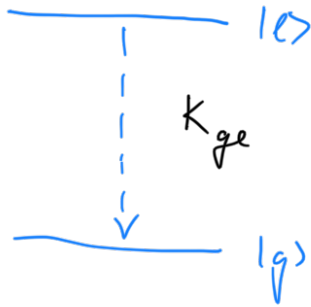
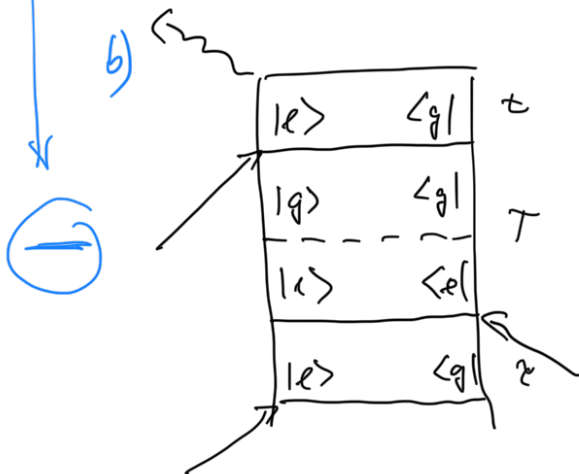
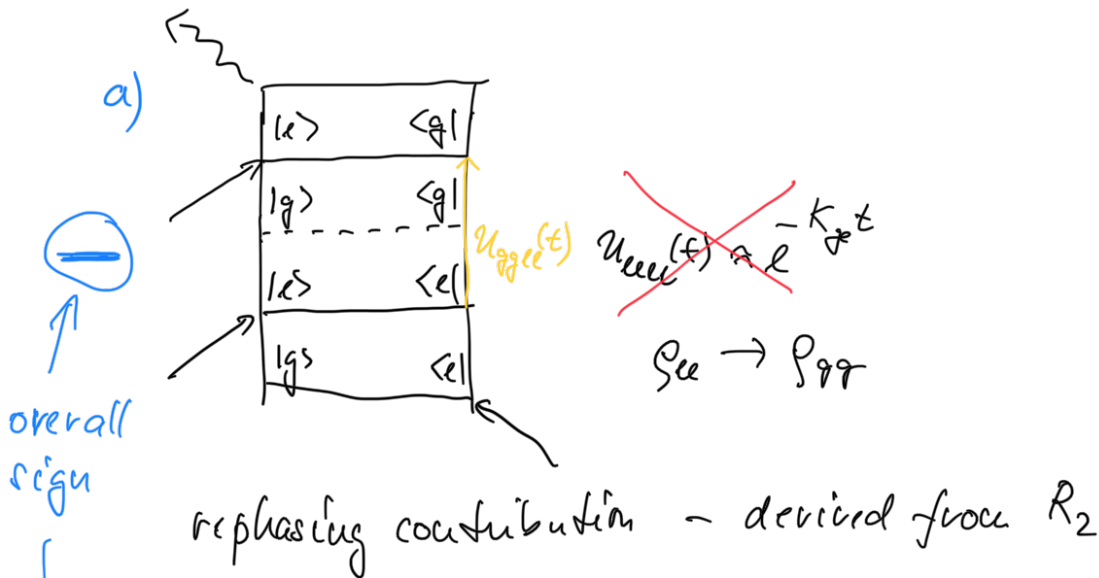


Liouville pathways with energy relaxation processes



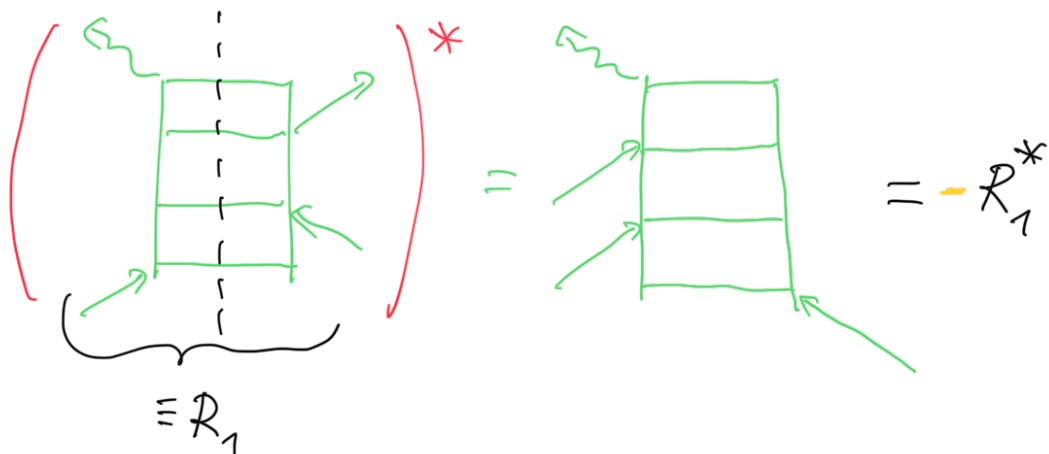
$$U_{ggee}(t) \quad \rho_{ee}(t) \rightarrow \rho_{gg}(t)$$

$$1 - e^{-k_{ge}t} = \int_0^t dt' k_{ge} U_{eeee}(t')$$

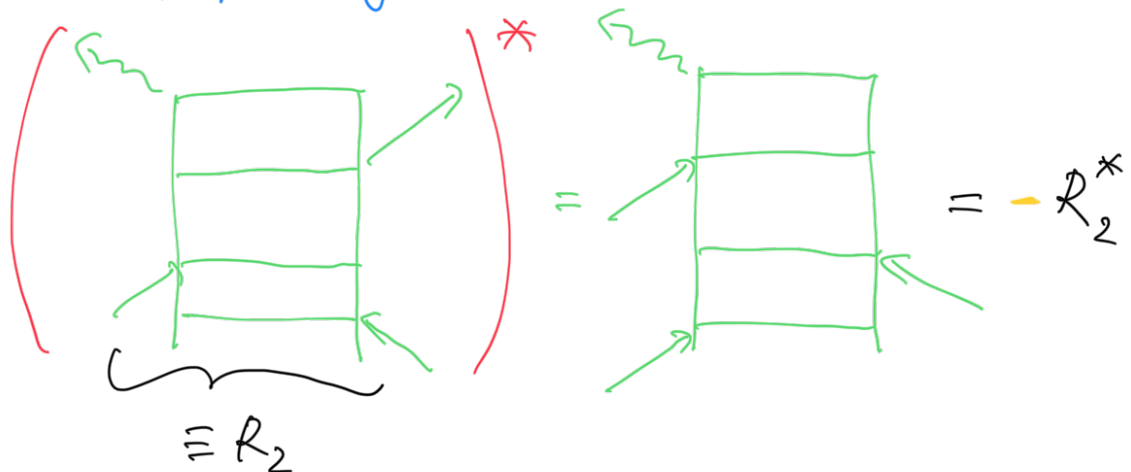


The above diagrams must correspond to R_1^*, R_2^*, R_3^* or R_4^* .

Identification of pathway a)



— of pathway b)



Summary of Liouville pathways for relaxing 2LS

$$R_1(t, T, \tau) = d^4 \mathcal{U}_{egeg}(t) \mathcal{U}_{eeec}(T) \mathcal{U}_{egeg}(\tau)$$

$$R_2(t, T, \tau) = d^4 \mathcal{U}_{egeg}(t) \mathcal{U}_{eeec}(T) \mathcal{U}_{gege}(\tau)$$

$$R_3(t, T, \tau) = d^4 \mathcal{U}_{egeg}(t) \mathcal{U}_{gggg}(T) \mathcal{U}_{goge}(\tau)$$

$$R_4(t, T, \tau) = d^4 \mathcal{U}_{egeg}(t) \mathcal{U}_{gggg}(T) \mathcal{U}_{egeg}(\tau)$$

$$R_1^*(t, T, \tau) = d^4 \mathcal{U}_{egeg}(t) \mathcal{U}_{ggee}(T) \mathcal{U}_{gege}(\tau)$$

$$R_2^*(t, T, \tau) = d^4 \mathcal{U}_{egeg}(t) \mathcal{U}_{geec}(T) \mathcal{U}_{egeg}(\tau)$$

$$\boxed{\text{FT}[\mathcal{U}_{egeg}(t)] = G_{eg}(\omega)}$$

Pump-probe

$$\Delta A(\omega) \approx -\omega \operatorname{Re} (R_1(\omega; T, z=0) + R_2 + R_3 + R_4 - R_1^* - R_2^*)$$

$$\Delta A(T=0) \approx -\omega d^4 \operatorname{Re} [G_{eg}^{R_1}(\omega) + G_{eg}^{R_L}(\omega) + G_{eg}(\omega) + G_{eg}(\omega)]$$

$$\begin{matrix} -0 & -0 \\ \uparrow & \uparrow \\ R_1^* & R_L^* \end{matrix}$$

$$\approx -4d^4 \chi_{eg}(\omega)$$

$$\Delta A(T) \approx -\omega d^4 \operatorname{Re} [(G_{eg}(\omega) + G_{eg}(\omega)) e^{-K_{ge}T}$$

$$+ (G_{eg}(\omega) + G_{eg}(\omega)) \cdot 1$$

$$- (G_{eg}(\omega) + G_{eg}(\omega)) (1 - e^{-K_{ge}T})]$$

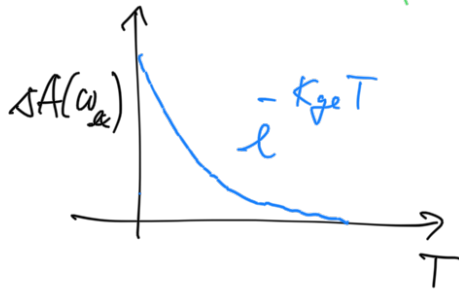
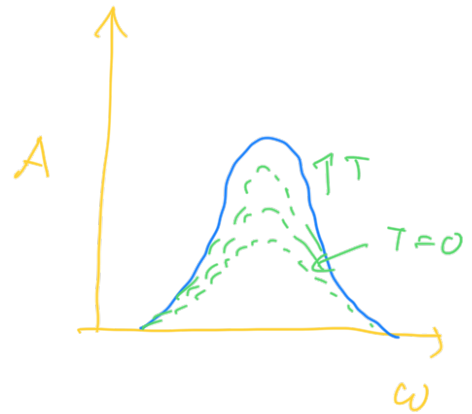
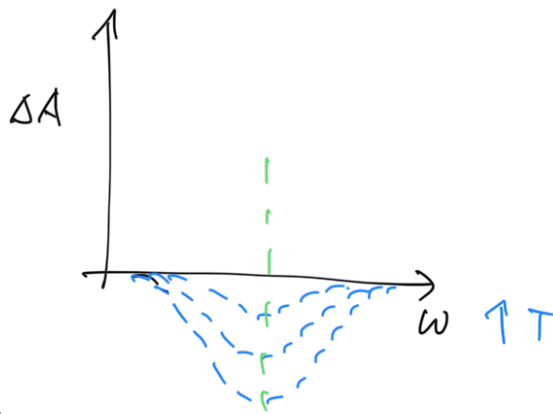
$\mathcal{N}_{exc}(T)$
↓

$\mathcal{N}_{ground}(T)$
←

↑
 $\mathcal{N}_{ground}(T)$

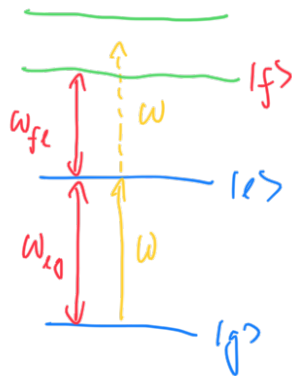
$$\Delta A(T) \approx -\omega d^4 \operatorname{Re} [4 G_{eg}(\omega)] e^{-K_{ge}T}$$

$$\Delta A(T \rightarrow \infty) \approx 0$$

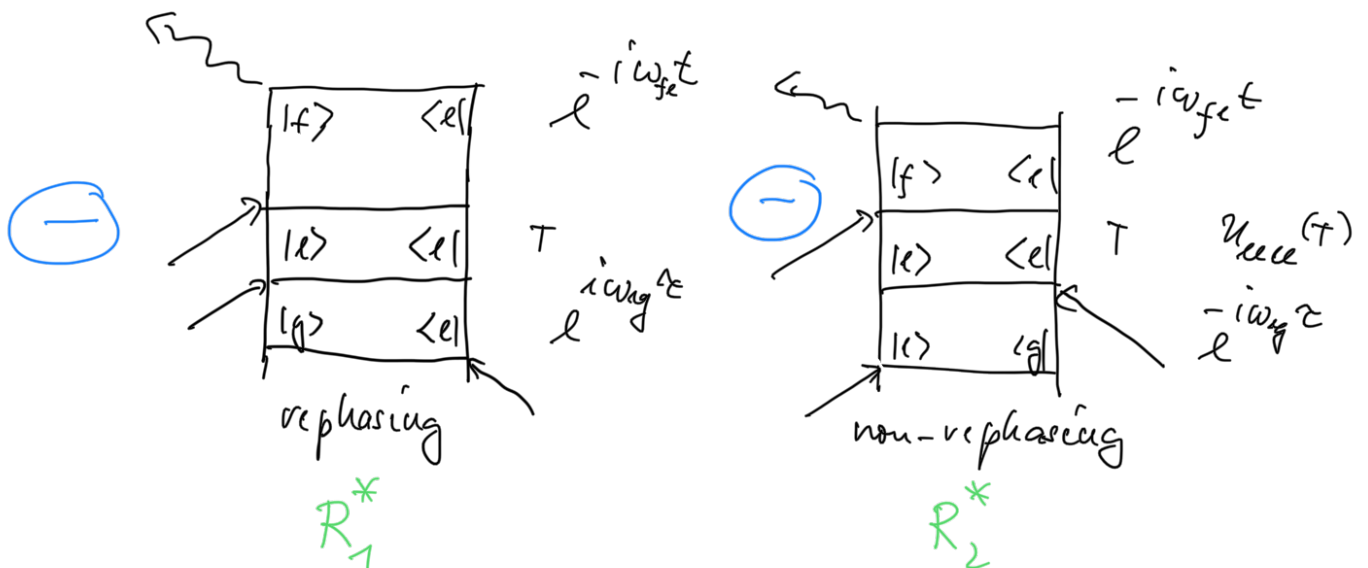


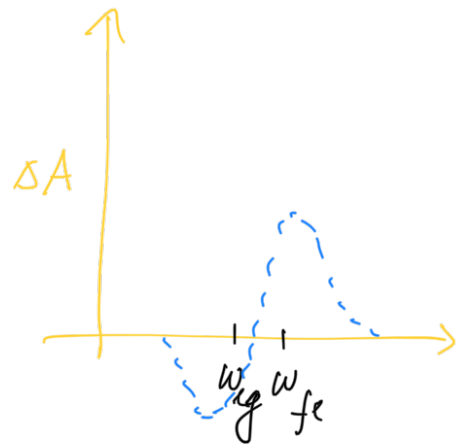
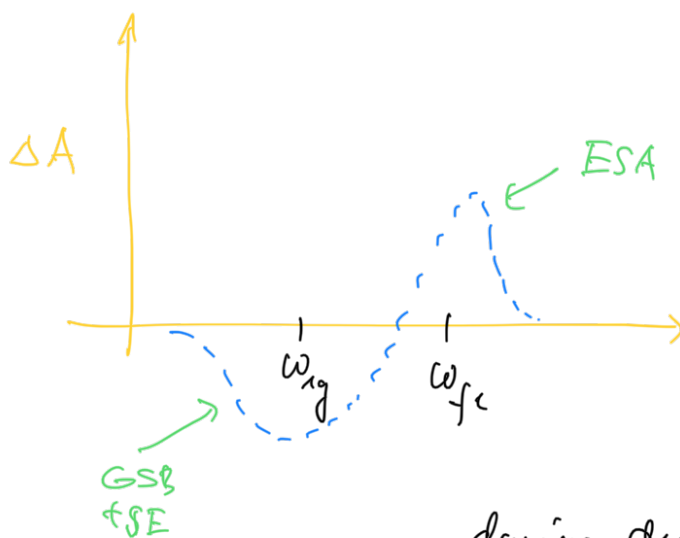
← pump-probe can measure relaxation rate k_{ge}

Systems with excited state absorption



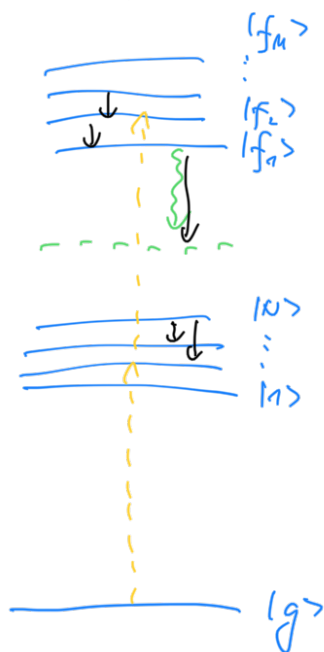
what pathways describe excited state absorption (ESA)





during decay of $|e\rangle$ also ESA decays
 $\sim e^{-k_{pe}t}$

General multi-level systems



$$U_{\text{multi}}(t)$$

Energy relaxation pathways

$$\sim \sum_m d_m^2 d_m^2 U_{ngng}(t) U_{mnmn}(\tau) U_{gmgn}(\tau)$$

Coherence decay pathways

$$\sim \sum_m d_m^2 d_m^2 U_{ngng}(t) U_{mnmn}(\tau) U_{gmgn}(\tau)$$

we assume secular approximation

$$U_{f_n f_n}(t) \rightarrow \text{into ESA}$$

How to calculate spectra

Absorption spectrum

$$\alpha(\omega) \approx \omega \sum_n \langle \vec{d}_{ng} \cdot \vec{e} \rangle_{\Omega} \text{Re } G_{ng}(\omega)$$



$$\begin{aligned} G_{ng}(\omega) &= \int_{-\infty}^{\infty} dt \theta(t) \mathcal{U}_{ngng}(t) e^{i\omega t} \\ &= \int_0^{\infty} dt \tilde{\mathcal{U}}_{ngng}(t) e^{-i\omega_{ng}t} e^{i\omega t} \end{aligned}$$

Pump-probe spectrum

$$\Delta A(\omega) \approx -\omega \sum_{nm} \sum_{\epsilon\epsilon'} \langle (\vec{d}_{g\epsilon} \cdot \vec{e}) (\vec{d}_{\epsilon g}) (\vec{d}_{ng} \cdot \vec{e}) (\vec{d}_{gn}) \rangle_{\Omega}$$

SE ~ R₂ + R₁

$$\times \text{Re} \left\{ 2 G_{\epsilon g}(\omega) \mathcal{U}_{\epsilon\epsilon' nm}(\tau) (\delta_{nm} \delta_{\epsilon\epsilon'} + (1 - \delta_{nm}) \delta_{\epsilon n} \delta_{\epsilon' m}) + \right.$$

FT[$\mathcal{U}_{\epsilon g \epsilon g}(t)$]

$$+ 2 G_{\epsilon g}(\omega) \delta_{nm} \delta_{\epsilon\epsilon'}$$

GSB ~ R₃ + R₄

$$- 2 G_{\epsilon g}(\omega) \mathcal{U}_{gg nn}(\tau) \delta_{nm} \delta_{\epsilon\epsilon'} \}$$

GSB filling $\propto R_1^* + R_2^*$ with relaxation to (g)

$$+ \omega \sum_{nm} \sum_{fkl} \langle (\vec{d}_{kf} \cdot \vec{e}) (\vec{d}_{fk}) (\vec{d}_{ng} \cdot \vec{e}) (\vec{d}_{gn}) \rangle_{\Omega}$$

$$\times 2 \text{Re} \left\{ G_{f\epsilon}(\omega) \mathcal{U}_{\epsilon\epsilon' nm}(\tau) [\delta_{nm} \delta_{\epsilon\epsilon'} + (1 - \delta_{nm}) \delta_{\epsilon n} \delta_{\epsilon' m}] \right\}$$

ESA

Rotational averaging over transition dipole moment directions



$$\hat{e} \cdot \vec{E} \approx \vec{d} \cdot \vec{E} = |\vec{d}| |\vec{E}| (\vec{n} \cdot \vec{e})$$

\vec{e} ... in the laboratory frame
fixed

\vec{n} ... different for each
molecule
- in the molecule frame
of reference

We assume isotropic
distribution of directions

$$\langle \vec{n} \rangle_{\Omega} \equiv 0$$

In absorption

scalar approximation

$$\langle (\vec{d}_n \cdot \vec{e})(\vec{d}_n \cdot \vec{e}) \rangle_{\Omega} = |\vec{d}|^2 \langle (\vec{n} \cdot \vec{e})^2 \rangle$$

In pump-probe (and other non-linear spectra)

$$\begin{aligned} & \langle (\vec{d}_4 \cdot \vec{e}_4)(\vec{d}_3 \cdot \vec{e}_3)(\vec{d}_2 \cdot \vec{e}_2)(\vec{d}_1 \cdot \vec{e}_1) \rangle_{\Omega} = \\ & = |\vec{d}_1| |\vec{d}_2| |\vec{d}_3| |\vec{d}_4| \langle (\vec{n}_4 \cdot \vec{e}_4)(\vec{n}_3 \cdot \vec{e}_3)(\vec{n}_2 \cdot \vec{e}_2)(\vec{n}_1 \cdot \vec{e}_1) \rangle_{\Omega} \end{aligned}$$

four different polarizations: $\vec{e}_1, \vec{e}_2, \vec{e}_3, \vec{e}_4$... fixed in space

four different orientations of \vec{d} , mutually fixed
but arbitrarily rotated

ad a)

$$\begin{aligned}
 (\vec{d}_1 \cdot \vec{r}_1) (\vec{d}_2 \cdot \vec{r}_2) &= d_1 d_2 (\vec{u}_1 \cdot \vec{r}_1) (\vec{u}_2 \cdot \vec{r}_2) \\
 &= d_1 d_2 \sum_{ij} (u_i^{(1)} r_i^{(1)}) (u_j^{(2)} r_j^{(2)}) \\
 &= d_1 d_2 \sum_{ij} (r_i^{(1)} r_j^{(2)}) \underbrace{(u_i^{(1)} u_j^{(2)})}_{\uparrow \text{tensor}}
 \end{aligned}$$

$$\langle (\vec{d}_1 \cdot \vec{r}_1) (\vec{d}_2 \cdot \vec{r}_2) \rangle_\Omega = d_1 d_2 \sum_{ij} (r_i^{(1)} r_j^{(2)}) \langle u_i^{(1)} u_j^{(2)} \rangle_\Omega$$

$$\begin{aligned}
 \sum_{ij} (\vec{S}^1)_{\epsilon i} A_{ij} S_{j\epsilon} &\stackrel{!}{=} a \delta_{ij} \\
 &= a \sum_{ij} (\vec{S}^1)_{\epsilon i} \delta_{ij} S_{j\epsilon} = \\
 &= a \sum_i (\vec{S}^1)_{\epsilon i} S_{i\epsilon} = a \delta_{\epsilon\epsilon}
 \end{aligned}$$

$$\langle u_i^{(1)} u_j^{(2)} \rangle = a \delta_{ij}$$

$$\begin{aligned}
 \langle (\vec{d}_1 \cdot \vec{r}_1) (\vec{d}_2 \cdot \vec{r}_2) \rangle_\Omega &= d_1 d_2 \sum_{ij} (r_i^{(1)} r_j^{(2)}) \left[\sum_{\alpha\beta} T_{i\alpha} u_\alpha^{(1)} T_{j\beta} u_\beta^{(2)} \right]_\Omega \\
 &= d_1 d_2 \sum_{ij} (r_i^{(1)} r_j^{(2)}) \underbrace{\langle T_{i\alpha} T_{j\beta} \rangle_\Omega}_{I_{ij\alpha\beta}} u_\alpha^{(1)} u_\beta^{(2)}
 \end{aligned}$$

$$I_{ij\alpha\beta} = a \delta_{ij} \dots ?$$

$$\sum_i u_i^{(1)} u_i^{(2)} = \text{const} = \sum_i \sum_{\alpha\beta} I_{i\alpha\beta} u_\alpha^{(1)} u_\beta^{(2)} = \sum_\alpha u_\alpha^{(1)} u_\alpha^{(2)}$$

$$\sum_\alpha \left(\sum_\beta \underbrace{\sum_i I_{i\alpha\beta}}_{\mathcal{H}_{\alpha\beta}} u_\beta^{(2)} - u_\alpha^{(2)} \right) u_\alpha^{(1)} = 0$$

$$\sum_p x_{\alpha p} \mu_p^{(2)} = \mu_\alpha^{(2)}$$

\parallel
 $\delta_{\alpha p}$

$$\sum_i I_{ii\alpha p} = \delta_{\alpha p} \rightarrow I_{ij\alpha p} = a \delta_{ij} \cdot \delta_{\alpha p}$$

$$\sum_i a \delta_{ii} \delta_{\alpha p} = \delta_{\alpha p}$$

$$3a \delta_{\alpha p} = \delta_{\alpha p}$$

$$\Rightarrow a = \frac{1}{3}$$

$$I_{ij\alpha p} = \frac{1}{3} \delta_{ij} \cdot \delta_{\alpha p}$$

$$\langle (\vec{d}_1 \cdot \vec{e}_1) (\vec{d}_2 \cdot \vec{e}_2) \rangle_\Omega = d_1 d_2 \sum_{\substack{i,j \\ \alpha}} (\vec{e}_i^{(1)} \cdot \vec{e}_j^{(2)}) \frac{1}{3} \delta_{ij} \cdot \delta_{\alpha p} (\mu_\alpha^{(1)} \mu_\alpha^{(2)})$$

$$= \frac{d_1 d_2}{3} \sum_{i\alpha} (\vec{e}_i^{(1)} \cdot \vec{e}_i^{(2)}) (\mu_\alpha^{(1)} \mu_\alpha^{(2)})$$

$$= \frac{d_1 d_2}{3} (\vec{e}_1 \cdot \vec{e}_2) (\vec{\mu}_1 \cdot \vec{\mu}_2)$$

In absorption spectroscopy

$$\vec{e}_1 = \vec{e}_2$$

$$\langle (\vec{d}_1 \cdot \vec{e}) (\vec{d}_2 \cdot \vec{e}) \rangle_\Omega = \frac{d_1 d_2}{3} (\vec{\mu}_1 \cdot \vec{\mu}_2)$$

secular approximation

$$\langle (\vec{d}_n \cdot \vec{e})^2 \rangle_{\Omega} = \frac{|\vec{d}_n|^2}{3}$$