

Instrumentation at Synchrotron Radiation Beamlines

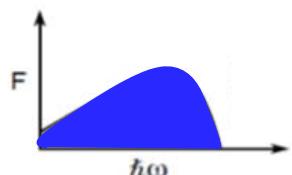
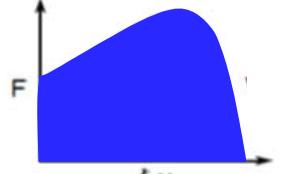
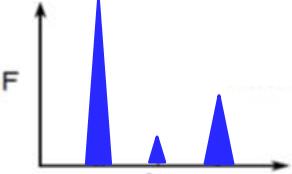
M. Altissimo, L. Raimondi

Elettra Sincrotrone Trieste SCpA
S.S. 14, km163.5, Basovizza (TS)



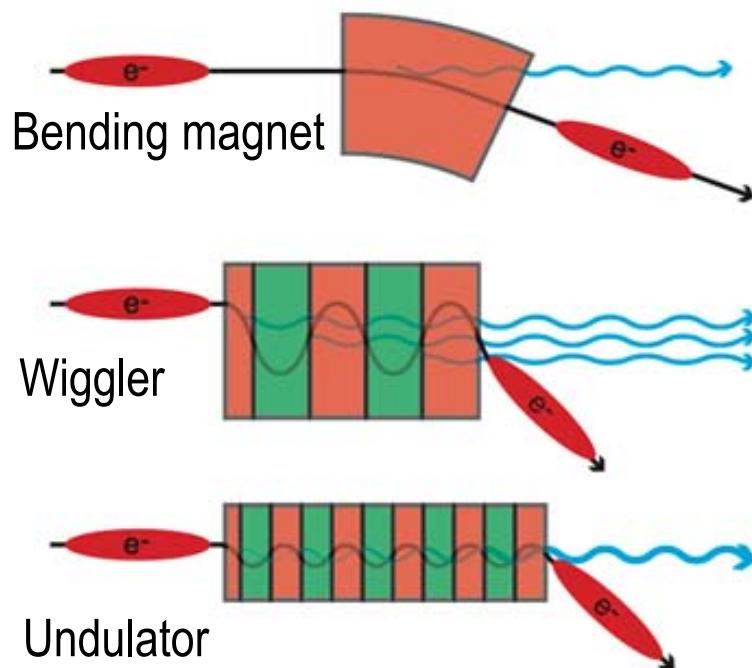
Synchrotron beam emitted by source

$$\gamma = 1957 E_e [\text{GeV}]$$

| Source | Spectrum | Divergence σ | Typically |
|----------------|--|----------------------------------|-----------------------------------|
| Bending magnet |  | $\sim \frac{I}{\gamma}$ | $\sim \text{mrad}$ |
| Wiggler |  | $\sim \frac{I}{\gamma}$ | $\sim \text{mrad}$ |
| Undulator |  | $\sim \frac{I}{\gamma \sqrt{N}}$ | $\sim \text{100's }\mu\text{rad}$ |

What does it mean in terms of beamsize?

Source



Beam size @ 20 m

~ 1 - 10's mm

~ 1 - 10's mm

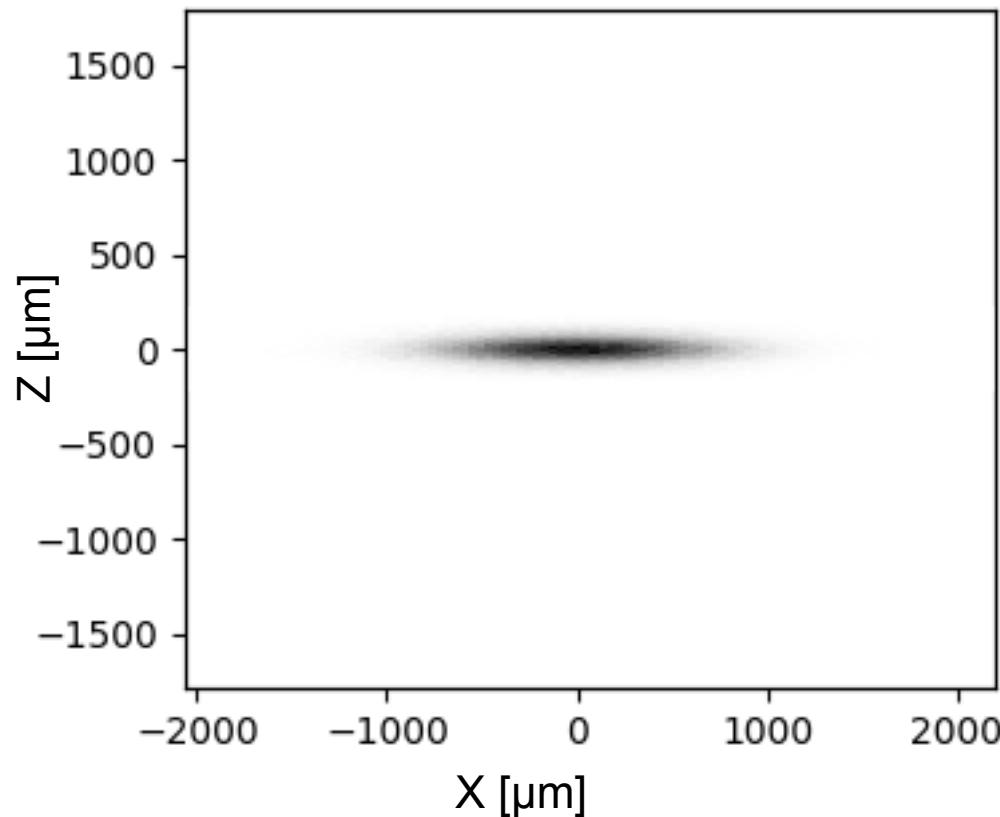
~ mm

For a typical experiment: required beam size ~ 0.1 to 10's of μm or more

A couple of undulator simulations (3rd generation storage ring)

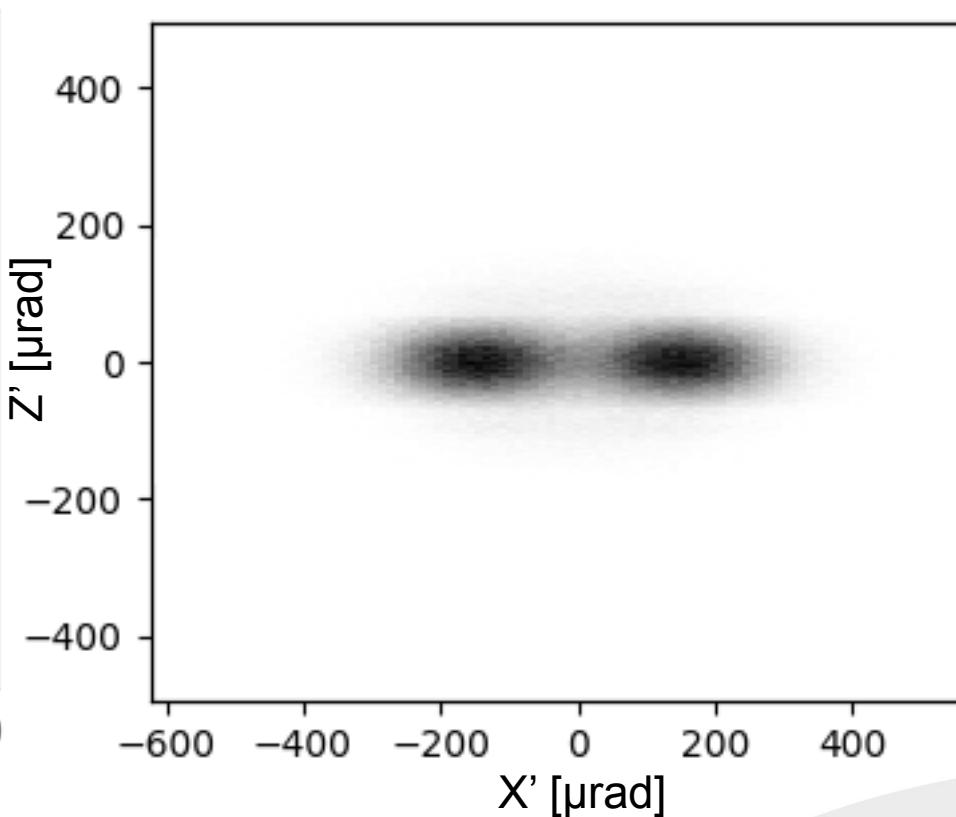
$E_e=2.4\text{GeV}$, $N = 17$, period = 56mm, first harmonic only

Source Size



FWHM: $1000 \times 100 \mu\text{m}^2$

Divergence

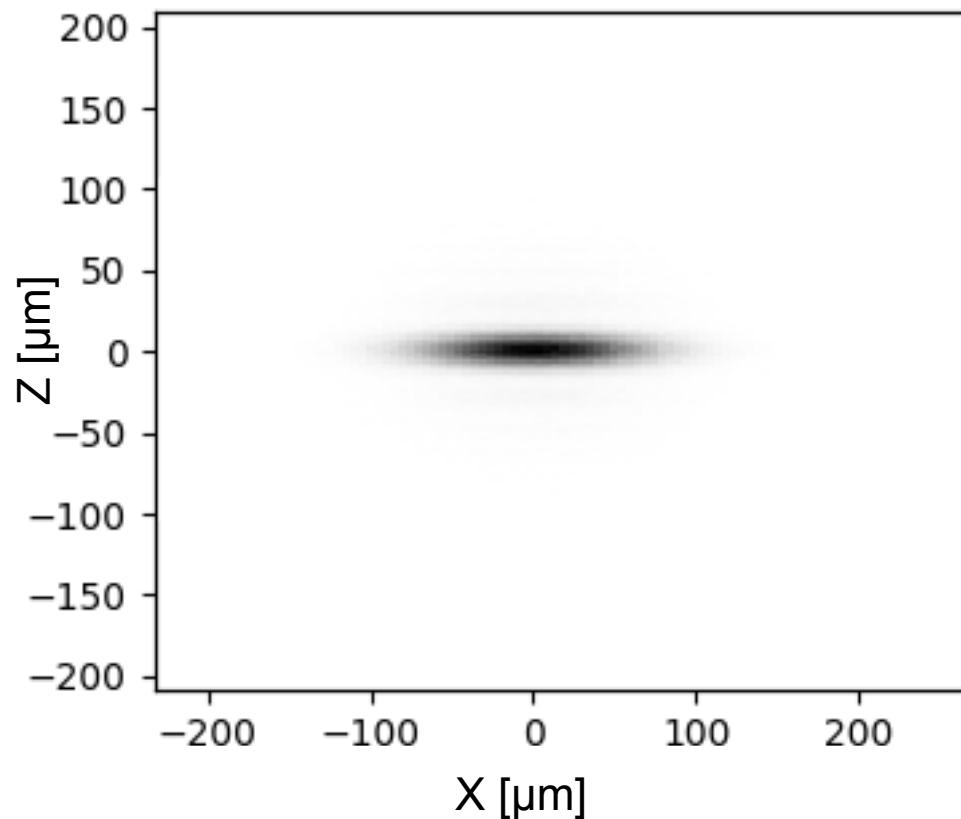


FWHM: $485 \times 80 \mu\text{rad}^2$

A couple of undulator simulations (next generation storage ring)

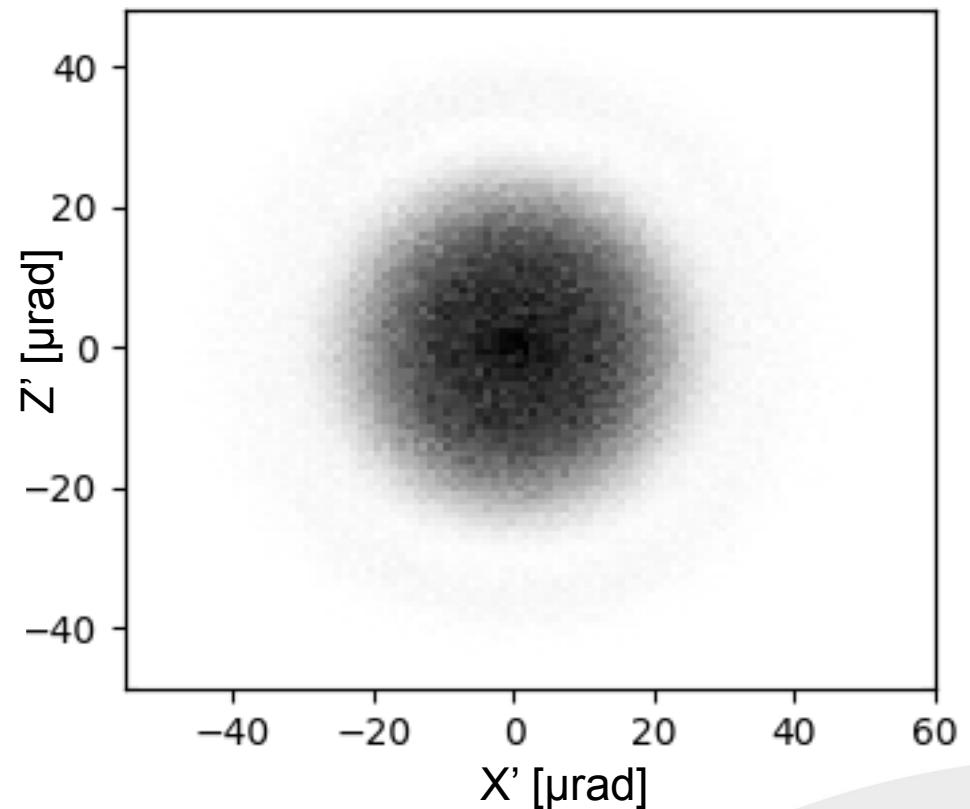
$E_e=2.4\text{GeV}$, $N = 17$, period = 56mm, first harmonic only

Source Size



FWHM: $100 \times 13 \mu\text{m}^2$

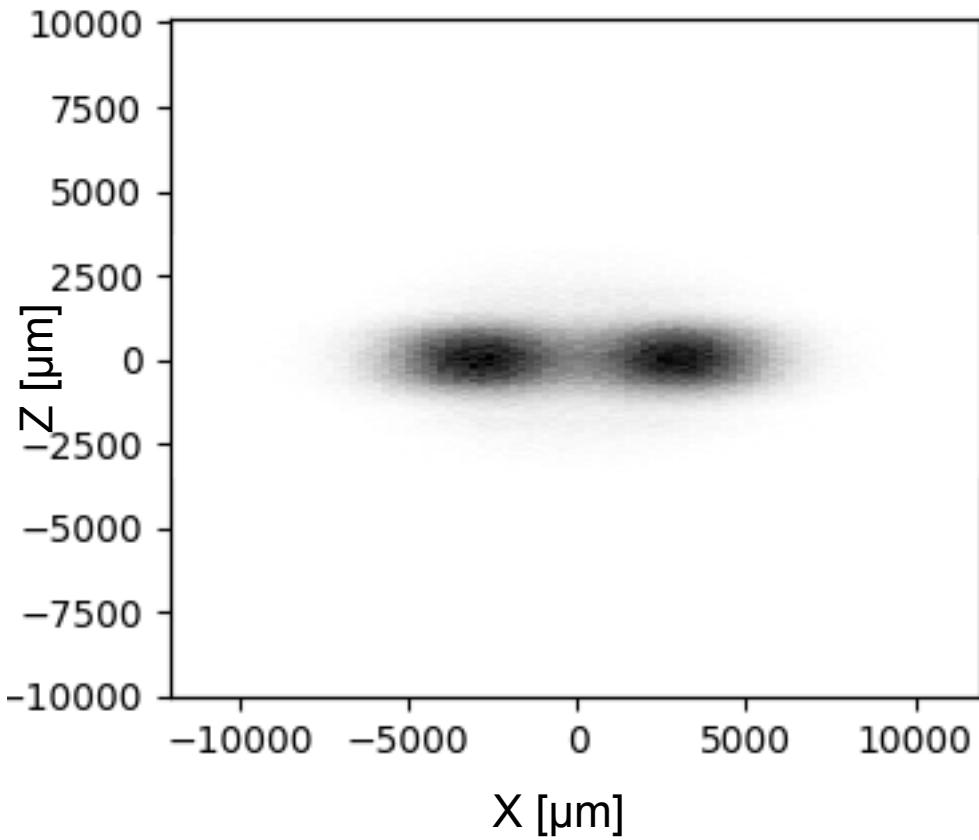
Divergence



FWHM: $33 \times 33 \mu\text{rad}^2$

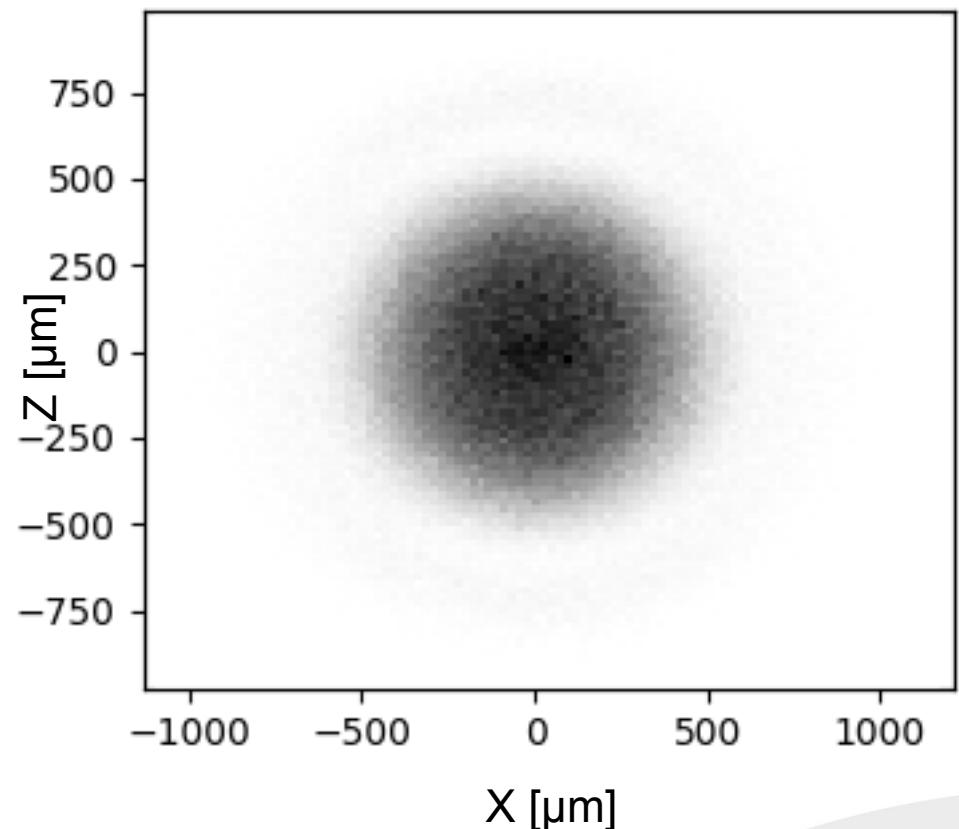
... and if we propagate those 20 m downstream

3rd generation storage ring



FWHM: 9600 x 1570 μm²

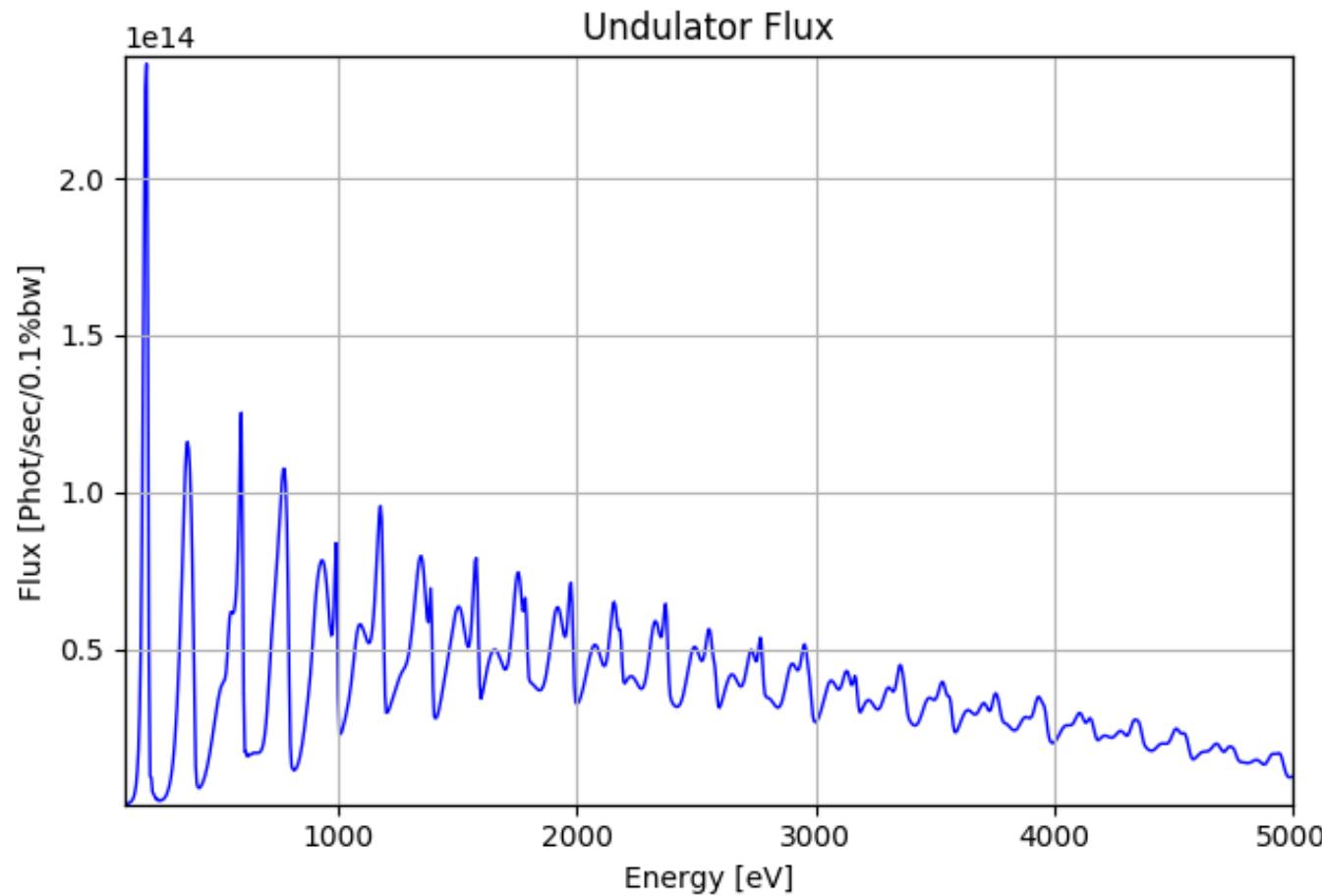
Next generation storage ring



FWHM: 680 x 670 μm²

It's even more complicated...

$E_e = 2.4\text{GeV}$, $N = 17$, period = 56mm



Total flux $\sim 10^{19}$ ph/s



What do we mean by “Instrumentation at Synchrotron Radiation beamlines”?

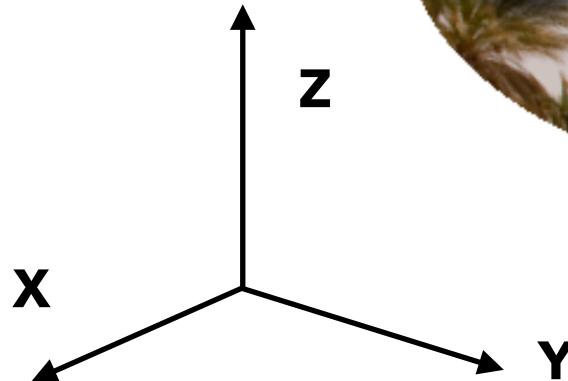
It's a set of interconnected devices designed and operated in order to:

- Deliver the beam produced in the storage ring to the sample
- Shape the beam size and divergence
- Select the photon energy appropriate for the experiment

Photon transport systems



Quick word about simulations



By SHADOW convention,
Y is the BEAM PROPAGATION DIRECTION

C. Welnak, P. Anderson, M. Khan, S. Singh, and F. Cerrina, "Recent developments in SHADOW," *Review of Scientific Instruments*, vol. 63, p. 865, 1992.

O. Chubar, P. E. P. O. T. E. Conference, 1998, "Accurate and efficient computation of synchrotron radiation in the near field region," accelconf.web.cern.ch

L. Rebuffi, M. Sanchez del Rio, "OASYS (OrAnge SYnchrotron Suite): an open-source graphical environment for x-ray virtual experiments", Proc. SPIE 10388, 103880S (2017) . DOI: 10.1117/12.2274263

L. Rebuffi, M. Sanchez del Rio, "ShadowOui: A new visual environment for X-ray optics and synchrotron beamline simulations", J. Synchrotron Rad. 23 (2016). DOI:10.1107/S1600577516013837



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Spoiler alert!

Introducing some optics relations



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Handles available for “manipulating” photons

Refraction

$$n = 1 - \delta + i\beta$$

$$\begin{aligned}\delta &= 10^{-1} \div 10^{-6} \\ \beta &= 10^{-1} \div 10^{-8}\end{aligned}$$

Usage

Focussing

Reflection

$$\sin\phi' = \frac{\sin\phi}{n} \cong \frac{\sin\phi}{1 - \delta}$$

$$\begin{aligned}\theta_C &\approx \sqrt{2\delta} \quad (\text{rad}) \\ \theta_C &\approx 81\sqrt{\delta} \quad (\text{degrees})\end{aligned}$$

Transport
Divergence corrections
Focussing

Diffraction

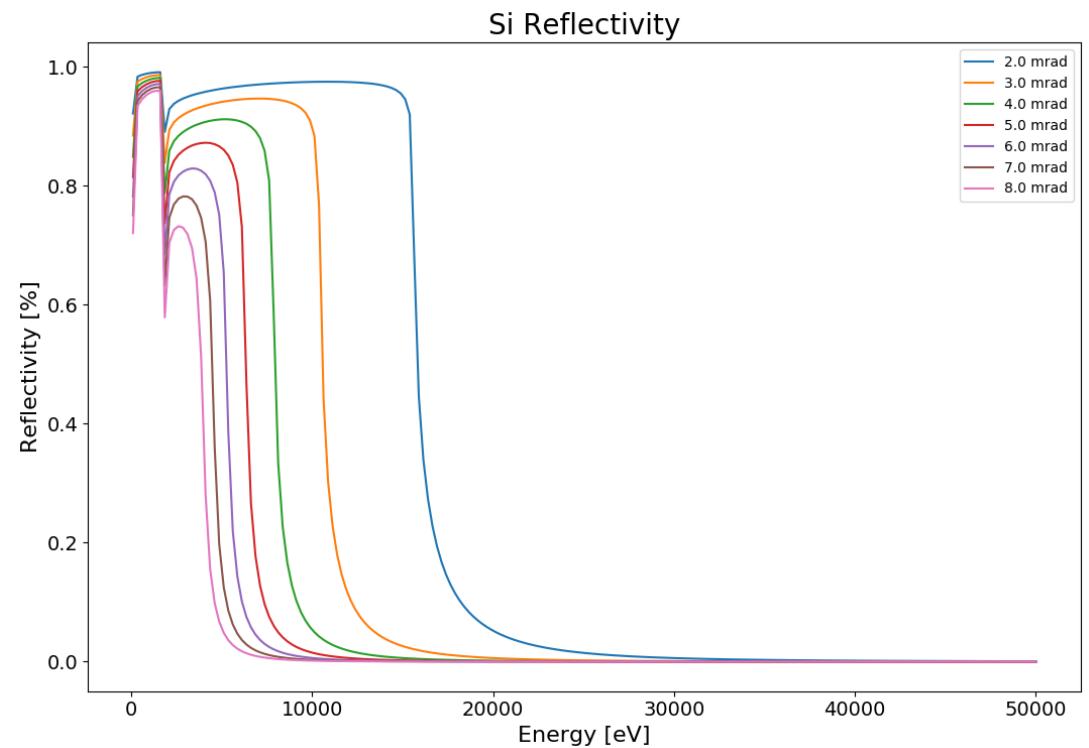
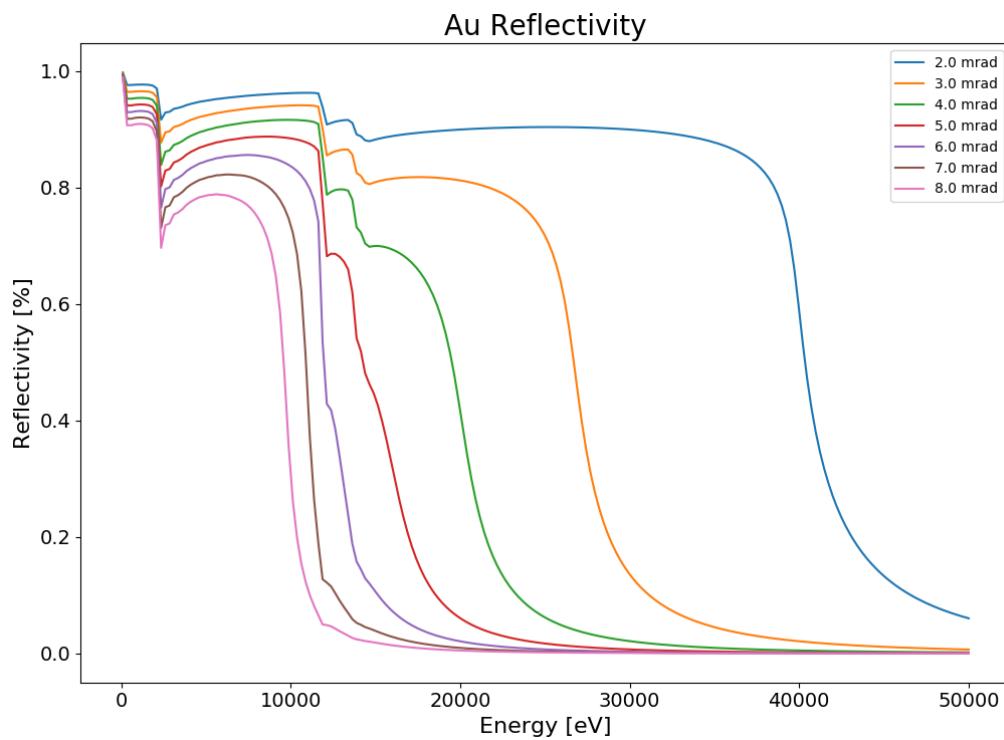
$$2d \cdot \sin\theta = m\lambda$$

$$d \cong \lambda$$

Monochromatization
Focussing

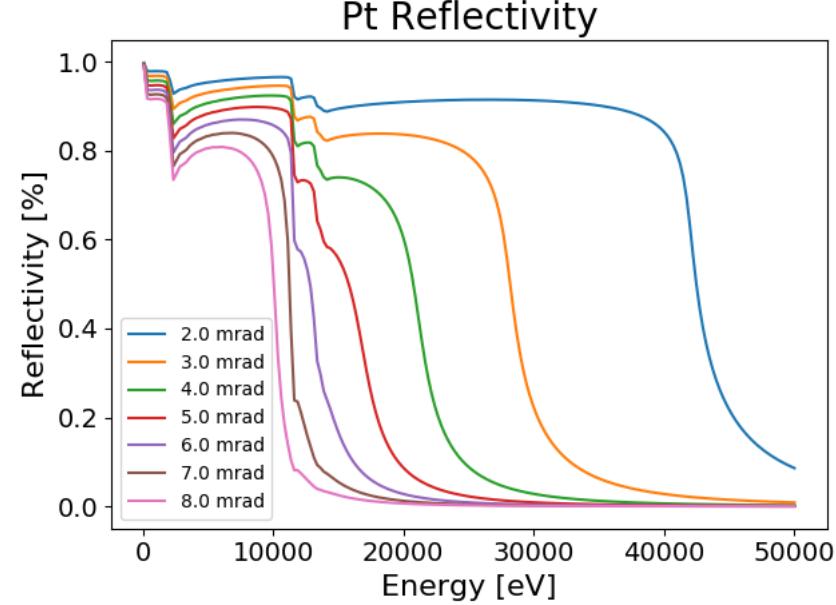
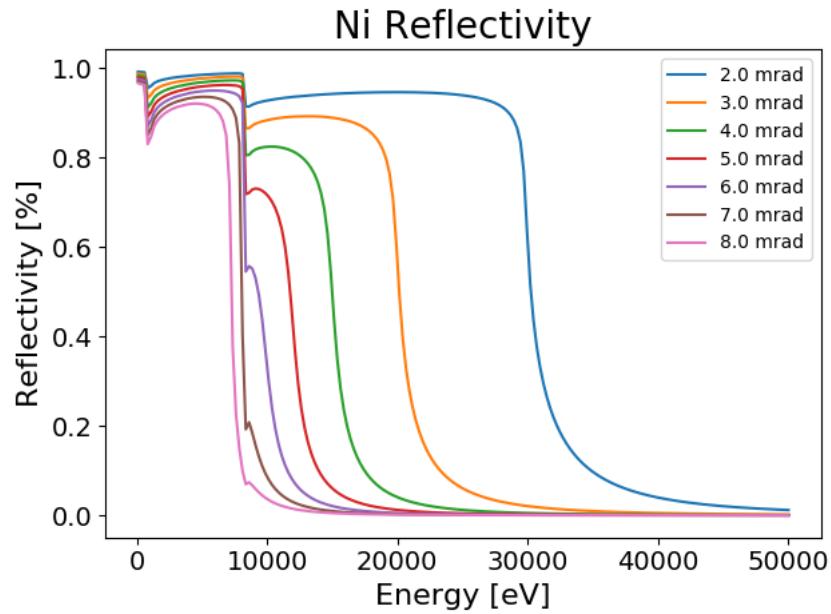
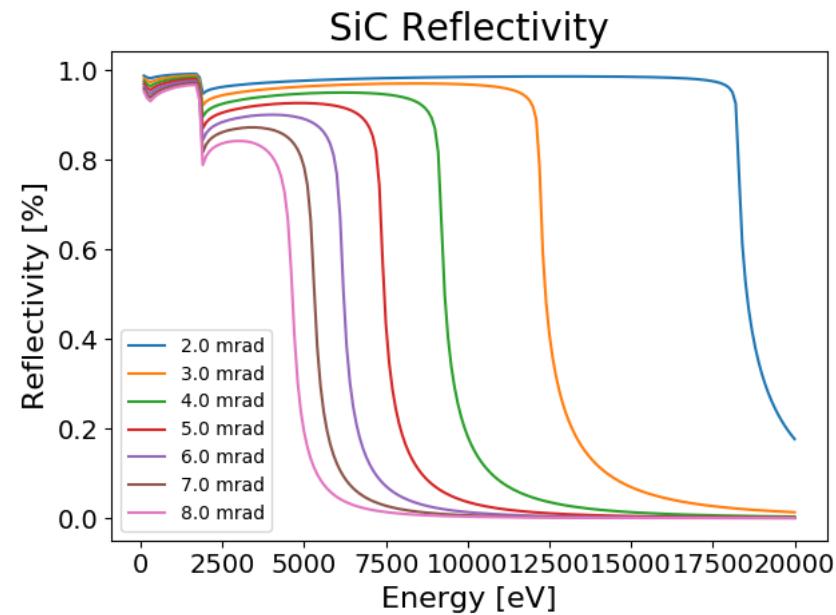
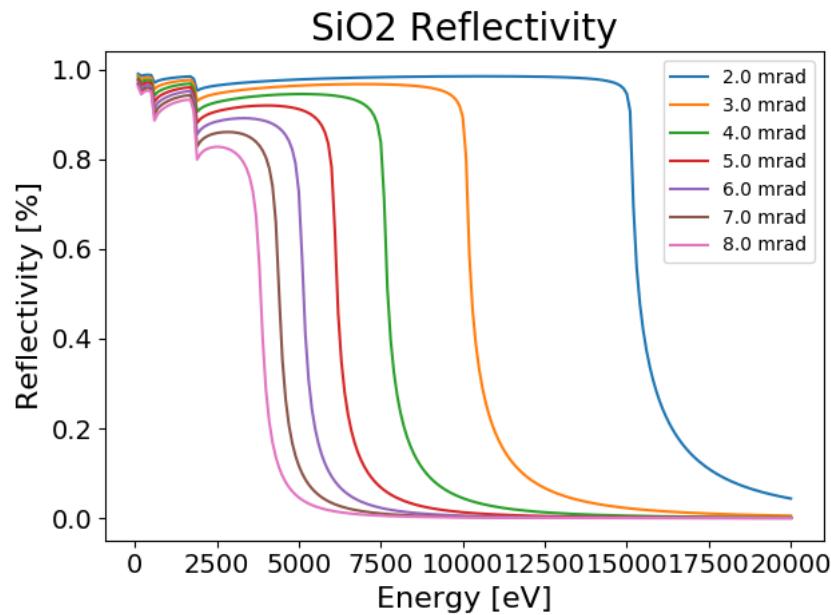
A quick look at reflectivities

$$\theta_c = \sqrt{2\delta} \propto \lambda \sqrt{Z}$$



The higher the energy, the more grazing the incidence angle

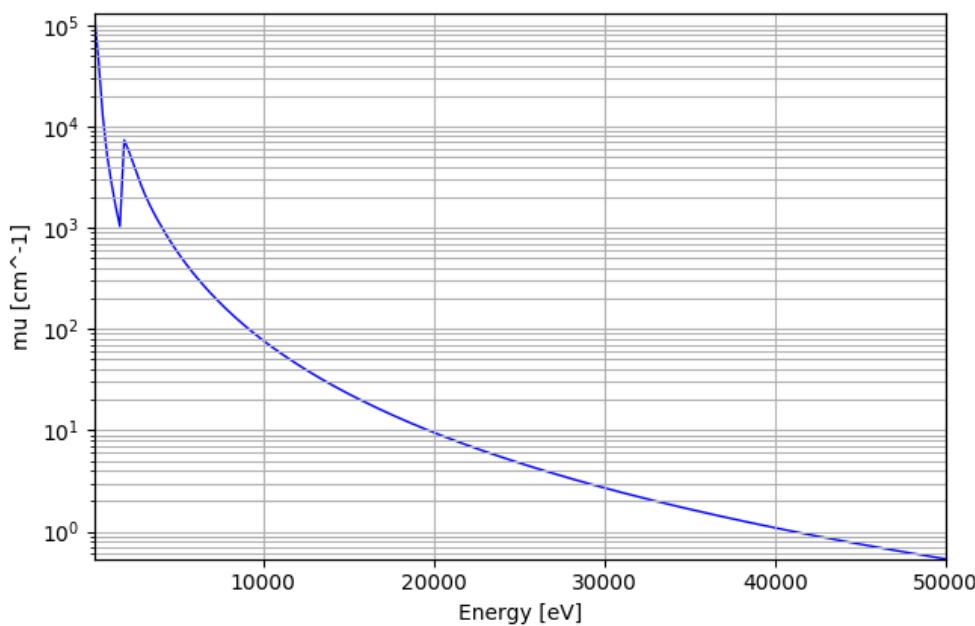
A quick look at reflectivities



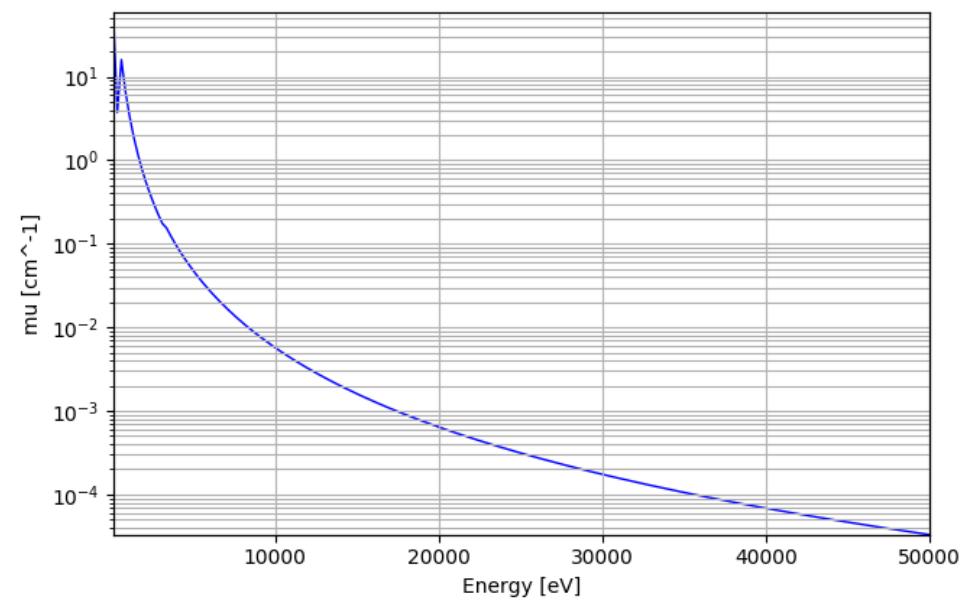
A quick look at attenuation coefficients

$$l_{att} \propto \beta$$

Silicon



Air, 1 atm, sea level



All beamlines must be in vacuum!



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Reflecting elements



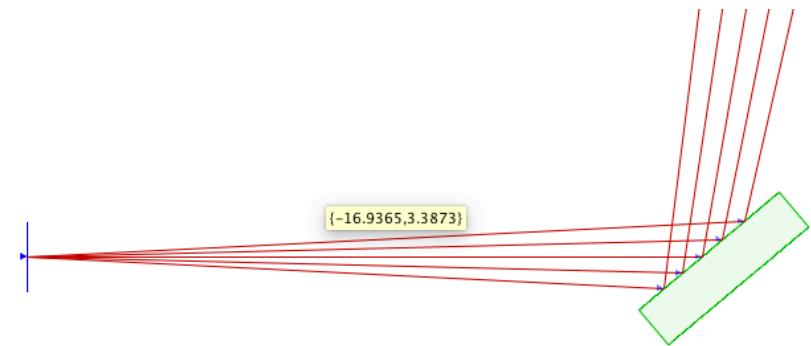
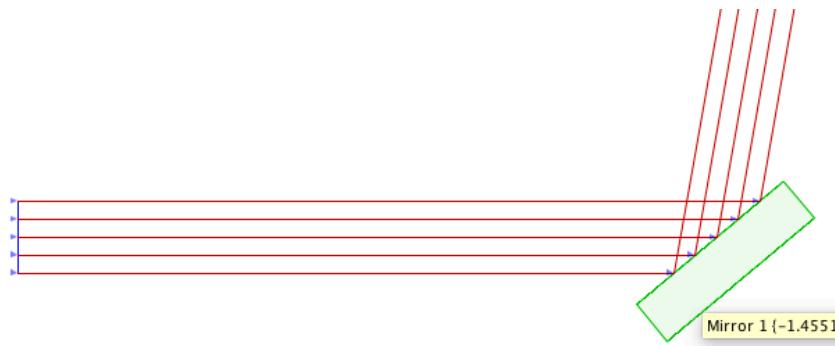
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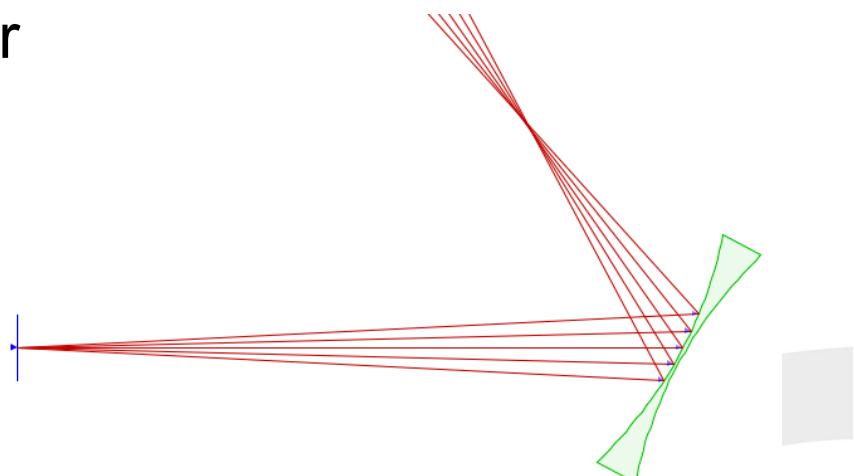
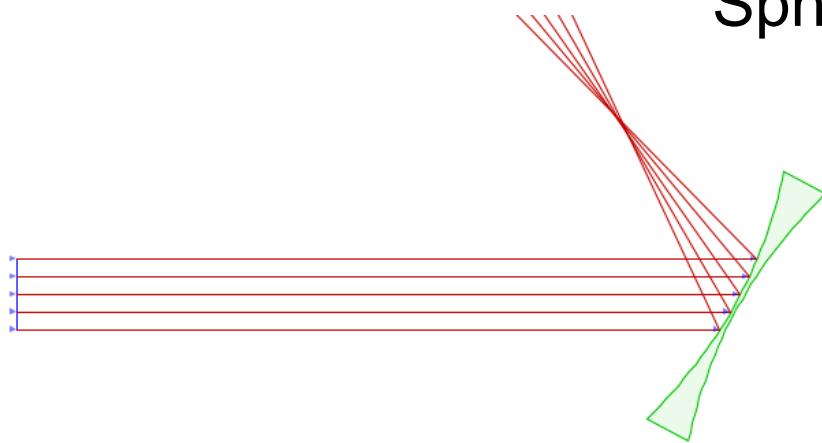
Reflection optics

$$\theta_{inc} = \theta_{refl}$$

Plane mirror

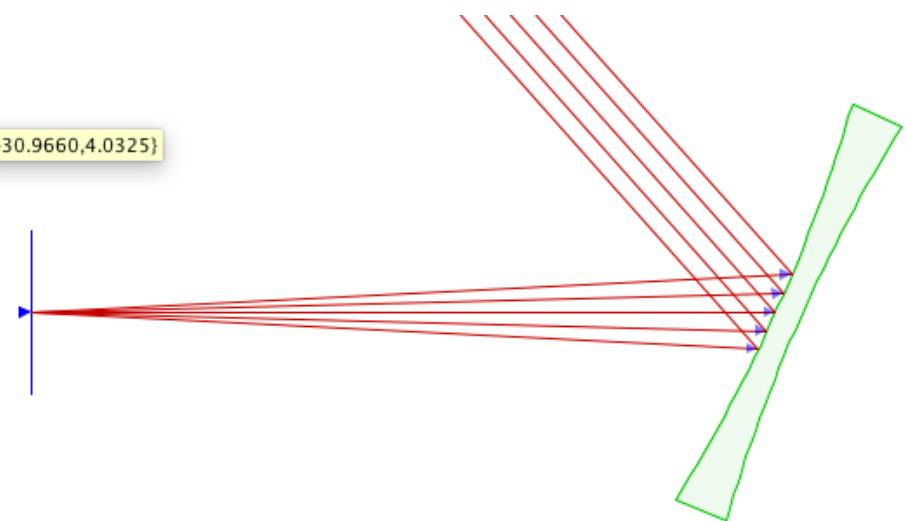
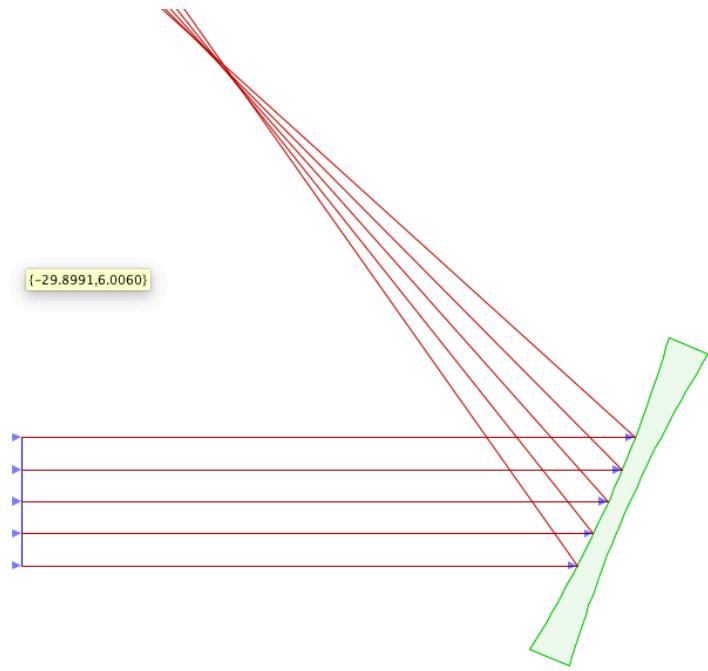


Spherical mirror

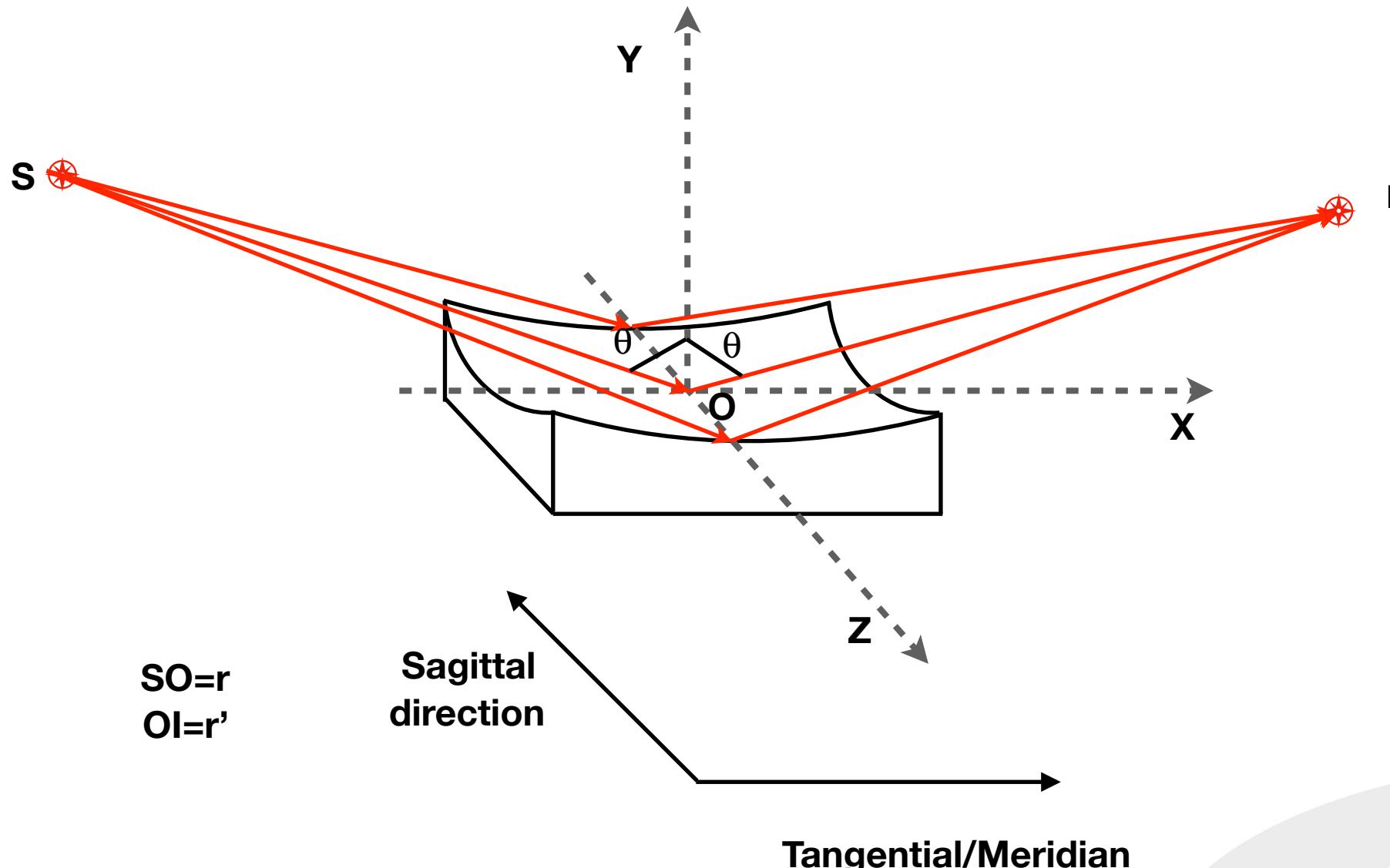


Reflection optics

Parabolic mirror



Some nomenclature

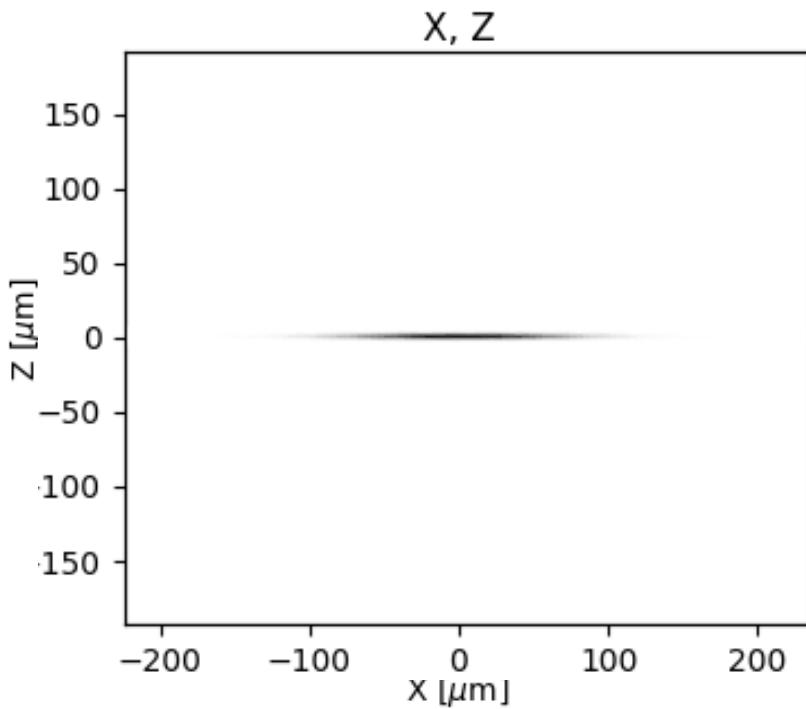


Source for examples

Spatial Dimensions:

$$\sigma_x = 48 \mu m \quad \sigma_z = 1.3 \mu m$$

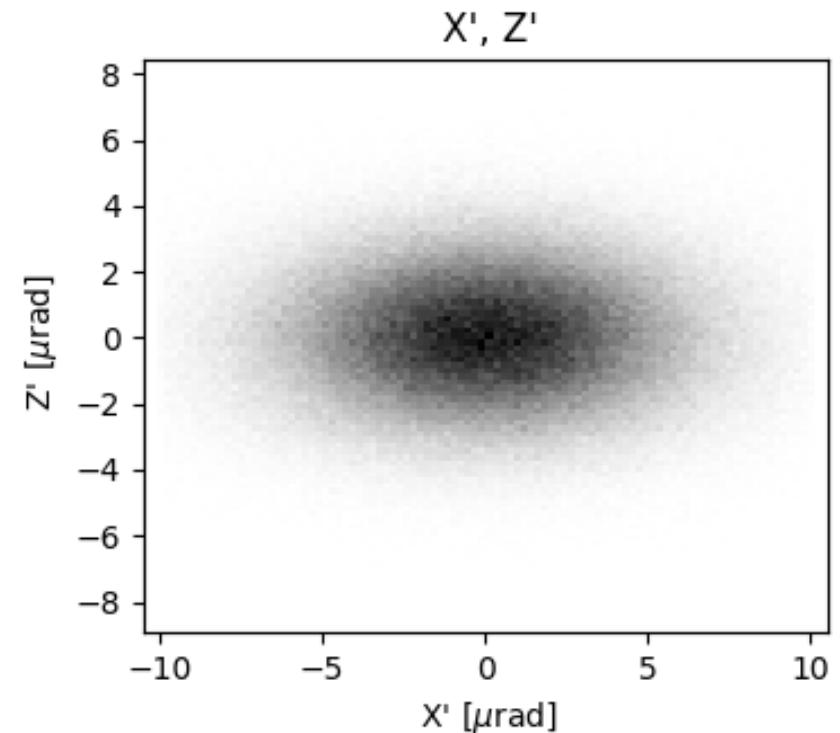
FWHM (X)=105 μm FWHM(Z)=3 μm



Angular dimensions:

$$\sigma'_x = 3.8 \mu rad \quad \sigma'_z = 1.82 \mu rad$$

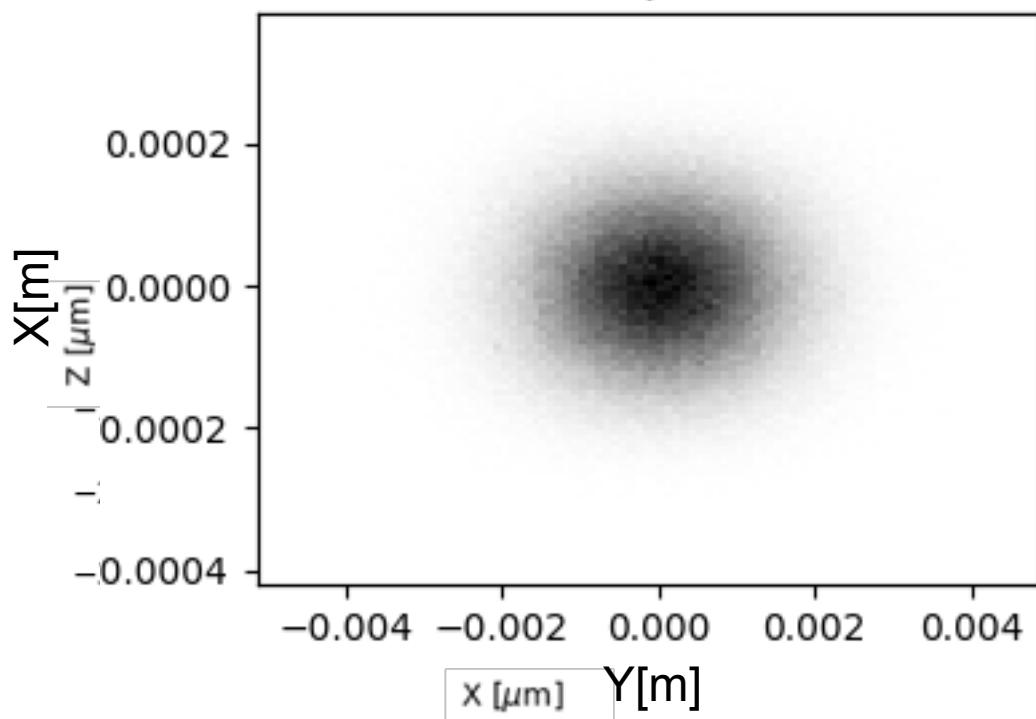
FWHM (X')=8.6 μrad FWHM(Z')=4.2 μrad



Plane mirror, $r = 20 \text{ m}$, $r' = 20 \text{ m}$, $\theta = 88^\circ$

Spatial Dimensions:

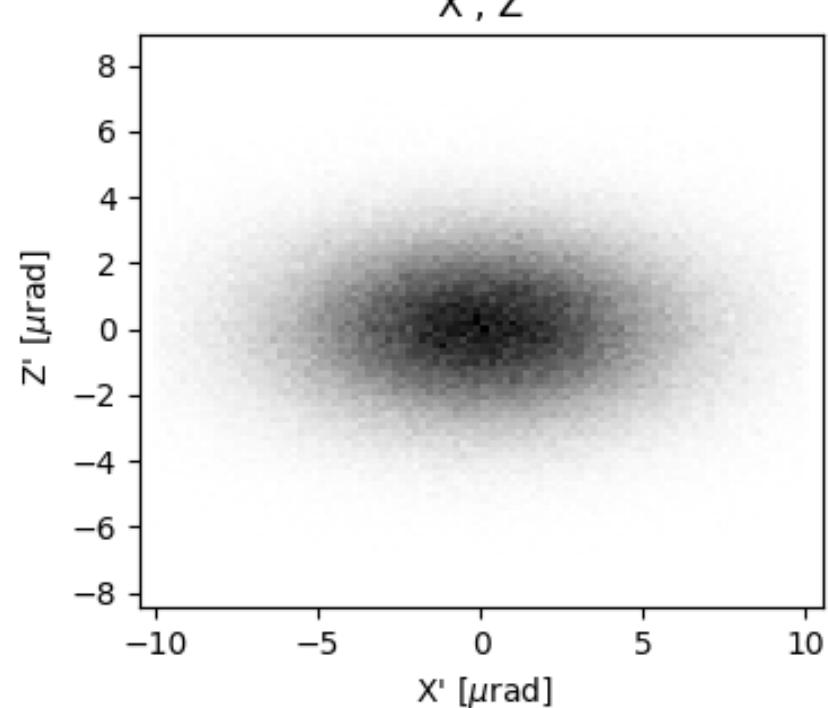
Footprint



$\text{FWHM}(X) = 250 \text{ nm}$ $\text{FWHM}(Z) = 254 \text{ } \mu\text{m}$

Angular dimensions:

X' , Z'

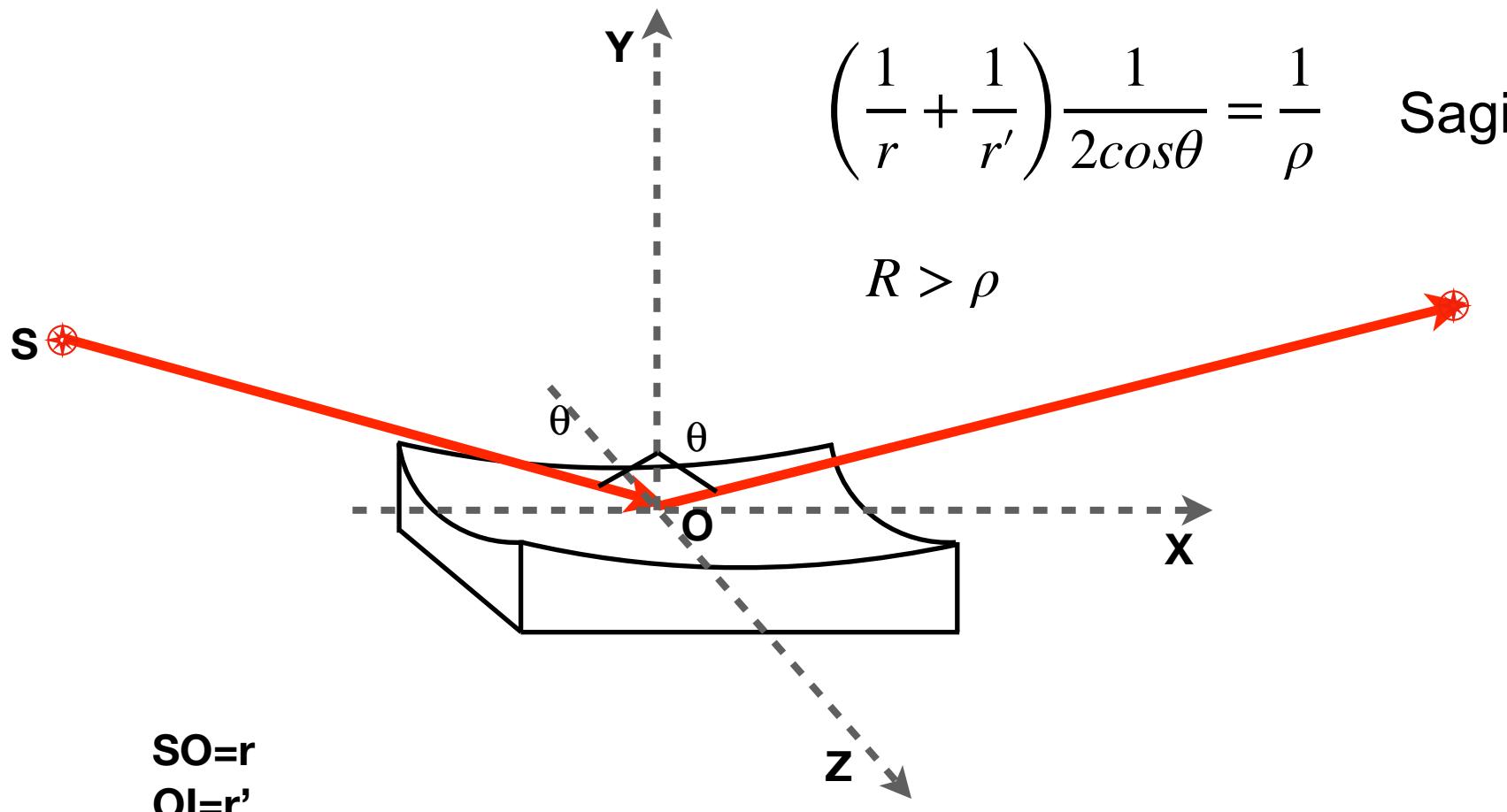


$\text{FWHM}(X') = 8.6 \text{ } \mu\text{rad}$ $\text{FWHM}(Z') = 4.2 \text{ } \mu\text{rad}$

Toroidal mirror

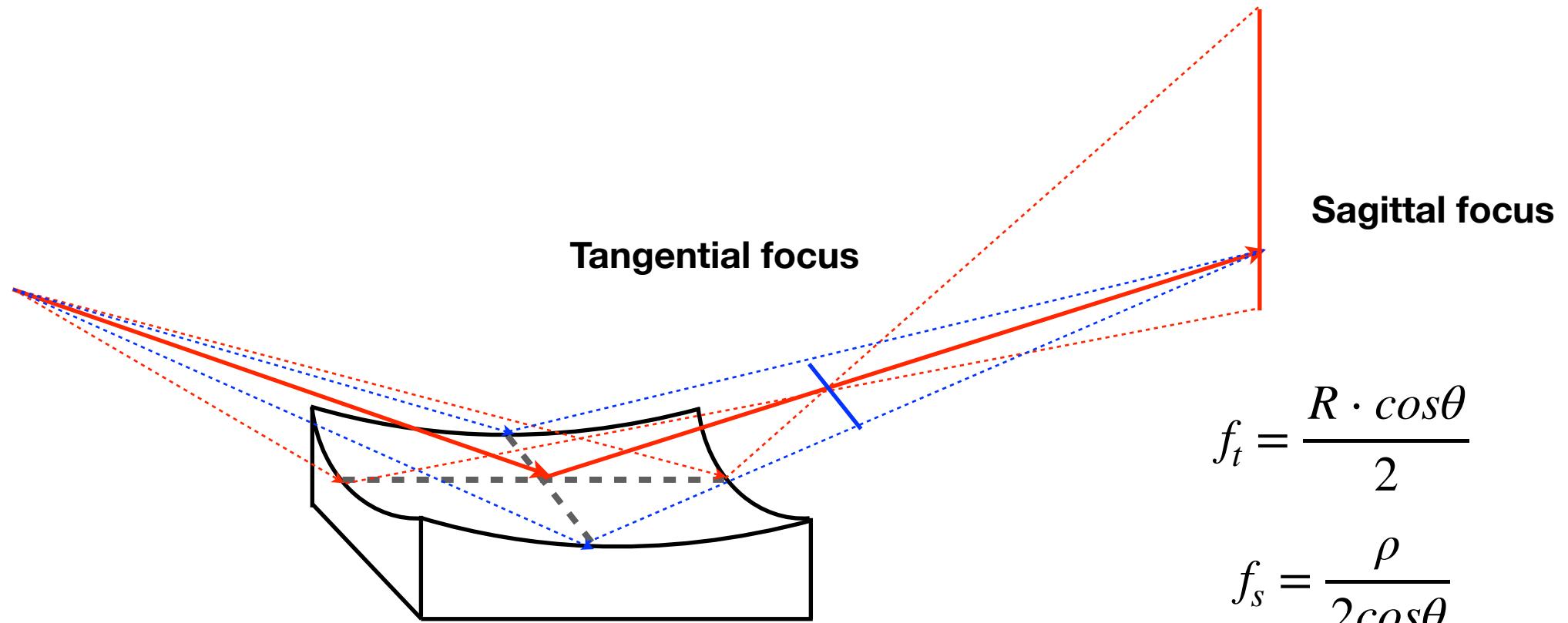
$$\left(\frac{1}{r} + \frac{1}{r'} \right) \frac{\cos\theta}{2} = \frac{1}{R} \quad \text{Tangential focus}$$

$$\left(\frac{1}{r} + \frac{1}{r'} \right) \frac{1}{2\cos\theta} = \frac{1}{\rho} \quad \text{Sagittal focus}$$



$$SO=r \\ OI=r'$$

Toroidal mirror: focussing properties



Condition for a stigmatic image of a point source:

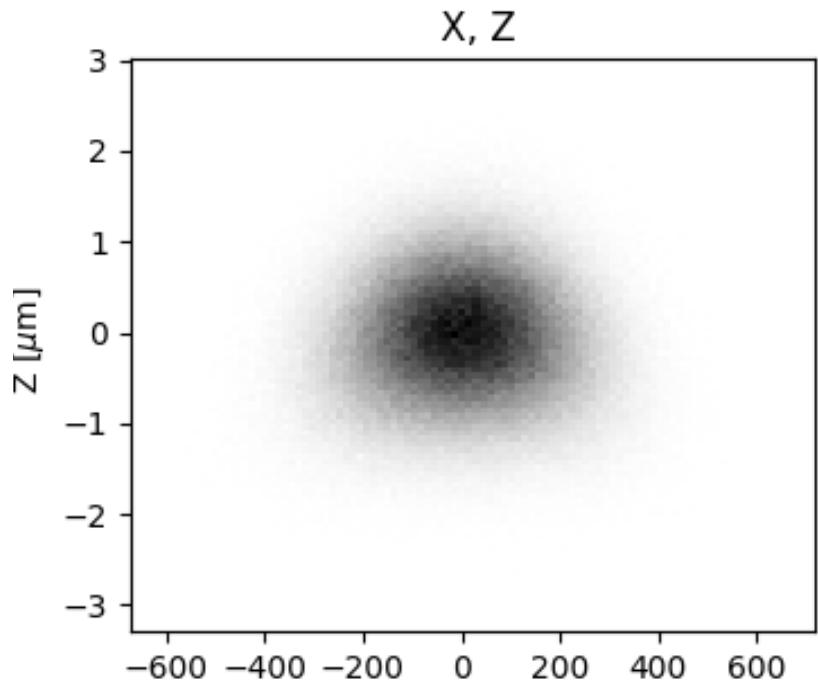
$$\frac{\rho}{R} = \cos^2\theta$$

Toroidal mirror, $r = 20 \text{ m}$, $r' = 10 \text{ m}$, $\theta = 88^\circ$

$$R = \left(\left(\frac{1}{r} + \frac{1}{r'} \right) \frac{\cos\theta}{2} \right)^{-1} = 382 \text{ m} \quad \rho = \left(\left(\frac{1}{r} + \frac{1}{r'} \right) \frac{1}{2\cos\theta} \right)^{-1} = 0.23 \text{ m}$$

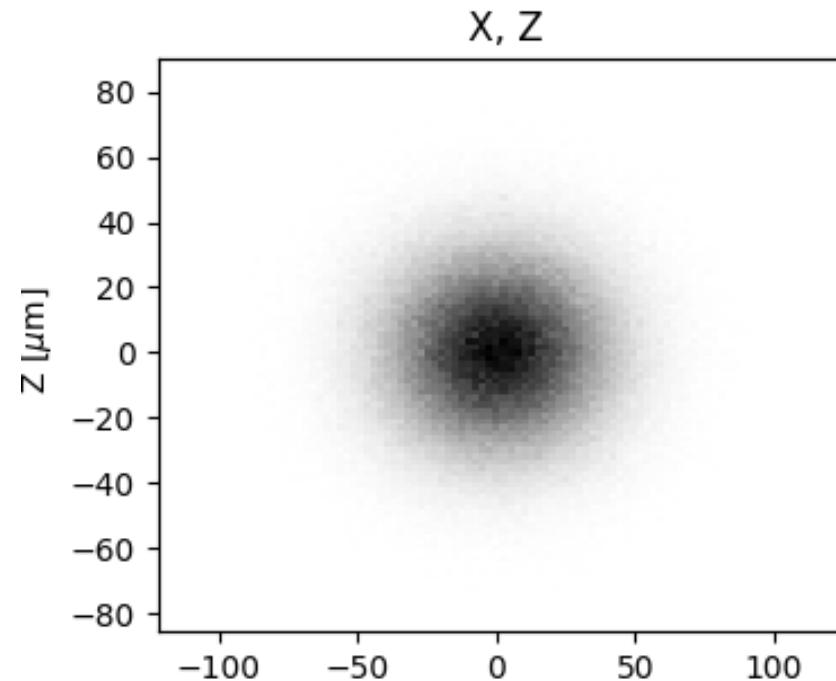
$$f_t = \frac{R \cdot \cos\theta}{2} = 6.6 \text{ m} \quad f_s = \frac{\rho}{2\cos\theta} = 3.3 \text{ m}$$

Tangential focus



FWHM (X)=333 μm FWHM(Z)=1.5μm

Sagittal focus



FWHM (X)=56 μm FWHM(Z)=42μm



Spherical mirrors

Same as toroidal mirrors with:

$$\left(\frac{1}{r} + \frac{1}{r'} \right) \frac{\cos\theta}{2} = \frac{1}{R}$$
$$\left(\frac{1}{r} + \frac{1}{r'} \right) \frac{1}{2\cos\theta} = \frac{1}{R}$$
$$R = \rho$$
$$f_t = \frac{R \cdot \cos\theta}{2}$$
$$f_s = \frac{R}{2\cos\theta}$$

A stigmatic image is only possible if:

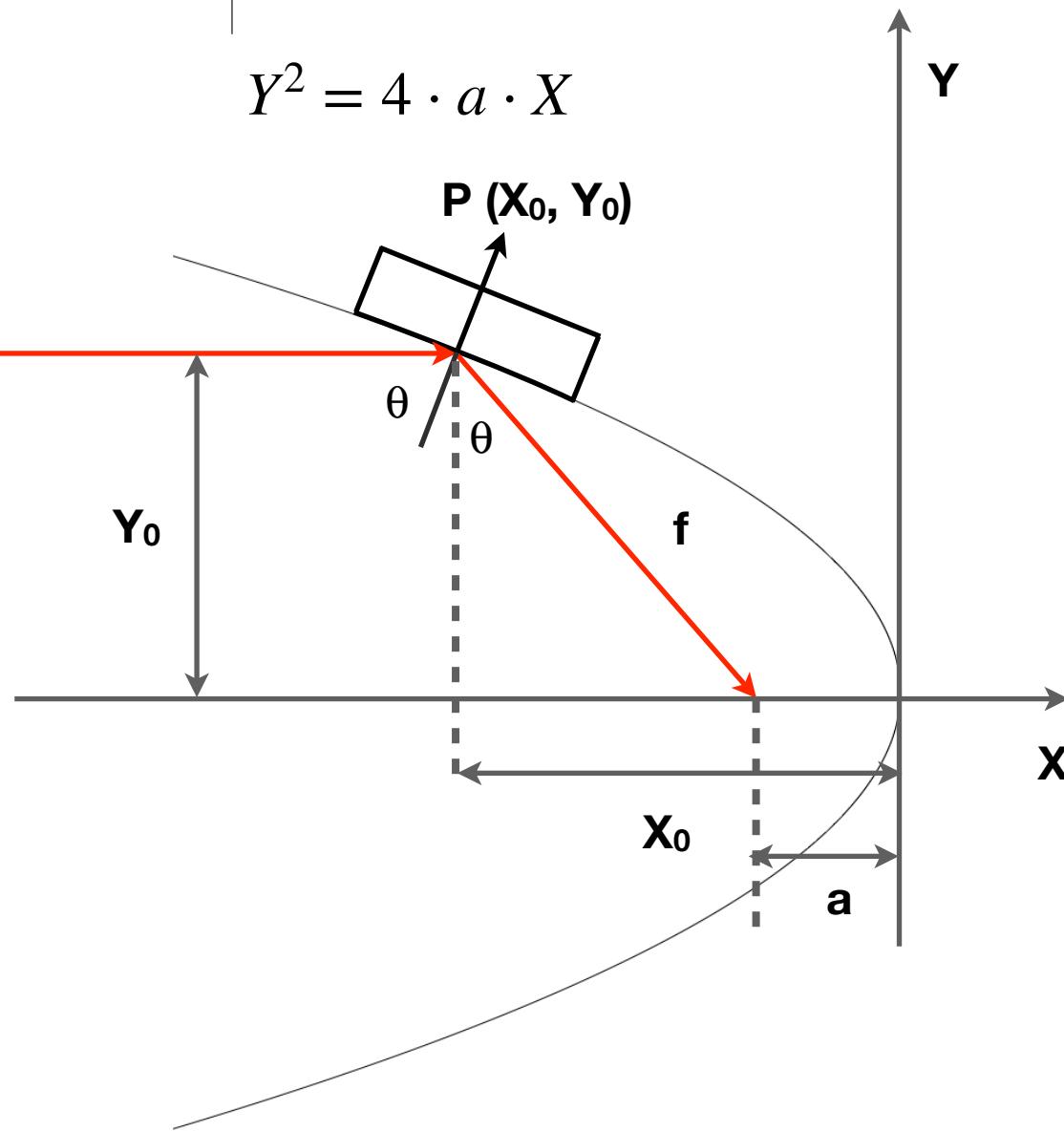
$$\frac{\rho}{R} = \cos^2\theta = 1$$

i.e. this is possible only for normal incidence!

Paraboloidal mirror

$$Y^2 = 4 \cdot a \cdot X$$

<- To source



$P(X_0, Y_0)$:

$$X_0 = a \cdot \tan^2\theta$$

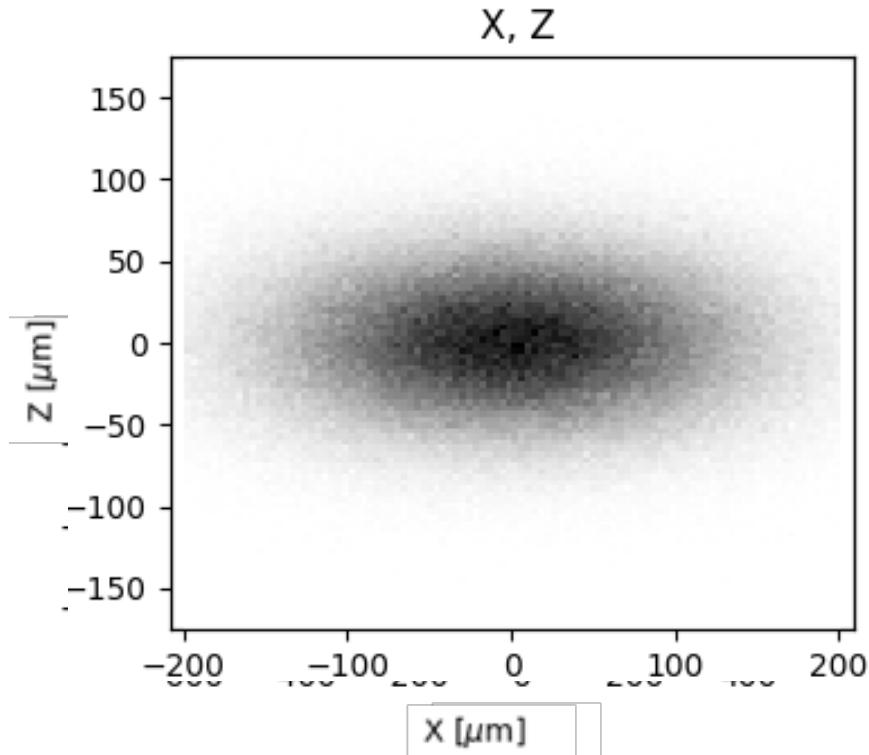
$$Y_0 = 2a \cdot \tan\theta$$

$$f = \frac{a}{\cos^2\theta}$$

Paraboloidal mirror, $r = 20$ m, $r' = 20$ m,
 $\theta = 88^\circ$

$$\text{Parabola parameter } a = f \cos^2 \theta = 0.02435\text{m}$$

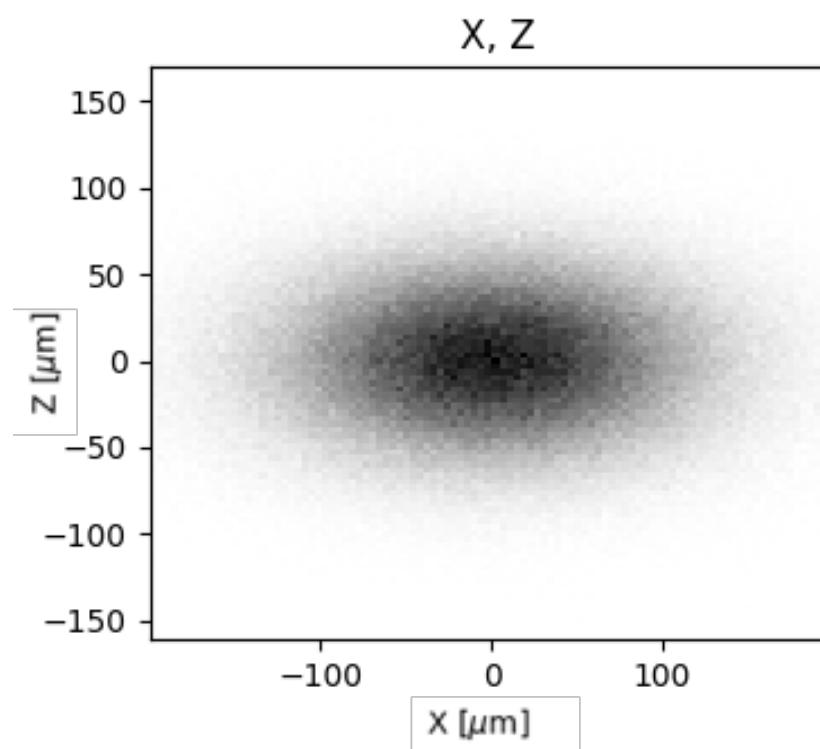
Source image @ 20 mt



FWHM (X)=360 μm FWHM(Z)=854 μm

FWHM (X')=8.6 μrad FWHM(Z')=4.2 μrad

Paraboloidal Mirror image

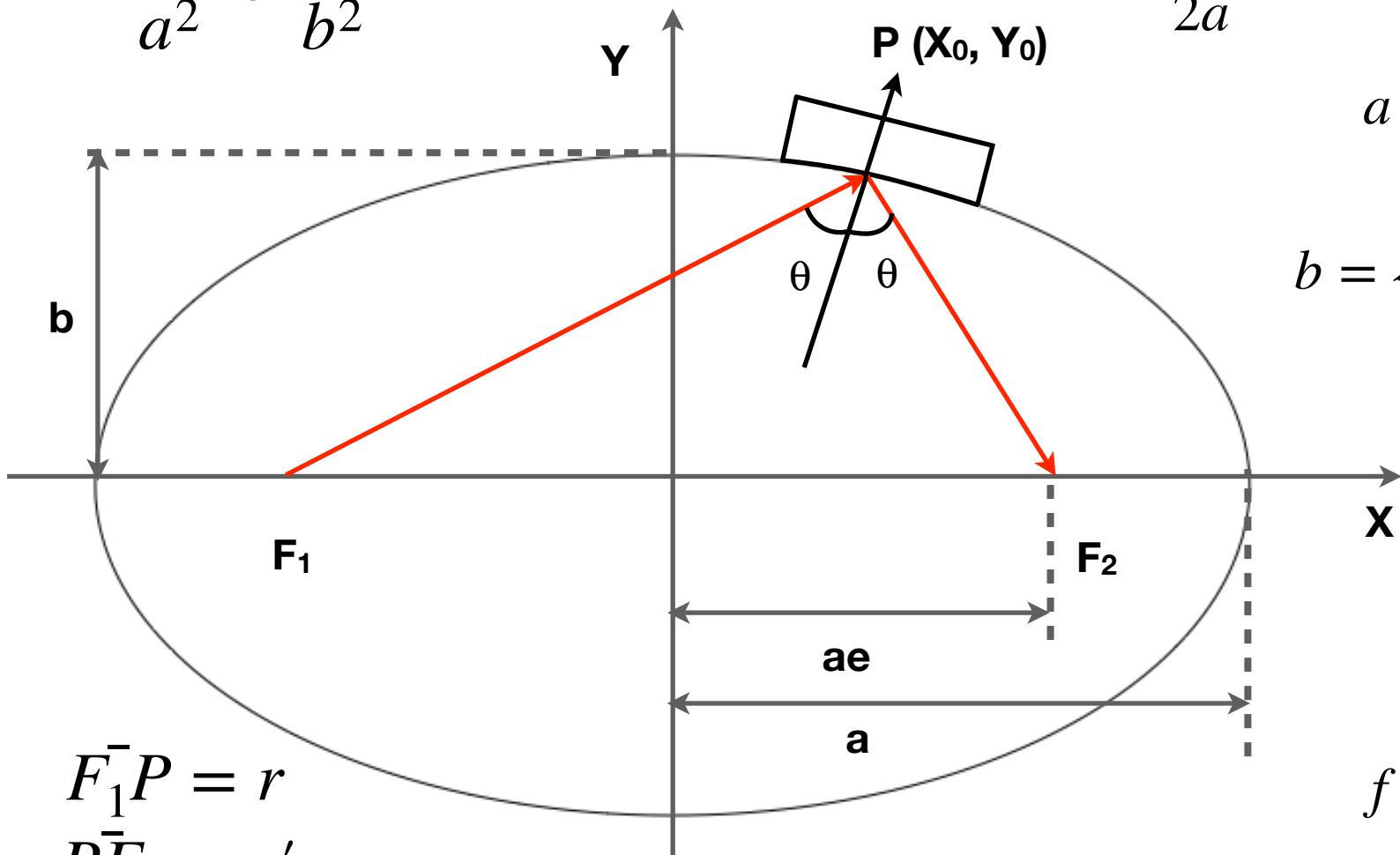


FWHM (X)=172 μm FWHM(Z)=83 μm

FWHM (X')=5.2 μrad FWHM(Z')=0.1 μrad

Ellipsoidal mirror

$$\frac{X^2}{a^2} + \frac{Y^2}{b^2} = 1$$



$$e = \frac{1}{2a} \sqrt{r^2 + r'^2 - 2rr' \cos 2\theta}$$

$$a = \frac{(r + r')}{2}$$

$$b = \sqrt{a^2 (1 - e^2)}$$

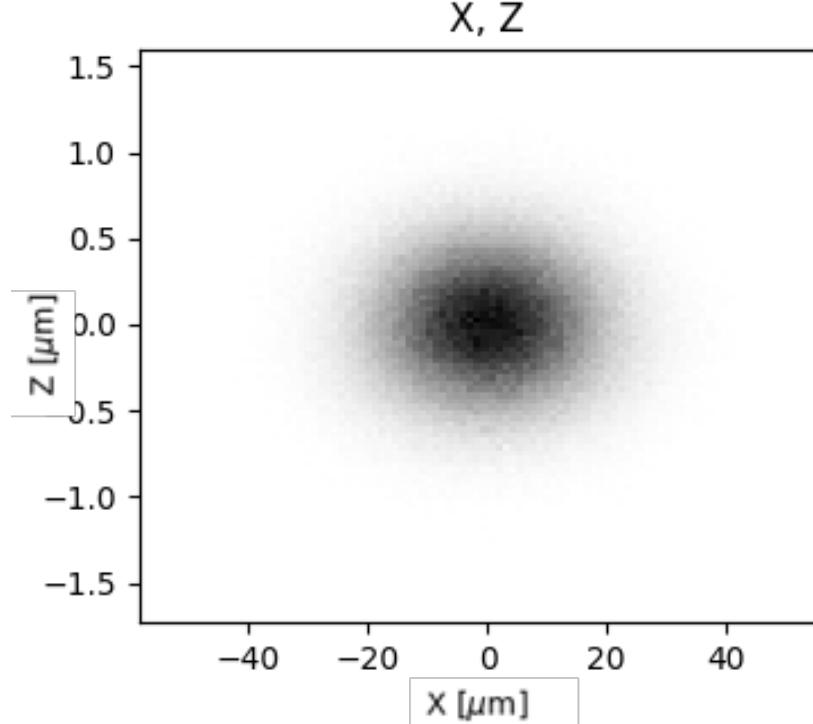
$$f = \left(\frac{1}{r} + \frac{1}{r'} \right)^{-1}$$

Ellipsoidal mirror, $r = 20 \text{ m}$, $r' = 5 \text{ m}$, $\theta = 88^\circ$

$$a = 12.5 \text{ m}, b = 0.349 \text{ m}, e = 0.999610$$

Our source dimensions are: FWHM (X)=105 μm FWHM(Z)=3 μm

$$M = \frac{r'}{r} = 0.25 \text{ i.e. we expect a focus of } \sim 26 \times 0.75 \mu\text{m (FWHM)}$$



$$\text{FWHM (X)}=26 \mu\text{m} \text{ FWHM(Z)}=0.7 \mu\text{m}$$



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WARNING!

All the simulations above are for educational purposes!

- Reflectivity set to 1, and independent of energy
- Ideal source
- No mirror errors (roughness, figure errors, etc)



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Mirror figures used in synchrotron beamlines

| | | Some numbers |
|--------------------|-------------------------------------|--|
| Plane | Re-direction | $R > 100\text{ km}$ |
| Cylindrical | 1D focusing | $R \sim 100's\text{ m}$ |
| Spherical | 2D focusing | $R \sim 100's\text{ m}$ |
| Paraboloid | Infinity to point (or viceversa) | $a \sim \text{cm}, f \sim \text{m}$ |
| Elliptical | Point to point focusing | $r \gg r'$ |
| Toroidal | Astigmatic focusing | $R \sim 100\text{m}, \rho \sim 10's\text{ cm}$ |

All this with an rms roughness $\sim \text{nm}$ or less



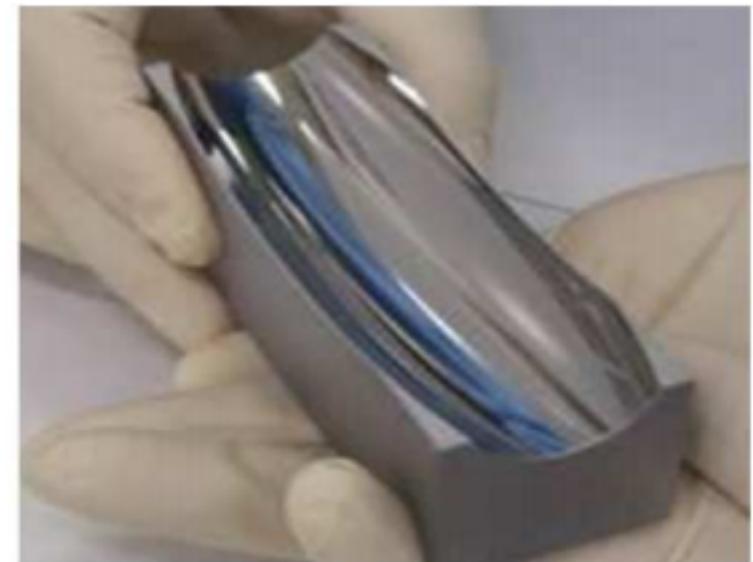
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[http://www.esrf.eu/home/UsersAndScience/
Experiments/CBS/ID09/OpticsHutch/mirror.html](http://www.esrf.eu/home/UsersAndScience/Experiments/CBS/ID09/OpticsHutch/mirror.html)



[http://www.crystal-scientific.com/
mirror_plano.html](http://www.crystal-scientific.com/mirror_plano.html)



R. Radhakirshnan et al,
DOI 10.1149/07711.1255ecst



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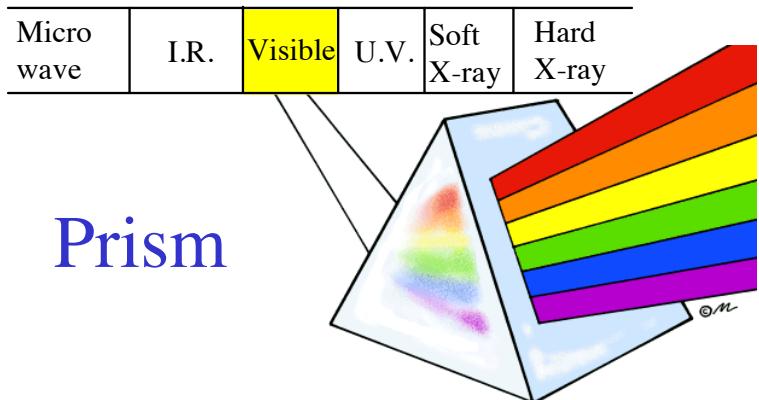
Diffracting elements



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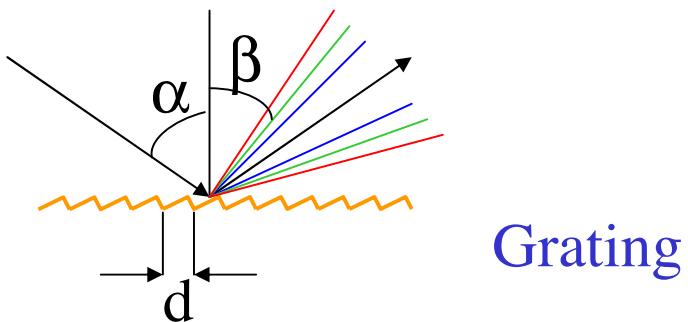
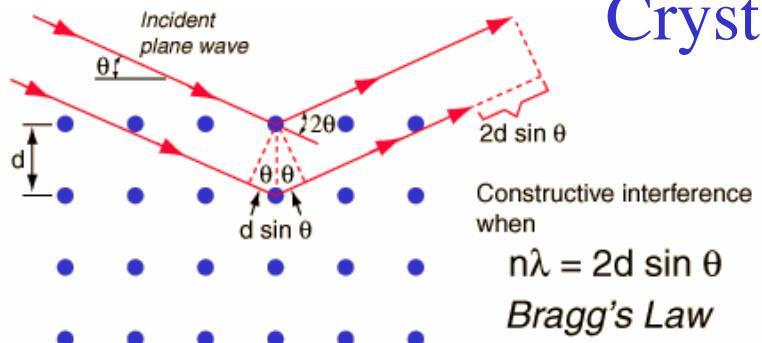
Usage: Overwhelmingly for monochromatization



| | | | | | |
|------------|------|---------|------|------------|------------|
| Micro wave | I.R. | Visible | U.V. | Soft X-ray | Hard X-ray |
|------------|------|---------|------|------------|------------|

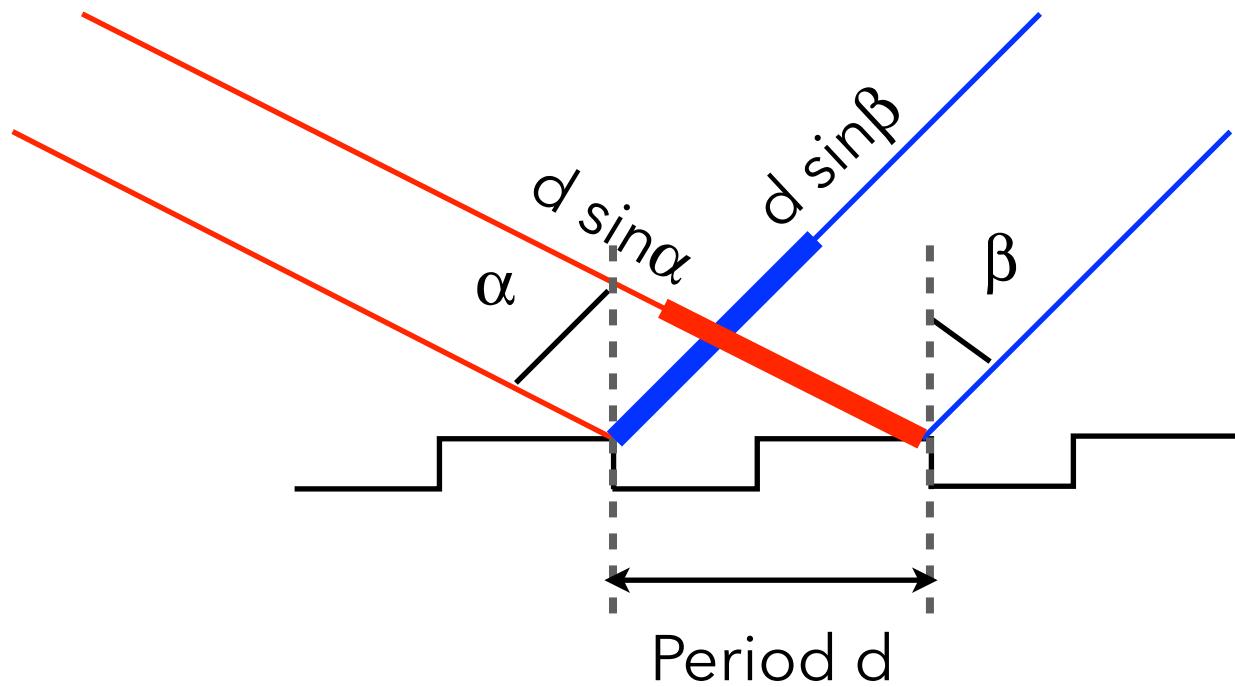
| | | | | | |
|------------|------|---------|------|------------|------------|
| Micro wave | I.R. | Visible | U.V. | Soft X-ray | Hard X-ray |
|------------|------|---------|------|------------|------------|

Crystal



Diffraction gratings

Artificial periodic structure, with a precisely defined period d .



Grating equation

$$\sin\alpha + \sin\beta = Lm\lambda$$

m is the diffraction order

α and β have opposite signs if on opposite side of the surface normal

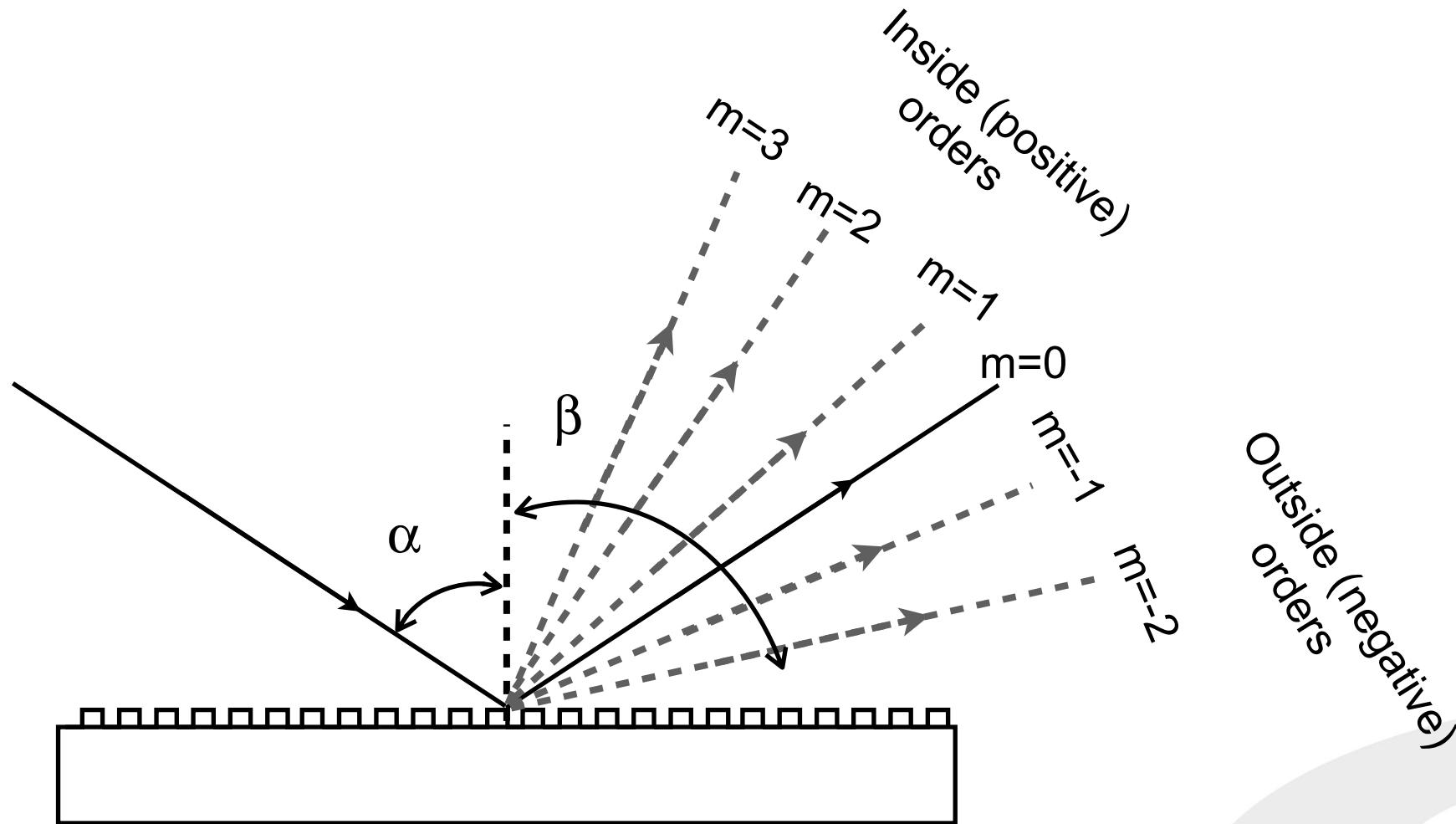
Line density

$$L = \frac{1}{d}$$

Diffraction gratings

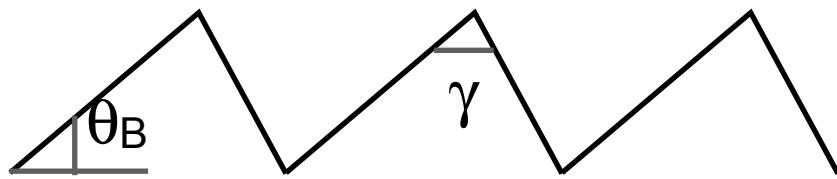
Grating equation

$$\sin\alpha + \sin\beta = Lm\lambda$$

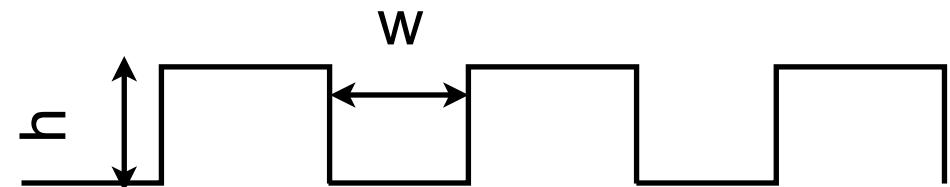


Grating profiles

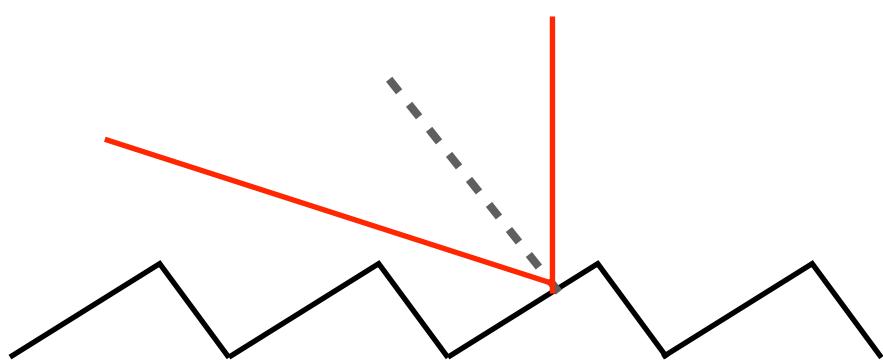
Blazed profile



Laminar profile

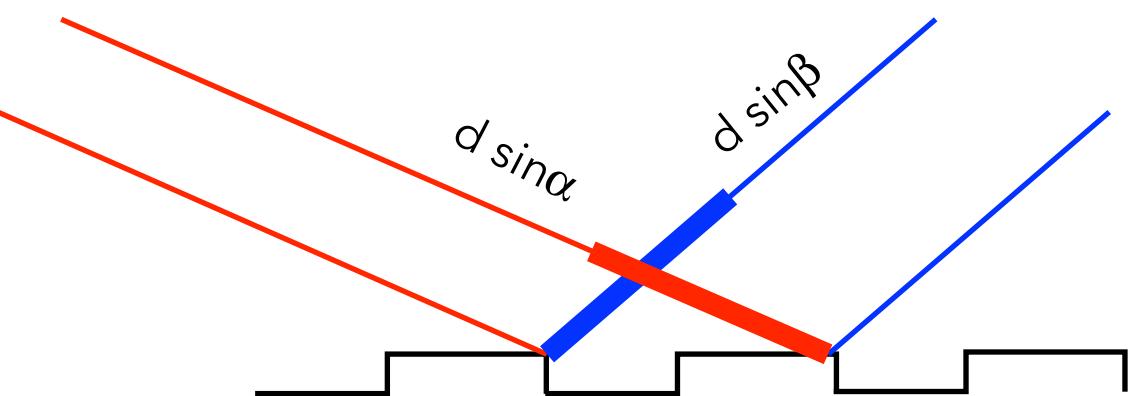


Blaze condition:



$$\theta_B = \frac{\alpha + \beta}{2}$$

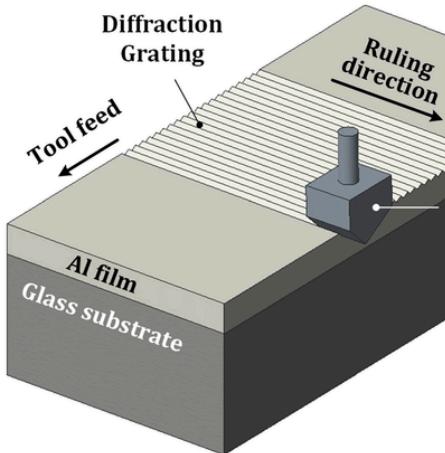
Higher efficiency



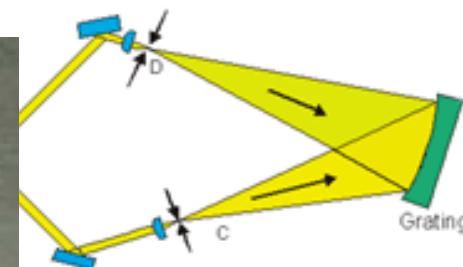
Higher spectral purity

Making of a grating

Mechanical ruling



Holographic ruling



Substrate

After development
and etching



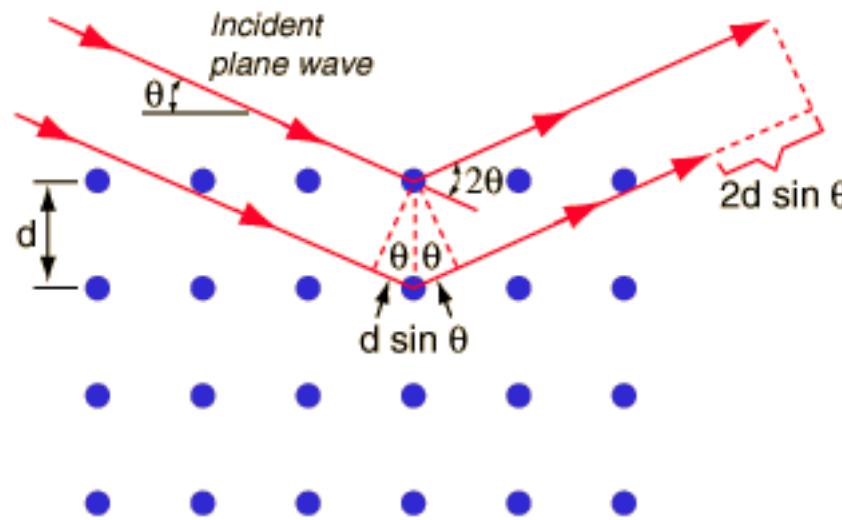
Resist removed



Final coating

Crystals

Based on Bragg's law: $2d \sin \theta = m\lambda$



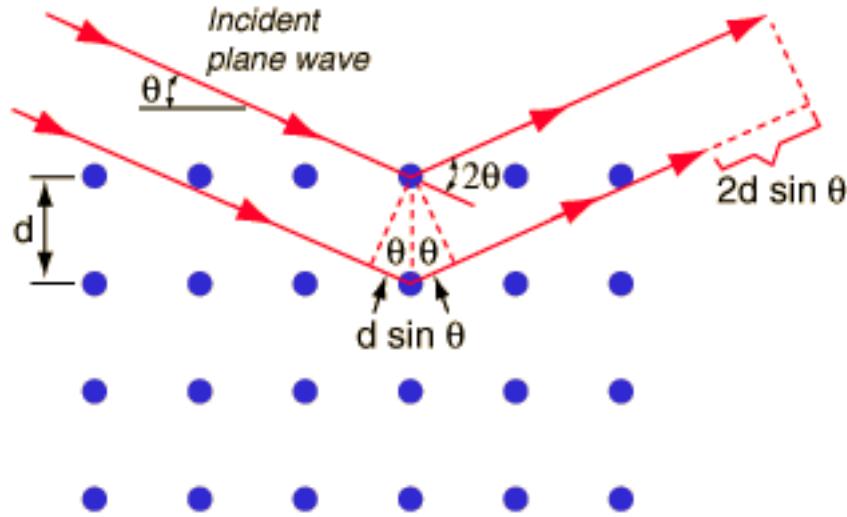
Since $\sin \theta \leq 1$, $\lambda \leq \lambda_{MAX}$ ($E \geq E_{MIN}$) = $2d$

Si(111): $d=3.13 \text{ \AA}$ ($E_{MIN} \sim 2 \text{ keV}$)

Si(311): $d=1.64 \text{ \AA}$ ($E_{MIN} \sim 3.8 \text{ keV}$)

InSb(111): $d=3.74 \text{ \AA}$ ($E_{MIN} \sim 1.7 \text{ keV}$)

Crystals' resolving power



$$\frac{\Delta E}{E} = \frac{\Delta \lambda}{\lambda} = \Delta \theta \frac{\cos \theta}{\sin \theta}$$

Angular spread of the beam

Where does $\Delta\theta$ come from?

$\Delta\theta_{beam}$ Angular divergence of the incoming beam *

$\omega_{crystal}$ Intrinsinc width of Bragg reflection,
the Darwin curve

* more on this later...

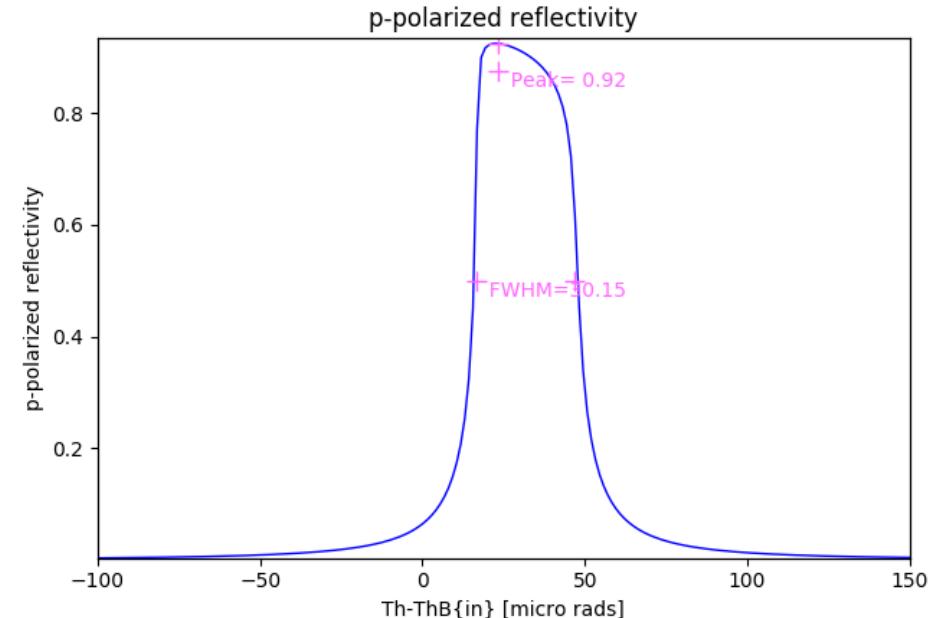
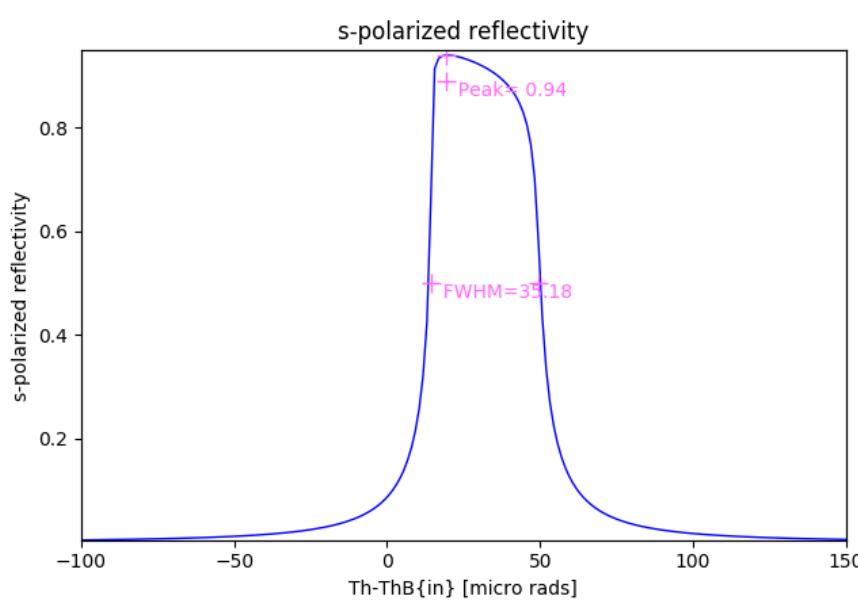
A quick word on the Darwin curve (ω_s)

Obtained by measuring the crystal reflectivity R with:

- 1) Strictly monochromatic beam which is
- 2) Perfectly collimated

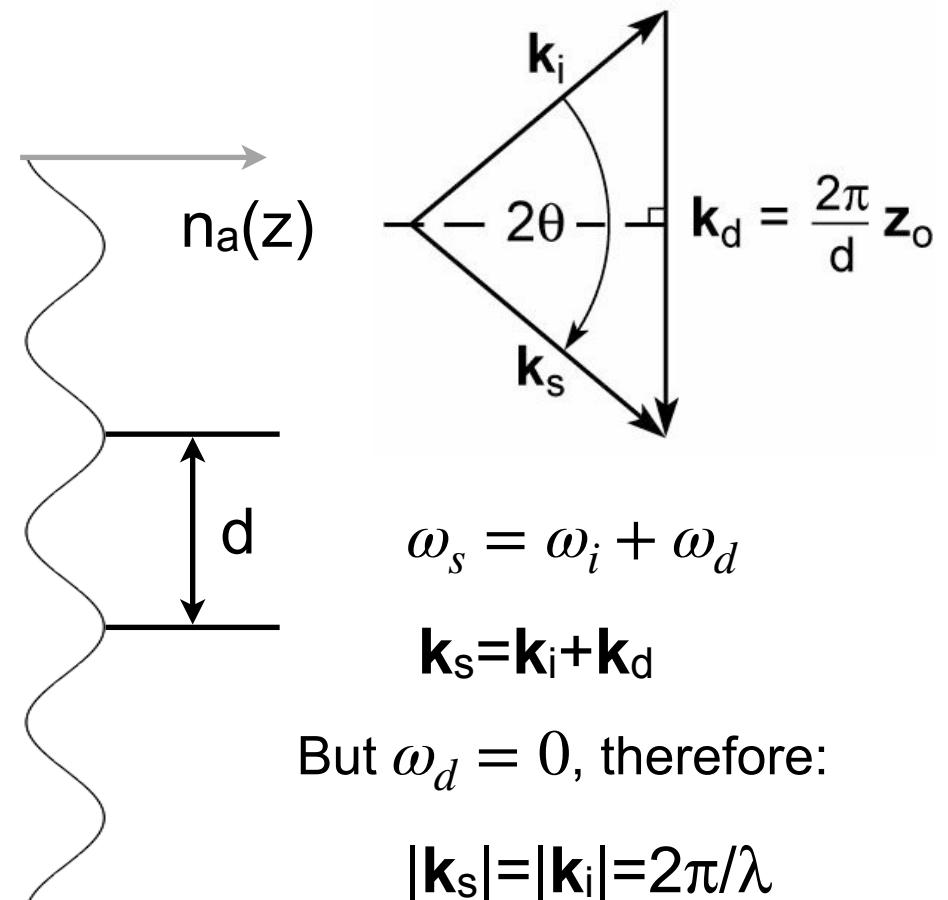
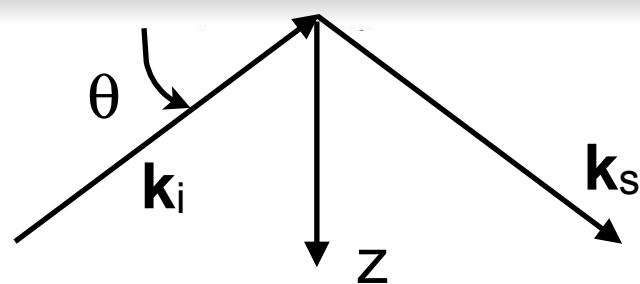
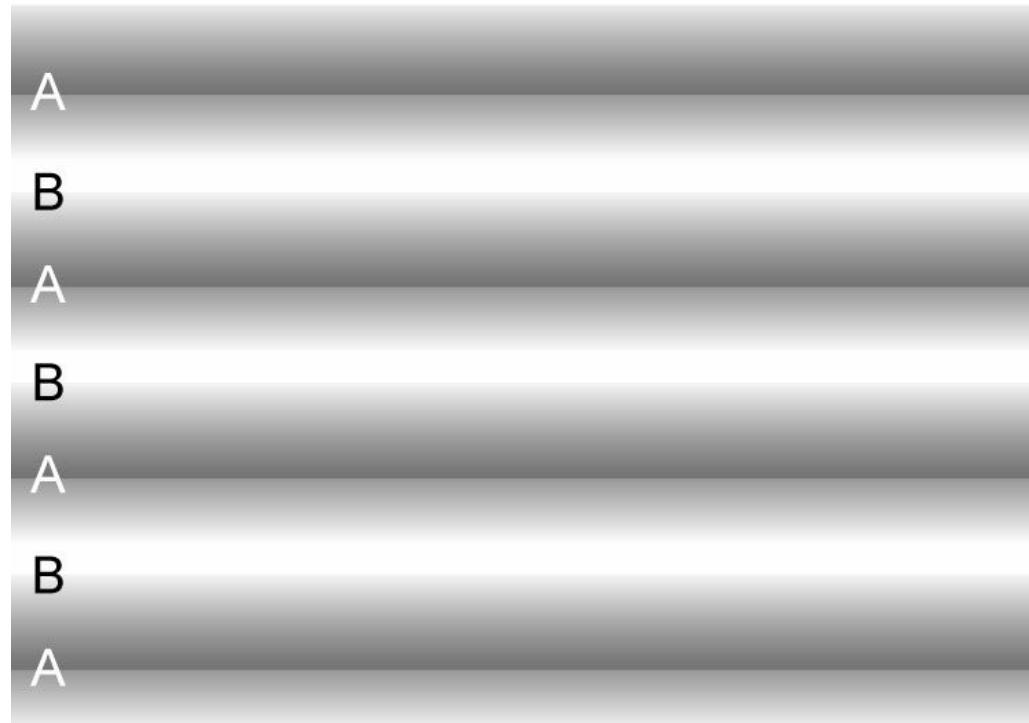
And as a function of $\Delta\theta = \theta - \theta_{Bragg}$

Its expression is derived by dynamic diffraction theory, which includes scattering, and is polarisation-dependent



Simulations for Si (111) at 8keV

Multi-layer mirrors



$$\omega_s = \omega_i + \omega_d$$

$$\mathbf{k}_s = \mathbf{k}_i + \mathbf{k}_d$$

But $\omega_d = 0$, therefore:

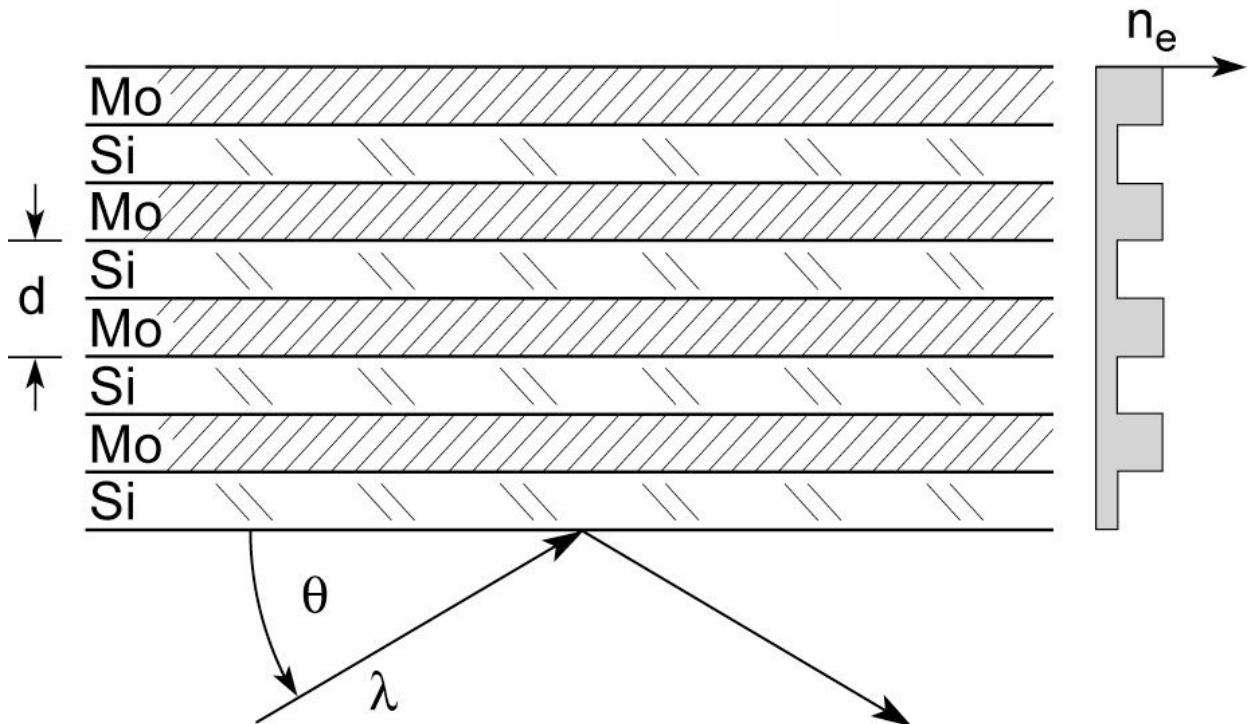
$$|\mathbf{k}_s| = |\mathbf{k}_i| = 2\pi/\lambda$$

$$\sin\theta = \frac{k_d/2}{k_i}$$

$$\lambda = 2ds\sin\theta$$

Multi-layer mirrors

What if $n_a(z)$ is still periodic, but not a simple sinusoid?



$$m\lambda = 2dsin\theta$$

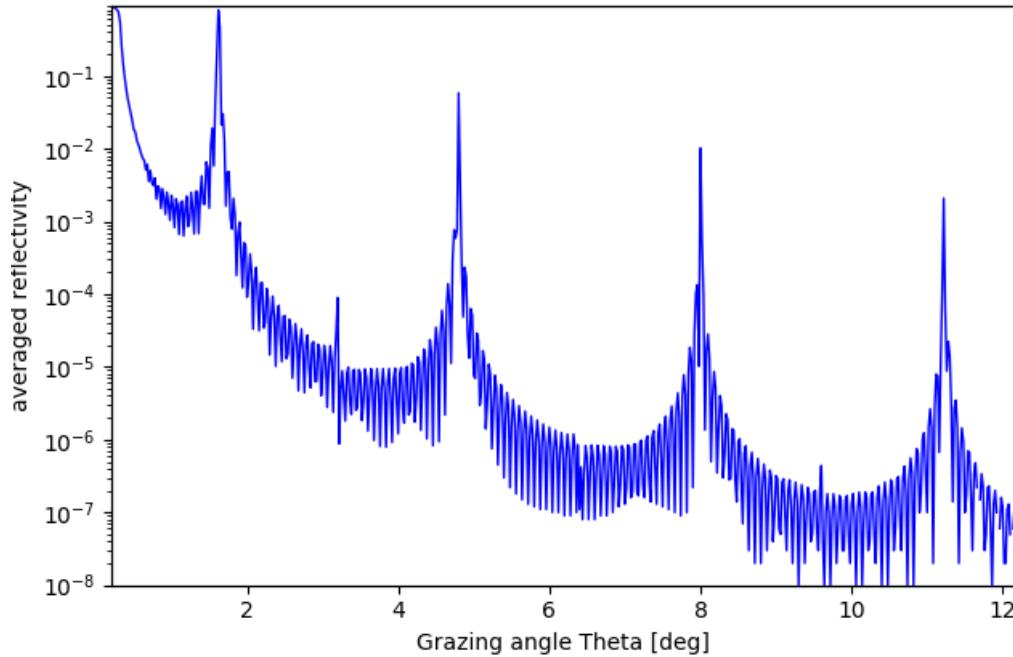
If $\theta = \pi/2$, and $m=1$:

$$\lambda = 2d$$

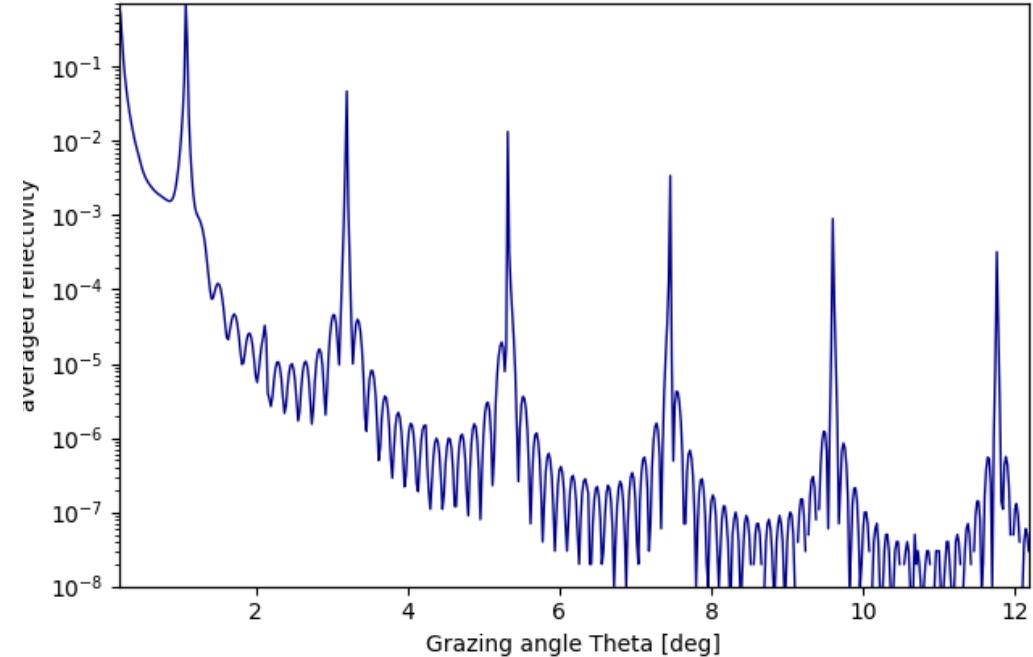
Multi-layer mirrors

W/C, $d=22.3 \text{ \AA}$, $\Gamma=0.5$, $N= 100$

$E=10\text{keV}$



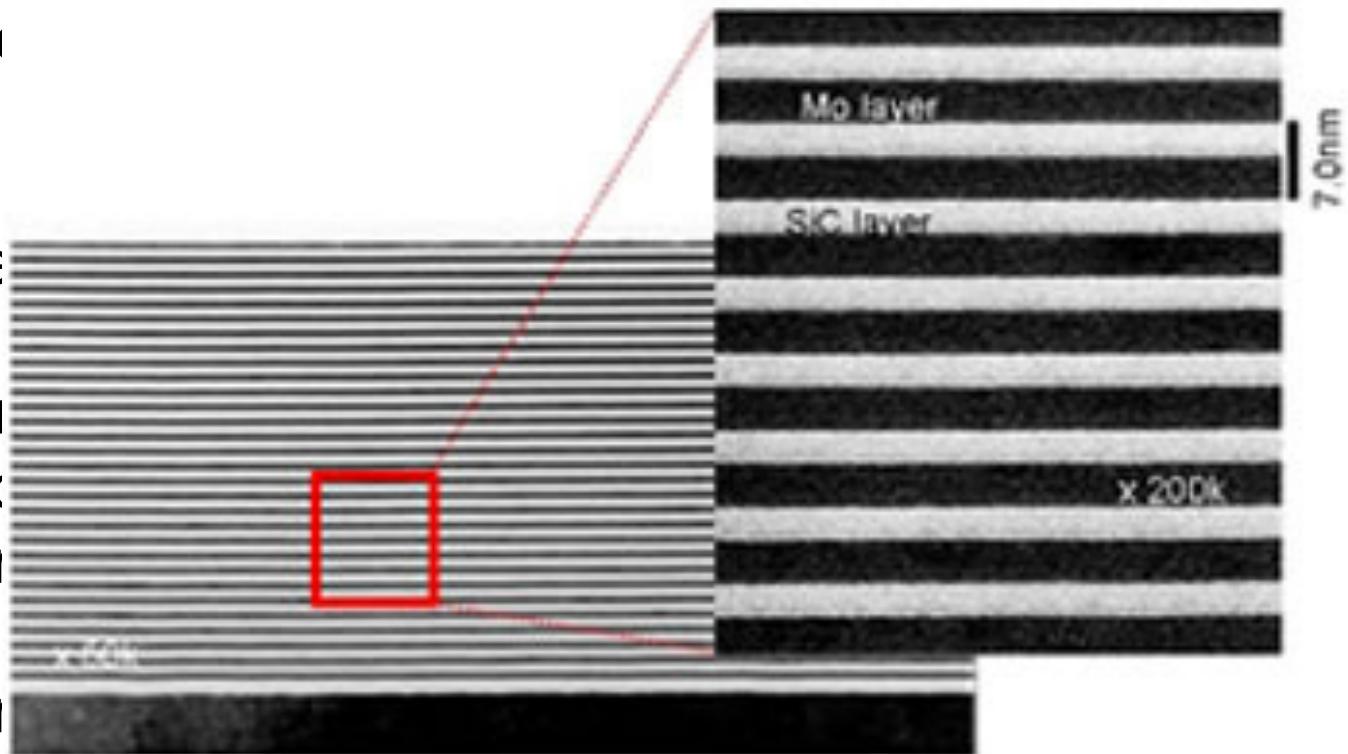
$E=15\text{keV}$



Making of Multi-layer mirrors

Typically by Magnetron Sputtering

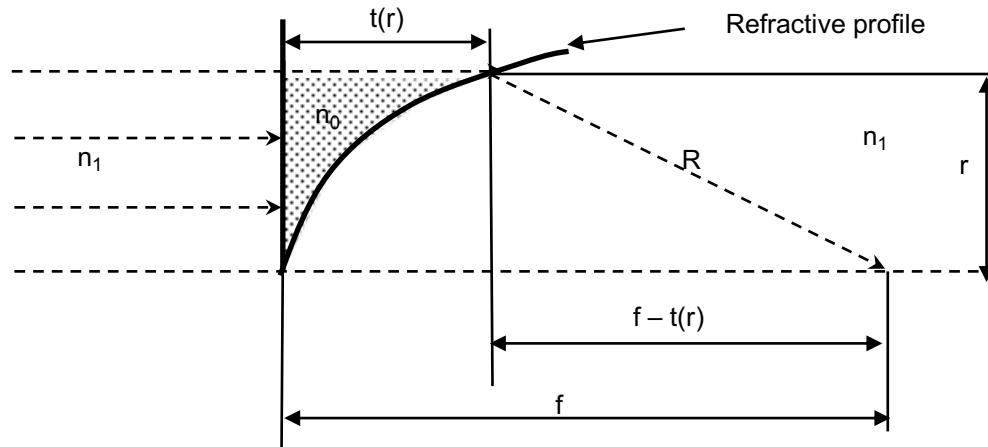
- Coating uniformity
- High refractive index
- Minimal absorption
- Chemically stable interfaces
- Minimal interdiffusion
- Minimal interfacial roughness
- Thermal stability
- Chemical stability in air



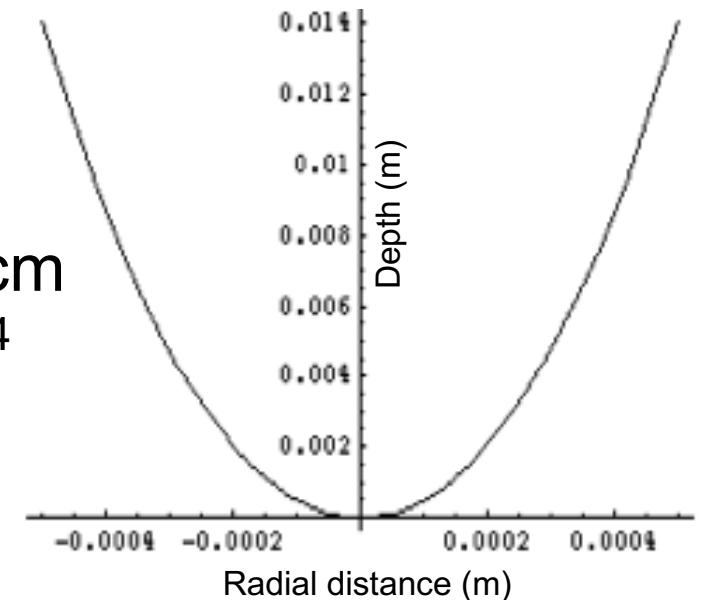
From NTT-AT website

Zone plates

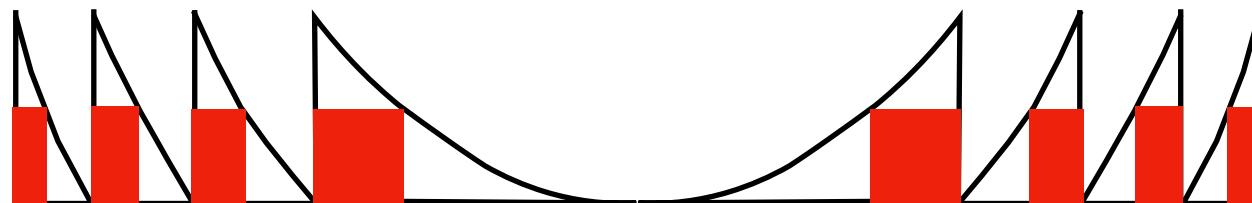
$$n = 1 - \delta + i\beta$$



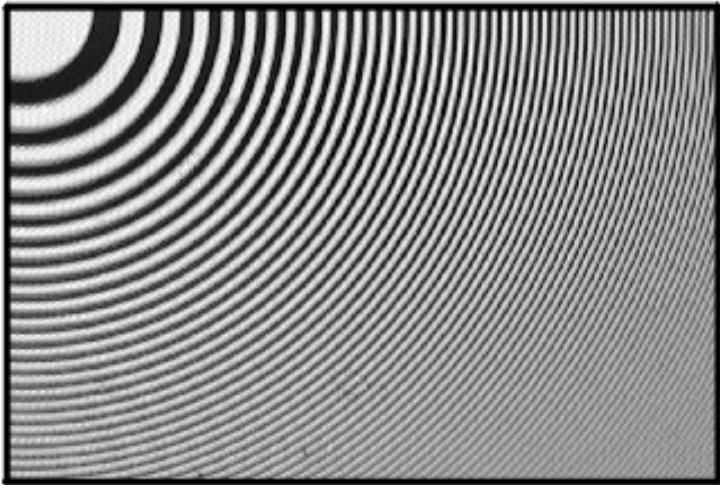
$$\begin{aligned} f &= 10 \text{ cm} \\ \delta &= 10^{-4} \end{aligned}$$



$$\Delta\phi = \delta \left(\frac{2\pi}{\lambda} \right) t = 2\pi$$

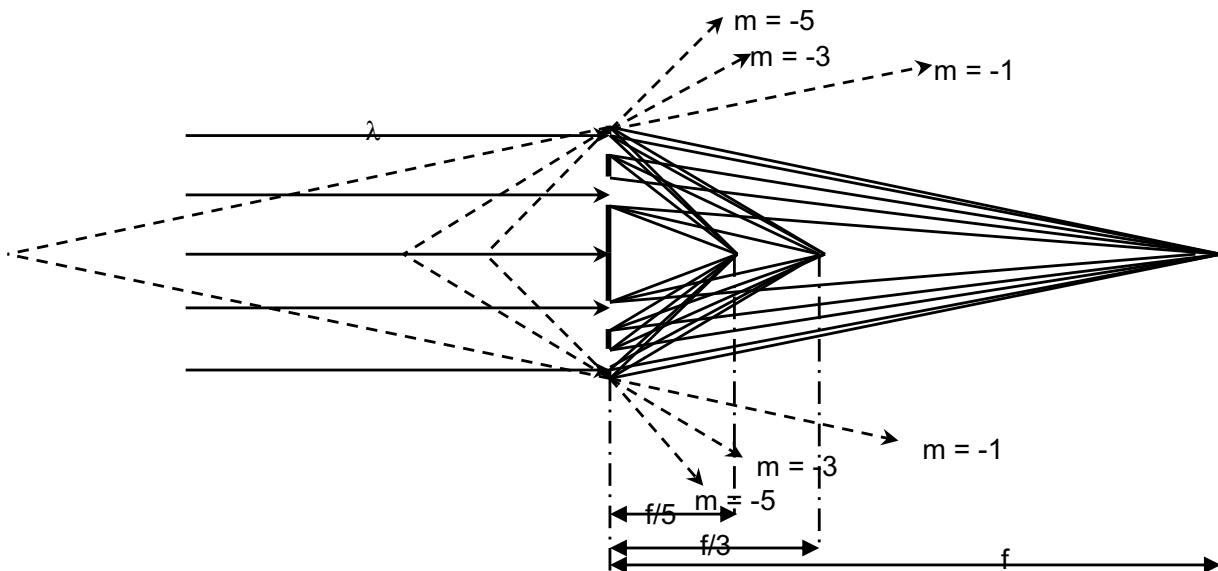


A couple of points on Zone Plates



A ZP is a circularly symmetric diffraction grating, with constant groove area, i.e. is periodic in r^2

$$\text{If } N > 100, \text{ then: } \frac{1}{f} = \frac{1}{p} + \frac{1}{q}$$

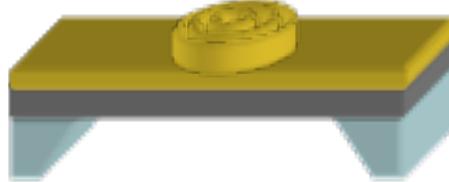
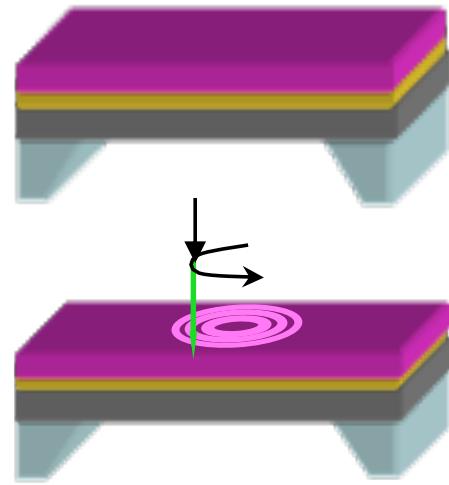


$$f_m = D \delta r_N / m \lambda$$

$$\delta_m = 1.22 \delta r_N / m$$

Making of Zone Plates

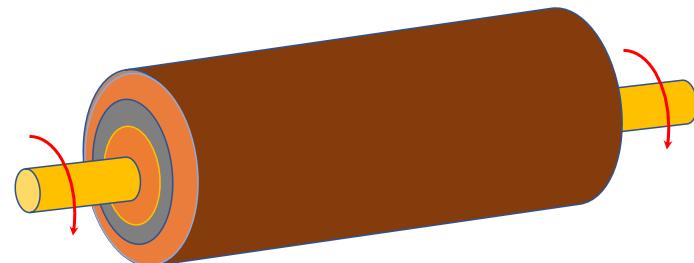
Electron beam lithography



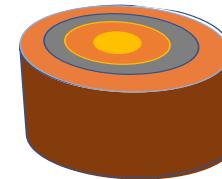
$D \sim 0.1$ to 1 mm, $\delta r_N = 6$ nm and up
 $t \sim 0.05$ to 1.5 μm

Sputtering and slicing

Cu/Al sputtering



Slicing and polishing



$t \sim 20$ -50 μm

$D \sim 80$ -100 μm



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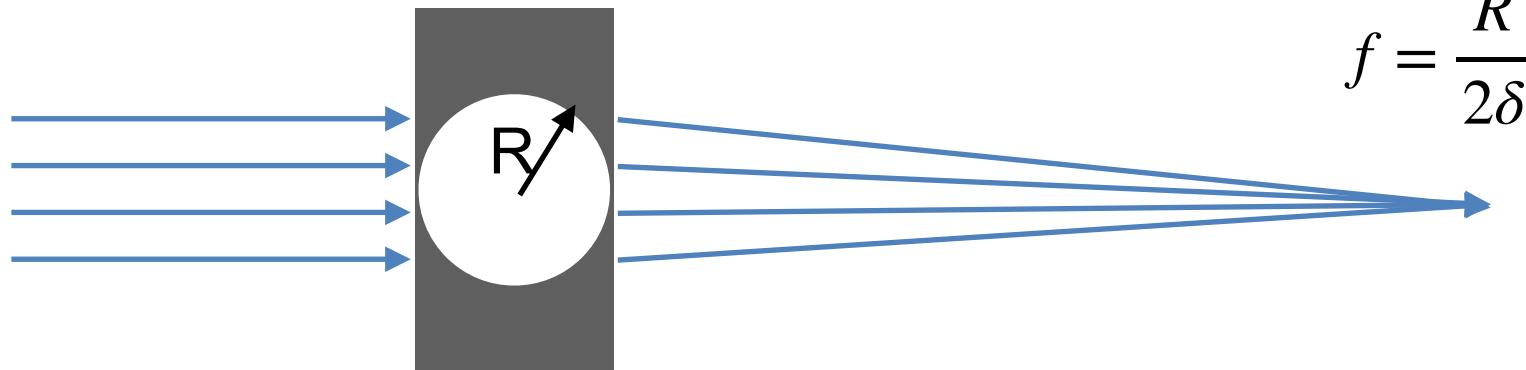
Refracting elements



XV School on Synchrotron Radiation:
Fundamentals, Methods and Applications

M. Altissimo, 16th September 2019

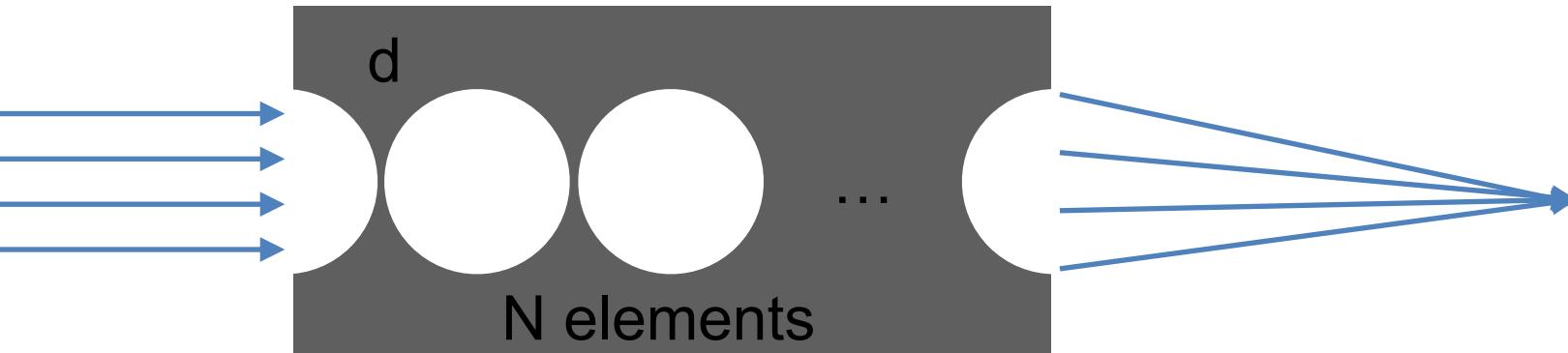
Compound refracting lenses



$$f = \frac{R}{2\delta}$$

$R \sim 100 \mu\text{m}$

$\delta \sim 10^{-5}$



$$f = \frac{R}{2N\delta}$$

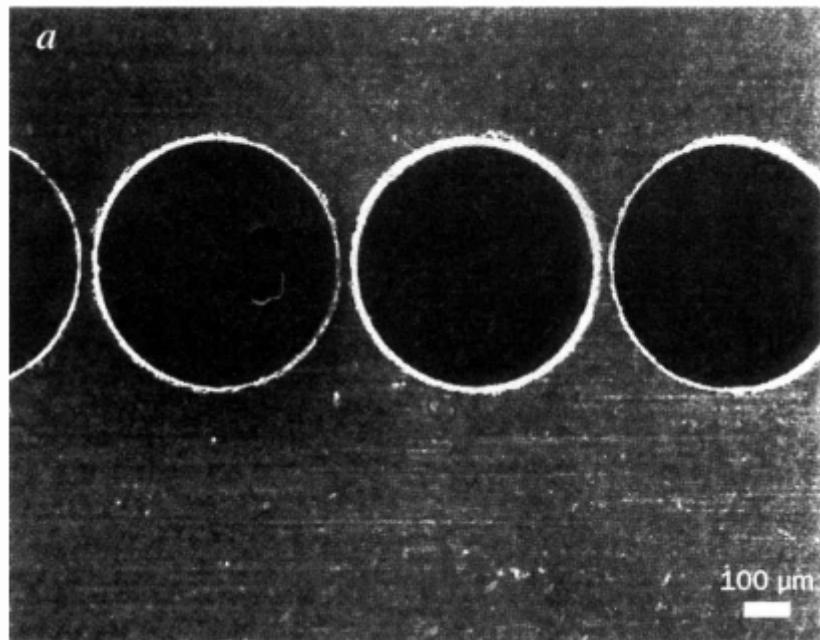
$$\sigma_{dl} \propto \frac{f\lambda}{A}$$

$$A = 2R\sqrt{\left(\frac{\lambda}{2\pi\beta RN}\right)}$$

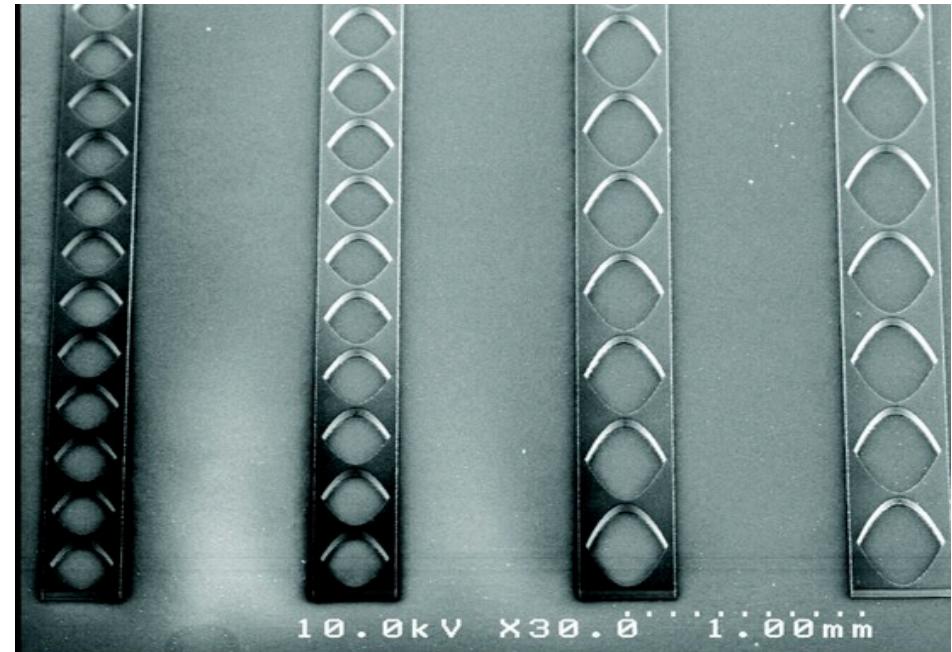
Compound refracting lenses

Characteristics:

- Must be made of low Z materials (Al, C, Be...)
- Focus only in 1D
- Paraboloid as ideal profile



Singirev A., et al, *Nature*, 384, (1996)



Alianelli, L., et al, *Journal of Applied Physics*. 108, 123107 (2010)



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Monochromator



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- Need for collimated illumination
- Crystal monochromators
- Plane grating monochromators



The need for collimated illumination

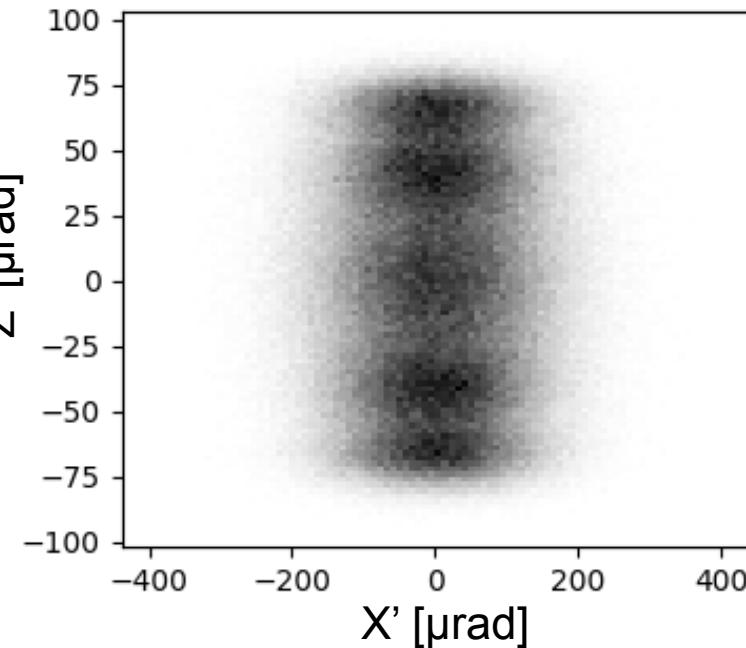
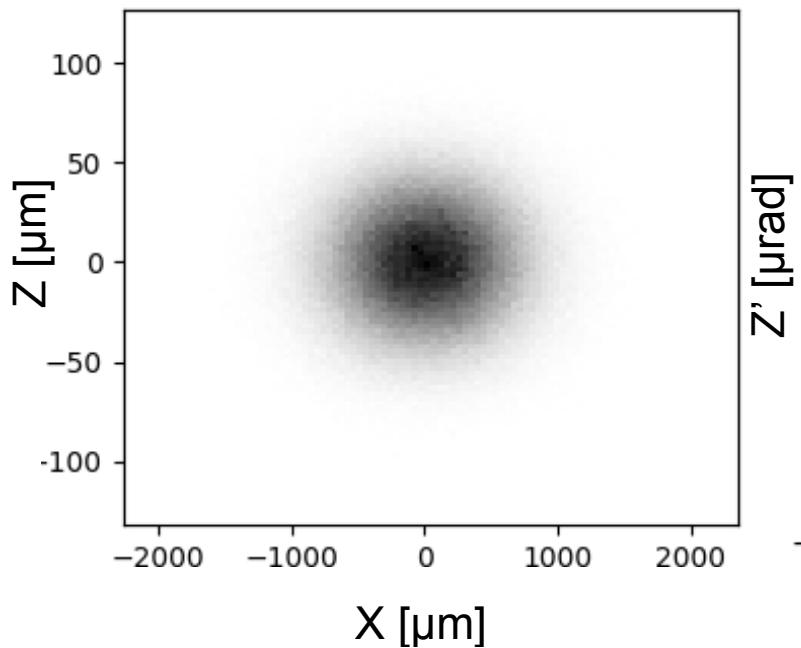
Crystals Energy resolution:

$$\frac{\Delta E}{E} = \frac{\Delta\lambda}{\lambda} = \boxed{\Delta\theta} \frac{\cos\theta}{\sin\theta}$$

Same for multilayers

Gratings

$$\sin\alpha + \sin\beta = Lm\lambda$$



Undulator

5th Harmonic (~1 keV)
 $\Delta E=500\text{eV}$

The need for collimated illumination

Collimating mirror before monochromator

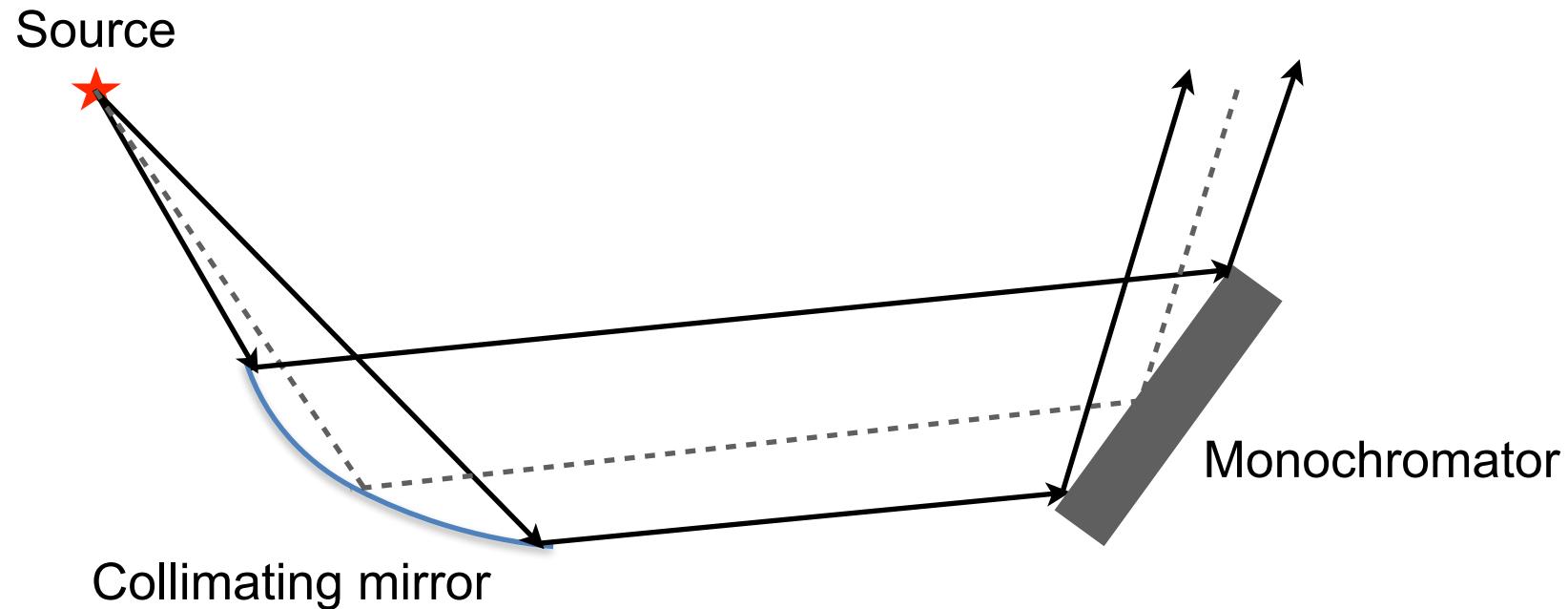


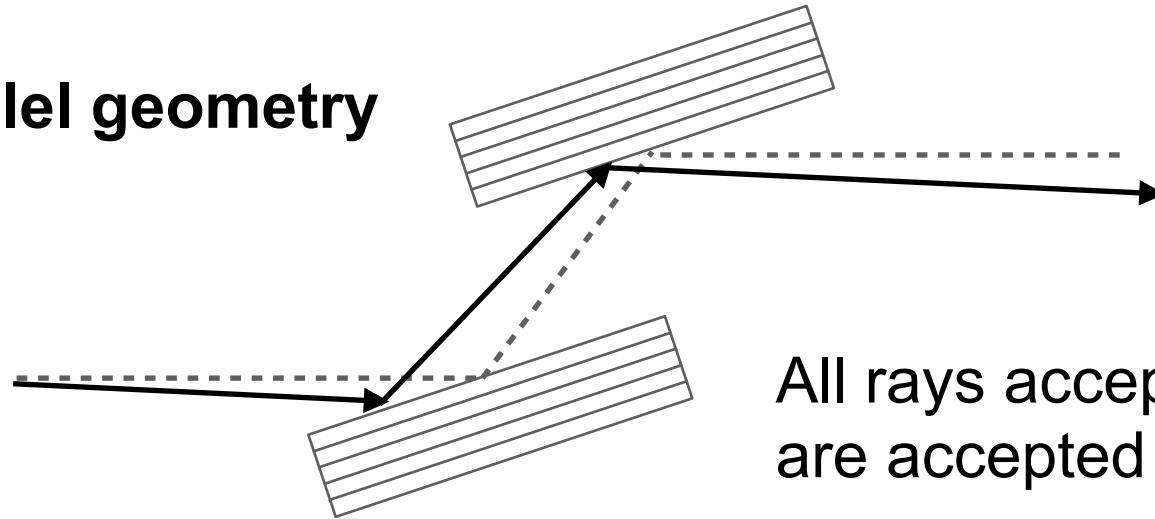
Figure calculated setting r' negative and very far ($\sim 100s$ m)

Double Crystal monochromator

$$2dsin\theta = m\lambda$$

$$\frac{\Delta E}{E} = \frac{\Delta\lambda}{\lambda} = \Delta\theta \frac{cos\theta}{sin\theta}$$

Parallel geometry



All rays accepted by first crystal
are accepted also by the second

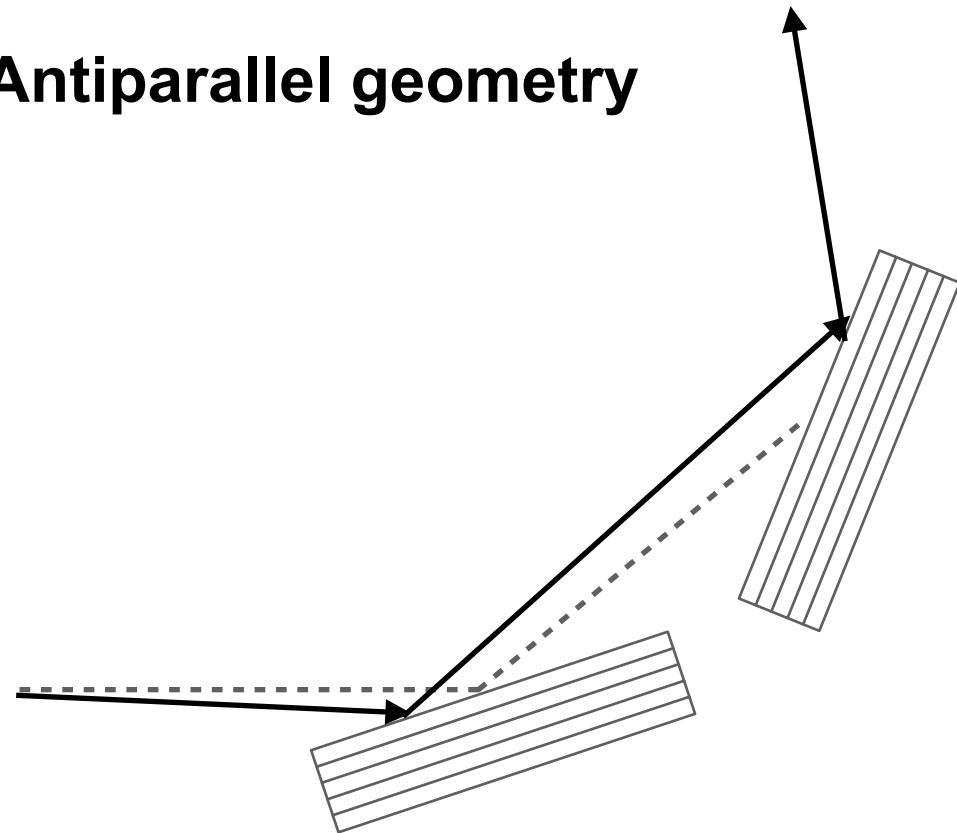
Second crystal acts merely as a mirror

Double Crystal monochromator

$$2dsin\theta = m\lambda$$

$$\frac{\Delta E}{E} = \frac{\Delta\lambda}{\lambda} = \Delta\theta \frac{cos\theta}{sin\theta}$$

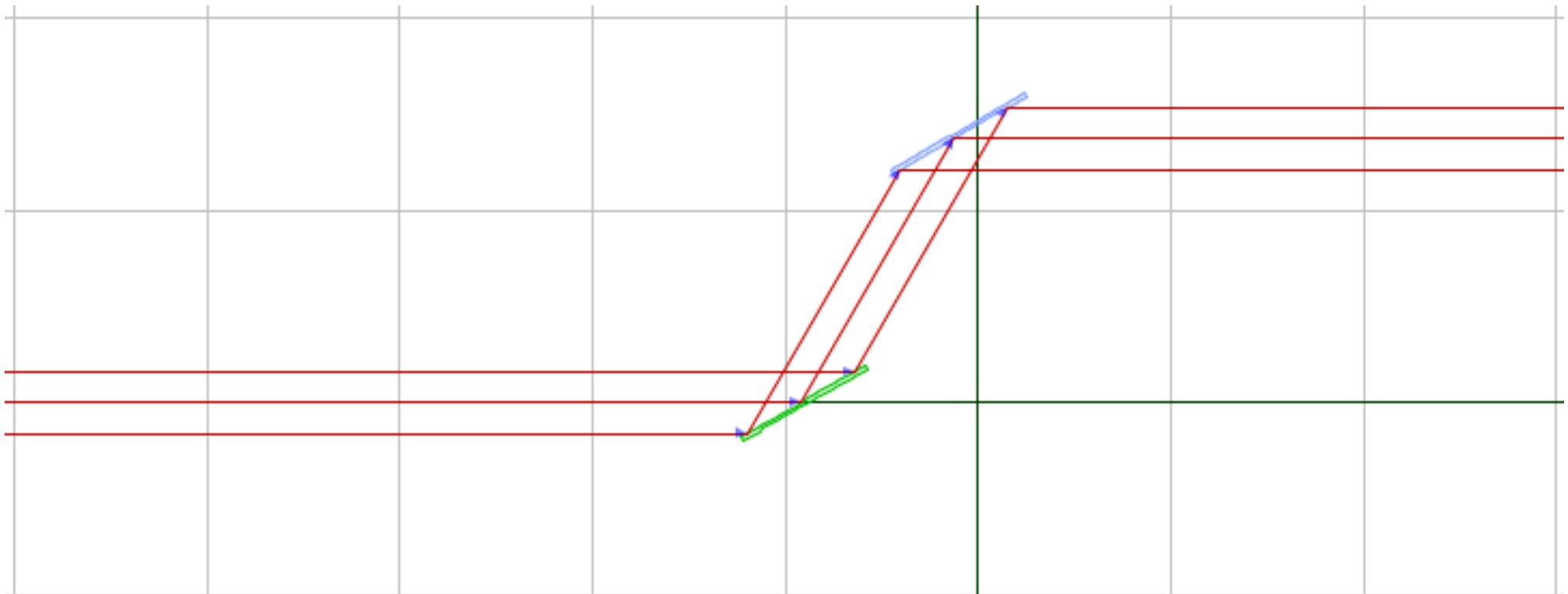
Antiparallel geometry



Second crystal is dispersive

Higher $\Delta E/E$
Lower flux

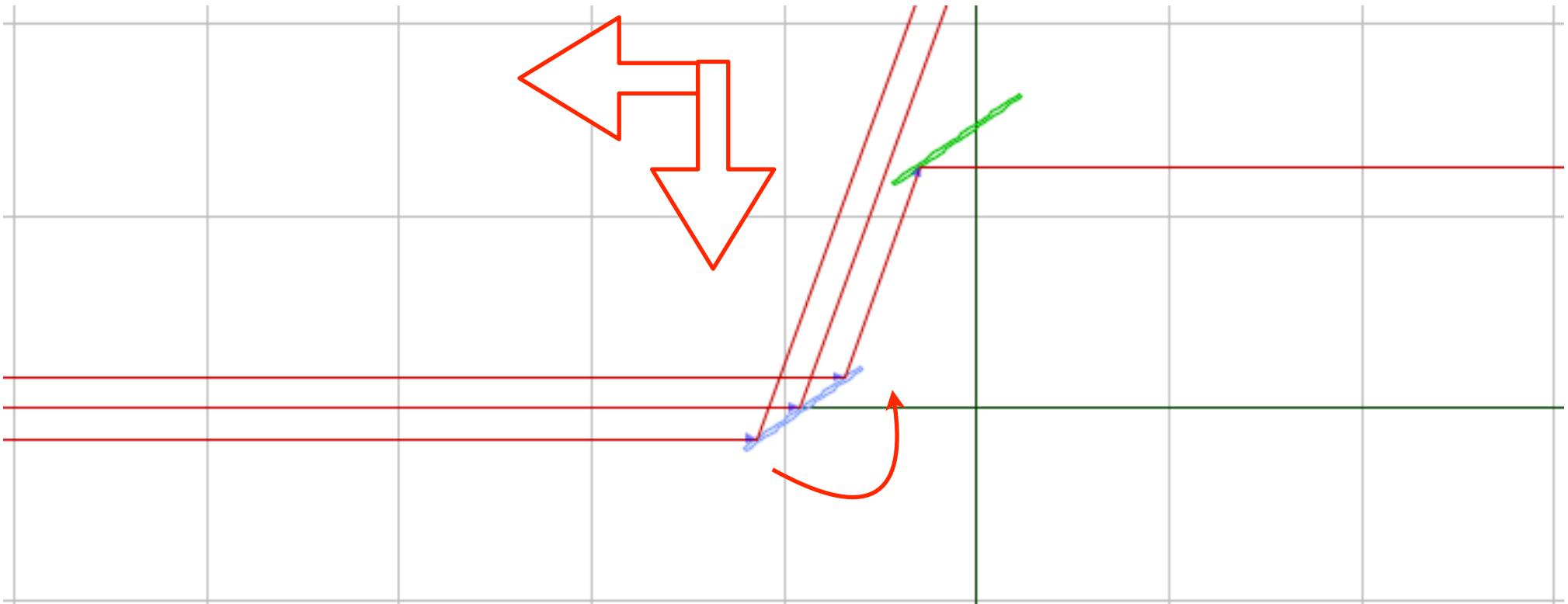
DCM in parallel configuration





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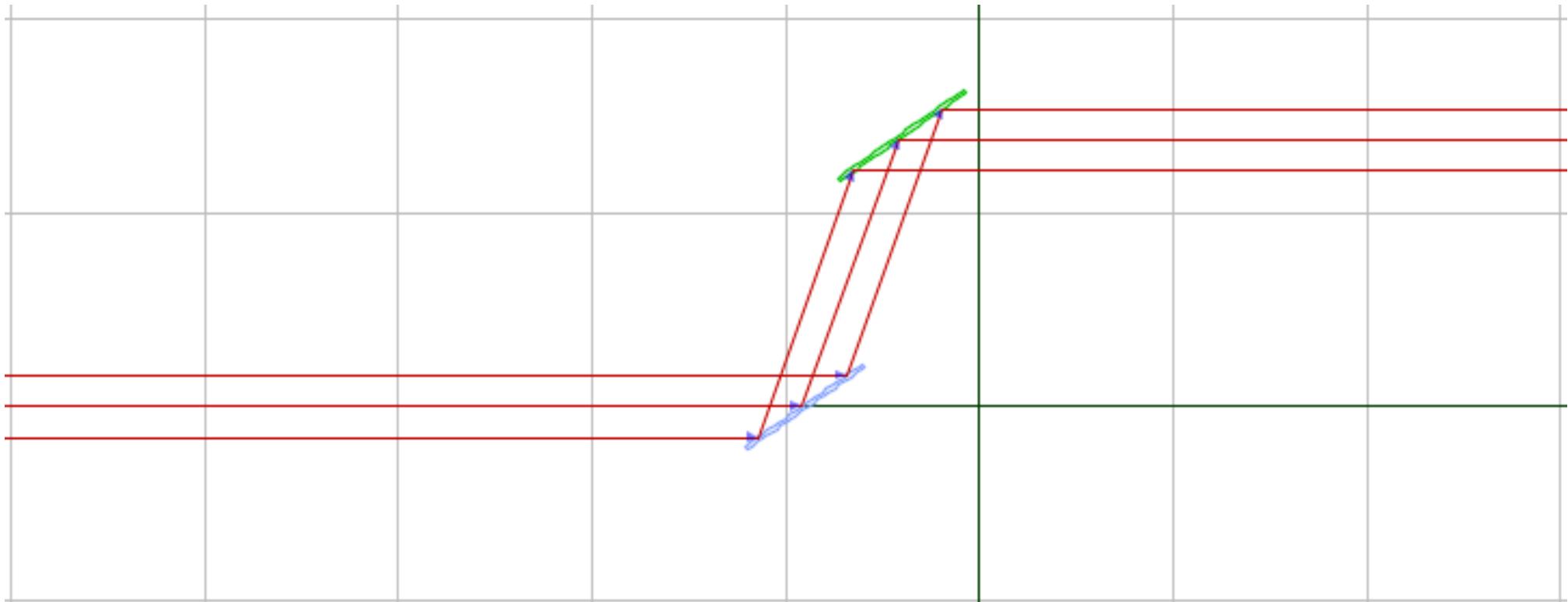
DCM in parallel configuration



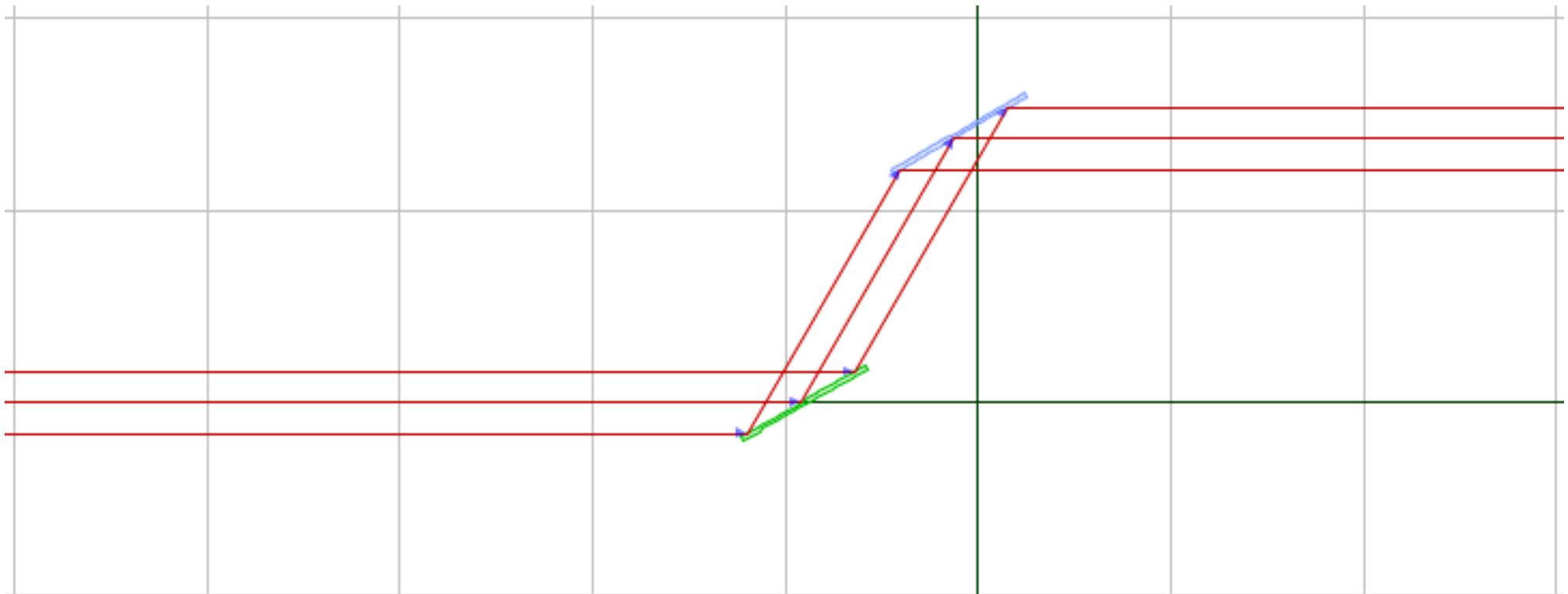
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DCM in parallel configuration



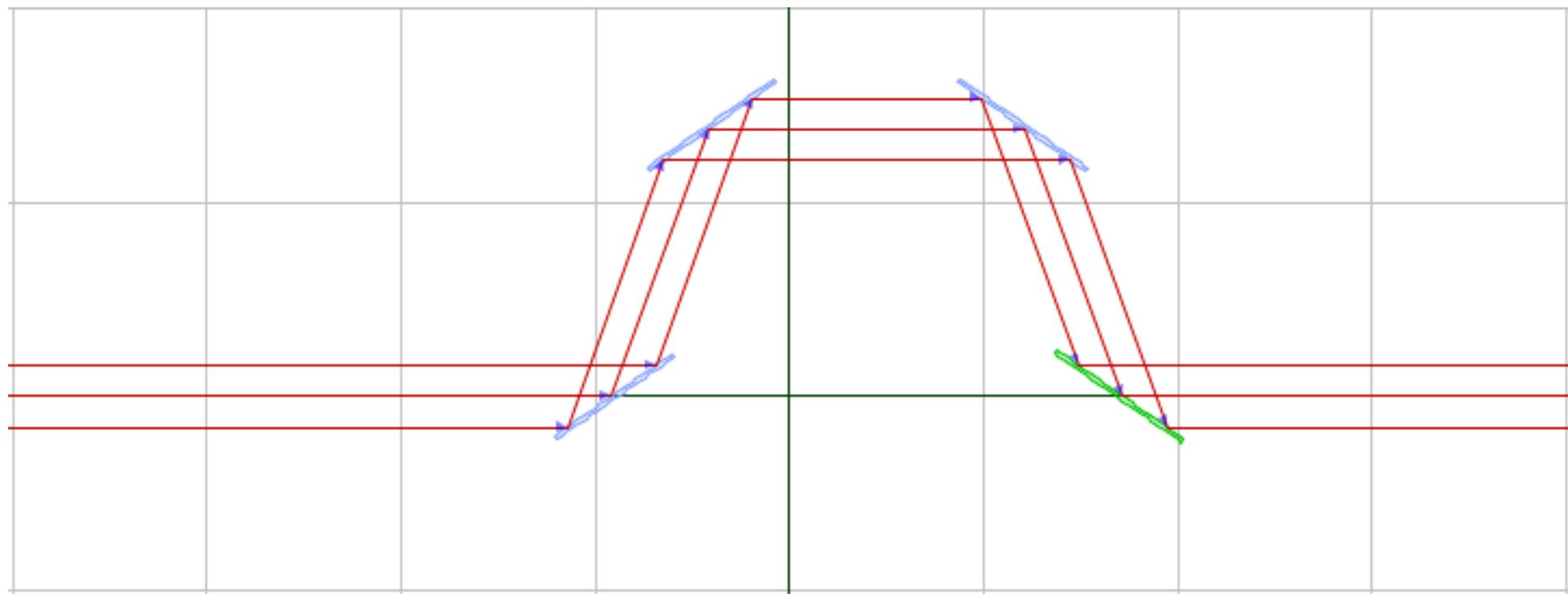
DCM in parallel configuration





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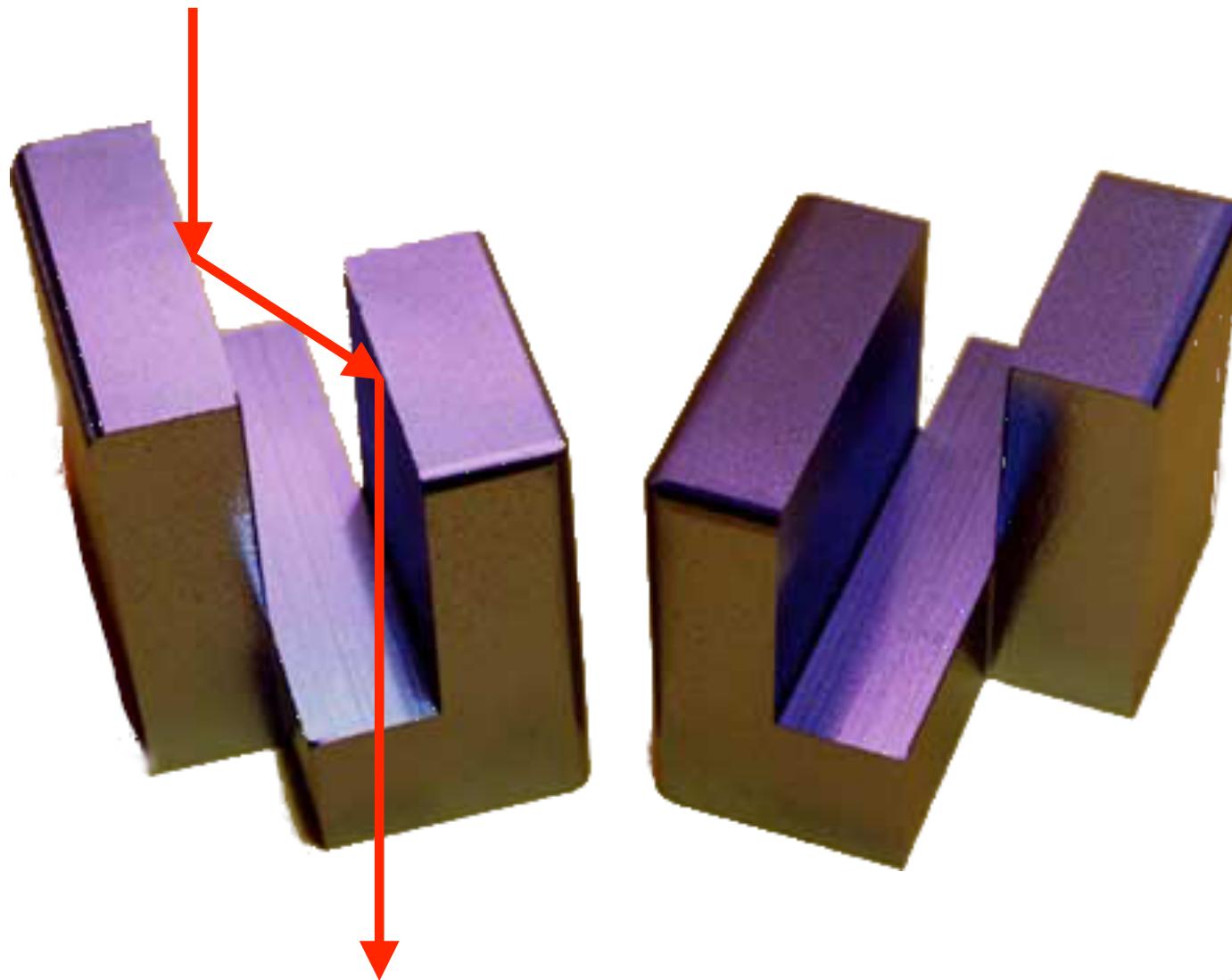
2 X DCM in parallel configuration





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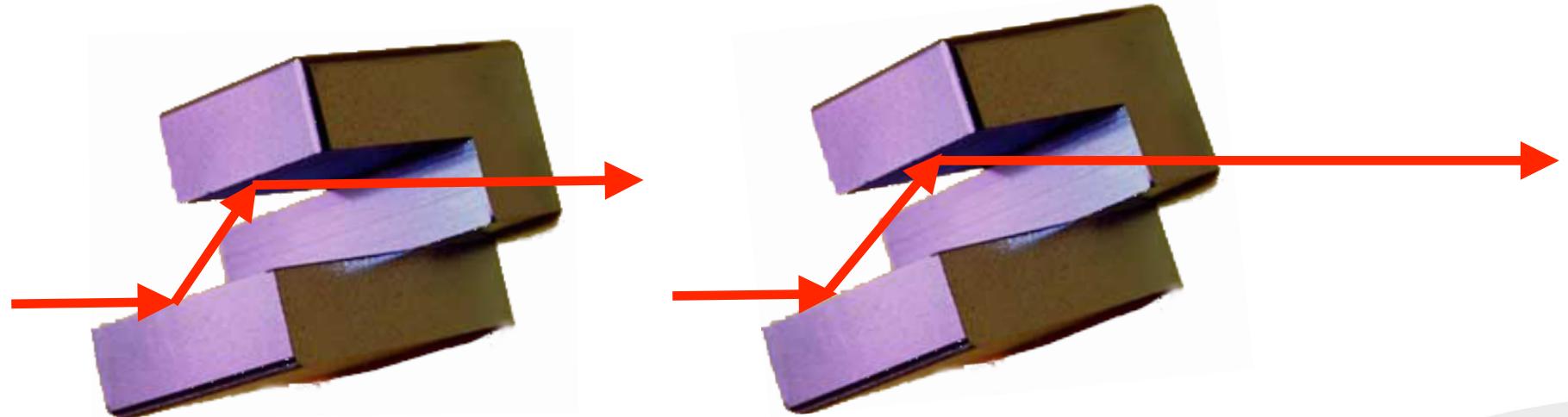
Channel cut Si monochromator



Channel cut Si monochromator

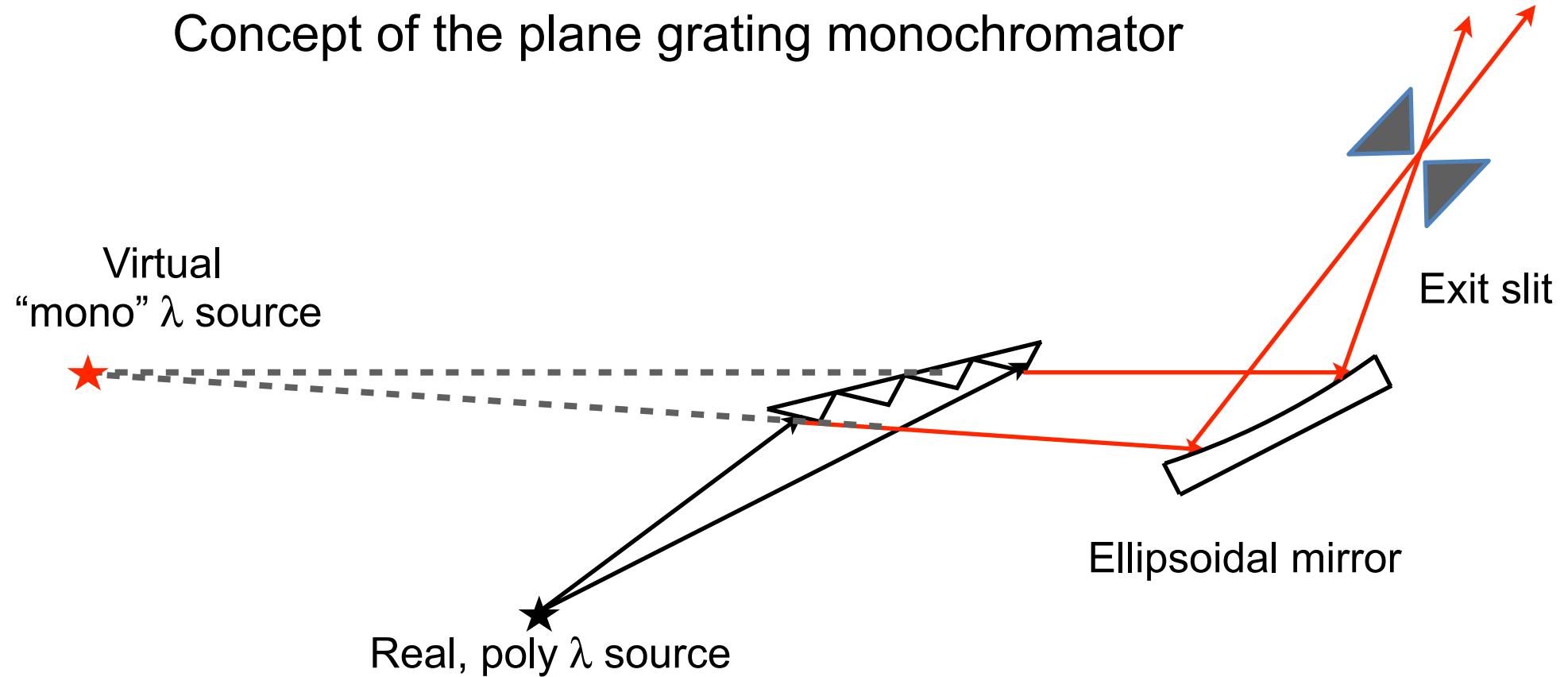
Easy to fabricate
Easy to align

Beam position not maintained if θ changes
Limited angular range



Plane grating monochromator

Concept of the plane grating monochromator



H. Petersen. O. Communication, vol. 40, no. 6. 1982, pp. 402–406.

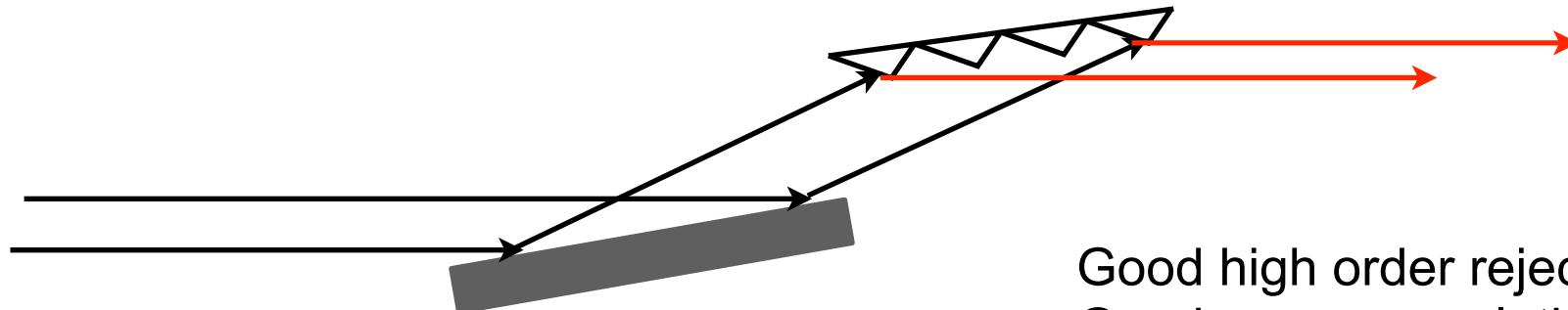
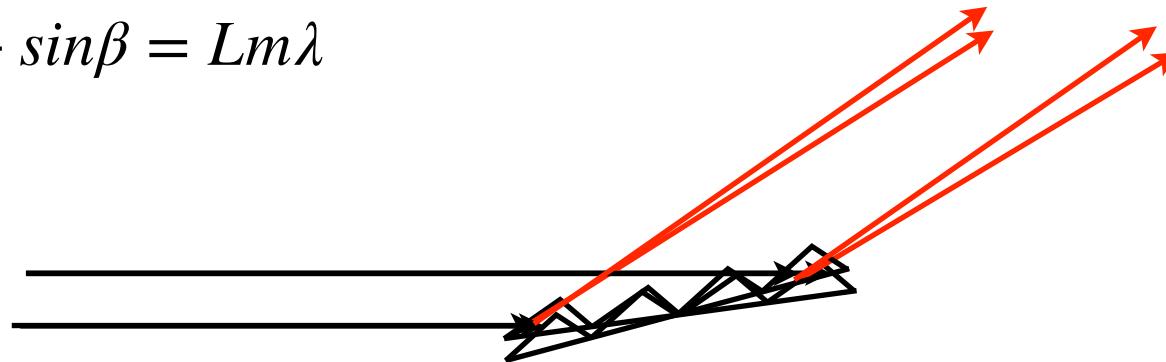


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Plane grating monochromator

$$\sin\alpha + \sin\beta = Lm\lambda$$



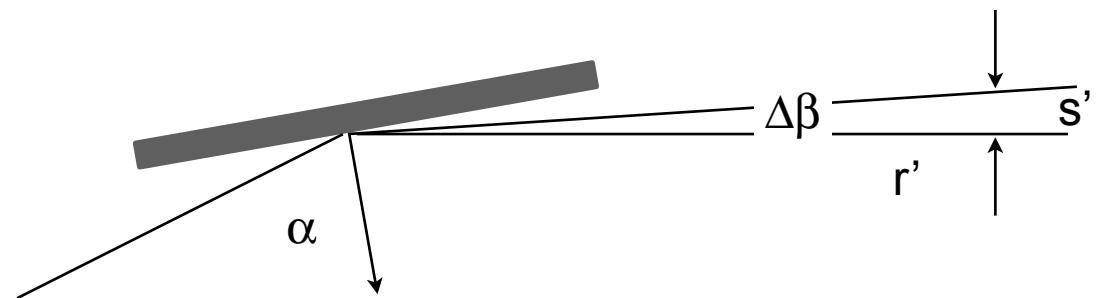
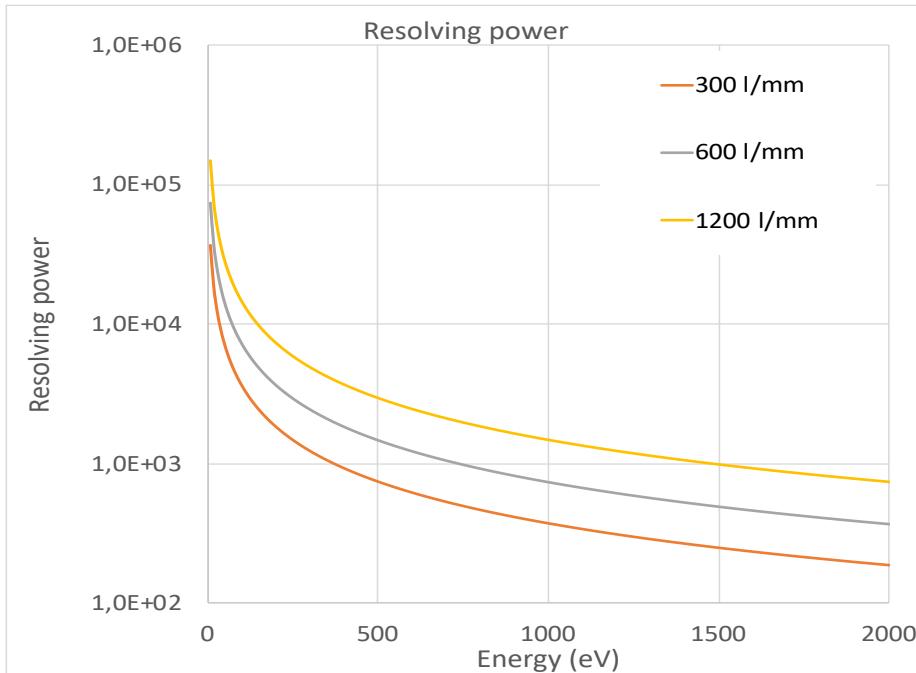
Plane pre-mirror

Good high order rejection
Good energy resolution
Fixed exit slit

Grating resolving power

Angular dispersion of a grating with line density L: $\Delta\lambda = \frac{\cos\beta}{Lm} \Delta\beta$

Resolving power R: $R = \frac{E}{\Delta E} = \frac{\lambda}{\Delta\lambda} = \frac{\lambda Lmr'}{s' \cos\beta}$



If $R=1000$, @ 100eV:
 $\Delta E = 100meV$

... finally a couple of examples



TwinMic Beamline @ Elettra

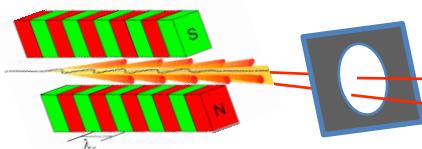
Photon energy: 400eV to 2 keV
X-ray microscopy and microFluorescence

Undulator source

N=17

Period=56mm

Min gap: 25 mm



Cylindrical mirror

16 m from source

Au coated

$\rho=0.56\text{m}$

$\theta=1\text{ deg}$

Toroidal mirror

20 m from source

Au Coated

$R=211\text{m}$

$\rho=0.07\text{m}$

$\theta=1\text{ deg}$

Zone Plate

$D=600\text{ }\mu\text{m}$

Res 50 nm

24 m from source

Entrance slit

12.6 m from source

$D=1\text{mm}$

PGM with plane pre-mirror

18 m from source

Blazed grating, $\theta_b=1.1\text{ deg}$

Au coated

$L= 600 \text{ l/mm}$

Secondary source

22 m from source

5-100 μm diameter

Experiment

$\sim 24.5\text{ m from source}$

Focal spot dia:

100 nm to 1.5 μm

Diffraction Beamline @ Elettra

Photon energy: 4 to 21 keV

Cylindrical mirror for vertical collimation

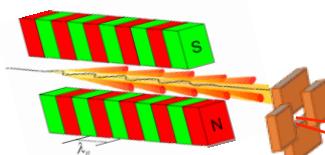
22 m from source

Pt Coated

R=14km

$\theta=0.172$ deg

Length: 1.4 m



Multi-pole wiggler

N=54, 1.5T mag field

Period=140mm

Critical Energy: 5.8keV @ 2.4GeV

5kW total power @ 140 mA



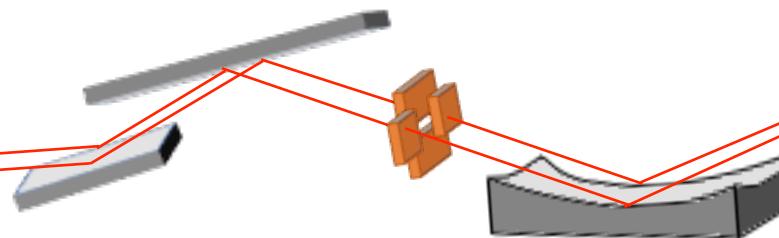
Toroidal focussing mirror

28 m from source

Sagittally cylindrical, bendable

R=9km (5km to ∞)

$p=0.055$ m



Experiment

41.5 m from source

Focal spot:

0.7×0.2 mm²

$\Delta E/E \sim 4000$

Double crystal mono

24 m from source

Si(111), $\omega_s=25$ μ rad @ 8keV

Thank you!

