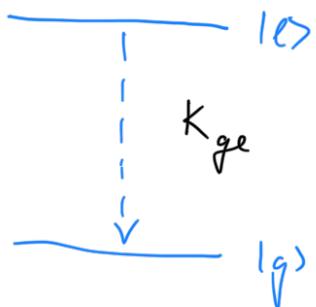
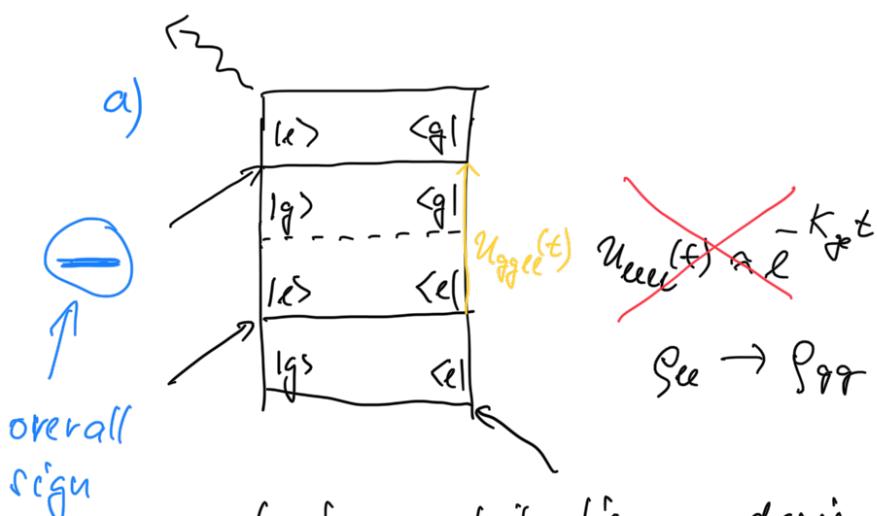


Liouville pathways with energy relaxation processes

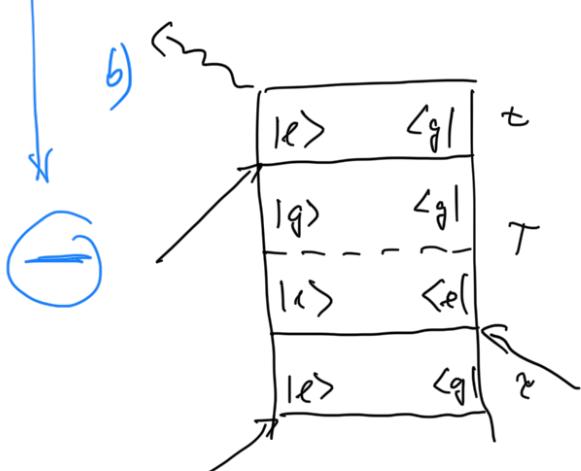


$$U_{\text{geee}}(t) \quad \rho_{ee}(t) \rightarrow \rho_{gg}(t)$$

$$\frac{1}{1 - e^{-K_{ge}t}} = \int_0^t dz K_{ge} U_{\text{eeee}}(z)$$



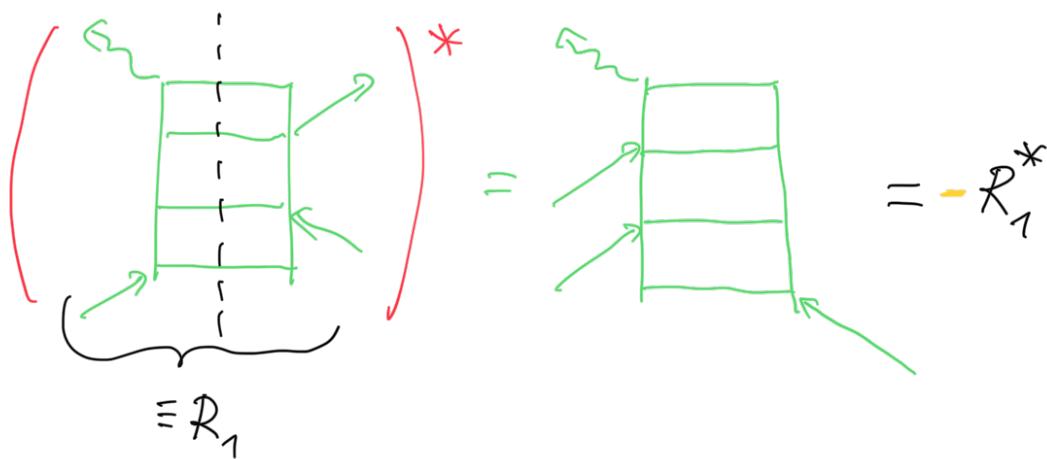
rephasing contribution - derived from R_2



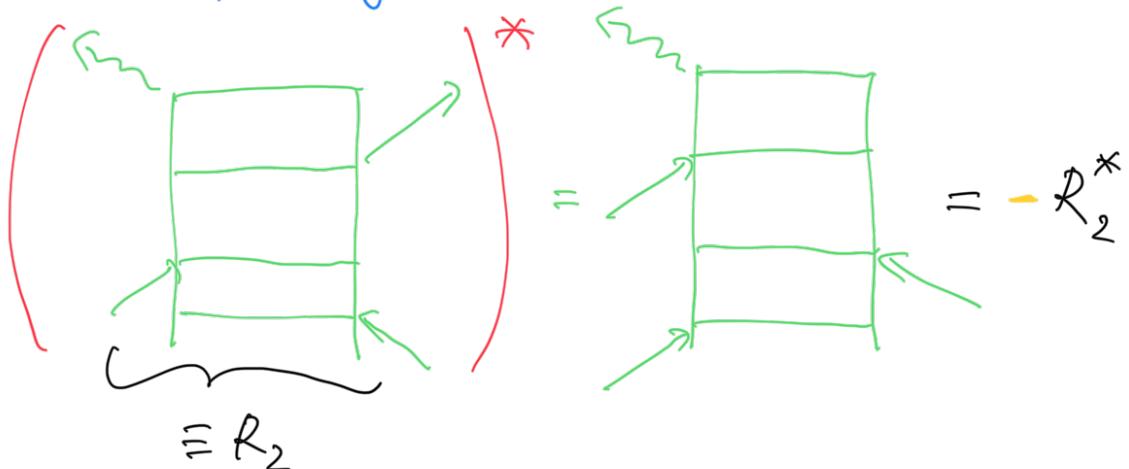
non-rephasing contribution - derived from R_1 pathway

The above diagrams must correspond to R_1^*, R_2^*, R_j^* or R_4^* .

Identification of pathway a)



—(i) of pathway b)



Summary of Liouville pathways for relaxing 2Ls

$$R_1(t, T, \tau) = d^4 \mathcal{U}_{egeg}(t) \mathcal{U}_{eeee}(T) \mathcal{U}_{egeg}(\tau)$$

$$R_2(t, T, \tau) = d^4 \mathcal{U}_{ggfg}(t) \mathcal{U}_{eeee}(T) \mathcal{U}_{gege}(\tau)$$

$$R_3(t, T, \tau) = d^4 \mathcal{U}_{egeg}(t) \mathcal{U}_{gggg}(T) \mathcal{U}_{gege}(\tau)$$

$$R_4(t, T, \tau) = d^4 \mathcal{U}_{egeg}(t) \mathcal{U}_{gggg}(T) \mathcal{U}_{egeg}(\tau)$$

$$R_1^*(t, T, \tau) = d^4 \mathcal{U}_{egeg}(t) \underline{\mathcal{U}_{ggge}(T)} \mathcal{U}_{gege}(\tau)$$

$$R_2^*(t, T, \tau) = d^4 \mathcal{U}_{egeg}(t) \underline{\mathcal{U}_{ggee}(T)} \mathcal{U}_{egeg}(\tau)$$

$\mathcal{F}\mathcal{T}[\mathcal{U}_{ggge}(T)] = G_{eg}(\omega)$

Pump-probe

$$\Delta A(\omega) \approx -\omega \operatorname{Re} (R_1(\omega; T, \tau=0) + R_2 + R_3 + R_4 - R_1^* - R_2^*)$$

$$\boxed{\Delta A(T=0) \approx -\omega d^4 \operatorname{Re} [G_{eg}(\omega) + G_{eq}(\omega) + G_{eg}(\omega) + G_{eq}(\omega)]}$$

$$-\overset{0}{\underset{R_1^*}{\uparrow}} \quad -\overset{0}{\underset{R_2^*}{\uparrow}}$$

$$\approx -4d^4 \chi_{eg}(\omega)$$

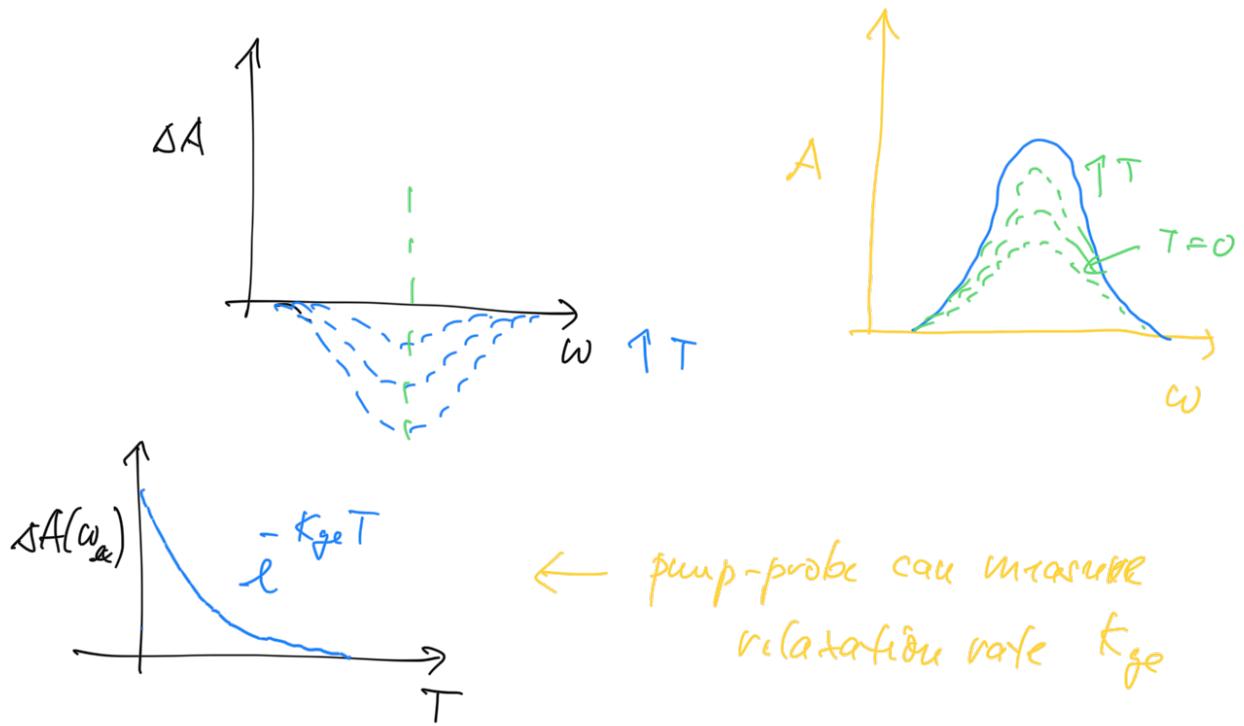
$$u_{ee}(T)$$

$$\begin{aligned} \Delta A(T) &\approx -\omega d^4 \operatorname{Re} \left[(G_{eg}(\omega) + G_{eq}(\omega)) e^{-K_{ge}T} \right. \\ &\quad + (G_{eg}(\omega) + G_{eq}(\omega)) \cdot 1 \\ &\quad \left. - (G_{eg}(\omega) + G_{eq}(\omega)) (1 - e^{-K_{ge}T}) \right] \end{aligned}$$

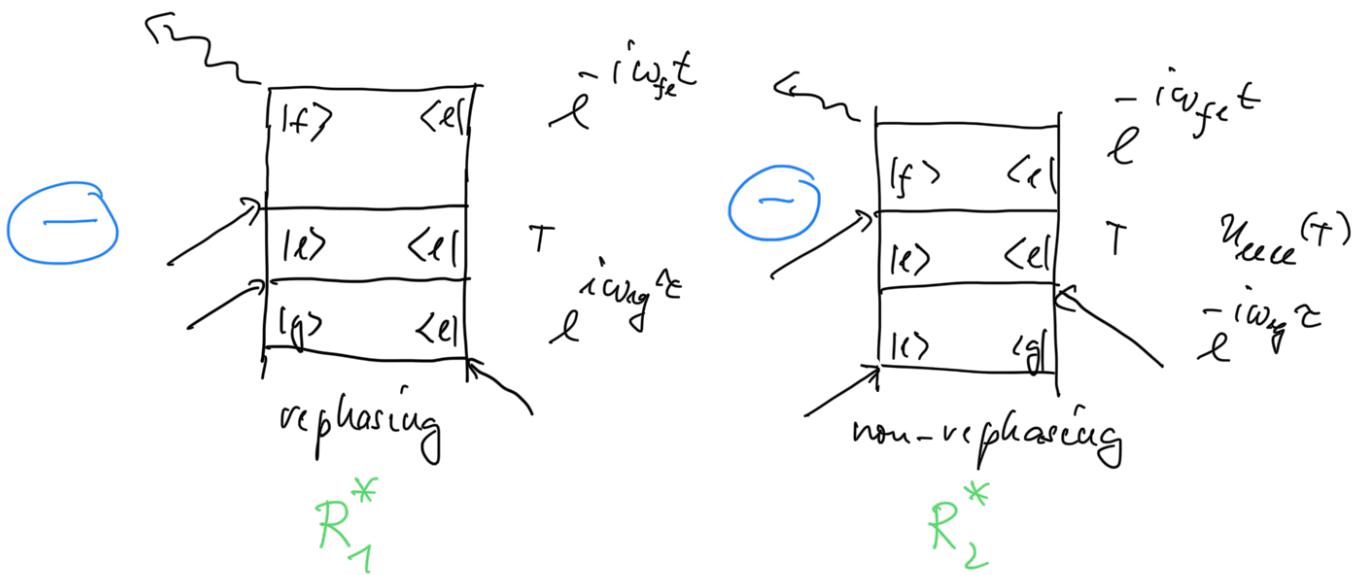
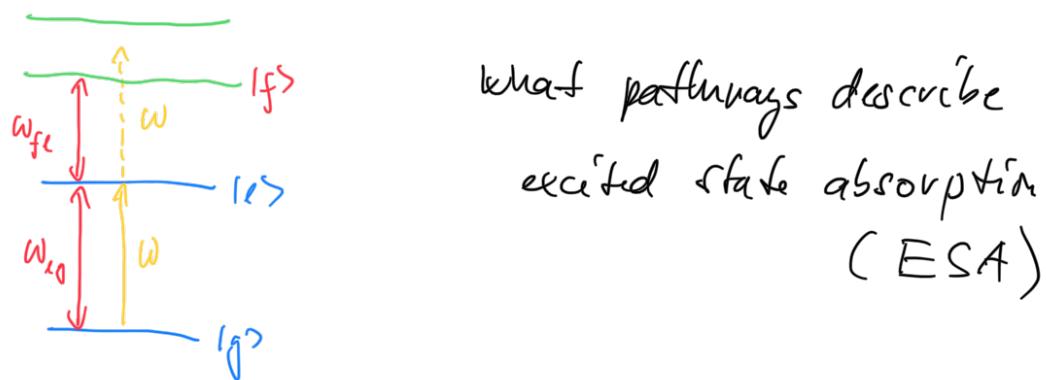
$$u_{ggg}(T) \quad u_{gge}(T)$$

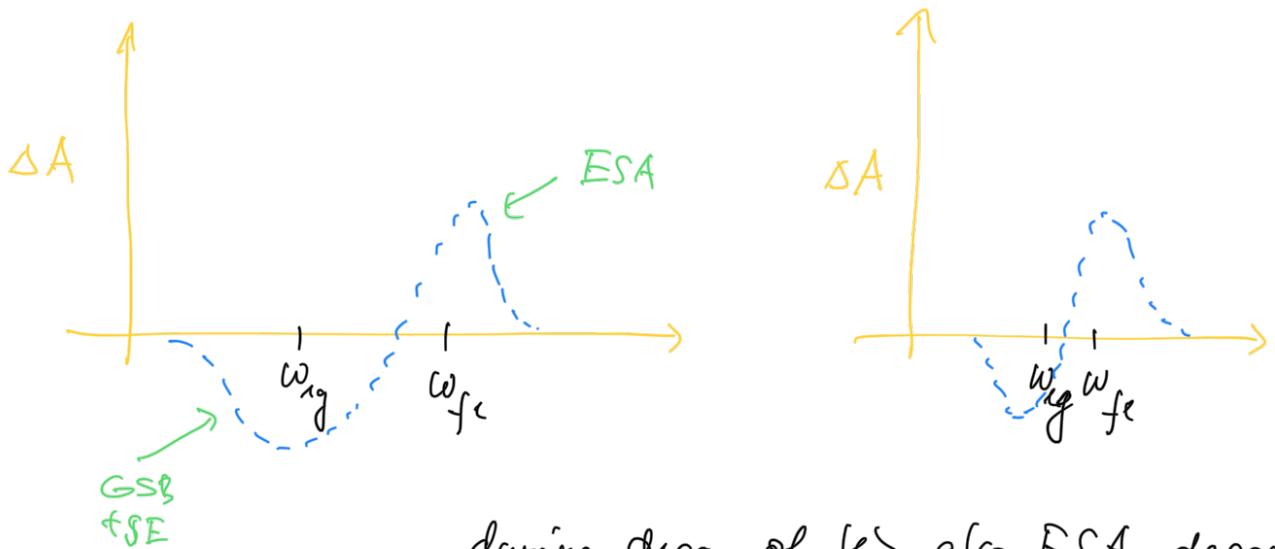
$$\boxed{\Delta A(T) \approx -\omega d^4 \operatorname{Re} [4 G_{eg}(\omega)] e^{-K_{ge}T}}$$

$$\Delta A(T \rightarrow \infty) \approx 0$$

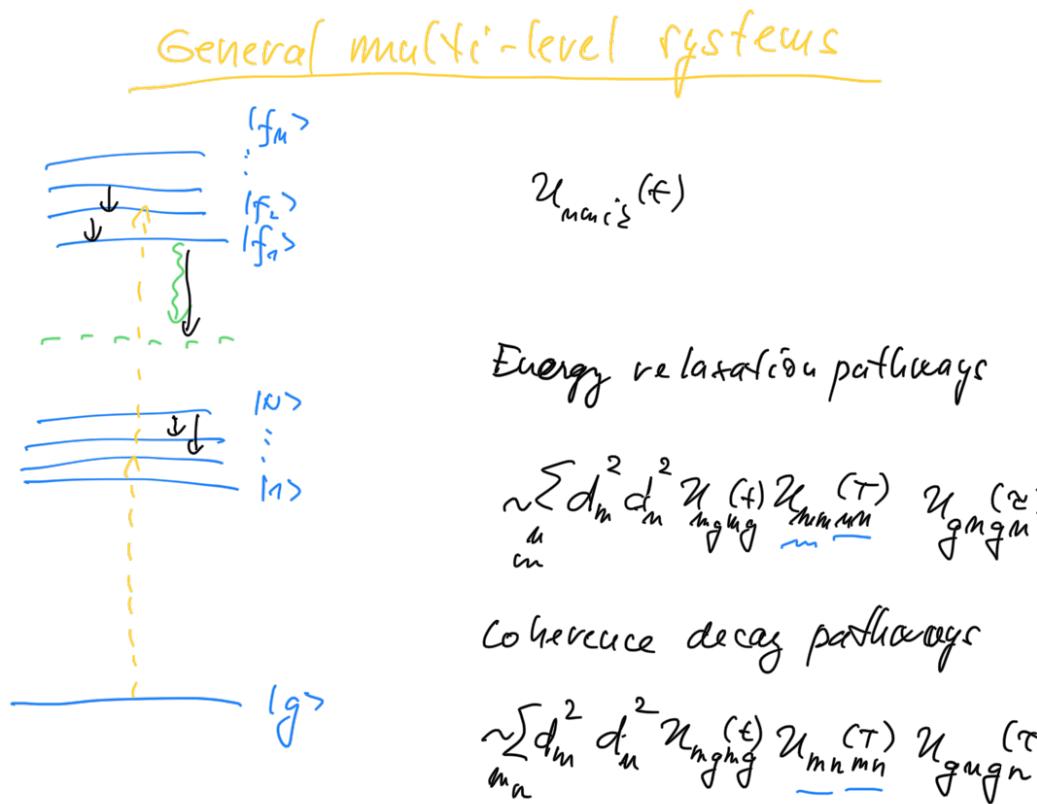


Systems with excited state absorption





dancing decay of (es) also ESA decay
 $\sim e^{-K_{fe} t}$



$U_{f_m \in f_n}(t) \rightarrow$ [auto] ESA

we assume secular approximation

How to calculate spectra

Absorption spectrum

$$\alpha(\omega) \approx \omega \sum_m \langle \vec{d}_{mg} \cdot \vec{e} \rangle \operatorname{Re} G_{mg}(\omega)$$



$$G_{mg}(\omega) = \int_0^\infty dt \Theta(t) U_{mgmg}(\epsilon) e^{i\omega t}$$

$$= \int_0^\infty dt \tilde{U}_{mgmg}(\epsilon) e^{-i\omega_{mg}t} e^{i\omega t}$$

Pump-probe spectrum

$$\Delta A(\omega) \approx -\omega \sum_{\substack{n m \\ \text{exc}}} \langle \vec{d}_{g\epsilon} \cdot \vec{e}_g \rangle \langle \vec{d}_{g\epsilon} \cdot \vec{e}_g \rangle \langle \vec{d}_{mg} \cdot \vec{e}_m \rangle \rangle_S$$

$SE \sim R_2 + R_1$

$$\times \operatorname{Re} \left\{ 2 G_{eg}(\omega) U_{ee}^{(T)} (\delta_{nm} \delta_{ke} + (1-\delta_{nm}) \delta_{ek} \delta_{km}) + \right.$$

$$\overbrace{\begin{aligned} & \text{FT}[U_{kgkg}^{(t)}] & GCB \sim R_3 + R_F \\ & + 2 G_{eg}(\omega) \delta_{nm} \delta_{ke} \end{aligned}}^{\text{GCB}}$$

$$\left. - 2 G_{eg}(\omega) U_{gggh}^{(T)} \delta_{nm} \delta_{ke} \right\}$$

GCB filling $\sqrt{R_1^* + R_2^*}$ with
related to (g)

$$+ \omega \sum_{\substack{n m \\ f k l}} \langle \vec{d}_{f\epsilon} \cdot \vec{e}_f \rangle \langle \vec{d}_{f\epsilon} \cdot \vec{e}_f \rangle \langle \vec{d}_{mg} \cdot \vec{e}_m \rangle \langle \vec{d}_{gm} \cdot \vec{e}_l \rangle \rangle_S$$

$$\times 2 \operatorname{Re} \left\{ G_{fe}(\omega) U_{ek}^{(T)} [\delta_{nm} \delta_{ke} + (1-\delta_{nm}) \delta_{en} \delta_{km}] \right\}$$

\uparrow
 ESA

Rotational averaging over transition dipole moment directions



We assume isotropic distribution of directions

$$\langle \vec{m} \rangle_Q = 0$$

\vec{l} ... in the laboratory frame fixed

\vec{m} ... different for each molecule
- in the molecular frame of reference

In absorption

secular approximation

$$\langle (\vec{d}_n \cdot \vec{l}) (\vec{d}_n^\dagger \cdot \vec{l}) \rangle_Q = |\vec{d}|^2 \langle (\vec{m} \cdot \vec{l})^2 \rangle$$

In pump-probe (and other non-linear spectra)

$$\begin{aligned} & \langle (\vec{d}_4^\dagger \cdot \vec{l}_4) (\vec{d}_3^\dagger \cdot \vec{l}_3) (\vec{d}_2^\dagger \cdot \vec{l}_2) (\vec{d}_1^\dagger \cdot \vec{l}_1) \rangle_Q = \\ & = |\vec{d}_1| |\vec{d}_2| |\vec{d}_3| |\vec{d}_4| \langle (\vec{m}_4^\dagger \cdot \vec{l}_4^\dagger) (\vec{m}_3^\dagger \cdot \vec{l}_3^\dagger) (\vec{m}_2^\dagger \cdot \vec{l}_2^\dagger) (\vec{m}_1^\dagger \cdot \vec{l}_1^\dagger) \rangle_Q \end{aligned}$$

four different polarizations: $\vec{l}_1^\dagger, \vec{l}_2^\dagger, \vec{l}_3^\dagger, \vec{l}_4^\dagger$... fixed in space

four different orientations of \vec{d}^\dagger , mutually fixed but arbitrarily rotated

ad a)

$$\begin{aligned}
 (\vec{d}_1 \cdot \vec{l}_1) (\vec{d}_2 \cdot \vec{l}_2) &= d_1 d_2 (\vec{u}_1 \cdot \vec{l}_1) (\vec{u}_2 \cdot \vec{l}_2) \\
 &= d_1 d_2 \sum_{ij} (M_i^{(1)} \cdot l_i^{(1)}) (M_j^{(2)} \cdot l_j^{(2)}) \\
 &= d_1 d_2 \sum_{ij} (l_i^{(1)} \cdot l_j^{(1)}) \underbrace{(M_i^{(1)} M_j^{(2)})}_{\text{factor}}
 \end{aligned}$$

$$\langle (\vec{d}_1 \cdot \vec{l}_1) (\vec{d}_2 \cdot \vec{l}_2) \rangle_S = d_1 d_2 \sum_{ij} (l_i^{(1)} \cdot l_j^{(1)}) \langle M_i^{(1)} M_j^{(2)} \rangle_S$$

$$\begin{aligned}
 \sum_{ij} (\hat{S})_{\varepsilon i} A_{ij} \cdot S_{j\varepsilon} &= \alpha \sum_{ij} (\hat{S})_{\varepsilon i} \delta_{ij} \cdot S_{j\varepsilon} = \\
 &\stackrel{\alpha \delta_{ij}}{=} \alpha \sum_i (\hat{S})_{\varepsilon i} S_{ie} = \alpha \delta_{\varepsilon e}
 \end{aligned}$$

$$\boxed{\langle M_i^{(1)} M_j^{(2)} \rangle = \alpha \delta_{ij}}$$

$$\begin{aligned}
 \langle (\vec{d}_1 \cdot \vec{l}_1) (\vec{d}_2 \cdot \vec{l}_2) \rangle_S &= d_1 d_2 \sum_{ij} (l_i^{(1)} \cdot l_j^{(2)}) \sum_{\alpha} \langle T_{i\alpha} M_\alpha^{(1)} T_{j\beta} M_\beta^{(2)} \rangle_S \\
 &= d_1 d_2 \sum_{ij} (l_i^{(1)} \cdot l_j^{(2)}) \underbrace{\langle T_{i\alpha} T_{j\beta} \rangle_S}_{I_{ij\alpha\beta}} M_\alpha^{(1)} M_\beta^{(2)}
 \end{aligned}$$

$I_{ij\alpha\beta}$

$$I_{ij\alpha\beta} = \alpha \delta_{ij} \dots ?$$

$$\sum_i M_i^{(1)} M_i^{(2)} = \text{const} = \underbrace{\sum_i \sum_{\alpha\beta} I_{i\alpha j\beta} M_\alpha^{(1)} M_\beta^{(2)}}_{\mathcal{H}_{\alpha\beta}} = \sum_{\alpha} M_{\alpha}^{(1)} M_{\alpha}^{(2)}$$

$$\sum_{\alpha} \left(\sum_{\beta} \underbrace{\sum_i I_{i\alpha j\beta} M_\beta^{(2)}}_{\mathcal{H}_{\alpha\beta}} - M_{\alpha}^{(2)} \right) M_{\alpha}^{(1)} = 0$$

$$\sum_{\alpha} \partial e_{\alpha p} m_{\alpha}^{(2)} = M_{\alpha}^{(2)}$$

$\delta_{\alpha p}$

$$\sum_i I_{ii\alpha p} = \delta_{\alpha p} \rightarrow \boxed{I_{ij\alpha p} = \alpha \delta_{ij} \cdot \delta_{\alpha p}}$$

$$\sum_i \alpha \delta_{ii} \cdot \delta_{\alpha p} = \delta_{\alpha p}$$

$$3\alpha \delta_{\alpha p} = \delta_{\alpha p}$$

$$\Rightarrow \boxed{\alpha = \frac{1}{3}}$$

$$\boxed{I_{ij\alpha p} = \frac{1}{3} \delta_{ij} \cdot \delta_{\alpha p}}$$

$$\begin{aligned} \langle (\vec{d}_1 \cdot \vec{e}_1) (\vec{d}_2 \cdot \vec{e}_2) \rangle_2 &= d_1 d_2 \sum_{\substack{i,j \\ \alpha}} \left(\vec{d}_i^{(1)} \vec{d}_j^{(2)} \right) \frac{1}{3} \delta_{ij} \cdot \delta_{\alpha p} (m_{\alpha}^{(1)} m_{\alpha}^{(2)}) \\ &= \frac{d_1 d_2}{3} \sum_{\alpha} \left(\vec{d}_i^{(1)} \vec{d}_i^{(2)} \right) (m_{\alpha}^{(1)} m_{\alpha}^{(2)}) \\ &= \frac{d_1 d_2}{3} (\vec{d}_1 \cdot \vec{d}_2) (\vec{m}_1 \cdot \vec{m}_2) \end{aligned}$$

In absorption spectroscopy

$$\vec{e}_1 = \vec{e}_2$$

$$\langle (\vec{d}_1 \cdot \vec{e}) (\vec{d}_2 \cdot \vec{e}) \rangle_2 = \frac{d_1 d_2}{3} (\vec{m}_1 \cdot \vec{m}_2)$$

secular approximation

$$\left\langle (\vec{d}_m \cdot \vec{\ell})^2 \right\rangle_2 = \frac{|\vec{d}_m|^2}{3}$$