## **Self-Consistent Field Hartree Fock Theory: Roothaan Equations**

LCAO approach:  $\Psi_a = \sum_{i=1}^{N} c_{ia} \phi_i$ 

 $\Psi_a$  = one-electron molecular orbital

 $\phi_i$  = atomic orbitals for N orbitals (Slater or Gaussian)

 $c_{ia}$  = orbital coefficients -- variational parameters

Molecule wavefunctions are Slater determinants of one-electron MO's

$$\Psi = 1/\!\sqrt{N!} \left| \begin{array}{ccccc} \Psi_a(1)\alpha(1) & \Psi_a(1)\beta(1) & \Psi_b(1)\alpha(1) & \Psi_b(1)\beta(1) & \dots \\ \Psi_a(2)\alpha(2) & \Psi_a(2)\beta(2) & \Psi_b(2)\alpha(2) & \Psi_b(2)\beta(2) & \dots \\ \Psi_a(3)\alpha(3) & \Psi_a(3)\beta(3) & \Psi_b(3)\alpha(3) & \Psi_b(3)\beta(3) & \dots \\ \vdots & \vdots & \vdots & \vdots & \dots \end{array} \right.$$

Spin wavefunction:  $\sigma(1) = \alpha(1)$  or  $\beta(1)$ 

**Variation Theorem: vary**  $c_{ia}$ 's to find the lowest energy:

Hartree-Fock Equations:

$$f_1 \, \Psi_a(1) \sigma(1) = \epsilon_a \, \Psi_a(1) \sigma(1)$$

 $\varepsilon_a$  = one-electron orbital energy

for MO a, b, c, ...

$$f_1 = h_1 + \sum_{j=1}^{n/2} \{2J_j(1) - K_j(1)\}$$
  $j = \text{all the other filled orbitals}$ 

for electron 1, 2, 3, ...

core Hamiltonian:

$$h_1 = \frac{-\overline{h}^2}{2m} \nabla_1^2 - \sum_{k=1}^m \frac{Z_k e^2}{4\pi \epsilon_0 r_{k_1}} \qquad k = \text{all nuclei}$$

electron-electron repulsion:

Coulomb operator:

$$J_{j}(1) \Psi_{a}(1) = \left( \int \Psi_{j}^{*}(2) \frac{e^{2}}{4\pi\epsilon_{0} r_{12}} \Psi_{j}(2) d\tau_{2} \right) \Psi_{a}(1)$$

Exchange operator: 
$$K_{j}(1) \Psi_{a}(1) = \left( \int \Psi_{j}^{*}(2) \frac{e^{2}}{4\pi\epsilon_{0} r_{12}} \Psi_{a}(2) d\tau_{2} \right) \Psi_{j}(1)$$

$$\begin{array}{|c|c|c|c|}\hline\\ \sum_{i=1}^{N} c_{ia} \int & \phi_{j}(1) \int_{1} & \phi_{i}(1) d\tau_{1} \\ \sum_{i=1}^{N} c_{ia} \int & \phi_{j}(1) & \phi_{i}(1) d\tau_{1} \\ \sum_{i=1}^{N} c_{ia} F_{ji} & = \epsilon_{a} \sum_{i=1}^{N} c_{ia} S_{ji} \\ F_{ji} & = \int & \phi_{j}(1) \int_{1} & \phi_{i}(1) d\tau_{1} \\ \hline\\ \end{array}$$

## Semi-Empirical Molecular Orbital Theory

- only valence electrons: the core electrons are included in the nuclear core.
- minimal basis set: basis set contains only occupied atomic orbitals (s, p and d), e.g., for carbon STO-3G only 2s and 2p functions.
- <u>parameterized integrals</u>: many integrals obtained by fitting the results to experimental data or *ab initio* calculations. Gives training set dependence.
- omission of some of the two electron integrals.

$$\int\!\! \varphi_i(1)\ \varphi_j(1)\, \frac{e^2}{4\pi\epsilon_0 r_{12}}\, \varphi_k(2)\ \varphi_l(2)\ d\tau_1\ d\tau_2)$$

 $\phi_i(1) = different atomic orbitals$ 

perhaps on different nuclei

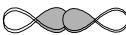
$$\int\!\! A(1)\; B(1)\; \frac{e^2}{4\pi\epsilon_0 r_{12}} \, C(2)\; D(2)\; d\tau_1 \, d\tau_2)$$

CNDO: complete neglect of differential overlap

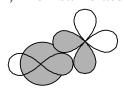
zero unless A=B

C=D

two-center integrals only



MNDO: modified neglect of differential overlap zero unless A, B on same atom C, D on same atom



neglect of diatomic differential overlap

AM1: more diverse training set; no d's for 3<sup>rd</sup> period or metals.

PM3: more heteroatoms in training set, use for 3<sup>rd</sup> period atoms and hydrogen bonding.

MNDO/d, PM3/TM: includes d orbitals for 3<sup>rd</sup> period elements, metals ZINDO: parameterized to reproduce electronic spectra

- If your molecule is *similar* to molecules in the training set, then the *results may be very good*.
- If your molecule is *significantly different* from anything in the training set, the *results may be very poor*.
- Semi-empirical methods are not as sensitive to the parameterization set as molecular mechanics methods.
- Much caution is warranted for transition states and unstable molecules and ions.