

5 Measuring ultrashort laser pulses

The problem: In order to measure an event in time, you need have a shorter one (or: a temporal resolution which is high enough)

The fastest photodetectors have rise times in the low ps range (≈ 100 GHz) \rightarrow too slow!

\Rightarrow for fs light pulses: all optical methods needed!

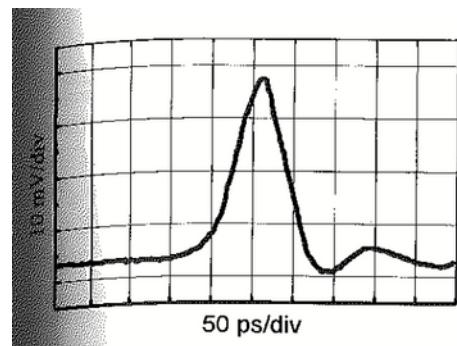
\rightarrow general approach:

1) map space on time ($1\text{fs} \hat{=} 0,3\mu\text{m}$)

2) make use of correlation function

Consider two time-dependent functions

$F(t)$, $F'(t)$ where $F'(t) < F(t)$, then



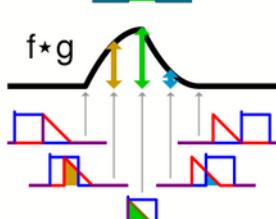
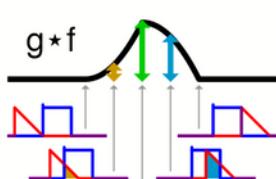
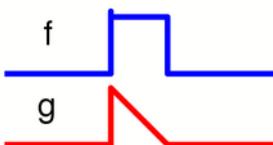
Measurement of a fs light pulse (from Ruillier)

$G(\tau)$ is

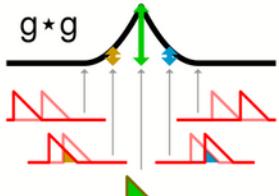
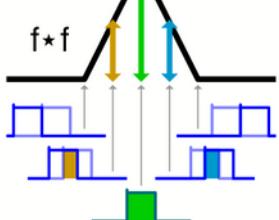
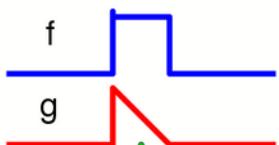
$$G(\tau) = \int_{-\infty}^{\infty} F'(t) F(t - \tau) dt$$

first order
correlation function

Cross-correlation



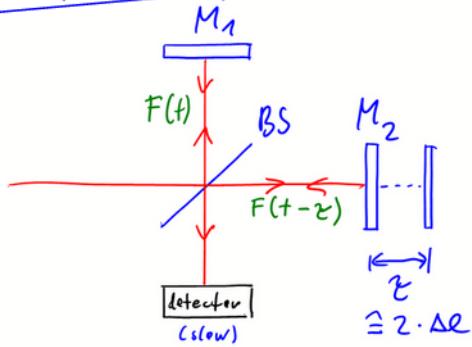
Autocorrelation



Cross- and autocorrelation /
Wikipedia /cmglee

If we have $F'(t)$ and $G(\tau)$, we also know $F(t - \tau)$. For fs lasers, we generally do not have a shorter test function \Rightarrow use of "autocorrelation", i.e., we correlate the pulse with itself. For two different pulses, we call this "cross correlation"

Exp. set up: Michelson interferometer



5.1. Optical correlation

$\vec{E}(t) \hat{=} \vec{F}(t)$, however our detector measures $I(t) = \langle \vec{E}(t) \cdot \vec{E}^*(t) \rangle$
 \Rightarrow Signal at detector:

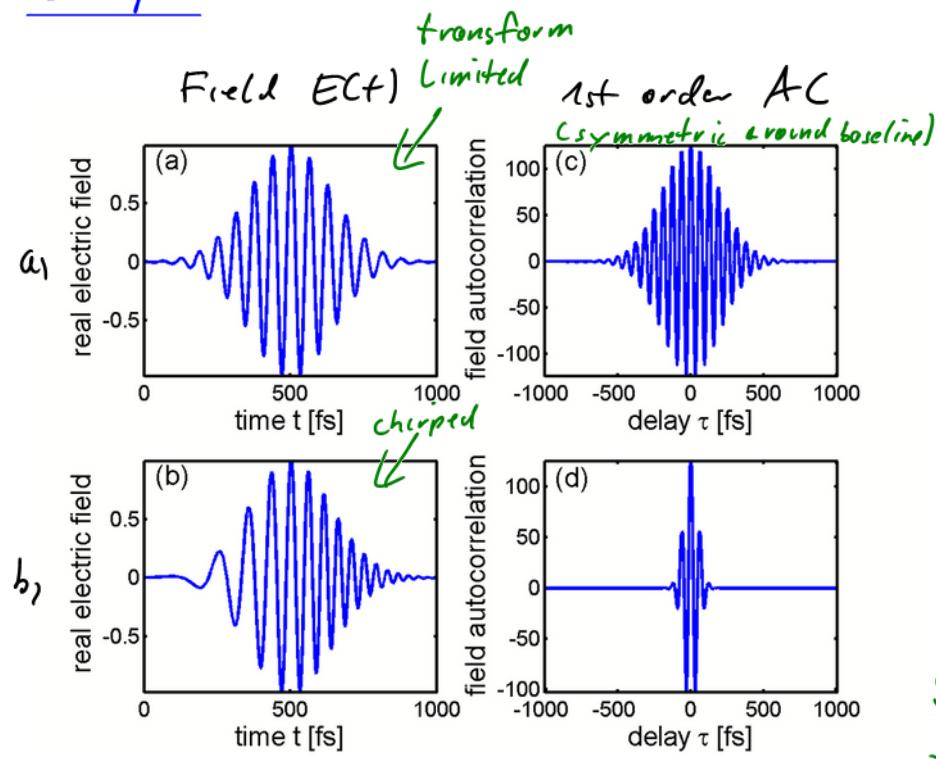
$$I_1(\tau) = \int_{-\infty}^{\infty} |\vec{E}(t) + E(t-\tau)|^2 dt$$

which, using $G(\tau)$, is $I_1(\tau) \propto 2 \int_{-\infty}^{\infty} I(t) dt + 2 G(\tau)$

$\underbrace{\int_{-\infty}^{\infty} |E(t)|^2 dt}_{\int_{-\infty}^{\infty} E(t) E(t-\tau) dt}$

Does the interference term for 1st order contain the pulse duration? A: NO! \rightarrow Measures only spectrum of E(t) (Wiener - Khinchin theorem) \rightarrow No phase information.

Example



Both pulses a, and b, have the same temporal duration, however the 1st order AC differs vastly!
 \rightarrow pulse b has higher bandwidth \Leftrightarrow shorter coherence length

Source : Wikipedia / Pgaboldi
 \rightarrow Article on optical correlation

\Rightarrow Need higher order correlation to learn about pulse duration, e.g.,

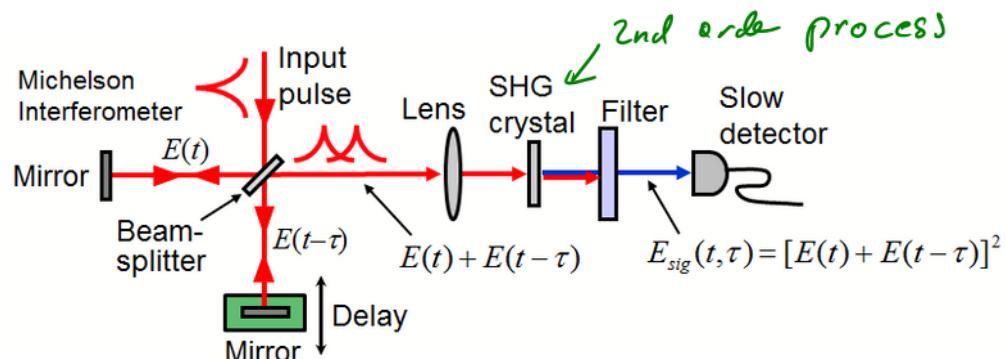
2nd order: Second harmonic generation, two photon absorption

$$h\nu < E_{gap}$$

3rd order: Kerr effect, third harmonic generation

Note: In principle, using only autocorrelation all order $G_n(\tau)$ need to be known to fully retrieve $\vec{E}(t)$, mostly (we hope) 2nd order is enough for realistic approximations.

5-2. 2nd order autocorrelator



Interferometric autocorrelator
Tresino / Ultrafast optics course

2nd order signal

$$I_2(\varepsilon) = \int_{-\infty}^{\infty} |[\vec{E}(t) + \vec{E}(t-\varepsilon)]|^2 dt = \dots$$

→ Ruillane,
Tresino

$$\begin{aligned} & \text{constant} \\ & \downarrow \\ & = \int_{-\infty}^{\infty} |2E^4 + 4E^2(t)E^2(t-\varepsilon) \\ & \quad \uparrow \text{intensity } A C \\ & + 4E(t)E(t-\varepsilon)[E^2(t) + E^2(t-\varepsilon)]\cos(\omega\varepsilon + \phi(t) - \phi(t-\varepsilon))| dt \end{aligned}$$

$$\vec{E}(t) = E(t) \cdot e^{i(\omega t + \phi(t))}$$

$$\begin{aligned} & + 4E(t)E(t-\varepsilon)[E^2(t) + E^2(t-\varepsilon)]\cos(\omega\varepsilon + \phi(t) - \phi(t-\varepsilon))| dt \\ & \quad \uparrow \text{modified interferogram of } E(t) \end{aligned}$$

$$\begin{aligned} & + 2E^2(t)E^2(t-\varepsilon)\cos[2(\omega\varepsilon + \phi(t) - \phi(t-\varepsilon))]| dt \\ & \quad \uparrow \text{interferogram of SHG} \end{aligned}$$

→ (possibly) very complicated signal, however

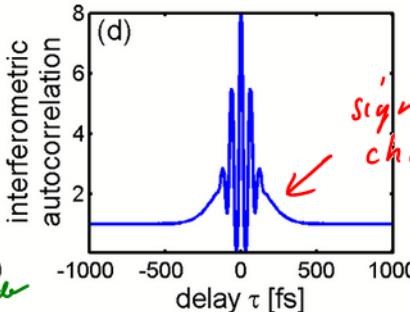
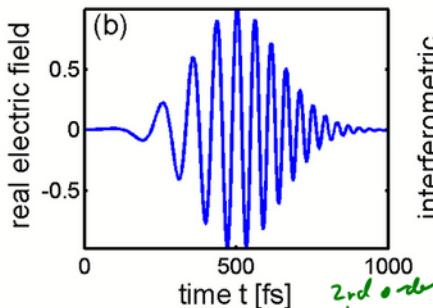
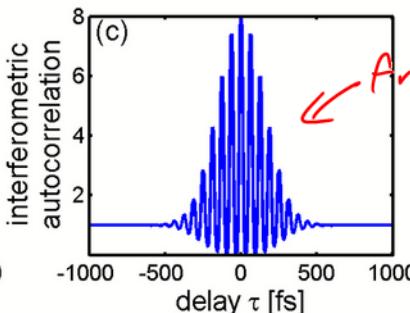
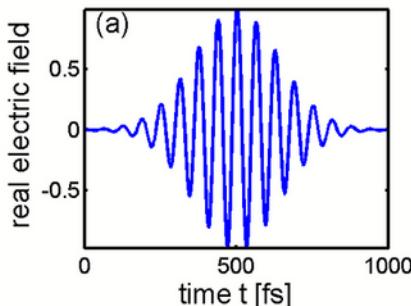
$$\begin{aligned} 1, I_2(\varepsilon=0) &= 2^4 \int \vec{E}^4(t) dt \\ 2, I_2(\varepsilon=\infty) &= 2 \int E^4(t) dt \end{aligned}$$

Note: $I_{\text{SHG}} \propto (I_{\text{Fund}})^2$

Field $E(t)$

Interferometric 2nd order AC

(7)



Interferometric AC / Wikipedia / PyroSOPP

Fringes only occur for equal optical frequencies of $E(t), E(t-\tau)$

→ chirp makes the pulse look shorter.

⇒ Determination of pulse duration easier using the so-called ..

Well aligned IAC (2nd order)

→ contrast 8:1

→ traces have to be
symmetric

→ Possibility to learn about
presence of chirp

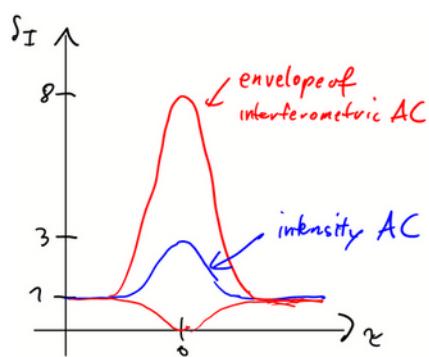
→ Can extract pulse duration
Signal analysis complicated,
as chirp changes fringe pattern

Intensity auto correlation

Either numerical filtering or experimentally washing signal out (e.g. sweep fast) yields intensity AC:

$$S_I = 2 \int I^2 dt + 4 \int \overset{\uparrow}{\text{background}} I(t) I(t-\tau) dt$$

2nd order T intensity AC
(no phase terms)

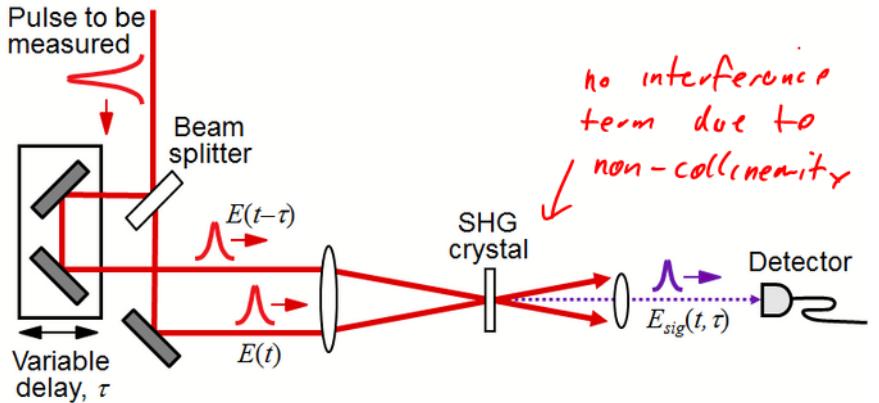


Contrast ratio 3:1, again symmetric in τ

Most common, because easy to use, pulse width given by "Full width at half maximum" (FWHM) and convolving factor.

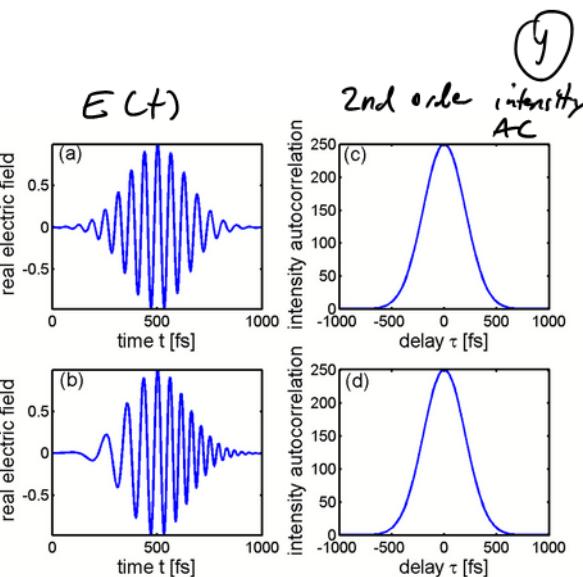
Alternatively, one may use a non-collinear arrangement, which directly yields the intensity AC without "background".

(drawback: no check for proper alignment)



2nd order intensity AC in
non-collinear geometry / Tresino

Here, we directly see that both pulses have some intensity AC,
in order to obtain a pulse duration, we now have to assume a pulse
shape!



Background Area
intensity AC / Wikipedia

Signal analysis

If we have a pulse with FWHM Δt , the width of the AC function $\Delta \tau$ is connected to this by a factor depending on the pulse shape. Today, the common choice is gaussian ($\propto e^{-t^2}$) and $\Delta \tau_{\text{Gauss}} = \sqrt{2} \Delta t$. In earlier literature often a sech² pulse is used with $\Delta \tau_{\text{sech}^2} = 1.54 \Delta t$ yielding the shortest pulses.

(end of lecture)
→ FR06 slides

Summary: Using ^{2nd order} AC, we can obtain a measure of our pulse duration by guessing a pulse shape.

Conclusions from Summary: 1) However, we do not know the accuracy of that guess
 → it might be terribly bad

2) We also know very little about the pulse phase, no quantitative analysis is possible from AC!

5.3 Frequency-resolved optical gating

(10)

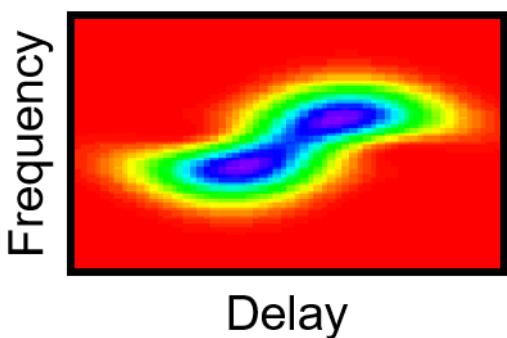
Better ways of characterizing ultrashort pulses measure the AC function spectrally resolved. This is called frequency resolved optical gating or short FROG. Basically, we replace the photodetector with a spectrometer.

For SHG-Frog, e.g., we obtain

$$I_{\text{Frog}}(w, z) = \left| \int_{-\infty}^{\infty} E_{\text{sig}}(t, z) \exp(iwt) dt \right|^2$$

where

$$E_{\text{sig}}^{\text{SHG}}(t, z) = E(t) E(t - z) \quad (\text{here } E(t - z) = G(z))$$



FROG trace of pulse with self phase modulation / Tresino

As we can see, the FROG trace is in the (ω, z) -plane,
a projection onto z yields the intensity AC.

What is the gain?

- For some FROG geometries, you immediately understand second order chirp (and possibly 3rd order) from traces

- In contrast to AC you can iteratively retrieve amplitude + phase from FROG traces, i.e., reconstruct E(t) using suitable algorithms for 2D phase retrieval

\Rightarrow See, e.g., frog.gatech.edu \Rightarrow Talks \Rightarrow Ultrafast Optics course
 \Rightarrow English lectures \Rightarrow UFD-M-FAOG.ppt

\rightarrow Some excerpts follow now

\rightarrow Advertisement: Master Lab course FM.ULP \rightarrow Spatio-temporal...