Variation Method- Helium

If you guess a solution to the Schrödinger Equation, how do you find out how good a guess you have made? $\mathcal{H}\Psi_i = E_i \Psi_i$ can't be solved exactly:

$$\frac{-\,\hbar^2}{2m}\left(\nabla_{_1}^2+\nabla_{_2}^2\right)\,\Psi_i + \frac{1}{4\pi\epsilon_o}\!\left(\!-\frac{Ze^2}{r_1}\!-\!\frac{Ze^2}{r_2}\!+\!\frac{e^2}{r_{12}}\right)\Psi_i = E_i\Psi_i$$

but you have:

 ϕ_1 guess

 ϕ_2 better guess

 ϕ_3 better yet $E^{\phi} \ge E$

The Variation Theorem guarantees that the trial energy is always greater than the exact energy.

$$\frac{\hat{\mathcal{H}}\Psi_{i} = E_{i} \Psi_{i}}{\mathcal{H}}$$

$$\int \Psi_i^* \, \hat{\mathcal{H}} \Psi_i \, d\tau = \int \Psi_i^* \, E_i \, \Psi_i \, d\tau$$

solve for E_i : $E_i = \frac{\int \Psi_i^* \, \hat{\mathcal{H}} \Psi_i \, d\tau}{\int \mathbf{w}^* \, \mathbf{w} \, d\tau}$ is exact

$$E^{\phi} = \frac{\int \phi^* \mathcal{H} \phi \ d\tau}{\int \phi^* \phi \ d\tau}$$

 $E^{\phi} \geq E_i$ by the variation theorem

 $\varphi_{gs} = \varphi_{1s}(r_1) \; \varphi_{1s}(r_2) = \frac{1}{\pi} \left(\frac{Z_{eff}}{a_o} \right)^3 \; e^{-Z_{eff} \; r_1/a_o} \; e^{-Z_{eff} \; r_2/a_o} \qquad \qquad Z_{eff} : \text{variation parameter}$

 $E^{\phi} = \frac{\int \phi_{gs}^* \, \hat{\mathcal{H}} \phi_{gs} \, d\tau}{\int \phi_{ss}^* \, \phi_{ss} \, d\tau} = \frac{e^2}{4\pi\epsilon_0 \, a_0} \left(Z_{eff}^2 - \frac{27}{8} \, Z_{eff} \right)$ see Karplus and Porter 4.1.5

minimize by changing Z_{eff} : $\frac{dE^{\phi}}{dZ_{eff}} = 0 = \frac{e^2}{4\pi\epsilon_0 a_0} (2 Z_{eff} - 27/8)$

$$Z_{eff} = \frac{27}{16} = 1.6875$$

$$E_{gs} = -13.6 \text{ eV} (Z_{eff}^2 / n_1^2 + Z_{eff}^2 / n_2^2) = -77.5 \text{ eV}$$
 (exp. -79.0 eV)

 $\phi_{1s}(r_1) \text{ and } \phi_{1s}(r_2) \text{ are normalized eigenfunctions of one-electron Hamiltonians with charge Z_{eff}:}$

$$\left(\frac{-\,\hbar^2}{2m} \nabla_1^2 - \frac{Z_{eff}\,e^2}{4\pi\epsilon_o r_{_1}} \right) \varphi_{1s}(r_{_1}) = E_1 \,\, \varphi_{1s}(r_{_1}) \qquad \qquad \left(\frac{-\,\hbar^2}{2m} \nabla_2^2 - \frac{Z_{eff}\,e^2}{4\pi\epsilon_o r_{_2}} \right) \varphi_{1s}(r_{_2}) = E_2 \,\, \varphi_{1s}(r_{_2}) \,\, \varphi_{1s}(r_{_2}) = E_3 \,\, \varphi_{1s}(r_{_2}) \,\, \varphi_{1s}(r_{_2}) = E_3 \,\, \varphi_{1s}(r_{_2}) \,\, \varphi_{1s}(r_{_2}) \,\, \varphi_{1s}(r_{_2}) = E_3 \,\, \varphi_{1s}(r_{_2}) \,\, \varphi_{1s}(r_{_2}) \,\, \varphi_{1s}(r_{_2}) = E_3 \,\, \varphi_{1s}(r_{_2}) \,\, \varphi_{1s}(r_{$$

$$\begin{array}{ll} \text{normalization:} & \int \varphi_{1s}^*(r_1) \; \varphi_{1s}(r_1) \; d\tau_1 = 1 \\ \text{and eigenvalues:} \; E_1 = -1/2 & \frac{e^2}{4\pi\epsilon_o \; a_o} \\ Z_{eff}^2 = -13.6 \; eV \; Z_{eff}^2 \\ \end{array} \qquad \begin{array}{ll} & \int \varphi_{1s}^*(r_2) \; \varphi_{1s}(r_2) \; d\tau_2 = 1 \\ E_2 = -1/2 & \frac{e^2}{4\pi\epsilon_o \; a_o} \\ Z_{eff}^2 = -13.6 \; eV \; Z_{eff}^2 \\ \end{array}$$

The exact Hamiltonian uses the full nuclear charge Z=2. We can split the Coulomb portion of the exact Hamiltonian using $Z=Z_{eff}+(Z-Z_{eff})$

$$\begin{split} &\left(\frac{-\,\hbar^2}{2m}\nabla_1^2 - \frac{Ze^2}{4\pi\epsilon_0 r_1}\right)\varphi_{1s}(r_1) = \left(\frac{-\,\hbar^2}{2m}\nabla_1^2 - \frac{Z_{eff}\,e^2}{4\pi\epsilon_0 r_1}\right)\varphi_{1s}(r_1) - \frac{(Z-Z_{eff})\,e^2}{4\pi\epsilon_0 r_1}\,\varphi_{1s}(r_1) \\ &\left(\frac{-\,\hbar^2}{2m}\nabla_2^2 - \frac{Ze^2}{4\pi\epsilon_0 r_2}\right)\varphi_{1s}(r_2) = \left(\frac{-\,\hbar^2}{2m}\nabla_2^2 - \frac{Z_{eff}\,e^2}{4\pi\epsilon_0 r_2}\right)\varphi_{1s}(r_2) - \frac{(Z-Z_{eff})\,e^2}{4\pi\epsilon_0 r_2}\,\varphi_{1s}(r_2) \end{split}$$

Both electrons are in the same 1s-orbital so these two previous equations give the same energies. The electrons are identical, except for spin. To calculate the expectation value of the exact Hamiltonian, $\int \phi_{1s}^*(r_1) \; \phi_{1s}^*(r_2) \; \hat{\mathcal{H}} \phi_{1s}(r_1) \; \phi_{1s}(r_2) \; d\tau, \text{ we find } \hat{\mathcal{H}} \phi_{1s}(r_1) \; \phi_{1s}(r_2) \text{ using the last two equations:}$

$$\begin{split} \hat{\mathcal{H}}\,\varphi_{gs} &= \varphi_{1s}(r_2) \left(\frac{-\,\hbar^2}{2m} \nabla_1^2 - \frac{Ze^2}{4\pi\epsilon_o r_1} \right) \varphi_{1s}(r_1) \, + \varphi_{1s}(r_1) \left(\frac{-\,\hbar^2}{2m} \nabla_2^2 - \frac{Ze^2}{4\pi\epsilon_o r_2} \right) \varphi_{1s}(r_2) \, + \frac{e^2}{4\pi\epsilon_o r_{12}} \, \varphi_{1s}(r_1) \varphi_{1s}(r_2) \\ &= - \left(\frac{e^2}{4\pi\epsilon_o} a_o \right) Z_{eff}^2 \, \varphi_{1s}(r_1) \, \varphi_{1s}(r_2) \, - 2 \, \frac{(Z - Z_{eff}) \, e^2}{4\pi\epsilon_o r_1} \, \varphi_{1s}(r_1) \, \varphi_{1s}(r_2) \, + \frac{e^2}{4\pi\epsilon_o r_{12}} \, \varphi_{1s}(r_1) \varphi_{1s}(r_2) \end{split}$$

The variational energy for our guessed wave function is then:

$$E^{\phi} = \int \phi_{1s}^*(r_1) \; \phi_{1s}^*(r_2) \; \hat{\mathcal{H}} \; \phi_{1s}(r_1) \; \phi_{1s}(r_2) \; d\tau = -\left(\frac{e^2}{4\pi\epsilon_o \, a_o}\right) Z_{eff}^2 - 2 \left\langle \frac{(Z-Z_{eff}) \, e^2}{4\pi\epsilon_o r_1} \right\rangle \\ + \left\langle \frac{e^2}{4\pi\epsilon_o r_{12}} \right\rangle = -\left(\frac{e^2}{4\pi\epsilon_o r_{12}}\right) Z_{eff}^2 - 2 \left\langle \frac{(Z-Z_{eff}) \, e^2}{4\pi\epsilon_o r_{12}} \right\rangle \\ + \left\langle \frac{e^2}{4\pi\epsilon_o r_{12}} \right\rangle + \left\langle \frac{e^2}{4\pi\epsilon_o r_{12}} \right\rangle + \left\langle \frac{e^2}{4\pi\epsilon_o r_{12}} \right\rangle \\ + \left\langle \frac{e^2}{4\pi\epsilon_o r_{12}} \right\rangle + \left\langle \frac{e^2}{4\pi\epsilon_o r_{12}} \right\rangle \\ + \left\langle \frac{e^2}{4\pi\epsilon_o r_{12}} \right\rangle + \left\langle \frac{e^2}{4\pi\epsilon_o r_{12}} \right\rangle \\ + \left\langle \frac{e^2}{4\pi\epsilon_o r_{12}} \right\rangle + \left\langle \frac{e^2}{4\pi\epsilon_o r_{12}} \right\rangle \\ + \left\langle \frac{e^2}{4\pi\epsilon_o r_{12}} \right\rangle + \left\langle \frac{e^2}{4\pi\epsilon_o r_{12}} \right\rangle \\ + \left\langle \frac{e^2}{4\pi\epsilon_o r_{12}} \right\rangle + \left\langle \frac{e^2}{4\pi\epsilon_o r_{12}} \right\rangle \\ + \left\langle \frac{e^2}{4$$

The first expectation value is the difference in the Coulomb attraction of an electron with a nucleus of charge Z and a nucleus of charge Z_{eff} :

The first integral is the normalization for electron-2. The second integral is the expectation value of the potential energy for a one-electron atom with nuclear charge Z_{eff} , $\langle V \rangle$, see Problem 25.5.

$$-\left\langle \frac{(Z-Z_{\text{eff}})e^{2}}{4\pi\epsilon_{0}r_{1}}\right\rangle = -(Z-Z_{\text{eff}})\left(\frac{Z_{\text{eff}}e^{2}}{4\pi\epsilon_{0}a_{0}}\right)$$

We evaluated the expectation value of the electron-electron repulsion for our perturbation treatment of the helium atom. We only need to substitute Z_{eff} for Z:

$$\big\langle \frac{e^2}{4\pi\epsilon_o r_{12}} \big\rangle = \int \varphi_{1s}^*(r_1) \; \varphi_{1s}^*(r_2) \, \frac{e^2}{4\pi\epsilon_o r_{12}} \, \varphi_{1s}(r_1) \; \varphi_{1s}(r_2) \; d\tau_1 d\tau_2 = \frac{5Z_{eff}}{8} \bigg(\frac{e^2}{4\pi\epsilon_o \, a_o} \bigg)$$

Setting Z = 2 gives the final result:

$$E^{\phi} = \left(\frac{e^2}{4\pi\epsilon_o \, a_o}\right) (\, - \, Z_{eff}^2 - 2 \, Z Z_{eff} + 2 \, Z_{eff}^2 + \frac{5 Z_{eff}}{8} \,)$$

$$E^{\phi} = \frac{e^2}{4\pi\epsilon_o \, a_o} \left(Z_{eff}^2 - \frac{27}{8} \, Z_{eff} \right) = 27.211 \, \, eV \left(Z_{eff}^2 - \frac{27}{8} \, Z_{eff} \right)$$