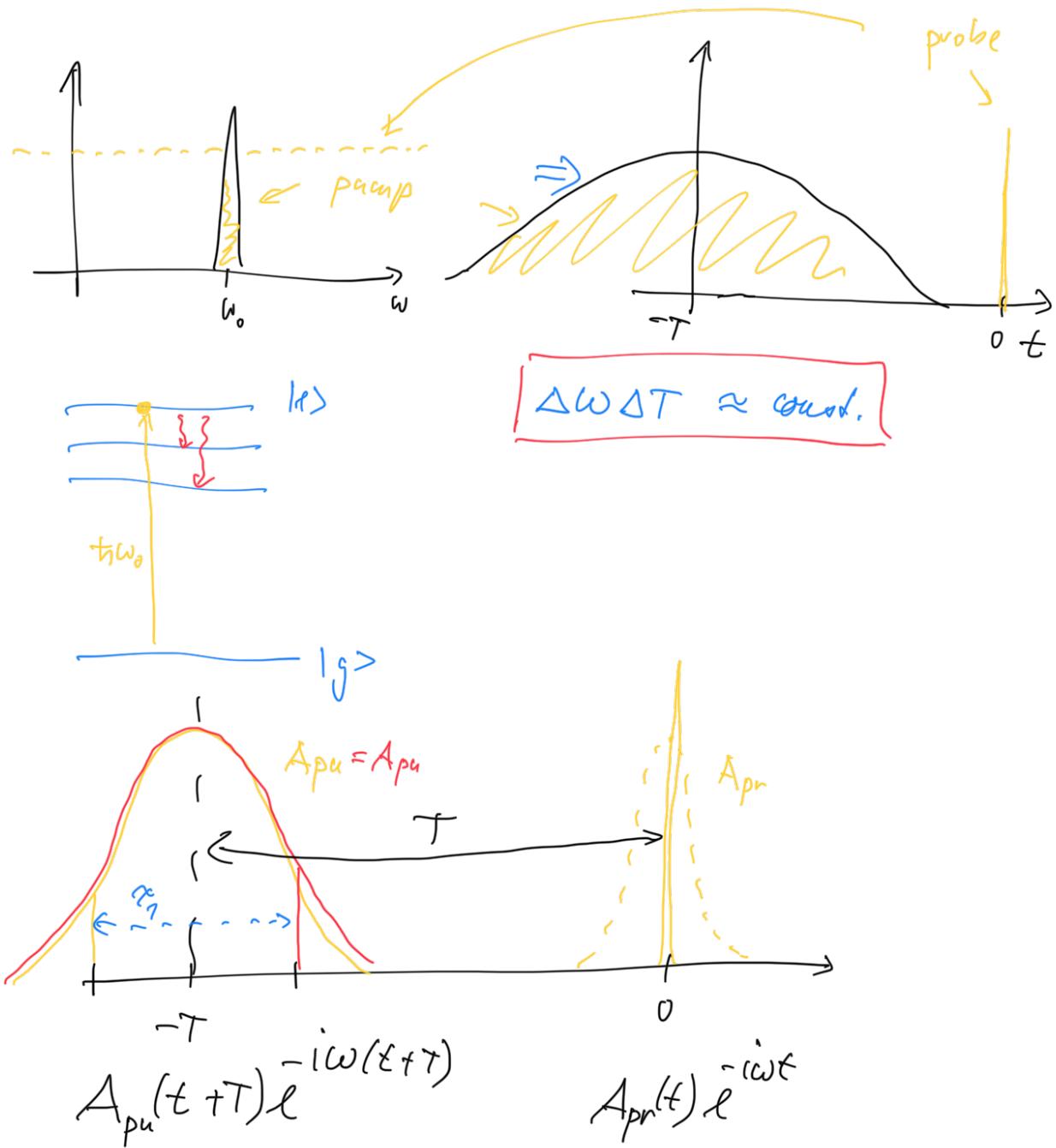


## Pump-probe with finite pulses



- Pump pulse occurs  $2\pi$  in the perturbation theory
- Pump precedes probe

$$P^{(3)}(t, T, \epsilon=0) \approx \int_0^\infty d\tau_3 \int_0^\infty d\tau_2 \int_0^\infty d\tau_1 S^{(3)}(\tau_3, \tau_2, \tau_1) E(t-\tau_3) E(t-\tau_3-\tau_2) + E(t-\tau_3-\tau_2-\tau_1)$$

- the response is a sum of Liouville pathways
- we look into the direction  $-\vec{\epsilon}_1 + \vec{\epsilon}_2 + \vec{\epsilon}_3 = -\vec{\epsilon}_{pr} + \vec{\epsilon}_{pu} + \vec{\epsilon}_{pr} = \vec{\epsilon}_{pr}$
- in this direction we consider interaction order

1-2-3  $\leftarrow$  rephasing pathways  
 2-1-3  $\leftarrow$  non-rephasing pathways

$$P^{(3)}(t, T, 0) \propto (i)^3 \frac{-i\omega(t-\tau)}{e} \int_0^{\infty} d\tau_3 \int_0^{\infty} d\tau_2 \int_0^{\infty} d\tau_1 R_{(reph)}(\tau_3, \tau_2, \tau_1)$$

$\boxed{-\vec{\epsilon}_1 + \vec{\epsilon}_2 + \vec{\epsilon}_3}$

1-2-3

$$\times e^{+i\omega(\tau_3 - \tau_1)} A_3(t - \tau_1) A_1(t + T - \tau_1 - \tau_2) A_2^*(t + T - \tau_1 - \tau_2 - \tau_3)$$

$\uparrow_{pr}$

$$+ (i)^3 \frac{-i\omega(t-\tau)}{e} \int_0^{\infty} d\tau_3 \int_0^{\infty} d\tau_2 \int_0^{\infty} d\tau_1 R_{(non-ph)}(\tau_3, \tau_2, \tau_1)$$

2-1-3

$$\times e^{+i\omega(\tau_3 + \tau_1)} A_3(t - \tau_1) A_1^*(t + T - \tau_1 - \tau_2) A_2(t + T - \tau_1 - \tau_2 - \tau_3)$$

$A_3(t) = A_{pr}(t) \approx A_{pr} \delta(t)$

$R(\tau_1, \tau_2, \tau_3) = U(t) U(\tau_2) U(\tau_3)$

$$P^{(3)}(t, T) \propto (i)^3 \cancel{e^{-i\omega t}} A_{pr} \int_0^{\infty} d\tau_2 \int_0^{\infty} d\tau_1 R_{(uph)}(t, \tau_2, \tau_1) e^{+i\omega(t-\tau_3)}$$

$$\times A_2(t + T - \tau_1 - \tau_2) A_1^*(t + T - \tau_1 - \tau_2 - \tau_3)$$

+ ...

$$= (i)^3 A_{pr} \mathcal{U}(t) \int_0^\infty d\tau_2 \mathcal{U}(\tau_2) A_2(t-\tau_2) \int_0^\infty d\tau_1 \mathcal{U}(\tau_1) e^{-i\omega \tau_1} f$$

$\times A_1^*(t-\tau_2-\tau_1) g^*$

$$\int_0^\infty d\tau_1 f(\tau_1) g(x-\tau_1)$$

$x = t - \tau_2$

Convolution

$$[f * g](x) = \int_0^\infty d\tau f(\tau) g(x-\tau)$$

$$\int_{-\infty}^\infty dx e^{i\omega x} [f * g](x) = f(\omega) g(\omega)$$

$$f(\omega) = \int_{-\infty}^\infty d\tau f(\tau) e^{i\omega \tau}$$

$$[f * g](x) = \int_{-\infty}^\infty d\omega e^{-i\omega x} f(\omega) g(\omega)$$

Rephrasing diagrams

### Complex lineshape

$$\mathcal{U}_{gege}(\tau_1) e^{-i\omega \tau_1}$$

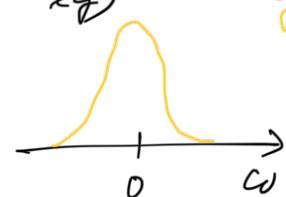
$$= \tilde{\mathcal{U}}_{gege}(\tau_1) e^{i(\tilde{\omega}_{eg} - \omega) \tau_1}$$

$$\int_{-\infty}^\infty d\omega \mathcal{U}_{gege}(\tau) e^{i\omega \tau} = \int_{-\infty}^\infty d\omega \tilde{\mathcal{U}}_{gege}(\tau) e^{-i\omega_{eg}\tau + i\omega \tau}$$

centered around  $\omega_{eg}$

$$= G_{eg}(\omega - \omega_{eg})$$

$$G_{eg}(\omega) = \int_{-\infty}^\infty d\omega \tilde{\mathcal{U}}_{gege}(\tau) e^{i\omega \tau} =$$



We have

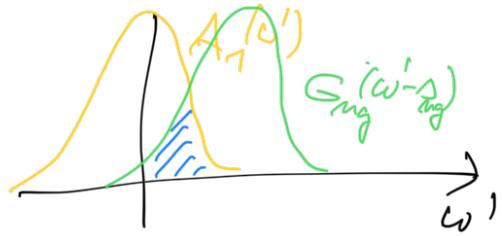
$$\begin{aligned} U_{gege}(\tau_1) e^{-i\omega \tau_1} &\xrightarrow{iFT} \int_{-\infty}^{\infty} d\tau_1 \left[ \tilde{U}_{gege}(\tau_1) e^{i\omega_{eq}\tau_1 - i\omega \tau_1} \right] e^{-i\omega' \tau_1} \\ &= \int_{-\infty}^{\infty} d\tau_1 \tilde{U}_{gege}(\tau_1) e^{i(\omega_{eq} - \omega - \omega') \tau_1} \\ &= \int_{-\infty}^{\infty} d\tau_1 \left[ \tilde{U}_{gege}(\tau_1) e^{i(\omega' + \omega - \omega_{eq}) \tau_1} \right]^* \\ &= G_{eq}^*(\omega' + \omega - \omega_{eq}) = \\ &= G_{eq}^*(\omega' - \Delta) \quad \boxed{\Delta = \omega_{eq} - \omega} \end{aligned}$$

$$A_1^*(T - \tau_2 - \tau_1) \xrightarrow{iFT} \int_{-\infty}^{\infty} d\tau_1 A_1^*(\tau_1) e^{-i\omega' \tau_1} = \int_{-\infty}^{\infty} d\omega' \left[ A_1(\omega) e^{i\omega' \tau_1} \right]^*$$
$$= A_1^*(\omega')$$

$$\begin{aligned} P^{(1)}(t, T) &\approx (i)^3 A_{pr} \mathcal{H}(t) \int_0^{\infty} d\tau_2 U(\tau_2) A_2(T - \tau_2) \\ &\times \int_{-\infty}^{\infty} d\omega' e^{-i\omega' (T - \tau_2)} G_{eq}^*(\omega' - \underbrace{(\omega_{eq} - \omega)}_{\Delta_{eq}}) A_1^*(\omega') \\ &+ \dots \end{aligned}$$

... defining  
of pulse train transition

$$G_{eq}^*(\omega' - \Delta_{eq}) A_1^*(\omega')$$



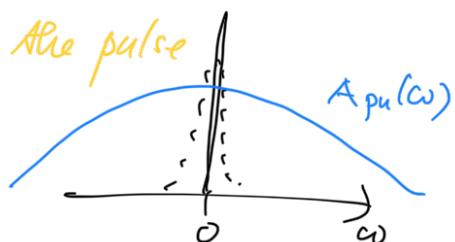
Non-rephasing part

The only difference is  $e^{i\omega\tau_1}$   
pulse  $A_2$

$$\begin{aligned} P^{(2)}(\epsilon, T, \tau=0) /_{\text{non-reph}} &\approx (i)^3 A_{pr} \mathcal{U}(t) \int d\tau_2 \mathcal{U}(\tau_2) A_1^*(\tau - \tau_2) \\ &\quad \times \int_0^\infty d\tau_1 G_{eq}(\tau_1) e^{i\omega\tau_1} A_2(T - \tau_2 - \tau_1) \\ &= (i)^3 A_{pr} \mathcal{U}(t) \int d\tau_2 \mathcal{U}(\tau_2) A_1^*(\tau - \tau_2) \int_{-\infty}^\infty d\omega' e^{i\omega'(\tau - \tau_2)} \underbrace{G_{eq}(\omega' + \Delta_{eq}) A_2(\omega')}_{\text{green circle}} \end{aligned}$$

Case: Line-shape narrower than the pulse

$$G_{eq}(\omega) \approx \delta(\omega)$$



$$\begin{aligned} \int_{-\infty}^\infty d\omega' e^{i\omega'(\tau - \tau_2)} \delta(\omega' + \omega - \omega_{eq}) A_1^*(\omega') &= \\ &= e^{-i \underbrace{(\omega_{eq} - \omega)(\tau - \tau_2)}_{\Delta_{eq}}} A_1^*(\Delta_{eq}) \end{aligned}$$

$$\begin{aligned}
P^{(G)}(\epsilon, T) &\approx (i)^3 A_{pr} \mathcal{U}(f) \int_0^\infty d\tau_2 \mathcal{U}(\tau_2) A_2(T-\tau_2) A_1^*(\Delta_{eq}) e^{-i\Delta_{eq}(T-\tau_2)} \\
&+ \dots \quad \mathcal{U}(\tau_2) \approx \text{const.} = \mathcal{U}(T) \quad \begin{array}{c} P_\alpha(f) \\ \rightarrow t \\ \mathcal{U}(T) \end{array} \\
&= (i)^3 A_{pr} \mathcal{U}(f) \mathcal{U}(\tau_2) \underbrace{A_1^*(\Delta_{eq}) \int_0^\infty d\tau_2 A_2(T-\tau_2) e^{-i\Delta_{eq}(T-\tau_2)}}_{\zeta' = T - \tau_2, d\tau' = -d\tau_2, \zeta'(\infty) = -\infty, \zeta'(0) = T} \\
&+ \dots \quad A_2(\Delta_{eq})
\end{aligned}$$

$$\begin{aligned}
P^{(G)}(\epsilon, T) &= \\
&\approx (i)^3 A_{pr} \mathcal{U}(f) \mathcal{U}(\tau_2) A_1^*(\Delta_{eq}) \int_{-\infty}^T d\tau' A_2(\tau') e^{-i\Delta_{eq}\tau'} \\
&+ (i)^3 A_{pr} \mathcal{U}(f) \mathcal{U}(\tau_2) A_2(\Delta_{eq}) \int_{-\infty}^T d\tau' A_1^*(\tau') e^{i\Delta_{eq}\tau'}
\end{aligned}$$

$A_1$  and  $A_2$  are real and equal to  $A$

$$\approx (i)^3 A_{\text{pr}} \mathcal{U}(t) \mathcal{U}(\frac{T}{2}) A(\Delta_{\text{eg}}) \left[ \int_{-\infty}^T d\tau' A(\tau') e^{-i\Delta_{\text{eg}} \tau'} + \int_{-\infty}^T d\tau' A(\tau') e^{i\Delta_{\text{eg}} \tau'} \right]$$

*tab.*

$$\gamma'' = -\tau'$$

$$d\tau' = -d\tau$$

$$\tau''(-\infty) = \infty$$

$$\tau''(T) = -T$$

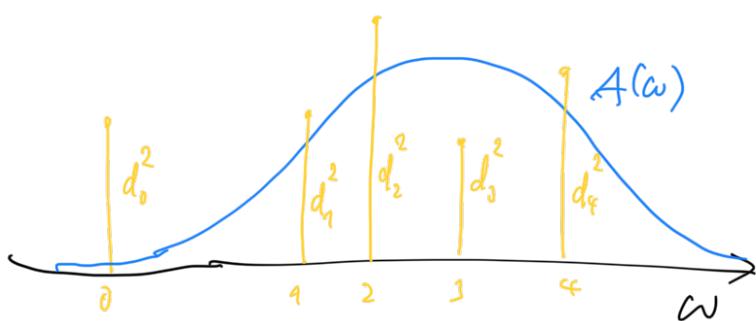
$$\int_{-\infty}^T (-d\tau') \rightarrow \int_{-T}^{\infty} d\tau'' A(-\tau'') e^{i\Delta_{\text{eg}} \tau''}$$

*A(τ'')* ← symmetric pulse

$$P^{(3)}(t, T) \approx (i)^3 A_{\text{pr}} \mathcal{U}(t) \mathcal{U}(\frac{T}{2}) A(\Delta_{\text{eg}}) \left[ \int_{-\infty}^{\infty} d\tau' A(\tau') e^{i\Delta_{\text{eg}} \tau'} + \int_{-T}^T d\tau' A(\tau') e^{i\Delta_{\text{eg}} \tau'} \right]$$

$$T > \nearrow$$

$$\approx (i)^3 A_{\text{pr}} \mathcal{U}(t) \mathcal{U}(\frac{T}{2}) \left| A(\Delta_{\text{eg}}) \right|^2$$



Pump-probe spectrum can be derived by  $|A(\omega)|^2 = |E_{\text{pu}}(\omega)|^2$  to correct for pulse width (pulse spectrum)!

Erratum

$\tau_2 \rightarrow T$