18.307: Integral Equations

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Solutions to Homework 6

(18.)
$$u(x) = \lambda \int_{-\infty}^{\infty} dy e^{-ixy} u(y)$$
, $-\infty < x < +\infty$.

(a) We write this equation symbolically as
$$u(x) = \lambda \mathcal{F} u(x) \qquad \text{where} \qquad \mathcal{F} u = \int_{-\infty}^{\infty} dy \ e^{-ixy} \ u(y)$$
integral operator

If follows that

$$\mathcal{F}^2 \mathbf{u}(\mathbf{x}) = \mathcal{F} \cdot \mathcal{F} \mathbf{u}(\mathbf{x}) = 2\pi \cdot \mathbf{u}(-\mathbf{x})$$

Then, $\mathcal{F}^{4}u(x) = \mathcal{F}^{2}\mathcal{F}^{2}u(x) = (2\pi)^{2}u(x)$ $\frac{1}{2}u(x) = \mathcal{F}^{4}u(x) \Rightarrow \mathcal{F}^{4}u(x) = \frac{1}{2}u(x)$ $\frac{1}{2}u(x) = \mathcal{F}^{4}u(x) \Rightarrow \mathcal{F}^{4}u(x) = \frac{1}{2}u(x)$ $u(x) \neq 0$

$$\Rightarrow \quad \lambda^{4} = \frac{1}{(2\pi)^{2}} \Rightarrow \quad \lambda = \frac{\pm 1}{\sqrt{2\pi}}, \quad \frac{\pm i}{\sqrt{2\pi}}$$

(b) Consider
$$u_n(x) = e^{-x^2/2} H_n(x)$$
, $H_n(x) = E_{11}^{\gamma} e^{x^2} \frac{d^{\gamma}}{dx^{\gamma}} e^{-x^2}$

$$\tilde{u}_{n}(k) = \int_{-\infty}^{\infty} dx \ e^{-ikx} \ e^{-x^{2}/2} H_{n}(x) = \int_{-\infty}^{\infty} dx \ e^{-ikx} \ e^{-x^{2}/2} G_{n}^{n} e^{-x^{2}} dx^{n} e^{-x^{2}/2}$$

$$\tilde{u}_{n}(k) = (-1)^{n} \int_{0}^{+\infty} dx = e^{ikx} e^{x^{2}/2} \frac{d^{n}}{dx^{n}} e^{-x^{2}}.$$

We move the differential operators to the right as follows.

$$e^{x^{2}/2} \frac{d^{n}}{dx^{n}} e^{-x^{2}} = e^{x^{2}/2} \frac{d^{n}}{dx^{n}} \left(e^{-x^{2}/2} e^{-x^{2}/2} \right) = e^{x^{2}/2} \frac{d^{n-1}}{dx^{n-1}} \frac{d}{dx} \left(e^{-x^{2}/2} e^{-x^{2}/2} \right)$$

$$= e^{x^{2}/2} \frac{d^{n-1}}{dx^{n-1}} e^{-x^{2}/2} \left(\frac{d}{dx} - x \right) e^{-x^{2}/2} = e^{x^{2}/2} \frac{d^{n-2}}{dx^{n-2}} \frac{d}{dx} e^{-x^{2}/2} \left(\frac{d}{dx} - x \right) e^{-x^{2}/2}$$

$$= e^{x^{2}/2} \frac{d^{n-2}}{dx^{n-2}} e^{-x^{2}/2} \left(\frac{d}{dx} - x \right) \left(\frac{d}{dx} - x \right) e^{-x^{2}/2}$$

$$= \cdots = e^{x^{2}/2} e^{-x^{2}/2} \left(\frac{d}{dx} - x \right) \left(\frac{d}{dx} - x \right) \cdots \left(\frac{d}{dx} - x \right) e^{-x^{2}/2}$$

$$= \left(\frac{d}{dx} - x \right)^{n} e^{-x^{2}/2}$$

$$= \left(\frac{d}{dx} - x \right)^{n} e^{-x^{2}/2}$$

by repeated use of the identity
$$\frac{d}{dx} e^{g(x)} f(x) = e^{g(x)} \left[\frac{d}{dx} + g'(x) \right] f(x)$$

It follows that
$$\widetilde{u}_n(k) = (-1)^n \int_{-\infty}^{\infty} dx \ e^{-ikx} \left(\frac{d}{dx} - x\right)^n \ e^{-x^2/2}$$

Integration by parts gives
$$\int_{-\infty}^{\infty} dx \ f(x) \left(\frac{d}{dx} - x\right) g(x) = f(x) g(x) \Big|_{x=-\infty}^{\infty}$$

$$+ \int_{-\infty}^{\infty} dx \ g(x) \left(-\frac{d}{dx} - x\right) f(x) = \int_{-\infty}^{\infty} dx \ g(x) \left(-\frac{d}{dx} - x\right) f(x),$$

assuming that
$$\lim_{|x|\to+\infty} [f(x)g(x)] = 0$$
.

Hence, we get
$$\tilde{u}_n(k) = (-1)^n \int_0^\infty dx \ e^{-x^2/2} \left(-\frac{d}{dx} - x \right) \left(-\frac{d}{dx} - x \right) \dots \left(-\frac{d}{dx} - x \right) e^{-ikx} ,$$

where
$$\left(-\frac{d}{dx} - x\right) e^{-ikx} = \left(ik - i\frac{d}{dk}\right) e^{-ikx} = i\left(k - \frac{d}{dk}\right) e^{-ikx}$$
.
We thus obtain

$$\tilde{u}_{n}(k) = (-1)^{n} i^{n} \int_{-\infty}^{+\infty} dx e^{-x^{2}/2} \left(k - \frac{d}{dk}\right) \left(k - \frac{d}{dk}\right) \dots \left(k - \frac{d}{dk}\right) e^{-ikx}$$

$$= (-i)^{n} \left(k - \frac{d}{dk}\right) \dots \left(k - \frac{d}{dk}\right) \int_{-\infty}^{\infty} dx e^{-x^{2}/2} e^{-ikx}$$

$$= i^{n} \left(\frac{d}{dk} - k\right) \dots \left(\frac{d}{dk} - k\right) \cdot \sqrt{2\pi} e^{k^{2}/2} = i^{n} \sqrt{2\pi} e^{k^{2}/2} \frac{d^{n}}{dk^{n}} e^{-k^{2}}$$

$$= (-i)^{n} \sqrt{2n^{i}} (-1)^{n} e^{k^{2}/2} \frac{d^{n}}{dk^{n}} e^{-k^{2}} = \sqrt{2\pi} (-i)^{n} u_{n}(k)$$

$$= \tilde{v}_{n}(k) = \sqrt{2n^{i}} (-i)^{n} u_{n}(k).$$

The given equation for
$$u(x) = u_n(x)$$
 gives

$$u_n(x) = \lambda \tilde{u}_n(x) = \lambda \sqrt{2\pi} (-i)^n u_n(x) \Rightarrow \lambda_n = \frac{i^n}{\sqrt{2\pi}}$$

For
$$n=0,1,2,3$$
, $\lambda=\frac{1}{\sqrt{2n}}$, $\frac{i}{\sqrt{2n}}$, $\frac{-i}{\sqrt{2n}}$, $\frac{-i}{\sqrt{2n}}$.

Clearly, the eigenvalues are infinitely degenerate.

The $u_n(x) = e^{-x^2/2}$ $H_n(x)$ are solutions of the eigenvalue problem of the quantum harmonic oscillator:

$$\left(-\frac{d^2}{dx^2}+x^2\right) u_n(x)=E_n u_n(x), \qquad u_n\left(|x|\to +\infty\right)=0.$$

Notice that the F.T. of both sides of this equation gives the same ode:

- (6) Suppose that u(x) is a solution of the given integral equation. There are four possibilities:
 - (i) u(x) corresponds to eigenvalue $\lambda_0 = \frac{1}{\sqrt{2\pi}}$, and hence is a linear combination of $u_n(x)$ with u=0,4,8,...=4m (m=0,1,...), with arbitrary coefficients.

$$u(x) = \sum_{n=0}^{\infty} C_n u_n(x) = \sum_{n=0}^{\infty} \frac{1}{2} c_n [1+(-i)^n] u_n(x), \text{ by extending } n$$

$$u(x) = \sum_{n=0}^{\infty} C_n u_n(x) = \sum_{n=0}^{\infty} \frac{1}{2} c_n [1+(-i)^n] u_n(x), \text{ by extending } n$$

arbitrary (admissible) wefficients C_2 , C_6 ,..., C_{2m} , C_{2m} .

[Since u(x) has to be square integrable, we need $\sum_{m=0}^{\infty} |C_{4m}|^2 < D0$;

similarly $\sum_{m=0}^{\infty} |C_{2m}|^2 < \infty$.

We thus get $u(x) = \sum_{\substack{n=0\\ n=2m}}^{\infty} \frac{1}{2} c_n u_n(x) + \sum_{\substack{n=0\\ n=2m}}^{\infty} \frac{c_n}{2} (-i)^n u_n(x)$

$$= \sum_{n:\text{even}} \frac{c_n}{2} u_n(x) + \sum_{n:\text{even}} \frac{c_n}{2\sqrt{2n}} \tilde{u}_n(x)$$

$$u(x) = \sum_{n:\text{even}} \frac{c_n}{2} u_n(x) + \frac{1}{\sqrt{2n'}} FT \left\{ \sum_{n:\text{even}} \frac{c_n}{2} u_n(x) \right\}$$

$$= f(x) + \frac{1}{\sqrt{2n'}} f(x),$$

where f(x): arbitrary <u>even</u> sq. integrable function (hence, $\tilde{f}(x)$ is also even and u(x): even)

Thus,
$$u(x) = f(x) + G f(x)$$
, $G = \frac{1}{\sqrt{2\pi}}$, $\lambda = \lambda_0 = \frac{1}{\sqrt{2\pi}}$.

(ii)
$$\lambda = \lambda_1 = \frac{i}{\sqrt{2\pi}}$$
, $n = 1, 5, 9, ... = 4m+1, m=0,1,...$

$$u(x) = \sum_{\substack{n=0\\n=4m+1}}^{\infty} c_n u_n(x) = \sum_{\substack{n=0\\n:odd}}^{\infty} \frac{c_n}{2} \left[| \bullet i(-i)^n \right] u_n(x) = \sum_{\substack{n:odd\\n:odd}}^{\infty} \frac{c_n}{2} u_n^{n+1} i \sum_{\substack{n:odd\\n:odd}}^{\infty} \frac{c_n}{2} (-i)^n u_n(x)$$

$$= \sum_{n:odd} \frac{c_n}{2} u_n(x) + \frac{i}{\sqrt{2\pi}} \sum_{n:odd} \frac{c_n}{2} \widetilde{u}_n(x) = f(x) + \frac{i}{\sqrt{2\pi}} \widetilde{f}(x)$$

where f(x) is any arbitrary odd square integrable function.

(iii)
$$\lambda = \lambda_2 = \frac{-1}{\sqrt{2\pi i}}$$
, $m = 2,6,10,... = 4m+2$, $m = 0,1,...$

$$u(x) = \sum_{n=0}^{\infty} c_n u_n(x) = \sum_{n=0}^{\infty} \frac{c_n}{2} \left[1 - (-i)^n \right] u_n(x) = \sum_{n: \text{even}} \frac{c_n}{2} u_n(x) - \sum_{n: \text{even}} \frac{c_n}{2} (-i)^n u_n(x)$$

$$n: \text{even}$$

=
$$f(x) - \frac{1}{\sqrt{2n}} \hat{f}(x)$$
, $f(x)$: even sq. integrable function

(iv)
$$\lambda = \lambda_3 = \frac{-i}{\sqrt{2\pi}}$$
, $n = 3, 7, 11, ... = 4m+3, m=0,1,2,...$

$$u(x) = \sum_{\substack{n=0\\ n=4m+3}}^{\infty} c_n u_n(x) = \sum_{\substack{n=0\\ n:odd}}^{\infty} \frac{c_n}{2} \left[1 - i (-i)^n \right] u_n(x) = f(x) - \frac{i}{\sqrt{2n'}} f(x),$$

f(x): odd sq. integrable function.

In stanmary,
$$u(x) = f(x) + C \cdot \tilde{f}(x)$$
 where $C = \begin{cases} \frac{1}{\sqrt{2n}}, & \text{for } \lambda = \frac{1}{\sqrt{2n}} \\ \frac{\pm i}{\sqrt{2n}}, & \lambda = \frac{\pm i}{\sqrt{2n}} \end{cases}$

(d) Let
$$f(x) = e^{-ax^2/2}$$
 (f: even).

The F.T. of
$$f(x)$$
 is $\tilde{f}(k) = \int_{-\infty}^{+\infty} dx e^{-ikx} e^{-ax^2/2} \frac{\tilde{a}x = y}{\sqrt{a}} \int_{-\infty}^{+\infty} dy e^{-i\sqrt{a}y} e^{-y^2/2}$

$$= \frac{1}{\sqrt{a}} \cdot \sqrt{2\pi} e^{-\left(\frac{k}{\sqrt{a}}\right)^2 \cdot \frac{1}{2}} = \sqrt{\frac{2\pi}{a}} e^{-k^2/2a}$$

The desired solution u(x) is

$$u(x) = f(x) \pm \frac{1}{\sqrt{2n}} \quad \tilde{f}(x) = e^{-ax^2/2} \pm \frac{1}{\sqrt{2n}} \quad \sqrt{\frac{2n}{a}} \quad e^{-x^2/2a}$$

$$= e^{-ax^2/2} \pm \frac{1}{\sqrt{a}} \quad e^{-x^2/2a}$$
which satisfies the given equation with eigenvalue $\lambda = \frac{\pm 1}{\sqrt{2n}} \quad (n=0,2)$.

Note that for $\underline{a=1}$, one of the two solutions becomes $\underline{=0}$, while the other one gives the Gaussian as solution, with (sole) eigenvalue $\lambda = \frac{1}{\sqrt{2}n}$.

In the following problem (Prob. 19) the kernel has a discontinuity across the line y=x. For such kernels the eigenfunctions Un(x) are continuous, i.e., such discontinuities in the kernel do not affect similarly the analytic properties of their eigenfunctions.

In contrast, if the kernel had a discontinuities at fixed lines, say $x=x_0$ or $y=y_0$, then the corresponding eigenfunctions would be discontinuous at fixed points (see example that I gave in class: K(x,y)=1, $0 \le x,y \le \frac{1}{2}$ and K(x,y)=2 otherwise, where $0 \le x,y \le 1$ in the entire region of definition.)

(19)
$$u(x) = f(x) + \lambda \int_{0}^{1} dy \ K(x,y) \ u(y)$$
, where $K(x,y) = \begin{cases} 3, \ 0 \le y < x \le 1, \\ 2, \ 0 \le x < y \le 1. \end{cases}$

(a) Homogeneous equation:

$$u(x) = \lambda \int_{0}^{x} dy \ 3 u(y) + \lambda \int_{x}^{1} dy \ 2 u(y) = 2\lambda \int_{0}^{1} dy u(y) + \lambda \int_{0}^{x} dy u(y)$$

$$\Rightarrow u'(x) = \lambda u(x) \Rightarrow u(x) = C e^{\lambda x}, 0 \le x \le 1., C \ne 0.$$

From the original equation we get $\begin{cases} u(0) = 2\lambda \int dy \ u(y) \\ u(1) = 3\lambda \int dy \ u(y) \end{cases} = u(1) = \frac{3}{2} u(0).$

Hence,
$$C_1 e^{\lambda} = \frac{3}{2} C_1 \implies e^{\lambda} = \frac{3}{2} = e^{\ln(3/2) + i2n\pi}$$
, $n: integer$

$$\Rightarrow \lambda = \lambda_n = \ln(\frac{3}{2}) + i2n\pi \quad , \quad n = 0, \pm 1, \pm 2, \dots$$

The corresponding eigenfunctions are $u_n(x) = G_n e^{\lambda_n x}$. Take G_n : real.

Finally:
$$u_n(x) = G_n e^{\left[\ln\left(\frac{3}{2}\right) + i2n\pi\right]x}$$
, n: integer.

(b) Clearly,
$$K(x,y) \neq K(y,x)$$
: K is not symmetric.

$$K^{T}(x,g) = K(y,x) = \begin{cases} 3, & 0 \le x < y \le 1 \\ 2, & 0 \le y < x \le 1 \end{cases}$$

Homogeneous equation for
$$K^T$$
: $\mathbf{v}(x) = \lambda \int dy K^T(x,y) V(y)$

$$\Rightarrow$$
 $V'(x) = -\lambda V(x) \Rightarrow V(x) = Fe^{-\lambda x}$, $0 \le x \le 1$, $F: const. \ne 0$. Take $F: real$.

From original equation,
$$V(0) = 3\lambda \int_{0}^{1} dy \ v(y) = \sqrt{1} =$$

$$\Rightarrow Fe^{-\lambda} = \frac{2}{3}F \Rightarrow \lambda = \lambda_n = \ln\left(\frac{3}{2}\right) + i2n\pi, \quad n: \text{ integer} \quad \text{(same as for K)}.$$

Eigenfunctions:
$$V_m(x) = F_m e^{-[ln(3/2) + i2m\pi]x}$$
, m: integer.

(d)
$$K(x,y) = \sum_{N=-\infty}^{\infty} \frac{u_n(x) v_n(y)}{\lambda_n} = \sum_{N=-\infty}^{\infty} \frac{e^{i2n(x-y)\pi + \ln(3/2) \cdot (x-y)}}{\ln(\frac{3}{2}) + i2n\pi}$$

by taking
$$C_n = 1 = F_n$$
 so that $\int dx \, U_n(x) \, V_m(x) = \delta_{nm}$.

In conclusion,
$$K(x,y) = \sum_{n=-\infty}^{\infty} \frac{\exp\{[\ln(3/2) + i 2n\pi](x-y)\}}{\ln(3/2) + i 2n\pi}$$
.

20 In order to make more clear the connection to integral equations, allow me to take $[E=\lambda]$. Suppose that $\lambda=\lambda n$ is any eigenvalue of $Au=\lambda u$, with using eigenfunction.

$$Au(x) = \lambda u(x) \implies \int_{0}^{1} dx \ u^{*}(x) A u(x) = \lambda \int_{0}^{1} dx \ |u(x)|^{2}$$

$$\implies \int_{0}^{1} dx \ u^{*}(x) \left\{ -\frac{d^{2}u}{dx^{2}} + \int_{0}^{1} dy \ xy \ u(y) \right\} = \lambda \int_{0}^{1} dx \ |u(x)|^{2}$$

$$\implies \int_{0}^{1} dx \ u^{*}(x) \left(-\frac{d^{2}u}{dx^{2}} \right) + \int_{0}^{1} dx \int_{0}^{1} dy \ u^{*}(x) \ u(y) \cdot xy \implies \lambda \int_{0}^{1} dx \cdot |u(x)|^{2}$$

$$\implies -u^{*}(x) \frac{du}{dx} \Big|_{0}^{1} + \int_{0}^{1} dx \cdot \left| \frac{du}{dx} \right|^{2} + \left| \int_{0}^{1} dx \ u(x) \cdot x \right|^{2} = \lambda \int_{0}^{1} dx \cdot |u(x)|^{2}$$

$$= 0, \text{ be cause}$$

$$u^{*}(x) = 0 = u'(x)$$

$$\Rightarrow \lambda = \left[\int_{0}^{1} dx \ |u(x)|^{2} \right]^{-1} \left\{ \int_{0}^{1} dx \ |u(x)|^{2} \right\}$$

It is inferred that 1= 2>0.

Suppose In, Im are two different eigenvalues (n≠m), with eigenfunctions $u_n(x)$ and $u_m(x)$.

$$Au_{m}(x) = \lambda_{n} u_{n}(x) \implies u_{m}^{*}(x) A u_{n}(x) = \lambda_{n} u_{m}^{*}(x) u_{n}(x)$$

$$Au_{m}(x) = \lambda_{m} u_{m}(x) \implies Au_{m}^{*}(x) = \lambda_{m} u_{m}^{*}(x) \implies u_{n}(x) A u_{m}^{*}(x) = \lambda_{m} u_{m}^{*}(x) u_{n}(x)$$

$$\implies \int_{0}^{1} dx \left[u_{n}^{*}(x) A u_{n}(x) - u_{n}(x) A u_{m}^{*}(x) \right] = (\lambda_{n} - \lambda_{m}) \int_{0}^{1} dx u_{m}^{*}(x) u_{n}(x)$$

$$\implies \int_{0}^{1} dx \left[u_{n}^{*}(x) A u_{n}(x) - u_{n}(x) A u_{m}^{*}(x) \right] = (\lambda_{n} - \lambda_{m}) \int_{0}^{1} dx u_{m}^{*}(x) u_{n}(x)$$

$$\implies \int_{0}^{1} dx \left[u_{n}^{*}(x) A u_{n}(x) - u_{n}(x) A u_{m}^{*}(x) \right] = (\lambda_{n} - \lambda_{m}) \int_{0}^{1} dx u_{m}^{*}(x) u_{n}(x)$$

$$\implies \int_{0}^{1} dx \left[u_{n}^{*}(x) A u_{n}(x) - u_{n}(x) A u_{n}(x) \right] = \int_{0}^{1} dx u_{n}^{*}(x) \left[-\frac{d^{2}u_{n}}{dx^{2}} + \int_{0}^{1} dy u_{n}(y) \right]$$

$$= \int_{0}^{1} dx \frac{du_{m}^{*}}{dx} \frac{du_{n}^{*}}{dx} + \int_{0}^{1} dx x u_{m}^{*}(x) \cdot \int_{0}^{1} dy u_{n}(y).$$

Hence,
$$\int dx \left[u_{m}^{\dagger}(x) A u_{n}(x) - u_{n}(x) A u_{m}^{\dagger}(x)\right]$$

$$= \int dx \frac{du_{m}^{\dagger}}{dx} \frac{du_{m}}{dx} + \int dx \times u_{m}^{\dagger}(x) \cdot \int dy \ y u_{n}(y)$$

$$- \int dx \frac{du_{n}}{dx} \cdot \frac{du_{m}^{\dagger}}{dx} - \int dx \times u_{n}(x) \cdot \int dy \ y u_{m}^{\dagger}(y) = 0.$$

$$= \partial \left(\lambda_{n} - \lambda_{m}\right) \int u_{m}^{\dagger}(x) u_{n}(x) dx = 0 \qquad \lim_{\lambda_{n} \neq \lambda_{m}} \int u_{m}^{\dagger}(x) u_{n}(x) dx = 0.$$

Since A is real $(A^* = A)$,

$$Au = \lambda u \qquad \begin{cases} A(u+u^*) = \lambda(u+u^*) \\ A(u-u^*) = \lambda(u-u^*) \end{cases}$$
 and
$$Au^* = \lambda u^* \qquad \begin{cases} A(u-u^*) = \lambda(u-u^*) \\ A(u-u^*) = \lambda(u-u^*) \end{cases}$$

i.e., if u(x) is an eigenfunction, so is $u(x) \pm u^*(x)$, with the same eigenvalue. Because $u(x) + u^*(x)$ is real and $-i [u(x) - u^*(x)]$ is real, we can always take the eigenfunctions to be real without any loss of generality.

(b) This is a tricky question (my original solution was wrong, since I kept finding a non-symmetric kernel that depends on E!). I will use the Green's function method to find a 1-independent, symmetric kernel.

Define
$$G(x,x')$$
 so that
$$\begin{cases} -G_{xx}(x,x') = S(x-x'), & 0 < x, x' < 1. \\ G(0,x') = 0 = G_{x}(1,x') \end{cases}$$

It follows that
$$G(x,x') = \begin{cases} A_1x + B_1, & 0 < x < x' < 1 \\ A_2x + B_2, & 0 < x' < x < 1 \end{cases}$$

$$G(0, k') = 0$$
 \Rightarrow $B_1 = 0$
 $G_{k}(1, k') = 0$ \Rightarrow $A_{k} = 0$

Two further conditions on
$$G(x,x')$$
 are:
$$\begin{cases} G(x=x'^{\dagger},x') = G(x=x'^{\dagger},x') \\ G_{x}(x=x'^{\dagger},x') = G_{x}(x=x'^{\dagger},x') = 1 \end{cases}$$

$$\Rightarrow \begin{cases} B_2 = A_1 x' \\ A_1 = 1 \end{cases} \Rightarrow \begin{cases} A_1 = 1 \\ B_2 = x' \end{cases} \qquad G(x, x') = \begin{cases} x, & 0 < x < x' < 1 \\ x', & 0 < x' < x < 1 \end{cases} = \min\{x, x'\}$$

The equation for u(x) is cast in the following form "inhomogeneous term"

$$\begin{cases} -u''(x) = g(x) - bx \\ u(0) = 0 = u'(1) \end{cases}$$
, where $g(x) = \lambda u(x)$, $b = \int dy \ y u(y) = const.$

By the Green's function method, it follows that

$$u(x) = \int dx' G(x,x') - [g(x') - bx']$$

$$= \int dx' G(x,x')g(x') - b \int dx' x' G(x,x')$$

We use this equation to determine to in terms of integrals of G:

$$b = \int dy \quad yu(y) = \int dy \cdot y \left[\int dx' \quad G(y,x') \quad g(x') - b \int dx' \quad x' \quad G(y,x') \right]$$

$$= \int dy \int dx' \quad y \cdot g(x') \quad G(y,x') - b \int dy \int dx' \quad x'y \quad G(y,x')$$

$$= \int dy \int dx' \quad y \cdot g(x') \quad G(y,x') = \frac{\int dy \int dx' \quad x'y \quad G(y,x')}{1 + \int dy \int dx' \quad x'y \quad G(y,x')} = \frac{B}{1 + A},$$

where
$$A = \int_{0}^{1} dy \int_{0}^{1} dx' \quad x'y \quad G(ty,x') = 2 \int_{0}^{1} dy \int_{0}^{y} dx' \quad x'y \quad min_{1}\{y,x'\} = 2 \int_{0}^{1} dy \int_{0}^{y} dx' \quad x'^{2}y$$

$$= 2 \int_{0}^{1} dy \cdot y \quad y^{3} = \frac{2}{3} \cdot \frac{1}{5} = \frac{2}{15} \implies A = \frac{2}{15},$$

$$B = \int_{0}^{1} dy \int_{0}^{1} dx' \quad yg(x') \quad G(y,x') = \int_{0}^{1} dy \int_{0}^{1} dx' \quad g(x') \quad yG(y,x') = \int_{0}^{1} dx' \quad g(x') \left[\int_{0}^{1} dy \quad yG(y,x')\right]$$

$$= \int_{0}^{1} dx' \quad g(x') \left[\int_{0}^{1} dy \cdot y^{2} + \int_{x}^{1} dy \cdot yx'\right] = \int_{0}^{1} dx' \quad g(x') \left[\frac{x^{3}}{3} + x' \cdot \frac{1 - x'^{2}}{2}\right] = \int_{0}^{1} dx' \quad g(x') \left(\frac{x'}{2} - \frac{x'^{3}}{6}\right)$$

$$Hence, \quad b = \frac{15}{17} \int_{0}^{1} dx' \quad g(x') \quad \left(\frac{x'}{2} - \frac{x'^{3}}{6}\right) = \frac{15}{17} \int_{0}^{1} dy \quad g(y) \quad \left(\frac{y}{2} - \frac{y^{3}}{6}\right)$$

$$Accordingly, \quad u(x) = \int_{0}^{1} dx' \quad G(x,x') \quad g(x') - b \int_{0}^{1} dx' \cdot x' G(x,x')$$

$$= \int_{0}^{1} dy \quad G(x,y) \quad g(y) - \int_{0}^{1} \frac{15}{17} \left(\frac{x}{2} - \frac{x^{3}}{6}\right) \left(\frac{y}{2} - \frac{y^{3}}{6}\right) \quad J(y) \quad J(y)$$

$$= \int_{0}^{1} dy \quad \left[\min_{1} \{x,y\} - \frac{15}{17} \left(\frac{x}{2} - \frac{x^{3}}{6}\right) \left(\frac{y}{2} - \frac{y^{3}}{6}\right) \right] \quad u(y)$$

Notice that the kennel K of this equation is symmetric! This fact is consistent with Au= hu having real eigenvalues.

(c) The original equation,
$$Au = \lambda u$$
, reads as $u''(x) + \lambda u(x) = bx$, $b = \int dy yu(y)$.

It follows that

$$u(x) = A \cos(\sqrt{\lambda}x) + B \sin(\sqrt{\lambda}x) + \frac{b}{\lambda}x \qquad (\lambda \neq 0).$$

.
$$u(0) = 0 \implies A = 0$$
; $u(x) = B \sin(\sqrt{\lambda} x) + \frac{b}{\lambda} x$

•
$$u'(1) = 0 = 0$$
 $B\sqrt{\lambda} \cos(\sqrt{\lambda}) + \frac{b}{\lambda} = 0 = 0$ $b = -B\lambda^{3/2} \cos(\sqrt{\lambda})$

$$b = \int dy \ y u(y) = \int dy \ y \left[B \sin(\sqrt{\lambda}y) + \frac{b}{\lambda}y \right] = -\frac{B}{\sqrt{\lambda}} \ y \cos(\sqrt{\lambda}y) \Big|_{\sqrt{\lambda}} + \frac{B}{\lambda} \int dy \cos(\sqrt{\lambda}y) \Big|_{\sqrt{\lambda}} + \frac{B}{\lambda} \Big|_{\sqrt{\lambda}} + \frac{B}{\lambda} \Big|_{\sqrt{\lambda}} + \frac{B}$$

$$\Rightarrow -B\lambda^{3/2}\cos(\sqrt{\lambda})\cdot\left(1-\frac{1}{3\lambda}\right)=-\frac{B}{\sqrt{\lambda}}\cos(\sqrt{\lambda})\cdot\left[1-\frac{1}{\sqrt{\lambda}}\tan(\sqrt{\lambda})\right]$$

Assume cos (IX) \$0 and B\$0:

$$\lambda^{3/2} \left(1 - \frac{1}{3\lambda} \right) = \frac{1}{\sqrt{\lambda}} \left[1 - \frac{1}{\sqrt{\lambda}} \tan(\sqrt{\lambda}) \right]$$

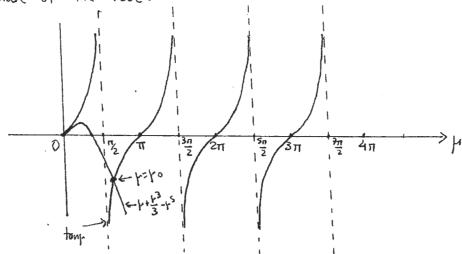
$$\Rightarrow \lambda^{5/2} \left(1 - \frac{1}{3\lambda} \right) = \sqrt{\lambda} - \tan(\sqrt{\lambda}).$$
 Set $\sqrt{\lambda} = \gamma \rightarrow \gamma = \lambda^2$:

$$\mu^{5}\left(1-\frac{1}{3\mu^{2}}\right)=\mu-\tanh\iff \tanh=\mu+\frac{\mu^{3}}{3}-\mu^{5}$$
, $\mu>0$.

For
$$B=0$$
, one gets $b=0$, and $u(x)=0$ (trivial).

For
$$\omega s(\sqrt{\lambda}) = 0$$
, one gets $b = 0$ and $B = 0$. Hence, $B \neq 0 \neq \omega s(\sqrt{\lambda})$

(d) My sketch is pretty awful and I hope that students will make a good estimate of the root!



The polynomial $P(p) = p + \frac{\mu^3}{3} - p^5$ has externa at

$$P'(\mu)=0 \Rightarrow 1+\mu^2-5\mu^4=0 \Rightarrow \mu=\frac{1+\sqrt{21}}{10} \Rightarrow \mu=\pm\sqrt{\frac{1+\sqrt{21}}{10}} \simeq \pm 0.7$$

At f=f+, P(+) has a maximum, and then it decreases rapidly to -a as f++0.

Note that P(0)=0, P'(0)=1 and $\frac{d}{dr}(tant)|_{r=0}=1$, while

$$P'(0) = 0$$
 and $\frac{d^2}{dt^2} \left(\frac{d}{dt} \right) \Big|_{t=0} = 0$. More precisely, tank $\sim \mu + \frac{\mu^3}{3} + \frac{\mu^5}{8}$,

which means that P(4) and tank do not intersect in (0, 1/2).

I therefore expect $\rho = \rho_0$ to

An estimate for $\lambda_0 = \rho_0^2$ is $\lambda = \frac{\pi^2}{4} \approx 2.5$ (e) $E_0 = \lambda_0 = \min_{v \in D_A} \frac{\int_0^1 dx \left[v'(x) \right]^2 + \left[\int_0^1 dx \, x \, v(x) \right]^2}{\int_0^1 dx \, v(x)^2}$

For the trial function v(x)=x(c-x) to belong to D_A , we have to choose c=2. Then v(x)=x(2-x), v(0)=0, $v'(x)=2-2x \Rightarrow v'(1)=0$.

$$\int_{0}^{1} dx \quad v(x)^{2} = \int_{0}^{1} dx \quad x^{2}(2-x)^{2} = \int_{0}^{1} dx \quad x^{2}(4+x^{2}-4x) = 4\frac{1}{3} + \frac{1}{5} - 4 \cdot \frac{1}{4}$$

$$= \frac{4}{3} + \frac{1}{5} - 1 = \frac{20+3-15}{15} = \frac{8}{15},$$

 $\int_{0}^{1} dx \left[v'(x) \right]^{2} = \int_{0}^{1} dx \ 4 \left(1-x \right)^{2} = 4 \int_{0}^{1} dx \left(1+x^{2}-2x \right) = 4 \left(1+\frac{1}{3}-2\cdot\frac{1}{2} \right) = \frac{4}{3},$ $\int_{0}^{1} dx \ x v(x) = \int_{0}^{1} dx \ x^{2}(2-x) = 2\frac{1}{3} - \frac{1}{4} = \frac{2}{3} - \frac{1}{4} = \frac{8-3}{12} = \frac{5}{12}.$

$$\lambda_0$$
 $\frac{4/3 + \frac{25}{144}}{8/15} = \frac{\frac{192 + 25}{144}}{\frac{8}{15}} = \frac{\frac{5}{15 \cdot 217}}{8.147} = \frac{1085}{384} \approx 2.82$

According to the trace inequality,

$$\lambda_0 \gg \left[\int_0^1 dx \ K_2(x,x)\right]^{-\nu_2}$$

where $K_2(x,y) = \int_0^1 d\xi \ K(x,\xi) \ K(\xi,y) = \int_0^1 d\xi \ K(x,\xi) \ K(y,\xi)$ $\int_0^1 dx \ K_2(x,x) = \int_0^1 dx \int_0^1 d\xi \ K(x,\xi)^2 = ||K||^2,$

and $K(x,y) = \min\{x,y\} - \frac{15}{17} \left(\frac{x}{2} - \frac{x^3}{6}\right) \left(\frac{y}{2} - \frac{y^3}{6}\right)$ from p. 1).

$$\|K\|^{2} = \int_{0}^{1} dx \int_{0}^{1} dy |K(x,y)|^{2} = 2 \int_{0}^{1} dx \int_{0}^{1} dy \cdot \left[y - \frac{15}{17} \left(\frac{x}{2} - \frac{x^{3}}{6} \right) \left(\frac{y}{2} - \frac{y^{3}}{6} \right) \right]^{2}$$

$$=2\int_{0}^{1}dx\int_{0}^{x}dy\left[y^{2}+\frac{15^{2}}{17^{2}}\left(\frac{x}{2}-\frac{x^{3}}{6}\right)^{2}\left(\frac{y}{2}-\frac{y^{3}}{6}\right)^{2}-\frac{30}{17}\left(\frac{x}{2}-\frac{x^{3}}{6}\right)\left(\frac{y^{2}}{2}-\frac{y^{4}}{6}\right)\right]$$

$$= 2 \int_{0}^{1} dx \frac{x^{3}}{3} + 2 \cdot \frac{15^{2}}{17^{2}} \int_{0}^{1} dx \left(\frac{x}{2} - \frac{x^{3}}{8} \right)^{2} \int_{0}^{x} dy \left(\frac{y}{2} - \frac{y^{3}}{6} \right)^{2} - \frac{60}{17} \int_{0}^{1} dx \left(\frac{x}{2} - \frac{x^{3}}{6} \right) \int_{0}^{x} dy \left(\frac{y^{2}}{2} - \frac{y^{3}}{6} \right)^{2} dy$$

$$= 2 \int_{3}^{1} \frac{1}{4} + 2 \cdot \frac{15^{2}}{17^{2}} \int_{0}^{1} dx \cdot \left(\frac{x}{2} - \frac{x^{3}}{6} \right)^{2} \cdot \left(\frac{x^{3}}{12} + \frac{x^{4}}{36 \cdot 7} - \frac{x^{5}}{30} \right) - \frac{60}{17} \int_{0}^{1} dx \left(\frac{x}{2} - \frac{x^{3}}{6} \right) \left(\frac{x^{3}}{6} - \frac{x^{5}}{30} \right) dy$$

$$= \frac{1}{6} + 2 \cdot \frac{15^{2}}{17^{2}} \int_{0}^{1} dx \left(\frac{x^{4}}{4} + \frac{x^{6}}{36} - \frac{x^{4}}{6} \right) \cdot \left(\frac{x^{3}}{12} + \frac{x^{7}}{36 \cdot 7} - \frac{x^{5}}{30} \right) - \frac{60}{17} \int_{0}^{1} dx \left(\frac{x^{4}}{12} - \frac{x^{6}}{60} - \frac{x^{6}}{36} + \frac{x^{8}}{180} \right)$$

$$= \frac{1}{6} + 2 \cdot \frac{15^{2}}{17^{2}} \int_{0}^{1} dx \left(\frac{x^{5}}{48} + \frac{x^{9}}{4 \cdot 36 \cdot 7} - \frac{x^{7}}{120} + \frac{x^{9}}{36 \cdot 12} + \frac{x^{13}}{36 \cdot 7} - \frac{x^{11}}{36 \cdot 30} - \frac{x^{7}}{6 \cdot 12} - \frac{x^{11}}{36 \cdot 6 \cdot 7} + \frac{x^{9}}{6 \cdot 30} \right)$$

$$- \frac{60}{17} \int_{0}^{1} dx \left(\frac{x^{4}}{12} - \frac{x^{6}}{60} - \frac{x^{6}}{36} + \frac{x^{8}}{180} \right)$$

$$= \frac{1}{16} + 2 \cdot \frac{15^{2}}{17^{2}} \left(\frac{1}{6.48} + \frac{1}{10.4.36.7} - \frac{1}{3.120} + \frac{1}{10.36.12} + \frac{1}{14.36^{2}.7} - \frac{1}{12.36.30} - \frac{1}{8.6.12} - \frac{1}{12.36.42} + \frac{1}{60.30} \right) - \frac{60}{77} \left(\frac{1}{60} - \frac{1}{420} - \frac{1}{7.36} + \frac{1}{9.180} \right) = 0.13034$$

Hence,
$$\lambda_0 > \frac{1}{\sqrt{0.13034}} \simeq 2.77$$