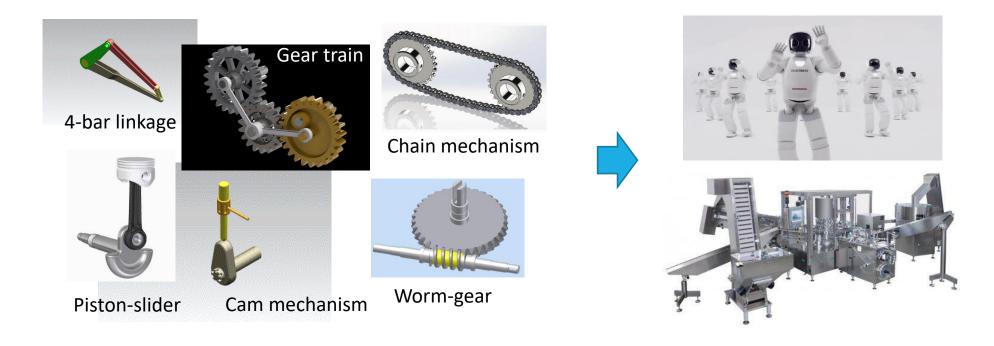
ME2115/TME2115 Mechanics of Machines

REVIEW ON BASIC MECHANICS

Study of Mechanism



A mechanical system integrates <u>power and control systems</u> with <u>basic mechanisms</u> to perform specific tasks.



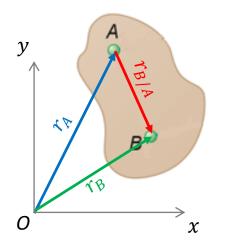
It involves either <u>analyses</u> and/or <u>synthesis</u> of the system via <u>Kinematic</u> and <u>Kinetic</u> studies. In either cases, they involves <u>vector equations</u>.

Kinematics of Rigid Body - Position



- Position vector represents the position of a point in space in relation to a given reference frame, e.g. \vec{r}_A and \vec{r}_B . For rigid body, the important points of interest are the center of mass and joints at which it is connected to other bodies.
- It is also common to determine the <u>relative position</u> vector

$$\vec{r}_{B/A} = \vec{r}_B - \vec{r}_A = -\vec{r}_{A/B}$$



- For 2D problem, the following trigonometry relations can be used to determine the geometry information, namely:
 - Sine rule:

$$\frac{\sin\theta}{a} = \frac{\sin\beta}{b} = \frac{\sin\gamma}{c}$$

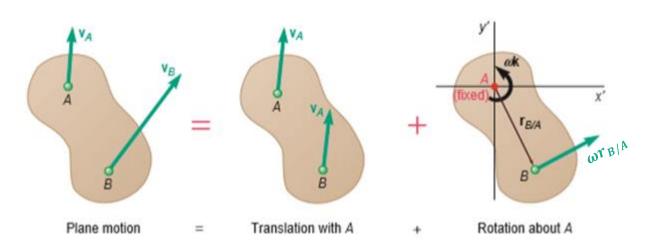
Cosine rule:

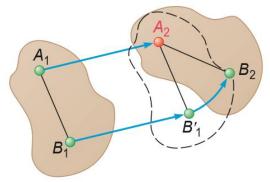
$$a^2 = b^2 + c^2 - 2bccos(\theta)$$

Kinematics of Rigid Body – Velocity



A general 2D plane motion of a rigid body comprises a translation and a rotation components.





At any given instant, the velocities of any two points A and B on the rigid body are related by the following relation:

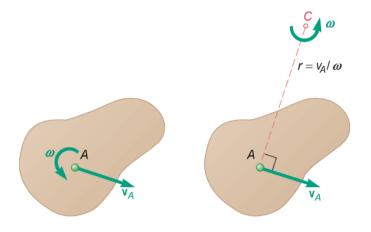
$$ec{v}_B = ec{v}_A + \overrightarrow{\omega} imes ec{r}_{B/A}$$
 Tangential to $ec{r}_{B/A}$

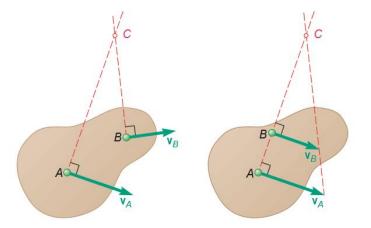
Kinematics of Rigid Body – Velocity...



- At any given instant, there exist a point at which the rigid body appears to be rotating about. This point is called the <u>instantaneous center of rotation</u> (I.C.R) in plane motion.
 - It can be determined if \vec{v}_A and ω are given.

 It can also be determined if two velocities on the rigid body were defined.





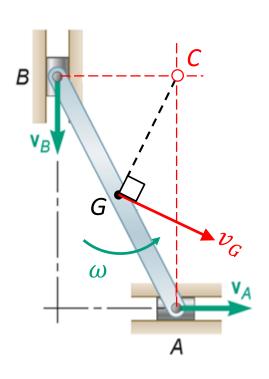
Once it is found, velocity at any other point can be found as

$$\vec{v}_B = \vec{\omega} \times \vec{r}_{B/C}$$

Kinematics of Rigid Body – Velocity...



This graphical approach provides an alternate way to solve for velocity, instead of using the vector equation. For example,



Using velocity equation,

$$\vec{v}_B = \vec{v}_A + \vec{\omega} \times \vec{r}_{B/A}$$

And then

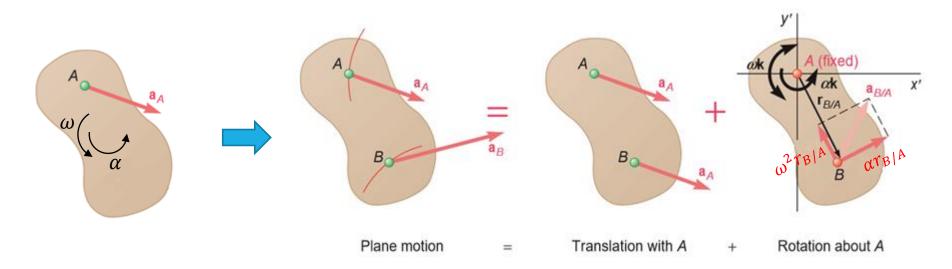
$$\vec{v}_G = \vec{v}_B + \vec{\omega} \times \vec{r}_{G/B}$$
, or $\vec{v}_G = \vec{v}_A + \vec{\omega} \times \vec{r}_{G/A}$

Based on instantaneous center of rotation

$$\vec{v}_G = \vec{\omega} \times \vec{r}_{G/C}$$

Kinematics of Rigid Body – Acceleration





At the given instant, the accelerations of any two points A and B on the rigid body are related by the following relation:

$$\vec{a}_B = \vec{a}_A + \vec{\alpha} \times \vec{r}_{B/A} + \vec{\omega} \times (\vec{\omega} \times \vec{r}_{B/A})$$
Tangential to $\vec{r}_{B/A}$ Centripetal $(-\omega^2 \vec{r}_{B/A} \ or \ \omega^2 \vec{r}_{A/B})$

Note: The concept of instantaneous center of rotation *CANNOT* be applied for acceleration.

Vector Calculus



ightharpoonup Given three vectors, $\vec{p}=p_x\vec{\imath}+p_y\vec{\jmath};\ \vec{q}=q_x\vec{\imath}+q_y\vec{\jmath};\ \ {\rm and}\ \ \vec{r}=r\vec{k}.$

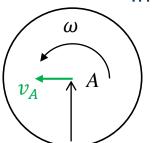
Summary of vector operations for dynamics in 2D planar motion

Vector addition	$\vec{p} + \vec{q} = \vec{q} + \vec{p} = (p_x + q_x)\vec{i} + (p_y + q_y)\vec{j}$
Scalar (Dot) product	$\vec{p}.\vec{q} = \vec{q}.\vec{p} = p_x q_x + p_y q_y$
	$\vec{p}.\vec{r} = \vec{q}.\vec{r} = 0$
Vector (Cross) product	$\vec{k} \times \vec{\iota} = \vec{j}; \ \vec{k} \times \vec{j} = -\vec{\iota}$
	$\vec{p} \times \vec{q} = -\vec{q} \times \vec{p} = (p_x q_y - p_y q_x) \vec{k}$
	$\vec{r} \times \vec{p} = -rp_y \vec{\iota} + rp_x \vec{J}$
	$\vec{p} \times \vec{p} = \vec{q} \times \vec{q} = \vec{r} \times \vec{r} = 0$

Kinematics of Rigid Body – Rolling on ground

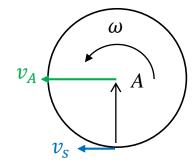


- For a <u>circular disc</u> or <u>sphere</u> rolling (with or without slipping) on a flat ground, its <u>center of</u> <u>geometric</u> is always moving parallel to the ground.
- v_A a_A
- This means that the velocity and acceleration at its geometric centre A must also be parallel to the ground.
- Suppose it is <u>rolling without sliding</u> (slipping), the point in contact with the ground is the <u>instantaneous centre of rotation</u>.



In this case,

$$\vec{v}_A = \vec{\omega} \times \vec{r}_{A/C}$$



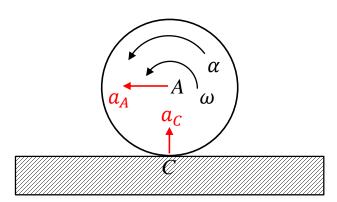
But if the body <u>slips</u> at \vec{v}_s , then

$$\vec{v}_A = \vec{v}_S + \vec{\omega} \times \vec{r}_{A/C}$$

which means that \vec{v}_s and $\vec{\omega}$ are independent quantities.

Kinematics of Rigid Body – Rolling on ground...





Likewise, in terms of acceleration, we can imposed the relation for <u>non-slip rolling</u>

$$\vec{a}_A = \vec{\alpha} \times \vec{r}_{A/C}$$

For any other point, it is given by

$$\vec{a}_B = \vec{a}_A + \vec{\alpha} \times \vec{r}_{B/A} + \vec{\omega} \times (\vec{\omega} \times \vec{r}_{B/A})$$

At the point of contact, it will experience centripetal acceleration, as given by

$$\vec{a}_C = \vec{a}_A + \vec{\alpha} \times \vec{r}_{C/A} + \vec{\omega} \times (\vec{\omega} \times \vec{r}_{C/A})$$

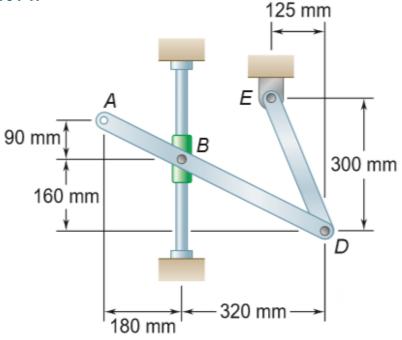
$$\vec{a}_C = (\vec{a}_A - \vec{\alpha} \times \vec{r}_{A/C}) - \omega^2 \vec{r}_{C/A} = \omega^2 \vec{r}_{A/C}$$

Example 1

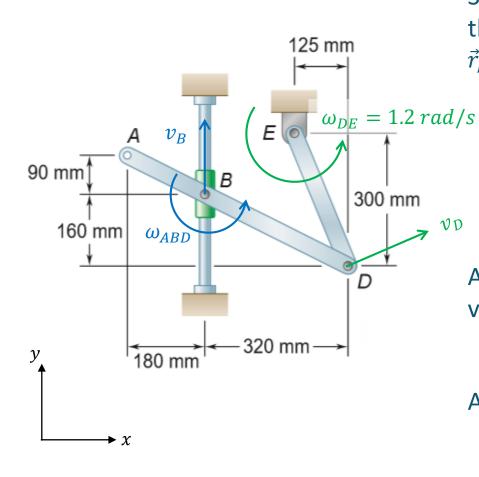


Arm *ABD* is connected by pins to a collar at *B* and to crank *DE*. Knowing that the angular velocity of crank *DE* is 1.2 *rad/s* counter-clockwise, determine at this given instant:

- (a) the angular velocity of arm ABD, and
- (b) the velocity of point A.







Since the crank DE rotates about a fixed point E, the velocity at D is expected to be tangential to $\vec{r}_{D/E}$, which is

$$\vec{v}_D = \vec{\omega}_{DE} \times \vec{r}_{D/E}$$

$$= 1.2\vec{k} \times (0.125\vec{\imath} - 0.3\vec{\jmath})$$

$$= 0.36\vec{\imath} + 0.15\vec{\jmath}$$

And the velocity at *B* is constrained to move vertically along the rod, i.e.

$$\vec{v}_B = v_B \vec{j}$$

And angular velocity of rod ABD is given by

$$\vec{\omega}_{ABD} = \omega_{ABD} \vec{k}$$



(a) Applying the velocity equation at point *D* and *B*, we have

$$\vec{v}_D = \vec{v}_B + \vec{\omega}_{ABD} \times \vec{r}_{D/B}$$

$$\Rightarrow \vec{v}_D = v_B \vec{j} + \omega_{ABD} \vec{k} \times (0.32\vec{\iota} - 0.16\vec{j})$$

$$\Rightarrow 0.36\vec{\iota} + 0.15\vec{j} = 0.16\omega_{ABD}\vec{\iota} + (v_B + 0.32\omega_{ABD})\vec{j}$$

Now, consider the *i*- and *j*- components separately, gives

$$0.36 = 0.16\omega_{ABD} \Rightarrow \omega_{ABD} = 2.25 \, rad/s$$
 $(\vec{\imath} - component)$
 $0.15 = v_B + 0.32\omega_{ABD} \Rightarrow v_B = -0.57 \, m/s$ $(\vec{\jmath} - component)$

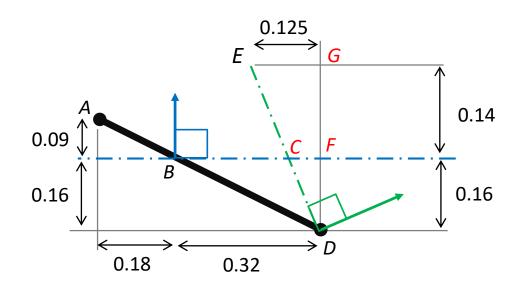
(b) Finally, the velocity of A can be determined as:

$$\vec{v}_A = \vec{v}_B + \vec{\omega}_{ABD} \times \vec{r}_{A/B} = -0.57\vec{j} + (2.25\vec{k}) \times (-0.18\vec{i} + 0.09\vec{j})$$

$$\Rightarrow \vec{v}_A = -\mathbf{0}. \, \mathbf{2025}\vec{i} - \mathbf{0}. \, \mathbf{975}\vec{j}$$



Alternatively: Based on these velocity information, we can determine the *I.C.R* for rod *ABD*, which reveals the following geometrical information.



Using similar triangles of *EGD* and *CFD*, we have

$$\frac{CF}{EG} = \frac{DF}{DG}$$

$$\Rightarrow CF = \frac{0.16}{0.30}(0.125) = 0.06667$$

Hence, position vector of point D w.r.t. C is

$$\vec{r}_{D/C} = 0.06667 \ \vec{\iota} - 0.16 \ \vec{\jmath}$$



(a) Now, consider the velocity of D from ABD, which is

$$\vec{v}_{D} = \vec{\omega}_{ABD} \times \vec{r}_{D/C}$$

$$\vec{v}_{D} = \omega_{ABD}\vec{k} \times (0.06667 \ \vec{\iota} - 0.16 \ \vec{\jmath})$$

$$\Rightarrow 0.36\vec{\iota} + 0.15\vec{\jmath} = 0.16\omega_{ABD}\vec{\iota} + 0.06667\omega_{ABD}\vec{\jmath}$$

$$\omega_{ABD} = \frac{0.36}{0.16} = 2.25 \ rad/s \ (\vec{\iota} - component), or$$

$$\omega_{ABD} = \frac{0.15}{0.06667} = 2.25 \ rad/s, \ (\vec{\jmath} - component)$$

(b) The velocity of A can be determined as:

$$\vec{v}_A = \vec{\omega}_{ABD} \times \vec{r}_{A/C}$$

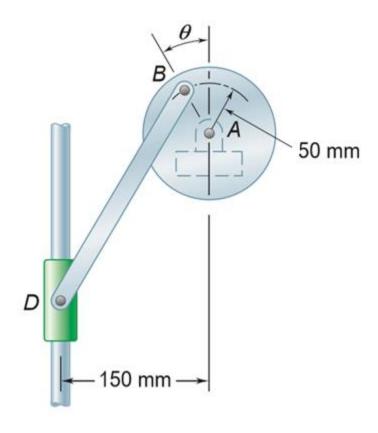
$$= (2.25\vec{k}) \times (-(0.5 - 0.06667)\vec{i} + 0.09\vec{j})$$

$$= -0.2025\vec{i} - 0.975\vec{j}$$

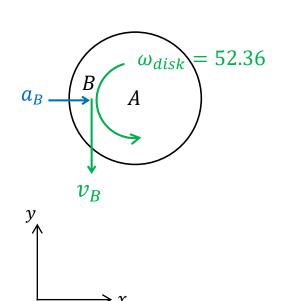
Example 2



The disk shown has a constant angular velocity of 500 rpm counter-clockwise. Knowing that rod BD is 250 mm long, determine the acceleration of collar D when $\theta = 90^{\circ}$.







Consider the disk rotating anti-clockwise at a <u>constant</u> angular velocity of 500 *rpm*, which means

$$\vec{\omega}_{disk} = 500 \left(\frac{2\pi}{60}\right) \vec{k} = 52.36 \vec{k} \, rad/s \,,$$

and

$$\vec{\alpha}_{disk} = 0 \ \vec{k}$$

Since the disk is fixed to rotate about A, this gives

$$\vec{v}_B = \vec{\omega}_{disk} \times \vec{r}_{B/A} = -\omega_{disk} r_{B/A} \vec{j}$$

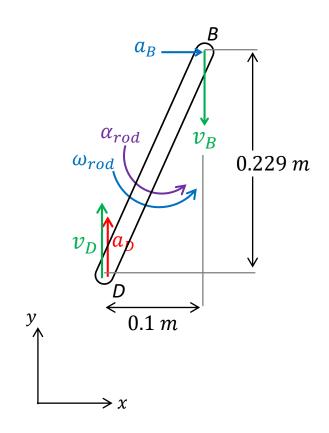
= $-(52.36)(0.05)\vec{j} = -2.618\vec{j}$

And the acceleration at B is only due to the centripetal component, i.e.

$$\vec{a}_B = \vec{\alpha}_{disk} \times \vec{r}_{B/A} + \vec{\omega}_{disk} \times (\vec{\omega}_{disk} \times \vec{r}_{B/A}) = \omega_{disk}^2 r_{B/A} \vec{\iota}$$

$$= (52.36)^2 (0.05) = 137.08 \vec{\iota}$$





Now, consider the kinematics of rod BD.

Since collar D is guided to slide along the vertical rod, hence its velocity \vec{v}_D and acceleration \vec{a}_D can be expressed as $v_D \vec{j}$ and $a_D \vec{j}$, respectively.

Based on the velocity analysis of point *B* and *D*, we have

$$\vec{v}_{D} = \vec{v}_{B} + \vec{\omega}_{rod} \times \vec{r}_{D/B}$$

$$\Rightarrow v_{D} \vec{j} = -2.618 \vec{j} + \omega_{rod} \vec{k} \times (-0.1 \vec{\iota} - 0.229 \vec{j})$$

$$= 0.229 \omega_{rod} \vec{\iota} - (2.618 + 0.229 \omega_{rod}) \vec{j}$$

From the *i*-component equation, we see that $\omega_{rod} = 0$.

Note that you can also infer this by looking at the ICR, which does not intercept and hence implies $\omega_{rod} = 0$.



Now, looking at the acceleration equation at the two points, we have

$$\vec{a}_{D} = \vec{a}_{B} + \vec{\alpha}_{rod} \times \vec{r}_{D/B} + \vec{\omega}_{rod} \times (\vec{\omega}_{rod} \times \vec{r}_{D/B})$$

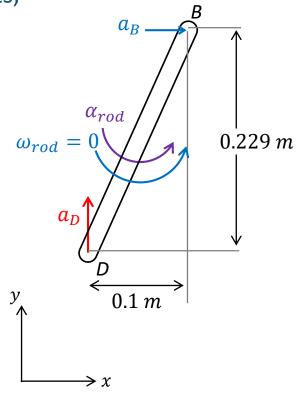
$$\Rightarrow a_{D} \vec{j} = 137.08 \vec{i} + \alpha_{rod} \vec{k} \times (-0.1 \vec{i} - 0.229 \vec{j})$$

$$= (137.08 + 0.229 \alpha_{rod}) \vec{i} - 0.1 \alpha_{rod} \vec{j}$$

Again, by looking at the i-component equation, we get

$$137.08 + 0.229\alpha_{rod} = 0$$

 $\Rightarrow \alpha_{rod} = -598.6 \ rad/s^2$

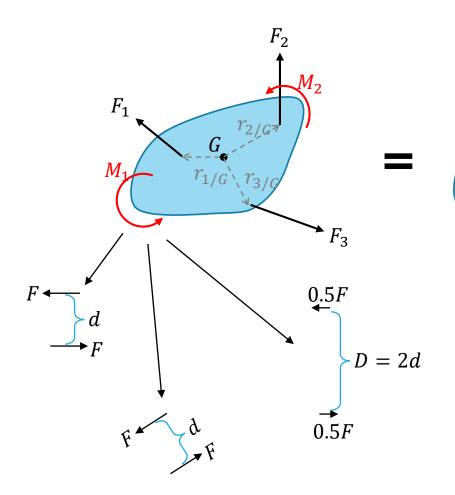


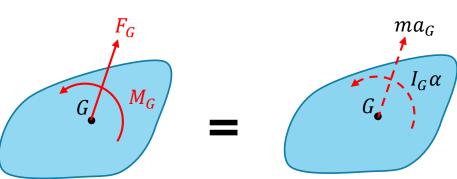
And finally substituting α into the j- component equation, gives

$$a_D = -0.1\alpha_{rod} = -0.1(-598.6) = 59.9 \text{ m/s}^2$$

Kinetics of Rigid Body - Newton's 2nd Law







Force equation

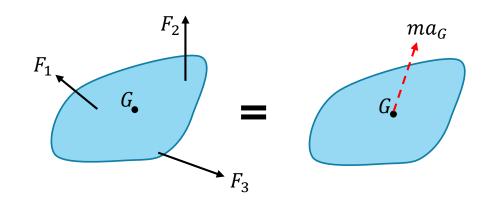
$$\sum \vec{F}_i = \vec{F}_G = m\vec{a}_G$$

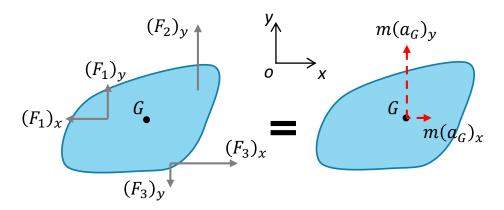
Moment equation about center of mass

$$\sum M_i + \sum \vec{r}_{i/G} \times \vec{F}_i = M_G = I_G \alpha$$

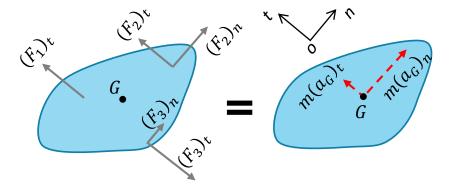
Newton's 2nd Law – Force equation







$$\sum (\vec{F}_i)_x = m(\vec{a}_G)_x$$
, and $\sum (\vec{F}_i)_y = m(\vec{a}_G)_y$

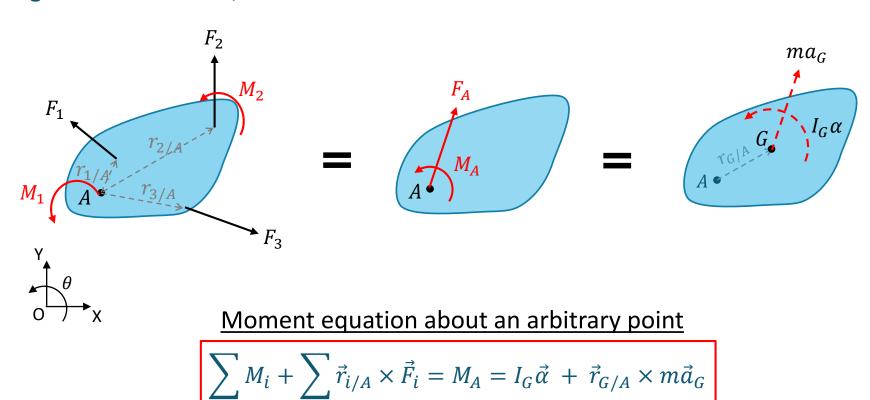


$$\sum (\vec{F}_i)_n = m(\vec{a}_G)_n, \text{ and } \sum (\vec{F}_i)_t = m(\vec{a}_G)_t$$

Newton's 2nd Law – Moment equation



Taking moment about A,

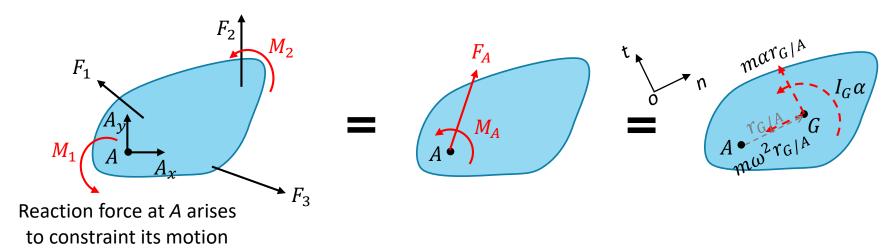


Note that $I_G \vec{\alpha}$ is always presented on the RHS, which is independent on where the moment is taken about.

Newton's 2nd Law – Moment equation...



Suppose the body rotates about a <u>fixed axis</u> at A, then the acceleration at the mass centre \vec{a}_G can be expressed in the normal and tangent coordinate.



Moment equation about a fixed point

$$M_{A} = I_{G}\vec{\alpha} + \vec{r}_{G/A} \times m\vec{a}_{t} = I_{G}\vec{\alpha} + \left(r_{G/A}\vec{i'}\right) \times \left(m\alpha r_{G/A}\vec{j'}\right)$$

$$M_{A} = \left(I_{G} + mr_{G/A}^{2}\right)\vec{\alpha} = I_{A}\vec{\alpha}$$

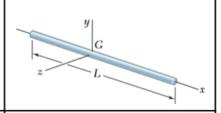
Moment of Inertia



 \triangleright <u>Centriodal Moment of inertia</u> I_G measures the extent to which an object resists rotational acceleration about its *centre of mass (C.M)*.

Slender rod,

$$I_y = I_z = \frac{1}{12}mL^2$$

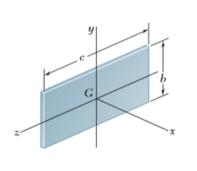


Thin rectangular plate,

$$I_{x} = \frac{1}{12}m(b^{2} + c^{2})$$

$$I_{y} = \frac{1}{12}mc^{2}$$

$$I_{z} = \frac{1}{12}mb^{2}$$



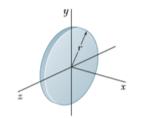
In general, it can be expressed as,

$$I_G = mk^2$$

where k is the <u>radius of gyration</u>.

$$I_x = \frac{1}{2}mr^2$$

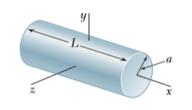
$$I_y = I_z = \frac{1}{4}mr^2$$



Circular cylinder,

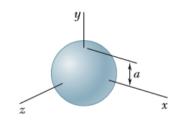
$$I_x = \frac{1}{2}ma^2$$

$$I_y = I_z = \frac{1}{12}m(3a^2 + L^2)$$



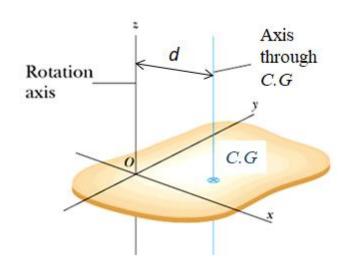
Sphere,

$$I_x = I_y = I_z = \frac{2}{5}ma^2$$



Moment of inertia – Parallel-axis theorem





Parallel-axis theorem is given by

$$I_O = I_G + md^2$$

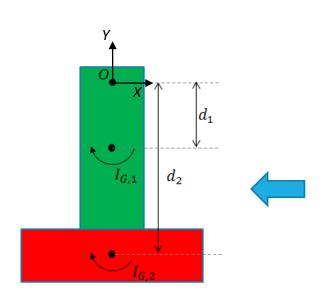
where *d* is the perpendicular distance between the two axes.

Parallel-axis theorem can be used to calculate the moment of inertia of more complex bodies that are made up of simple geometries with known *C.M.s* and their respective centroidal moment of inertia.



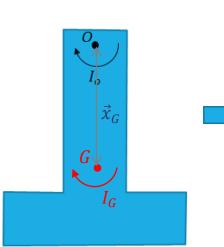
Moment of inertia - Parallel-axis theorem...

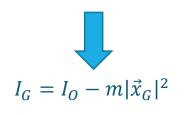


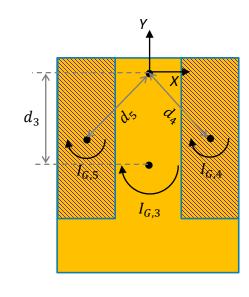


$$\vec{x}_G = \frac{M_1 \vec{x}_{G,1} + M_2 \vec{x}_{G,2}}{M_1 + M_2}$$

$$I_O = (I_{G,1} + m_1 d_1^2) + (I_{G,2} + m_2 d_2^2)$$







$$\vec{x}_G = \frac{M_3 \vec{x}_{G,3} - M_4 \vec{x}_{G,4} - M_5 \vec{x}_{G,5}}{M_3 - M_4 - M_5}$$

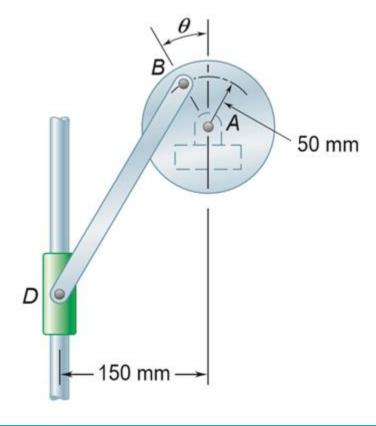
$$I_O = (I_{G,3} + m_3 d_3^2) - (I_{G,4} + m_4 d_4^2)$$

$$-(I_{G,5} + m_5 d_5^2)$$

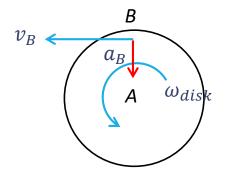
Example 3

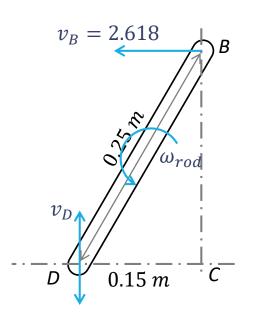


The 250-mm uniform rod BD, of mass 5 kg, is connected as shown to disk A and to a collar of negligible mass, that may slide freely along a vertical rod. Knowing that disk A rotates counter-clockwise at a constant rate of 500 rpm, determine the reactions at D when $\theta = 0$.









Kinematics analysis of the disk.

Given that $\omega_{disk} = 500 \ rpm = 52.36 \ rad/s$, and $\alpha_{disk} = 0$, we get

$$\vec{v}_B = \vec{\omega}_{disk} \times \vec{r}_{B/A} = 52.36\vec{k} \times 0.05\vec{j} = -2.618\vec{i}$$
$$\vec{a}_B = -(\omega_{disk})^2 \vec{r}_{B/A} = -52.36^2 (0.05\vec{j}) = -137.08\vec{j}$$

Velocity analysis of the rod, using I.C.R, gives

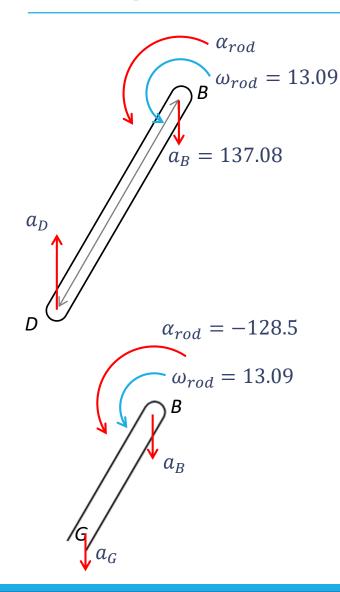
$$BC = \sqrt{0.25^2 - 0.15^2} = 0.20 m$$

$$\vec{v}_B = \vec{\omega}_{rod} \times \vec{r}_{B/C}$$

$$\Rightarrow -2.618\vec{\iota} = \omega_{rod}\vec{k} \times 0.20\vec{\jmath}$$

$$\Rightarrow \omega_{rod} = 13.09 \ rad/s$$





Acceleration analysis of the rod,

$$\vec{a}_{D} = \vec{a}_{B} + (\vec{\alpha}_{rod} \times \vec{r}_{D/B}) - \omega_{rod}^{2} \vec{r}_{D/B}$$

$$\Rightarrow a_{D} \vec{j} = -137.08 \vec{j} + \alpha_{rod} \vec{k} \times (-0.15 \vec{i} - 0.20 \vec{j})$$

$$-(13.09)^{2} (-0.15 \vec{i} - 0.20 \vec{j})$$

$$\Rightarrow a_{D} \vec{j} = (0.2\alpha_{rod} + 25.7) \vec{i} + (-102.8 - 0.15\alpha_{rod}) \vec{j}$$

$$0.2\alpha_{rod} + 25.7 = 0 \Rightarrow \alpha_{rod} = -128.5 \ rad/s^{2}$$

The acceleration at the mass centre of the rod G is

$$\vec{a}_{G} = \vec{a}_{B} + (\vec{\alpha}_{rod} \times \vec{r}_{G/B}) - \omega_{rod}^{2} \vec{r}_{G/B}$$

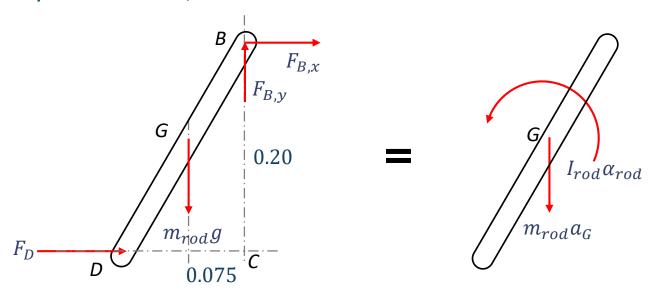
$$= -137.08\vec{j} - 128.5\vec{k} \times (-0.075\vec{i} - 0.10\vec{j})$$

$$-(13.09)^{2} (-0.075\vec{i} - 0.10\vec{j})$$

$$\Rightarrow \vec{a}_{G} = -110.3\vec{j}$$



Kinetic analysis of the rod,



From the FBD, and taking moment about B, we get

$$F_D(BC) + m_{rod}g(0.075) = I_{rod}\alpha_{rod} + m_{rod}a_G(0.075)$$

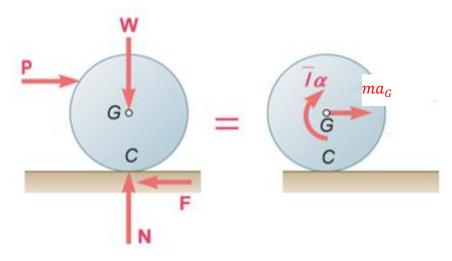
$$\Rightarrow 0.2F_D + 5(9.81)(0.075) = \left(\frac{1}{12}(5)(0.25)^2\right)(-128.5) + (5)(110.3)(0.075)$$

$$\Rightarrow F_D = \mathbf{172} N$$

Kinetics of Rolling Body



For a circular body in rolling motion,



Rolling without sliding:

$$F \leq \mu_S N$$
, and $\vec{a}_G = \vec{\alpha} \times \vec{r}_{G/C}$

Rotating and sliding:

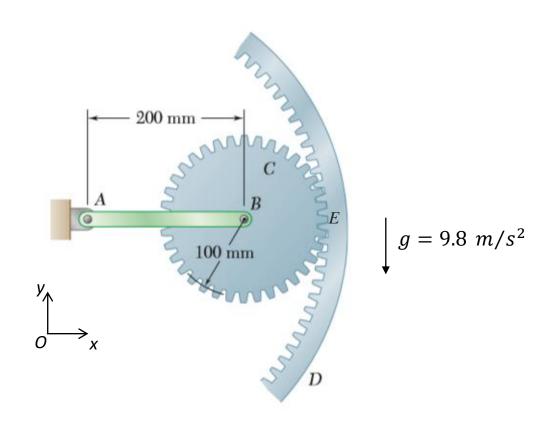
$$F = \mu_k N$$
, but \vec{a}_G and $\vec{\alpha}$ are independent.

where μ_{s} and μ_{k} are static and kinetic friction coefficients, respectively.

Example 4



Gear *C* has a mass of 5 *kg* and a centroidal radius of gyration of 75 *mm*. The uniform bar *AB* has a mass of 3 *kg* and gear *D* is stationary. If the system is released from rest in the position shown, determine the angular acceleration of gear *C*.



and



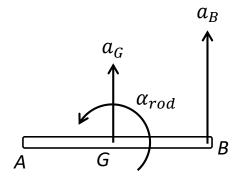
Consider the kinematics of the system.

Acceleration analysis of the bar AB.

$$\vec{a}_B = \vec{\alpha}_{rod} \times \vec{r}_{B/A}$$

$$\Rightarrow \vec{a}_B = \left(\alpha_{rod}\vec{k}\right) \times (0.2\vec{\iota}) = 0.2\alpha_{rod}\vec{\jmath}$$

$$\vec{a}_G = \left(\alpha_{rod}\vec{k}\right) \times (0.1\vec{\iota}) = 0.1\alpha_{rod}\vec{\jmath}$$

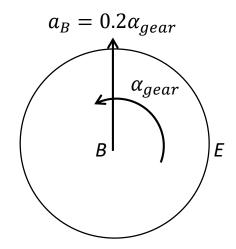


Acceleration analysis of the gear C.

$$\vec{a}_E = \vec{a}_B + (\vec{\alpha}_{gear} \times \vec{r}_{E/B})$$

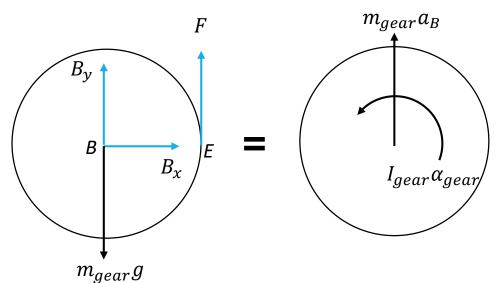
$$\Rightarrow 0 = 0.2\alpha_{rod} \vec{j} + (\alpha_{gear} \vec{k}) \times (0.1 \vec{i})$$

$$\Rightarrow \alpha_{rod} = -\frac{1}{2}\alpha_{gear}$$





Kinetic analysis of gear C.



Force in *x*-direction,

$$B_{\chi}=0$$

Force in *y*-direction,

$$B_y + F - m_{gear}g = m_{gear}a_B$$

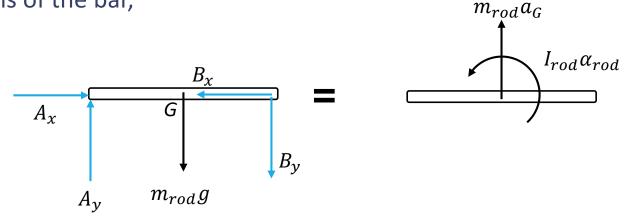
$$\Rightarrow B_y + F = m_{gear}(g + 0.2\alpha_{rod})$$

Taking moment about *B*, gives
$$F(BE) = I_{gear}\alpha_{gear} \Rightarrow F = \frac{(5)(0.075)^2}{0.1}\alpha_{gear}$$

$$\Rightarrow F = 0.28125\alpha_{gear}$$



Kinetics analysis of the bar,



Force in *x*-direction,

$$A_x - B_x = 0 \Rightarrow A_x = 0$$

Force in *y*-direction,

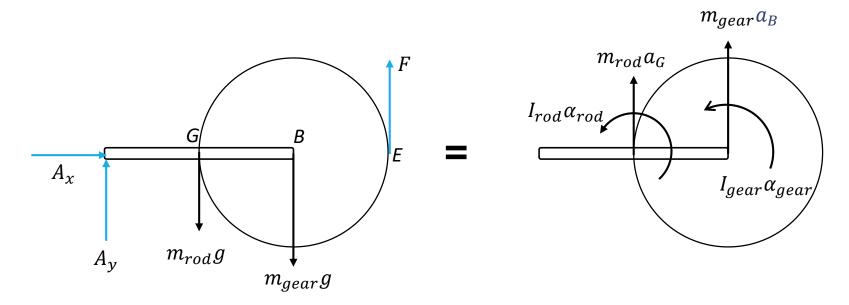
$$A_y - B_y - m_{rob}g = m_{rod}a_G$$

$$\Rightarrow A_y - B_y = m_{rod}(g + 0.1\alpha_{rod})$$

Moment equation would involve reaction force at either A and/or B.



Kinetics analysis of the bar and gear sub-system,



Taking the moment about A,

$$F(AE) - m_{rod}g(AG) - m_{gear}g(AB) = (m_{rod}a_G)(AG) + I_{rod}\alpha_{rod} + (m_{gear}a_B)(AB) + I_{gear}\alpha_{gear}$$



Using the numeric and further evaluating it, gives

$$0.3F - (3)(9.81)(0.1) - (5)(9.81)(0.2) = 3(0.1\alpha_{rod})(0.1) + \left(\frac{1}{12}(3)(0.2)^2\right)\alpha_{rod}$$
$$+ 5(0.2\alpha_{rod})(0.2) + \left((5)(0.075)^2\right)\alpha_{gear}$$
$$\Rightarrow 0.3F - 12.753 = 0.24\left(-\frac{1}{2}\alpha_{gear}\right) + 0.028125\alpha_{gear}$$
$$\Rightarrow 0.3F - 12.753 = -0.09188\alpha_{gear}$$

Finally, substituting $F = 0.28125\alpha_{gear}$ into the above equation, gives

$$0.3(0.28125\alpha_{gear}) - 12.753 = -0.09188\alpha_{gear}$$

 $\Rightarrow \alpha_{gear} = 72.36 \ rad/s^2$