

Master Advanced Lab Course
Universität Göttingen – Fakultät für Physik

Report on
the experiment KT.WZE

W/Z experiment at the Tevatron

Name:	Eric Bertok
Email:	eric.bertok@stud.uni-goettingen.de
Conducted on	24th January 2018
Assistant:	Dr. J. Veatch
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1 Introduction

In goal of this experiment is the determination the of the branching ratio of the W boson $\text{BR}(W \rightarrow \mu\nu)$. First, W and Z bosons are reconstructed using data provided by the Tevatron collider at Fermilab. By comparing with Monte Carlo simulations, selection parameters are obtained, which allow for clean cuts for filtering out background events (jets and cosmic source). The mass and the transverse mass is then determined for the Z and W boson respectively. Finally the branching ratio is calculated from the number of selected events, the trigger efficiencies, as well as the reconstruction efficiencies.

2 Theory

2.1 Electroweak interaction

The GWS theory (Glashow, Weinberg, Salam) is the unified description of both the electromagnetic force mediated by the photon and the weak interaction mediated by the massive W^+ , W^- and the neutral neutral Z boson. It was confirmed experimentally in the 1970s [3]. The gauge bosons are introduced by means of a local $\text{SU}(2)_L$ gauge symmetry in a weak isospin space. The weak isospin doublets are formed by fermions differing by one unit of charge [4, p.;416]. By also replacing the $\text{U}(1)$ symmetry by a new $\text{U}(1)_Y$ symmetry with the “hypercharge” Y , the neutral Z boson can be identified by a linear combination of the neutral $W^{(3)}$ boson and the B boson coupling to the hypercharge. More details can be found in [4, p. 418ff]. Being a charged boson, the W bosons couple to fermions differing by one unit of charge. Furthermore it maximally violates parity as it only couples to left-handed particles and right-handed antiparticles. The vertex factor is given by [4, p.;409]

$$-i \frac{g_W}{\sqrt{2}} \frac{1}{2} \gamma^\mu (1 - \gamma^5), \quad (2.1)$$

where g_W is the weak coupling constant and γ^μ are the gamma matrices. The Z boson however, couples to any pair of identical fermions, albeit coupling more strongly to left handed ones. This becomes apparent in the form of the vertex factor: [4, p. 432]

$$-i \frac{1}{2} g_Z \gamma^\mu (c_V - c_A \gamma^5), \quad (2.2)$$

with the vector and axial vector couplings c_V and c_A .

2.2 Matrix elements and Decay rates

The matrix elements for the electroweak interaction can be calculated with the appropriate Feynman rules. After averaging over the three possible polarizations, the spin-averaged matrix element squares is obtained for both the W and the Z boson decaying to a lepton and its neutrino or a lepton- anti-lepton pair, respectively [4, p.;242,411]:

$$\langle |\mathcal{M}_W^2| \rangle = \frac{1}{3} g_W^2 m_W^2 \quad (2.3)$$

$$\langle |\mathcal{M}_Z^2| \rangle = \frac{1}{3} (c_V^2 + c_A^2) g_Z^2 m_Z^2. \quad (2.4)$$

These can be inserted into the decay rate formula: [4, p. 411]

$$\Gamma = \frac{p^*}{32\pi^2 m^2} \int \langle |\mathcal{M}^2| \rangle d\Omega = \frac{p^*}{8\pi m^2} \langle |\mathcal{M}^2| \rangle, \quad (2.5)$$

where m is the mass of the boson and p^* is the momentum of the lepton in the center of mass frame. One can argue that $p^* = m_Z/2$, as the decay happens in the centre of mass frame of the decaying particle. Therefore the decay rate is

$$\Gamma(W^- \rightarrow e^- \bar{\nu}_e) = \frac{g_W^2 m_W}{48\pi}. \quad (2.6)$$

$$\Gamma(Z \rightarrow e^- e^+) = \frac{g_Z^2 m_Z}{48\pi} (c_V^2 + c_A^2). \quad (2.7)$$

Lepton universality tells us that this is the same for all three leptonic channels when neglecting masses. For hadronic processes, the CKM matrix has to be considered, while excluding the top quark, as it is too massive. For the W boson, one obtains for the decay width [cite]

$$\Gamma_W = (3 + 6\kappa)\Gamma(W^- \rightarrow e^- \bar{\nu}_e) \approx 9.2 \frac{g_W^2 m_W}{48\pi} = 2.1 \text{ GeV}. \quad (2.8)$$

$\kappa \approx 1.038$ is a correction factor that accounts for second order QCD processes. Similarly, for the Z boson, one obtains

$$\Gamma_Z \approx 2.5 \text{ GeV}. \quad (2.9)$$

The branching ratios for the muon channel are therefore

$$BR(W \rightarrow \mu \bar{\nu}_\mu) = 10.8\%, \quad (2.10)$$

$$BR(Z \rightarrow \mu^+ \mu^-) = 3.5\%. \quad (2.11)$$

2.3 Invariant and transverse mass

For the Z boson one can calculate the functional form of the invariant mass peak by taking into account its finite lifetime. The cross section for a $q\bar{q} \rightarrow \mu^+ \mu^-$ event is proportional to [cite]

$$\sigma \propto |\mathcal{M}|^2 \propto \left| \frac{1}{q^2 - m_Z^2 + im_Z \gamma_Z} \right|^2 = \frac{1}{(q^2 - m_Z^2)^2 + m_Z^2 \gamma_Z^2}, \quad (2.12)$$

which is a Breit-Wigner curve. q is the invariant mass of both muons. As both can be detected in such an event, the Breit-Wigner-curve can be fitted directly to the selected data to obtain the mass of the Z boson. For the W boson, things are more complicated. Due to the W events only having one muon, the undetectable neutrino has to be reconstructed from the missing momentum. For a hadronic collider such as the tevatron, the total centre of mass energy cannot be known on an event to event basis [cite] due to the composite nature of the hadrons. More specifically, the z -momentum of the interacting partons are unknown, making the invariant mass reconstruction impossible. However, one can define the transverse mass M_T , which can be calculated from the reconstructed transverse momentum of the neutrino \vec{p}_T^μ . First, the missing transverse energy MET is determined as

$$MET \approx |\mathbf{p}_T^\nu| = |-\mathbf{p}_T^\mu - \mathbf{u}_T|, \quad (2.13)$$

where \vec{u}_T is the transverse momentum of the hadrons [cite]. The transverse mass is then defined as

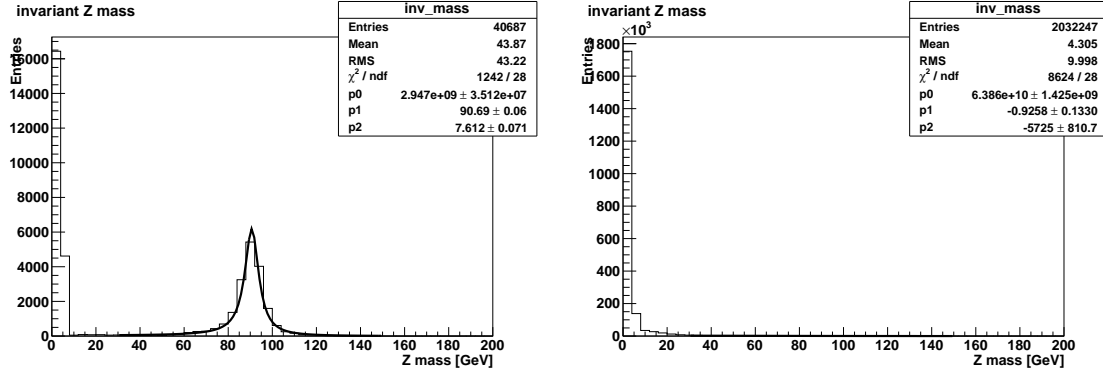
$$M_T = \sqrt{(MET + \mathbf{p}_T^\mu)^2 - (MET_x + p_x^\mu)^2 - (MET_y + p_y^\mu)^2}. \quad (2.14)$$

This quantity is lorentz invariant but does not peak at M_Z . However, the W mass can be read off from the position of the dropoff, as the longitudinal component of the invariant mass is then close to zero [cite].

3 Experimental setup and methods

For this experiment, 2 types of data have been provided: The first is a subset of real data from the DØ detector at Fermilab near Chicago. The second are two sets of simulated monte-carlo W and Z events generated by PYTHIA [cite].

At the DØ detector, muons are identified both in the muon detector and the tracking system. Whereas the tracking system directly surrounds the interaction point and allows for gauging the muon's momentum and direction precisely, the outer muon detector is mostly used to match the track in the tracking system. This is possible because muons are the only particles capable of reaching the muon detector, due to a combination of their relatively long lifetime and small calorimeter energy deposition. [cite?] For more details on the DØ detector see [2]. As the sheer amount of data from the collider is too much to analyse and save, both software and hardware triggers are used to decide, whether an event is worth investigating. The data that is provided is pre-filtered for at least one muon in every event that has a transverse momentum of at least 15 GeV/c. [??? zusätzlicher trigger] The monte-carlo data has been reconstructed, such that the events look like real data. Therefore the monte-carlo serves a benchmark for the analysis of the real data. To obtain invariant mass peaks for the Z - as well as the transverse mass peak for the W



boson, the right events have to be selected out of the 3 million events provided. This happens by first fixing a set of object level cuts, which define what is counted as a muon. These cuts are counted towards our trigger conditions for the muons and are the same for both W - and Z boson analysis [???]. Secondly, a range of object level cuts have to be performed to single out the right events for Z and W production respectively. These have to be physically motivated by keeping in mind what is expected for the muons in a W or Z decay and should also be compared to the monte-carlo simulation. By first comparing various muon parameters for the simulated data and the uncut experimental data, one can define appropriate cuts in these parameters that cut out most events not present in the simulation while also keeping events that do resemble the simulation mostly intact. Finally the real data is plotted with these cuts performed and compared with the simulated results. This process is repeated until a satisfactory isolation of Z or W events has been produced. All of the analysis is done with ROOT.

4 Analysis

4.1 Selection of $Z \rightarrow \mu\mu$ events

The Z boson mass distribution for the simulated monte-carlo data is shown in fig. 1(a). One can clearly see the mass peak of the Z boson. A large number of events, however, is also situated at the beginning of the mass spectrum. A fit performed by root with the function 4.1

$$\sigma \propto \frac{1}{(q^2 - m_Z^2)^2 + m_Z^2 \Gamma_Z^2} \quad (4.1)$$

is performed, with the parameters p_0 as a constant proportionality factor, p_1 being m_Z and p_2 being the decay width Γ . Apart from the large number of events at low mass, this clearly resembles the expected Z mass peak. In fig. 1(b), the uncut real data is plotted. From the 2 million events, almost all of them give a very small Z mass and the peak is not visible. Note that the trigger “TRIG_MUW_W_L2M3_TRK10” [1] is still included in the uncut case.

4.2 Cosmics

4.3 Zmass

4.4 Reconstructing and selecting W events

4.5 Determination of efficiencies

4.6 Determination of the $BR(W \rightarrow \mu\nu)$

5 Discussion

References

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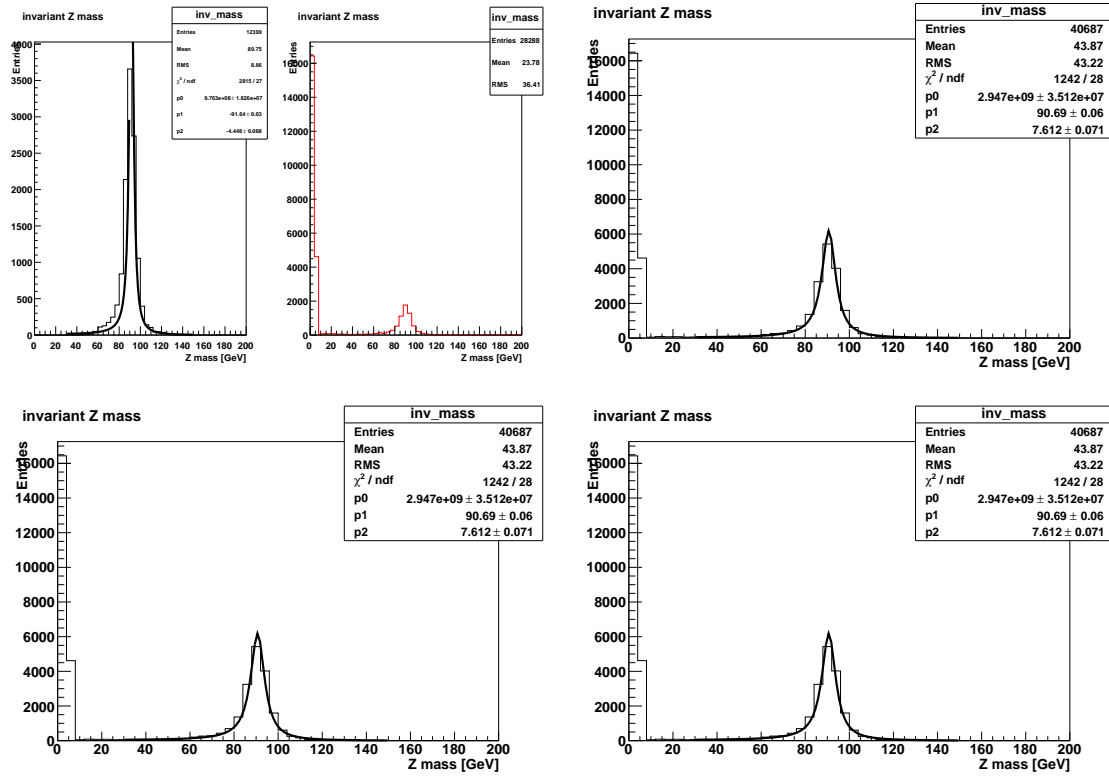


Figure 1: The l-o-n-g caption for all the subfigures (FirstFigure through FourthFigure) goes here.

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