



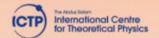
MaX School on Advanced Materials and Molecular Modelling with Quantum ESPRESSO

QE-2021: Hands-on session – Day-6 TDDFPT

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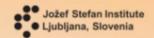
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Outline

1. Part I – Optical absorption in molecules

2. Part II – EELS in solids

(We will use HPC for this hands-on, because many examples are quite heavy)

Part I – Optical absorption in molecules

- 1. **Example 1:** turboDavidson from the Independent Particle Approximation (IPA) to interacting electrons
- 2. **Example 2:** turboLanczos total spectra at low cost for interacting electrons

turboDavidson program for calculations of absorption spectra

- The turbo_davidson.x program allows us to calculate absorption spectra of molecules using time-dependent density functional perturbation theory (TDDFpT).
- The interactions of electrons (Hartree and Exchange-Correlation effects) are taken into account fully ab initio and self-consistently (or they can be neglected if_dft_spectrum=.true.).
- The electronic transitions from occupied to empty states can be analyzed by selecting a frequency range in which the transitions occur.
- However, calculations of overall absorption spectra in a wide frequency range are computationally rather expensive, because many eigenvalues of the Hamiltonian must be computed.

turboDavidson program for calculations of absorption spectra

Coupled linear-response TDDFpT equations (resonant and anti-resonant):

$$(\hat{H}^{\circ} - \varepsilon_{v} - \hbar\omega)\tilde{\varphi}'_{v}(\mathbf{r}, \omega) + \hat{P}_{c}\hat{V}'_{HXC}[\tilde{\varphi}'_{v}(\mathbf{r}, \omega), \tilde{\varphi}'^{*}_{v}(\mathbf{r}, -\omega)]\varphi^{\circ}_{v}(\mathbf{r}) = 0$$

$$(\hat{H}^{\circ} - \varepsilon_{v} + \hbar\omega)\tilde{\varphi}'^{*}_{v}(\mathbf{r}, -\omega) + \hat{P}_{c}\hat{V}'_{HXC}[\tilde{\varphi}'_{v}(\mathbf{r}, \omega), \tilde{\varphi}'^{*}_{v}(\mathbf{r}, -\omega)]\varphi^{\circ}_{v}(\mathbf{r}) = 0$$

These equations can be re-written as Casida's matrix equation:

$$\begin{pmatrix} 0 & \hat{D} \\ \hat{D} + \hat{K} & 0 \end{pmatrix} \begin{pmatrix} Q \\ P \end{pmatrix} = \omega \begin{pmatrix} Q \\ P \end{pmatrix}$$
 interaction terms

where $\hat{D} \longrightarrow (\hat{H}^{\circ} - \varepsilon_{v})$ and $\hat{K} \longrightarrow \hat{P}_{c} \hat{V}'_{\text{HXC}} \varphi_{v}^{\circ}(\mathbf{r})$

Davidson algorithm is used (the same algorithm as in the ground state SCF calculation) to solve the Casida's equations and to obtain the eigenvalues which are used for a calculation of the absorption coefficient.

Independent Particle Approximation

The simplest approximation: Independent Particle Approximation (IPA) - a sum of independent excitations from occupied to empty states.

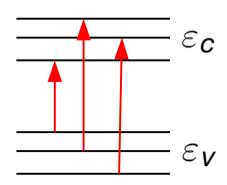
Fermi's golden rule

The transition probability per unit time from occupied states $|\varphi_v\rangle$ to empty states $|\varphi_c\rangle$ reads:

$$T(\omega) = \frac{2\pi}{\hbar} \sum_{v,c} \delta(\varepsilon_c - \varepsilon_v - \hbar\omega) \left| \langle \varphi_c | V'_{ext} | \varphi_v \rangle \right|^2$$

 $V_{ext}' = -e \mathbf{E} \cdot \mathbf{r}$ is the external potential induced by the electric field.

 $\varepsilon_{V}, \varepsilon_{c}$ and $|\varphi_{V}\rangle, |\varphi_{c}\rangle$ are the eigenvalues and the eigenfunctions of the groundstate Kohn-Sham equation

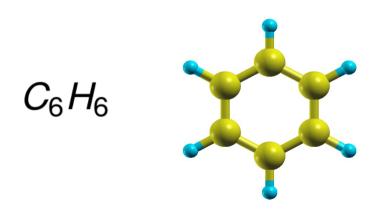


Example 1: Input file for PWscf

```
&CONTROL
    calculation = 'scf'
    restart mode='from scratch',
    prefix='Benzene'
    pseudo dir = '../../pseudo'
   outdir='./tmp'.
&SYSTEM
    ibrav = 6.
   celldm(1) = 32,
    celldm(3) = 0.83,
    nat = 12.
    ntyp = 2.
    ecutwfc = 30.
    nbnd = 20
&FLECTRONS
    conv thr = 1.0d-8
   mixing beta = 0.6,
ATOMIC SPECIES
H 1.0 H.blyp-vbc.UPF
C 12.01 C.blyp-mt.UPF
ATOMIC POSITIONS {Angstrom}
    5.633200899
                  6.320861303
                                5.000000000
    6.847051545
                  8.422621957
                                5.000000000
   8.060751351
                7.721904557
                                5.000000000
   8.060707879 6.320636665
                                5.000000000
   6.846898786
                5.620067381
                                5.0000000000
    5.633279551
                7.722134449
                                5.000000000
   6.847254360
                9.512254789
                                5.0000000000
   9.004364510
                8.266639340
                                5.000000000
   9.004297495
                 5.775895755
                                5.000000000
   6.846845929
                  4.530522778
                                5.000000000
   4.689556006
                  5.776237709
                                5.000000000
   4.689791688
                  8.267023318
                                5.000000000
K POINTS {gamma}
```

Step 1. Perform a Self-Consistent Field ground-state calculation for benzene at the equilibrium structure using the pw.x program.

Specify the total number of Kohn-Sham bands (occupied + some empty states)



Input

Output

Example 1: Output file of PWscf

The output SCF file pw.benzene.scf.out:

```
End of self-consistent calculation

k = 0.0000 0.0000 0.0000 ( 37796 PWs) bands (ev):

-21.1410 -18.3897 -18.3884 -14.6742 -14.6729 -12.5121 -11.0334 -10.4379
-9.8983 -9.8963 -8.5697 -7.8759 -7.8751 -5.8257 -5.8251 -0.7364
-0.7362 -0.3332 0.2501 0.3207

highest occupied, lowest unoccupied level (ev): -5.8251 -0.7364
```

HOMO LUMO

HOMO – Highest Occupied Molecular Orbital

LUMO – Lowest Unoccupied Molecular Orbital

Energy gap = LUMO - HOMO = 5.1 (eV)

Example 1: Calculation of absorption spectra using turboDavidson (IPA)

Perform the turbo Davidson calculation without the electronic interaction using the turbo davidson.x program.

```
&lr_input
    prefix = 'Benzene',
    outdir = './tmp'
/
&lr_dav
    if_dft_spectrum = .true.
    num_init = 30
/

Obsolete variable for the IPA case
(will not be needed in future version of
QE for the IPA case)
```

Perform a spectrum calculation using the turbo_spectrum.x program and using the eigenvalues computed in the previous step.

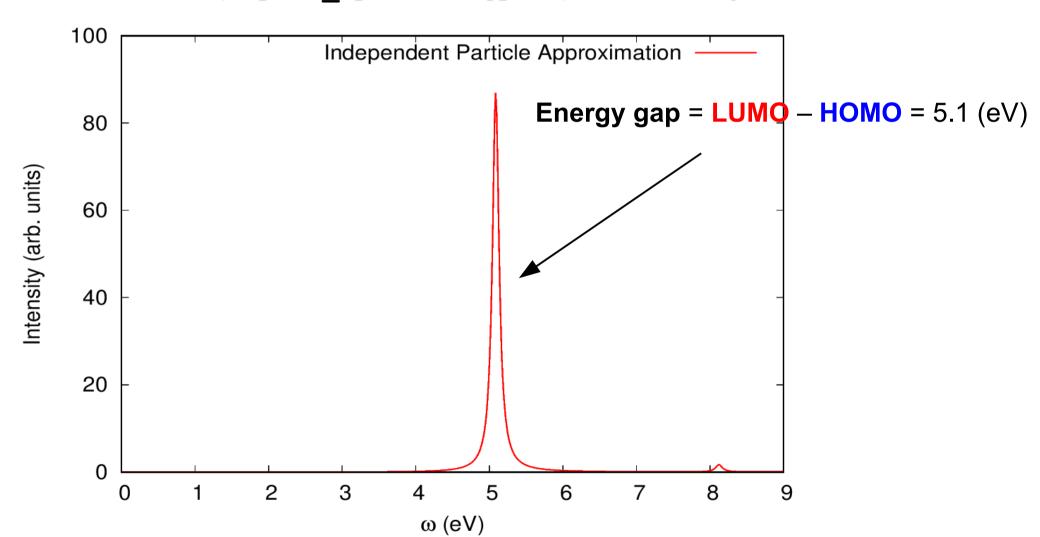
turbo_spectrum.Benzene.tddfpt_pp.in

```
&lr input
                                           The same prefix as in the SCF calculation
  prefix = 'Benzene', 
  outdir = './tmp ,
                                          - Directory for temporary files
                                          - Type of previous calculation
  td = 'davidson',
                                           The value of Lorenzian smearing in Ry
  epsil = 0.004,
                                           Minimum value of frequencies for a plot in Ry
  start = 0.0d0,
                                         — Maximum value of frequencies for a plot in Ry
         = 1.0d0,
  end
  increment = 0.0001d0,
                                         — Frequency step in Ry
  eign file = 'Benzene-dft.eigen' Frequency with Davidson eigenvalues
```

Example 1: Calculation of absorption spectra in IPA

Absorption spectrum of benzene in the Independent Particle Approximation (file Benzene spectrum.eps):

Use the script plot spectrum.gp to produce this figure.

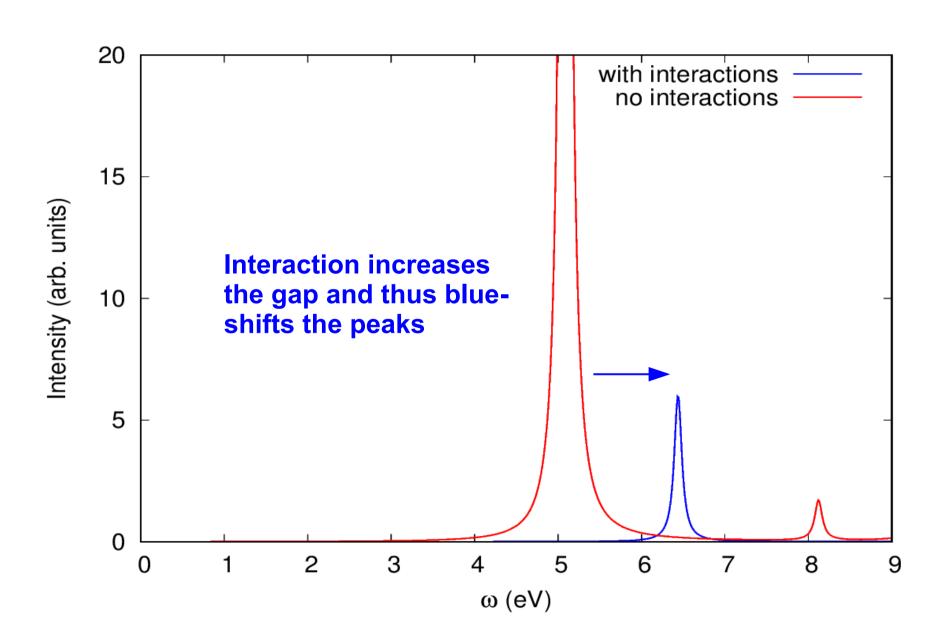


Example 1: Spectra including interactions

Perform the turboDavidson calculation without the electronic interaction using the turbo_davidson.x program.

```
&lr input
                                        The same prefix as in the SCF calculation
  prefix = 'Benzene',
  outdir='./tmp'
                                        Directory for temporary files
&lr dav
                                        Switch on the interactions
  if dft spectrum = .false.
                                        Number of eigenvalues to be calculated
  num eign = 15
  num init = 30
                                        Number of initial vectors
  num basis max = 90
                                        Maximum number of basis allowed for the sub-basis
  residue conv thr = 1.0E-4
                                        Convergence threshold
  start = 0.0
                                        Minimum value of frequencies for a plot in Ry
  finish = 1.0
                                        Maximum value of frequencies for a plot in Ry
  step = 0.001
                                        Frequency step in Ry
                                        Lorentzian broadening parameter in Ry
Reference frequency in Ry where the peak is expected
  broadening
                = 0.005
  reference
                = 0.3
```

Example 1: Comparison of spectra with and without interactions



turboLanczos program for calculation of absorption spectra

- The turbo_lanczos.x program allows us to calculate absorption spectra of molecules using time-dependent density functional perturbation theory (TDDFpT) without computing empty states!
- The interactions of electrons (Hartree and Exchange-Correlation effects) are taken into account fully ab initio and self-consistently.
- The electronic transitions from occupied to empty states cannot be analyzed (use turbo_davidson.x for this purpose).
- The overall absorption spectrum in a wide frequency range can be calculated at once!

turboLanczos program for calculation of absorption spectra

Coupled linear-response TDDFpT equations (resonant and anti-resonant):

$$(\hat{H}^{\circ} - \varepsilon_{v} - \hbar\omega)\tilde{\varphi}'_{v}(\mathbf{r}, \omega) + \hat{P}_{c}\hat{V}'_{HXC}[\tilde{\varphi}'_{v}(\mathbf{r}, \omega), \tilde{\varphi}'^{*}_{v}(\mathbf{r}, -\omega)]\varphi^{\circ}_{v}(\mathbf{r}) = -\hat{P}_{c}\hat{V}'_{ext}\varphi^{\circ}_{v}(\mathbf{r})$$

$$(\hat{H}^{\circ} - \varepsilon_{v} + \hbar\omega)\tilde{\varphi}'^{*}_{v}(\mathbf{r}, -\omega) + \hat{P}_{c}\hat{V}'_{HXC}[\tilde{\varphi}'_{v}(\mathbf{r}, \omega), \tilde{\varphi}'^{*}_{v}(\mathbf{r}, -\omega)]\varphi^{\circ}_{v}(\mathbf{r}) = -\hat{P}_{c}\hat{V}'_{ext}\varphi^{\circ}_{v}(\mathbf{r})$$

These equations can be re-written as a quantum Liouville equation: interaction terms

$$(\omega - \mathcal{L}) \cdot \rho'(\omega) = [\hat{V}'_{ext}, \rho^{\circ}]$$
 $\mathcal{L} \cdot \rho' \equiv [\hat{H}^{\circ}, \rho'] + [\hat{V}'_{HXC}, \rho^{\circ}]$

Absorption coefficient is computed as:

$$\alpha(\omega) = -(\mathbf{r}, (\omega - \mathcal{L})^{-1}[\mathbf{r}, \rho^{\circ}])$$

Lanczos algorithm is used to solve recursively quantum Liouville equation in the standard batch representation. This allows us to avoid inversions and multiplications of large matrices.

Example 2: Calculation of absorption spectra using turboLanczos

Step 1. Perform a SCF ground-state calculation (no nbnd in the input):

```
pw.x < pw.benzene.scf.in > pw.benzene.scf.out
```

Step 2. Perform Lanczos recursions using the turbo_lanczos.x program.

```
&lr_input
    prefix = 'Benzene',
    outdir='./tmp',
    restart_step = 100,
    restart = .false.

/
&lr_control
    itermax = 500,
    ipol = 1

/

The same prefix as in the SCF calculation
Directory for temporary files
The code writes restart files every restart_step iterations
Restart iterations after previous calculation

Number of Lanczos iterations
Polarization direction of incoming light, 1=x

/
```

```
turbo_lanczos.x < turbo_lanczos.benzene.lanczos.in >...
...turbo_lanczos.benzene.lanczos.out
```

Perform a spectrum calculation using the post-processing program turbo_spectrum.x and using the Lanczos coefficients computed in the previous step.

```
&lr_input
                              The same prefix as in the SCF calculation
  prefix = 'Benzene',
                               Directory for temporary files
  outdir = './tmp',

    Number of calculated Lanczos coefficient

  itermax0 = 500,
                                  Number of extrapolated Lanczos coefficient
  itermax = 500.
                              Type of extrapolation
  extrapolation = 'no',
                              ← The value of Lorenzian smearing in Ry
  epsil = 0.004,
                              ■ Minimum value of frequencies for a plot in Ry
  start = 0.0d0,
                              ■ Maximum value of frequencies for a plot in Ry
  end = 1.0d0,
                              Frequency step in Ry
  increment = 0.0001d0,
                               Polarization direction (same as in turbo lanczos.x)
  ipol = 1
```

```
turbo_spectrum.x < turbo_spectrum.benzene.in >...
...turbo_spectrum.benzene.out
```

No extrapolation of Lanczos coefficients

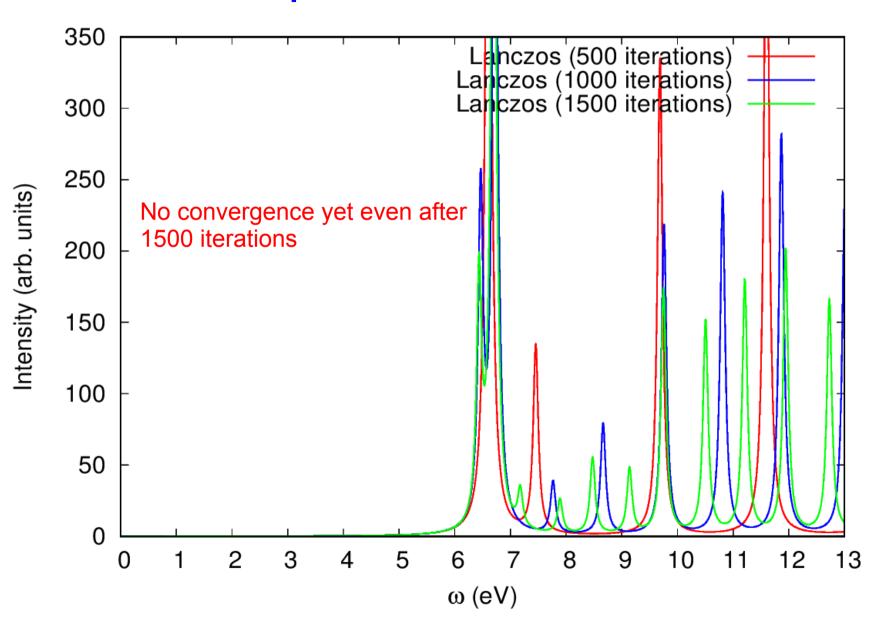
turbo_lanczos.benzene.lanczos.in turbo_spectrum.benzene.lanczos.in

```
&lr_input
&lr_input
                                             prefix = 'Benzene',
  prefix = 'Benzene',
                                             outdir = './tmp',
  outdir='./tmp',
                                             ttermax0 = 500,
  restart_step = 100,
                                             itermax = 500
  restart = .false.
                                             extrapolation = 'no',
                                             epsil = 0.004,
&lr control
                                             start = 0.0d0.
 itermax = 500
                                             end = 1.0d0,
  ipol = 1
                                             increment = 0.0001d0,
                                             ipol = 1
```

Must be equal!

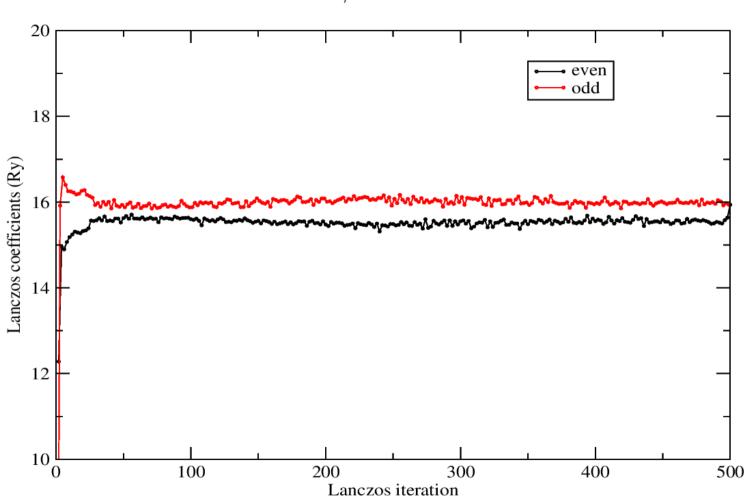
You can use **restart** = .**true**. option in order not to start from scratch, when you want to increase the number of Lanczos iterations.

No extrapolation of Lanczos coefficients



Example 2: Lanczos coefficients





Lanczos coefficients oscillate around half of the kinetic-energy cutoff (~30/2 = 15 Ry in this case). Hence, we can extrapolate these coefficients.

Example 2: Extrapolation of Lanczos coefficients

In the Liuoville-Lanczos approach, the absorption coefficient is computed as:

$$\alpha(\omega) = (\zeta_N, (\omega - \hat{T}_N)^{-1} \cdot \mathbf{e}_1^N)$$

where \hat{T}_N is the tridiagonal matrix composed of the Lanczos coefficients.

$$\hat{T}_{N} = \begin{pmatrix} 0 & \gamma_{2} & 0 & \cdots & 0 \\ \beta_{2} & 0 & \gamma_{3} & 0 & 0 \\ 0 & \beta_{3} & 0 & \cdots & 0 \\ \vdots & 0 & \cdots & 0 & \beta_{N} & 0 \end{pmatrix}$$
Extrapolation

$$\hat{T}_{N'} = \begin{pmatrix} 0 & \gamma_{2} & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ \beta_{2} & 0 & \gamma_{3} & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & \beta_{3} & 0 & \cdots & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & \cdots & \gamma_{N} & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & \beta_{N} & 0 & \langle \gamma \rangle & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & \langle \beta \rangle & 0 & \langle \gamma \rangle & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & \langle \beta \rangle & 0 & \langle \gamma \rangle \end{pmatrix}$$

$$\frac{\langle \beta \rangle}{N} = \frac{\beta_1 + \beta_2 + ... + \beta_N}{N} \qquad \frac{\langle \gamma \rangle}{N} = \frac{\gamma_1 + \gamma_2 + ... + \gamma_N}{N} \qquad \qquad N = itermax0, \quad N' = itermax$$

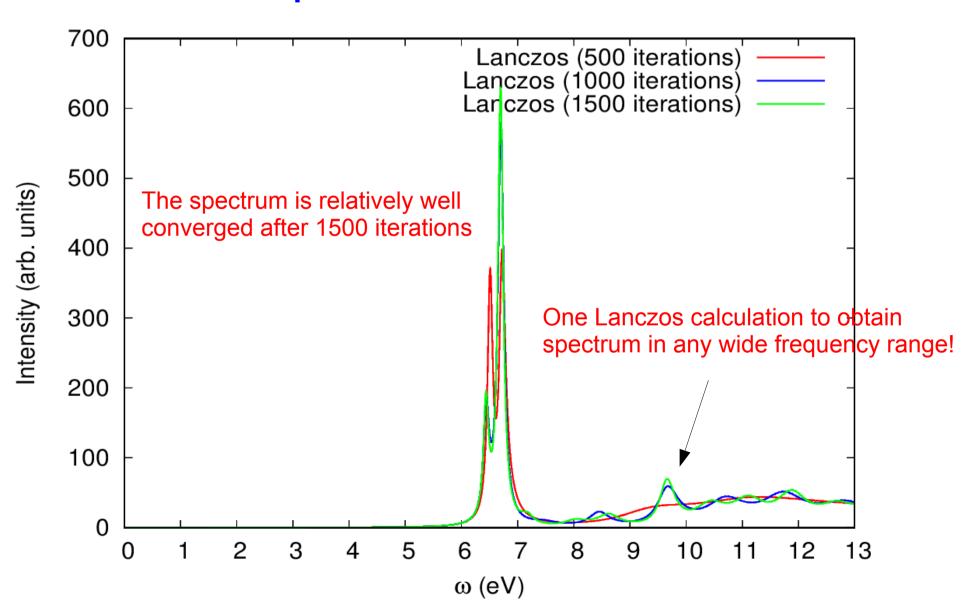
With extrapolation of Lanczos coefficients

turbo_lanczos.benzene.lanczos.in turbo_spectrum.benzene.lanczos.in

```
&lr input
                                           &lr input
  prefix = 'Benzene',
                                             prefix = 'Benzene',
  outdir='./tmp',
                                             outdir = './tmp',
                                            itermax0 = 500
  restart_step = 100,
                                            \mathsf{ttermax} = 20000.
  restart = .false.
                                            extrapolation = 'osc',
                                             epsil = 0.004,
&lr control
                                             start = 0.0d0,
 itermax = 500
                                             end = 1.0d0,
  ipol = 1
                                             increment = 0.0001d0,
                                             ipol =
                 Must be equal!
                                       Extrapolation until 20000
```

You can use **restart** = .true. option in order not to start from scratch, when you want to increase the number of Lanczos iterations.

With extrapolation of Lanczos coefficients



Part II – EELS in solids

1. **Example 3:** EELS spectrum of Si using the **turboEELS** code (Lanczos algorithm)

2. **Example 4:** EELS spectrum of Si using the **turboEELS** code (Sternheimer algorithm)

(very slow)

Basic equations

Coupled (resonant and anti-resonant) frequency-dependent Sternheimer equations:

$$\begin{split} (\hat{\mathbf{H}}_{\mathbf{k}+\mathbf{q}}^{\circ} - \varepsilon_{n,\mathbf{k}}^{\circ} - \omega) \, \tilde{\mathbf{u}}_{n,\mathbf{k}+\mathbf{q}}^{\prime}(\mathbf{r},\omega) \, + \, & \hat{\mathbf{P}}_{c}^{\mathbf{k}+\mathbf{q}} \, \tilde{\mathbf{v}}_{\mathrm{HXC},\mathbf{q}}^{\prime}(\mathbf{r},\omega) \, \mathbf{u}_{n,\mathbf{k}}^{\circ}(\mathbf{r}) \\ (\hat{\mathbf{H}}_{\mathbf{k}+\mathbf{q}}^{\circ} - \varepsilon_{n,\mathbf{k}}^{\circ} + \omega) \, \tilde{\mathbf{u}}_{n,-\mathbf{k}-\mathbf{q}}^{\prime \, *}(\mathbf{r},-\omega) \, + \, & \hat{\mathbf{P}}_{c}^{\mathbf{k}+\mathbf{q}} \, \tilde{\mathbf{v}}_{\mathrm{HXC},\mathbf{q}}^{\prime}(\mathbf{r},\omega) \, \mathbf{u}_{n,\mathbf{k}}^{\circ}(\mathbf{r}) \\ & & \text{interaction terms} \end{split} = -\hat{\mathbf{P}}_{c}^{\mathbf{k}+\mathbf{q}} \, \tilde{\mathbf{v}}_{\mathrm{ext},\mathbf{q}}^{\prime}(\mathbf{r},\omega) \, \mathbf{u}_{n,\mathbf{k}}^{\circ}(\mathbf{r}) \end{split}$$

These equations can be re-written as a quantum Liouville equation and solved using the **Lanczos** algorithm:

$$(\omega - \mathcal{L}) \cdot \hat{\rho}'_{\mathbf{q}}(\omega) = [\tilde{V}'_{ext,\mathbf{q}}(\omega), \hat{\rho}^{\circ}]$$

$$\mathcal{L} \cdot \hat{\rho}'_{\mathbf{q}} \equiv [\hat{H}^{\circ}, \hat{\rho}'_{\mathbf{q}}] + [\tilde{V}'_{HXC,\mathbf{q}}, \hat{\rho}^{\circ}]$$

$$-Im[\varepsilon^{-1}(\mathbf{q}, \omega)] = -\frac{4\pi e^{2}}{|\mathbf{q}|^{2}} Im[\chi(\mathbf{q}, \mathbf{q}; \omega)]$$
Loss function

Example 3

Coupled (resonant and anti-resonant) frequency-dependent Sternheimer equations:

These equations can be re-written as a quantum Liouville equation and solved using the **Lanczos** algorithm:

$$(\omega - \mathcal{L}) \cdot \hat{\rho}'_{\mathbf{q}}(\omega) = [\tilde{V}'_{ext,\mathbf{q}}(\omega), \hat{\rho}^{\circ}]$$

$$\mathcal{L} \cdot \hat{\rho}'_{\mathbf{q}} \equiv [\hat{H}^{\circ}, \hat{\rho}'_{\mathbf{q}}] + [\tilde{V}'_{HXC,\mathbf{q}}, \hat{\rho}^{\circ}]$$

$$-\mathrm{Im}[\varepsilon^{-1}(\mathbf{q}, \omega)] = -\frac{4\pi e^{2}}{|\mathbf{q}|^{2}} \operatorname{Im}[\chi(\mathbf{q}, \mathbf{q}; \omega)]$$
Loss function

Example 3: Input file for PWscf

```
&CONTROL
    calculation = 'scf'
    restart mode='from scratch',
    prefix='silicon'
    pseudo_dir = '../../pseudo',
    outdir='./tmp'
&SYSTEM
    ibrav = 2,
    celldm(1) = 10.26.
    nat = 2.
    ntvp = 1.
    ecutwfc = 16.0
&ELECTRONS
    conv thr = 1.0d-8
    mixing beta = 0.7.
ATOMIC SPECIES {alat}
 Si 28.08 Si.pz-vbc.UPF
ATOMIC POSITIONS {alat}
 Si 0.00 0.00 0.00
 Si 0.25 0.25 0.25
K POINTS {automatic}
10 10 10 1 1 1
```

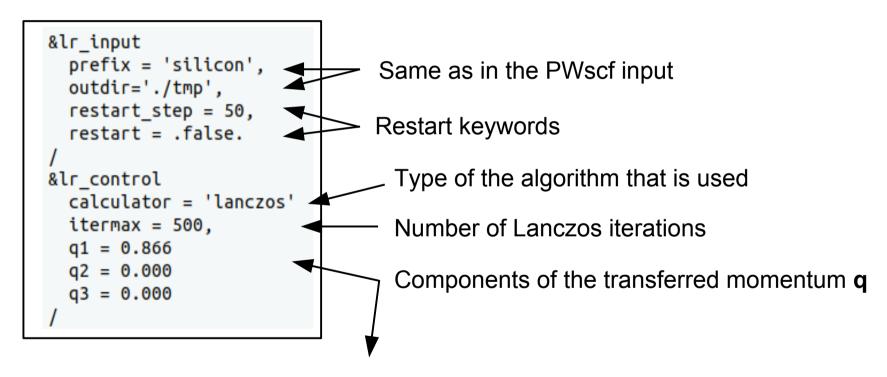
Step 1. Perform a Self-Consistent Field ground-state calculation for silicon using the pw.x program.

Input Output

pw.x < pw.si.scf.in > pw.si.scf.out

Example 3: Input file for turboEELS

Step 2. Perform Lanczos recursions using the turbo_eels.x program in order to compute Lanczos coefficients, which will be needed in Step 3.



How to specify q1, q2, q3? -> See next slide

Example 3: Transferred momentum specification

$$\mathbf{q} = \frac{2\pi}{a_0} \Big(\mathbf{q} 1, \mathbf{q} 2, \mathbf{q} 3 \Big)$$

$$a_0 = \text{celldm}(1) = 10.26 \text{ Bohr}$$
 - lattice parameter

Example:
$$|{\bf q}| = 0.53 \; {\rm Bohr}^{-1} \; {\rm and} \; {\bf q} \, || \, [100] \; .$$

Therefore, we obtain:

$$q1 = \frac{|\mathbf{q}| a_0}{2\pi} = 0.866$$

$$q2 = 0$$

$$q3 = 0$$

Step 3. Perform a spectrum calculation using the post-processing program turbo_spectrum.x and using the Lanczos coefficients computed in the previous step.

```
&lr input
  prefix = 'silicon',
                             The same prefix as in the SCF calculation
  outdir = './tmp',
                                   Directory for temporary files
                                ■ Must be .true. for EELS, otherwise absorption
  eels = .true.

    Number of calculated Lanczos coefficient.

  itermax0 = 500,
                              ■ Number of extrapolated Lanczos coefficient
  itermax = 500.
  extrapolation = 'no',
                             Type of extrapolation (bi-constant)
  epsil = 0.035,
                              The value of Lorentzian smearing in Ry
  units = 1.
                                  The units for start, end and increment
  start = 0.0d0,
                             ■ Minimum value of frequencies for a plot in eV
  end = 50.0d0,
                              Maximum value of frequencies for a plot in eV
  increment = 0.01d0

    Frequency step in eV
```

Example 3: Computed quantities

The code turbo_spectrum.x produces a file silicon.plot_eps.dat which contains the real and imaginary parts of the dielectric function and its inverse:

$$\operatorname{Re}[\varepsilon^{-1}(\mathbf{q},\omega)] = 1 + \frac{4\pi e^{2}}{|\mathbf{q}|^{2}} \operatorname{Re}[\chi(\mathbf{q},\mathbf{q};\omega)]$$

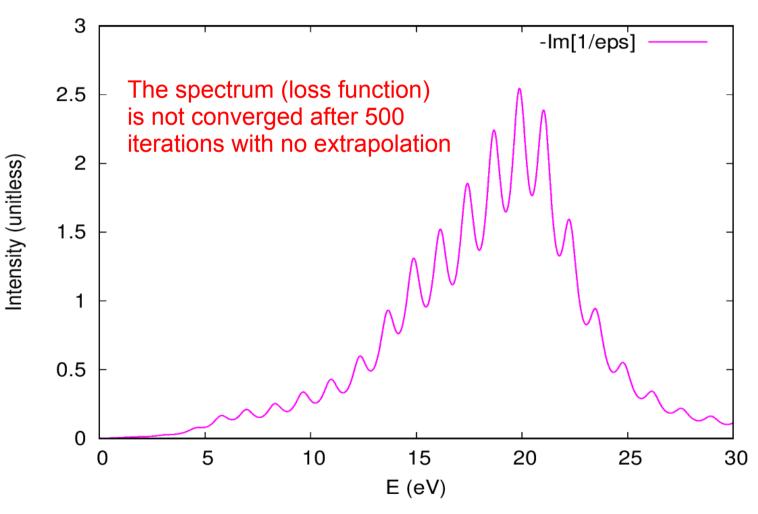
$$\varepsilon^{-1}(\mathbf{q},\omega) = 1 + \frac{4\pi e^{2}}{|\mathbf{q}|^{2}} \chi(\mathbf{q},\mathbf{q};\omega)$$

$$-\operatorname{Im}[\varepsilon^{-1}(\mathbf{q},\omega)] = -\frac{4\pi e^{2}}{|\mathbf{q}|^{2}} \operatorname{Im}[\chi(\mathbf{q},\mathbf{q};\omega)]$$

$$\varepsilon(\mathbf{q},\omega) = \frac{1}{\varepsilon^{-1}(\mathbf{q},\omega)} = \underbrace{\frac{\operatorname{Re}[\varepsilon^{-1}]}{(\operatorname{Re}[\varepsilon^{-1}])^2 + (\operatorname{Im}[\varepsilon^{-1}])^2}}_{\operatorname{Re}[\varepsilon(\mathbf{q},\omega)]} + i \underbrace{\frac{-\operatorname{Im}[\varepsilon^{-1}]}{(\operatorname{Re}[\varepsilon^{-1}])^2 + (\operatorname{Im}[\varepsilon^{-1}])^2}}_{\operatorname{Im}[\varepsilon(\mathbf{q},\omega)]}$$

Example 3: The computed EELS

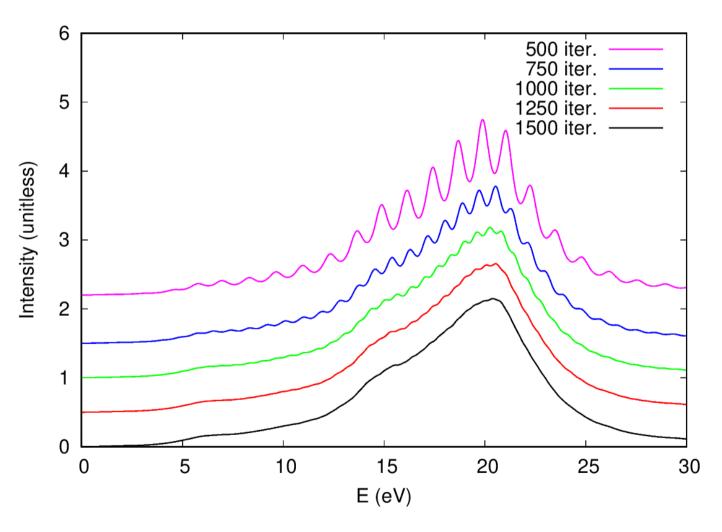
No extrapolation of Lanczos coefficients



Here, \mathbf{k} points mesh 10x10x10 1 1 1 is used.

Example 3: Convergence of EELS

No extrapolation of Lanczos coefficients



Here, \mathbf{k} points mesh 10x10x10 1 1 1 is used.

Step 3. Perform a spectrum calculation using the post-processing program turbo_spectrum.x and using the Lanczos coefficients computed in the previous step.

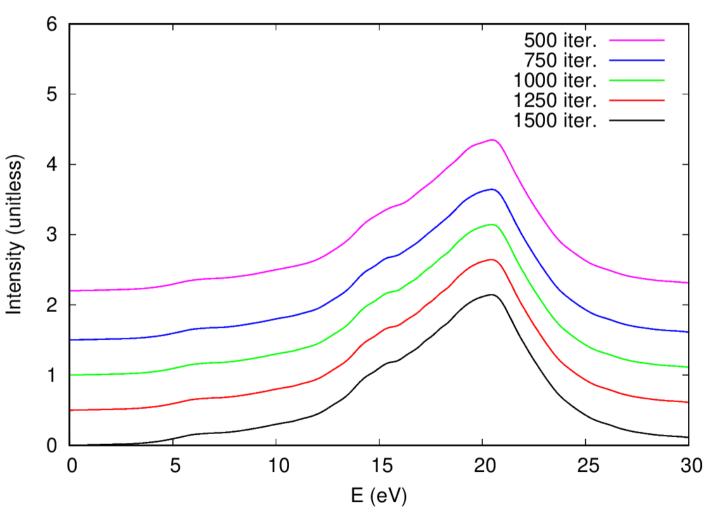
```
&lr input
  prefix = 'silicon',
                             ← The same prefix as in the SCF calculation
  outdir = './tmp',
                                  Directory for temporary files
                              Must be .true. for EELS, otherwise absorption
  eels = .true.

    Number of calculated Lanczos coefficient.

  itermax0 = 500,
                              ■ Number of extrapolated Lanczos coefficient
  itermax = 20000
  extrapolation = 'osc'
                              Type of extrapolation (bi-constant)
  epsil = 0.035,
                              The value of Lorentzian smearing in Ry
  units = 1.
                                  The units for start, end and increment
  start = 0.0d0,
                             ■ Minimum value of frequencies for a plot in eV
  end = 50.0d0,
                             ■ Maximum value of frequencies for a plot in eV
  increment = 0.01d0
                              Frequency step in eV
```

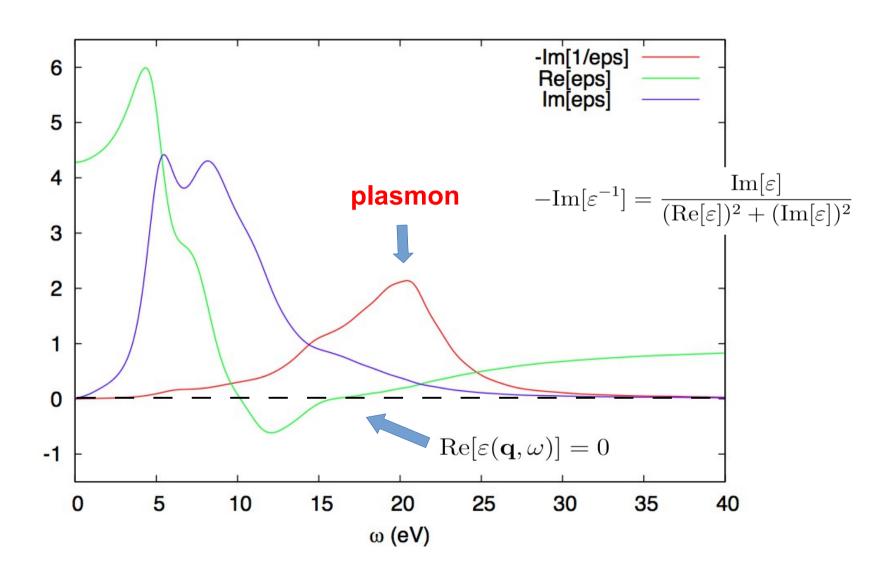
Example 3: Convergence of EELS

With extrapolation of Lanczos coefficients



Here, k points mesh 10x10x10 1 1 1 was used.

Example 3: Where is the plasmon peak?



Example 4

Coupled (resonant and anti-resonant) frequency-dependent **Sternheimer** equations:

$$\begin{split} (\hat{\mathbf{H}}_{\mathbf{k}+\mathbf{q}}^{\circ} - \varepsilon_{n,\mathbf{k}}^{\circ} - \omega) \, \tilde{\mathbf{u}}_{n,\mathbf{k}+\mathbf{q}}^{\prime}(\mathbf{r},\omega) \, + \, & \hat{\mathbf{P}}_{c}^{\mathbf{k}+\mathbf{q}} \, \tilde{\mathbf{v}}_{\mathrm{HXC},\mathbf{q}}^{\prime}(\mathbf{r},\omega) \, \mathbf{u}_{n,\mathbf{k}}^{\circ}(\mathbf{r}) \\ (\hat{\mathbf{H}}_{\mathbf{k}+\mathbf{q}}^{\circ} - \varepsilon_{n,\mathbf{k}}^{\circ} + \omega) \, \tilde{\mathbf{u}}_{n,-\mathbf{k}-\mathbf{q}}^{\prime \, *}(\mathbf{r},-\omega) \, + \, & \hat{\mathbf{P}}_{c}^{\mathbf{k}+\mathbf{q}} \, \tilde{\mathbf{v}}_{\mathrm{HXC},\mathbf{q}}^{\prime}(\mathbf{r},\omega) \, \mathbf{u}_{n,\mathbf{k}}^{\circ}(\mathbf{r}) \\ & & \text{interaction terms} \end{split} \end{split}$$

These equations can be re-written as a quantum Liouville equation and solved using the **Lanczos** algorithm:

Let's use the Sternheimer

$$(\omega - \mathcal{L}) \cdot \hat{\rho}_{q}^{\prime}$$
 algorithm $\hat{\rho}^{\circ}$

$$\mathcal{L} \cdot \hat{\rho}_{\mathbf{q}}' \equiv [\hat{H}^{\circ}, \hat{\rho}_{\mathbf{q}}'] + [\tilde{V}'_{\mathrm{HXC}, \mathbf{q}}, \hat{\rho}^{\circ}]$$

$$\chi(\mathbf{q},\mathbf{q};\omega)$$

Susceptibility



$$-\mathrm{Im}[\varepsilon^{-1}(\mathbf{q},\omega)] = -\frac{4\pi e^2}{|\mathbf{q}|^2}\,\mathrm{Im}[\chi(\mathbf{q},\mathbf{q};\omega)]$$

Loss function

Example 4: Input file for turboEELS

Step 1. Perform a SCF ground-state calculation for Si.

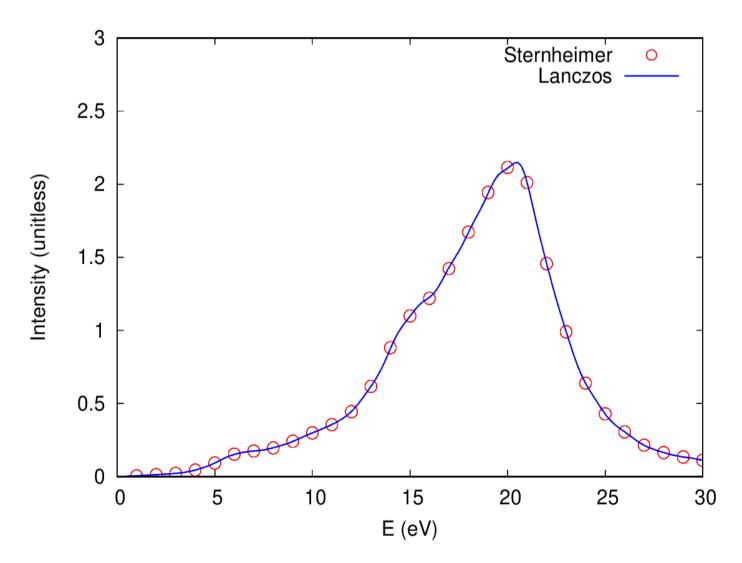
Step 2. Solve the Sternheimer equations using turbo_eels.x

```
&lr input
                                       Same as in the PWscf input
  prefix = 'silicon',
  outdir='./tmp'
&lr control
                                          Type of the algorithm that is used
  calculator = 'sternheimer'.
  q1 = 0.866.
  q2 = 0.000.
                                          Components of the transferred momentum q
  q3 = 0.000,
                                          The value of Lorenzian broadening in Ry
  epsil = 0.035.
                                          The units (0 \rightarrow Ry, 1 \rightarrow eV)
  start = 0.0d0.
  end = 30.0d0.
                                          Keywords that control the range in which the
  increment = 1.0d0
                                          spectrum is plotted
```

Example 4: The computed EELS

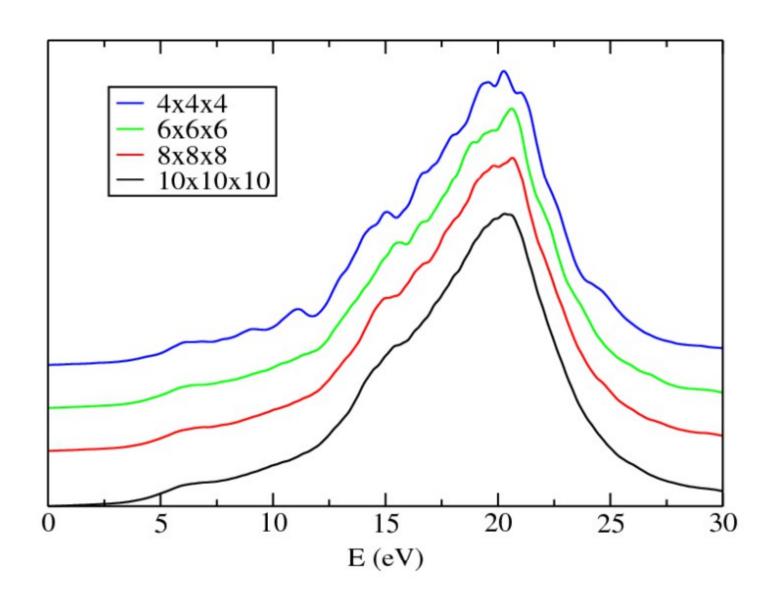
The Sternheimer and Lanczos algorithms give exactly the same spectrum.

Sternheimer is much slower than Lanczos!



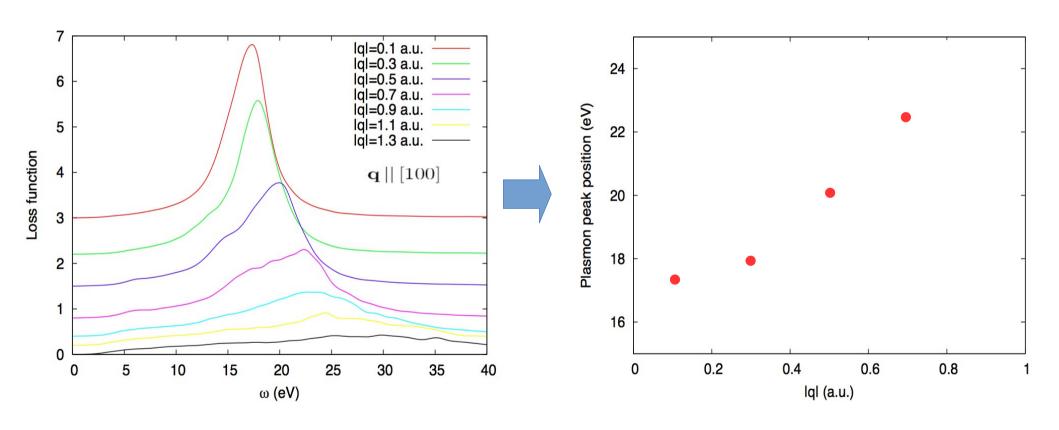
Example 4: Convergence w.r.t. k mesh

Convergence of EELS with respect to the size of the **k** points mesh



Example 4: Plasmon dispersion

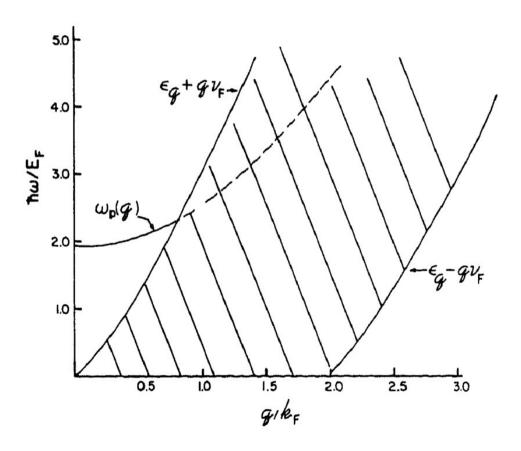
Plasmon dispersion in bulk silicon



By increasing **q**, the plasmon peak blue-shifts and broadens, because it enters in the electron-hole continuum.

Example 4: Plasmon dispersion and e-h continuum

The excitation region of the electron gas in (q,w) space



The plasmon line becomes strongly damped in a region of electron-hole pairs.

G. D. Mahan, Many-Particle Physics, Plenum Press, 2nd ed. New York, (1975)