

Internal Combustion Engines I: Fundamentals and Performance Metrics

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2018 Princeton-Combustion Institute Summer School on Combustion Course Length: 9 hrs

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Short course outline:

Internal Combustion (IC) engine fundamentals and performance metrics, computer modeling supported by in-depth understanding of fundamental engine processes and detailed experiments in engine design optimization.

Day 1 (Engine fundamentals)

Hour 1: IC Engine Review, Thermodynamics and 0-D modeling

Hour 2: 1-D modeling, Charge Preparation

Hour 3: Engine Performance Metrics, 3-D flow modeling

Day 2 (Computer modeling/engine processes)

Hour 4: Engine combustion physics and chemistry

Hour 5: Premixed Charge Spark-ignited engines

Hour 6: Spray modeling

Day 3 (Engine Applications and Optimization)

Hour 7: Heat transfer and Spray Combustion Research

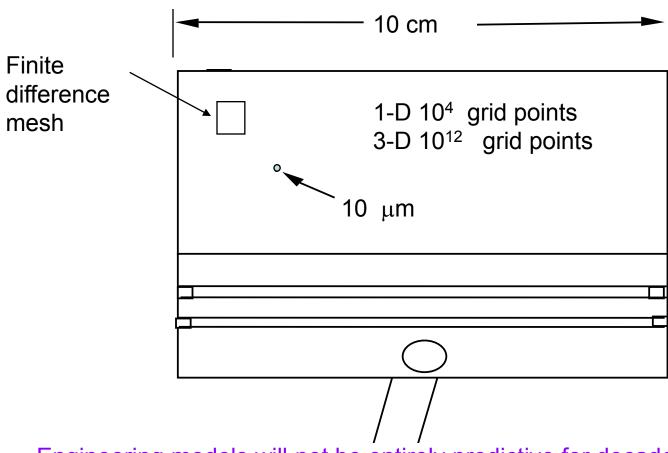
Hour 8: Diesel Combustion modeling

Hour 9: Optimization and Low Temperature Combustion





Resolution – predictive spray models?



Engineering models will not be entirely predictive for decades Accurate submodels will be needed for detailed spray processes

(e.g., drop drag, drop turbulence interaction, vaporization, atomization, drop breakup, collision and coalescence, and spray/wall interaction)





Governing Equations

Amsden, 1997

Gas phase Liquid phase Turbulence

Gas void fraction and drop number density

Lagrangian Drop, Eulerian Fluid (LDEF) models $\theta = 1 - \iint_{V_d} (\iiint_{A} \frac{4}{3} \pi r^3 f \, dr \, d\mathbf{v} \, dT_d) \, dVol / Vol$

Two-Phase Flow Regimes

Computational cell

Computational cell

Computational cell

Drop parcels

Intact

Churning

Thick

Thin

Very thin

Current LDEF spray

models:



LDEF Spray Modeling

Concept of using "drop parcels"

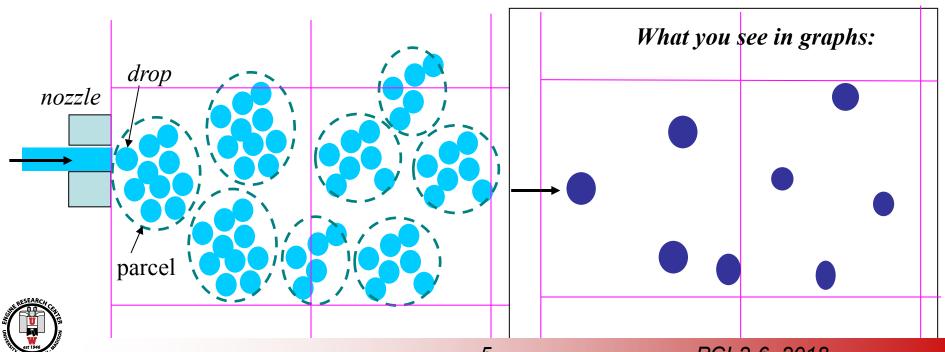
For typical heavy-duty diesel, injected fuel per cycle (75% load): 0.160 g

One spray plume: m_{fuel} =0.160/6=0.0267 g

If average SMD=10 μ m \rightarrow m_{drop} =3.8x10⁻¹⁰ g

of drops in the domain= $0.0267g/m_{drop}=7.1x10^7$

Impractical to track individual fuel drops – group identical drops into 'parcels'





Liquid Phase

Amsden, 1997

Spray drop number conservation $f = f(\mathbf{x}, \mathbf{v}, \mathbf{r}, T_d, \mathbf{y}, \mathbf{y}; \mathbf{t})$

$$\frac{\partial f}{\partial t} + \nabla_{\mathbf{x}} \cdot (f\mathbf{v}) + \nabla_{\mathbf{v}} \cdot (f\mathbf{F}) + \frac{\partial}{\partial r} (fR) + \frac{\partial}{\partial T_d} (f\dot{T}_d) + \frac{\partial}{\partial y} (fy) + \frac{\partial}{\partial y} (f\dot{y}) = \dot{f}_{coll} + \dot{f}_{bu}$$

 $F=d\mathbf{v}/dt$ drop drag

R = dr/dtVaporization and heating

Drop distortion Drop breakup, coalescence

Spray exchange functions

$$\mathbf{F}^{s} = -\int f \rho_{d} \left(4/3 \, \pi r^{3} \mathbf{F}' + 4\pi r^{2} R \mathbf{v} \right) d\mathbf{v} \, dr \, dT_{d} \, dy \, dy$$

$$\mathbf{Q}^{s} = -\int f \rho_{d} \left\{ 4\pi r^{2} R \left[I_{l} + \frac{1}{2} (\mathbf{v} - \mathbf{u})^{2} \right] + 4/3 \, \pi r^{3} \left[c_{l} T_{d} + \mathbf{F}' \cdot (\mathbf{v} - \mathbf{u} - \mathbf{u}') \right] \right\} d\mathbf{v} \, dr \, dT_{d} \, dy \, dy$$

Work done by drop drag forces $W^{s} = -\int f \rho_d 4/3 \pi r^3 \mathbf{F'} \cdot \mathbf{u'} d\mathbf{v} dr dT_d dy dy$



W

Lagrangian drop - liquid phase

Amsden, 1997

Discrete Drop Model

drop position

$$\frac{d\mathbf{x}}{dt} = \mathbf{v}$$

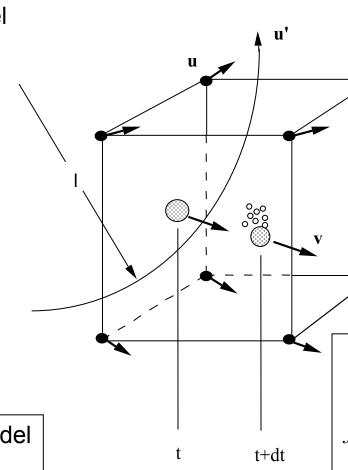
drop velocity

$$\frac{d\mathbf{v}}{dt} = \mathbf{F}$$

drop size

$$\frac{d\mathbf{r}}{dt} = \mathbf{R}$$

Turbulence model provides: *l*, *u* '



Spray submodels provide:

F - Drag, R - Vaporize

$$f_c$$
 + f_b - breakup/collide

Initial data:

v, *r*, T_d – Atomization model





Eulerian Gas Phase

Amsden,1997

Mass conservation (species)

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) = -\iiint \rho_l 4\pi r^2 R f \, dr \, d\mathbf{v} \, dT_d$$

$$R = \frac{dr}{dt} - \text{Vapor source}$$

Momentum conservation

$$\frac{\partial \rho \mathbf{u}}{\partial t} + \nabla \cdot (\rho \mathbf{u} \mathbf{u}) = -\nabla p - \nabla (\frac{2}{3} \rho k) + \nabla \tau + F^s + \rho \mathbf{g}$$
Turbulent and viscous stress

Rate of momentum gain due to spray – drop drag





Amsden, 1989

Gas Phase energy conservation

Internal energy conservation

$$\frac{\partial \rho I}{\partial t} + \nabla \cdot (\rho \mathbf{u} I) = -P \nabla \cdot \mathbf{u} - \nabla \cdot \mathbf{J} + \rho \varepsilon + Q^c + Q^s$$

Turbulence dissipation

Combustion

heat release

Energy due to Spray - vaporization

Heat flux

$$\mathbf{J} = -\lambda \nabla T - \rho D \sum_{m} h_{m} \nabla (\rho_{m} / \rho)$$

Equations of state

$$p = RT\sum_{m} \rho_{m} / W_{m}$$

Specific heat, enthalpy from JANAF data





Turbulence Model (RANS)

Amsden, 1989, 1997

Kinetic energy

Dissipation

$$\frac{\partial \rho k}{\partial t} + \nabla \cdot (\rho \mathbf{u} k) = -\frac{2}{3} \rho k \nabla \cdot \mathbf{u} + \mathbf{\tau} \cdot \mathbf{u} + \nabla \cdot \left[(\frac{\mu}{\Pr_k}) \nabla k \right] - \rho \varepsilon + W$$

Production due to mean flow

Rate of work to disperse drops

Dissipation rate

$$\frac{\partial \rho \varepsilon}{\partial t} + \nabla \cdot (\rho \mathbf{u} \varepsilon) = -\left(\frac{2}{3} \mathbf{C}_{\varepsilon 1} - \mathbf{C}_{\varepsilon 3}\right) \rho \varepsilon \nabla \cdot \mathbf{u} + \nabla \cdot \left[\left(\frac{\mu}{\mathbf{P} r_{\varepsilon}}\right) \nabla \varepsilon\right]
+ \frac{\varepsilon}{\mathbf{k}} \left[\mathbf{C}_{\varepsilon 1} \boldsymbol{\tau} : \nabla \mathbf{u} - \mathbf{C}_{\varepsilon 2} \rho \varepsilon + \mathbf{C}_{s} W^{s}\right]$$

Turbulence diffusivity

$$D = C_{\mu} k^2 / \varepsilon$$

Eddy size

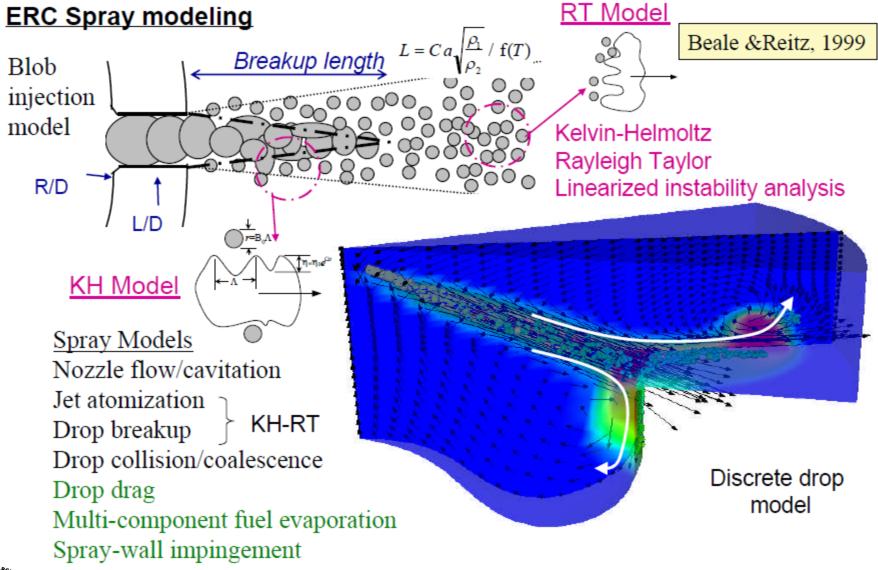
$$l = C \frac{3/2}{k}$$

Turbulence intensity

$$u'^2 = (2 k/3)$$



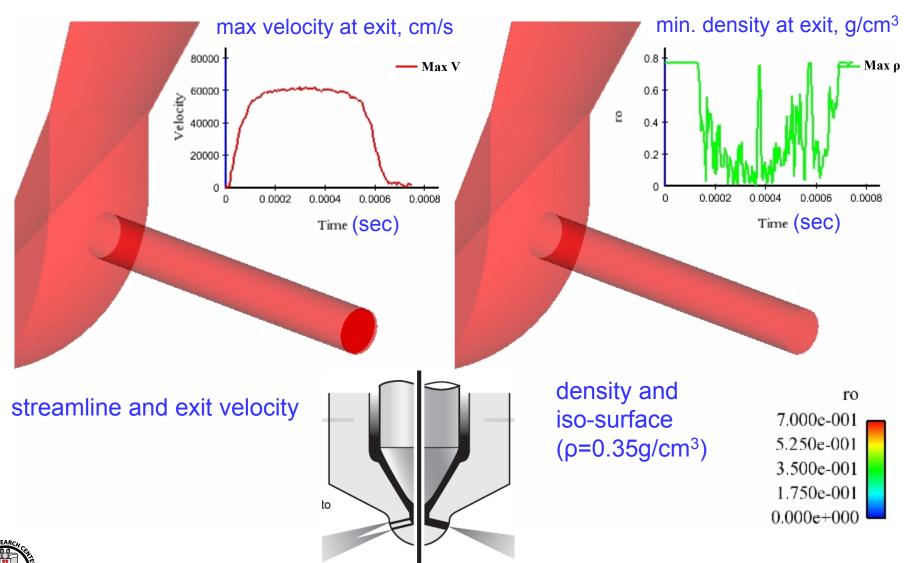








Fuel Injection and Nozzle Cavitation





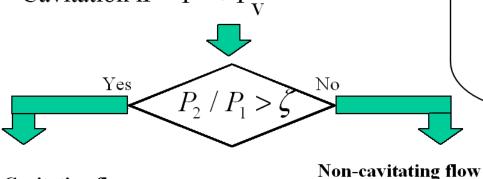
Initial



Cavitation inception

Account for effects of nozzle geometry

Cavitation if $P < P_V$

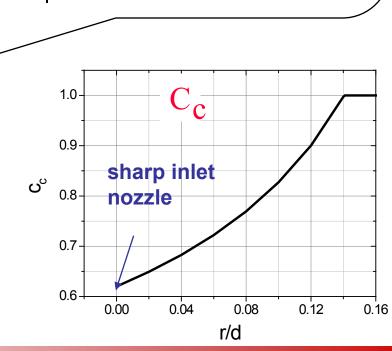


Cavitating flow

$$\zeta = \frac{1}{2(C_c - C_c^2)}$$

Contraction coefficient (Nurick (1976)

$$C_c = \left[\left(\frac{1}{0.62} \right)^2 - 11.4r / d \right]^{-1/2}$$



Cavitation

region

1/d

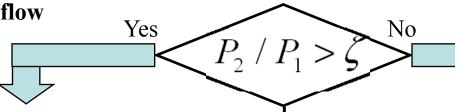
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Non-cavitating flow

ERC Nozzle Flow Model

Cavitating flow



Nozzle discharge coefficient

$$C_d = C_c \sqrt{\frac{p_1 - p_v}{p_1 - p_2}}$$

Effective injection velocity

$$u_{eff} = \frac{2C_c P_1 - P_2 + (1 - 2C_c)P_v}{C_c \sqrt{2\rho(P_1 - P_v)}}$$

Effective nozzle area

$$A_{eff} = \frac{2C_c^2(P_1 - P_v)}{2C_cP_1 - P_2 + (1 - 2C_c)P_v} A$$

Nozzle discharge coefficient

Lichtarowicz (1965)

$$C_d = 0.827 - 0.0085 \frac{l}{d}$$

Effective injection velocity

$$u_{eff} = C_d \sqrt{\frac{2(P_1 - P_2)}{\rho}}$$

Effective nozzle area

$$A_{eff} = A$$



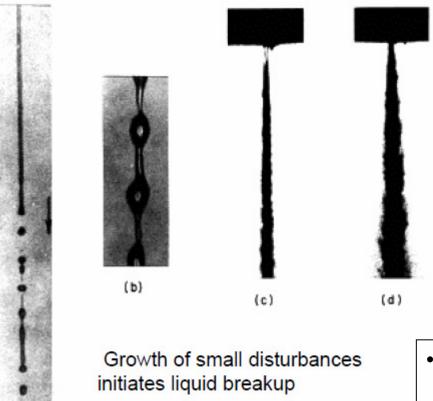


Atomization models (Single hole nozzle)

Reitz, 1982

Four main jet breakup regimes:

Rayleigh, first wind-induced, second wind-induced and atomization



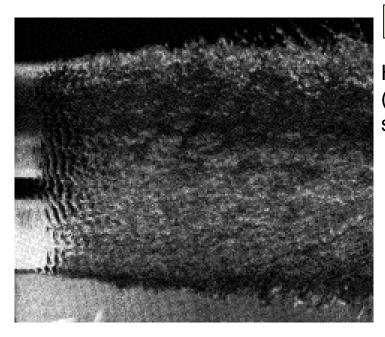
a.) Rayleigh breakup
Drop diameters > jet diameter.
Breakup far downstream nozzle
b.) First wind-induced regime
Drop diameter ~ jet diameter.
Breakup far downstream of nozzle

c.) Second wind-induced regime
Drop sizes < jet diameter.
Breakup starts close to nozzle exit
d.) Atomization regime
Drop sizes << jet diameter.
Breakup at nozzle exit.

- Jet breakup known to depend on nozzle design details.
- Need to start by considering flow in the injector nozzle passage

Atomization - "Wave" breakup model

Reitz, 1982



Taylor & Hoyt, 1983

High speed photograph of water jet close to nozzle exit (at top) in the second wind-induced breakup regime showing surface wave instability growth and breakup

Kelvin-Helmholtz Jet Breakup Model

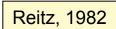
$$\begin{array}{c|c}
\uparrow & \lambda \\
\hline
 & \uparrow \\
 &$$

Linear Stability Theory:

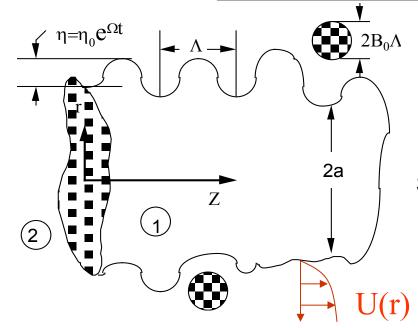
Cylindrical liquid jet issuing from a circular orifice into a stationary, incompressible gas.

Relate growth rate, ω , of perturbation to wavelength $\lambda = 2\pi/k$









Linearized analysis

U = Jet velocity

Surface waves breakup on jet or "blob"

$$\eta = R \left(\eta_0 e_{ikz} + \omega t \right)$$

Equation of liquid surface: $r = a + \eta$,

Axisymmetric fluctuating pressure, axial velocity, and radial velocity for both liquid and gas phases.

Fluctuations described by continuity equation $\frac{\partial u_i}{\partial z} + \frac{1}{r} \frac{\partial}{\partial r} (rv_i) = 0$

$$\frac{\partial u_i}{\partial z} + \frac{1}{r} \frac{\partial}{\partial r} (rv_i) = 0$$

plus linearized equations of motion for the liquid and the gas,



Axial:

$$\frac{\partial u}{\partial t} + U_{i}(r) \frac{\partial u}{\partial z} + v_{i} \frac{dU}{dr} = -\frac{1}{\rho} \frac{\partial p}{\partial z} + \frac{\mu}{\rho} \left[\frac{\partial^{2} u}{\partial z^{2}} + \frac{1}{r} \frac{\partial \left(\partial u \right)}{\partial r} \right]$$

Analysis (Cont.)

Reitz, 1982

Radial:
$$\frac{\partial v_i}{\partial t} + U_i(r) \frac{\partial v_i}{\partial z} = -\frac{1}{\rho_i} \frac{\partial p_i}{\partial r} + \frac{\mu_i}{\rho_i} \left[\frac{\partial^2 v_i}{\partial z^2} + \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial r v_i}{\partial r} \right) \right]$$

Gas is assumed to be inviscid U(r) = U - slip

With $\eta \ll a$, the gas equations give the pressure at the interface r=a

$$p_2 = -\rho_2 (U - i\frac{\omega}{k})^2 k \eta \frac{K_0(ka)}{K_1(ka)}$$

Boundary conditions-

Kinematic, tangential and normal stress at the interface:

$$\mathbf{v}_1 = \mathbf{w} = \frac{\partial \eta}{\partial t}, \qquad \frac{\partial \mathbf{u}_1}{\partial r} = -\frac{\partial \mathbf{v}_1}{\partial z}$$

$$-p_1 + 2v_1\rho_1 \frac{\partial v_1}{\partial r} - \frac{\sigma}{a^2} (\eta + a^2 \frac{\partial^2 \eta}{\partial z^2}) + p_2 = 0$$





Dispersion relationship

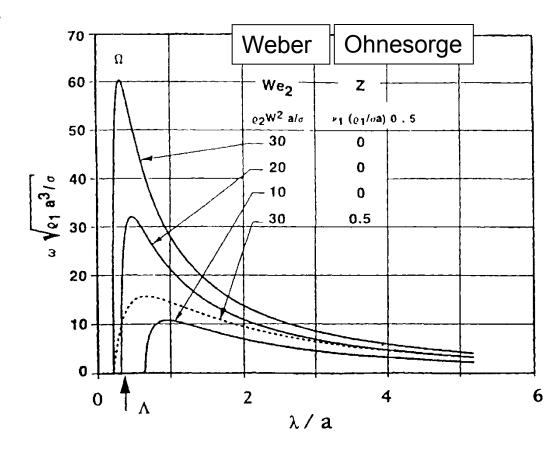
Reitz, 1988

$$\omega^{2} + 2v_{1}k^{2}\omega \left[\frac{I_{1}(ka)}{I_{0}(ka)} - \frac{2kl}{k^{2} + l^{2}}\frac{I_{1}(ka)}{I_{0}(ka)}\frac{I_{1}(la)}{I_{0}(la)}\right] = \frac{\sigma k}{\rho_{1}a^{2}}\left(1 - k^{2}a^{2}\right)\left(\frac{l^{2} - k^{2}}{l^{2} + k^{2}}\right)\frac{I_{1}(ka)}{I_{0}(ka)}$$

$$+\frac{\rho_2}{\rho_1} (U - i\omega/k)^2 k^2 \left(\frac{l^2 - k^2}{l^2 + k^2} \right) \frac{I_1(ka)K_0(ka)}{I_0(ka)K_1(ka)}$$

Maximum wave growth rate characterizes fastest growing waves which are responsible for breakup (as a function of Weber and Ohnesorge numbers)

Maximum wave growth rate and length scale: Ω and Λ







Curvefits of dispersion equation

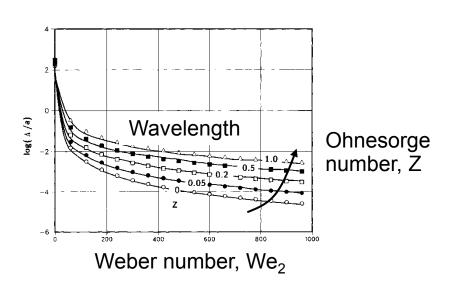
Reitz, 1988

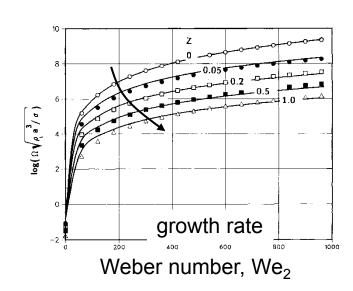
$$\frac{\Delta}{a} = 9.02 \frac{(1 + 0.45 Z^{0.5})(1 + 0.4 T^{0.7})}{(1 + 0.87 We_2^{1.67})^{0.6}}$$

$$\Omega \left(\frac{\rho_1 a^3}{\sigma} \right)^{0.5} = \frac{0.34 + 0.38 \ We_2^{1.5}}{(1+Z)(1+1.4T^{0.6})}$$

where
$$Z = \frac{We_1^{0.5}}{Re_1}$$
; $T = ZWe_2^{0.5}$; $We_1 = \frac{\rho_1 U^2 a}{\sigma}$; $We_2 = \frac{\rho_2 U^2 a}{\sigma}$; $Re_1 = \frac{Ua}{v_1}$

Maximum growth rate increases and wavelength decreases with We Increased viscosity reduces growth rate and increases wave length





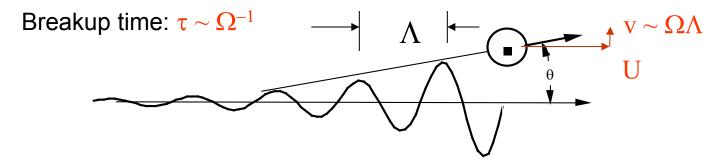




"Wave" atomization model

Reitz, 1988

Drop size: $r = B\Lambda$

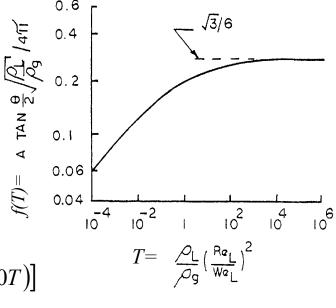


Spray angle prediction:

Tan
$$\theta = \frac{\mathbf{v}}{U} = \frac{1}{A} \frac{\rho}{4\pi (\frac{2}{\rho_1})^{1/2}} \mathbf{f}(T)$$

Breakup length of the core (Taylor, 1940):

$$L = C a \sqrt{\frac{\rho_1}{\rho_2}} / f(T) \quad \text{where} \quad f(T) = \frac{\sqrt{3}}{6} [1 - \exp(-10T)]$$







Drop breakup regimes

Breakup stages	Deformation or breakup regimes	Breakup process	Weber number	References	
First breakup stage	(1) Deformation and flattening	Air O	We < 12		
Second breakup stage	(b) Bag breakup	Air Bag growth Bag burst Rim burst	12 ≤ We ≤ 100 (including the Bag-and-Stamen breakup)	Pilch and Erdman	
	(c) Shear breakup	Air	We < 80	Ranger and Nicolls 1969	
	(d) Stretching and thinning breakup	Air Air Air	$100 \le \text{We} \le 350$	Liu and Reitz 1997	
	(e) Catastrophic breakup	Air Flattening and thinning waves KH waves	350 ≤ We	Hwang et al. 1996	



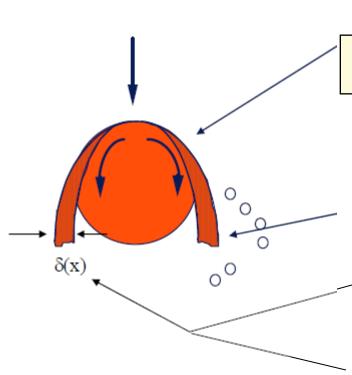
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High Speed Drop breakup

Liu, 1997

Mechanisms of drop breakup at high velocities poorly understood - Conflicting theories

Bag, 'Shear' and 'Catastrophic' breakup regimes



Breakup due to capillary surface waves

Hinze Chem Eng (1955) and Engel Nat. Bureau Stds (1958)

Boundary Layer Stripping due to Shear at the interface

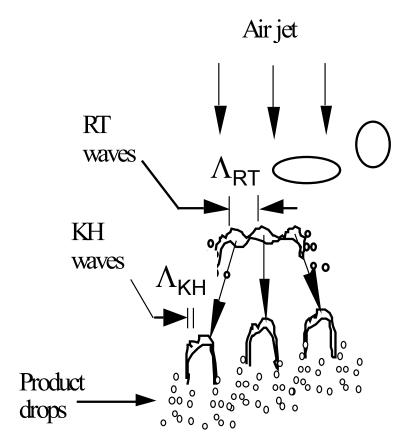
Ranger and Nicolls AIAA J. (1969)
Reinecke and Waldman AVCO Rep (1970)

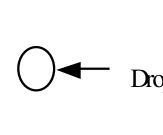
Delplanque & Sirignano Atom Sprays (1994)

Stretching and thinning – drop distortion - Liu and Reitz IJMF (1997)

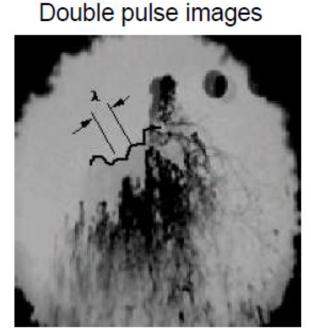


High speed drop breakup mechanisms





Rayleigh Taylor Breakup



$$\Omega_{RT} = \sqrt{\frac{2}{3\sqrt{\sigma}} \frac{\left[-g_t(\rho_l - \rho_g)\right]^{3/2}}{\rho_l + \rho_g}}$$

 g_t = acceleration

$$\Lambda_{RT} = \sqrt{\frac{-g_t(\rho_l - \rho_g)}{3\sigma}}$$

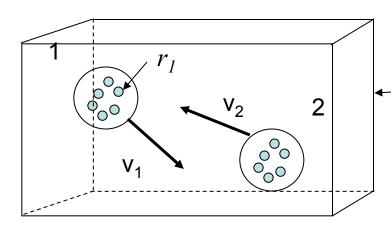


O'Rourke, 1981

Drop collision modeling

Collision frequency

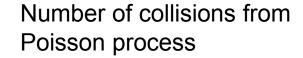
$$v_{12} = N_2 \pi (r_1 + r_2)^2 E_{12} |\mathbf{v}_1 - \mathbf{v}_2| / Vol$$



Collision efficiency

$$E_{12} = \left(\frac{K}{K + 1/2}\right)^2 \sim 1$$

$$K = \frac{2}{9} \frac{\rho_l |\mathbf{v}_1 - \mathbf{v}_2| r_2^2}{\mu_g r_1}$$



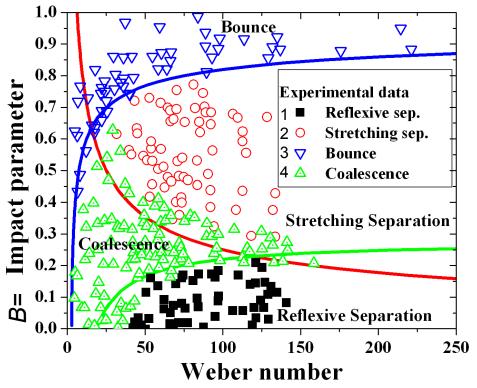
$$p(n) = e^{-v_{12}\Delta t} (v_{12}\Delta t)^n / n!$$

$$0 random number$$



Drop collision and coalescence

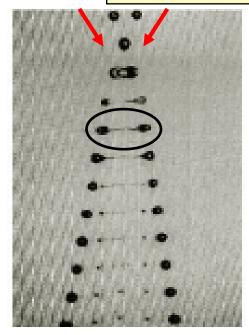
- 1. Reflexive vs. surface energy
- 2. Kinetic energy of unaffected part vs. surface energy
- 3. Drops cannot expel trapped gas film (bounce apart)
- 4. Drops form combined mass (coalesce)

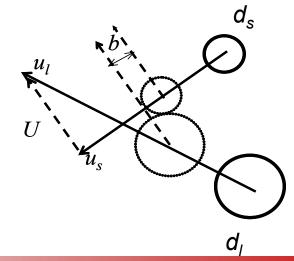


$$We = \frac{\rho_L U^2 d_S}{\sigma}$$
 , $B = \frac{2b}{(d_S + d_l)}$, $\Delta = \frac{d_S}{d_l}$



Munnannur, 2007 Georjon, 1999









Drop drag modeling

Liu et al. SAE 930072

Steady-state Stokes viscous drag, added-mass and Basset history integral

$$dv/dt = \mathbf{F} = 6 \pi r \mu_g \mathbf{v} + \frac{1}{2} \left(\frac{4}{3} \pi r^3 \rho_g\right) \frac{d\mathbf{v}}{dt} + 6r^2 \sqrt{\pi \mu \rho_g} \int_0^t \frac{d\mathbf{v}}{\sqrt{t - t'}} dt'$$
General form

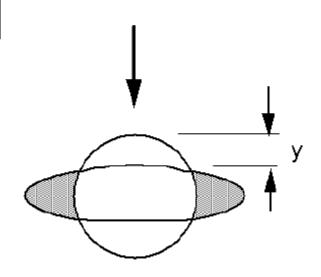
$$\rho_L V_d d\mathbf{v} / dt = C_D A_f \frac{\rho_g U^2}{2} \{ \mathbf{U} / |\mathbf{U}| \}$$

$$C_d = \begin{cases} 24 \operatorname{Re}_d^{-1} (1 + \operatorname{Re}_d^{2/3}/6), & \operatorname{Re}_d < 1000 \\ 0.424, & \operatorname{Re}_d \ge 1000 \end{cases}$$

Drop distortion (TAB model)

$$\ddot{y} = -5\frac{\mu_l}{\rho_l} \frac{\dot{y}}{r_d^2} - \frac{8\sigma y}{\rho_l r_d^3} + \frac{2}{3} \frac{\rho}{\rho_l} \frac{U_{rel}^2}{r_d^2}$$

$$C_d = C_{d,sphere}(1 + 2.632y)$$

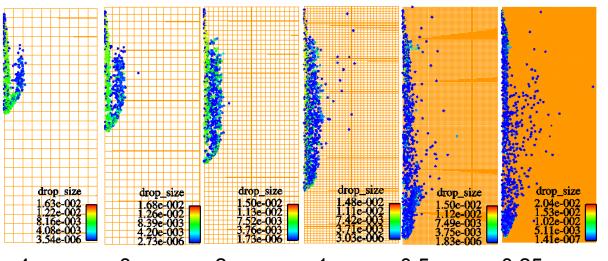






Grid independent spray models

Abani, 2008, Wang, 2013



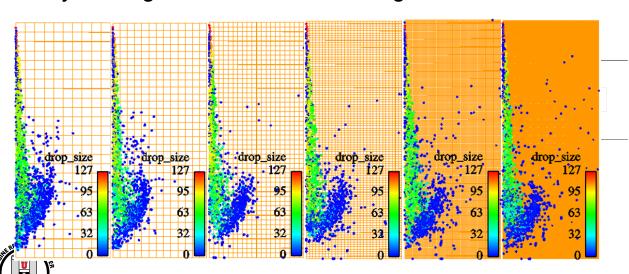
4 mm 3mm 2 mm 1 mm 0.5 mm 0.25n Gas-jet sub-grid momentum exchange near nozzle

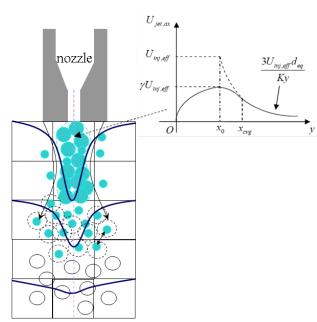
Coarse mesh:

Drop drag over-predicted

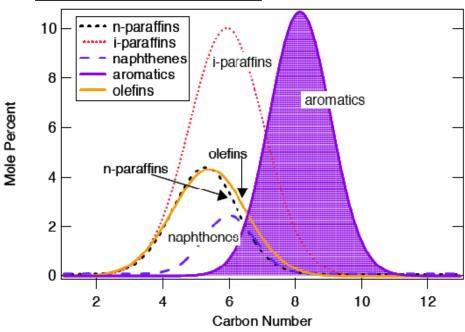
Fine mesh:

Drop coalescence underpredicted





Drop vaporization



No.	Physical property	No.	Physical property
p1	Liquid density	p7	Vapor heat capacity
p2	Vapor pressure	p8	Vapor diffusivity
рЗ	Surface tension	р9	Vapor viscosity
p4	Liquid viscosity	p10	Vapor thermal conductivity
p5	Liquid thermal conductivity	p11	Liquid Heat capacity
p6	Heat of vaporization	p12	Critical properties

- Physical properties used in sub-models
- Mixture properties calculated from pure components using appropriate equations

Discrete Multi-Component

Discrete system of liquid phase + discrete mixture system of vapor phase fuel and ambient gas:

$$G_p(I) = \sum_{F=1}^{N_F} x_F^{\ p} \delta(I - I_F) + \sum_{s=1}^{N_s} x_s^{\ p} \delta(I - I_s)$$
discrete phase of fuel

discrete phase of air/fuel mixture

Ra and Reitz, IJMF 2009

Vapor phase transport equation,

$$\frac{\partial}{\partial t} [\rho y_i] + \nabla \cdot [\rho y_i v] = \nabla \cdot (\rho D_i \nabla y_i) + s_{g,i}$$

$$\sum \Rightarrow \frac{\partial}{\partial t} [\rho y_F] + \nabla \cdot [\rho y_F v] = \nabla \cdot (\rho \overline{D} \nabla y_F) + S_g$$



177

190

Molecular weight [g/mol]

Ra, 2003



	diesel	Je	et-A	J	P-8	18 component DMC model						
cyC6	0.010000	0.02	22529	0.02	12529				-			
c7h8	0.037698	0.03	35873	0.02	25873							
c10h8	0.020000	0.00	09607	(0.0	alkan	es	650) 	⇒ diesel	model-diese	
c10h18	0.0	(0.0	0.03	30000		. •		4	∆ Jet-A	—model, Jet-A	
c10h22	0.113977	0.25	53797	0.72	21022	aromatics				JP-8	model, JP-8	<u> </u>
c12h26	0.124544	0.29	98704	0.13	18311	orvolo.	allran ag	600) +			
c14h10	0.0	(0.0	(0.0	cycloalkanes		<u> </u>				
c14h30	0.210805	0.20	02265	0.02	22265	PAH		⊌ 550	, [
c16h34	0.172593	0.06	60627	(0.0	1/11		E 330				
c18h38	0.085615	(0.0	(0.0]		erat		2		
c20h42	0.084268	(0.0	(0.0			<u>ē</u> 500)			<u> </u>
mxylene	0.010000	0.07	79442	(0.0]		g te				
mcymene	0.050355	(0.0	0.0				boiling temperature [K]				
c11h16	0.0	0.01	19881	9881								
tetralin	0.017362	0.01	17275		0.0			400	,			
c12h18	0.017362	(0.0		0.0			400	0	20	40 60	80 100
c13h20	0.045421	(0.0		0.0	7					n fraction [%]	00 100
nc7h16	0.0	(0.0	0.07	70000]						
				el	Mode		Jet-A	Mode		JP-8	Model-JP8	
	Density [g/cm^3]			'8	0.78		0.8101	0.7		0.7547	0.7333	← corrected
Viscosity [cSt]			2.71				1.55	1.3		1.15	1.01	
Surface tension [dynes/cm]			30 4252	26.91 6 42527.			29.1 43305		.53	25.2 44185	23.64 44199.4	
LHV [kJ/kg] Saturates [%]			4232	81.			84.4	43306.6 83.7		94.7	94.1	
Aromatics [%]					17.		14.0	14		1.2	2.2	
Olefins [%]							1.6	<u> </u>		4.1		
Naphthenes [%]					1.0	0		2.	.2		3.7	
C/H ratio			6.85	2	6.10	02	6.534	5.9	97	5.95	5.565	
						_ ı		1 4 4		1 4=4		

160

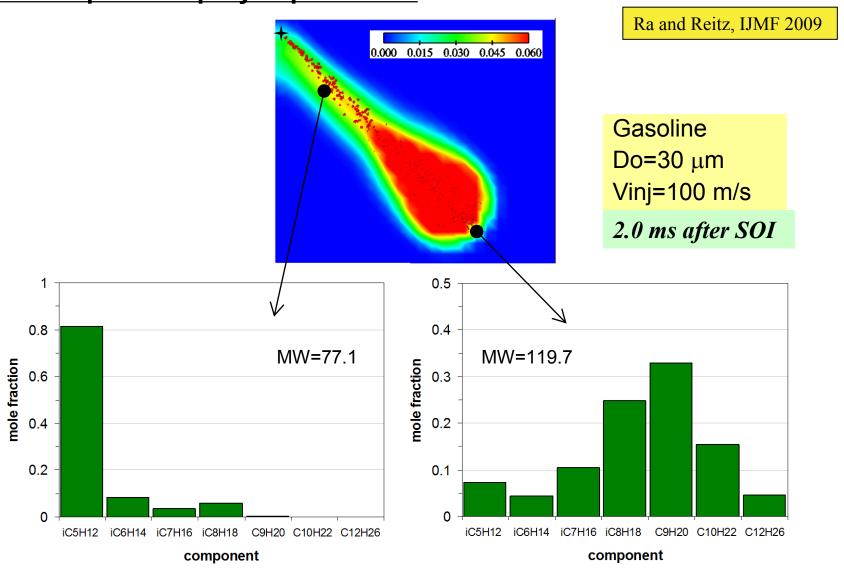
152

150

139 PCI-2-0, 2018



Multi-component spray vaporization





Summary

The Lagrangian Drop/Eulerian Fluid (LDEF) Discrete Drop model is the work-horse approach in commercial codes for simulating 2-phase flows.

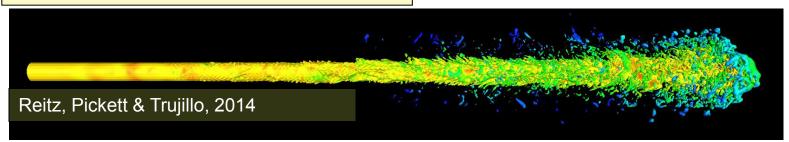
Detailed models are available for use in engine CFD models to describe the effects of injector nozzle flow, and liquid and gas properties on spray formation and drop breakup physics.

Significant progress is being made using LES/DNS spray modeling with high resolution experimental diagnostics to validate engine CFD spray models.

Ballistic imaging: Linne, 2009; X-Ray imaging: Liu SAE paper 2010-01-0877

LES: Villiers & Gosman, LES Primary Diesel Spray Atomization, SAE 2004-01-0100

DNS: Near field spray modeling (Trujillo - ERC)







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