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## V. ON THE EQUILIBRIUM OF HETEROGENEOUS SUBSTANCES.

BY J. WILLARD GIBBS.

“Die Energie der Welt ist constant.  
Die Entropie der Welt strebt einem Maximum zu.”  
CLAUSIUS.\*

THE comprehension of the laws which govern any material system is greatly facilitated by considering the energy and entropy of the system in the various states of which it is capable. As the difference of the values of the energy for any two states represents the combined amount of work and heat received or yielded by the system when it is brought from one state to the other, and the difference of entropy is the limit of all the possible values of the integral  $\int \frac{dQ}{t}$ , ( $dQ$  denoting the element of the heat received from external sources, and  $t$  the temperature of the part of the system receiving it,) the varying values of the energy and entropy characterize in all that is essential the effects producible by the system in passing from one state to another. For by mechanical and thermodynamic contrivances, supposed theoretically perfect, any supply of work and heat may be transformed into any other which does not differ from it either in the amount of work and heat taken together or in the value of the integral  $\int \frac{dQ}{t}$ . But it is not only in respect to the external relations of a system that its energy and entropy are of predominant importance. As in the case of simply mechanical systems, (such as are discussed in theoretical mechanics,) which are capable of only one kind of action upon external systems, viz., the performance of mechanical work, the function which expresses the capability of the system for this kind of action also plays the leading part in the theory of equilibrium, the condition of equilibrium being that the variation of this function shall vanish, so in a thermodynamic system, (such as all material systems actually are,) which is capable of two different kinds of action upon external systems, the two functions which express the twofold capabilities of the system afford an almost equally simple criterion of equilibrium.

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\*Pogg. Ann. Bd. exxv (1865), S. 400; or Mechanische Wärmetheorie, Abhand. ix., S. 44.

## CRITERIA OF EQUILIBRIUM AND STABILITY.

The criterion of equilibrium for a material system which is isolated from all external influences may be expressed in either of the following entirely equivalent forms:

I. *For the equilibrium of any isolated system it is necessary and sufficient that in all possible variations of the state of the system which do not alter its energy, the variation of its entropy shall either vanish or be negative.* If  $\varepsilon$  denote the energy, and  $\eta$  the entropy of the system, and we use a subscript letter after a variation to indicate a quantity of which the value is not to be varied, the condition of equilibrium may be written

$$(\delta\eta)_{\varepsilon} \leq 0. \quad (1)$$

II. *For the equilibrium of any isolated system it is necessary and sufficient that in all possible variations in the state of the system which do not alter its entropy, the variation of its energy shall either vanish or be positive.* This condition may be written

$$(\delta\varepsilon)_{\eta} \geq 0. \quad (2)$$

That these two theorems are equivalent will appear from the consideration that it is always possible to increase both the energy and the entropy of the system, or to decrease both together, viz., by imparting heat to any part of the system or by taking it away. For, if condition (1) is not satisfied, there must be some variation in the state of the system for which

$$\delta\eta > 0 \text{ and } \delta\varepsilon = 0;$$

therefore, by diminishing both the energy and the entropy of the system *in its varied state*, we shall obtain a state for which (considered as a variation from the original state)

$$\delta\eta = 0 \text{ and } \delta\varepsilon < 0;$$

therefore condition (2) is not satisfied. Conversely, if condition (2) is not satisfied, there must be a variation in the state of the system for which

$$\delta\varepsilon < 0 \text{ and } \delta\eta = 0;$$

hence there must also be one for which

$$\delta\varepsilon = 0 \text{ and } \delta\eta > 0;$$

therefore condition (1) is not satisfied.

The equations which express the condition of equilibrium, as also its statement in words, are to be interpreted in accordance with the general usage in respect to differential equations, that is, infinitesimals

of higher orders than the first relatively to those which express the amount of change of the system are to be neglected. But to distinguish the different kinds of equilibrium in respect to stability, we must have regard to the absolute values of the variations. We will use  $\Delta$  as the sign of variation in those equations which are to be construed *strictly*, i. e., in which infinitesimals of the higher orders are not to be neglected. With this understanding, we may express the necessary and sufficient conditions of the different kinds of equilibrium as follows;—for stable equilibrium

$$(\Delta\eta)_\varepsilon < 0, \text{ i. e., } (\Delta\varepsilon)_\eta > 0; \quad (3)$$

for neutral equilibrium there must be some variations in the state of the system for which

$$(\Delta\eta)_\varepsilon = 0, \text{ i. e., } (\Delta\varepsilon)_\eta = 0, \quad (4)$$

while in general

$$(\Delta\eta)_\varepsilon \leq 0, \text{ i. e., } (\Delta\varepsilon)_\eta \geq 0; \quad (5)$$

and for unstable equilibrium there must be some variations for which

$$(\Delta\eta)_\varepsilon > 0, \quad (6)$$

i. e., there must be some for which

$$(\Delta\varepsilon)_\eta < 0, \quad (7)$$

while in general

$$(\delta\eta)_\varepsilon \leq 0, \text{ i. e., } (\delta\varepsilon)_\eta \geq 0. \quad (8)$$

In these criteria of equilibrium and stability, account is taken only of *possible* variations. It is necessary to explain in what sense this is to be understood. In the first place, all variations in the state of the system which involve the transportation of any matter through any finite distance are of course to be excluded from consideration, although they may be capable of expression by infinitesimal variations of quantities which perfectly determine the state of the system. For example, if the system contains two masses of the same substance, not in contact, nor connected by other masses consisting of or containing the same substance or its components, an infinitesimal increase of the one mass with an equal decrease of the other is not to be considered as a possible variation in the state of the system. In addition to such cases of essential impossibility, if heat can pass by conduction or radiation from every part of the system to every other, only those variations are to be rejected as impossible, which involve changes which are prevented by passive forces or analogous resistances to change. But, if the system consist of parts between which there is supposed to be no thermal communication, it will be necessary to regard as impossible any diminution of the entropy of any of

these parts, as such a change can not take place without the passage of heat. This limitation may most conveniently be applied to the second of the above forms of the condition of equilibrium, which will then become

$$(\delta\varepsilon)_{\eta', \eta'', \text{etc.}} \geq 0, \quad (9)$$

$\eta', \eta'', \text{etc.}$ , denoting the entropies of the various parts between which there is no communication of heat. When the condition of equilibrium is thus expressed, the limitation in respect to the conduction of heat will need no farther consideration.

In order to apply to any system the criteria of equilibrium which have been given, a knowledge is requisite of its passive forces or resistances to change, in so far, at least, as they are capable of *preventing* change. (Those passive forces which only retard change, like viscosity, need not be considered.) Such properties of a system are in general easily recognized upon the most superficial knowledge of its nature. As examples, we may instance the passive force of friction which prevents sliding when two surfaces of solids are pressed together,—that which prevents the different components of a solid, and sometimes of a fluid, from having different motions one from another,—that resistance to change which sometimes prevents either of two forms of the same substance (simple or compound), which are capable of existing, from passing into the other,—that which prevents the changes in solids which imply plasticity, (in other words, changes of the form to which the solid tends to return,) when the deformation does not exceed certain limits.

It is a characteristic of all these passive resistances that they prevent a certain kind of motion or change, however the initial state of the system may be modified, and to whatever external agencies of force and heat it may be subjected, within limits, it may be, but yet within limits which allow finite variations in the values of all the quantities which express the initial state of the system or the mechanical or thermal influences acting on it, without producing the change in question. The equilibrium which is due to such passive properties is thus widely distinguished from that caused by the balance of the active tendencies of the system, where an external influence, or a change in the initial state, infinitesimal in amount, is sufficient to produce change either in the positive or negative direction. Hence the ease with which these passive resistances are recognized. Only in the ease that the state of the system lies so near the limit at which the resistances cease to be operative to prevent change, as to create a

doubt whether the case falls within or without the limit, will a more accurate knowledge of these resistances be necessary.

To establish the validity of the criterion of equilibrium, we will consider first the sufficiency, and afterwards the necessity, of the condition as expressed in either of the two equivalent forms.

In the first place, if the system is in a state in which its entropy is greater than in any other state of the same energy, it is evidently in equilibrium, as any change of state must involve either a decrease of entropy or an increase of energy, which are alike impossible for an isolated system. We may add that this is a case of *stable* equilibrium, as no infinitely small cause (whether relating to a variation of the initial state or to the action of any external bodies) can produce a finite change of state, as this would involve a finite decrease of entropy or increase of energy.

We will next suppose that the system has the greatest entropy consistent with its energy, and therefore the least energy consistent with its entropy, but that there are other states of the same energy and entropy as its actual state. In this case, it is impossible that any motion of masses should take place; for if any of the energy of the system should come to consist of *vis viva* (of sensible motions), a state of the system identical in other respects but without the motion would have less energy and not less entropy, which would be contrary to the supposition. (But we cannot apply this reasoning to the motion within any mass of its different components in different directions, as in diffusion, when the momenta of the components balance one another.) Nor, in the case supposed, can any conduction of heat take place, for this involves an increase of entropy, as heat is only conducted from bodies of higher to those of lower temperature. It is equally impossible that any changes should be produced by the transfer of heat by radiation. The condition which we have supposed is therefore sufficient for equilibrium, so far as the motion of masses and the transfer of heat are concerned, but to show that the same is true in regard to the motions of diffusion and chemical or molecular changes, when these can occur without being accompanied or followed by the motions of masses or the transfer of heat, we must have recourse to considerations of a more general nature. The following considerations seem to justify the belief that the condition is sufficient for equilibrium in every respect.

Let us suppose, in order to test the tenability of such a hypothesis, that a system may have the greatest entropy consistent with its energy without being in equilibrium. In such a case, changes in the

state of the system must take place, but these will necessarily be such that the energy and the entropy will remain unchanged and the system will continue to satisfy the same condition, as initially, of having the greatest entropy consistent with its energy. Let us consider the change which takes place in any time so short that the change may be regarded as uniform in nature throughout that time. This time must be so chosen that the change does not take place in it infinitely slowly, which is always easy, as the change which we suppose to take place cannot be infinitely slow except at particular moments. Now no change whatever in the state of the system, which does not alter the value of the energy, and which commences with the same state in which the system was supposed at the commencement of the short time considered, will cause an increase of entropy. Hence, it will generally be possible by some slight variation in the circumstances of the case to make all changes in the state of the system like or nearly like that which is supposed actually to occur, and not involving a change of energy, to involve a necessary decrease of entropy, which would render any such change impossible. This variation may be in the values of the variables which determine the state of the system, or in the values of the constants which determine the nature of the system, or in the form of the functions which express its laws,—only there must be nothing in the system as modified which is thermodynamically impossible. For example, we might suppose temperature or pressure to be varied, or the composition of the different bodies in the system, or, if no small variations which could be actually realized would produce the required result, we might suppose the properties themselves of the substances to undergo variation, subject to the general laws of matter. If, then, there is any tendency toward change in the system as first supposed, it is a tendency which can be entirely checked by an infinitesimal variation in the circumstances of the case. As this supposition cannot be allowed, we must believe that a system is always in equilibrium when it has the greatest entropy consistent with its energy, or, in other words, when it has the least energy consistent with its entropy.

The same considerations will evidently apply to any case in which a system is in such a state that  $\Delta\eta \leq 0$  for any possible infinitesimal variation of the state for which  $\Delta\varepsilon = 0$ , even if the entropy is not the least of which the system is capable with the same energy. (The term *possible* has here the meaning previously defined, and the character  $\Delta$  is used, as before, to denote that the equations are to be

constrained strictly, i. e., without neglect of the infinitesimals of the higher orders.)

The only case in which the sufficiency of the condition of equilibrium which has been given remains to be proved is that in which in our notation  $\delta\eta \leq 0$  for all possible variations not affecting the energy, but for some of these variations  $\Delta\eta > 0$ , that is, when the entropy has in some respects the characteristics of a minimum. In this case the considerations adduced in the last paragraph will not apply without modification, as the change of state may be infinitely slow at first, and it is only in the initial state that the condition  $\delta\eta_e \leq 0$  holds true. But the differential coefficients of all orders of the quantities which determine the state of the system, taken with respect of the time, must be functions of these same quantities. None of these differential coefficients can have any value other than 0, for the state of the system for which  $\delta\eta_e \leq 0$ . For otherwise, as it would generally be possible, as before, by some infinitely small modification of the case, to render impossible any change like or nearly like that which might be supposed to occur, this infinitely small modification of the case would make a finite difference in the value of the differential coefficients which had before the finite values, or in some of lower orders, which is contrary to that continuity which we have reason to expect. Such considerations seem to justify us in regarding such a state as we are discussing as one of theoretical equilibrium; although as the equilibrium is evidently unstable, it cannot be realized.

We have still to prove that the condition enunciated is in every case necessary for equilibrium. It is evidently so in all cases in which the active tendencies of the system are so balanced that changes of every kind, except those excluded in the statement of the condition of equilibrium, can take place *reversibly*, (i. e., both in the positive and the negative direction,) in states of the system differing infinitely little from the state in question. In this case, we may omit the sign of inequality and write as the condition of such a state of equilibrium

$$(\delta\eta)_e = 0, \quad \text{i. e.,} \quad (\delta\varepsilon)_\eta = 0 \quad (10)$$

But to prove that the condition previously enunciated is in every case necessary, it must be shown that whenever an isolated system remains without change, if there is any infinitesimal variation in its state, not involving a finite change of position of any (even an infinitesimal part) of its matter, which would diminish its energy by a

quantity which is not infinitely small relatively to the variations of the quantities which determine the state of the system, without altering its entropy,—or, if the system has thermally isolated parts, without altering the entropy of any such part,—this variation involves changes in the system which are prevented by its passive forces or analogous resistances to change. Now, as the described variation in the state of the system diminishes its energy without altering its entropy, it must be regarded as theoretically possible to produce that variation by some process, perhaps a very indirect one, so as to gain a certain amount of work (above all expended on the system). Hence we may conclude that the active forces or tendencies of the system favor the variation in question, and that equilibrium cannot subsist unless the variation is prevented by passive forces.

The preceding considerations will suffice, it is believed, to establish the validity of the criterion of equilibrium which has been given. The criteria of stability may readily be deduced from that of equilibrium. We will now proceed to apply these principles to systems consisting of heterogeneous substances and deduce the special laws which apply to different classes of phenomena. For this purpose we shall use the second form of the criterion of equilibrium, both because it admits more readily the introduction of the condition that there shall be no thermal communication between the different parts of the system, and because it is more convenient, as respects the form of the general equations relating to equilibrium, to make the entropy one of the independent variables which determine the state of the system, than to make the energy one of these variables.

#### THE CONDITIONS OF EQUILIBRIUM FOR HETEROGENEOUS MASSES IN CONTACT WHEN UNINFLUENCED BY GRAVITY, ELECTRICITY, DISTORTION OF THE SOLID MASSES, OR CAPILLARY TENSIONS.

In order to arrive as directly as possible at the most characteristic and essential laws of chemical equilibrium, we will first give our attention to a case of the simplest kind. We will examine the conditions of equilibrium of a mass of matter of various kinds enclosed in a rigid and fixed envelop, which is impermeable to and unalterable by any of the substances enclosed, and perfectly non-conducting to heat. We will suppose that the case is not complicated by the action of gravity, or by any electrical influences, and that in the solid portions of the mass the pressure is the same in every direction.

We will farther simplify the problem by supposing that the variations of the parts of the energy and entropy which depend upon the surfaces separating heterogeneous masses are so small in comparison with the variations of the parts of the energy and entropy which depend upon the quantities of these masses, that the former may be neglected by the side of the latter; in other words, we will exclude the considerations which belong to the theory of capillarity.

It will be observed that the supposition of a rigid and non-conducting envelop enclosing the mass under discussion involves no real loss of generality, for if any mass of matter is in equilibrium, it would also be so, if the whole or any part of it were enclosed in an envelop as supposed; therefore the conditions of equilibrium for a mass thus enclosed are the general conditions which must always be satisfied in case of equilibrium. As for the other suppositions which have been made, all the circumstances and considerations which are here excluded will afterward be made the subject of special discussion.

*Conditions relating to the Equilibrium between the initially existing Homogeneous Parts of the given Mass.*

Let us first consider the energy of any homogeneous part of the given mass, and its variation for any possible variation in the composition and state of this part. (By *homogeneous* is meant that the part in question is uniform throughout, not only in chemical composition, but also in physical state.) If we consider the amount and kind of matter in this homogeneous mass as fixed, its energy  $\varepsilon$  is a function of its entropy  $\eta$ , and its volume  $v$ , and the differentials of these quantities are subject to the relation

$$d\varepsilon = t d\eta - p dv, \quad (11)$$

$t$  denoting the (absolute) temperature of the mass, and  $p$  its pressure. For  $t dv$  is the heat received, and  $p dv$  the work done, by the mass during its change of state. But if we consider the matter in the mass as variable, and write  $m_1, m_2, \dots, m_n$  for the quantities of the various substances  $S_1, S_2, \dots, S_n$  of which the mass is composed,  $\varepsilon$  will evidently be a function of  $\eta, v, m_1, m_2, \dots, m_n$ , and we shall have for the complete value of the differential of  $\varepsilon$

$$d\varepsilon = t d\eta - p dv + \mu_1 dm_1 + \mu_2 dm_2 + \dots + \mu_n dm_n, \quad (12)$$

$\mu_1, \mu_2, \dots, \mu_n$  denoting the differential coefficients of  $\varepsilon$  taken with respect to  $m_1, m_2, \dots, m_n$ .

The substances  $S_1, S_2, \dots, S_n$  of which we consider the mass composed, must of course be such that the values of the differen-

tials  $dm_1$ ,  $dm_2$ , . . .  $dm_n$  shall be independent, and shall express every possible variation in the composition of the homogeneous mass considered, including those produced by the absorption of substances different from any initially present. It may therefore be necessary to have terms in the equation relating to component substances which do not initially occur in the homogeneous mass considered, provided, of course, that these substances, or their components, are to be found in some part of the whole given mass.

If the conditions mentioned are satisfied, the choice of the substances which we are to regard as the components of the mass considered, may be determined entirely by convenience, and independently of any theory in regard to the internal constitution of the mass. The number of components will sometimes be greater, and sometimes less, than the number of chemical elements present. For example, in considering the equilibrium in a vessel containing water and free hydrogen and oxygen, we should be obliged to recognize three components in the gaseous part. But in considering the equilibrium of dilute sulphuric acid with the vapor which it yields, we should have only two components to consider in the liquid mass, sulphuric acid (anhydrous, or of any particular degree of concentration) and (additional) water. If, however, we are considering sulphuric acid in a state of maximum concentration in connection with substances which might possibly afford water to the acid, it must be noticed that the condition of the independence of the differentials will require that we consider the acid in the state of maximum concentration as one of the components. The quantity of this component will then be capable of variation both in the positive and in the negative sense, while the quantity of the other component can increase but cannot decrease below the value 0.

For brevity's sake, we may call a substance  $S_a$  an *actual component* of any homogeneous mass, to denote that the quantity  $m_a$  of that substance in the given mass may be either increased or diminished (although we may have so chosen the other component substances that  $m_a = 0$ ); and we may call a substance  $S_b$  a *possible component* to denote that it may be combined with, but cannot be subtracted from the homogeneous mass in question. In this case, as we have seen in the above example, we must so choose the component substances that  $m_b = 0$ .

The units by which we measure the substances of which we regard the given mass as composed may each be chosen independently. To fix our ideas for the purpose of a general discussion, we may suppose

all substances measured by weight or mass. Yet in special cases, it may be more convenient to adopt chemical equivalents as the units of the component substances.

It may be observed that it is not necessary for the validity of equation (12) that the variations of nature and state of the mass to which the equation refers should be such as do not disturb its homogeneity, provided that in all parts of the mass the variations of nature and state are infinitely small. For, if this last condition be not violated, an equation like (12) is certainly valid for all the infinitesimal parts of the (initially) homogeneous mass; i. e., if we write  $D\varepsilon$ ,  $D\eta$ , etc., for the energy, entropy, etc., of any infinitesimal part,  $dD\varepsilon = t dD\eta - p dDv + \mu_1 dDm_1 + \mu_2 dDm_2 \dots + \mu_n dDm_n$ , (13) whence we may derive equation (12) by integrating for the whole initially homogeneous mass.

We will now suppose that the whole mass is divided into parts so that each part is homogeneous, and consider such variations in the energy of the system as are due to variations in the composition and state of the several parts remaining (at least approximately) homogeneous, and together occupying the whole space within the envelop. We will at first suppose the case to be such that the component substances are the same for each of the parts, each of the substances  $S_1$ ,  $S_2$ , . . .  $S_n$  being an actual component of each part. If we distinguish the letters referring to the different parts by accents, the variation in the energy of the system may be expressed by  $\delta\varepsilon' + \delta\varepsilon'' +$  etc., and the general condition of equilibrium requires that

$$\delta\varepsilon' + \delta\varepsilon'' + \text{etc.} \geq 0 \quad (14)$$

for all variations which do not conflict with the *equations of condition*. These equations must express that the entropy of the whole given mass does not vary, nor its volume, nor the total quantities of any of the substances  $S_1$ ,  $S_2$ , . . .  $S_n$ . We will suppose that there are no other equations of condition. It will then be necessary for equilibrium that

$$\begin{aligned} t' \delta\eta' - p' \delta v' + \mu_1' \delta m_1' + \mu_2' \delta m_2' \dots + \mu_n' \delta m_n' \\ + t'' \delta\eta'' - p'' \delta v'' + \mu_1'' \delta m_1'' + \mu_2'' \delta m_2'' \dots + \mu_n'' \delta m_n'' \\ + \text{etc.} \geq 0 \end{aligned} \quad (15)$$

for any values of the variations for which

$$\delta\eta' + \delta\eta'' + \delta\eta''' + \text{etc.} = 0, \quad (16)$$

$$\delta v' + \delta v'' + \delta v''' + \text{etc.} = 0, \quad (17)$$

$$\left. \begin{array}{l} \delta m_1' + \delta m_1'' + \delta m_1''' + \text{etc.} = 0, \\ \delta m_2' + \delta m_2'' + \delta m_2''' + \text{etc.} = 0, \\ \vdots \quad \vdots \quad \vdots \quad \vdots \\ \delta m_n' + \delta m_n'' + \delta m_n''' + \text{etc.} = 0. \end{array} \right\} \quad (18)$$

For this it is evidently necessary and sufficient that

$$t' = t'' = t''' = \text{etc.} \quad (19)$$

$$p' = p'' = p''' = \text{etc.} \quad (20)$$

$$\left. \begin{array}{l} \mu_1' = \mu_1'' = \mu_1''' = \text{etc.} \\ \mu_2' = \mu_2'' = \mu_2''' = \text{etc.} \\ \vdots \quad \vdots \quad \vdots \\ \mu_n' = \mu_n'' = \mu_n''' = \text{etc.} \end{array} \right\} \quad (21)$$

Equations (19) and (20) express the conditions of thermal and mechanical equilibrium, viz., that the temperature and the pressure must be constant throughout the whole mass. In equations (21) we have the conditions characteristic of chemical equilibrium. If we call a quantity  $\mu_s$ , as defined by such an equation as (12), the *potential* for the substance  $S_s$  in the homogeneous mass considered, these conditions may be expressed as follows:

*The potential for each component substance must be constant throughout the whole mass.*

It will be remembered that we have supposed that there is no restriction upon the freedom of motion or combination of the component substances, and that each is an actual component of all parts of the given mass.

The state of the whole mass will be completely determined (if we regard as immaterial the position and form of the various homogeneous parts of which it is composed), when the values are determined of the quantities of which the variations occur in (15). The number of these quantities, which we may call the independent variables, is evidently  $(n+2)\nu$ ,  $\nu$  denoting the number of homogeneous parts into which the whole mass is divided. All the quantities which occur in (19), (20), (21), are functions of these variables, and may be regarded as known functions, if the energy of each part is known as a function of its entropy, volume, and the quantities of its components. (See eq. (12).) Therefore, equations (19), (20), (21), may be regarded as  $(\nu-1)(n+2)$  independent equations between the independent variables. The volume of the whole mass and the total quantities of the various substances being known afford  $n+1$  additional equations. If we also know the total energy of the given mass, or its total entropy, we will have as many equations as there are independent variables.

But if any of the substances  $S_1, S_2 \dots S_n$  are only possible components of some parts of the given mass, the variation  $\delta m$  of the quantity of such a substance in such a part cannot have a negative value, so that the general condition of equilibrium (15) does not require that the potential for that substance in that part should be equal to the potential for the same substance in the parts of which it is an actual component, but only that it shall not be less. In this case instead of (21) we may write

$$\mu_1 = M_1$$

for all parts of which  $S_1$  is an actual component, and

$$\mu_1 \geq M_1$$

for all parts of which  $S_1$  is a possible (but not actual) component,

$$\mu_2 = M_2$$

for all parts of which  $S_2$  is an actual component, and

$$\mu_2 \geq M_2$$

for all parts of which  $S_2$  is a possible (but not actual) component,

etc.,

$M_1, M_2$ , etc., denoting constants of which the value is only determined by these equations.

If we now suppose that the components (actual or possible) of the various homogeneous parts of the given mass are not the same, the result will be of the same character as before, provided that all the different components are *independent*, (i. e., that no one can be made out of the others,) so that the total quantity of each component is fixed. The general condition of equilibrium (15) and the equations of condition (16), (17), (18) will require no change, except that, if any of the substances  $S_1, S_2 \dots S_n$  is not a component (actual or possible) of any part, the term  $\mu \delta m$  for that substance and part will be wanting in the former, and the  $\delta m$  in the latter. This will require no change in the form of the particular conditions of equilibrium as expressed by (19), (20), (22); but the number of single conditions contained in (22) is of course less than if all the component substances were components of all the parts. Whenever, therefore, each of the different homogeneous parts of the given mass may be regarded as composed of some or of all of the same set of substances, no one of which can be formed out of the others, the condition which (with equality of temperature and pressure) is necessary and sufficient for equilibrium between the different parts of the given mass may be expressed as follows:

The potential for each of the component substances must have a constant value in all parts of the given mass of which that substance is an actual component, and have a value not less than this in all parts of which it is a possible component.

The number of equations afforded by these conditions, after elimination of  $M_1, M_2, \dots, M_n$ , will be less than  $(n + 2)(\nu - 1)$  by the number of terms in (15) in which the variation of the form  $\delta m$  is either necessarily nothing or incapable of a negative value. The number of variables to be determined is diminished by the same number, or, if we choose, we may write an equation of the form  $m = 0$  for each of these terms. But when the substance is a possible component of the part concerned, there will also be a condition (expressed by  $\geq$ ) to show whether the supposition that the substance is not an actual component is consistent with equilibrium.

We will now suppose that the substances  $S_1, S_2, \dots, S_n$  are not all independent of each other, i. e., that some of them can be formed out of others. We will first consider a very simple case. Let  $S_3$  be composed of  $S_1$  and  $S_2$  combined in the ratio of  $a$  to  $b$ ,  $S_1$  and  $S_2$  occurring as actual components in some parts of the given mass, and  $S_3$  in other parts, which do not contain  $S_1$  and  $S_2$  as separately variable components. The general condition of equilibrium will still have the form of (15) with certain of the terms of the form  $\mu \delta m$  omitted. It may be written more briefly [23]

$\Sigma(t \delta \eta) - \Sigma(p \delta v) + \Sigma(\mu_1 \delta m_1) + \Sigma(\mu_2 \delta m_2) \dots + \Sigma(\mu_n \delta m_n) \geq 0$ ,  
the sign  $\Sigma$  denoting summation in regard to the different parts of the given mass. But instead of the three equations of condition,

$$\Sigma \delta m_1 = 0, \quad \Sigma \delta m_2 = 0, \quad \Sigma \delta m_3 = 0, \quad (24)$$

we shall have the two,

$$\left. \begin{aligned} \Sigma \delta m_1 + \frac{a}{a+b} \Sigma \delta m_3 &= 0, \\ \Sigma \delta m_2 + \frac{b}{a+b} \Sigma \delta m_3 &= 0. \end{aligned} \right\} \quad (25)$$

The other equations of condition,

$$\Sigma \delta \eta = 0, \quad \Sigma \delta v = 0, \quad \Sigma \delta m_4 = 0, \quad \text{etc.}, \quad (26)$$

will remain unchanged. Now as all values of the variations which satisfy equations (24) will also satisfy equations (25), it is evident that all the particular conditions of equilibrium which we have already deduced, (19), (20), (22), are necessary in this case also. When these are satisfied, the general condition (23) reduces to

$$M_1 \Sigma \delta m_1 + M_2 \Sigma \delta m_2 + M_3 \Sigma \delta m_3 \geq 0. \quad (27)$$

For, although it may be that  $\mu_1'$ , for example, is greater than  $M_1$ , yet it can only be so when the following  $\delta m_1'$  is incapable of a negative value. Hence, if (27) is satisfied, (23) must also be. Again, if (23) is satisfied, (27) must also be satisfied, so long as the variation of the quantity of every substance has the value 0 in all the parts of which it is not an actual component. But as this limitation does not affect the range of the possible values of  $\Sigma \delta m_1$ ,  $\Sigma \delta m_2$ , and  $\Sigma \delta m_3$ , it may be disregarded. Therefore the conditions (23) and (27) are entirely equivalent, when (19), (20), (22) are satisfied. Now, by means of the equations of condition (25), we may eliminate  $\Sigma \delta m_1$  and  $\Sigma \delta m_2$  from (27), which becomes

$$-aM_1 \geq \delta m_3 - bM_2 \geq \delta m_3 + (a+b)M_3 \geq \delta m_3 \geq 0, \quad (28)$$

i. e., as the value of  $\Sigma \delta m_3$  may be either positive or negative,

$$aM_1 + bM_2 = (a+b)M_3, \quad (29)$$

which is the additional condition of equilibrium which is necessary in this case.

The relations between the component substances may be less simple than in this case, but in any case they will only affect the equations of condition, and these may always be found without difficulty, and will enable us to eliminate from the general condition of equilibrium as many variations as there are equations of condition, after which the coefficients of the remaining variations may be set equal to zero, except the coefficients of variations which are incapable of negative values, which coefficients must be equal to or greater than zero. It will be easy to perform these operations in each particular case, but it may be interesting to see the form of the resultant equations in general.

We will suppose that the various homogeneous parts are considered as having in all  $n$  components,  $S_1, S_2, \dots, S_n$ , and that there is no restriction upon their freedom of motion and combination. But we will so far limit the generality of the problem as to suppose that each of these components is an actual component of some part of the given mass.\* If some of these components can be formed out of others, all such relations can be expressed by equations such as

$$\alpha \mathcal{E}_a + \beta \mathcal{E}_b + \text{etc.} = \kappa \mathcal{E}_k + \lambda \mathcal{E}_l + \text{etc.} \quad (30)$$

where  $\mathcal{E}_a, \mathcal{E}_b, \mathcal{E}_k$ , etc. denote the units of the substances  $S_a, S_b, S_k$ , etc.,

\* When we come to seek the conditions of equilibrium relating to the formation of masses unlike any previously existing, we shall take up *de novo* the whole problem of the equilibrium of heterogeneous masses enclosed in a non-conducting envelop, and give it a more general treatment, which will be free from this limitation.

(that is, of certain of the substances  $S_1, S_2, \dots, S_n$ ) and  $\alpha, \beta, \kappa$ , etc. denote numbers. These are not, it will be observed, equations between abstract quantities, but the sign  $\Sigma$  denotes qualitative as well as quantitative equivalence. We will suppose that there are  $r$  independent equations of this character. The equations of condition relating to the component substances may easily be derived from these equations, but it will not be necessary to consider them particularly. It is evident that they will be satisfied by any values of the variations which satisfy equations (18); hence, the particular conditions of equilibrium (19), (20), (22) must be necessary in this case, and, if these are satisfied, the general equation of equilibrium (15) or (23) will reduce to

$$M_1 \Sigma \delta m_1 + M_2 \Sigma \delta m_2 + \dots + M_n \Sigma \delta m_n \geq 0. \quad (31)$$

This will appear from the same considerations which were used in regard to equations (23) and (27). Now it is evidently possible to give to  $\Sigma \delta m_a, \Sigma \delta m_b, \Sigma \delta m_k$ , etc. values proportional to  $\alpha, \beta, -\kappa$ , etc. in equation (30), and also the same values taken negatively, making  $\Sigma \delta m = 0$  in each of the other terms; therefore

$$\alpha M_a + \beta M_b + \text{etc.} - \kappa M_k - \lambda M_l - \text{etc.} = 0, \quad (32)$$

or,

$$\alpha M_a + \beta M_b + \text{etc.} = \kappa M_k + \lambda M_l + \text{etc.} \quad (33)$$

It will be observed that this equation has the same form and coefficients as equation (30),  $M$  taking the place of  $\mathfrak{S}$ . It is evident that there must be a similar condition of equilibrium for every one of the  $r$  equations of which (30) is an example, which may be obtained simply by changing  $\mathfrak{S}$  in these equations into  $M$ . When these conditions are satisfied, (31) will be satisfied with any possible values of  $\Sigma \delta m_1, \Sigma \delta m_2, \dots, \Sigma \delta m_n$ . For no values of these quantities are possible, except such that the equation

$$(\Sigma \delta m_1) \mathfrak{S}_1 + (\Sigma \delta m_2) \mathfrak{S}_2 + \dots + (\Sigma \delta m_n) \mathfrak{S}_n = 0 \quad (34)$$

after the substitution of these values, can be derived from the  $r$  equations like (30), by the ordinary processes of the reduction of linear equations. Therefore, on account of the correspondence between (31) and (34), and between the  $r$  equations like (33) and the  $r$  equations like (30), the conditions obtained by giving any possible values to the variations in (31) may also be derived from the  $r$  equations like (33); that is, the condition (31) is satisfied, if the  $r$  equations like (33) are satisfied. Therefore the  $r$  equations like (33) are with (19), (20), and (22) the equivalent of the general condition (15) or (23).

For determining the state of a given mass when in equilibrium and having a given volume and given energy or entropy, the condition of equilibrium affords an additional equation corresponding to each of the  $r$  independent relations between the  $n$  component substances. But the equations which express our knowledge of the matter in the given mass will be correspondingly diminished, being  $n - r$  in number, like the equations of condition relating to the quantities of the component substances, which may be derived from the former by differentiation.

*Conditions relating to the possible Formation of Masses Unlike any Previously Existing.*

The variations which we have hitherto considered do not embrace every possible infinitesimal variation in the state of the given mass, so that the particular conditions already formed, although always necessary for equilibrium (when there are no other equations of condition than such as we have supposed), are not always sufficient. For, besides the infinitesimal variations in the state and composition of different parts of the given mass, infinitesimal masses may be formed entirely different in state and composition from any initially existing. Such parts of the whole mass in its varied state as cannot be regarded as parts of the initially existing mass which have been infinitesimally varied in state and composition, we will call *new parts*. These will necessarily be infinitely small. As it is more convenient to regard a vacuum as a limiting case of extreme rarefaction than to give a special consideration to the possible formation of empty spaces within the given mass, the term *new parts* will be used to include any empty spaces which may be formed, when such have not existed initially. We will use  $D\epsilon$ ,  $D\eta$ ,  $Dv$ ,  $Dm_1$ ,  $Dm_2$ , . . .  $Dm_n$  to denote the infinitesimal energy, entropy, and volume of any one of these new parts, and the infinitesimal quantities of its components. The component substances  $S_1$ ,  $S_2$ , . . .  $S_n$  must now be taken to include not only the independently variable components (actual or possible) of all parts of the given mass as initially existing, but also the components of all the new parts, the possible formation of which we have to consider. The character  $\delta$  will be used as before to express the infinitesimal variations of the quantities relating to those parts which are only infinitesimally varied in state and composition, and which for distinction we will call *original parts*, including under this term the empty spaces, if such exist initially, within the envelop bounding the system. As we may divide the given mass into as many parts as we choose, and as not only the

initial boundaries, but also the movements of these boundaries during any variation in the state of the system are arbitrary, we may so define the parts which we have called original, that we may consider them as initially homogeneous and remaining so, and as initially constituting the whole system.

The most general value of the energy of the whole system is evidently

$$\Sigma \delta \varepsilon + \Sigma D\varepsilon, \quad (35)$$

the first summation relating to all the original parts, and the second to all the new parts. (Throughout the discussion of this problem, the letter  $\delta$  or  $D$  following  $\Sigma$  will sufficiently indicate whether the summation relates to the original or to the new parts.) Therefore the general condition of equilibrium is

$$\Sigma \delta \varepsilon + \Sigma D\varepsilon \geq 0, \quad (36)$$

or, if we substitute the value of  $\delta \varepsilon$  taken from equation (12), [37]

$$\Sigma D\varepsilon + \Sigma (t\delta\eta) - \Sigma (\mu\delta v) + \Sigma (\mu_1\delta m_1) + \Sigma (\mu_2\delta m_2) \dots + \Sigma (\mu_n\delta m_n) \geq 0.$$

If any of the substances  $S_1, S_2, \dots, S_n$  can be formed out of others, we will suppose, as before (see page 122), that such relations are expressed by equations between the units of the different substances. Let these be

$$\left. \begin{array}{l} a_1 \mathcal{S}_1 + a_2 \mathcal{S}_2 \dots + a_n \mathcal{S}_n = 0 \\ b_1 \mathcal{S}_1 + b_2 \mathcal{S}_2 \dots + b_n \mathcal{S}_n = 0 \\ \text{etc.} \end{array} \right\} r \text{ equations.} \quad (38)$$

The equations of condition will be (if there is no restriction upon the freedom of motion and composition of the components)

$$\Sigma \delta\eta + \Sigma D\eta = 0, \quad (39)$$

$$\Sigma \delta v + \Sigma Dv = 0, \quad (40)$$

and  $n - r$  equations of the form

$$\left. \begin{array}{l} h_1 (\Sigma \delta m_1 + \Sigma Dm_1) + h_2 (\Sigma \delta m_2 + \Sigma Dm_2) \dots \\ \quad + h_n (\Sigma \delta m_n + \Sigma Dm_n) = 0 \\ i_1 (\Sigma \delta m_1 + \Sigma Dm_1) + i_2 (\Sigma \delta m_2 + \Sigma Dm_2) \dots \\ \quad + i_n (\Sigma \delta m_n + \Sigma Dm_n) = 0 \\ \text{etc.} \end{array} \right\} (41)^*$$

\* In regard to the relation between the coefficients in (41) and those in (38), the reader will easily convince himself that the coefficients of any one of equations (41) are such as would satisfy all the equations (38) if substituted for  $S_1, S_2, \dots, S_n$ ; and that this is the only condition which these coefficients must satisfy, except that the  $n - r$  sets of coefficients shall be independent, i.e., shall be such as to form independent equations; and that this relation between the coefficients of the two sets of equations is a reciprocal one.

Now, using Lagrange's "method of multipliers,"\* we will subtract  $T(\Sigma \delta\eta + \Sigma D\eta) - P(\Sigma \delta v + \Sigma Dv)$  from the first member of the general condition of equilibrium (37),  $T$  and  $P$  being constants of which the value is as yet arbitrary. We might proceed in the same way with the remaining equations of condition, but we may obtain the same result more simply in another way. We will first observe that

$$(\Sigma \delta m_1 + \Sigma Dm_1) \mathfrak{S}_1 + (\Sigma \delta m_2 + \Sigma Dm_2) \mathfrak{S}_2 \dots + (\Sigma \delta m_n + \Sigma Dm_n) \mathfrak{S}_n = 0, \quad (42)$$

which equation would hold identically for any possible values of the quantities in the parentheses, if for  $r$  of the letters  $\mathfrak{S}_1, \mathfrak{S}_2, \dots, \mathfrak{S}_n$  were substituted their values in terms of the others as derived from equations (38). (Although  $\mathfrak{S}_1, \mathfrak{S}_2, \dots, \mathfrak{S}_n$  do not represent abstract quantities, yet the operations necessary for the reduction of linear equations are evidently applicable to equations (38).) Therefore, equation (42) will hold true if for  $\mathfrak{S}_1, \mathfrak{S}_2, \dots, \mathfrak{S}_n$  we substitute  $n$  numbers which satisfy equations (38). Let  $M_1, M_2, \dots, M_n$  be such numbers, i. e., let

$$\left. \begin{array}{l} a_1 M_1 + a_2 M_2 \dots + a_n M_n = 0, \\ b_1 M_1 + b_2 M_2 \dots + b_n M_n = 0, \\ \text{etc.} \end{array} \right\} r \text{ equations, (43)}$$

then

$$M_1 (\Sigma \delta m_1 + \Sigma Dm_1) + M_2 (\Sigma \delta m_2 + \Sigma Dm_2) \dots + M_n (\Sigma \delta m_n + \Sigma Dm_n) = 0. \quad (44)$$

This expression, in which the values of  $n - r$  of the constants  $M_1, M_2, \dots, M_n$  are still arbitrary, we will also subtract from the first member of the general condition of equilibrium (37), which will then become

$$\begin{aligned} & \Sigma D\varepsilon + \Sigma (t \delta\eta) - \Sigma (p \delta v) + \Sigma (\mu_1 \delta m_1) \dots + \Sigma (\mu_n \delta m_n) \\ & - T \Sigma \delta\eta + P \Sigma \delta v - M_1 \Sigma \delta m_1 \dots + M_n \Sigma \delta m_n \\ & - T \Sigma D\eta + P \Sigma Dv - M_1 \Sigma Dm_1 \dots - M_n \Sigma Dm_n \geq 0. \end{aligned} \quad (45)$$

That is, having assigned to  $T, P, M_1, M_2, \dots, M_n$  any values consistent with (43), we may assert that it is necessary and sufficient for equilibrium that (45) shall hold true for any variations in the state of the system consistent with the equations of condition (39), (40), (41). But it will always be possible, in case of equilibrium, to assign such values to  $T, P, M_1, M_2, \dots, M_n$ , without violating equations (43),

\* On account of the sign  $\geq$  in (37), and because some of the variations are incapable of negative values, the successive steps in the reasoning will be developed at greater length than would be otherwise necessary.

that (45) shall hold true for all variations in the state of the system and in the quantities of the various substances composing it, even though these variations are not consistent with the equations of condition (39), (40), (41). For, when it is not possible to do this, it must be possible by applying (45) to variations in the system not necessarily restricted by the equations of condition (39), (40), (41) to obtain conditions in regard to  $T$ ,  $P$ ,  $M_1$ ,  $M_2$ , . . .  $M_n$ , some of which will be inconsistent with others or with equations (43). These conditions we will represent by

$$A \geq 0, \quad B \geq 0, \quad \text{etc.}, \quad (46)$$

$A$ ,  $B$ , etc. being linear functions of  $T$ ,  $P$ ,  $M_1$ ,  $M_2$ , . . .  $M_n$ . Then it will be possible to deduce from these conditions a single condition of the form

$$\alpha A + \beta B + \text{etc.} \geq 0, \quad (47)$$

$\alpha$ ,  $\beta$ , etc. being positive constants, which cannot hold true consistently with equations (43). But it is evident from the form of (47) that, like any of the conditions (46), it could have been obtained directly from (45) by applying this formula to a certain change in the system (perhaps not restricted by the equations of condition (39), (40), (41)). Now as (47) cannot hold true consistently with eqs. (43), it is evident, in the first place, that it cannot contain  $T$  or  $P$ , therefore in the change in the system just mentioned (for which (45) reduces to (47))

$$\Sigma \delta\eta + \Sigma D\eta = 0, \quad \text{and} \quad \Sigma \delta v + \Sigma Dv = 0,$$

so that the equations of condition (39) and (40) are satisfied. Again, for the same reason, the homogeneous function of the first degree of  $M_1$ ,  $M_2$ , . . .  $M_n$  in (47) must be one of which the value is fixed by eqs. (43). But the value thus fixed can only be zero, as is evident from the form of these equations. Therefore

$$\begin{aligned} (\Sigma \delta m_1 + \Sigma Dm_1) M_1 + (\Sigma \delta m_2 + \Sigma Dm_2) M_2 & \dots \\ + (\Sigma \delta m_n + \Sigma Dm_n) M_n & = 0 \end{aligned} \quad (48)$$

for any values of  $M_1$ ,  $M_2$  . . .  $M_n$  which satisfy eqs. (43), and therefore

$$\begin{aligned} (\Sigma \delta m_1 + \Sigma Dm_1) \mathfrak{S}_1 + (\Sigma \delta m_2 + \Sigma Dm_2) \mathfrak{S}_2 & \dots \\ + (\Sigma \delta m_n + \Sigma Dm_n) \mathfrak{S}_n & = 0 \end{aligned} \quad (49)$$

for any numerical values of  $\mathfrak{S}_1$ ,  $\mathfrak{S}_2$ , . . .  $\mathfrak{S}_n$  which satisfy eqs. (38). This equation (49) will therefore hold true, if for  $r$  of the letters  $\mathfrak{S}_1$ ,  $\mathfrak{S}_2$ , . . .  $\mathfrak{S}_n$  we substitute their values in terms of the others taken from eqs. (38), and therefore it will hold true when we use

$\mathfrak{S}_1, \mathfrak{S}_2, \dots, \mathfrak{S}_n$ , as before, to denote the units of the various components. Thus understood, the equation expresses that the values of the quantities in the parentheses are such as are consistent with the equations of condition (41). The change in the system, therefore, which we are considering, is not one which violates any of the equations of condition, and as (45) does not hold true for this change, and for all values of  $T, P, M_1, M_2, \dots, M_n$  which are consistent with eqs. (43), the state of the system cannot be one of equilibrium. Therefore it is necessary, and it is evidently sufficient for equilibrium, that it shall be possible to assign to  $T, P, M_1, M_2, \dots, M_n$  such values, consistent with eqs. (43), that the condition (45) shall hold true for any change in the system irrespective of the equations of condition (39), (40), (41).

For this it is necessary and sufficient that

$$t = T, \quad p = P, \quad (50)$$

$$\mu_1 \delta m_1 \geq M_1 \delta m_1, \quad \mu_2 \delta m_2 \geq M_2 \delta m_2, \dots, \quad \mu_n \delta m_n \geq M_n \delta m_n \quad (51)$$

for each of the *original parts* as previously defined, and that

$$D\varepsilon - TD\eta + PDv - M_1 Dm_1 - M_2 Dm_2 - \dots - M_n Dm_n \geq 0, \quad (52)$$

for each of the *new parts* as previously defined. If to these conditions we add equations (43), we may treat  $T, P, M_1, M_2, \dots, M_n$  simply as unknown quantities to be eliminated.

In regard to conditions (51), it will be observed that if a substance,  $S_1$ , is an actual component of the part of the given mass distinguished by a single accent,  $\delta m'_1$  may be either positive or negative, and we shall have  $\mu'_1 = M_1$ ; but if  $S_1$  is only a possible component of that part,  $\delta m'_1$  will be incapable of a negative value, and we will have  $\mu'_1 \geq M_1$ .

The formulae (50), (51), and (43) express the same particular conditions of equilibrium which we have before obtained by a less general process. It remains to discuss (52). This formula must hold true of any infinitesimal mass in the system in its varied state which is not approximately homogeneous with any of the surrounding masses, the expressions  $D\varepsilon, D\eta, Dv, Dm_1, Dm_2, \dots, Dm_n$  denoting the energy, entropy, and volume of this infinitesimal mass, and the quantities of the substances  $S_1, S_2, \dots, S_n$  which we regard as composing it, (not necessarily as *independently* variable components). If there is more than one way in which this mass may be considered as composed of these substances, we may choose whichever is most convenient. Indeed it follows directly from the relations existing between  $M_1, M_2, \dots, M_n$  that the result would be the same in

any case. Now, if we assume that the values of  $D\varepsilon$ ,  $D\eta$ ,  $Dv$ ,  $Dm_1$ ,  $Dm_2$ , ...,  $Dm_n$  are proportional to the values of  $\varepsilon$ ,  $\eta$ ,  $v$ ,  $m_1$ ,  $m_2$ , ...,  $m_n$  for any large homogeneous mass of similar composition, and of the same temperature and pressure, the condition is equivalent to this, that

$$\varepsilon - T\eta + Pv - M_1 m_1 - M_2 m_2 - \dots - M_n m_n \geq 0 \quad (53)$$

for any large homogeneous body which can be formed out of the substances  $S_1$ ,  $S_2$ , ...,  $S_n$ .

But the validity of this last transformation cannot be admitted without considerable limitation. It is assumed that the relation between the energy, entropy, volume, and the quantities of the different components of a very small mass surrounded by substances of different composition and state is the same as if the mass in question formed a part of a large homogeneous body. We started, indeed, with the assumption that we might neglect the part of the energy, etc., depending upon the surfaces separating heterogeneous masses. Now, in many cases, and for many purposes, as, in general, when the masses are large, such an assumption is quite legitimate, but in the case of these masses which are formed within or among substances of different nature or state, and which at their first formation must be infinitely small, the same assumption is evidently entirely inadmissible, as the surfaces must be regarded as infinitely large in proportion to the masses. We shall see hereafter what modifications are necessary in our formulae in order to include the parts of the energy, etc., which are due to the surfaces, but this will be on the assumption, which is usual in the theory of capillarity, that the radius of curvature of the surfaces is large in proportion to the radius of sensible molecular action, and also to the thickness of the lamina of matter at the surface which is not (sensibly) homogeneous in all respects with either of the masses which it separates. But although the formulae thus modified will apply with sensible accuracy to masses (occurring within masses of a different nature) much smaller than if the terms relating to the surfaces were omitted, yet their failure when applied to masses infinitely small in all their dimensions is not less absolute.

Considerations like the foregoing might render doubtful the validity even of (52) as the necessary and sufficient condition of equilibrium in regard to the formation of masses not approximately homogeneous with those previously existing, when the conditions of equilibrium between the latter are satisfied, unless it is shown that in establishing this formula there have been no quantities neglected relating to the

mutual action of the new and the original parts, which can affect the result. It will be easy to give such a meaning to the expressions  $D\varepsilon$ ,  $D\eta$ ,  $Dv$ ,  $Dm_1$ ,  $Dm_2$ , . . .  $Dm_n$  that this shall be evidently the case. It will be observed that the quantities represented by these expressions have not been perfectly defined. In the first place, we have no right to assume the existence of any surface of absolute discontinuity to divide the new parts from the original, so that the position given to the dividing surface is to a certain extent arbitrary. Even if the surface separating the masses were determined, the energy to be attributed to the masses separated would be partly arbitrary, since a part of the total energy depends upon the mutual action of the two masses. We ought perhaps to consider the case the same in regard to the entropy, although the entropy of a system never depends upon the mutual relations of parts at sensible distances from one another. Now the condition (52) will be valid if the quantities  $D\varepsilon$ ,  $D\eta$ ,  $Dv$ ,  $Dm_1$ ,  $Dm_2$  . . .  $Dm_n$  are so defined that none of the assumptions which have been made, tacitly or otherwise, relating to the formation of these new parts, shall be violated. These assumptions are the following:—that the relation between the variations of the energy, entropy, volume, etc., of any of the original parts is not affected by the vicinity of the new parts; and that the energy, entropy, volume, etc., of the system in its varied state are correctly represented by the sums of the energies, entropies, volumes, etc., of the various parts (original and new), so far at least as any of these quantities are determined or affected by the formation of the new parts. We will suppose  $D\varepsilon$ ,  $D\eta$ ,  $Dv$ ,  $Dm_1$ ,  $Dm_2$  . . .  $Dm_n$  to be so defined that these conditions shall not be violated. This may be done in various ways. We may suppose that the position of the surfaces separating the new and the original parts has been fixed in any suitable way. This will determine the space and the matter belonging to the parts separated. If this does not determine the division of the entropy, we may suppose this determined in any suitable arbitrary way. Thus we may suppose the total energy in and about any new part to be so distributed that equation (12) as applied to the original parts shall not be violated by the formation of the new parts. Or, it may seem more simple to suppose that the imaginary surface which divides any new part from the original is so placed as to include all the matter which is affected by the vicinity of the new formation, so that the part or parts which we regard as original may be left homogeneous in the strictest sense, including uniform *densities of energy and entropy*, up to the very

bounding surface. The homogeneity of the new parts is of no consequence, as we have made no assumption in that respect. It may be doubtful whether we can consider the new parts, *as thus bounded*, to be infinitely small even in their earliest stages of development. But if they are not infinitely small, the only way in which this can affect the validity of our formulae will be that in virtue of the equations of condition, i. e., in virtue of the evident necessities of the case, finite variations of the energy, entropy, volume, etc., of the original parts will be caused, to which it might seem that equation (12) would not apply. But if the nature and state of the mass be not varied, equation (12) will hold true of finite differences. (This appears at once, if we integrate the equation under the above limitation.) Hence, the equation will hold true for finite differences, provided that the nature and state of the mass be infinitely little varied. For the differences may be considered as made up of two parts, of which the first are for a constant nature and state of the mass, and the second are infinitely small. We may therefore regard the new parts to be bounded as supposed without prejudice to the validity of any of our results.

The condition (52) understood in either of these ways (or in others which will suggest themselves to the reader) will have a perfectly definite meaning, and will be valid as the necessary and sufficient condition of equilibrium in regard to the formation of new parts, when the conditions of equilibrium in regard to the original parts, (50), (51), and (43), are satisfied.

In regard to the condition (53), it may be shown that with (50), (51), and (43) it is always sufficient for equilibrium. To prove this, it is only necessary to show that when (50), (51), and (43) are satisfied, and (52) is not, (53) will also not be satisfied.

We will first observe that an expression of the form

$$-\varepsilon + T\eta - Pv + M_1 m_1 + M_2 m_2 \dots + M_n m_n \quad (54)$$

denotes the work obtainable by the formation (by a reversible process) of a body of which  $\varepsilon$ ,  $\eta$ ,  $v$ ,  $m_1$ ,  $m_2$ ,  $\dots$ ,  $m_n$  are the energy, entropy, volume, and the quantities of the components, within a medium having the pressure  $P$ , the temperature  $T$ , and the potentials  $M_1$ ,  $M_2$ ,  $\dots$ ,  $M_n$ . (The medium is supposed so large that its properties are not sensibly altered in any part by the formation of the body.) For  $\varepsilon$  is the energy of the body formed, and the remaining terms represent (as may be seen by applying equation (12) to the medium) the decrease of the energy of the medium, if, after the

formation of the body, the joint entropy of the medium and the body, their joint volumes and joint quantities of matter, were the same as the entropy, etc., of the medium before the formation of the body. This consideration may convince us that for any given finite values of  $v$  and of  $T, P, M_1$ , etc. this expression cannot be infinite when  $\varepsilon, \eta, m_1$ , etc. are determined by any real body, whether homogeneous or not, (but of the given volume), even when  $T, P, M_1$ , etc. do not represent the values of the temperature, pressure, and potentials of any real substance. (If the substances  $S_1, S_2, \dots, S_n$  are all actual components of any homogeneous part of the system of which the equilibrium is discussed, that part will afford an example of a body having the temperature, pressure, and potentials of the medium supposed.)

Now by integrating equation (12) on the supposition that the nature and state of the mass considered remain unchanged, we obtain the equation

$$\varepsilon = t\eta - Pv + \mu_1 m_1 + \mu_2 m_2 + \dots + \mu_n m_n \quad (55)$$

which will hold true of any homogeneous mass whatever. Therefore for any one of the original parts, by (50) and (51),

$$\varepsilon - T\eta + Pv - M_1 m_1 - M_2 m_2 - \dots - M_n m_n = 0. \quad (56)$$

If the condition (52) is not satisfied in regard to all possible new parts, let  $N$  be a new part occurring in an original part  $O$ , for which the condition is not satisfied. It is evident that the value of the expression

$$\varepsilon - T\eta + Pv - M_1 m_1 - M_2 m_2 - \dots - M_n m_n \quad (57)$$

applied to a mass like  $O$  including some very small masses like  $N$ , will be negative, and will decrease if the number of these masses like  $N$  is increased, until there remains within the whole mass no portion of any sensible size without these masses like  $N$ , which, it will be remembered, have no sensible size. But it cannot decrease without limit, as the value of (54) cannot become infinite. Now we need not inquire whether the least value of (57) (for constant values of  $T, P, M_1, M_2, \dots, M_n$ ) would be obtained by excluding entirely the mass like  $O$ , and filling the whole space considered with masses like  $N$ , or whether a certain mixture would give a smaller value,—it is certain that the least possible value of (57) per unit of volume, and that a negative value, will be realized by a mass having a certain homogeneity. If the new part  $N$  for which the condition (52) is not satisfied occurs between two different original parts  $O'$  and  $O''$ , the argument need not be essentially varied. We may consider the

value of (57) for a body consisting of masses like  $O'$  and  $O''$  separated by a lamina  $N$ . This value may be decreased by increasing the extent of this lamina, which may be done within a given volume by giving it a convoluted form; and it will be evident, as before, that the least possible value of (57) will be for a homogeneous mass, and that the value will be negative. And such a mass will be not merely an ideal combination, but a body capable of existing, for as the expression (57) has for this mass in the state considered its least possible value per unit of volume, the energy of the mass included in a unit of volume is the least possible for the same matter with the same entropy and volume,—hence, if confined in a non-conducting vessel, it will be in a state of not unstable equilibrium. Therefore when (50), (51), and (43) are satisfied, if the condition (52) is not satisfied in regard to all possible new parts, there will be some homogeneous body which can be formed out of the substances  $S_1, S_2, \dots, S_n$  which will not satisfy condition (53).

Therefore, if the initially existing masses satisfy the conditions (50), (51), and (43), and condition (53) is satisfied by every homogeneous body which can be formed out of the given matter, there will be equilibrium.

On the other hand, (53) is not a necessary condition of equilibrium. For we may easily conceive that the condition (52) shall hold true (for any very small formations within or between any of the given masses), while the condition (53) is not satisfied (for all large masses formed of the given matter), and experience shows that this is very often the case. Supersaturated solutions, superheated water, etc., are familiar examples. Such an equilibrium will, however, be *practically* unstable. By this is meant that, although, strictly speaking, an infinitely small disturbance or change may not be sufficient to destroy the equilibrium, yet a very small change in the initial state, perhaps a circumstance which entirely escapes our powers of perception, will be sufficient to do so. The presence of a small portion of the substance for which the condition (53) does not hold true, is sufficient to produce this result, when this substance forms a variable component of the original homogeneous masses. In other cases, when, if the new substances are formed at all, different kinds must be formed simultaneously, the initial presence of the different kinds, and that in immediate proximity, may be necessary.

It will be observed, that from (56) and (53) we can at once obtain (50) and (51), viz., by applying (53) to bodies differing infinitely little from the various homogeneous parts of the given mass. There-

fore, the condition (56) (relating to the various homogeneous parts of the given mass) and (53) (relating to any bodies which can be formed of the given matter) with (43) are always sufficient for equilibrium, and always necessary for an equilibrium which shall be practically stable. And, if we choose, we may get rid of limitation in regard to equations (43). For, if we compare these equations with (38), it is easy to see that it is always immaterial, in applying the tests (56) and (53) to any body, how we consider it to be composed. Hence, in applying these tests, we may consider all bodies to be composed of the *ultimate* components of the given mass. Then the terms in (56) and (53) which relate to other components than these will vanish, and we need not regard the equations (43). Such of the constants  $M_1, M_2 \dots M_n$  as relate to the ultimate components, may be regarded, like  $T$  and  $P$ , as unknown quantities subject only to the conditions (56) and (53).

These two conditions, which are sufficient for equilibrium and necessary for a practically stable equilibrium, may be united in one, viz., (if we choose the ultimate components of the given mass for the component substances to which  $m_1, m_2, \dots m_n$  relate) that it shall be possible to give such values to the constants  $T, P, M_1, M_2, \dots M_n$  in the expression (57) that the value of the expression for each of the homogeneous parts of the mass in question shall be as small as for any body whatever made of the same components.

#### *Effect of Solidity of any Part of the given Mass.*

If any of the homogeneous masses of which the equilibrium is in question are solid, it will evidently be proper to treat the proportion of their components as invariable in the application of the criterion of equilibrium, even in the case of *compounds of variable proportions*, i. e., even when bodies can exist which are compounded in proportions infinitesimally varied from those of the solids considered. (Those solids which are capable of absorbing fluids form of course an exception, so far as their fluid components are concerned.) It is true that a solid may be increased by the formation of new solid matter on the surface where it meets a fluid, which is not homogeneous with the previously existing solid, but such a deposit will properly be treated as a distinct part of the system, (viz., as one of the parts which we have called *new*). Yet it is worthy of notice that if a homogeneous solid which is a compound of variable proportions is in contact and equilibrium with a fluid, and the actual components of the solid (considered as of variable composition) are also actual com-

ponents of the fluid, and the condition (53) is satisfied in regard to all bodies which can be formed out of the actual components of the fluid, (which will always be the case unless the fluid is practically unstable,) all the conditions will hold true of the solid, which would be necessary for equilibrium if it were fluid.

This follows directly from the principles stated on the preceding pages. For in this case the value of (57) will be zero as determined either for the solid or for the fluid considered with reference to their ultimate components, and will not be negative for any body whatever which can be formed of these components; and these conditions are sufficient for equilibrium independently of the solidity of one of the masses. Yet the point is perhaps of sufficient importance to demand a more detailed consideration.

Let  $S_a, \dots, S_g$  be the actual components of the solid, and  $S_h, \dots, S_k$  its possible components (which occur as actual components in the fluid); then, considering the proportion of the components of the solid as variable, we shall have for this body by equation (12)

$$\begin{aligned} d\varepsilon' = t d\eta' - p' dv' + \mu_a' dm_a' &\dots + \mu_g' dm_g' \\ &+ \mu_h' dm_h' \dots + \mu_k' dm_k'. \end{aligned} \quad (58)$$

By this equation the potentials  $\mu_a', \dots, \mu_k'$  are perfectly defined. But the differentials  $dm_a', \dots, dm_k'$ , considered as independent, evidently express variations which are not *possible* in the sense required in the criterion of equilibrium. We might, however, introduce them into the general condition of equilibrium, if we should express the dependence between them by the proper equations of condition. But it will be more in accordance with our method hitherto, if we consider the solid to have only a single independently variable component  $S_x$ , of which the nature is represented by the solid itself. We may then write

$$\delta\varepsilon' = t' \delta\eta' - p' \delta v' + \mu_x' \delta m_x'. \quad (59)$$

In regard to the relation of the potential  $\mu_x'$  to the potentials occurring in equation (58) it will be observed, that as we have by integration of (58) and (59)

$$\varepsilon' = t' \eta' - p' v' + \mu_a' m_a' \dots + \mu_g' m_g', \quad (60)$$

and  $\varepsilon' = t' \eta' - p' v' + \mu_x' m_x'; \quad (61)$

therefore  $\mu_x' m_x' = \mu_a' m_a' \dots + \mu_g' m_g'. \quad (62)$

Now, if the fluid has besides  $S_a, \dots, S_g$  and  $S_h, \dots, S_k$  the actual components  $S_l, \dots, S_n$ , we may write for the fluid

$$\delta\varepsilon' = t'' \delta\eta'' - p'' \delta v'' + \mu_a'' \delta m_a'' + \dots + \mu_g'' \delta m_g'' + \mu_h'' \delta m_h'' + \dots + \mu_k'' \delta m_k'' + \mu_l'' \delta m_l'' + \dots + \mu_n'' \delta m_n'', \quad (63)$$

and as by supposition

$$m_x' \mathcal{Z}_x = m_a' \mathcal{Z}_a + \dots + m_g' \mathcal{Z}_g \quad (64)$$

equations (43), (50), and (51) will give in this case on elimination of the constants  $T, P$ , etc.,

$$t' = t'', \quad p' = p'', \quad (65)$$

and

$$m_x' \mu_x' = m_a' \mu_a'' + \dots + m_g' \mu_g''. \quad (66)$$

Equations (65) and (66) may be regarded as expressing the conditions of equilibrium between the solid and the fluid. The last condition may also, in virtue of (62), be expressed by the equation

$$m_a' \mu_a' + \dots + m_g' \mu_g' = m_a' \mu_a'' + \dots + m_g' \mu_g''. \quad (67)$$

But if condition (53) holds true of all bodies which can be formed of  $S_a, \dots, S_g, S_h, \dots, S_k, S_l, \dots, S_n$ , we may write for all such bodies

$$\begin{aligned} \varepsilon - t'' \eta + p'' v - \mu_a'' m_a - \dots - \mu_g'' m_g - \mu_h'' m_h \\ \dots - \mu_k'' m_k - \mu_l'' m_l - \dots - \mu_n'' m_n \geq 0. \end{aligned} \quad (68)$$

(In applying this formula to various bodies, it is to be observed that only the values of the unaccented letters are to be determined by the different bodies to which it is applied, the values of the accented letters being already determined by the given fluid.) Now, by (60), (65), and (67), the value of the first member of this condition is zero when applied to the solid in its given state. As the condition must hold true of a body differing infinitesimally from the solid, we shall have

$$\begin{aligned} d\varepsilon - t'' d\eta + p'' dv - \mu_a'' dm_a' - \dots - \mu_g'' dm_g' \\ - \mu_h'' dm_h' - \dots - \mu_k'' dm_k' \geq 0, \end{aligned} \quad (69)$$

or, by equations (58) and (65),

$$\begin{aligned} (\mu_a' - \mu_a'') dm_a' + (\mu_g' - \mu_g'') dm_g' \\ + (\mu_h' - \mu_h'') dm_h' + \dots + (\mu_k' - \mu_k'') dm_k' \geq 0. \end{aligned} \quad (70)$$

Therefore, as these differentials are all independent,

$$\mu_a' = \mu_a'', \dots, \mu_g' = \mu_g'', \quad \mu_h' \geq \mu_h'', \dots, \mu_k' \geq \mu_k''; \quad (71)$$

which with (65) are evidently the same conditions which we would have obtained if we had neglected the fact of the solidity of one of the masses.

We have supposed the solid to be homogeneous. But it is evident that in any case the above conditions must hold for every separate point where the solid meets the fluid. Hence, the temperature and pressure and the potentials for all the actual components of the solid must have a constant value in the solid at the surface where it meets the fluid. Now, these quantities are determined by the nature and state of the solid, and exceed in number the independent variations of which its nature and state are capable. Hence, if we reject as improbable the supposition that the nature or state of a body can vary without affecting the value of any of these quantities, we may conclude that a solid which varies (continuously) in nature or state at its surface cannot be in equilibrium with a stable fluid which contains, as independently variable components, the variable components of the solid. (There may be, however, in equilibrium with the same stable fluid, a *finite* number of different solid bodies, composed of the variable components of the fluid, and having their nature and state completely determined by the fluid.)\*

#### *Effect of Additional Equations of Condition.*

As the equations of condition, of which we have made use, are such as always apply to matter enclosed in a rigid, impermeable, and non-conducting envelop, the particular conditions of equilibrium which we have found will always be sufficient for equilibrium. But the number of conditions necessary for equilibrium, will be diminished, in a case otherwise the same, as the number of equations of condition is increased. Yet the problem of equilibrium which has been treated will sufficiently indicate the method to be pursued in all cases and the general nature of the results.

It will be observed that the position of the various homogeneous parts of the given mass, which is otherwise immaterial, may determine the existence of certain equations of condition. Thus, when different parts of the system in which a certain substance is a variable component are entirely separated from one another by parts of which this substance is not a component, the quantity of this substance will be invariable for each of the parts of the system which are thus separated, which will be easily expressed by equations of condition. Other equations of condition may arise from the passive forces (or resistances to change) inherent in the given masses. In the prob-

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\* The solid has been considered as subject only to isotropic stresses. The effect of other stresses will be considered hereafter.

lem which we are next to consider there are equations of condition due to a cause of a different nature.

*Effect of a Diaphragm (Equilibrium of Osmotic Forces).*

If the given mass, enclosed as before, is divided into two parts, each of which is homogeneous and fluid, by a diaphragm which is capable of supporting an excess of pressure on either side, and is permeable to some of the components and impermeable to others, we shall have the equations of condition

$$\delta\eta' + \delta\eta'' = 0, \quad (72)$$

$$\delta v' = 0, \quad \delta v'' = 0, \quad (73)$$

and for the components which cannot pass the diaphragm

$$\delta m_a' = 0, \quad \delta m_a'' = 0, \quad \delta m_b' = 0, \quad \delta m_b'' = 0, \quad \text{etc.}, \quad (74)$$

and for those which can

$$\delta m_h' + \delta m_h'' = 0, \quad \delta m_i' + \delta m_i'' = 0, \quad \text{etc.} \quad (75)$$

With these equations of condition, the general condition of equilibrium (see (15)) will give the following particular conditions:

$$t' = t'', \quad (76)$$

and for the components which can pass the diaphragm, if actual components of both masses,

$$\mu_h' = \mu_h'', \quad \mu_i' = \mu_i'', \quad \text{etc.}, \quad (77)$$

but not

$$p' = p'',$$

nor

$$\mu_a' = \mu_a'', \quad \mu_b' = \mu_b'', \quad \text{etc.}$$

Again, if the diaphragm is permeable to the components in certain proportions only, or in proportions not entirely determined yet subject to certain conditions, these conditions may be expressed by equations of condition, which will be linear equations between  $\delta m_1'$ ,  $\delta m_2'$ , etc., and if these be known the deduction of the particular conditions of equilibrium will present no difficulties. We will however observe that if the components  $S_1$ ,  $S_2$ , etc. (being actual components on each side) can pass the diaphragm simultaneously in the proportions  $a_1$ ,  $a_2$ , etc. (without other resistances than such as vanish with the velocity of the current), values proportional to  $a_1$ ,  $a_2$ , etc. are possible for  $\delta m_1'$ ,  $\delta m_2'$ , etc. in the general condition of equilibrium,  $\delta m_1''$ ,  $\delta m_2''$ , etc. having the same values taken negatively, so that we shall have for one particular condition of equilibrium

$$a_1 \mu_1' + a_2 \mu_2' + \text{etc.} = a_1 \mu_1'' + a_2 \mu_2'' + \text{etc.} \quad (78)$$

There will evidently be as many independent equations of this form

as there are independent combinations of the elements which can pass the diaphragm.

These conditions of equilibrium do not of course depend in any way upon the supposition that the volume of each fluid mass is kept constant, if the diaphragm is in any case supposed immovable. In fact, we may easily obtain the same conditions of equilibrium, if we suppose the volumes variable. In this case, as the equilibrium must be preserved by forces acting upon the external surfaces of the fluids, the variation of the energy of the sources of these forces must appear in the general condition of equilibrium, which will be

$$\delta\varepsilon' + \delta\varepsilon'' + P' \delta v' + P'' \delta v'' \geq 0, \quad (79)$$

$P'$  and  $P''$  denoting the external forces per unit of area. (Compare (14).) From this condition we may evidently derive the same internal conditions of equilibrium as before, and in addition the external conditions

$$p' = P', \quad p'' = P''. \quad (80)$$

In the preceding paragraphs it is assumed that the permeability of the diaphragm is perfect, and its impermeability absolute, i. e., that it offers no resistance to the passage of the components of the fluids in certain proportions, except such as vanishes with the velocity, and that in other proportions the components cannot pass at all. How far these conditions are satisfied in any particular case is of course to be determined by experiment.

If the diaphragm is permeable to all the  $n$  components without restriction, the temperature and the potentials for all the components must be the same on both sides. Now, as one may easily convince himself, a mass having  $n$  components is capable of only  $n+1$  independent variations in nature and state. Hence, if the fluid on one side of the diaphragm remains without change, that on the other side cannot (in general) vary in nature or state. Yet the pressure will not necessarily be the same on both sides. For, although the pressure is a function of the temperature and the  $n$  potentials, it may be a many-valued function (or any one of several functions) of these variables. But when the pressures are different on the two sides, the fluid which has the less pressure will be *practically unstable*, in the sense in which the term has been used on page 133. For

$$\varepsilon'' - t'' \eta'' + p'' v'' - \mu_1'' m_1'' - \mu_2'' m_2'' \dots - \mu_n'' m_n'' = 0, \quad (81)$$

as appears from equation (12) if integrated on the supposition that the nature and state of the mass remain unchanged. Therefore, if  $p' < p''$  while  $t' = t''$ ,  $\mu_1' = \mu_1''$ , etc.,

$$\varepsilon'' - t' \eta'' + p' v'' - \mu_1' m_1'' - \mu_2' m_2'' \dots - \mu_n' m_n'' < 0. \quad (82)$$

This relation indicates the instability of the fluid to which the single accents refer. (See page 133.)

But independently of any assumption in regard to the permeability of the diaphragm, the following relation will hold true in any case in which each of the two fluid masses may be regarded as uniform throughout in nature and state. Let the character  $\delta$  be used with the variables which express the nature, state, and quantity of the fluids to denote the increments of the values of these quantities actually occurring in a time either finite or infinitesimal. Then, as the heat received by the two masses cannot exceed  $t' d\eta' + t'' d\eta''$ , and as the increase of their energy is equal to the difference of the heat they receive and the work they do,

$$d\varepsilon' + d\varepsilon'' \leq t' d\eta' + t'' d\eta'' - p' dv' - p'' dv'', \quad (83)$$

i. e., by (12),

$$\mu_1' dm_1' + \mu_1'' dm_1'' + \mu_2' dm_2' + \mu_2'' dm_2'' + \text{etc.} \leq 0, \quad (84)$$

or

$$(\mu_1'' - \mu_1') dm_1'' + (\mu_2'' - \mu_2') dm_2'' + \text{etc.} \leq 0. \quad (85)$$

It is evident that the sign  $\leq$  holds true only in the limiting case in which no motion takes place.

#### DEFINITION AND PROPERTIES OF FUNDAMENTAL EQUATIONS.

The solution of the problems of equilibrium which we have been considering has been made to depend upon the equations which express the relations between the energy, entropy, volume, and the quantities of the various components, for homogeneous combinations of the substances which are found in the given mass. The nature of such equations must be determined by experiment. As, however, it is only *differences* of energy and of entropy that can be measured, or indeed, that have a physical meaning, the values of these quantities are so far arbitrary, that we may choose independently for each simple substance the state in which its energy and its entropy are both zero. The values of the energy and the entropy of any compound body in any particular state will then be fixed. Its energy will be the sum of the work and heat expended in bringing its components from the states in which their energies and their entropies are zero into combination and to the state in question; and its entropy is the value of the integral  $\int \frac{dQ}{t}$  for any *reversible* process

by which that change is effected ( $dQ$  denoting an element of the heat communicated to the matter thus treated, and  $t$  the temperature of the matter receiving it). In the determination both of the energy and of the entropy, it is understood that at the close of the process, all bodies which have been used, other than those to which the determinations relate, have been restored to their original state, with the exception of the sources of the work and heat expended, which must be used only as such sources.

We know, however, *a priori*, that if the quantity of any homogeneous mass containing  $n$  independently variable components varies and not its nature or state, the quantities  $\varepsilon, \eta, v, m_1, m_2, \dots, m_n$  will all vary in the same proportion; therefore it is sufficient if we learn from experiment the relation between all but any one of these quantities for a given constant value of that one. Or, we may consider that we have to learn from experiment the relation subsisting between the  $n+2$  ratios of the  $n+3$  quantities  $\varepsilon, \eta, v, m_1, m_2, \dots, m_n$ . To fix our ideas we may take for these ratios  $\frac{\varepsilon}{v}, \frac{\eta}{v}, \frac{m_1}{v}, \frac{m_2}{v}$ , etc., that is, the separate densities of the components, and the ratios  $\frac{\varepsilon}{v}$  and  $\frac{\eta}{v}$ , which may be called the *densities of energy and entropy*. But when there is but one component, it may be more convenient to choose  $\frac{\varepsilon}{m}, \frac{\eta}{m}, \frac{v}{m}$  as the three variables. In any case, it is only a function of  $n+1$  independent variables, of which the form is to be determined by experiment.

Now if  $\varepsilon$  is a known function of  $\eta, v, m_1, m_2, \dots, m_n$ , as by equation (12)

$$d\varepsilon = t d\eta - p dv + \mu_1 dm_1 + \mu_2 dm_2 \dots + \mu_n dm_n, \quad (86)$$

$t, p, \mu_1, \mu_2, \dots, \mu_n$  are functions of the same variables, which may be derived from the original function by differentiation, and may therefore be considered as known functions. This will make  $n+3$  independent known relations between the  $2n+5$  variables,  $\varepsilon, \eta, v, m_1, m_2, \dots, m_n, t, p, \mu_1, \mu_2, \dots, \mu_n$ . These are all that exist, for of these variables,  $n+2$  are evidently independent. Now upon these relations depend a very large class of the properties of the compound considered,—we may say in general, all its thermal, mechanical, and chemical properties, so far as *active tendencies* are concerned, in cases in which the form of the mass does not require consideration. A single equation from which all these relations may

be deduced we will call a *fundamental equation* for the substance in question. We shall hereafter consider a more general form of the fundamental equation for solids, in which the pressure at any point is not supposed to be the same in all directions. But for masses subject only to isotropic stresses an equation between  $\varepsilon$ ,  $\eta$ ,  $v$ ,  $m_1$ ,  $m_2$ , . . .  $m_n$  is a fundamental equation. There are other equations which possess this same property.\*

Let

$$\psi = \varepsilon - t\eta, \quad (87)$$

then by differentiation and comparison with (86) we obtain

$$d\psi = -\eta dt - p dv + \mu_1 dm_1 + \mu_2 dm_2 \dots + \mu_n dm_n. \quad (88)$$

If, then,  $\psi$  is known as a function of  $t$ ,  $v$ ,  $m_1$ ,  $m_2$ , . . .  $m_n$ , we can find  $\eta$ ,  $p$ ,  $\mu_1$ ,  $\mu_2$ , . . .  $\mu_n$  in terms of the same variables. If we then substitute for  $\psi$  in our original equation its value taken from eq. (87), we shall have again  $n+3$  independent relations between the same  $2n+5$  variables as before.

Let

$$\chi = \varepsilon + p v, \quad (89)$$

then by (86),

$$d\chi = t d\eta + v dp + \mu_1 dm_1 + \mu_2 dm_2 \dots + \mu_n dm_n. \quad (90)$$

If, then,  $\chi$  be known as a function of  $\eta$ ,  $p$ ,  $m_1$ ,  $m_2$ , . . .  $m_n$ , we can find  $t$ ,  $v$ ,  $\mu_1$ ,  $\mu_2$ , . . .  $\mu_n$  in terms of the same variables. By eliminating  $\chi$ , we may obtain again  $n+3$  independent relations between the same  $2n+5$  variables as at first.

Let

$$\zeta = \varepsilon - t\eta + p v, \quad (91)$$

then, by (86)

$$d\zeta = -\eta dt + v dp + \mu_1 dm_1 + \mu_2 dm_2 \dots + \mu_n dm_n. \quad (92)$$

If, then,  $\zeta$  is known as a function of  $t$ ,  $p$ ,  $m_1$ ,  $m_2$ , . . .  $m_n$ , we can

\* M. Massieu (Comptes Rendus, T. Ixix, 1869, p. 858 and p. 1057) has shown how all the properties of a fluid "which are considered in thermodynamics" may be deduced from a single function, which he calls a characteristic function of the fluid considered. In the papers cited, he introduces two different functions of this kind; viz., a function of the temperature and volume, which he denotes by  $\psi$ , the value of which in our notation would be  $\frac{-\varepsilon + t\eta}{t}$  or  $\frac{-\psi}{t}$ ; and a function of the temperature and pressure, which he denotes by  $\psi'$ , the value of which in our notation would be  $\frac{-\varepsilon + t\eta - pv}{t}$  or  $\frac{-\zeta}{t}$ . In both cases he considers a constant quantity (one kilogram) of the fluid, which is regarded as invariable in composition.

find  $\eta$ ,  $v$ ,  $\mu_1$ ,  $\mu_2$ , . . .  $\mu_n$  in terms of the same variables. By eliminating  $\zeta$ , we may obtain again  $n + 3$  independent relations between the same  $2n + 5$  variables as at first.

If we integrate (86), supposing the quantity of the compound substance considered to vary from zero to any finite value, its nature and state remaining unchanged, we obtain

$$\varepsilon = t\eta - pv + \mu_1 m_1 + \mu_2 m_2 \dots + \mu_n m_n, \quad (93)$$

and by (87), (89), (91)

$$\psi = -pv + \mu_1 m_1 + \mu_2 m_2 \dots + \mu_n m_n, \quad (94)$$

$$\chi = t\eta + \mu_1 m_1 + \mu_2 m_2 \dots + \mu_n m_n, \quad (95)$$

$$\zeta = \mu_1 m_1 + \mu_2 m_2 \dots + \mu_n m_n. \quad (96)$$

The last three equations may also be obtained directly by integrating (88), (90), and (92).

If we differentiate (93) in the most general manner, and compare the result with (86), we obtain

$$-vdp + \eta dt + m_1 d\mu_1 + m_2 d\mu_2 \dots + m_n d\mu_n = 0, \quad (97)$$

or

$$dp = \frac{\eta}{v} dt + \frac{m_1}{v} d\mu_1 + \frac{m_2}{v} d\mu_2 \dots + \frac{m_n}{v} d\mu_n = 0. \quad (98)$$

Hence, there is a relation between the  $n + 2$  quantities  $t$ ,  $p$ ,  $\mu_1$ ,  $\mu_2$ , . . .  $\mu_n$ , which, if known, will enable us to find in terms of these quantities all the ratios of the  $n + 2$  quantities  $\eta$ ,  $v$ ,  $m_1$ ,  $m_2$  . . .  $m_n$ . With (93), this will make  $n + 3$  independent relations between the same  $2n + 5$  variables as at first.

Any equation, therefore, between the quantities

$$\varepsilon, \quad \eta, \quad v, \quad m_1, \quad m_2, \dots m_n, \quad (99)$$

$$\text{or } \psi, \quad t, \quad v, \quad m_1, \quad m_2, \dots m_n, \quad (100)$$

$$\text{or } \chi, \quad \eta, \quad p, \quad m_1, \quad m_2, \dots m_n, \quad (101)$$

$$\text{or } \zeta, \quad t, \quad p, \quad m_1, \quad m_2, \dots m_n, \quad (102)$$

$$\text{or } t, \quad p, \quad \mu_1, \quad \mu_2, \dots \mu_n, \quad (103)$$

is a fundamental equation, and any such is entirely equivalent to any other.\* For any homogeneous mass whatever, considered (in general) as variable in composition, in quantity, and in thermodynamic state, and having  $n$  independently variable components, to which

\* The distinction between equations which are, and which are not, *fundamental*, in the sense in which the word is here used, may be illustrated by comparing an equation

the subscript numerals refer, (but not excluding the case in which  $n = 1$  and the composition of the body is invariable,) there is a relation between the quantities enumerated in any one of the above sets, from which, if known, with the aid only of *general* principles and relations, we may deduce all the relations subsisting for such a mass between the quantities  $\varepsilon, \psi, \chi, \zeta, \eta, v, m_1, m_2, \dots, m_n, t, p, \mu_1, \mu_2, \dots, \mu_n$ . It will be observed that, besides the equations which define  $\psi, \chi$ , and  $\zeta$ , there is one finite equation, (93), which subsists between these quantities independently of the form of the fundamental equation.

Other sets of quantities might of course be added which possess the same property. The sets (100), (101), (102) are mentioned on account of the important properties of the quantities  $\psi, \chi, \zeta$ , and because the equations (88), (90), (92), like (86), afford convenient definitions of the potentials, viz.,

$$\mu_1 = \left( \frac{d\varepsilon}{dm_1} \right)_{\eta, v, m} = \left( \frac{d\psi}{dm_1} \right)_{t, v, m} = \left( \frac{d\chi}{dm_1} \right)_{\eta, p, m} = \left( \frac{d\zeta}{dm_1} \right)_{t, p, m} \quad (104)$$

etc., where the subscript letters denote the quantities which remain constant in the differentiation,  $m$  being written for brevity for all the letters  $m_1, m_2, \dots, m_n$  except the one occurring in the denominator. It will be observed that the quantities in (103) are all independent of the quantity of the mass considered, and are those which must, in general, have the same value in contiguous masses in equilibrium.

### *On the quantities $\psi, \chi, \zeta$ .*

The quantity  $\psi$  has been defined for any homogeneous mass by the equation

$$\psi = \varepsilon - t\eta. \quad (105)$$

between

$$\varepsilon, \eta, v, m_1, m_2, \dots, m_n$$

with one between

$$\varepsilon, t, v, m_1, m_2, \dots, m_n.$$

As by (86),

$$t = \left( \frac{d\varepsilon}{d\eta} \right)_{vm},$$

the second equation may evidently be derived from the first. But the first equation cannot be derived from the second; for an equation between

$$\varepsilon, \left( \frac{d\varepsilon}{d\eta} \right)_{vm}, v, m_1, m_2, \dots, m_n$$

is equivalent to one between

$$\left( \frac{d\eta}{d\varepsilon} \right)_{vm}, \varepsilon, v, m_1, m_2, \dots, m_n.$$

which is evidently not sufficient to determine the value of  $\eta$  in terms of the other variables.

We may extend this definition to any material system whatever which has a uniform temperature throughout.

If we compare two states of the system of the same temperature, we have

$$\psi' - \psi'' = \varepsilon' - \varepsilon'' - t(\eta' - \eta''). \quad (106)$$

If we suppose the system brought from the first to the second of these states without change of temperature and by a reversible process in which  $W$  is the work done and  $Q$  the heat received by the system, then

$$\varepsilon' - \varepsilon'' = W - Q, \quad (107)$$

and

$$t(\eta'' - \eta') = Q. \quad (108)$$

Hence

$$\psi' - \psi'' = W; \quad (109)$$

and for an infinitely small reversible change in the state of the system, in which the temperature remains constant, we may write

$$-d\psi = dW. \quad (110)$$

Therefore,  $-\psi$  is the force function of the system for constant temperature, just as  $-\varepsilon$  is the force function for constant entropy. That is, if we consider  $\psi$  as a function of the temperature and the variables which express the distribution of the matter in space, for every different value of the temperature  $-\psi$  is the different force function required by the system if maintained at that special temperature.

From this we may conclude that when a system has a uniform temperature throughout, the additional conditions which are necessary and sufficient for equilibrium may be expressed by

$$(\delta\psi)_t \geq 0.* \quad (111)$$

\* This general condition of equilibrium might be used instead of (2) in such problems of equilibrium as we have considered and others which we shall consider hereafter with evident advantage in respect to the brevity of the formulæ, as the limitation expressed by the subscript  $t$  in (111) applies to every part of the system taken separately, and diminishes by one the number of independent variations in the state of these parts which we have to consider. The more cumbersome course adopted in this paper has been chosen, among other reasons, for the sake of deducing *all* the particular conditions of equilibrium from one general condition, and of having the quantities mentioned in this general condition such as are most generally used and most simply defined; and because in the longer formulæ as given, the reader will easily see in each case the form which they would take if we should adopt (111) as the general condition of equilibrium, which would be in effect to take the thermal condition of equilibrium for granted, and to seek only the remaining conditions. For example, in the problem treated on pages 116 ff., we would obtain from (111) by (88) a condition precisely like (15), except that the terms  $t\delta\eta'$ ,  $t\delta\eta''$ , etc. would be wanting.

When it is not possible to bring the system from one to the other of the states to which  $\psi'$  and  $\psi''$  relate by a reversible process without altering the temperature, it will be observed that it is not necessary for the validity of (107)–(109) that the temperature of the system should remain constant during the reversible process to which  $W$  and  $Q$  relate, provided that the only source of heat or cold used has the same temperature as the system in its initial or final state. Any external bodies may be used in the process in any way not affecting the condition of reversibility, if restored to their original condition at the close of the process; nor does the limitation in regard to the use of heat apply to such heat as may be restored to the source from which it has been taken.

It may be interesting to show directly the equivalence of the conditions (111) and (2) when applied to a system of which the temperature in the given state is uniform throughout.

If there are any variations in the state of such a system which do not satisfy (2), then for these variations

$$\delta\varepsilon < 0 \quad \text{and} \quad \delta\eta = 0.$$

If the temperature of the system in its varied state is not uniform, we may evidently increase its entropy without altering its energy by supposing heat to pass from the warmer to the cooler parts. And the state having the greatest entropy for the energy  $\varepsilon + \delta\varepsilon$  will necessarily be a state of uniform temperature. For this state (regarded as a variation from the original state)

$$\delta\varepsilon < 0 \quad \text{and} \quad \delta\eta > 0.$$

Hence, as we may diminish both the energy and the entropy by cooling the system, there must be a state of uniform temperature for which (regarded as a variation of the original state)

$$\delta\varepsilon < 0 \quad \text{and} \quad \delta\eta = 0.$$

From this we may conclude that for systems of initially uniform temperature condition (2) will not be altered if we limit the variations to such as do not disturb the uniformity of temperature.

Confining our attention, then, to states of uniform temperature, we have by differentiation of (105)

$$\delta\varepsilon - t\delta\eta = \delta\psi + \eta\delta t. \quad (112)$$

Now there are evidently changes in the system (produced by heating or cooling) for which

$$\delta\varepsilon - t\delta\eta = 0 \quad \text{and therefore} \quad \delta\psi + \eta\delta t = 0, \quad (113)$$

neither  $\delta\eta$  nor  $\delta t$  having the value zero. This consideration is sufficient to show that the condition (2) is equivalent to

$$\delta\varepsilon - t \delta\eta \geq 0. \quad (114)$$

and that the condition (111) is equivalent to

$$\delta\psi + \eta \delta t \geq 0 \quad (115)$$

and by (112) the two last conditions are equivalent.

In such cases as we have considered on pages 115–137, in which the form and position of the masses of which the system is composed is immaterial, uniformity of temperature and pressure are always necessary for equilibrium, and the remaining conditions, when these are satisfied, may be conveniently expressed by means of the function  $\zeta$ , which has been defined for a homogeneous mass on page 142, and which we will here define for any mass of uniform temperature and pressure by the same equation

$$\zeta = \varepsilon - t \eta + p v. \quad (116)$$

For such a mass, the condition of (internal) equilibrium is

$$(\delta\zeta)_{t,p} \geq 0. \quad (117)$$

That this condition is equivalent to (2) will easily appear from considerations like those used in respect to (111).

Hence, it is necessary for the equilibrium of two contiguous masses identical in composition that the values of  $\zeta$  as determined for equal quantities of the two masses should be equal. Or, when one of three contiguous masses can be formed out of the other two, it is necessary for equilibrium that the value of  $\zeta$  for any quantity of the first mass should be equal to the sum of the values of  $\zeta$  for such quantities of the second and third masses as together contain the same matter. Thus, for the equilibrium of a solution composed of  $a$  parts of water and  $b$  parts of a salt which is in contact with vapor of water and crystals of the salt, it is necessary that the value of  $\zeta$  for the quantity  $a+b$  of the solution should be equal to the sum of the values of  $\zeta$  for the quantities  $a$  of the vapor and  $b$  of the salt. Similar propositions will hold true in more complicated cases. The reader will easily deduce these conditions from the particular conditions of equilibrium given on page 128.

In like manner we may extend the definition of  $\chi$  to any mass or combination of masses in which the pressure is everywhere the same, using  $\varepsilon$  for the energy and  $v$  for the volume of the whole and setting as before

$$\chi = \varepsilon + p v. \quad (118)$$

If we denote by  $Q$  the heat received by the combined masses from external sources in any process in which the pressure is not varied, and distinguish the initial and final states of the system by accents we have

$$\chi'' - \chi' = \varepsilon'' - \varepsilon' + p(v'' - v') = Q. \quad (119)$$

This function may therefore be called the *heat function for constant pressure* (just as the energy might be called the heat function for constant volume), the diminution of the function representing in all cases in which the pressure is not varied the heat given out by the system. In all cases of chemical action in which no heat is allowed to escape the value of  $\chi$  remains unchanged.

#### POTENTIALS.

In the definition of the potentials  $\mu_1, \mu_2$ , etc., the energy of a homogeneous mass was considered as a function of its entropy, its volume, and the quantities of the various substances composing it. Then the potential for one of these substances was defined as the differential coefficient of the energy taken with respect to the variable expressing the quantity of that substance. Now, as the manner in which we consider the given mass as composed of various substances is in some degree arbitrary, so that the energy may be considered as a function of various different sets of variables expressing quantities of component substances, it might seem that the above definition does not fix the value of the potential of any substance in the given mass, until we have fixed the manner in which the mass is to be considered as composed. For example, if we have a solution obtained by dissolving in water a certain salt containing water of crystallization, we may consider the liquid as composed of  $m_s$  weight-units of the hydrate and  $m_w$  of water, or as composed of  $m_s$  of the anhydrous salt and  $m_w$  of water. It will be observed that the values of  $m_s$  and  $m_w$  are not the same, nor those of  $m_{w'}$  and  $m_w$ , and hence it might seem that the potential for water in the given liquid considered as composed of the hydrate and water, viz.,

$$\left( \frac{d\varepsilon}{dm_w} \right)_{\eta, v, m_s}$$

would be different from the potential for water in the same liquid considered as composed of anhydrous salt and water, viz.,

$$\left( \frac{d\varepsilon}{dm_w} \right)_{\eta, v, m_s}$$

The value of the two expressions is, however, the same, for, although  $m_w$  is not equal to  $m_e$ , we may of course suppose  $dm_w$  to be equal to  $dm_e$ , and then the numerators in the two fractions will also be equal, as they each denote the increase of energy of the liquid, when the quantity  $dm_w$  or  $dm_e$  of water is added without altering the entropy and volume of the liquid. Precisely the same considerations will apply to any other case.

In fact, we may give a definition of a potential which shall not presuppose any choice of a particular set of substances as the components of the homogeneous mass considered.

*Definition.*—If to any homogeneous mass we suppose an infinitesimal quantity of any substance to be added, the mass remaining homogeneous and its entropy and volume remaining unchanged, the increase of the energy of the mass divided by the quantity of the substance added is the *potential* for that substance in the mass considered. (For the purposes of this definition, any chemical element or combination of elements in given proportions may be considered a substance, whether capable or not of existing by itself as a homogeneous body.)

In the above definition we may evidently substitute for entropy, volume, and energy, respectively, either temperature, volume, and the function  $\psi$ ; or entropy, pressure, and the function  $\chi$ ; or temperature, pressure, and the function  $\zeta$ . (Compare equation (104).)

In the same homogeneous mass, therefore, we may distinguish the potentials for an indefinite number of substances, each of which has a perfectly determined value.

Between the potentials for different substances in the same homogeneous mass the same equations will subsist as between the units of these substances. That is, if the substances,  $S_a, S_b$ , etc.,  $S_k, S_l$ , etc., are components of any given homogeneous mass, and are such that

$$\alpha \mathfrak{S}_a + \beta \mathfrak{S}_b + \text{etc.} = \pi \mathfrak{S}_k + \lambda \mathfrak{S}_l + \text{etc.}, \quad (120)$$

$\mathfrak{S}_a, \mathfrak{S}_b$ , etc.,  $\mathfrak{S}_k, \mathfrak{S}_l$ , etc. denoting the units of the several substances, and  $\alpha, \beta$ , etc.,  $\pi, \lambda$ , etc. denoting numbers, then if  $\mu_a, \mu_b$ , etc.,  $\mu_k, \mu_l$ , etc. denote the potentials for these substances in the homogeneous mass,

$$\alpha \mu_a + \beta \mu_b + \text{etc.} = \pi \mu_k + \lambda \mu_l + \text{etc.} \quad (121)$$

To show this, we will suppose the mass considered to be very large. Then, the first member of (121) denotes the increase of the energy of the mass produced by the addition of the matter represented by the first member of (120), and the second member of (121) denotes the

increase of energy of the same mass produced by the addition of the matter represented by the second member of (120), the entropy and volume of the mass remaining in each case unchanged. Therefore, as the two members of (120) represent the same matter in kind and quantity, the two members of (121) must be equal.

But it must be understood that equation (120) is intended to denote equivalence of the substances represented *in the mass considered*, and not merely chemical identity; in other words, it is supposed that there are no passive resistances to change in the mass considered which prevent the substances represented by one member of (120) from passing into those represented by the other. For example, in respect to a mixture of vapor of water and free hydrogen and oxygen (at ordinary temperatures), we may not write

$$9\mathfrak{S}_{Aq} = 1\mathfrak{S}_H + 8\mathfrak{S}_O,$$

but water is to be treated as an independent substance, and no necessary relation will subsist between the potential for water and the potentials for hydrogen and oxygen.

The reader will observe that the relations expressed by equations (43) and (51) (which are essentially relations between the potentials for actual components in different parts of a mass in a state of equilibrium) are simply those which by (121) would necessary subsist between the same potentials in any homogeneous mass containing as variable components all the substances to which the potentials relate.

In the case of a body of invariable composition, the potential for the single component is equal to the value of  $\zeta$  for one unit of the body, as appears from the equation

$$\zeta = \mu_m \tag{122}$$

to which (96) reduces in this case. Therefore, when  $n = 1$ , the fundamental equation between the quantities in the set (102) (see page 143) and that between the quantities in (103) may be derived either from the other by simple substitution. But, with this single exception, an equation between the quantities in one of the sets (99)–(103) cannot be derived from the equation between the quantities in another of these sets without differentiation.

Also in the case of a body of variable composition, when all the quantities of the components except one vanish, the potential for that one will be equal to the value of  $\zeta$  for one unit of the body. We may make this occur for any given composition of the body by

choosing as one of the components the matter constituting the body itself, so that the value of  $\zeta$  for one unit of a body may always be considered as a potential. Hence the relations between the values of  $\zeta$  for contiguous masses given on page 147 may be regarded as relations between potentials.

The two following propositions afford definitions of a potential which may sometimes be convenient.

The potential for any substance in any homogeneous mass is equal to the amount of mechanical work required to bring a unit of the substance by a reversible process from the state in which its energy and entropy are both zero into combination with the homogeneous mass, which at the close of the process must have its original volume, and which is supposed so large as not to be sensibly altered in any part. All other bodies used in the process must by its close be restored to their original state, except those used to supply the work, which must be used only as the source of the work. For, in a reversible process, when the entropies of other bodies are not altered, the entropy of the substance and mass taken together will not be altered. But the original entropy of the substance is zero; therefore the entropy of the mass is not altered by the addition of the substance. Again, the work expended will be equal to the increment of the energy of the mass and substance taken together, and therefore equal, as the original energy of the substance is zero, to the increment of energy of the mass due to the addition of the substance, which by the definition on page 149 is equal to the potential in question.

The potential for any substance in any homogeneous mass is equal to the work required to bring a unit of the substance by a reversible process from a state in which  $\psi = 0$  and the temperature is the same as that of the given mass into combination with this mass, which at the close of the process must have the same volume and temperature as at first, and which is supposed so large as not to be sensibly altered in any part. A source of heat or cold of the temperature of the given mass is allowed, with this exception, other bodies are to be used only on the same conditions as before. This may be shown by applying equation (109) to the mass and substance taken together.

The last proposition enables us to see very easily, how the value of the potential is affected by the arbitrary constants involved in the definition of the energy and the entropy of each elementary substance. For we may imagine the substance brought from the state in which  $\psi = 0$  and the temperature is the same as that of the given

mass, first to any specified state of the same temperature, and then into combination with the given mass. In the first part of the process the work expended is evidently represented by the value of  $\psi$  for the unit of the substance in the state specified. Let this be denoted by  $\psi'$ , and let  $\mu$  denote the potential in question, and  $W$  the work expended in bringing a unit of the substance from the specified state into combination with the given mass as aforesaid; then

$$\mu = \psi' + W. \quad (123)$$

Now as the state of the substance for which  $\varepsilon = 0$  and  $\eta = 0$  is arbitrary, we may simultaneously increase the energies of the unit of the substance in all possible states by any constant  $C$ , and the entropies of the substance in all possible states by any constant  $K$ . The value of  $\psi$ , or  $\varepsilon - t\eta$ , for any state would then be increased by  $C - tK$ ,  $t$  denoting the temperature of the state. Applying this to  $\psi'$  in (123) and observing that the last term in this equation is independent of the values of these constants, we see that the potential would be increased by the same quantity  $C - tK$ ,  $t$  being the temperature of the mass in which the potential is to be determined.

#### ON COEXISTENT PHASES OF MATTER.

In considering the different homogeneous bodies which can be formed out of any set of component substances, it will be convenient to have a term which shall refer solely to the composition and thermodynamic state of any such body without regard to its quantity or form. We may call such bodies as differ in composition or state different *phases* of the matter considered, regarding all bodies which differ only in quantity and form as different examples of the same phase. Phases which can exist together, the dividing surfaces being plane, in an equilibrium which does not depend upon passive resistances to change, we shall call *coexistent*.

If a homogeneous body has  $n$  independently variable components, the phase of the body is evidently capable of  $n + 1$  independent variations. A system of  $r$  coexistent phases, each of which has the same  $n$  independently variable components is capable of  $n + 2 - r$  variations of phase. For the temperature, the pressure, and the potentials for the actual components have the same values in the different phases, and the variations of these quantities are by (97) subject to as many conditions as there are different phases. Therefore, the number of independent variations in the values of these quantities, i. e., the number of independent variations of phase of the system, will be  $n + 2 - r$ .

Or, when the  $r$  bodies considered have not the same independently variable components, if we still denote by  $n$  the number of independently variable components of the  $r$  bodies taken as a whole, the number of independent variations of phase of which the system is capable will still be  $n+2-r$ . In this case, it will be necessary to consider the potentials for more than  $n$  component substances. Let the number of these potentials be  $n+h$ . We shall have by (97), as before,  $r$  relations between the variations of the temperature, of the pressure, and of these  $n+h$  potentials, and we shall also have by (43) and (51)  $h$  relations between these potentials, of the same form as the relations which subsist between the units of the different component substances.

Hence, if  $r=n+2$ , no variation in the phases (remaining coexistent) is possible. It does not seem probable that  $r$  can ever exceed  $n+2$ . An example of  $n=1$  and  $r=3$  is seen in the coexistent solid, liquid, and gaseous forms of any substance of invariable composition. It seems not improbable that in the case of sulphur and some other simple substances there is more than one triad of coexistent phases; but it is entirely improbable that there are four coexistent phases of any simple substance. An example of  $n=2$  and  $r=4$  is seen in a solution of a salt in water in contact with vapor of water and two different kinds of crystals of the salt.

#### *Concerning $n+1$ Coexistent Phases.*

We will now seek the differential equation which expresses the relation between the variations of the temperature and the pressure in a system of  $n+1$  coexistent phases ( $n$  denoting, as before, the number of independently variable components in the system taken as a whole).

In this case we have  $n+1$  equations of the general form of (97) (one for each of the coexistent phases), in which we may distinguish the quantities  $\eta$ ,  $v$ ,  $m_1$ ,  $m_2$ , etc. relating to the different phases by accents. But  $t$  and  $p$  will each have the same value throughout, and the same is true of  $\mu_1$ ,  $\mu_2$ , etc., so far as each of these occurs in the different equations. If the total number of these potentials is  $n+h$ , there will be  $h$  independent relations between them, corresponding to the  $h$  independent relations between the units of the component substances to which the potentials relate, by means of which we may eliminate the variations of  $h$  of the potentials from the equations of the form of (97) in which they occur.

Let one of these equations be

$$v' dp = \eta' dt + m_a' d\mu_a + m_b' d\mu_b + \text{etc.}, \quad (124)$$

and by the proposed elimination let it become

$$v' dp = \eta' dt + A_1' d\mu_1 + A_2' d\mu_2 \dots + A_n' d\mu_n. \quad (125)$$

It will be observed that  $\mu_a$ , for example, in (124) denotes the potential in the mass considered for a substance  $S_a$  which may or may not be identical with any of the substances  $S_1, S_2$ , etc. to which the potentials in (125) relate. Now as the equations between the potentials by means of which the elimination is performed are similar to those which subsist between the units of the corresponding substances, (compare equations (38), (43), and (51).) if we denote these units by  $\mathfrak{S}_a, \mathfrak{S}_b$ , etc.,  $\mathfrak{S}_1, \mathfrak{S}_2$ , etc., we must also have

$$m_a' \mathfrak{S}_a + m_b' \mathfrak{S}_b + \text{etc.} = A_1' \mathfrak{S}_1 + A_2' \mathfrak{S}_2 \dots + A_n' \mathfrak{S}_n. \quad (126)$$

But the first member of this equation denotes (in kind and quantity) the matter in the body to which equations (124) and (125) relate. As the same must be true of the second member, we may regard this same body as composed of the quantity  $A_1'$  of the substance  $S_1$ , with the quantity  $A_2'$  of the substance  $S_2$ , etc. We will therefore, in accordance with our general usage, write  $m_1', m_2'$ , etc. for  $A_1', A_2'$ , etc. in (125), which will then become

$$v' dp = \eta' dt + m_1' d\mu_1 + m_2' d\mu_2 \dots + m_n' d\mu_n. \quad (127)$$

But we must remember that the components to which the  $m_1', m_2'$ , etc. of this equation relate are not necessarily independently variable, as are the components to which the similar expressions in (97) and (124) relate. The rest of the  $n+1$  equations may be reduced to a similar form, viz.,

$$v'' dp = \eta'' dt + m_1'' d\mu_1 + m_2'' d\mu_2 \dots + m_n'' d\mu_n, \quad (128)$$

etc.

By elimination of  $d\mu_1, d\mu_2, \dots, d\mu_n$  from these equations we obtain

$$\begin{vmatrix} v' & m_1' & m_2' & \dots & m_n' \\ v'' & m_1'' & m_2'' & \dots & m_n'' \\ v''' & m_1''' & m_2''' & \dots & m_n''' \\ \cdot & \cdot & \cdot & \dots & \cdot \\ \cdot & \cdot & \cdot & \dots & \cdot \end{vmatrix} dp = \begin{vmatrix} \eta' & m_1' & m_2' & \dots & m_n' \\ \eta'' & m_1'' & m_2'' & \dots & m_n'' \\ \eta''' & m_1''' & m_2''' & \dots & m_n''' \\ \cdot & \cdot & \cdot & \dots & \cdot \\ \cdot & \cdot & \cdot & \dots & \cdot \end{vmatrix} dt. \quad (129)$$

In this equation we may make  $v', v'',$  etc. equal to unity. Then  $m_1', m_2', m_1'',$  etc. will denote the separate densities of the components in the different phases, and  $\eta', \eta'',$  etc. the densities of entropy.

When  $n = 1$ ,

$$(m'' v' - m' v'') dp = (m'' \eta' - m' \eta'') dt, \quad (130)$$

or, if we make  $m' = 1$  and  $m'' = 1$ , we have the usual formula

$$\frac{dp}{dt} = \frac{\eta' - \eta''}{v' - v''} = \frac{Q}{t(v'' - v')}, \quad (131)$$

in which  $Q$  denotes the heat absorbed by a unit of the substance in passing from one state to the other without change of temperature or pressure.

*Concerning Cases in which the Number of Coexistent Phases is less than  $n + 1$ .*

When  $n > 1$ , if the quantities of all the components  $S_1, S_2, \dots, S_n$  are proportional in two coexistent phases, the two equations of the form of (127) and (128) relating to these phases will be sufficient for the elimination of the variations of all the potentials. In fact, the condition of the coexistence of the two phases together with the condition of the equality of the  $n - 1$  ratios of  $m_1', m_2', \dots, m_n'$  with the  $n - 1$  ratios of  $m_1'', m_2'', \dots, m_n''$  is sufficient to determine  $p$  as a function of  $t$  if the fundamental equation is known for each of the phases. The differential equation in this case may be expressed in the form of (130),  $m'$  and  $m''$  denoting either the quantities of any one of the components or the total quantities of matter in the bodies to which they relate. Equation (131) will also hold true in this case, if the total quantity of matter in each of the bodies is unity. But this case differs from the preceding in that the matter which absorbs the heat  $Q$  in passing from one state to another, and to which the other letters in the formula relate, although the same in quantity, is not in general the same in kind at different temperatures and pressures. Yet the case will often occur that one of the phases is essentially invariable in composition, especially when it is a crystalline body, and in this case the matter to which the letters in (131) relate will not vary with the temperature and pressure.

When  $n = 2$ , two coexistent phases are capable, when the temperature is constant, of a single variation in phase. But as (130) will hold true in this case when  $m_1' : m_2' :: m_1'' : m_2''$ , it follows that for constant temperature the pressure is in general a maximum or a minimum when the composition of the two phases is identical. In like manner, the temperature of the two coexistent phases is in general a maximum or a minimum, for constant pressure, when the composition of the two phases is identical. Hence, the series of simultaneous values of  $t$  and  $p$  for which the composition of two coexistent phases

is identical separates those simultaneous values of  $t$  and  $p$  for which no coexistent phases are possible from those for which there are two pair of coexistent phases. This may be applied to a liquid having two independently variable components in connection with the vapor which it yields, or in connection with any solid which may be formed in it.

When  $n = 3$ , we have for three coexistent phases three equations of the form of (127), from which we may obtain the following,

$$\begin{vmatrix} v' & m_1' & m_2' \\ v'' & m_1'' & m_2'' \\ v''' & m_1''' & m_2''' \end{vmatrix} dp = \begin{vmatrix} \eta' & m_1' & m_2' \\ \eta'' & m_1'' & m_2'' \\ \eta''' & m_1''' & m_2''' \end{vmatrix} dt + \begin{vmatrix} m_1' & m_2' & m_3' \\ m_1'' & m_2'' & m_3'' \\ m_1''' & m_2''' & m_3''' \end{vmatrix} d\mu_3. \quad (132)$$

Now the value of the last of these determinants will be zero, when the composition of one of the three phases is such as can be produced by combining the other two. Hence, the pressure of three coexistent phases will in general be a maximum or minimum for constant temperature, and the temperature a maximum or minimum for constant pressure, when the above condition in regard to the composition of the coexistent phases is satisfied. The series of simultaneous values of  $t$  and  $p$  for which the condition is satisfied separates those simultaneous values of  $t$  and  $p$  for which three coexistent phases are not possible, from those for which there are two triads of coexistent phases. These propositions may be extended to higher values of  $n$ , and illustrated by the boiling temperatures and pressures of saturated solutions of  $n - 2$  different solids in solvents having two independently variable components.

#### INTERNAL STABILITY OF HOMOGENEOUS FLUIDS AS INDICATED BY FUNDAMENTAL EQUATIONS.

We will now consider the stability of a fluid enclosed in a rigid envelop which is non-conducting to heat and impermeable to all the components of the fluid. The fluid is supposed initially homogeneous in the sense in which we have before used the word, i. e., uniform in every respect throughout its whole extent. Let  $S_1, S_2, \dots, S_n$  be the *ultimate* components of the fluid; we may then consider every body which can be formed out of the fluid to be composed of  $S_1, S_2, \dots, S_n$ , and that in only one way. Let  $m_1, m_2, \dots, m_n$  denote the quantities of these substances in any such body, and let  $\epsilon, \eta, v$ , denote its energy, entropy, and volume. The fundamental equation for compounds of  $S_1, S_2, \dots, S_n$ , if completely determined, will give us all possible sets of simultaneous values of these variables for homogeneous bodies.

Now, if it is possible to assign such values to the constants  $T$ ,  $P$ ,  $M_1$ ,  $M_2$ , . . .  $M_n$  that the value of the expression

$$\varepsilon = T\eta + Pv - M_1 m_1 - M_2 m_2 - \dots - M_n m_n \quad (133)$$

shall be zero for the given fluid, and shall be positive for every other phase of the same components, i. e., for every homogeneous body\* not identical in nature and state with the given fluid (but composed entirely of  $S_1$ ,  $S_2$ , . . .  $S_n$ ), the condition of the given fluid will be stable.

For, in any condition whatever of the given mass, whether or not homogeneous, or fluid, if the value of the expression (133) is not negative for any homogeneous part of the mass, its value for the whole mass cannot be negative; and if its value cannot be zero for any homogeneous part which is not identical in phase with the mass in its given condition, its value cannot be zero for the whole except when the whole is in the given condition. Therefore, in the case supposed, the value of this expression for any other than the given condition of the mass is positive. (That this conclusion cannot be invalidated by the fact that it is not entirely correct to regard a composite mass as made up of homogeneous parts having the same properties in respect to energy, entropy, etc., as if they were parts of larger homogeneous masses, will easily appear from considerations similar to those adduced on pages 131–133.) If, then, the value of the expression (133) for the mass considered is less when it is in the given condition than when it is in any other, the energy of the mass in its given condition must be less than in any other condition in which it has the same entropy and volume. The given condition is therefore stable. (See page 110.)

Again, if it is possible to assign such values to the constants in (133) that the value of the expression shall be zero for the given fluid mass, and shall not be negative for any phase of the same components, the given condition will be evidently not unstable. (See page 110.) It will be stable unless it is possible for the given matter in the given volume and with the given entropy to consist of homogeneous parts for all of which the value of the expression (133) is zero, but which are not all identical in phase with the mass in its given condition. (A mass consisting of such parts would be in equilibrium, as we have already seen on pages 133, 134.) In this case, if we disregard the quantities connected with the surfaces which divide the

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\* A vacuum is throughout this discussion to be regarded as a limiting case of an extremely rarified body. We may thus avoid the necessity of the specific mention of a vacuum in propositions of this kind.

homogeneous parts, we must regard the given condition as one of neutral equilibrium. But in regard to these homogeneous parts, which we may evidently consider to be all different phases, the following conditions must be satisfied. (The accents distinguish the letters referring to the different parts, and the unaccented letters refer to the whole mass.)

$$\left. \begin{array}{l} \eta' + \eta'' + \text{etc.} = \eta, \\ v' + v'' + \text{etc.} = v, \\ m_1' + m_1'' + \text{etc.} = m_1, \\ m_2' + m_2'' + \text{etc.} = m_2, \\ \text{etc.} \end{array} \right\}. \quad (134)$$

Now the values of  $\eta$ ,  $v$ ,  $m_1$ ,  $m_2$ , etc. are determined by the whole fluid mass in its given state, and the values of  $\frac{\eta'}{v}$ ,  $\frac{\eta''}{v''}$ , etc.,  $\frac{m_1'}{v'}$ ,  $\frac{m_1''}{v''}$ , etc.,  $\frac{m_2'}{v'}$ ,  $\frac{m_2''}{v''}$ , etc., etc., etc., are determined by the phases of the various parts. But the phases of these parts are evidently determined by the phase of the fluid as given. They form, in fact, the whole set of coexistent phases of which the latter is one. Hence, we may regard (134) as  $n+2$  linear equations between  $v'$ ,  $v''$ , etc. (The values of  $v'$ ,  $v''$ , etc. are also subject to the condition that none of them can be negative.) Now one solution of these equations must give us the given condition of the fluid; and it is not to be expected that they will be capable of any other solution, unless the number of different homogeneous parts, that is, the number of different coexistent phases, is greater than  $n+2$ . We have already seen (page 153) that it is not probable that this is ever the case.

We may, however, remark that in a certain sense an infinitely large fluid mass will be in neutral equilibrium in regard to the formation of the substances, if such there are, other than the given fluid, for which the value of (133) is zero (when the constants are so determined that the value of the expression is zero for the given fluid, and not negative for any substance); for the tendency of such a formation to be reabsorbed will diminish indefinitely as the mass out of which it is formed increases.

When the substances  $S_1$ ,  $S_2$ , . . .  $S_n$  are all independently variable components of the given mass, it is evident from (86) that the conditions that the value of (133) shall be zero for the mass as given, and shall not be negative for any phase of the same components, can only be fulfilled when the constants  $T$ ,  $P$ ,  $M_1$ ,  $M_2$ , . . .  $M_n$  are equal to the temperature, the pressure, and the several potentials in the given

mass. If we give these values to the constants, the expression (133) will necessarily have the value zero for the given mass and we shall only have to inquire whether its value is positive for all other phases. But when  $S_1, S_2, \dots, S_n$  are not all independently variable components of the given mass, the values which it will be necessary to give to the constants in (133) cannot be determined entirely from the properties of the given mass; but  $T$  and  $P$  must be equal to its temperature and pressure, and it will be easy to obtain as many equations connecting  $M_1, M_2, \dots, M_n$  with the potentials in the given mass as it contains independently variable components.

When it is not possible to assign such values to the constants in (133) that the value of the expression shall be zero for the given fluid, and either zero or positive for any phase of the same components, we have already seen (pages 129–134) that if equilibrium subsists without passive resistances to change, it must be in virtue of properties which are peculiar to small masses surrounded by masses of different nature, and which are not indicated by fundamental equations. In this case, the fluid will necessarily be unstable, if we extend this term to embrace all cases in which an initial disturbance confined to a small part of an indefinitely large fluid mass will cause an ultimate change of state not indefinitely small in degree throughout the whole mass. In the discussion of stability as indicated by fundamental equations it will be convenient to use the term in this sense.\*

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\* If we wish to know the stability of the given fluid when exposed to a constant temperature, or to a constant pressure, or to both, we have only to suppose that there is enclosed in the same envelop with the given fluid another body (which cannot combine with the fluid) of which the fundamental equation is  $\epsilon = T\eta$ , or  $\epsilon = -Pv$ , or  $\epsilon = T\eta - Pv$ , as the case may be, ( $T$  and  $P$  denoting the constant temperature and pressure, which of course must be those of the given fluid,) and to apply the criteria of page 110 to the whole system. When it is possible to assign such values to the constants in (133) that the value of the expression shall be zero for the given fluid and positive for every other phase of the same components, the value of (133) for the whole system will be less when the system is in its given condition than when it is in any other. (Changes of form and position of the given fluid are of course regarded as immaterial.) Hence the fluid is stable. When it is not possible to assign such values to the constants that the value of (133) shall be zero for the given fluid and zero or positive for any other phase, the fluid is of course unstable. In the remaining case, when it is possible to assign such values to the constants that the value of (133) shall be zero for the given fluid and zero or positive for every other phase, but not without the value zero for some other phase, the state of equilibrium of the fluid as stable or neutral will be determined by the possibility of satisfying, for any other than the given condition of the fluid, equations like (134), in which, however, the first or the second or both are to be stricken out, according as we are considering the

In determining for any given positive values of  $T$  and  $P$  and any given values whatever of  $M_1, M_2, \dots, M_n$  whether the expression (133) is capable of a negative value for any phase of the components  $S_1, S_2, \dots, S_n$ , and if not, whether it is capable of the value zero for any other phase than that of which the stability is in question, it is only necessary to consider phases having the temperature  $T$  and pressure  $P$ . For we may assume that a mass of matter represented by any values of  $m_1, m_2, \dots, m_n$  is capable of at least one state of not unstable equilibrium (which may or may not be a homogeneous state) at this temperature and pressure. It may easily be shown that for such a state the value of  $\varepsilon - T\eta + Pv$  must be as small as for any other state of the same matter. The same will therefore be true of the value of (133). Therefore if this expression is capable of a negative value for any mass whatever, it will have a negative value for that mass at the temperature  $T$  and pressure  $P$ . And if this mass is not homogeneous, the value of (133) must be negative for at least one of its homogeneous parts. So also, if the expression (133) is not capable of a negative value for any phase of the components, any phase for which it has the value zero must have the temperature  $T$  and the pressure  $P$ .

It may easily be shown that the same must be true in the limiting cases in which  $T=0$  and  $P=0$ . For negative values of  $P$ , (133) is always capable of negative values, as its value for a vacuum is  $Pv$ .

For any body of the temperature  $T$  and pressure  $P$ , the expression (133) may by (91) be reduced to the form

$$\zeta - M_1 m_1 - M_2 m_2 \dots - M_n m_n. \quad (135)$$

We have already seen (pages 131, 132) that an expression like (133), when  $T, P, M_1, M_2, \dots, M_n$  and  $v$  have any given finite values, cannot have an infinite negative value as applied to any real body. Hence, in determining whether (133) is capable of a negative value for any phase of the components  $S_1, S_2, \dots, S_n$ , and if not, whether it is capable of the value zero for any other phase than that of which the stability is in question, we have only to consider the least value of which it is capable for a constant value of  $v$ . Any body giving this value must satisfy the condition that for constant volume

$$d\varepsilon - T d\eta - M_1 dm_1 - M_2 dm_2 \dots - M_n dm_n \geq 0, \quad (136)$$

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stability of the fluid for constant temperature, or for constant pressure, or for both. The number of coexistent phases will sometimes exceed by one or two the number of the remaining equations, and then the equilibrium of the fluid will be neutral in respect to one or two independent changes.

or, if we substitute the value of  $d\varepsilon$  taken from equation (86), using subscript  $a \dots g$  for the quantities relating to the actual components of the body, and subscript  $h \dots k$  for those relating to the possible,

$$t d\eta + \mu_a dm_a \dots + \mu_g dm_g + \mu_h dm_h \dots + \mu_k dm_k - T d\eta - M_1 dm_1 - M_2 dm_2 \dots - M_n dm_n \geq 0. \quad (137)$$

That is, the temperature of the body must be equal to  $T$ , and the potentials of its components must satisfy the same conditions as if it were in contact and in equilibrium with a body having potentials  $M_1, M_2, \dots, M_n$ . Therefore the same relations must subsist between  $\mu_a \dots \mu_g$ , and  $M_1 \dots M_n$  as between the units of the corresponding substances, so that

$$m_a \mu_a \dots + m_g \mu_g = m_1 M_1 \dots + m_n M_n; \quad (138)$$

and as we have by (93)

$$\varepsilon = t \eta - p v + \mu_a m_a \dots + \mu_g m_g, \quad (139)$$

the expression (133) will reduce (for the body or bodies for which it has the least value per unit of volume) to

$$(P - p) v, \quad (140)$$

the value of which will be positive, null, or negative, according as the value of

$$P - p \quad (141)$$

is positive, null, or negative.

Hence, the conditions in regard to the stability of a fluid of which all the ultimate components are independently variable admit a very simple expression. If the pressure of the fluid is greater than that of any other phase of the same components which has the same temperature and the same values of the potentials for its actual components, the fluid is stable without coexistent phases; if its pressure is not as great as that of some other such phase, it will be unstable; if its pressure is as great as that of any other such phase, but not greater than that of every other, the fluid will certainly not be unstable, and in all probability it will be stable (when enclosed in a rigid envelop which is impermeable to heat and to all kinds of matter), but it will be one of a set of coexistent phases of which the others are the phases which have the same pressure.

The considerations of the last two pages, by which the tests relating to the stability of a fluid are simplified, apply to such bodies as actually exist. But if we should form arbitrarily any equation as a fundamental equation, and ask whether a fluid of which the proper-

ties were given by that equation would be stable, the tests of stability last given would be insufficient, as some of our assumptions might not be fulfilled by the equation. The test, however, as first given (pages 156–159) would in all cases be sufficient.

*Stability in respect to Continuous Changes of Phase.*

In considering the changes which may take place in any mass, we have already had occasion to distinguish between infinitesimal changes in existing phases, and the formation of entirely new phases. A phase of a fluid may be stable in regard to the former kind of change, and unstable in regard to the latter. In this case it may be capable of continued existence in virtue of properties which prevent the commencement of discontinuous changes. But a phase which is unstable in regard to continuous changes is evidently incapable of permanent existence on a large scale except in consequence of passive resistances to change. We will now consider the conditions of stability in respect to continuous changes of phase, or, as it may also be called, stability in respect to adjacent phases. We may use the same general test as before, except that the expression (133) is to be applied only to phases which differ infinitely little from the phase of which the stability is in question. In this case the component substances to be considered will be limited to the independently variable components of the fluid, and the constants  $M_1, M_2$ , etc. must have the values of the potentials for these components in the given fluid. The constants in (133) are thus entirely determined and the value of the expression for the given phase is necessarily zero. If for any infinitely small variation of the phase, the value of (133) can become negative, the fluid will be unstable; but if for every infinitely small variation of the phase the value of (133) becomes positive, the fluid will be stable. The only remaining case, in which the phase can be varied without altering the value of (133) can hardly be expected to occur. The phase concerned would in such a case have coexistent adjacent phases.

This condition, which for brevity's sake we will call the condition of stability, may be written in the form

$$\epsilon'' - t' \eta'' + p' v'' - \mu_1' m_1'' - \dots - \mu_n' m_n'' > 0, \quad (142)$$

in which the quantities relating to the phase of which the stability is in question are distinguished by single accents, and those relating to

the other phase by double accents. This condition is by (93) equivalent to

$$\begin{aligned} \varepsilon'' - t' \eta'' + p' v'' - \mu_1' m_1'' \dots - \mu_n' m_n'' \\ - \varepsilon' + t' \eta' - p' v' + \mu_1' m_1' \dots + \mu_n' m_n' > 0, \end{aligned} \quad (143)$$

and to

$$\begin{aligned} -t' \eta'' + p' v'' - \mu_1' m_1'' \dots - \mu_n' m_n'' \\ + t'' \eta'' - p'' v'' + \mu_1'' m_1'' \dots + \mu_n'' m_n'' > 0. \end{aligned} \quad (144)$$

The condition (143) may be expressed more briefly in the form

$$\Delta \varepsilon > t \Delta \eta - p \Delta v + \mu_1 \Delta m_1 \dots + \mu_n \Delta m_n, \quad (145)$$

if we use the character  $\Delta$  to signify that the condition, although relating to infinitesimal differences, is not to be interpreted in accordance with the usual convention in respect to differential equations with neglect of infinitesimals of higher orders than the first, but is to be interpreted *strictly*, like an equation between finite differences. In fact, when a condition like (145) (interpreted strictly) is satisfied for infinitesimal differences, it must be possible to assign limits within which it shall hold true of finite differences. But it is to be remembered that the condition is not to be applied to any arbitrary values of  $\Delta \eta$ ,  $\Delta v$ ,  $\Delta m_1$ ,  $\dots$ ,  $\Delta m_n$ , but only to such as are determined by a change of phase. (If only the quantity of the body which determines the value of the variables should vary and not its phase, the value of the first member of (145) would evidently be zero.) We may free ourselves from this limitation by making  $v$  constant, which will cause the term  $-p \Delta v$  to disappear. If we then divide by the constant  $v$ , the condition will become

$$\Delta \frac{\varepsilon}{v} > t \Delta \frac{\eta}{v} + \mu_1 \Delta \frac{m_1}{v} \dots + \mu_n \Delta \frac{m_n}{v}, \quad (146)$$

in which form it will not be necessary to regard  $v$  as constant. As we may obtain from (86)

$$\frac{d \frac{\varepsilon}{v}}{v} = t d \frac{\eta}{v} + \mu_1 d \frac{m_1}{v} \dots + \mu_n d \frac{m_n}{v}, \quad (147)$$

we see that *the stability of any phase in regard to continuous changes depends upon the same conditions in regard to the second and higher differential coefficients of the density of energy regarded as a function of the density of entropy and the densities of the several components, which would make the density of energy a minimum, if the necessary conditions in regard to the first differential coefficients were fulfilled.* When  $n=1$ , it may be more convenient to regard  $m$  as constant

in (145) than  $v$ . Regarding  $m$  a constant, it appears that the stability of a phase depends upon the same conditions in regard to the second and higher differential coefficients of the energy of a unit of mass regarded as a function of its entropy and volume, which would make the energy a minimum, if the necessary conditions in regard to the first differential coefficients were fulfilled.

The formula (144) expresses the condition of stability for the phase to which  $t'$ ,  $p'$ , etc. relate. But it is evidently the necessary and sufficient condition of the stability of all phases of certain kinds of matter, or of all phases within given limits, that (144) shall hold true of any two infinitesimally differing phases within the same limits, or, as the case may be, in general. For the purpose, therefore, of such *collective* determinations of stability, we may neglect the distinction between the two states compared, and write the condition in the form

$$-\eta \frac{\partial t}{\partial p} + v \frac{\partial p}{\partial t} - m_1 \frac{\partial \mu_1}{\partial t} \dots - m_n \frac{\partial \mu_n}{\partial t} > 0, \quad (148)$$

or

$$\frac{\partial p}{\partial t} > \frac{\eta}{v} \frac{\partial t}{\partial p} + \frac{m_1}{v} \frac{\partial \mu_1}{\partial t} \dots + \frac{m_n}{v} \frac{\partial \mu_n}{\partial t}. \quad (149)$$

Comparing (98), we see that it is necessary and sufficient for the stability in regard to continuous changes of all the phases within any given limits, that within those limits the same conditions should be fulfilled in respect to the second and higher differential coefficients of the pressure regarded as a function of the temperature and the several potentials, which would make the pressure a minimum, if the necessary conditions with respect to the first differential coefficients were fulfilled.

By equations (87) and (94), the condition (142) may be brought to the form

$$\begin{aligned} & \psi'' + t'' \eta'' + p' v'' - \mu_1' m_1'' \dots - \mu_n' m_n'' \\ & - \psi' - t' \eta'' - p' v' + \mu_1' m_1' \dots + \mu_n' m_n' > 0. \end{aligned} \quad (150)$$

For the stability of all phases within any given limits it is necessary and sufficient that within the same limits this condition shall hold true of any two phases which differ infinitely little. This evidently requires that when  $v' = v'', m_1' = m_1'', \dots, m_n' = m_n''$ ,

$$\psi'' - \psi' + (t'' - t') \eta'' > 0; \quad (151)$$

and that when  $t' = t''$

$$\begin{aligned} & \psi'' + p' v'' - \mu_1' m_1'' \dots + \mu_n' m_n'' \\ & - \psi' - p' v' + \mu_1' m_1' \dots + \mu_n' m_n' > 0. \end{aligned} \quad (152)$$

These conditions may be written in the form

$$[\Delta\psi + \eta \Delta t]_{v,m} < 0, \quad (153)$$

$$[\Delta\psi + p \Delta v - \mu_1 \Delta m_1 - \dots - \mu_n \Delta m_n]_t > 0, \quad (154)$$

in which the subscript letters indicate the quantities which are to be regarded as constant,  $m$  standing for all the quantities  $m_1 \dots m_n$ . If these conditions hold true within any given limits, (150) will also hold true of any two infinitesimally differing phases within the same limits. To prove this, we will consider a third phase, determined by the equations

$$t''' = t', \quad (155)$$

and

$$v''' = v'', \quad m_1''' = m_1'', \quad \dots \quad m_n''' = m_n''. \quad (156)$$

Now by (153),

$$\psi''' - \psi'' + (t''' - t'') \eta'' < 0; \quad (157)$$

and by (154),

$$\begin{aligned} \psi''' + p' v''' - \mu_1' m_1''' - \dots - \mu_n' m_n''' \\ - \psi' - p' v' + \mu_1' m_1' - \dots + \mu_n' m_n' > 0. \end{aligned} \quad (158)$$

Hence,

$$\begin{aligned} \psi'' + t'' \eta'' + p' v'' - \mu_1' m_1'' - \dots - \mu_n' m_n'' \\ - \psi' - t''' \eta'' - p' v' + \mu_1' m_1' - \dots + \mu_n' m_n' > 0, \end{aligned} \quad (159)$$

which by (155) and (156) is equivalent to (150). Therefore, the conditions (153) and (154) in respect to the phases within any given limits are necessary and sufficient for the stability of all the phases within those limits. It will be observed that in (153) we have the condition of thermal stability of a body considered as unchangeable in composition and in volume, and in (154), the condition of mechanical and chemical stability of the body considered as maintained at a constant temperature. Comparing equation (88), we see that the condition (153) will be satisfied, if  $\frac{d^2\psi}{dt^2} < 0$ , i. e., if  $\frac{d\eta}{dt}$  or  $t \frac{d\eta}{dt}$  (the specific heat for constant volume) is positive. When  $n = 1$ , i. e., when the composition of the body is invariable, the condition (154) will evidently not be altered, if we regard  $m$  as constant, by which the condition will be reduced to

$$[\Delta\psi + p \Delta v]_{t,m} > 0. \quad (160)$$

This condition will evidently be satisfied if  $\frac{d^2\psi}{dv^2} > 0$ , i. e., if  $-\frac{dp}{dv}$  or  $-v \frac{dp}{dv}$  (the elasticity for constant temperature) is positive. But when  $n > 1$ , (154) may be abbreviated more symmetrically by making  $v$  constant.

Again, by (91) and (96), the condition (142) may be brought to the form

$$\begin{aligned} \zeta'' + t'' \eta'' - p'' v'' - \mu_1' m_1'' - \dots - \mu_n' m_n'' \\ - \zeta' - t' \eta' + p' v' + \mu_1' m_1' - \dots + \mu_n' m_n' > 0. \end{aligned} \quad (161)$$

Therefore, for the stability of all phases within any given limits it is necessary and sufficient that within the same limits

$$[\mathcal{A}\zeta + \eta \mathcal{A}t - v \mathcal{A}p]_m < 0, \quad (162)$$

and

$$[\mathcal{A}\zeta - \mu_1 \mathcal{A}m_1 - \dots - \mu_n \mathcal{A}m_n]_{t,p} > 0, \quad (163)$$

as may easily be proved by the method used with (153) and (154). The first of these formulæ expresses the thermal and mechanical conditions of stability for a body considered as unchangeable in composition, and the second the conditions of chemical stability for a body considered as maintained at a constant temperature and pressure. If  $\eta = 1$ , the second condition falls away, and as in this case  $\zeta = m\mu$ , condition (162) becomes identical with (148).

The foregoing discussion will serve to illustrate the relation of the general condition of stability in regard to continuous changes to some of the principal forms of fundamental equations. It is evident that each of the conditions (146), (149), (154), (162), (163) involve in general several particular conditions of stability. We will now give our attention to the latter. Let

$$\Phi = \varepsilon - t' \eta + p' v - \mu_1' m_1 - \dots - \mu_n' m_n, \quad (164)$$

the accented letters referring to one phase and the unaccented to another. It is by (142) the necessary and sufficient condition of the stability of the first phase that, for constant values of the quantities relating to that phase and of  $v$ , the value of  $\Phi$  shall be a minimum when the second phase is identical with the first. Differentiating (164), we have by (86)

$$\begin{aligned} d\Phi = (t - t') d\eta - (p - p') dv + (\mu_1 - \mu_1') dm_1 \\ \dots - (\mu_n - \mu_n') dm_n. \end{aligned} \quad (165)$$

Therefore, the above condition requires that if we regard  $v, m_1, \dots, m_n$  as having the constant values indicated by accenting these letters,  $t$  shall be an increasing function of  $\eta$ , when the variable phase differs sufficiently little from the fixed. But as the fixed phase may be any one within the limits of stability,  $t$  must be an increasing function of  $\eta$  (within these limits) for any constant values of  $v, m_1, \dots, m_n$ . This condition may be written

$$\left( \frac{\mathcal{A}t}{\mathcal{A}\eta} \right)_{v, m_1, \dots, m_n} > 0. \quad (166)$$

When this condition is satisfied, the value of  $\Phi$ , for any given values of  $v, m_1, \dots, m_n$ , will be a minimum when  $t = t'$ . And therefore, in applying the general condition of stability relating to the value of  $\Phi$ , we need only consider the phases for which  $t = t'$ .

We see again by (165) that the general condition requires that if we regard  $t, v, m_2, \dots, m_n$  as having the constant values indicated by accenting these letters,  $\mu_1$  shall be an increasing function of  $m_1$ , when the variable phase differs sufficiently little from the fixed. But as the fixed phase may be any one within the limits of stability,  $\mu_1$  must be an increasing function of  $m_1$  (within these limits) for any constant values of  $t, v, m_2, \dots, m_n$ . That is,

$$\left(\frac{d\mu_1}{dm_1}\right)_{t, v, m_2, \dots, m_n} > 0. \quad (167)$$

When this condition is satisfied, as well as (166),  $\Phi$  will have a minimum value, for any constant values of  $v, m_2, \dots, m_n$ , when  $t = t'$  and  $\mu_1 = \mu_1'$ ; so that in applying the general condition of stability we need only consider the phases for which  $t = t'$  and  $\mu_1 = \mu_1'$ .

In this way we may also obtain the following particular conditions of stability:

$$\left(\frac{d\mu_2}{dm_2}\right)_{t, v, m_1, m_3, \dots, m_n} > 0, \quad (168)$$

$$\cdot \quad \cdot \quad \cdot \quad \cdot \quad \cdot \quad \left(\frac{d\mu_n}{dm_n}\right)_{t, v, m_1, \dots, m_{n-1}} > 0. \quad (169)$$

When the  $n+1$  conditions (166)–(169) are all satisfied, the value of  $\Phi$ , for any constant value of  $v$ , will be a minimum when the temperature and the potentials of the variable phase are equal to those of the fixed. The pressures will then also be equal and the phases will be entirely identical. Hence, the general condition of stability will be completely satisfied, when the above particular conditions are satisfied.

From the manner in which these particular conditions have been derived, it is evident that we may interchange in them  $\eta, m_1, \dots, m_n$  in any way, provided that we also interchange in the same way  $t, \mu_1, \dots, \mu_n$ . In this way we may obtain different sets of  $n+1$  conditions which are necessary and sufficient for stability. The quantity  $v$  might be included in the first of these lists, and  $-p$  in the second, except in cases when, in some of the phases considered, the entropy or the quantity of one of the components has the value zero.

Then the condition that that quantity shall be constant would create a restriction upon the variations of the phase, and cannot be substituted for the condition that the volume shall be constant in the statement of the general condition of stability relative to the minimum value of  $\Phi$ .

To indicate more distinctly all these particular conditions at once, we observe that the condition (144), and therefore also the condition obtained by interchanging the single and double accents, must hold true of any two infinitesimally differing phases within the limits of stability. Combining these two conditions we have

$$(t'' - t')(\eta'' - \eta') - (p'' - p')(v'' - v') + (\mu_1'' - \mu_1') (m_1'' - m_1') \dots (\mu_n'' - \mu_n') (m_n'' - m_n') > 0, \quad (170)$$

which may be written more briefly

$$\Delta t \Delta \eta - \Delta p \Delta v + \Delta \mu_1 \Delta m_1 \dots + \Delta \mu_n \Delta m_n > 0. \quad (171)$$

This must hold true of any two infinitesimally differing phases within the limits of stability. If, then, we give the value zero to one of the differences in every term except one, but not so as to make the phases completely identical, the values of the two differences in the remaining term will have the same sign, except in the case of  $\Delta p$  and  $\Delta v$ , which will have opposite signs. (If both states are stable this will hold true even on the limits of stability.) Therefore, within the limits of stability, either of the two quantities occurring (after the sign  $\Delta$ ) in any term of (171) in an increasing function of the other, —except  $p$  and  $v$ , of which the opposite is true,—when we regard as constant one of the quantities occurring in each of the other terms, but not such as to make the phases identical.

If we write  $d$  for  $\Delta$  in (166)–(169), we obtain conditions which are always sufficient for stability. If we also substitute  $\geq$  for  $>$ , we obtain conditions which are necessary for stability. Let us consider the form which these conditions will take when  $\eta, v, m_1, \dots, m_n$  are regarded as independent variables. When  $dv = 0$ , we shall have

$$\left. \begin{aligned} dt &= \frac{dt}{d\eta} d\eta + \frac{dt}{dm_1} dm_1 \dots + \frac{dt}{dm_n} dm_n \\ d\mu_1 &= \frac{d\mu_1}{d\eta} d\eta + \frac{d\mu_1}{dm_1} dm_1 \dots + \frac{d\mu_1}{dm_n} dm_n \\ &\vdots \\ d\mu_n &= \frac{d\mu_n}{d\eta} d\eta + \frac{d\mu_n}{dm_1} dm_1 \dots + \frac{d\mu_n}{dm_n} dm_n \end{aligned} \right\} \quad (172)$$

Let us write  $R_{n+1}$  for the determinant of the order  $n+1$

$$\left| \begin{array}{ccc} \frac{d^2\varepsilon}{d\eta^2} & \frac{d^2\varepsilon}{dm_1 d\eta} & \cdots & \frac{d^2\varepsilon}{dm_n d\eta} \\ \frac{d^2\varepsilon}{d\eta dm_1} & \frac{d^2\varepsilon}{dm_1^2} & \cdots & \frac{d^2\varepsilon}{dm_n dm_1} \\ \cdot & \cdot & \cdots & \cdot \\ \frac{d^2\varepsilon}{d\eta dm_n} & \frac{d^2\varepsilon}{dm_1 dm_n} & \cdots & \frac{d^2\varepsilon}{dm_n^2} \end{array} \right| \quad (173)$$

of which the constituents are by (86) the same as the coefficients in equations (172), and  $R_n$ ,  $R_{n-1}$ , etc. for the minors obtained by erasing the last column and row in the original determinant and in the minors successively obtained, and  $R_1$  for the last remaining constituent. Then if  $dt$ ,  $d\mu_1$ ,  $\dots$ ,  $d\mu_{n-1}$ , and  $dv$  all have the value zero, we have by (172)

$$R_n d\mu_n = R_{n+1} dm_n, \quad (174)$$

that is,

$$\left( \frac{d\mu_n}{dm_n} \right)_{t, v, \mu_1, \dots, \mu_{n-1}} = \frac{R_{n+1}}{R_n}. \quad (175)$$

In like manner we obtain

$$\left( \frac{d\mu_{n-1}}{dm_{n-1}} \right)_{t, v, \mu_1, \dots, \mu_{n-2}, m_n} = \frac{R_n}{R_{n-1}}, \left. \right\} \quad (176)$$

Therefore, the conditions obtained by writing  $d$  for  $\mathcal{A}$  in (166)–(169) are equivalent to this, that the determinant given above with the  $n$  minors obtained from it as above mentioned and the last remaining constituent  $\frac{d^2\varepsilon}{d\eta^2}$  shall all be positive. Any phase for which this condition is satisfied will be stable, and no phase will be stable for which any of these quantities has a negative value. But the conditions (166)–(169) will remain valid, if we interchange in any way  $\eta, m_1, \dots, m_n$  (with corresponding interchange of  $t, \mu_1, \dots, \mu_n$ ). Hence the order in which we erase successive columns with the corresponding rows in the determinant is immaterial. Therefore none of the minors of the determinant (173) which are formed by erasing corresponding rows and columns, and none of the constituents of the principal diagonal, can be negative for a stable phase.

We will now consider the conditions which characterize the *limits of stability* (i. e., the limits which divide stable from unstable phases)

with respect to continuous changes.\* Here, evidently, one of the conditions (166)–(169) must cease to hold true. Therefore, one of the differential coefficients formed by changing  $\mathcal{A}$  into  $d$  in the first members of these conditions must have the value zero. (That it is the numerator and not the denominator in the differential coefficient which vanishes at the limit appears from the consideration that the denominator is in each case the differential of a quantity which is necessarily capable of progressive variation, so long at least as the phase is capable of variation at all under the conditions expressed by the subscript letters.) The same will hold true of the set of differential coefficients obtained from these by interchanging in any way  $\eta, m_1, \dots, m_n$ , and simultaneously interchanging  $t, \mu_1, \dots, \mu_n$  in the same way. But we may obtain a more definite result than this.

Let us give to  $\eta$  or  $t$ , to  $m_1$  or  $\mu_1$ , . . . to  $m_{n-1}$  or  $\mu_{n-1}$ , and to  $v$ , the constant values indicated by these letters when accented. Then by (165)

$$d\Phi = (\mu_n - \mu'_n) dm_n. \quad (177)$$

Now

$$\mu_n - \mu'_n = \left( \frac{d\mu_n}{dm_n} \right)' (m_n - m'_n) \quad (178)$$

approximately, the differential coefficient being interpreted in accordance with the above assignment of constant values to certain variables, and its value being determined for the phase to which the accented letters refer. Therefore,

$$d\Phi = \left( \frac{d\mu_n}{dm_n} \right)' (m_n - m'_n) dm_n, \quad (179)$$

and

$$\Phi = \frac{1}{2} \left( \frac{d\mu_n}{dm_n} \right)' (m_n - m'_n)^2. \quad (180)$$

The quantities neglected in the last equation are evidently of the same order as  $(m_n - m'_n)^3$ . Now this value of  $\Phi$  will of course be different (the differential coefficient having a different meaning) according as we have made  $\eta$  or  $t$  constant, and according as we have made  $m_1$  or  $\mu_1$  constant, etc.; but since, within the limits of stability, the value of  $\Phi$ , for any constant values of  $m_n$  and  $v$ , will be the least when  $t, p, \mu_1 \dots \mu_{n-1}$  have the values indicated by accenting these letters, the value of the differential coefficient will be at least as small

\* The limits of stability with respect to discontinuous changes are formed by phases which are coexistent with other phases. Some of the properties of such phases have already been considered. See pages 152–156.

when we give these variables these constant values, as when we adopt any other of the suppositions mentioned above in regard to the quantities remaining constant. And in all these relations we may interchange in any way  $\eta, m_1, \dots, m_n$ , if we interchange in the same way  $t, \mu_1, \dots, \mu_n$ . It follows that, within the limits of stability, when we choose for any one of the differential coefficients

$$\frac{dt}{d\eta}, \frac{d\mu_1}{dm_1}, \dots, \frac{d\mu_n}{dm_n} \quad (181)$$

the quantities following the sign  $d$  in the numerators of the others together with  $v$  as those which are to remain constant in differentiation, the value of the differential coefficient as thus determined will be at least as small as when one or more of the constants in differentiation are taken from the denominators, one being still taken from each fraction, and  $v$  as before being constant.

Now we have seen that none of these differential coefficients, as determined in any of these ways, can have a negative value within the limit of stability, and that some of them must have the value zero at that limit. Therefore, in virtue of the relations just established, one at least of these differential coefficients determined by considering constant the quantities occurring in the numerators of the others together with  $v$ , will have the value zero. But if one such has the value zero, all such will in general have the same value. For if

$$\left( \frac{d\mu_n}{dm_n} \right)_{t, v, \mu_1, \dots, \mu_{n-1}} \quad (182)$$

for example, has the value zero, we may change the density of the component  $S_n$  without altering (if we disregard infinitesimals of higher orders than the first) the temperature or the potentials, and therefore, by (98), without altering the pressure. That is, we may change the phase without altering any of the quantities  $t, p, \mu_1, \dots, \mu_n$ . (In other words, the phases adjacent to the limits of stability exhibit *approximately* the relations characteristic of neutral equilibrium.) Now this change of phase, which changes the density of one of the components, will in general change the density of the others and the density of entropy. Therefore, all the other differential coefficients formed after the analogy of (182), i. e., formed from the fractions in (181) by taking as constants for each the quantities in the numerators of the others together with  $v$ , will in general have the value zero at the limit of stability. And the relation which characterizes the limit of stability may be expressed, in general, by setting any one of these differential coefficients equal to zero. Such

an equation, when the fundamental equation is known, may be reduced to the form of an equation between the independent variables of the fundamental equation.

Again, as the determinant (173) is equal to the product of the differential coefficients obtained by writing  $d$  for  $\Delta$  in the first members of (166)–(169), the equation of the limit of stability may be expressed by setting this determinant equal to zero. The form of the differential equation as thus expressed will not be altered by the interchange of the expressions  $\eta, m_1, \dots, m_n$ , but it will be altered by the substitution of  $v$  for any one of these expressions, which will be allowable whenever the quantity for which it is substituted has not the value zero in any of the phases to which the formula is to be applied.

The condition formed by setting the expression (182) equal to zero is evidently equivalent to this, that

$$\left[ \frac{d\mu_n}{d\frac{m_n}{v}} \right]_{t, \mu_1, \dots, \mu_{n-1}} = 0, \quad (183)$$

that is, that

$$\left[ \frac{d\frac{m_n}{v}}{d\mu_n} \right]_{t, \mu_1, \dots, \mu_{n-1}} = \infty, \quad (184)$$

or by (98), if we regard  $t, \mu_1, \dots, \mu_n$  as the independent variables,

$$\left( \frac{d^2 p}{d\mu_n^2} \right) = \infty. \quad (185)$$

In like manner we may obtain

$$\frac{d^2 p}{dt^2} = \infty, \quad \frac{d^2 p}{d\mu_1^2} = \infty, \quad \dots, \quad \frac{d^2 p}{d\mu_{n-1}^2} = \infty. \quad (186)$$

Any one of these equations, (185), (186), may be regarded, in general, as the equation of the limit of stability. We may be certain that at every phase at that limit one at least of these equations will hold true.

#### GEOMETRICAL ILLUSTRATIONS.

*Surfaces in which the Composition of the Body represented is Constant.*

In vol. ii, p. 382, of the Trans. Conn. Acad., a method is described of representing the thermodynamic properties of substances of invariable composition by means of surfaces. The volume, entropy, and energy

of a constant quantity of a substance are represented by rectangular co-ordinates. This method corresponds to the first kind of fundamental equation described on pages 140–144. Any other kind of fundamental equation for a substance of invariable composition will suggest an analogous geometrical method. Thus, if we make  $m$  constant, the variables in any one of the sets (99)–(103) are reduced to three, which may be represented by rectangular co-ordinates. This will, however, afford but four different methods, for, as has already (page 150) been observed, the two last sets are essentially equivalent when  $n = 1$ .

The method described in the preceding volume has certain advantages, especially for the purposes of theoretical discussion, but it may often be more advantageous to select a method in which the properties represented by two of the co-ordinates shall be such as best serve to identify and describe the different states of the substance. This condition is satisfied by temperature and pressure as well, perhaps, as by any other properties. We may represent these by two of the co-ordinates and the potential by the third. (See page 143.) It will not be overlooked that there is the closest analogy between these three quantities in respect to their parts in the general theory of equilibrium. (A similar analogy exists between volume, entropy, and energy.) If we give  $m$  the constant value unity, the third co-ordinate will also represent  $\zeta$ , which then becomes equal to  $\mu$ .

Comparing the two methods, we observe that in one

$$v = x, \quad \eta = y, \quad \varepsilon = z, \quad (187)$$

$$p = -\frac{dz}{dx}, \quad t = \frac{dz}{dy}, \quad \mu = \zeta = z - \frac{dz}{dx}x - \frac{dz}{dy}y; \quad (188)$$

and in the other

$$t = x, \quad p = y, \quad \mu = \zeta = z, \quad (189)$$

$$\eta = -\frac{dz}{dx}, \quad v = \frac{dz}{dy}, \quad \varepsilon = z - \frac{dz}{dx}x - \frac{dz}{dy}y. \quad (190)$$

Now  $\frac{dz}{dx}$  and  $\frac{dz}{dy}$  are evidently determined by the inclination of the tangent plane, and  $z - \frac{dz}{dx}x - \frac{dz}{dy}y$  is the segment which it cuts off on the axis of  $Z$ . The two methods, therefore, have this reciprocal relation, that the quantities represented in one by the position of a point in a surface are represented in the other by the position of a tangent plane.

The surfaces defined by equations (187) and (189) may be distinguished as the  $v\text{-}\eta\text{-}\epsilon$  surface, and the  $t\text{-}p\text{-}\zeta$  surface, of the substance to which they relate.

In the  $t\text{-}p\text{-}\zeta$  surface a line in which one part of the surface cuts another represents a series of pairs of coexistent states. A point through which pass three different parts of the surface represents a triad of coexistent states. Through such a point will evidently pass the three lines formed by the intersection of these sheets taken two by two. The perpendicular projection of these lines upon the  $p\text{-}t$  plane will give the curves which have recently been discussed by Professor J. Thomson.\* These curves divide the space about the projection of the triple point into six parts which may be distinguished as follows: Let  $\zeta^{(r)}$ ,  $\zeta^{(L)}$ ,  $\zeta^{(S)}$  denote the three ordinates determined for the same values of  $p$  and  $t$  by the three sheets passing through the triple point, then in one of the six spaces

$$\zeta^{(r)} < \zeta^{(L)} < \zeta^{(S)}, \quad (191)$$

in the next space, separated from the former by the line for which  $\zeta^{(L)} = \zeta^{(S)}$ ,

$$\zeta^{(r)} < \zeta^{(S)} < \zeta^{(L)}, \quad (192)$$

in the third space, separated from the last by the line for which  $\zeta^{(r)} = \zeta^{(S)}$ ,

$$\zeta^{(S)} < \zeta^{(r)} < \zeta^{(L)}, \quad (193)$$

in the fourth

$$\zeta^{(S)} < \zeta^{(L)} < \zeta^{(r)}, \quad (194)$$

in the fifth

$$\zeta^{(L)} < \zeta^{(S)} < \zeta^{(V)}, \quad (195)$$

in the sixth

$$\zeta^{(L)} < \zeta^{(V)} < \zeta^{(S)}. \quad (196)$$

The sheet which gives the least values of  $\zeta$  is in each case that which represents the stable states of the substance. From this it is evident that in passing around the projection of the triple point we pass through lines representing alternately coexistent stable and coexistent unstable states. But the states represented by the intermediate values of  $\zeta$  may be called stable *relatively* to the states represented by the highest. The differences  $\zeta^{(L)} - \zeta^{(r)}$ , etc. represent the amount of work obtained in bringing the substance by a reversible process from one to the other of the states to which these quantities relate, in a medium having the temperature and pressure common to the two states. To illustrate such a process, we may suppose a plane perpendicular to the axis of temperature to pass through the points

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\* See the Reports of the British Association for 1871 and 1872; and Philosophical Magazine, vol. xlvi. (1874), p. 447.

representing the two states. This will in general cut the double line formed by the two sheets to which the symbols ( $L$ ) and ( $V$ ) refer. The intersections of the plane with the two sheets will connect the double point thus determined with the points representing the initial and final states of the process, and thus form a *reversible path* for the body between those states.

The geometrical relations which indicate the stability of any state may be easily obtained by applying the principles stated on pp. 156 ff. to the case in which there is but a single component. The expression (133) as a test of stability will reduce to

$$\varepsilon - t' \eta + p' v - \mu' m, \quad (197)$$

the accented letters referring to the state of which the stability is in question, and the unaccented letters to any other state. If we consider the quantity of matter in each state to be unity, this expression may be reduced by equations (91) and (96) to the form

$$\zeta - \zeta' + (t - t') \eta - (p - p') v, \quad (198)$$

which evidently denotes the distance of the point  $(t', p', \zeta')$  below the tangent plane for the point  $(t, p, \zeta)$ , measured parallel to the axis of  $\zeta$ . Hence if the tangent plane for every other state passes above the point representing any given state, the latter will be stable. If any of the tangent planes pass below the point representing the given state, that state will be unstable. Yet it is not always necessary to consider these tangent planes. For, as has been observed on page 160, we may assume that (in the case of any real substance) there will be at least one not unstable state for any given temperature and pressure, except when the latter is negative. Therefore the state represented by a point in the surface on the positive side of the plane  $p = 0$  will be unstable only when there is a point in the surface for which  $t$  and  $p$  have the same values and  $\zeta$  a less value. It follows from what has been stated, that where the surface is doubly convex upwards (in the direction in which  $\zeta$  is measured) the states represented will be stable in respect to adjacent states. This also appears directly from (162). But where the surface is concave upwards in either of its principal curvatures the states represented will be unstable in respect to adjacent states.

When the number of component substances is greater than unity, it is not possible to represent the fundamental equation by a single surface. We have therefore to consider how it may be represented by an infinite number of surfaces. A natural extension of either of the methods already described will give us a series of surfaces in

which every one is the  $v$ - $\eta$ - $\varepsilon$  surface, or every one the  $t$ - $p$ - $\zeta$  surface for a body of constant composition, the proportion of the components varying as we pass from one surface to another. But for a simultaneous view of the properties which are exhibited by compounds of two or three components without change of temperature or pressure, we may more advantageously make one or both of the quantities  $t$  or  $p$  constant in each surface.

*Surfaces and Curves in which the Composition of the Body represented is Variable and its Temperature and Pressure are Constant.*

When there are three components, the position of a point in the  $X$ - $Y$  plane may indicate the composition of a body most simply, perhaps, as follows. The body is supposed to be composed of the quantities  $m_1, m_2, m_3$  of the substances  $S_1, S_2, S_3$ , the value of  $m_1 + m_2 + m_3$  being unity. Let  $P_1, P_2, P_3$  be any three points in the plane, which are not in the same straight line. If we suppose masses equal to  $m_1, m_2, m_3$  to be placed at these three points, the center of gravity of these masses will determine a point which will indicate the value of these quantities. If the triangle is equiangular and has the height unity, the distances of the point from the three sides will be equal numerically to  $m_1, m_2, m_3$ . Now if for every possible phase of the components, of a given temperature and pressure, we lay off from the point in the  $X$ - $Y$  plane which represents the composition of the phase a distance measured parallel to the axis of  $Z$  and representing the value of  $\zeta$  (when  $m_1 + m_2 + m_3 = 1$ ), the points thus determined will form a surface, which may be designated as the  $m_1-m_2-m_3-\zeta$  surface of the substances considered, or simply as their  $m$ - $\zeta$  surface, for the given temperature and pressure. In like manner, when there are but two component substances, we may obtain a curve, which we will suppose in the  $X$ - $Z$  plane. The coordinate  $y$  may then represent temperature or pressure. But we will limit ourselves to the consideration of the properties of the  $m$ - $\zeta$  surface for  $n=3$ , or the  $m$ - $\zeta$  curve for  $n=2$ , regarded as a surface, or curve, which varies with the temperature and pressure.

As by (96) and (92)

$$\zeta = \mu_1 m_1 + \mu_2 m_2 + \mu_3 m_3,$$

and (for constant temperature and pressure)

$$d\zeta = \mu_1 dm_1 + \mu_2 dm_2 + \mu_3 dm_3,$$

if we imagine a tangent plane for the point to which these letters relate, and denote by  $\zeta'$  the ordinate for any point in the plane,

and by  $m_1'$ ,  $m_2'$ ,  $m_3'$ , the distances of the foot of this ordinate from the three sides of the triangle  $P_1 P_2 P_3$ , we may easily obtain

$$\xi' = \mu_1 m_1' + \mu_2 m_2' + \mu_3 m_3', \quad (199)$$

which we may regard as the equation of the tangent plane. Therefore the ordinates for this plane at  $P_1$ ,  $P_2$ , and  $P_3$  are equal respectively to the potentials  $\mu_1$ ,  $\mu_2$ , and  $\mu_3$ . And in general, the ordinate for any point in the tangent plane is equal to the potential (in the phase represented by the point of contact) for a substance of which the composition is indicated by the position of the ordinate. (See page 149.) Among the bodies which may be formed of  $S_1$ ,  $S_2$ , and  $S_3$ , there may be some which are incapable of variation in composition, or which are capable only of a single kind of variation. These will be represented by single points and curves in vertical planes. Of the tangent plane to one of these curves only a single line will be fixed, which will determine a series of potentials of which only two will be independent. The phase represented by a separate point will determine only a single potential, viz., the potential for the substance of the body itself, which will be equal to  $\xi$ .

The points representing a set of coexistent phases have in general a common tangent plane. But when one of these points is situated on the edge where a sheet of the surface terminates, it is sufficient if the plane is tangent to the edge and passes below the surface. Or, when the point is at the end of a separate line belonging to the surface, or at an angle in the edge of a sheet, it is sufficient if the plane pass through the point and below the line or sheet. If no part of the surface lies below the tangent plane, the points where it meets the plane will represent a stable (or at least not unstable) set of coexistent phases.

The surface which we have considered represents the relation between  $\xi$  and  $m_1$ ,  $m_2$ ,  $m_3$  for homogeneous bodies when  $t$  and  $p$  have any constant values and  $m_1 + m_2 + m_3 = 1$ . It will often be useful to consider the surface which represents the relation between the same variables for bodies which consist of parts in different but coexistent phases. We may suppose that these are stable, at least in regard to adjacent phases, as otherwise the case would be devoid of interest. The point which represents the state of the composite body will evidently be at the center of gravity of masses equal to the parts of the body placed at the points representing the phases of these parts. Hence from the surface representing the properties of homogeneous bodies, which may be called the primitive surface, we

may easily construct the surface representing the properties of bodies which are in equilibrium but not homogeneous. This may be called the secondary or derived surface. It will consist, in general, of various portions or sheets. The sheets which represent a combination of two phases may be formed by rolling a double tangent plane upon the primitive surface: the part of the envelop of its successive positions which lies between the curves traced by the points of contact will belong to the derived surface. When the primitive surface has a triple tangent plane or one of higher order, the triangle in the tangent plane formed by joining the points of contact, or the smallest polygon without re-entrant angles which includes all the points of contact, will belong to the derived surface, and will represent masses consisting in general of three or more phases.

Of the whole thermodynamic surface as thus constructed for any temperature and any positive pressure, that part is especially important which gives the least value of  $\zeta$  for any given values of  $m_1, m_2, m_3$ . The state of a mass represented by a point in this part of the surface is one in which no dissipation of energy would be possible if the mass were enclosed in a rigid envelop impermeable both to matter and to heat; and the state of any mass composed of  $S_1, S_2, S_3$  in any proportions, in which the dissipation of energy has been completed, so far as internal processes are concerned, (i. e., under the limitations imposed by such an envelop as above supposed,) would be represented by a point in the part which we are considering of the  $m\text{-}\zeta$  surface for the temperature and pressure of the mass. We may therefore briefly distinguish this part of the surface as the *surface of dissipated energy*. It is evident that it forms a continuous sheet, the projection of which upon the  $X\text{-}Y$  plane coincides with the triangle  $P_1 P_2 P_3$ , (except when the pressure for which the  $m\text{-}\zeta$  surface is constructed is negative, in which case there is no surface of dissipated energy,) that it nowhere has any convexity upward, and that the states which it represents are in no case unstable.

The general properties of the  $m\text{-}\zeta$  lines for two component substances are so similar as not to require separate consideration. We now proceed to illustrate the use of both the surfaces and the lines by the discussion of several particular cases.

Three coexistent phases of two component substances may be represented by the points A, B, and C, in figure 1, in which  $\zeta$  is measured toward the top of the page from  $P_1 P_2$ ,  $m_1$  toward the left from  $P_2 Q_2$ , and  $m_2$  toward the right from  $P_1 Q_1$ . It is supposed that  $P_1 P_2 = 1$ . Portions of the curves to which these points belong

are seen in the figure, and will be denoted by the symbols (A), (B), (C). We may, for convenience, speak of these as separate curves, without implying anything in regard to their possible continuity in parts of the diagram remote from their common tangent AC. The *line of dissipated energy* includes the straight line AC and portions of the primitive curves (A) and (C). Let us first consider how the

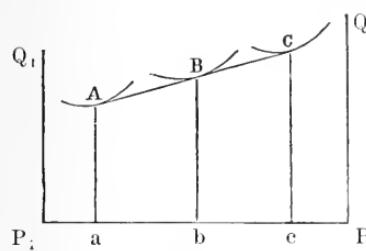


FIG. 1.

diagram will be altered, if the temperature is varied while the pressure remains constant. If the temperature receives the increment  $dt$ , an ordinate of which the position is fixed will receive the increment  $\left(\frac{d\zeta}{dt}\right)_{p, m} dt$ , or

$$P_2 - \eta dt. \quad (\text{The reader will easily convince himself that this is true of the ordinates for the secondary line } AC, \text{ as well as of the ordinates for the primitive curves.})$$

Now if we denote by  $\eta'$  the entropy of the phase represented by the point B considered as belonging to the curve (B), and by  $\eta''$  the entropy of the composite state of the same matter represented by the point B considered as belonging to the tangent to the curves (A) and (C),  $t(\eta' - \eta'')$  will denote the heat yielded by a unit of matter in passing from the first to the second of these states. If this quantity is positive, an elevation of temperature will evidently cause a part of the curve (B) to protrude below the tangent to (A) and (C), which will no longer form a part of the line of dissipated energy. This line will then include portions of the three curves (A), (B), and (C), and of the tangents to (A) and (B) and to (B) and (C). On the other hand, a lowering of the temperature will cause the curve (B) to lie entirely above the tangent to (A) and (C), so that all the phases of the sort represented by (B) will be unstable. If  $t(\eta' - \eta'')$  is negative, these effects will be produced by the opposite changes of temperature.

The effect of a change of pressure while the temperature remains constant may be found in a manner entirely analogous. The variation of any ordinate will be  $\left(\frac{d\zeta}{dp}\right)_{t, m} dp$  or  $v dp$ . Therefore, if the volume of the homogeneous phase represented by the point B is a greater than the volume of the same matter divided between the the phases represented by A and C, an increase of pressure will give diagram indicating that all phases of the sort represented by curve (B) are unstable, and a decrease of pressure will give a diagram indi-

eating two stable pairs of coexistent phases, in each of which one of the phases is of the sort represented by the curve (B). When the relation of the volumes is the reverse of that supposed, these results will be produced by the opposite changes of pressure.

When we have four coexistent phases of three component substances, there are two cases which must be distinguished. In the first, one of the points of contact of the primitive surface with the quadruple tangent plane lies within the triangle formed by joining the other three; in the second, the four points may be joined so as to form a quadrilateral without re-entrant angles. Figure 2 represents the projection upon the  $X-Y$  plane (in which  $m_1, m_2, m_3$  are measured) of a part of the surface of dissipated energy, when one of the points of contact D falls within the triangle formed by the other three A, B, C. This surface includes the triangle ABC in the quadruple tangent plane, portions of the three sheets of the primitive surface which touch the triangle at its vertices, EAF, GBH, ICK, and portions of the three developable surfaces formed by a tangent plane rolling upon each pair of these sheets. These developable surfaces are repre-

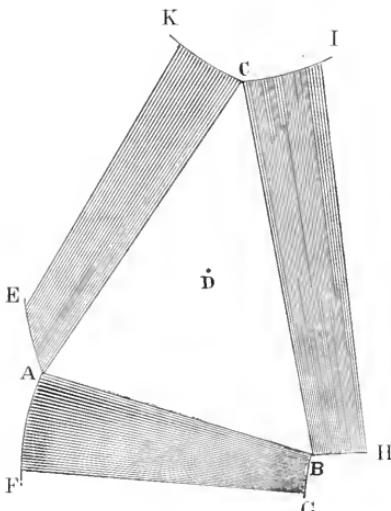


FIG. 2.

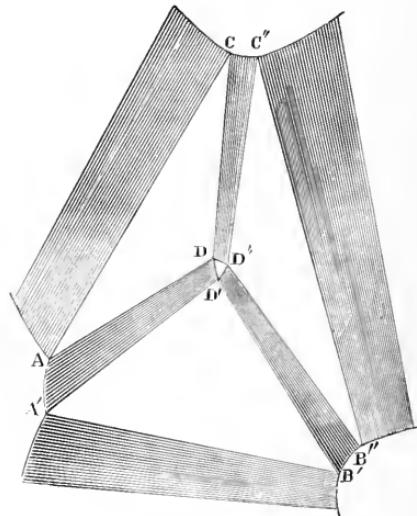


FIG. 3.

sented in the figure by ruled surfaces, the lines indicating the direction of their rectilinear elements. A point within the triangle ABC represents a mass of which the matter is divided, in general, between three or four different phases, in a manner not entirely determined by the position of a point. (The quantities of matter in these phases are such that if placed at the corresponding points, A, B, C, D, their center of gravity would be at the point representing the total mass.)

Such a mass, if exposed to constant temperature and pressure, would be in neutral equilibrium. A point in the developable surfaces represents a mass of which the matter is divided between two coexisting phases, which are represented by the extremities of the line in the figure passing through that point. A point in the primitive surface represents of course a homogeneous mass.

To determine the effect of a change of temperature without change of pressure upon the general features of the surface of dissipated energy, we must know whether heat is absorbed or yielded by a mass in passing from the phase represented by the point D *in the primitive surface* to the composite state consisting of the phases A, B, and C which is represented by the same point. If the first is the case, an increase of temperature will cause the sheet (D) (i. e., the sheet of the primitive surface to which the point D belongs) to separate from the plane tangent to the three other sheets, so as to be situated entirely above it, and a decrease of temperature, will cause a part of the sheet (D) to protrude through the plane tangent to the other sheets. These effects will be produced by the opposite changes of temperature, when heat is yielded by a mass passing from the homogeneous to the composite state above mentioned.

In like manner, to determine the effect of a variation of pressure without change of temperature, we must know whether the volume for the homogeneous phase represented by D is greater or less than the volume of the same matter divided between the phases A, B, and C. If the homogeneous phase has the greater volume, an increase of pressure will cause the sheet (D) to separate from the plane tangent to the other sheets, and a diminution of pressure will cause a part of the sheet (D) to protrude below that tangent plane. And these effects will be produced by the opposite changes of pressure, if the homogeneous phase has the less volume. All this appears from precisely the same considerations which were used in the analogous case for two component substances.

Now when the sheet (D) rises above the plane tangent to the other sheets, the general features of the surface of dissipated energy are not altered, except by the disappearance of the point D. But when the sheet (D) protrudes below the plane tangent to the other sheets, the surface of dissipated energy will take the form indicated in figure 3. It will include portions of the four sheets of the primitive surface, portions of the six developable surfaces formed by a double tangent plane rolling upon these sheets taken two by two, and portions of three triple tangent planes for these sheets taken by threes, the sheet (D) being always one of the three.

But when the points of contact with the quadruple tangent plane which represent the four coexistent phases can be joined so as to form a quadrilateral ABCD (fig. 4) without reentrant angles, the surface of dissipated energy will include this plane quadrilateral, portions of the four sheets of the primitive surface which are tangent to it, and portions of the four developable surfaces formed by double

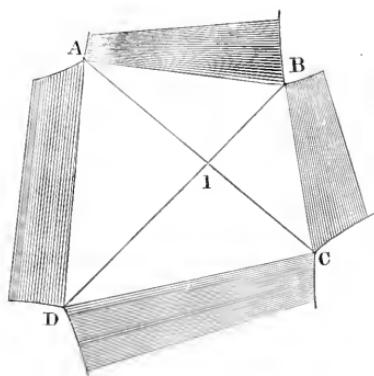


FIG. 4.

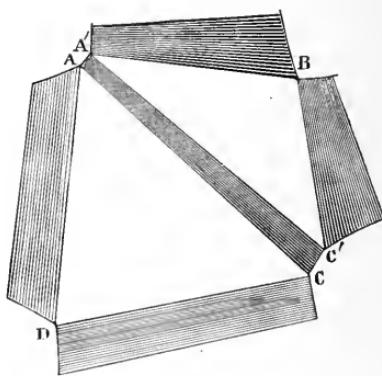


FIG. 5.

tangent planes rolling upon the four pairs of these sheets which correspond to the four sides of the quadrilateral. To determine the general effect of a variation of temperature upon the surface of dissipated energy, let us consider the composite states represented by the point I at the intersection of the diagonals of the quadrilateral. Among these states (which all relate to the same kind and quantity of matter) there is one which is composed of the phases A and C, and another which is composed of the phases B and D. Now if the entropy of the first of these states is greater than that of the second, (i. e., if heat is given out by a body in passing from the first to the second state at constant temperature and pressure,) which we may suppose without loss of generality, an elevation of temperature while the pressure remains constant will cause the triple tangent planes to (B), (D), and (A), and to (B), (D), and (C), to rise above the triple tangent planes to (A), (C), and (B), and to (A), (C), and (D), in the vicinity of the point I. The surface of dissipated energy will therefore take the form indicated in figure 5, in which there are two plane triangles and five developable surfaces besides portions of the four primitive sheets. A diminution of temperature will give a different but entirely analogous form to the surface of dissipated energy. The quadrilateral ABCD will in this case break into two triangles along the diameter BD. The effects produced by

variation of the pressure while the temperature remains constant will of course be similar to those described. By considering the difference of volume instead of the difference of entropy of the two states represented by the point I in the quadruple tangent plane, we may distinguish between the effects of increase and diminution of pressure.

It should be observed that the points of contact of the quadruple tangent plane with the primitive surface may be at isolated points or curves belonging to the latter. So also, in the case of two component substances, the points of contact of the triple tangent line may be at isolated points belonging to the primitive curve. Such cases need not be separately treated, as the necessary modifications in the preceding statements, when applied to such cases, are quite evident. And in the remaining discussion of this geometrical method, it will generally be left to the reader to make the necessary limitations or modifications in analogous cases.

The necessary condition in regard to simultaneous variations of temperature and pressure, in order that four coexistent phases of three components, or three coexistent phases of two components, shall remain possible, has already been deduced by purely analytical processes. (See equation (129).)

We will next consider the case of two coexistent phases of identical composition, and first, when the number of components is two. The coexistent phases, if each is variable in composition, will be represented by the point of contact of two curves. One of the curves will in general lie above the other except at the point of contact; therefore, when the temperature and pressure remain constant, one phase cannot be varied in composition without becoming unstable, while the other phase will be stable if the proportion of either component is increased. By varying the temperature or pressure, we may cause the upper curve to protrude below the other, or to rise (relatively) entirely above it. (By comparing the volumes or the entropies of the two coexistent phases, we may easily determine which result would be produced by an increase of temperature or of pressure.) Hence, the temperatures and pressures for which two coexistent phases have the same composition form the limit to the temperatures and pressures for which such coexistent phases are possible. It will be observed that as we pass this limit of temperature and pressure, the pair of coexistent phases does not simply become unstable, like pairs and triads of coexistent phases which we have considered before, but there ceases to be any such pair of coexistent phases. The same result has already been obtained analytically on

page 155. But on that side of the limit on which the coexistent phases are possible, there will be two pairs of coexistent phases for the same values of  $t$  and  $p$ , as seen in figure 6. If the curve AA' represents vapor, and the curve BB' liquid, a liquid (represented by) B may exist in contact with a vapor A, and (at the same temperature and pressure) a liquid B' in contact with a vapor A'. If we compare these phases in respect to their composition, we see that in one case the vapor is richer than the liquid in a certain component, and in the other case poorer. Therefore, if these liquids are made to boil, the effect on their composition will be opposite. If the boiling is continued under constant pressure, the temperature will rise as the liquids approach each other in composition, and the curve BB' will rise *relatively* to the curve AA', until the curves are tangent to each other, when the two liquids become identical in nature, as also the vapors which they yield. In composition, and in the value of  $\zeta$  per unit of mass, the vapor will then agree with the liquid. But if the curve BB' (which has the greater curvature) represents vapor, and AA' represents liquid, the effect of boiling will make the liquids A and A' differ more in composition. In this case, the relations indicated in the figure will hold for a temperature higher than that for which (with the same pressure) the curves are tangent to one another.

When two coexistent phases of three component substances have the same composition, they are represented by the point of contact of two sheets of the primitive surface. If these sheets do not intersect at the point of contact, the case is very similar to that which we have just considered. The upper sheet except at the point of contact represents unstable phases. If the temperature or pressure are so varied that a part of the upper sheet protrudes through the lower, the points of contact of a double tangent plane rolling upon the two sheets will describe a closed curve on each, and the surface of dissipated energy will include a portion of each sheet of the primitive surface united by a ring-shaped developable surface.

If the sheet having the greater curvatures represents liquid, and the other sheet vapor, the boiling temperature for any given pressure will be a maximum, and the pressure of saturated vapor for any given temperature will be a minimum, when the coexistent liquid and vapor have the same composition.

But if the two sheets, constructed for the temperature and pressure of the coexistent phases which have the same composition, intersect

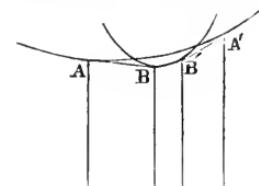


FIG. 6.

at the point of contact, the whole primitive surface as seen from below will in general present four re-entrant furrows, radiating from the point of contact, for each of which a developable surface may be formed by a rolling double tangent plane. The different parts of the surface of dissipated energy in the vicinity of the point of contact are represented in figure 7. ATB, ETF are parts of one sheet of the primitive surface, and CTD, GTH are parts of the other. These are united by the developable surfaces BTC, DTE, FTG, HTA. Now we may make either sheet of the primitive surface sink relatively to the other by the proper variation of temperature or pressure. If the sheet to which ATB, ETF belong is that which sinks relatively, these

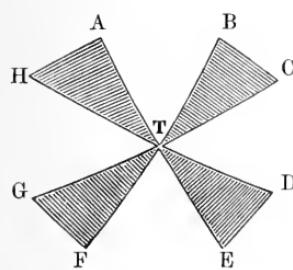


FIG. 7.

parts of the surface of dissipated energy will be merged in one, as well as the developable surfaces BTC, DTE, and also FTG, HTA. (The lines CTD, BTE, ATF, HTG will separate from one another at T, each forming a continuous curve.) But if the sheet of the primitive surface which sinks relatively is that to which CTD and GTH belong, then these parts will be merged in one in the surface of dissipated energy, as will be the developable surfaces BTC, ATH, and also DTE, FTG.

It is evident that this is not a case of maximum or minimum temperature for coexistent phases under constant pressure, or of maximum or minimum pressure for coexistent phases at constant temperature.

Another case of interest is when the composition of one of three coexistent phases is such as can be produced by combining the other two. In this case, the primitive surface must touch the same plane in three points in the same straight line. Let us distinguish the parts of the primitive surface to which these points belong as the sheets (A), (B), and (C), (C) denoting that which is intermediate in position. The sheet (C) is evidently tangent to the developable surface formed upon (A) and (B). It may or it may not intersect it at the point of contact. If it does not, it must lie above the developable surface, (unless it represents states which are unstable in regard to continuous changes,) and the surface of dissipated energy will include parts of the primitive sheets (A) and (B), the developable surface joining them, and the single point of the sheet (C) in which it meets this developable surface. Now, if the temperature or pressure is varied so as to make the sheet (C) rise above the

developable surface formed on the sheets (A) and (B), the surface of dissipated energy will be altered in its general features only by the removal of the single point of the sheet (C). But if the temperature or pressure is altered so as to make a part of the sheet (C) protrude through the developable surface formed on (A) and (B), the surface

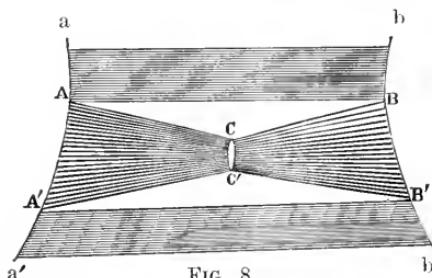


FIG. 8.

of dissipated energy will have the form indicated in figure 8. It will include two plane triangles ABC and A'B'C', a part of each of the sheets (A) and (B), represented in the figure by the spaces on the left of the line aAA'a' and on the right of the line bBB'b', a small part CC' of the sheet (C), and developable surfaces formed upon these sheets taken by pairs ACC'A', BCC'B', aABb, a'A'B'b', the last two being different portions of the same developable surface.

But if, when the primitive surface is constructed for such a temperature and pressure that it has three points of contact with the same plane in the same straight line, the sheet (C) (which has the middle position) at its point of contact with the triple tangent plane intersects the developable surface formed upon the other sheets (A) and (B), the surface of dissipated energy will not include this developable surface, but will consist of portions of the three primitive sheets with two developable surfaces formed on (A) and (C) and on (B) and (C). These developable surfaces meet one another at the point of contact of (C) with the triple tangent plane, dividing the

portion of this sheet which belongs to the surfaces of dissipated energy into two parts. If now the temperature or pressure are varied so as to make the sheet (C) sink relatively to the developable surface formed on (A) and (B), the only alteration in the general features of the surface of dissipated energy will be that the developable surfaces

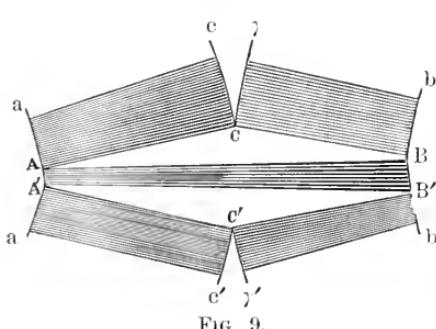


FIG. 9.

formed on (A) and (C) and on (B) and (C) will separate from one another, and the two parts of the sheet (C) will be merged in one. But a contrary variation of temperature or pressure will give a

surface of dissipated energy such as is represented in figure (9), containing two plane triangles ABC, A'B'C' belonging to triple tangent planes, a portion of the sheet (A) on the left of the line  $aAA'a'$ , a portion of the sheet (B) on the right of the line  $bBB'b'$ , two separate portions  $eC\gamma$  and  $c'C'\gamma'$  of the sheet (C), two separate portions  $aACCe$  and  $a'A'C'c'$  of the developable surface formed on (A) and (C), two separate portions  $bBC\gamma$  and  $b'B'C'\gamma'$  of the developable surface formed on (B) and (C), and the portion  $A'ABB'$  of the developable surface formed on (A) and (B).

From these geometrical relations it appears that (in general) the temperature of three coexistent phases is a maximum or minimum for constant pressure, and the pressure of three coexistent phases a maximum or minimum for constant temperature, when the composition of the three coexistent phases is such that one can be formed by combining the other two. This result has been obtained analytically on page 156.

The preceding examples are amply sufficient to illustrate the use of the  $m\text{-}\zeta$  surfaces and curves. The physical properties indicated by the nature of the surface of dissipated energy have been only occasionally mentioned, as they are often far more distinctly indicated by the diagrams than they could be in words. It will be observed that a knowledge of the lines which divide the various different portions of the surface of dissipated energy and of the direction of the rectilinear elements of the developable surfaces, *as projected upon the X-Y plane*, without a knowledge of the form of the  $m\text{-}\zeta$  surface in space, is sufficient for the determination (in respect to the quantity and composition of the resulting masses) of the combinations and separations of the substances, and of the changes in their states of aggregation, which take place when the substances are exposed to the temperature and pressure to which the projected lines relate, except so far as such transformations are prevented by passive resistances to change.

#### CRITICAL PHASES.

It has been ascertained by experiment that the variations of two coexistent states of the same substance are in some cases limited in one direction by a terminal state at which the distinction of the coexistent states vanishes.\* This state has been called the *critical state*. Analogous properties may doubtless be exhibited by compounds of variable composition without change of temperature or

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\* See Dr. Andrews "On the continuity of the gaseous and liquid states of matter." Phil. Trans., vol. 159, p. 575.

pressure. For if, at any given temperature and pressure, two liquids are capable of forming a stable mixture in any ratio  $m_1 : m_2$  less than  $a$ , and in any greater than  $b$ ,  $a$  and  $b$  being the values of that ratio for two coexistent phases, while either can form a stable mixture with a third liquid in all proportions, and any small quantities of the first and second can unite at once with a great quantity of the third to form a stable mixture, it may easily be seen that two coexistent mixtures of the three liquids may be varied in composition, the temperature and pressure remaining the same, from initial phases in each of which the quantity of the third liquid is nothing, to a terminal phase in which the distinction of the two phases vanishes.

In general, we may define a *critical phase* as one at which the distinction between coexistent phases vanishes. We may suppose the coexistent phases to be stable in respect to continuous changes, for although relations in some respects analogous might be imagined to hold true in regard to phases which are unstable in respect to continuous changes, the discussion of such cases would be devoid of interest. But if the coexistent phases and the critical phase are unstable only in respect to the possible formation of phases entirely different from the critical and adjacent phases, the liability to such changes will in no respect affect the relations between the critical and adjacent phases, and need not be considered in a theoretical discussion of these relations, although it may prevent an experimental realization of the phases considered. For the sake of brevity, in the following discussion, phases in the vicinity of the critical phase will generally be called stable, if they are unstable only in respect to the formation of phases entirely different from any in the vicinity of the critical phase.

Let us first consider the number of independent variations of which a critical phase (while remaining such) is capable. If we denote by  $n$  the number of independently variable components, a pair of coexistent phases will be capable of  $n$  independent variations, which may be expressed by the variations of  $n$  of the quantities  $t, p, \mu_1, \mu_2, \dots, \mu_n$ . If we limit these variations by giving to  $n-1$  of the quantities the constant values which they have for a certain critical phase, we obtain a linear\* series of pairs of coexistent phases terminated by the critical phase. If we now vary infinitesimally the values of these  $n-1$  quantities, we shall have for the new set of values considered constant a new linear series of pairs of coexistent phases. Now for every pair of phases in the first series, there must be pairs of phases in the

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\* This term is used to characterize a series having a *single* degree of extension.

second series differing infinitely little from the pair in the first, and *vice versa*, therefore the second series of coexistent phases must be terminated by a critical phase which differs, but differs infinitely little, from the first. We see, therefore, that if we vary arbitrarily the values of any  $n-1$  of the quantities  $t, p, \mu_1, \mu_2, \dots, \mu_n$ , as determined by a critical phase, we obtain one and only one critical phase for each set of varied values; i. e., a critical phase is capable of  $n-1$  independent variations.

The quantities  $t, p, \mu_1, \mu_2, \dots, \mu_n$ , have the same values in two coexistent phases, but the ratios of the quantities  $\eta, v, m_1, m_2, \dots, m_n$ , are in general different in the two phases. Or, if for convenience we compare equal volumes of the two phases (which involves no loss of generality), the quantities  $\eta, m_1, m_2, \dots, m_n$  will in general have different values in two coexistent phases. Applying this to coexistent phases indefinitely near to a critical phase, we see that in the immediate vicinity of a critical phase, if the values of  $n$  of the quantities  $t, p, \mu_1, \mu_2, \dots, \mu_n$ , are regarded as constant (as well as  $v$ ), the variations of either of the others will be infinitely small compared with the variations of the quantities  $\eta, m_1, m_2, \dots, m_n$ . This condition, which we may write in the form

$$\left( \frac{d\mu_n}{dm_n} \right)_{t, v, \mu_1, \dots, \mu_{n-1}} = 0, \quad (200)$$

characterizes, as we have seen on page 171, the limits which divide stable from unstable phases in respect to continuous changes.

In fact, if we give to the quantities  $t, \mu_1, \mu_2, \dots, \mu_{n-1}$  constant values determined by a pair of coexistent phases, and to  $\frac{m_n}{v}$  a series of values increasing from the less to the greater of the values which it has in these coexistent phases, we determine a linear series of phases connecting the coexistent phases, in some part of which  $\mu_n$ —since it has the same value in the two coexistent phases, but not a uniform value throughout the series (for if it had, which is theoretically improbable, all these phases would be coexistent)—must be a decreasing function of  $\frac{m_n}{v}$ , or of  $m_n$ , if  $v$  also is supposed constant. Therefore, the series must contain phases which are unstable in respect to continuous changes. (See page 168.) And as such a pair of coexistent phases may be taken indefinitely near to any critical phase, the unstable phases (with respect to continuous changes) must approach indefinitely near to this phase.

Critical phases have similar properties with reference to stability as determined with regard to discontinuous changes. For as every stable phase which has a coexistent phase lies upon the limit which separates stable from unstable phases, the same must be true of any stable critical phase. (The same may be said of critical phases which are unstable in regard to discontinuous changes if we leave out of account the liability to the particular kind of discontinuous change in respect to which the critical phase is unstable.)

The linear series of phases determined by giving to  $n$  of the quantities  $t, p, \mu_1, \mu_2, \dots, \mu_n$  the constant values which they have in any pair of coexistent phases consists of unstable phases in the part between the coexistent phases, but in the part beyond these phases in either direction it consists of stable phases. Hence, if a critical phase is varied in such a manner that  $n$  of the quantities  $t, p, \mu_1, \mu_2, \dots, \mu_n$  remain constant, it will remain stable in respect both to continuous and to discontinuous changes. Therefore,  $\mu_n$  is an increasing function of  $m_n$  when  $t, v, \mu_1, \mu_2, \dots, \mu_{n-1}$  have constant values determined by any critical phase. But as equation (200) holds true at the critical phase, the following conditions must also hold true at that phase:

$$\left( \frac{d^2 \mu_n}{dm_n^2} \right)_{t, v, \mu_1, \dots, \mu_{n-1}} = 0, \quad (201)$$

$$\left( \frac{d^3 \mu_n}{dm_n^3} \right)_{t, v, \mu_1, \dots, \mu_{n-1}} \geq 0. \quad (202)$$

If the sign of equality holds in the last condition, additional conditions, concerning the differential coefficients of higher orders, must be satisfied.

Equations (200) and (201) may in general be called the equations of critical phases. It is evident that there are only two independent equations of this character, as a critical phase is capable of  $n - 1$  independent variations.

We are not, however, absolutely certain that equation (200) will always be satisfied by a critical phase. For it is possible that the denominator in the fraction may vanish as well as the numerator for an infinitesimal change of phase in which the quantities indicated are constant. In such a case, we may suppose the subscript  $n$  to refer to some different component substance, or use another differential coefficient of the same general form (such as are described on page 171 as characterizing the limits of stability in respect to continuous changes), making the corresponding changes in (201) and (202). We may be certain that some of the formulæ thus formed will not fail. But for a perfectly rigorous method there is an advan-

tage in the use of  $\eta, v, m_1, m_2, \dots, m_n$  as independent variables. The condition that the phase may be varied without altering any of the quantities  $t, \mu_1, \mu_2, \dots, \mu_n$  will then be expressed by the equation

$$R_{n+1} = 0, \quad (203)$$

in which  $R_{n+1}$  denotes the same determinant as on page 169. To obtain the second equation characteristic of critical phases, we observe that as a phase which is critical cannot become unstable when varied so that  $n$  of the quantities  $t, p, \mu_1, \mu_2, \dots, \mu_n$  remain constant, the differential of  $R_{n+1}$  for constant volume, viz.,

$$\frac{dR_{n+1}}{d\eta} d\eta + \frac{dR_{n+1}}{dm_1} dm_1 \dots + \frac{dR_{n+1}}{dm_n} dm_n \quad (204)$$

cannot become negative when  $n$  of the equations (172) are satisfied. Neither can it have a positive value, for then its value might become negative by a change of sign of  $d\eta, dm_1$ , etc. Therefore the expression (204) has the value zero, if  $n$  of the equations (172) are satisfied. This may be expressed by an equation

$$S = 0, \quad (205)$$

in which  $S$  denotes a determinant in which the constituents are the same as in  $R_{n+1}$ , except in a single horizontal line, in which the differential coefficients in (204) are to be substituted. In whatever line this substitution is made, the equation (205), as well as (203), will hold true of every critical phase without exception.

If we choose  $t, p, m_1, m_2, \dots, m_n$  as independent variables, and write  $U$  for the determinant

$$\begin{vmatrix} \frac{d^2\zeta}{dm_1^2} & \frac{d^2\zeta}{dm_2 dm_1} & \cdots & \frac{d^2\zeta}{dm_{n-1} dm_1} \\ \frac{d^2\zeta}{dm_1 dm_2} & \frac{d^2\zeta}{dm_2^2} & \cdots & \frac{d^2\zeta}{dm_{n-1} dm_2} \\ \vdots & \vdots & \ddots & \vdots \\ \frac{d^2\zeta}{dm_1 dm_{n-1}} & \frac{d^2\zeta}{dm_2 dm_{n-1}} & \cdots & \frac{d^2\zeta}{dm_{n-1}^2} \end{vmatrix} \quad (206)$$

and  $V$  for the determinant formed from this by substituting for the constituents in any horizontal line the expressions

$$\frac{dU}{dm_1}, \quad \frac{dU}{dm_2}, \quad \cdots \quad \frac{dU}{dm_{n-1}}, \quad (207)$$

the equations of critical phases will be

$$U = 0, \quad V = 0. \quad (208)$$

It results immediately from the definition of a critical phase, that an infinitesimal change in the condition of a mass in such a phase

may cause the mass, if it remains in a state of dissipated energy (i. e., in a state in which the dissipation of energy by internal processes is complete), to cease to be homogeneous. In this respect a critical phase resembles any phase which has a coexistent phase, but differs from such phases in that the two parts into which the mass divides when it ceases to be homogeneous differ infinitely little from each other and from the original phase, and that neither of these parts is in general infinitely small. If we consider a change in the mass to be determined by the values of  $d\eta$ ,  $dv$ ,  $dm_1$ ,  $dm_2$ , . . .  $dm_n$ , it is evident that the change in question will cause the mass to cease to be homogeneous whenever the expression

$$\frac{dR_{n+1}}{d\eta} d\eta + \frac{dR_{n+1}}{dv} dv + \frac{dR_{n+1}}{dm_1} dm_1 + \dots + \frac{dR_{n+1}}{dm_n} dm_n \quad (209)$$

has a negative value. For if the mass should remain homogeneous, it would become unstable, as  $R_{n+1}$  would become negative. Hence, in general, any change thus determined, or its reverse (determined by giving to  $d\eta$ ,  $dv$ ,  $dm_1$ ,  $dm_2$ , . . .  $dm_n$  the same values taken negatively), will cause the mass to cease to be homogeneous. The condition which must be satisfied with reference to  $d\eta$ ,  $dv$ ,  $dm_1$ ,  $dm_2$ , . . .  $dm_n$ , in order that neither the change indicated, nor the reverse, shall destroy the homogeneity of the mass, is expressed by equating the above expression to zero.

But if we consider the change in the state of the mass (supposed to remain in a state of dissipated energy) to be determined by arbitrary values of  $n+1$  of the differentials  $dt$ ,  $d\rho$ ,  $d\mu_1$ ,  $d\mu_2$ , . . .  $d\mu_n$ , the case will be entirely different. For, if the mass ceases to be homogeneous, it will consist of two coexistent phases, and as applied to these only  $n$  of the quantities  $t$ ,  $\rho$ ,  $\mu_1$ ,  $\mu_2$ , . . .  $\mu_n$  will be independent. Therefore, for arbitrary variations of  $n+1$  of these quantities, the mass must in general remain homogeneous.

But if, instead of supposing the mass to remain in a state of dissipated energy, we suppose that it remains homogeneous, it may easily be shown that to certain values of  $n+1$  of the above differentials there will correspond three different phases, of which one is stable with respect both to continuous and to discontinuous changes, another is stable with respect to the former and unstable with respect to the latter, and the third is unstable with respect to both.

In general, however, if  $n$  of the quantities  $p$ ,  $t$ ,  $\mu_1$ ,  $\mu_2$ , . . .  $\mu_n$ , or  $n$  arbitrary functions of these quantities, have the same constant values as at a critical phase, the linear series of phases thus determined will be stable, in the vicinity of the critical phase. But if less

than  $n$  of these quantities or functions of the same together with certain of the quantities  $\eta, v, m_1, m_2, \dots, m_n$ , or arbitrary functions of the latter quantities, have the same values as at a critical phase, so as to determine a linear series of phases, the differential of  $R_{n+1}$  in such a series of phases will not in general vanish at the critical phase, so that in general a part of the series will be unstable.

We may illustrate these relations by considering separately the cases in which  $n=1$  and  $n=2$ . If a mass of invariable composition is in a critical state, we may keep its volume constant, and destroy its homogeneity by changing its entropy (i. e., by adding or subtracting heat—probably the latter), or we may keep its entropy constant and destroy its homogeneity by changing its volume; but if we keep its pressure constant we cannot destroy its homogeneity by any thermal action, nor if we keep its temperature constant can we destroy its homogeneity by any mechanical action.

When a mass having two independently variable components is in a critical phase, and either its volume or its pressure is maintained constant, its homogeneity may be destroyed by a change of entropy or temperature. Or, if either its entropy or its temperature is maintained constant, its homogeneity may be destroyed by a change of volume or pressure. In both these cases it is supposed that the quantities of the components remain unchanged. But if we suppose both the temperature and the pressure to be maintained constant, the mass will remain homogeneous, however the proportion of the components be changed. Or, if a mass consists of two coexistent phases, one of which is a critical phase having two independently variable components, and either the temperature or the pressure of the mass is maintained constant, it will not be possible by mechanical or thermal means, or by changing the quantities of the components, to cause the critical phase to change into a pair of coexistent phases, so as to give three coexistent phases in the whole mass. The statements of this paragraph and of the preceding have reference only to infinitesimal changes.\*

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\* A brief abstract (which came to the author's notice after the above was in type) of a memoir by M. Duclaux, "Sur la séparation des liquides mélangés, etc." will be found in *Comptes Rendus*, vol. lxxxi. (1875), p. 815.

## ON THE VALUES OF THE POTENTIALS WHEN THE QUANTITY OF ONE OF THE COMPONENTS IS VERY SMALL.

If we apply equation (97) to a homogeneous mass having two independently variable components  $S_1$  and  $S_2$ , and make  $t$ ,  $p$ , and  $m_1$ , constant, we obtain

$$m_1 \left( \frac{d\mu_1}{dm_2} \right)_{t, p, m_1} + m_2 \left( \frac{d\mu_2}{dm_2} \right)_{t, p, m_1} = 0. \quad (210)$$

Therefore, for  $m_2=0$ , either

$$\left( \frac{d\mu_1}{dm_2} \right)_{t, p, m_1} = 0, \quad (211)$$

or

$$\left( \frac{d\mu_2}{dm_2} \right)_{t, p, m_1} = \infty. \quad (212)$$

Now, whatever may be the composition of the mass considered, we may always so choose the substance  $S_1$  that the mass shall consist solely of that substance, and in respect to any other variable component  $S_2$ , we shall have  $m_2=0$ . But equation (212) cannot hold true *in general* as thus applied. For it may easily be shown (as has been done with regard to the potential on pages 148, 149) that the value of a differential coefficient like that in (212) for any given mass, when the substance  $S_2$  (to which  $m_2$  and  $\mu_2$  relate) is determined, is independent of the particular substance which we may regard as the other component of the mass; so that, if equation (212) holds true when the substance denoted by  $S_1$  has been so chosen that  $m_2=0$ , it must hold true without such a restriction, which cannot generally be the case.

In fact, it is easy to prove directly that equation (211) will hold true of any phase which is stable in regard to continuous changes and in which  $m_2=0$ , if  $m_2$  is capable of negative as well as positive values. For by (171), in any phase having that kind of stability,  $\mu_1$  is an increasing function of  $m_1$  when  $t$ ,  $p$ , and  $m_2$  are regarded as constant. Hence,  $\mu_1$  will have its greatest value when the mass consists wholly of  $S_1$ , i. e., when  $m_2=0$ . Therefore, if  $m_2$  is capable of negative as well as positive values, equation (211) must hold true for  $m_2=0$ . (This appears also from the geometrical representation of potentials in the  $m$ - $\xi$  curve. See page 177.)

But if  $m_2$  is capable only of positive values, we can only conclude from the preceding considerations that the value of the differential coefficient in (211) cannot be positive. Nor, if we consider the physical significance of this case, viz., that an increase of  $m_2$  denotes an

addition to the mass in question of a substance not before contained in it, does any reason appear for supposing that this differential coefficient has generally the value zero. To fix our ideas, let us suppose that  $S_1$  denotes water, and  $S_2$  a salt (either anhydrous or any particular hydrate). The addition of the salt to water, previously in a state capable of equilibrium with vapor or with ice, will destroy the possibility of such equilibrium at the same temperature and pressure. The liquid will dissolve the ice, or condense the vapor, which is brought in contact with it under such circumstances, which shows that  $\mu_1$  (the potential for water in the liquid mass) is diminished by the addition of the salt, when the temperature and pressure are maintained constant. Now there seems to be no *a priori* reason for supposing that the ratio of this diminution of the potential for water to the quantity of the salt which is added vanishes with this quantity. We should rather expect that, for small quantities of the salt, an effect of this kind would be proportional to its cause, i. e., that the differential coefficient in (211) would have a finite negative value for an infinitesimal value of  $m_2$ . That this is the case with respect to numerous watery solutions of salts is distinctly indicated by the experiments of Wüllner\* on the tension of the vapor yielded by such solutions, and of Rüdorff† on the temperature at which ice is formed in them; and unless we have experimental evidence that cases are numerous in which the contrary is true, it seems not unreasonable to assume, as a *general* law, that when  $m_2$  has the value zero and is incapable of negative values, the differential coefficient in (211) will have a finite negative value, and that equation (212) will therefore hold true. But this case must be carefully distinguished from that in which  $m_2$  is capable of negative values, which also may be illustrated by a solution of a salt in water. For this purpose let  $S_1$  denote a hydrate of the salt which can be crystallized, and let  $S_2$  denote water, and let us consider a liquid consisting entirely of  $S_1$  and of such temperature and pressure as to be in equilibrium with crystals of  $S_1$ . In such a liquid, an increase or a diminution of the quantity of water would alike cause crystals of  $S_1$  to dissolve, which requires that the differential coefficient in (211) shall vanish at the particular phase of the liquid for which  $m_2 = 0$ .

Let us return to the case in which  $m_2$  is incapable of negative values, and examine, without other restriction in regard to the substances

\* Pogg. Ann., vol. ciii. (1858), p. 529; vol. cv. (1858), p. 85; vol. cx. (1860), p. 564.

† Pogg. Ann., vol. cxiv. (1861), p. 63.

denoted by  $S_1$  and  $S_2$ , the relation between  $\mu_2$  and  $\frac{m_2}{m_1}$  for any constant temperature and pressure and for such small values of  $\frac{m_2}{m_1}$  that the differential coefficient in (211) may be regarded as having the same constant value as when  $m_2 = 0$ , the values of  $t$ ,  $p$ , and  $m_1$  being unchanged. If we denote this value of the differential coefficient by  $\frac{-A}{m_1}$ , the value of  $A$  will be positive, and will be independent of  $m_1$ .

Then for small values of  $\frac{m_2}{m_1}$ , we have by (210), approximately,

$$m_2 \left( \frac{d\mu_2}{dm_2} \right)_{t, p, m_1} = A, \quad (213)$$

i. e.,

$$\left( \frac{d\mu_2}{d \log m_2} \right)_{t, p, m_1} = A. \quad (214)$$

If we write the integral of this equation in the form

$$\mu_2 = A \log \frac{Bm_2}{m_1}, \quad (215)$$

$B$  like  $A$  will have a positive value depending only upon the temperature and pressure. As this equation is to be applied only to cases in which the value of  $m_2$  is very small compared with  $m_1$ , we may regard  $\frac{m_1}{v}$  as constant, when temperature and pressure are constant, and write

$$\mu_2 = A \log \frac{Cm_2}{v}, \quad (216)$$

$C$  denoting a positive quantity, dependent only upon the temperature and pressure.

We have so far considered the composition of the body as varying only in regard to the proportion of two components. But the argument will be in no respect invalidated, if we suppose the composition of the body to be capable of other variations. In this case, the quantities  $A$  and  $C$  will be functions not only of the temperature and pressure but also of the quantities which express the composition of the substance of which together with  $S_2$  the body is composed. If the quantities of any of the components besides  $S_2$  are very small (relatively to the quantities of others), it seems reasonable to assume that the value of  $\mu_2$ , and therefore the values of  $A$  and  $C$ , will be nearly the same as if these components were absent.

Hence, if the independently variable components of any body are  $S_a, \dots, S_g$ , and  $S_h, \dots, S_k$ , the quantities of the latter being very small as compared with the quantities of the former, and are incapable of negative values, we may express approximately the values of the potentials for  $S_h, \dots, S_k$  by equations (subject of course to the uncertainties of the assumptions which have been made) of the form

$$\mu_h = A_h \log \frac{C_h m_h}{v}, \quad (217)$$

$$\mu_k = A_k \log \frac{C_k m_k}{v}, \quad (218)$$

in which  $A_h, C_h, \dots, A_k, C_k$  denote functions of the temperature, the pressure, and the ratios of the quantities  $m_a, \dots, m_g$ .

We shall see hereafter, when we come to consider the properties of gases, that these equations may be verified experimentally in a very large class of cases, so that we have considerable reason for believing that they express a general law in regard to the limiting values of potentials.\*

#### ON CERTAIN POINTS RELATING TO THE MOLECULAR CONSTITUTION OF BODIES.

It not unfrequently occurs that the number of proximate components which it is necessary to recognize as independently variable in a body exceeds the number of components which would be sufficient to express its ultimate composition. Such is the case, for example, as has been remarked on page 117, in regard to a mixture at ordinary temperatures of vapor of water and free hydrogen and oxygen. This case is explained by the existence of three sorts of molecules in the gaseous mass, viz., molecules of hydrogen, of oxygen, and of hydrogen and oxygen combined. In other cases, which are essentially the same in principle, we suppose a greater number of different sorts of molecules, which differ in composition, and the relations between

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\* The reader will not fail to remark that, if we could assume the universality of this law, the statement of the conditions necessary for equilibrium between different masses in contact would be much simplified. For, as the potential for a substance which is only a *possible* component (see page 117) would always have the value  $-\infty$ , the case could not occur that the potential for any substance should have a greater value in a mass in which that substance is only a possible component, than in another mass in which it is an actual component; and the conditions (22) and (51) might be expressed with the sign of equality without exception for the case of possible components.

these may be more complicated. Other cases are explained by molecules which differ in the quantity of matter which they contain, but not in the kind of matter, nor in the proportion of the different kinds. In still other cases, there appear to be different sorts of molecules, which differ neither in the kind nor in the quantity of matter which they contain, but only in the manner in which they are constituted. What is essential in the cases referred to is that a certain number of some sort or sorts of molecules shall be equivalent to a certain number of some other sort or sorts in respect to the kinds and quantities of matter which they collectively contain, and yet the former shall never be transformed into the latter within the body considered, nor the latter into the former, however the proportion of the numbers of the different sorts of molecules may be varied, or the composition of the body in other respects, or its thermodynamic state as represented by temperature and pressure or any other two suitable variables, provided, it may be, that these variations do not exceed certain limits. Thus, in the example given above, the temperature must not be raised beyond a certain limit, or molecules of hydrogen and of oxygen may be transformed into molecules of water.

The differences in bodies resulting from such differences in the constitution of their molecules are capable of continuous variation, in bodies containing the same matter and in the same thermodynamic state as determined, for example, by pressure and temperature, as the numbers of the molecules of the different sorts are varied. These differences are thus distinguished from those which depend upon the manner in which the molecules are combined to form sensible masses. The latter do not cause an increase in the number of variables in the fundamental equation; but they may be the cause of different values of which the function is sometimes capable for one set of values of the independent variables, as, for example, when we have several different values of  $\xi$  for the same values of  $t, p, m_1, m_2, \dots, m_n$ , one perhaps being for a gaseous body, one for a liquid, one for an amorphous solid, and others for different kinds of crystals, and all being invariable for constant values of the above mentioned independent variables.

But it must be observed that when the differences in the constitution of the molecules are entirely determined by the quantities of the different kinds of matter in a body with the two variables which express its thermodynamic state, these differences will not involve any increase in the number of variables in the fundamental equation. For example, if we should raise the temperature of the mixture of

vapor of water and free hydrogen and oxygen, which we have just considered, to a point at which the numbers of the different sorts of molecules are entirely determined by the temperature and pressure and the total quantities of hydrogen and of oxygen which are present, the fundamental equation of such a mass would involve but four independent variables, which might be the four quantities just mentioned. The fact of a certain part of the matter present existing in the form of vapor of water would, of course, be one of the facts which determine the nature of the relation between  $\zeta$  and the independent variables, which is expressed by the fundamental equation.

But in the case first considered, in which the quantities of the different sorts of molecules are *not* determined by the temperature and pressure and the quantities of the different kinds of matter in the body as determined by its ultimate analysis, the components of which the quantities or the potentials appear in the fundamental equation must be those which are determined by the proximate analysis of the body, so that the variations in their quantities, with two variations relating to the thermodynamic state of the body, shall include all the variations of which the body is capable.\* Such cases present no especial difficulty; there is indeed nothing in the physical and chemical properties of such bodies, so far as a certain range of experiments is concerned, which is different from what might be, if the proximate components were incapable of farther reduction or transformation. Yet among the various phases of the kinds of matter concerned, represented by the different sets of values of the variables which satisfy the fundamental equation, there is a certain class which merit especial attention. These are the phases for which the entropy has a maximum value for the same matter, as determined by the ultimate analysis of the body, with the same energy and volume. To fix our ideas let us call the proximate components  $S_1, \dots, S_n$ , and the ultimate components  $S_a, \dots, S_b$ ; and let  $m_1, \dots, m_n$  denote the quantities of the former, and  $m_a, \dots, m_b$ , the quantities of the latter. It is evident that  $m_a \dots m_b$  are homogeneous functions of the first degree of  $m_1, \dots, m_n$ ; and that the relations between the substances  $S_1, \dots, S_n$  might be expressed by homogeneous equations of the first degree between the units of these substances, equal in number to the difference of the numbers of the proximate and of the ultimate com-

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\* The terms proximate or ultimate are not necessarily to be understood in an absolute sense. All that is said here and in the following paragraphs will apply to many cases in which components may conveniently be regarded as proximate or ultimate, which are such only in a relative sense.

ponents. The phases in question are those for which  $\eta$  is a maximum for constant values of  $\epsilon, v, m_a, \dots m_b$ ; or, as they may also be described, those for which  $\epsilon$  is a minimum for constant values of  $\eta, v, m_a, \dots m_b$ ; or for which  $\zeta$  is a minimum for constant values of  $t, p, m_a, \dots m_b$ . The phases which satisfy this condition may be readily determined when the fundamental equation (which will contain the quantities  $m_1, \dots m_n$  or  $\mu_1, \dots \mu_n$ ) is known. Indeed it is easy to see that we may express the conditions which determine these phases by substituting  $\mu_1, \dots \mu_n$  for the letters denoting the units of the corresponding substances in the equations which express the equivalence in ultimate analysis between these units.

These phases may be called, with reference to the kind of change which we are considering, phases of dissipated energy. That we have used a similar term before, with reference to a different kind of changes, yet in a sense entirely analogous, need not create confusion.

It is characteristic of these phases that we cannot alter the values of  $m_1, \dots m_n$  in any real mass in such a phase, while the volume of the mass as well as its matter remain unchanged, without diminishing the energy or increasing the entropy of some other system. Hence, if the mass is large, its equilibrium can be but slightly disturbed by the action of any small body, or by a single electric spark, or by any cause which is not in some way proportioned to the effect to be produced. But when the proportion of the proximate components of a mass taken in connection with its temperature and pressure is not such as to constitute a phase of dissipated energy, it may be possible to cause great changes in the mass by the contact of a very small body. Indeed it is possible that the changes produced by such contact may only be limited by the attainment of a phase of dissipated energy. Such a result will probably be produced in a fluid mass by contact with another fluid which contains molecules of all the kinds which occur in the first fluid (or at least all those which contain the same kinds of matter which also occur in other sorts of molecules), but which differs from the first fluid in that the quantities of the various kinds of molecules are entirely determined by the ultimate composition of the fluid and its temperature and pressure. Or, to speak without reference to the molecular state of the fluid, the result considered would doubtless be brought about by contact with another fluid, which absorbs all the proximate components of the first,  $S_1, \dots S_n$  (or all those between which there exist relations of equivalence in respect to their ultimate analysis), independently, and without passive resistances, but for which the phase is completely deter-

mined by its temperature and pressure and its ultimate composition (in respect at least to the particular substances just mentioned). By the absorption of the substances  $S_1, \dots, S_n$  independently and without passive resistances, it is meant that when the absorbing body is in equilibrium with another containing these substances, it shall be possible by *infinitesimal* changes in these bodies to produce the exchange of all these substances in either direction and independently. An exception to the preceding statement may of course be made for cases in which the result in question is prevented by the occurrence of some other kinds of change; in other words, it is assumed that the two bodies can remain in contact preserving the properties which have been mentioned.

The term *catalysis* has been applied to such action as we are considering. When a body has the property of reducing another, without limitation with respect to the proportion of the two bodies, to a phase of dissipated energy, in regard to a certain kind of molecular change, it may be called a *perfect catalytic agent* with respect to the second body and the kind of molecular change considered.

It seems not improbable that in some cases in which molecular changes take place slowly in homogeneous bodies, a mass of which the temperature and pressure are maintained constant will be finally brought to a state of equilibrium which is entirely determined by its temperature and pressure and the quantities of its ultimate components, while the various transitory states through which the mass passes, (which are evidently not completely defined by the quantities just mentioned,) may be completely defined by the quantities of certain proximate components with the temperature and pressure, and the matter of the mass may be brought by processes approximately reversible from permanent states to these various transitory states. In such cases, we may form a fundamental equation with reference to all possible phases, whether transitory or permanent; and we may also form a fundamental equation of different import and containing a smaller number of independent variables, which has reference solely to the final phases of equilibrium. The latter are the phases of dissipated energy (with reference to molecular changes), and when the more general form of the fundamental equation is known, it will be easy to derive from it the fundamental equation for these permanent phases alone.

Now, as these relations, theoretically considered, are independent of the rapidity of the molecular changes, the question naturally arises, whether in cases in which we are not able to distinguish such transi-

tory phases, they may not still have a theoretical significance. If so, the consideration of the subject from this point of view, may assist us, in such cases, in discovering the form of the fundamental equation with reference to the ultimate components, which is the only equation required to express all the properties of the bodies which are capable of experimental demonstration. Thus, when the phase of a body is completely determined by the quantities of  $n$  independently variable components, with the temperature and pressure, and we have reason to suppose that the body is composed of a greater number  $n'$  of proximate components, which are therefore not independently variable (while the temperature and pressure remain constant), it seems quite possible that the fundamental equation of the body may be of the same form as the equation for the phases of dissipated energy of analogous compounds of  $n'$  proximate and  $n$  ultimate components, in which the proximate components are capable of independent variation (without variation of temperature or pressure). And if such is found to be the case, the fact will be of interest as affording an indication concerning the proximate constitution of the body.

Such considerations seem to be especially applicable to the very common case in which at certain temperatures and pressures, regarded as constant, the quantities of certain proximate components of a mass are capable of independent variations, and all the phases produced by these variations are permanent in their nature, while at other temperatures and pressures, likewise regarded as constant, the quantities of these proximate components are not capable of independent variation, and the phase may be completely defined by the quantities of the ultimate components with the temperature and pressure. There may be, at certain intermediate temperatures and pressures, a condition with respect to the independence of the proximate components intermediate in character, in which the quantities of the proximate components are independently variable when we consider all phases, the essentially transitory as well as the permanent, but in which these quantities are not independently variable when we consider the permanent phases alone. Now we have no reason to believe that the passing of a body in a state of dissipated energy from one to another of the three conditions mentioned has any necessary connection with any discontinuous change of state. Passing the limit which separates one of these states from another will not therefore involve any discontinuous change in the values of any of the quantities enumerated in (99)–(103) on page 143, if  $m_1, m_2, \dots, m_n, \mu_1, \mu_2, \dots, \mu_n$  are

understood as always relating to the ultimate components of the body. Therefore, if we regard masses in the different conditions mentioned above as having different fundamental equations, (which we may suppose to be of any one of the five kinds described on page 143,) these equations will agree at the limits dividing these conditions not only in the values of all the variables which appear in the equations, but also in all the differential coefficients of the first order involving these variables. We may illustrate these relations by supposing the values of  $t$ ,  $p$ , and  $\xi$  for a mass in which the quantities of the ultimate components are constant to be represented by rectilinear coordinates. Where the proximate composition of such a mass is not determined by  $t$  and  $p$ , the value of  $\xi$  will not be determined by these variables, and the points representing connected values of  $t$ ,  $p$ , and  $\xi$  will form a solid. This solid will be bounded in the direction opposite to that in which  $\xi$  is measured, by a surface which represents the phases of dissipated energy. In a part of the figure, all the phases thus represented may be permanent, in another part only the phases in the bounding surface, and in a third part there may be no such solid figure (for any phases of which the existence is experimentally demonstrable), but only a surface. This surface together with the bounding surfaces representing phases of dissipated energy in the parts of the figure mentioned above forms a continuous sheet, without discontinuity in regard to the direction of its normal at the limits dividing the different parts of the figure which have been mentioned. (There may, indeed, be different sheets representing liquid and gaseous states, etc., but if we limit our consideration to states of one of these sorts, the case will be as has been stated.)

We shall hereafter, in the discussion of the fundamental equations of gases, have an example of the derivation of the fundamental equation for phases of dissipated energy (with respect to the molecular changes on which the proximate composition of the body depends) from the more general form of the fundamental equation.

#### THE CONDITIONS OF EQUILIBRIUM FOR HETEROGENEOUS MASSES UNDER THE INFLUENCE OF GRAVITY.

Let us now seek the conditions of equilibrium for a mass of various kinds of matter subject to the influence of gravity. It will be convenient to suppose the mass enclosed in an immovable envelop which is impermeable to matter and to heat, and in other respects, except in regard to gravity, to make the same suppositions as on pages 115, 116. The energy of the mass will now consist of two parts, one of

which depends upon its intrinsic nature and state, and the other upon its position in space. Let  $Dm$  denote an element of the mass,  $D\varepsilon$  the intrinsic energy of this element,  $h$  its height above a fixed horizontal plane, and  $g$  the force of gravity; then the total energy of the mass (when without sensible motions) will be expressed by the formula

$$\int D\varepsilon + \int gh Dm, \quad (219)$$

in which the integrations include all the elements of the mass; and the general condition of equilibrium will be

$$\delta \int D\varepsilon + \delta \int gh Dm \geq 0, \quad (220)$$

the variations being subject to certain equations of condition. These must express that the entropy of the whole mass is constant, that the surface bounding the whole mass is fixed, and that the total quantities of each of the component substances is constant. We shall suppose that there are no other equations of condition, and that the independently variable components are the same throughout the whole mass; and we shall at first limit ourselves to the consideration of the conditions of equilibrium with respect to the changes which may be expressed by infinitesimal variations of the quantities which define the initial state of the mass, without regarding the possibility of the formation at any place of infinitesimal masses entirely different from any initially existing in the same vicinity.

Let  $D\eta$ ,  $Dv$ ,  $Dm_1, \dots, Dm_n$  denote the entropy of the element  $Dm$ , its volume, and the quantities which it contains of the various components. Then

$$Dm = Dm_1 \dots + Dm_n, \quad (221)$$

and

$$\delta Dm = \delta Dm_1 \dots + \delta Dm_n. \quad (222)$$

Also, by equation (12),

$$\delta D\varepsilon = t \delta D\eta - p \delta Dv + \mu_1 \delta Dm_1 \dots + \mu_n \delta Dm_n. \quad (223)$$

By these equations the general condition of equilibrium may be reduced to the form

$$\begin{aligned} & \int t \delta D\eta - \int p \delta Dv + \int \mu_1 \delta Dm_1 \dots + \int \mu_n \delta Dm_n \\ & + \int gh \delta h Dm + \int gh \delta Dm_1 \dots + \int gh \delta Dm_n \geq 0. \end{aligned} \quad (224)$$

Now it will be observed that the different equations of condition affect different parts of this condition, so that we must have, separately,

$$\int t \delta D\eta \geq 0, \quad \text{if} \quad \int \delta D\eta = 0; \quad (225)$$

$$-\int p \delta Dv + \int g \delta h Dm \geq 0, \quad (226)$$

if the bounding surface is unvaried;

$$\left. \begin{array}{l} \int \mu_1 \delta Dm_1 + \int g h \delta Dm_1 \geq 0, \text{ if } \int \delta Dm_1 = 0; \\ \vdots \quad \vdots \\ \int \mu_n \delta Dm_n + \int g h \delta Dm_n \geq 0, \text{ if } \int \delta Dm_n = 0. \end{array} \right\} \quad (227)$$

From (225) we may derive the condition of thermal equilibrium,

$$t = \text{Const.} \quad (328)$$

Condition (226) is evidently the ordinary mechanical condition of equilibrium, and may be transformed by any of the usual methods. We may, for example, apply the formula to such motions as might take place longitudinally within an infinitely narrow tube, terminated at both ends by the external surface of the mass, but otherwise of indeterminate form. If we denote by  $m$  the mass, and by  $v$  the volume, included in the part of the tube between one end and a transverse section of variable position, the condition will take the form

$$-\int p \delta dv + \int g \delta h dm \geq 0, \quad (229)$$

in which the integrations include the whole contents of the tube. Since no motion is possible at the ends of the tube,

$$\int p \delta dv + \int \delta v dp = \int d(p \delta v) = 0. \quad (230)$$

Again, if we denote by  $\gamma$  the density of the fluid,

$$\int g \delta h dm = \int g \frac{dh}{dv} \delta v \gamma dv = \int g \gamma \delta v dh. \quad (231)$$

By these equations condition (229) may be reduced to the form

$$\int \delta v (dp + g \gamma dh) \geq 0. \quad (232)$$

Therefore, since  $\delta v$  is arbitrary in value,

$$dp = -g \gamma dh, \quad (233)$$

which will hold true at any point in the tube, the differentials being taken with respect to the direction of the tube at that point. Therefore, as the form of the tube is indeterminate, this equation must hold true, without restriction, throughout the whole mass. It evidently requires that the pressure shall be a function of the height alone, and that the density shall be equal to the first derivative of this function, divided by  $-g$ .

Conditions (227) contain all that is characteristic of chemical equilibrium. To satisfy these conditions it is necessary and sufficient that

$$\left. \begin{array}{l} \mu_1 + gh = \text{Const.} \\ \vdots \quad \vdots \quad \vdots \quad \vdots \quad \vdots \quad \vdots \\ \mu_n + gh = \text{Const.} \end{array} \right\} \quad (234)$$

The expressions  $\mu_1, \dots, \mu_n$  denote quantities which we have called the potentials for the several components, and which are entirely determined at any point in a mass by the nature and state of the mass about that point. We may avoid all confusion between these quantities and the potential of the force of gravity, if we distinguish the former, when necessary, as *intrinsic* potentials. The relations indicated by equations (234) may then be expressed as follows:

*When a fluid mass is in equilibrium under the influence of gravity, and has the same independently variable components throughout, the intrinsic potentials for the several components are constant in any given level, and diminish uniformly as the height increases, the difference of the values of the intrinsic potential for any component at two different levels, being equal to the work done by the force of gravity when a unit of matter falls from the higher to the lower level.*

The conditions expressed by equations (228), (233), (234) are necessary and sufficient for equilibrium, except with respect to the possible formation of masses which are not approximately identical in phase with any previously existing about the points where they may be formed. The possibility of such formations at any point is evidently independent of the action of gravity, and is determined entirely by the phase or phases of the matter about that point. The conditions of equilibrium in this respect have been discussed on pages 128–134.

But equations (228), (233), and (234) are not entirely independent. For with respect to any mass in which there are no surfaces of discontinuity (i. e., surfaces where adjacent elements of mass have finite differences of phase), one of these equations will be a consequence of the others. Thus by (228) and (234), we may obtain from (97), which will hold true of any continuous variations of phase, the equation

$$v dp = -g(m_1 + \dots + m_n) dh; \quad (235)$$

$$\text{or} \quad dp = -g \gamma dh; \quad (236)$$

which will therefore hold true in any mass in which equations (228) and (234) are satisfied, and in which there are no surfaces of discontinuity. But the condition of equilibrium expressed by equation (233) has no exception with respect to surfaces of discontinuity; therefore in any mass in which such surfaces occur, it will be necessary for equilibrium, in addition to the relations expressed by equations (228) and (234), that there shall be no discontinuous change of pressure at these surfaces.

This superfluity in the particular conditions of equilibrium which we have found, as applied to a mass which is everywhere continuous

in phase, is due to the fact that we have made the elements of volume variable in position and size, while the matter initially contained in these elements is not supposed to be confined to them. Now, as the different components may move in different directions when the state of the system varies, it is evidently impossible to define the elements of volume so as always to include the same matter; we must, therefore, suppose the matter contained in the elements of volume to vary; and therefore it would be allowable to make these elements fixed in space. If the given mass has no surfaces of discontinuity, this would be much the simplest plan. But if there are any surfaces of discontinuity, it will be possible for the state of the given mass to vary, not only by infinitesimal changes of phase in the fixed elements of volume, but also by movements of the surfaces of discontinuity. It would therefore be necessary to add to our general condition of equilibrium terms relating to discontinuous changes in the elements of volume about these surfaces,—a necessity which is avoided if we consider these elements movable, as we can then suppose that each element remains always on the same side of the surface of discontinuity.

*Method of treating the preceding problem, in which the elements of volume are regarded as fixed.*

It may be interesting to see in detail how the particular conditions of equilibrium may be obtained if we regard the elements of volume as fixed in position and size, and consider the possibility of finite as well as infinitesimal changes of phase in each element of volume. If we use the character  $\Delta$  to denote the differences determined by such finite differences of phase, we may express the variation of the intrinsic energy of the whole mass in the form

$$\int' \delta D\varepsilon + \int' \Delta D\varepsilon, \quad (237)$$

in which the first integral extends over all the elements which are infinitesimally varied, and the second over all those which experience a finite variation. We *may* regard both integrals as extending throughout the whole mass, but their values will be zero except for the parts mentioned.

If we do not wish to limit ourselves to the consideration of masses so small that the force of gravity can be regarded as constant in direction and in intensity, we may use  $\Upsilon$  to denote the potential of the force of gravity, and express the variation of the part of the energy which is due to gravity in the form

$$-\int' \Upsilon \delta Dm - \int' \Upsilon \Delta Dm. \quad (238)$$

We shall then have, for the general condition of equilibrium,

$$\int \delta D\varepsilon + \int \Delta D\varepsilon - \int \Gamma \delta Dm - \int \Gamma \Delta Dm \geq 0; \quad (239)$$

and the equations of condition will be

$$\int \delta D\eta + \int \Delta D\eta = 0, \quad (240)$$

$$\left. \begin{array}{l} \int \delta Dm_1 + \int \Delta Dm_1 = 0, \\ \cdot \quad \cdot \quad \cdot \quad \cdot \quad \cdot \quad \cdot \end{array} \right\} \quad (241)$$

$$\int \delta Dm_n + \int \Delta Dm_n = 0.$$

We may obtain a condition of equilibrium independent of these equations of condition, by subtracting these equations, multiplied each by an indeterminate constant, from condition (239). If we denote these indeterminate constants by  $T, M_1, \dots, M_n$ , we shall obtain after arranging the terms

$$\begin{aligned} & \int \delta D\varepsilon - \Gamma \delta Dm - T \delta D\eta - M_1 \delta Dm_1 \dots - M_n \delta Dm_n \\ & + \int \Delta D\varepsilon - \Gamma \Delta Dm - T \Delta D\eta - M_1 \Delta Dm_1 \dots - M_n \Delta Dm_n \geq 0. \end{aligned} \quad (242)$$

The variations, both infinitesimal and finite, in this condition are independent of the equations of condition (240) and (241), and are only subject to the condition that the varied values of  $D\varepsilon, D\eta, Dm_1, \dots, Dm_n$  for each element are determined by a certain change of phase. But as we do not suppose the same element to experience both a finite and an infinitesimal change of phase, we must have

$$\delta D\varepsilon - \Gamma \delta Dm - T \delta D\eta - M_1 \delta Dm_1 \dots - M_n \delta Dm_n \geq 0, \quad (243)$$

and

$$\Delta D\varepsilon - \Gamma \Delta Dm - T \Delta D\eta - M_1 \Delta Dm_1 \dots - M_n \Delta Dm_n \geq 0. \quad (244)$$

By equation (12), and in virtue of the necessary relation (222), the first of these conditions reduces to

$$\begin{aligned} & (t - T) \delta D\eta + (\mu_1 - \Gamma - M_1) \delta Dm_1 \dots \\ & + (\mu_n - \Gamma - M_n) \delta Dm_n \geq 0; \end{aligned} \quad (245)$$

for which it is necessary and sufficient that

$$t = T, \quad (246)$$

$$\left. \begin{array}{l} \mu_1 - \Gamma = M_1, \\ \cdot \quad \cdot \quad \cdot \quad \cdot \quad \cdot \quad \cdot \end{array} \right\} * \quad (247)$$

$$\mu_n - \Gamma = M_n.$$

\* The gravitation potential is here supposed to be defined in the usual way. But if it were defined so as to decrease when a body falls, we would have the sign + instead of - in these equations; i. e., for each component, the sum of the gravitation and intrinsic potentials would be constant throughout the whole mass.

Condition (244) may be reduced to the form

$$\Delta D\varepsilon - T \Delta D\eta - (T + M_1) \Delta Dm_1 \dots - (T + M_n) \Delta Dm_n \geq 0; \quad (248)$$

and by (246) and (247) to

$$\Delta D\varepsilon - t \Delta D\eta - \mu_1 \Delta Dm_1 \dots - \mu_n \Delta Dm_n \geq 0. \quad (249)$$

If values determined subsequently to the change of phase are distinguished by accents, this condition may be written

$$\begin{aligned} D\varepsilon' - t D\eta' - \mu_1' Dm_1' \dots - \mu_n' Dm_n' \\ - D\varepsilon + t D\eta + \mu_1 Dm_1 \dots + \mu_n Dm_n \geq 0, \end{aligned} \quad (250)$$

which may be reduced by (93) to

$$D\varepsilon' - t D\eta' - \mu_1' Dm_1' \dots - \mu_n' Dm_n' + p Dv \geq 0. \quad (251)$$

Now if the element of volume  $Dv$  is adjacent to a surface of discontinuity, let us suppose  $D\varepsilon'$ ,  $D\eta'$ ,  $Dm_1'$ , ...,  $Dm_n'$  to be determined (for the same element of volume) by the phase existing on the other side of the surface of discontinuity. As  $t$ ,  $\mu_1$ , ...,  $\mu_n$  have the same values on both sides of this surface, the condition may be reduced by (93) to

$$-p'Dv + pDv \geq 0. \quad (252)$$

That is, the pressure must not be greater on one side of a surface of discontinuity than on the other.

Applied more generally, (251) expresses the condition of equilibrium with respect to the possibility of discontinuous changes of phases at any point. As  $Dv' = Dv$ , the condition may also be written

$$D\varepsilon' - t D\eta' + p Dv' - \mu_1' Dm_1' \dots - \mu_n' Dm_n' \geq 0, \quad (253)$$

which must hold true when  $t$ ,  $p$ ,  $\mu_1$ , ...,  $\mu_n$  have values determined by any point in the mass, and  $D\varepsilon'$ ,  $D\eta'$ ,  $Dv'$ ,  $Dm_1'$ , ...,  $Dm_n'$  have values determined by any possible phase of the substances of which the mass is composed. The application of the condition is, however, subject to the limitations considered on pages 128–134. It may easily be shown (see pages 160, 161) that for constant values of  $t$ ,  $\mu_1$ , ...,  $\mu_n$ , and of  $Dv'$ , the first member of (253) will have the least possible value when  $D\varepsilon'$ ,  $D\eta'$ ,  $Dm_1'$ , ...,  $Dm_n'$  are determined by a phase for which the temperature has the value  $t$ , and the potentials the values  $\mu_1$ , ...,  $\mu_n$ . It will be sufficient, therefore, to consider the condition as applied to such phases, in which case it may be reduced by (93) to

$$p - p' \geq 0. \quad (254)$$

That is, the pressure at any point must be as great as that of any phase of the same components, for which the temperature and the

potentials have the same values as at that point. We may also express this condition by saying that the pressure must be as great as is consistent with equations (246), (247). This condition with the equations mentioned will always be sufficient for equilibrium; when the condition is not satisfied, if equilibrium subsists, it will be at least practically unstable.

Hence, the phase at any point of a fluid mass, which is in stable equilibrium under the influence of gravity (whether this force is due to external bodies or to the mass itself), and which has throughout the same independently variable components, is completely determined by the phase at any other point and the difference of the values of the gravitation potential for the two points.

#### FUNDAMENTAL EQUATIONS OF IDEAL GASES AND GAS-MIXTURES.—

For a constant quantity of a perfect or ideal gas, the product of the volume and pressure is proportional to the temperature, and the variations of energy are proportional to the variations of temperature. For a unit of such a gas we may write

$$\begin{aligned} p v &= a t, \\ d\varepsilon &= c dt, \end{aligned}$$

$a$  and  $c$  denoting constants. By integration, we obtain the equation  
 $\varepsilon = c t + E,$

in which  $E$  also denotes a constant. If by these equations we eliminate  $t$  and  $p$  from (11), we obtain

$$d\varepsilon = \frac{\varepsilon - E}{c} d\eta - \frac{a}{v} \frac{\varepsilon - E}{c} dv,$$

or

$$c \frac{d\varepsilon}{\varepsilon - E} = d\eta - a \frac{dv}{v}.$$

The integral of this equation may be written in the form

$$c \log \frac{\varepsilon - E}{c} = \eta - a \log v - H,$$

where  $H$  denotes a fourth constant. We may regard  $E$  as denoting the energy of a unit of the gas for  $t=0$ ;  $H$  its entropy for  $t=1$  and  $v=1$ ;  $a$  its pressure in the latter state, or its volume for  $t=1$  and  $p=1$ ;  $c$  its specific heat at constant volume. We may extend the application of the equation to any quantity of the gas, without altering the values of the constants, if we substitute  $\frac{\varepsilon}{m}, \frac{\eta}{m}, \frac{v}{m}$  for  $\varepsilon, \eta, v$ , respectively. This will give

$$c \log \frac{\varepsilon - Em}{cm} = \frac{\eta}{m} - H + a \log \frac{m}{v}. \quad (255)$$

This is a fundamental equation (see pages 140–144) for an ideal gas of invariable composition. It will be observed that if we do not have to consider the properties of the matter which forms the gas as appearing in any other form or combination, but solely as constituting the gas in question (in a state of purity), we may without loss of generality give to  $E$  and  $H$  the value zero, or any other arbitrary values. But when the scope of our investigations is not thus limited, we may have determined the states of the substance of the gas for which  $\varepsilon=0$  and  $\eta=0$  with reference to some other form in which the substance appears, or, if the substance is compound, the states of its components for which  $\varepsilon=0$  and  $\eta=0$  may be already determined; so that the constants  $E$  and  $H$  cannot in general be treated as arbitrary.

We obtain from (255) by differentiation

$$\frac{c}{\varepsilon - Em} d\varepsilon = \frac{1}{m} d\eta - \frac{a}{v} dv + \left( \frac{cE}{\varepsilon - Em} + \frac{c+a}{m} - \frac{\eta}{m^2} \right) dm, \quad (256)$$

whence, in virtue of the general relation expressed by (86),

$$t = \frac{\varepsilon - Em}{cm}, \quad (257)$$

$$p = a \frac{\varepsilon - Em}{cv}, \quad (258)$$

$$\mu = E + \frac{\varepsilon - Em}{cm^2}(cm + am - \eta). \quad (259)$$

We may obtain the fundamental equation between  $\psi$ ,  $t$ ,  $v$ , and  $m$  from equations (87), (255), and (257). Eliminating  $\varepsilon$  we have

$$\psi = Em + cm t - t \eta,$$

$$\text{and } c \log t = \frac{\eta}{m} - H + a \log \frac{m}{v};$$

and eliminating  $\eta$ , we have the fundamental equation

$$\psi = Em + mt \left( c - H - c \log t + a \log \frac{m}{v} \right). \quad (260)$$

Differentiating this equation, we obtain

$$d\psi = -m \left( H + c \log t + a \log \frac{v}{m} \right) dt - \frac{amt}{v} dv \\ + \left( E + t \left( c + a - H - c \log t + a \log \frac{m}{v} \right) \right) dm; \quad (261)$$

whence, by the general equation (88),

$$\eta = m \left( H + c \log t + a \log \frac{v}{m} \right), \quad (262)$$

$$p = \frac{amt}{v}, \quad (263)$$

$$\mu = E + t \left( c + a - H - c \log t + a \log \frac{m}{v} \right). \quad (264)$$

From (260), by (87) and (91), we obtain

$$\zeta = Em + mt \left( c - H - c \log t + a \log \frac{m}{v} \right) + p v,$$

and eliminating  $v$  by means of (263), we obtain the fundamental equation

$$\zeta = Em + mt \left( c + a - H - (c+a) \log t + a \log \frac{p}{a} \right). \quad (265)$$

From this, by differentiation and comparison with (92), we may obtain the equations

$$\eta = m \left( H + (c+a) \log t + a \log \frac{p}{a} \right), \quad (266)$$

$$v = \frac{amt}{p}, \quad (267)$$

$$\mu = E + t \left( c + a - H - (c+a) \log t + a \log \frac{p}{a} \right). \quad (268)$$

The last is also a fundamental equation. It may be written in the form

$$\log \frac{p}{a} = \frac{H-c-a}{a} + \frac{c+a}{a} \log t + \frac{\mu-E}{at}, \quad (269)$$

or, if we denote by  $e$  the base of the Naperian system of logarithms,

$$p = a e^{\frac{H-c-a}{a}} t^{\frac{c+a}{a}} e^{\frac{\mu-E}{at}} \quad (270)$$

The fundamental equation between  $\chi$ ,  $\eta$ ,  $p$ , and  $m$  may also be easily obtained; it is

$$(c+a) \log \frac{\chi - Em}{(c+a)m} = \frac{\eta}{m} - H + a \log \frac{p}{a}, \quad (271)$$

which can be solved with respect to  $\chi$ .

Any one of the fundamental equations (255), (260), (265), (270), and (271), which are entirely equivalent to one another, may be

regarded as defining an ideal gas. It will be observed that most of these equations might be abbreviated by the use of different constants. In (270), for example, a single constant might be used for

$\frac{H-c-a}{a}e^{-\frac{a}{a}}$ , and another for  $\frac{c+a}{a}$ . The equations have been given

in the above form, in order that the relations between the constants occurring in the different equations might be most clearly exhibited. The sum  $c+a$  is the specific heat for constant pressure, as appears if we differentiate (266) regarding  $p$  and  $m$  as constant.\*

\* We may easily obtain the equation between the temperature and pressure of a saturated vapor, if we know the fundamental equations of the substance both in the gaseous, and in the liquid or solid state. If we suppose that the density and the specific heat at constant pressure of the liquid may be regarded as constant quantities (for such moderate pressures as the liquid experiences while in contact with the vapor), and denote this specific heat by  $k$ , and the volume of a unit of the liquid by  $V$ , we shall have for a unit of the liquid

$$t \, dy = k \, dt,$$

whence

$$y = k \log t + H',$$

where  $H'$  denotes a constant. Also, from this equation and (97),

$$d\mu = -(k \log t + H') \, dt + V \, dp,$$

whence

$$\mu = kt - kt \log t - H' t + Vp + E', \quad (\text{A})$$

where  $E'$  denotes another constant. This is a fundamental equation for the substance in the liquid state. If (268) represents the fundamental equation for the same substance in the gaseous state, the two equations will both hold true of coexistent liquid and gas. Eliminating  $\mu$  we obtain

$$\log \frac{p}{a} = \frac{H - H' + k - c - a}{a} - \frac{k - c - a}{a} \log t - \frac{E - E'}{at} + \frac{V}{a} \cdot \frac{p}{t}.$$

If we neglect the last term, which is evidently equal to the density of the vapor divided by the density of the liquid, we may write

$$\log p = A - B \log t - \frac{C}{t},$$

$A$ ,  $B$ , and  $C$  denoting constants. If we make similar suppositions in regard to the substance in the solid state, the equation between the pressure and temperature of coexistent solid and gaseous phases will of course have the same form.

A similar equation will also apply to the phases of an ideal gas which are coexistent with two different kinds of solids, one of which can be formed by the combination of the gas with the other, each being of invariable composition and of constant specific heat and density. In this case we may write for one solid

$$\mu_1 = k't - k't \log t - Ht + V'p + E,$$

and for the other

$$\mu_2 = k''t - k''t \log t - H''t + V''p + E'',$$

and for the gas

$$\mu_3 = E + t \left( c + a - H - (c + a) \log t + a \log \frac{p}{a} \right).$$

The preceding fundamental equations all apply to gases of *constant composition*, for which the matter is entirely determined by a single

Now if a unit of the gas unites with the quantity  $\lambda$  of the first solid to form the quantity  $1+\lambda$  of the second it will be necessary for equilibrium (see pages 121, 122) that

$$\mu_3 + \lambda \mu_1 = (1+\lambda) \mu_2.$$

Substituting the values of  $\mu_1$ ,  $\mu_2$ ,  $\mu_3$ , given above, we obtain after arranging the terms and dividing by  $at$

$$\log \frac{p}{a} = A - B \log t - \frac{C}{t} + D \frac{p}{t},$$

when

$$A = \frac{H + \lambda H' - (1+\lambda) H'' - c - a - \lambda k' + (1+\lambda) k''}{a},$$

$$B = \frac{(1+\lambda) k'' - \lambda k' - c - a}{a},$$

$$C = \frac{E + \lambda E' - (1+\lambda) E''}{a}, \quad D = \frac{(1+\lambda) V'' - \lambda V}{a}.$$

We may conclude from this that an equation of the same form may be applied to an ideal gas in equilibrium with a liquid of which it forms an independently variable component, when the specific heat and density of the liquid are entirely determined by its composition, except that the letters  $A$ ,  $B$ ,  $C$ , and  $D$  must in this case be understood to denote quantities which vary with the composition of the liquid. But to consider the case more in detail, we have for the liquid by (a)

$$\frac{\zeta}{m} = \mu - kt - kt \log t - H't + Vp + E',$$

where  $k$ ,  $H'$ ,  $V$ ,  $E'$  denote quantities which depend only upon the composition of the liquid. Hence, we may write

$$\zeta = \mathbf{k}t - \mathbf{k}t \log t - \mathbf{H}t + \mathbf{V}p + \mathbf{E},$$

where  $\mathbf{k}$ ,  $\mathbf{H}$ ,  $\mathbf{V}$ , and  $\mathbf{E}$  denote functions of  $m_1$ ,  $m_2$ , etc. (the quantities of the several components of the liquid). Hence, by (92),

$$\mu_1 = \frac{d\mathbf{k}}{dm_1} t - \frac{d\mathbf{k}}{dm_1} t \log t - \frac{d\mathbf{H}}{dm_1} t + \frac{d\mathbf{V}}{dm_1} p + \frac{d\mathbf{E}}{dm_1}.$$

If the component to which this potential relates is that which also forms the gas, we shall have by (269)

$$\log \frac{p}{a} = \frac{H - c - a}{a} + \frac{c + a}{a} \log t + \frac{\mu_1 - E}{at}.$$

Eliminating  $\mu_1$ , we obtain the equation

$$\log \frac{p}{a} = A - B \log t - \frac{C}{t} + D \frac{p}{t},$$

in which  $A$ ,  $B$ ,  $C$ , and  $D$  denote quantities which depend only upon the composition of the liquid, viz:

$$A = \frac{1}{a} \left( H - \frac{d\mathbf{H}}{dm_1} - c - a + \frac{d\mathbf{k}}{dm_1} \right),$$

$$B = \frac{1}{a} \left( \frac{d\mathbf{k}}{dm_1} - c - a \right).$$

variable ( $m$ ). We may obtain corresponding fundamental equations for a mixture of gases, in which the proportion of the components shall be variable, from the following considerations.

It is a rule which admits of a very general and in many cases very exact experimental verification, that if several liquid or solid substances which yield different gases or vapors are simultaneously in equilibrium with a mixture of these gases (cases of chemical action between the gases being excluded,) the pressure in the gas-mixture is equal to the sum of the pressures of the gases yielded at the same temperature by the various liquid or solid substances taken separately. Now the potential in any of the liquids or solids for the substance which it yields in the form of gas has very nearly the same value when the liquid or solid is in equilibrium with the gas-mixture as when it is in equilibrium with its own gas alone. The difference of the pressure in the two cases will cause a certain difference in the values of the potential, but that this difference will be small, we may infer from the equation

$$\left( \frac{d\mu_1}{dp} \right)_{t, m} = \left( \frac{dv}{dm_1} \right)_{t, p, m}, \quad (272)$$

which may be derived from equation (92). In most cases, there will be a certain absorption by each liquid of the gases yielded by the others, but as it is well known that the above rule does not apply to cases in which such absorption takes place to any great extent, we may conclude that the effect of this circumstance in the cases with which we have to do is of secondary importance. If we neglect the slight differences in the values of the potentials due to these circumstances, the rule may be expressed as follows:

*The pressure in a mixture of different gases is equal to the sum of the pressures of the different gases as existing each by itself at the same temperature and with the same value of its potential.*

To form a precise idea of the practical significance of the law as thus stated with reference to the equilibrium of two liquids with a mixture of the gases which they emit, when neither liquid absorbs the gas emitted by the other, we may imagine a long tube closed at each end and bent in the form of a W to contain in each of the descending

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$$C = \frac{1}{a} \left( E - \frac{d\mathbf{E}}{dm_1} \right) \quad D = \frac{1}{a} \frac{d\mathbf{V}}{dm_1}.$$

With respect to some of the equations which have here been deduced, the reader may compare Professor Kirchhoff "Ueber die Spannung des Dampfes von Mischungen aus Wasser und Schwefelsäure," Pogg. Ann., vol. civ. (1858), p. 612; and Dr. Rankine "On Saturated Vapors," Phil. Mag., vol. xxxi. (1866), p. 199.

loops one of the liquids, and above these liquids the gases which they emit, viz., the separate gases at the ends of the tube, and the mixed gases in the middle. We may suppose the whole to be in equilibrium, the difference of the pressures of the gases being balanced by the proper heights of the liquid columns. Now it is evident from the principles established on pages 203–210 that the potential for either gas will have the same value in the mixed and in the separate gas *at the same level*, and therefore according to the rule in the form which we have given, the pressure in the gas-mixture is equal to the sum of the pressures in the separate gases, *all these pressures being measured at the same level*. Now the experiments by which the rule has been established relate rather to the gases in the vicinity of the surfaces of the liquids. Yet, although the differences of level in these surfaces may be considerable, the corresponding differences of pressure in the columns of gas will certainly be very small in all cases which can be regarded as falling under the laws of ideal gases, for which very great pressures are not admitted.

If we apply the above law to a mixture of ideal gases and distinguish by subscript numerals the quantities relating to the different gases, and denote by  $\Sigma_1$  the sum of all similar terms obtained by changing the subscript numerals, we shall have by (270)

$$p = \Sigma_1 \left( a_1 e^{\frac{H_1 - c_1 - a_1}{a_1}} t^{\frac{c_1 + a_1}{a_1}} e^{\frac{\mu_1 - E_1}{a_1 t}} \right), \quad (273)$$

It will be legitimate to assume this equation provisionally as the fundamental equation defining an ideal gas-mixture, and afterwards to justify the suitableness of such a definition by the properties which may be deduced from it. In particular, it will be necessary to show that an ideal gas-mixture as thus defined, when the proportion of its components remains constant, has all the properties which have already been assumed for an ideal gas of invariable composition; it will also be desirable to consider more rigorously and more in detail the equilibrium of such a gas-mixture with solids and liquids, with respect to the above rule.

By differentiation and comparison with (98) we obtain

$$\frac{\eta}{v} = \Sigma_1 \left( (c_1 + a_1 - \frac{\mu_1 - E_1}{t}) e^{\frac{H_1 - c_1 - a_1}{a_1}} t^{\frac{c_1}{a_1}} e^{\frac{\mu_1 - E_1}{a_1 t}} \right), \quad (274)$$

$$\left. \begin{aligned} \frac{m_1}{v} &= e^{\frac{H_1 - c_1 - a_1}{a_1}} t^{\frac{c_1}{a_1}} e^{\frac{\mu_1 - E_1}{a_1 t}}, \\ \frac{m_2}{v} &= e^{\frac{H_2 - c_2 - a_2}{a_2}} t^{\frac{c_2}{a_2}} e^{\frac{\mu_2 - E_2}{a_2 t}}, \\ &\text{etc.} \end{aligned} \right\} \quad (275)$$

Equations (275) indicate that the relation between the temperature, the density of any component, and the potential for that component, is not affected by the presence of the other components. They may also be written

$$\left. \begin{aligned} \mu_1 &= E_1 + t \left( c_1 + a_1 - H_1 - c_1 \log t + a_1 \log \frac{m_1}{v} \right), \\ &\text{etc.} \end{aligned} \right\} \quad (276)$$

Eliminating  $\mu_1$ ,  $\mu_2$ , etc. from (273) and (274) by means of (275) and (276), we obtain

$$p = \Sigma_1 \frac{a_1 m_1 t}{v}, \quad (277)$$

$$\eta = \Sigma_1 \left( m_1 H_1 + m_1 c_1 \log t + m_1 a_1 \log \frac{v}{m_1} \right). \quad (278)$$

Equation (277) expresses the familiar principle that the pressure in a gas-mixture is equal to the sum of the pressures which the component gases would possess if existing separately with the same volume at the same temperature. Equation (278) expresses a similar principle in regard to the entropy of the gas-mixture.

From (276) and (277) we may easily obtain the fundamental equation between  $\psi$ ,  $t$ ,  $v$ ,  $m_1$ ,  $m_2$ , etc. For by substituting in (94) the values of  $p$ ,  $\mu_1$ ,  $\mu_2$ , etc. taken from these equations, we obtain

$$\psi = \Sigma_1 \left( E_1 m_1 + m_1 t \left( c_1 - H_1 - c_1 \log t + a_1 \log \frac{m_1}{v} \right) \right). \quad (279)$$

If we regard the proportion of the various components as constant, this equation may be simplified by writing

$$\begin{aligned} m &\text{ for } \Sigma_1 m_1, \\ c m &\text{ for } \Sigma_1 (c_1 m_1), \\ a m &\text{ for } \Sigma_1 (a_1 m_1), \\ E m &\text{ for } \Sigma_1 (E_1 m_1), \end{aligned}$$

and  $H m - a m \log m$  for  $\Sigma_1 (H_1 m_1 - a_1 m_1 \log m_1)$ .

The values of  $c$ ,  $a$ ,  $E$ , and  $H$ , will then be constant and  $m$  will denote the total quantity of gas. As the equation will thus be reduced to the

form of (260), it is evident that an ideal gas-mixture, as defined by (273) or (279), when the proportion of its components remains unchanged, will have all the properties which we have assumed for an ideal gas of invariable composition. The relations between the specific heats of the gas-mixture at constant volume and at constant pressure and the specific heats of its components are expressed by the equations

$$c = \sum_1 \frac{m_1 c_1}{m}, \quad (280)$$

and

$$c + a = \sum_1 \frac{m_1 (c_1 + a_1)}{m}. \quad (281)$$

We have already seen that the values of  $t$ ,  $v$ ,  $m_1$ ,  $\mu_1$  in a gas-mixture are such as are possible for the component  $G_1$  (to which  $m_1$  and  $\mu_1$  relate) existing separately. If we denote by  $p_1$ ,  $\eta_1$ ,  $\psi_1$ ,  $\varepsilon_1$ ,  $\chi_1$ ,  $\zeta_1$  the connected values of the several quantities which the letters indicate determined for the gas  $G_1$  as thus existing separately, and extend this notation to the other components, we shall have by (273), (274), and (279)

$$p = \sum_1 p_1, \quad \eta = \sum_1 \eta_1, \quad \psi = \sum_1 \psi_1; \quad (282)$$

whence by (87), (89), and (91)

$$\varepsilon = \sum_1 \varepsilon_1, \quad \chi = \sum_1 \chi_1, \quad \zeta = \sum_1 \zeta_1. \quad (283)$$

The quantities  $p$ ,  $\eta$ ,  $\psi$ ,  $\varepsilon$ ,  $\chi$ ,  $\zeta$  relating to the gas-mixture may therefore be regarded as consisting of parts which may be attributed to the several components in such a manner that between the parts of these quantities which are assigned to any component, the quantity of that component, the potential for that component, the temperature, and the volume, the same relations shall subsist as if that component existed separately. It is in this sense that we should understand the law of Dalton, that every gas is as a vacuum to every other gas.

It is to be remarked that these relations are consistent and possible for a mixture of gases which are not ideal gases, and indeed without any limitation in regard to the thermodynamic properties of the individual gases. They are all consequences of the law that the pressure in a mixture of different gases is equal to the sum of the pressures of the different gases as existing each by itself at the same temperature and with the same value of its potential. For let  $p_1$ ,  $\eta_1$ ,  $\varepsilon_1$ ,  $\psi_1$ ,  $\chi_1$ ,  $\zeta_1$ ;  $p_2$ , etc.; etc. be defined as relating to the different gases existing each by itself with the same volume, temperature, and potential as in the gas-mixture; if

$$p = \sum_1 p_1,$$

then

$$\left(\frac{dp}{d\mu_1}\right)_{t, \mu_2 \dots \mu_n} = \left(\frac{dp_1}{d\mu_1}\right)_t;$$

and therefore, by (98), the quantity of any component gas  $G_1$  in the gas-mixture, and in the separate gas to which  $p_1$ ,  $\eta_1$ , etc. relate, is the same and may be denoted by the same symbol  $m_1$ . Also

$$\eta = v \left(\frac{dp}{dt}\right)_{\mu_1, \dots, \mu_n} = v \Sigma_1 \left(\frac{dp}{dt}\right)_{\mu_1} = \Sigma_1 \eta_1;$$

whence also, by (93)–(96),

$$\varepsilon = \Sigma_1 \varepsilon_1, \quad \psi = \Sigma_1 \psi_1, \quad \chi = \Sigma_1 \chi_1, \quad \zeta = \Sigma_1 \zeta_1.$$

All the same relations will also hold true whenever the value of  $\psi$  for the gas-mixture is equal to the sum of the values of this function for the several component gases existing each by itself in the same quantity as in the gas-mixture and with the temperature and volume of the gas-mixture. For if  $p_1$ ,  $\eta_1$ ,  $\varepsilon_1$ ,  $\psi_1$ ,  $\chi_1$ ,  $\zeta_1$ ;  $p_2$ , etc.; etc. are defined as relating to the components existing thus by themselves, we shall have

$$\psi = \Sigma_1 \psi_1,$$

whence

$$\left(\frac{d\psi}{dm_1}\right)_{t, v, m} = \left(\frac{d\psi_1}{dm_1}\right)_{t, v}^*$$

Therefore, by (88), the potential  $\mu_1$  has the same value in the gas-mixture and in the gas  $G_1$  existing separately as supposed. Moreover,

$$\eta = - \left(\frac{d\psi}{dt}\right)_{v, m} = - \Sigma_1 \left(\frac{d\psi_1}{dt}\right)_{v, m} = \Sigma_1 \eta_1,$$

$$\text{and } p = - \left(\frac{d\psi}{dv}\right)_{t, m} = - \Sigma_1 \left(\frac{d\psi_1}{dv}\right)_{t, m} = \Sigma_1 p_1,$$

whence

$$\varepsilon = \Sigma_1 \varepsilon_1, \quad \chi = \Sigma_1 \chi_1, \quad \zeta = \Sigma_1 \zeta_1.$$

Whenever different bodies are combined without communication of work or heat between them and external bodies, the energy of the body formed by the combination is necessarily equal to the sum of the energies of the bodies combined. In the case of ideal gas-mixtures, when the initial temperatures of the gas-masses which are combined

\* A subscript  $m$  after a differential coefficient relating to a body having several independently variable components is used here and elsewhere in this paper to indicate that each of the quantities  $m_1$ ,  $m_2$ , etc., unless its differential occurs in the expression to which the suffix is applied, is to be regarded as constant in the differentiation.

are the same, (whether these gas-masses are entirely different gases, or gas-mixtures differing only in the proportion of their components,) the condition just mentioned can only be satisfied when the temperature of the resultant gas-mixture is also the same. In such combinations, therefore, the final temperature will be the same as the initial.

If we consider a vertical column of an ideal gas-mixture which is in equilibrium, and denote the densities of one of its components at two different points by  $\gamma_1$  and  $\gamma_1'$ , we shall have by (275) and (234)

$$\frac{\gamma_1'}{\gamma_1} = e^{\frac{\mu_1 - \mu_1'}{a_1 t}} = e^{\frac{g(h' - h)}{a_1 t}}. \quad (284)$$

From this equation, in which we may regard the quantities distinguished by accents as constant, it appears that the relation between the density of any one of the components and the height is not affected by the presence of the other components.

The work obtained or expended in any reversible process of combination or separation of ideal gas-mixtures at constant temperature, or when the temperatures of the initial and final gas-masses and of the only external source of heat or cold which is used are all the same, will be found by taking the difference of the sums of the values of  $\psi$  for the initial, and for the final gas-masses. (See pages 145, 146). It is evident from the form of equation (279) that this work is equal to the sum of the quantities of work which would be obtained or expended in producing in each different component existing separately the same changes of density which that component experiences in the actual process for which the work is sought.\*

We will now return to the consideration of the equilibrium of a liquid with the gas which it emits as affected by the presence of different gases, when the gaseous mass in contact with the liquid may be regarded as an ideal gas-mixture.

It may first be observed, that the density of the gas which is emitted by the liquid will not be affected by the presence of other gases which are not absorbed by the liquid, when the liquid is protected in any way from the pressure due to these additional gases. This may be accomplished by separating the liquid and gaseous

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\* This result has been given by Lord Rayleigh, (Phil. Mag., vol. xlix, 1875, p. 311). It will be observed that equation (279) might be deduced immediately from this principle in connection with equation (260) which expresses the properties ordinarily assumed for perfect gases.

masses by a diaphragm which is permeable to the liquid. It will then be easy to maintain the liquid at any constant pressure which is not greater than that in the gas. The potential in the liquid for the substance which it yields as gas will then remain constant, and therefore the potential for the same substance in the gas and the density of this substance in the gas and the part of the gaseous pressure due to it will not be affected by the other components of the gas.

But when the gas and liquid meet under ordinary circumstances, i. e., in a free plane surface, the pressure in both is necessarily the same, as also the value of the potential for any common component  $S_1$ . Let us suppose the density of an insoluble component of the gas to vary, while the composition of the liquid and the temperature remain unchanged. If we denote the increments of pressure and of the potential for  $S_1$  by  $d\mu$  and  $d\mu_1$ , we shall have by (272)

$$d\mu_1 = \left( \frac{d\mu_1}{dp} \right)_{t, m}^{(L)} dp = \left( \frac{dv}{dm_1} \right)_{t, p, m}^{(L)} dp,$$

the index (L) denoting that the expressions to which it is affixed refer to the liquid. (Expressions without such an index will refer to the gas alone or to the gas and liquid in common.) Again, since the gas is an ideal gas-mixture, the relation between  $p_1$  and  $\mu_1$  is the same as if the component  $S_1$  existed by itself at the same temperature, and therefore by (268)

$$d\mu_1 = \alpha_1 t d \log p_1.$$

Therefore

$$\alpha_1 t d \log p_1 = \left( \frac{dv}{dm_1} \right)_{t, p, m}^{(L)} dp. \quad (285)$$

This may be integrated at once if we regard the differential coefficient in the second member as constant, which will be a very close approximation. We may obtain a result more simple, but not quite so accurate, if we write the equation in the form

$$dp_1 = \gamma_1 \left( \frac{dv}{dm_1} \right)_{t, p, m}^{(L)} dp, \quad (286)$$

where  $\gamma_1$  denotes the density of the component  $S_1$  in the gas, and integrate regarding this quantity also as constant. This will give

$$p_1 - p_1' = \gamma_1 \left( \frac{dv}{dm_1} \right)_{t, p, m}^{(L)} (p - p'). \quad (287)$$

where  $p_1'$  and  $p'$  denote the values of  $p_1$  and  $p$  when the insoluble component of the gas is entirely wanting. It will be observed that  $p - p'$  is nearly equal to the pressure of the insoluble component, in the phase of the gas-mixture to which  $p_1$  relates.  $S_1$  is not neces-

sarily the only common component of the gas and liquid. If there are others, we may find the increase of the part of the pressure in the gas-mixture belonging to any one of them by equations differing from the last only in the subscript numerals.

Let us next consider the effect of a gas which is absorbed to some extent, and which must therefore in strictness be regarded as a component of the liquid. We may commence by considering in general the equilibrium of a gas-mixture of two components  $S_1$  and  $S_2$  with a liquid formed of the same components. Using a notation like the previous, we shall have by (98) for constant temperature,

$$dp = \gamma_1 d\mu_1 + \gamma_2 d\mu_2,$$

and

$$dp = \gamma_1^{(L)} d\mu_1 + \gamma_2^{(L)} d\mu_2;$$

whence

$$(\gamma_1^{(L)} - \gamma_1) d\mu_1 = (\gamma_2 - \gamma_2^{(L)}) d\mu_2.$$

Now if the gas is an ideal gas-mixture,

$$d\mu_1 = \frac{a_1 t}{p_1} dp_1 = \frac{dp_1}{\gamma_1}, \quad \text{and} \quad d\mu_2 = \frac{a_2 t}{p_2} dp_2 = \frac{dp_2}{\gamma_2},$$

therefore

$$\left( \frac{\gamma_1^{(L)}}{\gamma_1} - 1 \right) dp_1 = \left( 1 - \frac{\gamma_2^{(L)}}{\gamma_2} \right) dp_2. \quad (288)$$

We may now suppose that  $S_1$  is the principal component of the liquid, and  $S_2$  is a gas which is absorbed in the liquid to a slight extent. In such cases it is well known that the ratio of the densities of the substance  $S_2$  in the liquid and in the gas is for a given temperature approximately constant. If we denote this constant by  $A$ , we shall have

$$\left( \frac{\gamma_1^{(L)}}{\gamma_1} - 1 \right) dp_1 = (1 - A) dp_2. \quad (289)$$

It would be easy to integrate this equation regarding  $\gamma_1$  as variable, but as the variation in the value of  $p_1$  is necessarily very small we shall obtain sufficient accuracy if we regard  $\gamma_1$  as well as  $\gamma_1^{(L)}$  as constant. We shall thus obtain

$$\left( \frac{\gamma_1^{(L)}}{\gamma_1} - 1 \right) (p_1 - p_1') = (1 - A) p_2, \quad (290)$$

where  $p_1'$  denotes the pressure of the saturated vapor of the pure liquid consisting of  $S_1$ . It will be observed that when  $A = 1$ , the presence of the gas  $S_2$  will not affect the pressure or density of the gas  $S_1$ . When  $A < 1$ , the pressure and density of the gas  $S_1$  are greater than if  $S_2$  were absent, and when  $A > 1$ , the reverse is true.

The properties of an ideal gas-mixture (according to the definition which we have assumed) when in equilibrium with liquids or solids have been developed at length, because it is only in respect to these properties that there is any variation from the properties usually attributed to perfect gases. As the pressure of a gas saturated with vapor is usually given as a little less than the sum of the pressure of the gas calculated from its density and that of saturated vapor in a space otherwise empty, while our formulæ would make it a little more, when the gas is insoluble, it would appear that in this respect our formulæ are less accurate than the rule which would make the pressure of the gas saturated with vapor equal to the sum of the two pressures mentioned. Yet the reader will observe that the magnitude of the quantities concerned is not such that any stress can be laid upon this circumstance.

It will also be observed that the statement of Dalton's law which we have adopted, while it serves to complete the theory of gas-mixtures (with respect to a certain class of properties), asserts nothing with reference to any solid or liquid bodies. But the common rule that the density of a gas necessary for equilibrium with a solid or liquid is not altered by the presence of a different gas which is not absorbed by the solid or liquid, if construed *strictly*, will involve consequences in regard to solids and liquids which are entirely inadmissible. To show this, we will assume the correctness of the rule mentioned. Let  $S_1$  denote the common component of the gaseous and liquid or solid masses, and  $S_2$  the insoluble gas, and let quantities relating to the gaseous mass be distinguished when necessary by the index (G), and those relating to the liquid or solid by the index (L). Now while the gas is in equilibrium with the liquid or solid, let the quantity which it contains of  $S_2$  receive the increment  $dm_2$ , its volume and the quantity which it contains of the other component, as well as the temperature, remaining constant. The potential for  $S_1$  in the gaseous mass will receive the increment

$$\left(\frac{d\mu_1}{dm_2}\right)_{t, v, m}^{(G)} dm_2$$

and the pressure will receive the increment

$$\left(\frac{dp}{dm_2}\right)_{t, v, m}^{(G)} dm_2.$$

Now the liquid or solid remaining in equilibrium with the gas must experience the same variations in the values of  $\mu_1$  and  $p$ . But by (272)

$$\left(\frac{d\mu_1}{dp}\right)_{t, m}^{(L)} = \left(\frac{dv}{dm_1}\right)_{t, p, m}^{(L)}.$$

Therefore,

$$\left(\frac{dv}{dm_1}\right)_{t, p, m}^{(L)} = \frac{\left(\frac{d\mu_1}{dm_2}\right)_{t, v, m}^{(G)}}{\left(\frac{dp}{dm_2}\right)_{t, v, m}^{(G)}}.$$

It will be observed that the first member of this equation relates solely to the liquid or solid, and the second member solely to the gas. Now we may suppose the same gaseous mass to be capable of equilibrium with several different liquids or solids, and the first member of this equation must therefore have the same value for all such liquids or solids; which is quite inadmissible. In the simplest case, in which the liquid or solid is identical in substance with the vapor which it yields, it is evident that the expression in question denotes the reciprocal of the density of the solid or liquid. Hence, when a gas is in equilibrium with one of its components both in the solid and liquid states (as when a moist gas is in equilibrium with ice and water), it would be necessary that the solid and liquid should have the same density.

The foregoing considerations appear sufficient to justify the definition of an ideal gas-mixture which we have chosen. It is of course immaterial whether we regard the definition as expressed by equation (273), or by (279), or by any other fundamental equation which can be derived from these.

The fundamental equations for an ideal gas-mixture corresponding to (255), (265), and (271) may easily be derived from these equations by using inversely the substitutions given on page 217. They are

$$\Sigma_1(c_1 m_1) \log \frac{\varepsilon - \Sigma_1(E_1 m_1)}{\Sigma_1(c_1 m_1)} = \eta + \Sigma_1(a_1 m_1 \log \frac{m_1}{v} - H_1 m_1), \quad (291)$$

$$\begin{aligned} \Sigma_1(c_1 m_1 + a_1 m_1) \log \frac{\chi - \Sigma_1(E_1 m_1)}{\Sigma_1(c_1 m_1 + a_1 m_1)} \\ = \eta + \Sigma_1 \left( a_1 m_1 \log \frac{p m_1}{\Sigma_1(a_1 m_1)} - H_1 m_1 \right), \end{aligned} \quad (292)$$

$$\begin{aligned} \xi = \Sigma_1 \left( E_1 m_1 + m_1 t (c_1 + a_1 - H_1) \right) \\ - \Sigma_1(c_1 m_1 + a_1 m_1) t \log t + \Sigma_1 \left( a_1 m_1 t \log \frac{p m_1}{\Sigma_1(a_1 m_1)} \right). \end{aligned} \quad (293)$$

The components to which the fundamental equations (273), (279), (291), (292), 293 refer, may themselves be gas-mixtures. We may for example apply the fundamental equations of a binary gas-mixture

to a mixture of hydrogen and air, or to any ternary gas-mixture in which the proportion of two of the components is fixed. In fact, the form of equation (279) which applies to a gas-mixture of any particular number of components may easily be reduced, when the proportions of some of these components are fixed, to the form which applies to a gas-mixture of a smaller number of components. The necessary substitutions will be analogous to those given on page 217. But the components must be entirely different from one another with respect to the gases of which they are formed by mixture. We cannot, for example, apply equation (279) to a gas-mixture in which the components are oxygen and air. It would indeed be easy to form a fundamental equation for such a gas-mixture with reference to the designated gases as components. Such an equation might be derived from (279) by the proper substitutions. But the result would be an equation of more complexity than (279). A *chemical* compound, however, with respect to Dalton's law, and with respect to all the equations which have been given, is to be regarded as entirely different from its components. Thus, a mixture of hydrogen, oxygen, and vapor of water is to be regarded as a ternary gas-mixture, having the three components mentioned. This is certainly true when the quantities of the compound gas and of its components are all independently variable in the gas-mixture, without change of temperature or pressure. Cases in which these quantities are not thus independently variable will be considered hereafter.

#### *Inferences in regard to Potentials in Liquids and Solids.*

Such equations as (264), (268), (276), by which the values of potentials in pure or mixed gases may be derived from quantities capable of direct measurement, have an interest which is not confined to the theory of gases. For as the potentials of the independently variable components which are common to coexistent liquid and gaseous masses have the same values in each, these expressions will generally afford the means of determining for liquids, at least approximately, the potential for any independently variable component which is capable of existing in the gaseous state. For although every state of a liquid is not such as can exist in contact with a gaseous mass, it will always be possible, when any of the components of the liquid are volatile, to bring it by a change of pressure alone, its temperature and composition remaining unchanged, to a state for which there is a coexistent phase of vapor, in which

the values of the potentials of the volatile components of the liquid may be estimated from the density of these substances in the vapor. The variations of the potentials in the liquid due to the change of pressure will in general be quite trifling as compared with the variations which are connected with changes of temperature or of composition, and may moreover be readily estimated by means of equation (272). The same considerations will apply to volatile solids with respect to the determination of the potential for the substance of the solid.

As an application of this method of determining the potentials in liquids, let us make use of the law of Henry in regard to the absorption of gases by liquids to determine the relation between the quantity of the gas contained in any liquid mass and its potential. Let us consider the liquid as in equilibrium with the gas, and let  $m_1^{(G)}$  denote the quantity of the gas existing as such,  $m_1^{(L)}$  the quantity of the same substance contained in the liquid mass,  $\mu_1$  the potential for this substance common to the gas and liquid,  $v^{(G)}$  and  $v^{(L)}$  the volumes of the gas and liquid. When the absorbed gas forms but a very small part of the liquid mass, we have by Henry's law

$$\frac{m_1^{(L)}}{v^{(L)}} = A \frac{m_1^{(G)}}{v^{(G)}}, \quad (294)$$

where  $A$  is a function of the temperature; and by (276)

$$\mu_1 = B + C \log \frac{m_1^{(G)}}{v^{(G)}}, \quad (295)$$

$B$  and  $C$  also denoting functions of the temperature. Therefore

$$\mu_1 = B + C \log \frac{m_1^{(L)}}{A v^{(L)}}. \quad (296)$$

It will be seen (if we disregard the difference of notation) that this equation is equivalent in form to (216), which was deduced from *a priori* considerations as a probable relation between the quantity and the potential of a small component. When a liquid absorbs several gases at once, there will be several equations of the form of (296), which will hold true simultaneously, and which we may regard as equivalent to equations (217), (218). The quantities  $A$  and  $C$  in (216), with the corresponding quantities in (217), (218), were regarded as functions of the temperature and pressure, but since the potentials in liquids are but little affected by the pressure, we might anticipate that these quantities in the case of liquids might be regarded as functions of the temperature alone.

In regard to equations (216), (217), (218), we may now observe that by (264) and (276) they are shown to hold true in ideal gases or gas-mixtures, not only for components which form only a small part of the whole gas-mixture, but without any such limitation, and not only approximately but absolutely. It is noticeable that in this case quantities  $A$  and  $C$  are functions of the temperature alone, and do not even depend upon the nature of the gaseous mass, except upon the particular component to which they relate. As all gaseous bodies are generally supposed to approximate to the laws of ideal gases when sufficiently rarefied, we may regard these equations as approximately valid for gaseous bodies in general when the density is sufficiently small. When the density of the gaseous mass is very great, but the separate density of the component in question is small, the equations will probably hold true, but the values of  $A$  and  $C$  may not be entirely independent of the pressure, or of the composition of the mass in respect to its principal components. These equations will also apply, as we have just seen, to the potentials in liquid bodies for components of which the density in the liquid is very small, whenever these components exist also in the gaseous state, and conform to the law of Henry. This seems to indicate that the law expressed by these equations has a very general application.

*Considerations relating to the Increase of Entropy due to the Mixture of Gases by Diffusion.*

From equation (278) we may easily calculate the increase of entropy which takes place when two different gases are mixed by diffusion, at a constant temperature and pressure. Let us suppose that the quantities of the gases are such that each occupies initially one half of the total volume. If we denote this volume by  $V$ , the increase of entropy will be

$$m_1 a_1 \log V + m_2 a_2 \log V - m_1 a_1 \log \frac{V}{2} - m_2 a_2 \log \frac{V}{2},$$

or

$$(m_1 a_1 + m_2 a_2) \log 2.$$

Now  $m_1 a_1 = \frac{pV}{2t}$ , and  $m_2 a_2 = \frac{pV}{2t}$ .

Therefore the increase of entropy may be represented by the expression

$$\frac{pV}{t} \log 2. \quad (297)$$

It is noticeable that the value of this expression does not depend upon the kinds of gas which are concerned, if the quantities are such as has been supposed, except that the gases which are mixed must be

of different kinds. If we should bring into contact two masses of the same kind of gas, they would also mix, but there would be no increase of entropy. But in regard to the relation which this case bears to the preceding, we must bear in mind the following considerations. When we say that when two different gases mix by diffusion, as we have supposed, the energy of the whole remains constant, and the entropy receives a certain increase, we mean that the gases could be separated and brought to the same volume and temperature which they had at first by means of certain changes in external bodies, for example, by the passage of a certain amount of heat from a warmer to a colder body. But when we say that when two gas-masses of the same kind are mixed under similar circumstances there is no change of energy or entropy, we do not mean that the gases which have been mixed can be separated without change to external bodies. On the contrary, the separation of the gases is entirely impossible. We call the energy and entropy of the gas-masses when mixed the same as when they were unmixed, because we do not recognize any difference in the substance of the two masses. So when gases of different kinds are mixed, if we ask what changes in external bodies are necessary to bring the system to its original state, we do not mean a state in which each particle shall occupy more or less exactly the same position as at some previous epoch, but only a state which shall be undistinguishable from the previous one in its sensible properties. It is to states of systems thus incompletely defined that the problems of thermodynamics relate.

But if such considerations explain why the mixture of gas-masses of the same kind stands on a different footing from the mixture of gas-masses of different kinds, the fact is not less significant that the increase of entropy due to the mixture of gases of different kinds, in such a case as we have supposed, is independent of the nature of the gases.

Now we may without violence to the general laws of gases which are embodied in our equations suppose other gases to exist than such as actually do exist, and there does not appear to be any limit to the resemblance which there might be between two such kinds of gas. But the increase of entropy due to the mixing of given volumes of the gases at a given temperature and pressure would be independent of the degree of similarity or dissimilarity between them. We might also imagine the case of two gases which should be absolutely identical in all the properties (sensible and molecular) which come into play while they exist as gases either pure or mixed with each other,

but which should differ in respect to the attractions between their atoms and the atoms of some other substances, and therefore in their tendency to combine with such substances. In the mixture of such gases by diffusion an increase of entropy would take place, although the process of mixture, dynamically considered, might be absolutely identical in its minutest details (even with respect to the precise path of each atom) with processes which might take place without any increase of entropy. In such respects, entropy stands strongly contrasted with energy. Again, when such gases have been mixed, there is no more impossibility of the separation of the two kinds of molecules in virtue of their ordinary motions in the gaseous mass without any especial external influence, than there is of the separation of a homogeneous gas into the same two parts into which it has once been divided, after these have once been mixed. In other words, the impossibility of an uncompensated decrease of entropy seems to be reduced to improbability.

There is perhaps no fact in the molecular theory of gases so well established as that the number of molecules in a given volume at a given temperature and pressure is the same for every kind of gas when in a state to which the laws of ideal gases apply. Hence the quantity  $\frac{pV}{t}$  in (297) must be entirely determined by the number of molecules which are mixed. And the increase of entropy is therefore determined by the number of these molecules and is independent of their dynamical condition and of the degree of difference between them.

The result is of the same nature when the volumes of the gases which are mixed are not equal, and when more than two kinds of gas are mixed. If we denote by  $v_1, v_2$ , etc., the initial volumes of the different kinds of gas, and by  $V$  as before the total volume, the increase of entropy may be written in the form

$$\Sigma_1 (m_1 a_1) \log V - \Sigma_1 (m_1 a_1 \log v_1).$$

And if we denote by  $r_1, r_2$ , etc., the numbers of the molecules of the several different kinds of gas, we shall have

$$r_1 = C m_1 a_1, \quad r_2 = C m_2 a_2, \quad \text{etc.,}$$

where  $C$  denotes a constant. Hence

$$r_1 : V :: m_1 a_1 : \Sigma_1 (m_1 a_1) :: r_1 : \Sigma_1 r_1;$$

and the increase of entropy may be written

$$\frac{\Sigma_1 r_1 \log \Sigma_1 r_1 - \Sigma_1 (r_1 \log r_1)}{C}. \quad (298)$$

*The Phases of Dissipated Energy of an Ideal Gas-mixture with Components which are Chemically Related.*

We will now pass to the consideration of the phases of dissipated energy (see page 200) of an ideal gas-mixture, in which the number of the proximate components exceeds that of the ultimate.

Let us first suppose that an ideal gas-mixture has for proximate components the gases  $G_1$ ,  $G_2$ , and  $G_3$ , the units of which are denoted by  $\mathfrak{G}_1$ ,  $\mathfrak{G}_2$ ,  $\mathfrak{G}_3$ , and that in ultimate analysis

$$\mathfrak{G}_3 = \lambda_1 \mathfrak{G}_1 + \lambda_2 \mathfrak{G}_2, \quad (299)$$

$\lambda_1$  and  $\lambda_2$  denoting positive constants, such that  $\lambda_1 + \lambda_2 = 1$ . The phases which we are to consider are those for which the energy of the gas-mixture is a minimum for constant entropy and volume and constant quantities of  $G_1$  and  $G_2$ , as determined in ultimate analysis. For such phases, by (86),

$$\mu_1 \delta m_1 + \mu_2 \delta m_2 + \mu_3 \delta m_3 \geq 0. \quad (300)$$

for such values of the variations as do not affect the quantities of  $G_1$  and  $G_2$  as determined in ultimate analysis. Values of  $\delta m_1$ ,  $\delta m_2$ ,  $\delta m_3$  proportional to  $\lambda_1$ ,  $\lambda_2$ ,  $-1$ , and only such, are evidently consistent with this restriction: therefore

$$\lambda_1 \mu_1 + \lambda_2 \mu_2 = \mu_3. \quad (301)$$

If we substitute in this equation values of  $\mu_1$ ,  $\mu_2$ ,  $\mu_3$  taken from (276), we obtain, after arranging the terms and dividing by  $t$ ,

$$\lambda_1 a_1 \log \frac{m_1}{v} + \lambda_2 a_2 \log \frac{m_2}{v} - a_3 \log \frac{m_3}{v} = A + B \log t - \frac{C}{t}, \quad (302)$$

where

$$A = \lambda_1 H_1 + \lambda_2 H_2 - H_3 - \lambda_1 c_1 - \lambda_2 c_2 + c_3 - \lambda_1 a_1 - \lambda_2 a_2 + a_3, \quad (303)$$

$$B = \lambda_1 c_1 + \lambda_2 c_2 - c_3, \quad (304)$$

$$C = \lambda_1 E_1 + \lambda_2 E_2 - E_3. \quad (305)$$

If we denote by  $\beta_1$  and  $\beta_2$  the volumes (determined under standard conditions of temperature and pressure) of the quantities of the gases  $G_1$  and  $G_2$  which are contained in a unit of volume of the gas  $G_3$ , we shall have

$$\beta_1 = \frac{\lambda_1 a_1}{a_3}, \text{ and } \beta_2 = \frac{\lambda_2 a_2}{a_3}, \quad (306)$$

and (302) will reduce to the form

$$\log \frac{\frac{\beta_1}{m_1} - \frac{\beta_2}{m_2}}{m_3 v^{\frac{\beta_1 + \beta_2}{\beta_1 + \beta_2 - 1}}} = \frac{A}{a_3} + \frac{B}{a_3} \log t - \frac{C}{a_3 t}. \quad (307)$$

Moreover, as by (277)

$$p v = (a_1 m_1 + a_2 m_2 + a_3 m_3) t, \quad (308)$$

we have on eliminating  $v$

$$\log \frac{m_1^{\beta_1} m_2^{\beta_2} p^{\beta_1 + \beta_2 - 1}}{m_3 (a_1 m_1 + a_2 m_2 + a_3 m_3)^{\beta_1 + \beta_2 - 1}} = \frac{A}{a_3} + \frac{B'}{a_3} \log t - \frac{C}{a_3 t}, \quad (309)$$

where

$$B' = \lambda_1 c_1 + \lambda_2 c_2 - c_3 + \lambda_1 a_1 + \lambda_2 a_2 - a_3. \quad (310)$$

It will be observed that the quantities  $\beta_1$ ,  $\beta_2$  will always be positive and have a simple relation to unity, and that the value of  $\beta_1 + \beta_2 - 1$  will be positive or zero, according as gas  $G_3$  is formed of  $G_1$  and  $G_2$  with or without condensation. If we should assume, according to the rule often given for the specific heat of compound gases, that the thermal capacity at constant volume of any quantity of the gas  $G_3$  is equal to the sum of the thermal capacities of the quantities which it contains of the gases  $G_1$  and  $G_2$ , the value of  $B$  would be zero. The heat evolved in the formation of a unit of the gas  $G_3$  out of the gases  $G_1$  and  $G_2$ , without mechanical action, is by (283) and (257)

$$\lambda_1 (c_1 t + E_1) + \lambda_2 (c_2 t + E_2) - (c_3 t + E_3),$$

or

$$Bt + C,$$

which will reduce to  $C$  when the above relation in regard to the specific heats is satisfied. In any case the quantity of heat thus evolved divided by  $a_3 t^2$  will be equal to the differential coefficient of the second member of equation (307) with respect to  $t$ . Moreover, the heat evolved in the formation of a unit of the gas  $G_3$  out of the gases  $G_1$  and  $G_2$  under constant pressure is

$$Bt + C + \lambda_1 a_1 t + \lambda_2 a_2 t - a_3 t = B't + C,$$

which is equal to the differential coefficient of the second member of (309) with respect to  $t$ , multiplied by  $a_3 t^2$ .

It appears by (307) that, except in the case when  $\beta_1 + \beta_2 = 1$ , for any given finite values of  $m_1$ ,  $m_2$ ,  $m_3$ , and  $t$  (infinitesimal values being excluded as well as infinite), it will always be possible to assign such a finite value to  $v$  that the mixture shall be in a state of dissipated energy. Thus, if we regard a mixture of hydrogen, oxygen, and vapor of water as an ideal gas-mixture, for a mixture containing any given quantities of these three gases at any given tem-

perature there will be a certain volume at which the mixture will be in a state of dissipated energy. In such a state no such phenomenon as explosion will be possible, and no formation of water by the action of platinum. (If the mass should be expanded beyond this volume, the only possible action of a catalytic agent would be to resolve the water into its components.) It may indeed be true that at ordinary temperatures, except when the quantity either of hydrogen or of oxygen is very small compared with the quantity of water, the state of dissipated energy is one of such extreme rarefaction as to lie entirely beyond our power of experimental verification. It is also to be noticed that a state of great rarefaction is so unfavorable to any condensation of the gases, that it is quite probable that the catalytic action of platinum may cease entirely at a degree of rarefaction far short of what is necessary for a state of dissipated energy. But with respect to the theoretical demonstration, such states of great rarefaction are precisely those to which we should suppose that the laws of ideal gas-mixtures would apply most perfectly.

But when the compound gas  $G_3$  is formed of  $G_1$  and  $G_2$  without condensation, (i. e., when  $\beta_1 + \beta_2 = 1$ ,) it appears from equation (307) that the relation between  $m_1$ ,  $m_2$ , and  $m_3$  which is necessary for a phase of dissipated energy is determined by the temperature alone.

In any case, if we regard the total quantities of the gases  $G_1$  and  $G_2$  (as determined by the ultimate analysis of the gas-mixture), and also the volume, as constant, the quantities of these gases which appear uncombined in a phase of dissipated energy will increase with the temperature, if the formation of the compound  $G_3$  *without change of volume* is attended with evolution of heat. Also, if we regard the total quantities of the gases  $G_1$  and  $G_2$ , and also the pressure, as constant, the quantities of these gases which appear uncombined in a phase of dissipated energy, will increase with the temperature, if the formation of the compound  $G_3$  *under constant pressure* is attended with evolution of heat. If  $B = 0$ , (a case, as has been seen, of especial importance), the heat obtained by the formation of a unit of  $G_3$  out of  $G_1$  and  $G_2$  without change of volume or of temperature will be equal to  $C$ . If this quantity is positive, and the total quantities of the gases  $G_1$  and  $G_2$  and also the volume have given finite values, for an infinitesimal value of  $t$  we shall have (for a phase of dissipated energy) an infinitesimal value either of  $m_1$  or of  $m_2$ , and for an infinite value of  $t$  we shall have finite (neither infinitesimal nor infinite) values of  $m_1$ ,  $m_2$ , and  $m_3$ . But if we suppose the pressure instead of the volume to have a given finite value (with

suppositions otherwise the same), we shall have for infinitesimal values of  $t$  an infinitesimal value either of  $m_1$  or  $m_2$ , and for infinite values of  $t$  finite or infinitesimal values of  $m_3$  according as  $\beta_1 + \beta_2$  is equal to or greater than unity.

The case which we have considered is that of a ternary gas-mixture, but our results may easily be generalized in this respect. In fact, whatever the number of component gases in a gas-mixture, if there are relations of equivalence in ultimate analysis between these components, such relations may be expressed by one or more equations of the form

$$\lambda_1 \mathfrak{G}_1 + \lambda_2 \mathfrak{G}_2 + \lambda_3 \mathfrak{G}_3 + \text{etc.} = 0, \quad (311)$$

where  $\mathfrak{G}_1, \mathfrak{G}_2$ , etc. denote the units of the various component gases, and  $\lambda_1, \lambda_2$ , etc. denote positive or negative constants such that  $\Sigma_1 \lambda_i = 0$ . From (311) with (86) we may derive for phases of dissipated energy,

$$\lambda_1 \mu_1 + \lambda_2 \mu_2 + \lambda_3 \mu_3 + \text{etc.} = 0,$$

$$\text{or} \quad \Sigma_1 (\lambda_1 \mu_1) = 0. \quad (312)$$

Hence, by (276),

$$\Sigma_1 \left( \lambda_1 a_1 \log \frac{m_1}{v} \right) = A + B \log t - \frac{C}{t}, \quad (313)$$

where  $A, B$  and  $C$  are constants determined by the equations

$$A = \Sigma_1 (\lambda_1 H_1 - \lambda_1 c_1 - \lambda_1 a_1), \quad (314)$$

$$B = \Sigma_1 (\lambda_1 c_1), \quad (315)$$

$$C = \Sigma_1 (\lambda_1 E_1). \quad (316)$$

Also, since

$$p v = \Sigma_1 (a_1 m_1) t,$$

$$\begin{aligned} \Sigma_1 (\lambda_1 a_1 \log m_1) - \Sigma (\lambda_1 a_1) \log \Sigma_1 (a_1 m_1) \\ + \Sigma (\lambda_1 a_1) \log p = A + B' \log t - \frac{C}{t}, \end{aligned} \quad (317)$$

where

$$B' = \Sigma_1 (\lambda_1 c_1 + \lambda_1 a_1). \quad (318)$$

If there is more than one equation of the form (311), we shall have more than one of each of the forms (313) and (317), which will hold true simultaneously for phases of dissipated energy.

It will be observed that the relations necessary for a phase of dissipated energy between the volume and temperature of an ideal gas-mixture, and the quantities of the components which take part in the chemical processes, and the pressure due to these components, are not affected by the presence of neutral gases in the gas-mixture.

From equations (312) and (324) it follows that if there is a phase of dissipated energy at any point in an ideal gas-mixture in equilibrium under the influence of gravity, the whole gas-mixture must consist of such phases.

The equations of the phases of dissipated energy of a binary gas-mixture, the components of which are identical in substance, are comparatively simple in form. In this case the two components have the same potential, and if we write  $\beta$  for  $\frac{a_1}{a_2}$  (the ratio of the volumes of equal quantities of the two components under the same conditions of temperature and pressure), we shall have

$$\log \frac{m_1^\beta}{m_2 v^{\beta-1}} = \frac{A}{a_2} + \frac{B}{a_2} \log t - \frac{C}{a_2 t}, \quad (319)$$

$$\log \frac{m_1^\beta p^{\beta-1}}{m_2 (a_1 m_1 + a_2 m_2)^{\beta-1}} = \frac{A}{a_2} + \frac{B'}{a_2} \log t - \frac{C}{a_2 t}; \quad (320)$$

where

$$A = H_1 - H_2 - c_1 + c_2 - a_1 + a_2, \quad (321)$$

$$B = c_1 - c_2, \quad B' = c_1 - c_2 + a_1 - a_2, \quad (322)$$

$$C = E_1 - E_2. \quad (323)$$

#### *Gas-mixtures with Convertible Components.*

The equations of the phases of dissipated energy of ideal gas-mixtures which have components of which some are identical in ultimate analysis to others have an especial interest in relation to the theory of gas-mixtures in which the components are not only thus equivalent, but are actually transformed into each other within the gas-mixture on variations of temperature and pressure, so that quantities of these (proximate) components are entirely determined, at least in any permanent phase of the gas-mixture, by the quantities of a smaller number of ultimate components, with the temperature and pressure. Such gas-mixtures may be distinguished as having *convertible components*. The very general considerations adduced on pages 197–203, which are not limited in their application to gaseous bodies, suggest the hypothesis that the equations of the phases of dissipated energy of ideal gas-mixtures may apply to such gas-mixtures as have been described. It will, however, be desirable to consider the matter more in detail.

In the first place, if we consider the case of a gas-mixture which only differs from an ordinary ideal gas-mixture for which some of the components are equivalent in that there is perfect freedom in regard to the transformation of these components, it follows at once from the general formulae of equilibrium (1) or (2) that equilibrium is only possible for such phases as we have called phases of dissipated energy, for which some of the characteristic equations have been deduced in the preceding pages.

If it should be urged, that regarding a gas-mixture which has convertible components as an ideal gas-mixture of which, for some reason, only a part of the phases are actually capable of existing, we might still suppose the particular phases which alone can exist to be determined by some other principle than that of the free convertibility of the components (as if, perhaps, the case were analogous to one of *constraint* in mechanics), it may easily be shown that such a hypothesis is entirely untenable, when the quantities of the proximate components may be varied independently by suitable variations of the temperature and pressure, and of the quantities of the ultimate components, and it is admitted that the relations between the energy, entropy, volume, temperature, pressure, and the quantities of the several proximate components in the gas-mixture are the same as for an ordinary ideal gas-mixture, in which the components are not convertible. Let us denote the quantities of the  $n'$  proximate components of a gas-mixture  $A$  by  $m_1, m_2$ , etc., and the quantities of its  $n$  ultimate components by  $\mathbf{m}_1, \mathbf{m}_2$ , etc. ( $n$  denoting a number less than  $n'$ ), and let us suppose that for this gas-mixture the quantities  $\varepsilon, \eta, v, t, p, m_1, m_2$ , etc. satisfy the relations characteristic of an ideal gas-mixture, while the phase of the gas-mixture is entirely determined by the values of  $\mathbf{m}_1, \mathbf{m}_2$ , etc., with two of the quantities  $\varepsilon, \eta, v, t, p$ . We may evidently imagine such an ideal gas-mixture  $B$  having  $n'$  components (not convertible), that every phase of  $A$  shall correspond with one of  $B$  in the values of  $\varepsilon, \eta, v, t, p, m_1, m_2$ , etc. Now let us give to the quantities  $\mathbf{m}_1, \mathbf{m}_2$ , etc. in the gas-mixture  $A$  any fixed values, and for the body thus defined let us imagine the  $v\text{-}\eta\text{-}\varepsilon$  surface (see page 174) constructed; likewise for the ideal gas-mixture  $B$  let us imagine the  $v\text{-}\eta\text{-}\varepsilon$  surface constructed for every set of values of  $m_1, m_2$ , etc. which is consistent with the given values of  $\mathbf{m}_1, \mathbf{m}_2$ , etc., i. e., for every body of which the *ultimate* composition would be expressed by the given values of  $\mathbf{m}_1, \mathbf{m}_2$ , etc. It follows immediately from our supposition, that every point in the  $v\text{-}\eta\text{-}\varepsilon$  surface relating to  $A$  must coincide with some point of one of the  $v\text{-}\eta\text{-}\varepsilon$  surfaces relating

to  $B$  not only in respect to position but also in respect to its tangent plane (which represents temperature and pressure); therefore the  $v\text{-}\eta\text{-}\varepsilon$  surface relating to  $A$  must be tangent to the various  $v\text{-}\eta\text{-}\varepsilon$  surfaces relating to  $B$ , and therefore must be an envelop of these surfaces. From this it follows that the points which represent phases common to both gas-mixtures must represent the phases of dissipated energy of the gas-mixture  $B$ .

The properties of an ideal gas-mixture which are assumed in regard to the gas-mixture of convertible components in the above demonstration are expressed by equations (277) and (278) with the equation

$$\varepsilon = \sum_1 (c_1 m_1 t + m_1 E_1). \quad (324)$$

It is usual to assume in regard to gas-mixtures having convertible components that the convertibility of the components does not affect the relations (277) and (324). The same cannot be said of the equation (278). But in a very important class of cases it will be sufficient if the applicability of (277) and (324) is admitted. The cases referred to are those in which in certain phases of a gas-mixture the components are convertible, and in other phases of the same proximate composition the components are not convertible, and the equations of an ideal gas-mixture hold true.

If there is only a single degree of convertibility between the components, (i. e., if only a single kind of conversion, with its reverse, can take place among the components,) it will be sufficient to assume, in regard to the phases in which conversion takes place, the validity of equation (277) and of the following, which can be derived from (324) by differentiation, and comparison with equation (11), which expresses a necessary relation,

$$[t d\eta - p dv - \sum_1 (c_1 m_1) dt]_m = 0.* \quad (325)$$

We shall confine our demonstration to this case. It will be observed that the physical signification of (325) is that if the gas-mixture is subjected to such changes of volume and temperature as do not alter its proximate composition, the heat absorbed or yielded may be calculated by the same formula as if the components were not convertible.

Let us suppose the thermodynamic state of a gaseous mass  $M$ , of such a kind as has just been described, to be varied while within the limits within which the components are not convertible. (The quantities of the proximate components, therefore, as well as of the ultimate

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\* This notation is intended to indicate that  $m_1$ ,  $m_2$ , etc. are regarded as constant

mate, are supposed constant). If we use the same method of geometrical representation as before, the point representing the volume, entropy, and energy of the mass will describe a line in the  $v\text{-}\eta\text{-}\epsilon$  surface of an ideal gas-mixture of inconvertible components, the form and position of this surface being determined by the proximate composition of  $M$ . Let us now suppose the same mass to be carried beyond the limit of inconvertibility, the variations of state after passing the limit being such as not to alter its proximate composition. It is evident that this will in general be possible. Exceptions can only occur when the limit is formed by phases in which the proximate composition is uniform. The line traced in the region of convertibility must belong to the same  $v\text{-}\eta\text{-}\epsilon$  surface of an ideal gas-mixture of inconvertible components as before, continued beyond the limit of inconvertibility for the components of  $M$ , since the variations of volume, entropy and energy are the same as would be possible if the components were not convertible. But it must also belong to the  $v\text{-}\eta\text{-}\epsilon$  surface of the body  $M$ , which is here a gas-mixture of convertible components. Moreover, as the inclination of each of these surfaces must indicate the temperature and pressure of the phases through which the body passes, these two surfaces must be tangent to each other along the line which has been traced. As the  $v\text{-}\eta\text{-}\epsilon$  surface of the body  $M$  in the region of convertibility must thus be tangent to all the surfaces representing ideal gas-mixtures of every possible proximate composition consistent with the ultimate composition of  $M$ , continued beyond the region of inconvertibility, in which alone their form and position may be capable of experimental demonstration, the former surface must be an envelop of the latter surfaces, and therefore a continuation of the surface of the phases of dissipated energy in the region of inconvertibility.

The foregoing considerations may give a measure of *a priori* probability to the results which are obtained by applying the ordinary laws of ideal gas-mixtures to cases in which the components are convertible. It is only by experiments upon gases in phases in which their components are convertible that the validity of any of these results can be established.

The very accurate determinations of density which have been made for the peroxide of nitrogen enable us to subject some of our equations to a very critical test. That this substance in the gaseous state is properly regarded as a mixture of different gases can hardly be doubted, as the proportion of the components derived from its density on the supposition that one component has the molecular formula

$\text{NO}_2$  and the other the formula  $\text{N}_2\text{O}_4$  is the same as that derived from the depth of the color on the supposition that the absorption of light is due to one of the components alone, and is proportioned to the separate density of that component.\*

MM. Sainte-Claire Deville and Troost† have given a series of determinations of what we shall call the *relative densities* of peroxide of nitrogen at various temperatures under atmospheric pressure. We use the term *relative density* to denote what it is usual in treatises on chemistry to denote by the term *density*, viz., the actual density of a gas divided by the density of a standard perfect gas at the same pressure and temperature, the standard gas being air, or more strictly, an ideal gas which has the same density as air at the zero of the centigrade scale and the pressure of one atmosphere. In order to test our equations by these determinations, it will be convenient to transform equation (320), so as to give directly the relation between the relative density, the pressure, and the temperature.

As the density of the standard gas at any given temperature and pressure may by (263) be expressed by the formula  $\frac{P}{a_s t}$ , the relative density of a binary gas-mixture may be expressed by

$$D = (m_1 + m_2) \frac{a_s t}{P v}. \quad (326)$$

Now by (263)

$$a_1 m_1 + a_2 m_2 = \frac{P v}{t}. \quad (327)$$

By giving to  $m_2$  and  $m_1$  successively the value zero in these equations, we obtain

$$D_1 = \frac{a_s}{a_1}, \quad D_2 = \frac{a_s}{a_2}, \quad (328)$$

where  $D_1$  and  $D_2$  denote the values of  $D$  when the gas consists wholly of one or of the other component. If we assume that

$$D_2 = 2 D_1, \quad (329)$$

we shall have

$$a_1 = 2 a_2. \quad (330)$$

From (326) we have

$$m_1 + m_2 = D \frac{P v}{a_s t},$$

\* Salet, "Sur la coloration du peroxyde d'azote," Comptes Rendus, vol. lxvii, p. 488.

† Comptes Rendus, vol. lxiv, p. 237.

and from (327), by (328) and (330),

$$2m_1 + m_2 = D_2 \frac{p v}{a_s t} = 2D_1 \frac{p v}{a_s t},$$

whence

$$m_1 = (D_2 - D) \frac{p v}{a_s t} \quad (331)$$

$$m_2 = 2(D - D_1) \frac{p v}{a_s t} \quad (332)$$

By (327), (331), and (332) we obtain from (320)

$$\log \frac{(D_2 - D)^2 p}{2(D - D_1) a_s} = \frac{A}{a_2} + \frac{B'}{a_2} \log t - \frac{C}{t}. \quad (333)$$

This formula will be more convenient for purposes of calculation if we introduce common logarithms (denoted by  $\log_{10}$ ) instead of hyperbolic, the temperature of the ordinary contigraide scale  $t_c$  instead of the absolute temperature  $t$ , and the pressure in atmospheres  $p_{at}$  instead of  $p$  the pressure in a rational system of units. If we also add the logarithm of  $a_s$  to both sides of the equation, we obtain

$$\log_{10} \frac{(D_2 - D)^2 p_{at}}{2(D - D_1)} = \mathbf{A} + \frac{B'}{a_2} \log_{10} (t_c + 273) - \frac{\mathbf{C}}{t_c + 273}, \quad (334)$$

where  $\mathbf{A}$  and  $\mathbf{C}$  denote constants, the values of which are closely connected with those of  $A$  and  $C$ .

From the molecular formulae of peroxide of nitrogen  $\text{NO}_2$  and  $\text{N}_2\text{O}_4$ , we may calculate the relative densities

$$D_1 = \frac{14 + 32}{2} \cdot 0.691 = 1.589, \text{ and } D_2 = \frac{28 + 64}{2} \cdot 0.691 = 3.178. \quad (335)$$

The determinations of MM. Deville and Troost are satisfactorily represented by the equation

$$\log_{10} \frac{(3.178 - D)^2 p_{at}}{2(D - 1.589)} = 9.47056 - \frac{3118.6}{t_c + 273}, \quad (336)$$

which gives

$$D = 3.178 + \Theta - \sqrt{\Theta(3.178 + \Theta)}$$

$$\text{where } \log_{10} \Theta = 9.47056 - \frac{3118.6}{t_c + 273} - \log_{10} p_{at}.$$

In the first part of the following table are given in successive columns the temperature and pressure of the gas in the several experiments of MM. Deville and Troost, the relative densities calculated from these numbers by equation (336), the relative densities as observed, and the difference of the observed and calculated relative

densities. It will be observed that these differences are quite small, in no case reaching .03, and on the average scarcely exceeding .01. The significance of such correspondence in favor of the hypothesis by means of which equation (336) has been established is of course diminished by the fact that two constants in the equation have been determined from these experiments. If the same equation can be shown to give correctly the relative densities at other pressures than that for which the constants have been determined, such correspondence will be much more decisive.

$t_c$	$p_{at}$	$D$ calculated by eq. (336).	$D$ observed.	diff.	Observers.
26.7	1	2.676	2.65	-.026	D. & T.
35.4	1	2.524	2.53	+.006	D. & T.
39.8	1	2.443	2.46	+.017	D. & T.
49.6	1	2.256	2.27	+.014	D. & T.
60.2	1	2.067	2.08	+.013	D. & T.
70.0	1	1.920	1.92	.000	D. & T.
80.6	1	1.801	1.80	-.001	D. & T.
90.0	1	1.728	1.72	-.008	D. & T.
100.1	1	1.676	1.68	+.004	D. & T.
111.3	1	1.641	1.65	+.009	D. & T.
121.5	1	1.622	1.62	-.002	D. & T.
135.0	1	1.607	1.60	-.007	D. & T.
154.0	1	1.597	1.58	-.017	D. & T.
183.2	1	1.592	1.57	-.022	D. & T.
<hr/>					
97.5	1	1.687			
97.5	$\frac{1.0450}{2.6397}$	1.631	1.783	+.152	P. & W.
24.5	1	2.711			
24.5	$\frac{1.8090}{4.2529}$	2.524	2.52	-.004	P. & W.
11.3	1	2.891			
11.3	$\frac{9.2655}{4.4205}$	2.620	2.645	+.025	P. & W.
4.2	1	2.964			
4.2	$\frac{6.023}{3.5438}$	2.708	2.588	-.120	P. & W.

Messrs. Playfair and Wanklyn have published\* four determinations of the relative density of peroxide of nitrogen at various temperatures when diluted with nitrogen. Since the relations expressed by equations (319) and (320) are not affected by the presence of a third gas which is different from the gases  $G_1$  and  $G_2$  (to which  $m_1$  and  $m_2$  relate) and neutral to them, (see the remark at the foot of page 233), —provided that we take  $p$  to denote the pressure which we attribute to the gases  $G_1$  and  $G_2$ , i. e., the total pressure diminished by the pressure which the third gas would exert if occupying alone the same space at the same temperature,—it follows that the relations expressed for

\* Transactions of the Royal Society of Edinburgh, vol. xxii, p. 441.

peroxide of nitrogen by (333), (334), and (336) will not be affected by the presence of free nitrogen, if the pressure expressed by  $p$  or  $p_{at}$  and contained implicitly in the symbol  $D$  (see equation (326) by which  $D$  is defined) is understood to denote the total pressure diminished by the pressure due to the free nitrogen. The determinations of Playfair and Wanklyn are given in the latter part of the above table. The pressures given are those obtained by subtracting the pressure due to the free nitrogen from the total pressure. We may suppose such reduced pressures to have been used in the reduction of the observations by which the numbers in the column of observed relative densities were obtained. Besides the relative densities calculated by equation (336) for the temperatures and (reduced) pressures of the observations, the table contains the relative densities calculated for the same temperatures and the pressure of one atmosphere.

The reader will observe that in the second and third experiments of Playfair and Wanklyn there is a very close accordance between the calculated and observed values of  $D$ , while in the second and fourth experiments there is a considerable difference. Now the weight to be attributed to the several determinations is very different. The quantities of peroxide of nitrogen which were used in the several experiments were respectively .2410, .5893, .3166, and .2016 grammes. For a rough approximation, we may assume that the probable errors of the relative densities are inversely proportional to these numbers. This would make the probable error of the first and fourth observations two or three times as great as that of the second and considerably greater than that of the third. We must also observe that in the first of these experiments, the observed relative density 1.783 is greater than 1.687, the relative density calculated by equation (336) for the temperature of the experiment and the pressure of one atmosphere. Now the number 1.687 we may regard as established directly by the experiments of Deville and Troost. For in seven successive experiments in this part of the series the calculated relative densities differ from the observed by less than .01. If then we accept the numbers given by experiment, the effect of diluting the gas with nitrogen is to increase its relative density. As this result is entirely at variance with the facts observed in the case of other gases, and in the case of this gas at lower temperatures, as appears from the three other determinations of Playfair and Wanklyn, it cannot possibly be admitted on the strength of a single

observation. The first experiment of this series cannot therefore properly be used as a test of our equations. Similar considerations apply with somewhat less force to the last experiment. By comparing the temperatures and pressures of the three last experiments with the observed relative densities, the reader may easily convince himself that if we admit the substantial accuracy of the determinations in the two first of these experiments (the second and third of the series, which have the greatest weight), the last determination of relative density 2.588 must be too small. In fact, it should evidently be greater than the number in the preceding experiment 2.645.

If we confine our attention to the second and third experiments of the series, the agreement is as good as could be desired. Nor will the admission of errors of .152 and .120 (certainly not large in determinations of this kind) in the first and fourth experiments involve any serious doubt of the substantial accuracy of the second and third, when the difference of weight of the determinations is considered. Yet it is much to be desired that the relation expressed by (336), or with more generality by (334), should be tested by more numerous experiments.

It should be stated that the numbers in the column of pressures are not quite accurate. In the experiments of Deville and Troost the gas was subject to the actual atmospheric pressure at the time of the experiment. This varied from 747 to 764 millimeters of mercury. The precise pressure for each experiment is not given. In the experiments of Playfair and Wanklyn the mixture of nitrogen and peroxide of nitrogen was subject to the actual atmospheric pressure at the time of the experiment. The numbers in the column of pressures express the fraction of the whole pressure which remains after subtracting the part due to the free nitrogen. But no indication is given in the published account of the experiments in regard to the height of the barometer. Now it may easily be shown that a variation of  $\frac{1}{760}$  in the value of  $p$  can in no case cause a variation of more than .005 in the value of  $D$  as calculated by equation (336). In any of the experiments of Playfair and Wanklyn a variation of more than 30<sup>mm</sup> in the height of the barometer would be necessary to produce a variation of .01 in the value of  $D$ . The errors due to this source cannot therefore be very serious. They might have been avoided altogether in the discussion of the experiments of Deville and Troost by using instead of (336) a formula expressing the relation between the relative density, the temperature, and the actual density, as the reciprocal of the latter quantity is given for each ex-

periment of this series. It seemed best, however, to make a trifling sacrifice of accuracy for the sake of simplicity.

It might be thought that the experiments under discussion would be better represented by a formula in which the term containing  $\log t$  (see equation (333)) was retained. But an examination of the figures in the table will show that nothing important can be gained in this respect, and there is hardly sufficient motive for adding another term to the formula of calculation. Any attempt to determine the *real* values of  $A$ ,  $B'$ , and  $C$  in equation (333), (assuming the absolute validity of such an equation for peroxide of nitrogen,) from the experiments under discussion would be entirely misleading, as the reader may easily convince himself.

From equation (336), however, the following conclusions may be deduced. By comparison with (334) we obtain

$$A + \frac{B'}{\alpha_2} \log_{10} t - \frac{C}{t} = 9.47056 - \frac{3118.6}{t},$$

which must hold true approximately between the temperatures  $11^{\circ}$  and  $90^{\circ}$ . (At higher temperatures the relative densities vary too slowly with the temperatures to afford a critical test of the accuracy of this relation.) By differentiation we obtain

$$\frac{MB'}{\alpha_2 t} + \frac{C}{t^2} = \frac{3118.6}{t^2},$$

where  $M$  denotes the modulus of the common system of logarithms. Now by comparing equations (333) and (334) we see that

$$C = \frac{MC}{\alpha_2} = .43429 \frac{C}{\alpha_2}.$$

Hence

$$B't + C = 7181 \alpha_2 = 3590 \alpha_1,$$

which may be regarded as a close approximation at  $40^{\circ}$  or  $50^{\circ}$ , and a tolerable approximation between the limits of temperature above mentioned. Now  $B't + C$  represents the heat evolved by the conversion of a unit of  $\text{NO}_2$  into  $\text{N}_2\text{O}_4$  under constant pressure. Such conversion cannot take place at constant pressure without change of temperature, which renders the experimental verification of the last equation less simple. But since by equations (322)

$$B' = B + \alpha_1 - \alpha_2 = B + \frac{1}{2} \alpha_1,$$

we shall have for the temperature of  $40^{\circ}$

$$Bt + C = 3434 \alpha_1.$$

Now  $Bt + C$  represents the decrease of energy when a unit of  $\text{NO}_2$  is transformed into  $\text{N}_2\text{O}_4$  without change of temperature. It therefore

represents the excess of the heat evolved over the work done by external forces when a mass of the gas is compressed at constant temperature until a unit of  $\text{NO}_2$  has been converted into  $\text{N}_2\text{O}_4$ . This quantity will be constant if  $B = 0$ , i. e., if the specific heats at constant volume of  $\text{NO}_2$  and  $\text{N}_2\text{O}_4$  are the same. This assumption would be more simple from a theoretical stand-point and perhaps safer than the assumption that  $B' = 0$ . If  $B = 0$ ,  $B' = a_2$ . If we wish to embody this assumption in the equation between  $D$ ,  $p$ , and  $t$ , we may substitute

$$6.5228 + \log_{10} (t_c + 273) - \frac{2977.4}{t_c + 273}$$

for the second member of equation (336). The relative densities calculated by the equation thus modified from the temperatures and pressures of the experiments under discussion will not differ from those calculated from the unmodified equation by more than .002 in any case, or by more than .001 in the first series of experiments.

It is to be noticed that if we admit the validity of the volumetrical relation expressed by equation (333), which is evidently equivalent to an equation between  $p$ ,  $t$ ,  $v$ , and  $m$  (this letter denoting the quantity of the gas without reference to its molecular condition), or if we admit the validity of the equation only between certain limits of temperature and for densities less than a certain limit of density, and also admit that between the given limits of temperature the specific heat of the gas at constant volume may be regarded as a constant quantity when the gas is sufficiently rarefied to be regarded as consisting wholly of  $\text{NO}_2$ ,—or, to speak without reference to the molecular state of the gas, when it is rarefied until its relative density  $D$  approximates to its limiting value  $D_1$ ,—we must also admit the validity (within the same limits of temperature and density) of all the calorimetical relations which belong to ideal gas-mixtures with convertible components. The premises are evidently equivalent to this,—that we may imagine an ideal gas with convertible components such that between certain limits of temperature and above a certain limit of density the relation between  $p$ ,  $t$ , and  $v$  shall be the same for a unit of this ideal gas as for a unit of peroxide of nitrogen, and for a very great value of  $v$  (within the given limits of temperature) the thermal capacity at constant volume of the ideal and actual gases shall be the same. Let us regard  $t$  and  $v$  as independent variables; we may let these letters and  $p$  refer alike to the ideal and real gases, but we must distinguish the entropy  $\eta'$  of the ideal gas from the entropy  $\eta$  of the real gas. Now by (88)

$$\frac{d\eta}{dv} = \frac{dp}{dt}, \quad (337)$$

therefore

$$\frac{d}{dv} \frac{d\eta}{dt} = \frac{d}{dt} \frac{d\eta}{dv} = \frac{d}{dt} \frac{dp}{dt} = \frac{d^2 p}{dt^2}. \quad (338)$$

Since a similar relation will hold true for  $\eta'$ , we obtain

$$\frac{d}{dv} \frac{d\eta}{dt} = \frac{d}{dv} \frac{d\eta'}{dt}, \quad (339)$$

which must hold true within the given limits of temperature and density. Now it is granted that

$$\frac{d\eta}{dt} = \frac{d\eta'}{dt} \quad (340)$$

for very great values of  $v$  at any temperature within the given limits, (for the two members of the equation represent the thermal capacities at constant volume of the real and ideal gases divided by  $t$ ,) hence, in virtue of (339), this equation must hold true in general within the given limits of temperature and density. Again, as an equation like (337) will hold true of  $\eta'$ , we shall have

$$\frac{d\eta}{dv} = \frac{d\eta'}{dv}. \quad (341)$$

From the two last equations it is evident that in all calorimetric relations the ideal and real gases are identical. Moreover the energy and entropy of the ideal gas are evidently so far arbitrary that we may suppose them to have the same values as in the real gas for any given values of  $t$  and  $v$ . Hence the entropies of the two gases are the same within the given limits; and on account of the necessary relation

$$d\varepsilon = t d\eta - p dv,$$

the energies of the two gases are in like manner identical. Hence the fundamental equation between the energy, entropy, volume, and quantity of matter must be the same for the ideal gas as for the actual.

We may easily form a fundamental equation for an ideal gas-mixture with convertible components, which shall relate only to the phases of equilibrium. For this purpose, we may use the equations of the form (312) to eliminate from the equation of the form (273), which expresses the relation between the pressure, the temperature, and the potentials for the proximate components, as many of the potentials as there are equations of the former kind, leaving the

potentials for those components which it is convenient to regard as the ultimate components of the gas-mixture.

In the case of a binary gas-mixture with convertible components, the components will have the same potential, which may be denoted by  $\mu$ , and the fundamental equation will be

$$p = a_1 L_1 t^{\frac{c_1 + a_1}{a_1}} e^{\frac{\mu - E_1}{a_1 t}} + a_2 L_2 t^{\frac{c_2 + a_2}{a_2}} e^{\frac{\mu - E_2}{a_2 t}}, \quad (342)$$

where

$$L_1 = e^{\frac{H_1 - c_1 - a_1}{a_1}}, \quad L_2 = e^{\frac{H_2 - c_2 - a_2}{a_2}}. \quad (343)$$

From this equation, by differentiation and comparison with (98), we obtain

$$\begin{aligned} \frac{\eta}{v} &= L_1 \left( c_1 + a_1 - \frac{\mu - E_1}{t} \right) t^{\frac{c_1}{a_1}} e^{\frac{\mu - E_1}{a_1 t}} \\ &\quad + L_2 \left( c_2 + a_2 - \frac{\mu - E_2}{t} \right) t^{\frac{c_2}{a_2}} e^{\frac{\mu - E_2}{a_2 t}}, \end{aligned} \quad (244)$$

$$\frac{m}{v} = L_1 t^{\frac{c_1}{a_1}} e^{\frac{\mu - E_1}{a_1 t}} + L_2 t^{\frac{c_2}{a_2}} e^{\frac{\mu - E_2}{a_2 t}}. \quad (345)$$

From the general equation (93) with the preceding equations the following is easily obtained,—

$$\frac{\varepsilon}{v} = L_1 (c_1 t + E_1) t^{\frac{c_1}{a_1}} e^{\frac{\mu - E_1}{a_1 t}} + L_2 (c_2 t + E_2) t^{\frac{c_2}{a_2}} e^{\frac{\mu - E_2}{a_2 t}}. \quad (346)$$

We may obtain the relation between  $p$ ,  $t$ ,  $v$ , and  $m$  by eliminating  $\mu$  from (342) and (345). For this purpose we may proceed as follows. From (342) and (345) we obtain

$$p - a_2 t \frac{m}{v} = (a_1 - a_2) L_1 t^{\frac{c_1 + a_1}{a_1}} e^{\frac{\mu - E_1}{a_1 t}} \quad (347)$$

$$a_1 t \frac{m}{v} - p = (a_1 - a_2) L_2 t^{\frac{c_2 + a_2}{a_2}} e^{\frac{\mu - E_2}{a_2 t}}; \quad (348)$$

and from these equations we obtain

$$\begin{aligned} a_1 \log \left( p - a_2 t \frac{m}{v} \right) - a_2 \log \left( a_1 t \frac{m}{v} - p \right) &= (a_1 - a_2) \log (a_1 - a_2) \\ + a_1 \log L_1 - a_2 \log L_2 + (c_2 - c_1 + a_1 - a_2) \log t - \frac{E_1 - E_2}{t}. \end{aligned} \quad (349)$$

(In the particular case when  $a_1 = 2a_2$  this equation will be equivalent to (333)). By (347) and (348) we may easily eliminate  $\mu$  from (346).

The reader will observe that the relations thus deduced from the fundamental equation (342) without any reference to the different components of the gaseous mass are equivalent to those which relate to the phases of dissipated energy of a binary gas-mixture with components which are equivalent in substance but not convertible, except that the equations derived from (342) do not give the quantities of the proximate components, but relate solely to those properties which are capable of direct experimental verification without the aid of any theory of the constitution of the gaseous mass.

The practical application of these equations is rendered more simple by the fact that the ratio  $a_1 : a_2$  will always bear a simple relation to unity. When  $a_1$  and  $a_2$  are equal, if we write  $a$  for their common value, we shall have by (342) and (345)

$$pv = amt, \quad (350)$$

and by (345) and (346)

$$\frac{\varepsilon}{m} = \frac{\frac{c_2 - c_1}{a} \frac{E_1 - E_2}{e^{-at}}}{\frac{L_1(c_1 t + E_1) + L_2(c_2 t + E_2)t}{L_1 + L_2 t} \frac{c_2 - c_1}{a} \frac{E_1 - E_2}{e^{-at}}}. \quad (351)$$

By this equation we may calculate directly the amount of heat required to raise a given quantity of the gas from one given temperature to another at constant volume. The equation shows that the amount of heat will be independent of the volume of the gas. The heat necessary to produce a given change of temperature in the gas at constant pressure, may be found by taking the difference of the values of  $\chi$ , as defined by equation (89), for the initial and final states of the gas. From (89), (350), and (351) we obtain

$$\frac{\chi}{m} = \frac{\frac{c_2 - c_1}{a} \frac{E_1 - E_2}{e^{-at}}}{\frac{L_1(c_1 t + at + E_1) + L_2(c_2 t + at + E_2)t}{L_1 + L_2 t} \frac{c_2 - c_1}{a} \frac{E_1 - E_2}{e^{-at}}}. \quad (352)$$

By differentiation of the two last equations we may obtain directly the specific heats of the gas at constant volume and at constant pressure.

The fundamental equation of an ideal ternary gas-mixture with a single relation of convertibility between its components is

$$\begin{aligned}
 p = & a_1 e^{\frac{H_1 - c_1 - a_1}{a_1}} t^{\frac{c_1 + a_1}{a_1}} e^{\frac{\mu_1 - E_1}{a_1 t}} \\
 & + a_2 e^{\frac{H_2 - c_2 - a_2}{a_2}} t^{\frac{c_2 + a_2}{a_2}} e^{\frac{\mu_2 - E_2}{a_2 t}} \\
 & + a_3 e^{\frac{H_3 - c_3 - a_3}{a_3}} t^{\frac{c_3 + a_3}{a_3}} e^{\frac{\lambda_1 \mu_1 + \lambda_2 \mu_2 - E_3}{a_3 t}}, \quad (353)
 \end{aligned}$$

where  $\lambda_1$  and  $\lambda_2$  have the same meaning as on page 230.

(To be continued.)

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#### ERRATA.

Page 167, formula (168), for  $m_1$  read  $\mu_1$ .

" formula (169), for  $m_1, \dots, m_{n-1}$  read  $\mu_1, \dots, \mu_{n-1}$ .

Page 239, formula (333), for  $\frac{C}{t}$  read  $\frac{C}{a_2 t}$ .

## IX. ON THE EQUILIBRIUM OF HETEROGENEOUS SUBSTANCES.

BY J. WILLARD GIBBS.

(Continued from page 248).

THE CONDITIONS OF INTERNAL AND EXTERNAL EQUILIBRIUM FOR SOLIDS IN CONTACT WITH FLUIDS WITH REGARD TO ALL POSSIBLE STATES OF STRAIN OF THE SOLIDS.

In treating of the physical properties of a solid, it is necessary to consider its *state of strain*. A body is said to be *strained* when the relative position of its parts is altered, and by its *state of strain* is meant its state in respect to the relative position of its parts. We have hitherto considered the equilibrium of solids only in the case in which their state of strain is determined by pressures having the same values in all directions about any point. Let us now consider the subject without this limitation.

If  $x', y', z'$  are the rectangular co-ordinates of a point of a solid body in any completely determined state of strain, which we shall call the *state of reference*, and  $x, y, z$ , the rectangular co-ordinates of the same point of the body in the state in which its properties are the subject of discussion, we may regard  $x, y, z$  as functions of  $x', y', z'$ , the form of the functions determining the second state of strain. For brevity, we may sometimes distinguish the variable state, to which  $x, y, z$  relate, and the constant state (state of reference), to which  $x', y', z'$  relate, as the *strained* and the *unstrained* states; but it must be remembered that these terms have reference merely to the change of form or *strain* determined by the functions which express the relations of  $x, y, z$  and  $x', y', z'$ , and do not imply any particular physical properties in either of the two states, nor prevent their possible coincidence. The axes to which the co-ordinates  $x, y, z$ , and  $x', y', z'$  relate will be distinguished as the axes of  $X, Y, Z$ , and  $X', Y', Z'$ . It is not necessary, nor always convenient, to regard these systems of axes as identical, but they should be similar, i. e., capable of superposition.

The state of strain of any element of the body is determined by the values of the differential coefficients of  $x, y$ , and  $z$  with respect to  $x', y',$  and  $z'$ ; for changes in the values of  $x, y, z$ , when the differential coefficients remain the same, only cause motions of translation of the

body. When the differential coefficients of the first order do not vary sensibly except for distances greater than the radius of sensible molecular action, we may regard them as completely determining the state of strain of any element. There are nine of these differential coefficients, viz.,

$$\left. \begin{array}{c} \frac{dx}{dx'} \\ \frac{dy}{dx'} \\ \frac{dz}{dx'} \\ \hline \end{array} \quad \begin{array}{c} \frac{dx}{dy'} \\ \frac{dy}{dy'} \\ \frac{dz}{dy'} \\ \hline \end{array} \quad \begin{array}{c} \frac{dx}{dz'} \\ \frac{dy}{dz'} \\ \frac{dz}{dz'} \\ \hline \end{array} \right\} \quad (354)$$

It will be observed that these quantities determine the orientation of the element as well as its strain, and both these particulars must be given in order to determine the nine differential coefficients. Therefore, since the orientation is capable of three independent variations, which do not affect the strain, the strain of the element, considered without regard to directions in space, must be capable of six independent variations.

The physical state of any given element of a solid in any unvarying state of strain is capable of one variation, which is produced by addition or subtraction of heat. If we write  $\varepsilon_v$  and  $\eta_v$  for the energy and entropy of the element divided by its volume in the state of reference, we shall have for any constant state of strain

$$\delta\varepsilon_v = t \delta\eta_v$$

But if the strain varies, we may consider  $\varepsilon_v$  as a function of  $\eta_v$ , and the nine quantities in (354), and may write

$$\begin{aligned} \delta\varepsilon_v = & t \delta\eta_v + X_{xv} \delta \frac{dx}{dx'} + X_{yv} \delta \frac{dy}{dy'} + X_{zv} \delta \frac{dz}{dz'} \\ & + Y_{xv} \delta \frac{dy}{dx'} + Y_{yv} \delta \frac{dy}{dy'} + Y_{zv} \delta \frac{dy}{dz'} \\ & + Z_{xv} \delta \frac{dz}{dx'} + Z_{yv} \delta \frac{dz}{dy'} + Z_{zv} \delta \frac{dz}{dz'}, \end{aligned} \quad (355)$$

where  $X_{xv}, \dots, Z_{zv}$  denote the differential coefficients of  $\varepsilon_v$ , taken with respect to  $\frac{dx}{dx'}, \dots, \frac{dz}{dz'}$ . The physical signification of these quantities will be apparent, if we apply the formula to an element which in the state of reference is a right parallelopiped having the edges  $dx', dy', dz'$ , and suppose that in the strained state the face in which  $x'$  has the smaller constant value remains fixed, while the opposite face is moved parallel to the axis of  $X$ . If we also suppose

no heat to be imparted to the element, we shall have, on multiplying by  $dx' dy' dz'$ ,

$$\delta \varepsilon_v dx' dy' dz' = X_x \delta \frac{dx}{dx'} dx' dy' dz'.$$

Now the first member of this equation evidently represents the work done upon the element by the surrounding elements; the second member must therefore have the same value. Since we must regard the forces acting on opposite faces of the elementary parallelopiped as equal and opposite, the whole work done will be zero except for the face which moves parallel to  $X$ . And since  $\delta \frac{dx}{dx'} dx'$  represents the distance moved by this face,  $X_x dy' dz'$  must be equal to the component parallel to  $X$  of the force acting upon this face. In general, therefore, if by the positive side of a surface for which  $x'$  is constant we understand the side on which  $x'$  has the greater value, we may say that  $X_x$  denotes the component parallel to  $X$  of the force exerted by the matter on the positive side of a surface for which  $x'$  is constant upon the matter on the negative side of that surface per unit of the surface measured in the state of reference. The same may be said, *mutatis mutandis*, of the other symbols of the same type.

It will be convenient to use  $\Sigma$  and  $\Sigma'$  to denote summation with respect to quantities relating to the axes  $X, Y, Z$ , and to the axes  $X', Y', Z'$ , respectively. With this understanding we may write

$$\delta \varepsilon_v = t \delta \eta_v + \Sigma \Sigma' \left( X_x \delta \frac{dx}{dx'} \right). \quad (356)$$

This is the complete value of the variation of  $\varepsilon_v$  for a given element of the solid. If we multiply by  $dx' dy' dz'$ , and take the integral for the whole body, we shall obtain the value of the variation of the total energy of the body, when this is supposed invariable in substance. But if we suppose the body to be increased or diminished in substance at its surfaces (the increment being continuous in nature and state with the part of the body to which it is joined), to obtain the complete value of the variation of the energy of the body, we must add the integral

$$\int' \varepsilon_v \delta N' Ds'$$

in which  $Ds'$  denotes an element of the surface measured in the state of reference, and  $\delta N'$  the change in position of this surface (due to the substance added or taken away) measured normally and outward in the state of reference. The complete value of the variation of the intrinsic energy of the solid is therefore

$$\int \int \int t \delta \eta_v dx' dy' dz' + \int \int \int \Sigma \Sigma' \left( X_{xv} \delta \frac{dx'}{dx} \right) dx' dy' dz' + \int \varepsilon_v \delta N' D s'. \quad (357)$$

This is entirely independent of any supposition in regard to the homogeneity of the solid.

To obtain the conditions of equilibrium for solid and fluid masses in contact, we should make the variation of the energy of the whole equal to or greater than zero. But since we have already examined the conditions of equilibrium for fluids, we need here only seek the conditions of equilibrium for the interior of a solid mass and for the surfaces where it comes in contact with fluids. For this it will be necessary to consider the variations of the energy of the fluids only so far as they are immediately connected with the changes in the solid. We may suppose the solid with so much of the fluid as is in close proximity to it to be enclosed in a fixed envelop, which is impermeable to matter and to heat, and to which the solid is firmly attached wherever they meet. We may also suppose that in the narrow space or spaces between the solid and the envelop, which are filled with fluid, there is no motion of matter or transmission of heat across any surfaces which can be generated by moving normals to the surface of the solid, since the terms in the condition of equilibrium relating to such processes may be cancelled on account of the internal equilibrium of the fluids. It will be observed that this method is perfectly applicable to the case in which a fluid mass is entirely enclosed in a solid. A detached portion of the envelop will then be necessary to separate the great mass of the fluid from the small portion adjacent to the solid, which alone we have to consider. Now the variation of the energy of the fluid mass will be, by equation (13),

$$\int^F t \delta D\eta - \int^F p \delta Dv + \Sigma_1 \int^F \mu_1 \delta Dm_1, \quad (358)$$

where  $\int^F$  denotes an integration extending over all the elements of the fluid (within the envelop), and  $\Sigma_1$  denotes a summation with regard to those independently variable components of the fluid of which the solid is composed. Where the solid does not consist of substances which are components, actual or possible (see page 117), of the fluid, this term is of course to be cancelled.

If we wish to take account of gravity, we may suppose that it acts in the negative direction of the axis of  $Z$ . It is evident that the variation of the energy due to gravity for the whole mass considered is simply

$$\int \int \int g I'' \delta z dx' dy' dz', \quad (359)$$

where  $g$  denotes the force of gravity, and  $I''$  the density of the

element in the state of reference, and the triple integration, as before, extends throughout the solid.

We have, then, for the general condition of equilibrium,

$$\begin{aligned} & \int \int \int t \delta \eta_v dx' dy' dz' + \int \int \int \Sigma \Sigma' \left( X_x, \delta \frac{dx}{dx'} \right) dx' dy' dz' \\ & + \int \int \int g \Gamma' \delta z dx' dy' dz' + \int \varepsilon_v \delta N' Ds' \\ & + \int^F t \delta D\eta - \int^F p \delta Dv + \Sigma_1 \int^F \mu_1 \delta Dm_1 \geq 0. \end{aligned} \quad (360)$$

The equations of condition to which these variations are subject are: (1) that which expresses the constancy of the total entropy,

$$\int \int \int \delta \eta_v dx' dy' dz' + \int \eta_v \delta N' Ds' + \int^F \delta D\eta = 0; \quad (361)$$

(2) that which expresses how the value of  $\delta Dv$  for any element of the fluid is determined by changes in the solid,

$$\delta Dv = -(\alpha \delta x + \beta \delta y + \gamma \delta z) Ds - v_v \delta N' Ds', \quad (362)$$

where  $\alpha, \beta, \gamma$  denote the direction cosines of the normal to the surface of the body in the state to which  $x, y, z$  relate,  $Ds$  the element of the surface in this state corresponding to  $Ds'$  in the state of reference, and  $v_v$ , the volume of an element of the solid divided by its volume in the state of reference;

(3) those which express how the values of  $\delta Dm_1, \delta Dm_2$ , etc. for any element of the fluid are determined by the changes in the solid,

$$\left. \begin{aligned} \delta Dm_1 &= -\Gamma'_1 \delta N' Ds', \\ \delta Dm_2 &= -\Gamma'_2 \delta N' Ds', \\ \text{etc.,} \end{aligned} \right\} \quad (363)$$

where  $\Gamma'_1, \Gamma'_2$ , etc. denote the separate densities of the several components in the solid in the state of reference.

Now, since the variations of entropy are independent of all the other variations, the condition of equilibrium (360), considered with regard to the equation of condition (361), evidently requires that throughout the whole system

$$t = \text{const.} \quad (364)$$

We may therefore use (361) to eliminate the first and fifth integrals from (360). If we multiply (362) by  $p$ , and take the integrals for the whole surface of the solid and for the fluid in contact with it, we obtain the equation

$$\int^F p \delta Dv = -\int^F p (\alpha \delta x + \beta \delta y + \gamma \delta z) Ds - \int^F p v_v \delta N' Ds', \quad (365)$$

by means of which we may eliminate the sixth integral from (360). If we add equations (363) multiplied respectively by  $\mu_1, \mu_2$ , etc., and take the integrals, we obtain the equation

$$\Sigma_1 f^r \mu_1 \delta Dm_1 = - f' \Sigma_1 (\mu_1 \Gamma_1') \delta N' Ds', \quad (366)$$

by means of which we may eliminate the last integral from (360).

The condition of equilibrium is thus reduced to the form

$$\begin{aligned} & \int \int \int \Sigma \Sigma' \left( X_{x'} \delta \frac{dx}{dx'} \right) dx' dy' dz' + \int \int \int g \Gamma' \delta z dx' dy' dz' \\ & + f' \varepsilon_v \delta N' Ds' - f' t \eta_v \delta N' Ds' + f' p (\alpha \delta x + \beta \delta y + \gamma \delta z) Ds \\ & + f' p v_v \delta N' Ds' - f' \Sigma_1 (\mu_1 \Gamma_1') \delta N' Ds' \geq 0, \end{aligned} \quad (367)$$

in which the variations are independent of the equations of condition, and in which the only quantities relating to the fluids are  $p$  and  $\mu_1$ ,  $\mu_2$ , etc.

Now by the ordinary method of the calculus of variations, if we write  $\alpha'$ ,  $\beta'$ ,  $\gamma'$  for the direction-cosines of the normal to the surface of the solid in the state of reference, we have

$$\begin{aligned} & \int \int \int X_{x'} \delta \frac{dx}{dx'} dx' dy' dz' \\ & = f' \alpha' X_{x'} \delta x Ds' - \int \int \int \frac{dX_{x'}}{dx'} \delta x dx' dy' dz', \end{aligned} \quad (368)$$

with similar expressions for the other parts into which the first integral in (367) may be divided. The condition of equilibrium is thus reduced to the form

$$\begin{aligned} & - \int \int \int \Sigma \Sigma' \left( \frac{dX_{x'}}{dx'} \delta x \right) dx' dy' dz' + \int \int \int g \Gamma' \delta z dx' dy' dz' \\ & + f \Sigma \Sigma' (\alpha' X_{x'} \delta x) Ds' + f' p \Sigma (\alpha \delta x) Ds \\ & + f' [\varepsilon_v - t \eta_v + p v_v - \Sigma_1 (\mu_1 \Gamma_1')] \delta N' Ds' \geq 0. \end{aligned} \quad (369)$$

It must be observed that if the solid mass is not continuous throughout in nature and state, the surface-integral in (368), and therefore the first surface-integral in (369), must be taken to apply not only to the external surface of the solid, but also to every surface of discontinuity within it, and that with reference to each of the two masses separated by the surface. To satisfy the condition of equilibrium, as thus understood, it is necessary and sufficient that throughout the solid mass

$$\Sigma \Sigma' \left( \frac{dX_{x'}}{dx'} \delta x \right) - g \Gamma' \delta z = 0; \quad (370)$$

that throughout the surfaces where the solid meets the fluid

$$Ds' \Sigma \Sigma' (\alpha' X_{x'} \delta x) + Dsp \Sigma (\alpha \delta x) = 0, \quad (371)$$

and

$$[\varepsilon_v - t \eta_v + p v_v - \Sigma_1 (\mu_1 \Gamma_1')] \delta N' \geq 0; \quad (372)$$

and that throughout the internal surfaces of discontinuity

$$\Sigma \Sigma' (\alpha' X_{x'} \delta x) + \Sigma \Sigma' (\alpha' Y_{x'} \delta x) = 0, \quad (373)$$

where the suffixed numerals distinguish the expressions relating to the masses on opposite sides of a surface of discontinuity.

Equation (370) expresses the mechanical conditions of internal equilibrium for a continuous solid under the influence of gravity. If we expand the first term, and set the coefficients of  $\delta x$ ,  $\delta y$ , and  $\delta z$  separately equal to zero, we obtain

$$\left. \begin{aligned} \frac{dX_{x'}}{dx'} + \frac{dX_{y'}}{dy'} + \frac{dX_{z'}}{dz'} &= 0, \\ \frac{dY_{x'}}{dx'} + \frac{dY_{y'}}{dy'} + \frac{dY_{z'}}{dz'} &= 0, \\ \frac{dZ_{x'}}{dx'} + \frac{dZ_{y'}}{dy'} + \frac{dZ_{z'}}{dz'} &= g I'. \end{aligned} \right\} \quad (374)$$

The first member of any one of these equations multiplied by  $dx' dy' dz'$  evidently represents the sum of the components parallel to one of the axes  $X$ ,  $Y$ ,  $Z$  of the forces exerted on the six faces of the element  $dx' dy' dz'$  by the neighboring elements.

As the state which we have called the state of reference is arbitrary, it may be convenient for some purposes to make it coincide with the state to which  $x$ ,  $y$ ,  $z$  relate, and the axes  $X'$ ,  $Y'$ ,  $Z'$  with the axes  $X$ ,  $Y$ ,  $Z$ . The values of  $X_{x'}$ , . . .  $Z_{z'}$ , on this particular supposition may be represented by the symbols  $X_x$ , . . .  $Z_z$ . Since

$$X_y = \frac{d\varepsilon_{v_x}}{d\frac{dx}{dy}}, \quad \text{and} \quad Y_x = \frac{d\varepsilon_{v_y}}{d\frac{dy}{dx}},$$

and since, when the states  $x$ ,  $y$ ,  $z$  and  $x'$ ,  $y'$ ,  $z'$  coincide, and the axes  $X$ ,  $Y$ ,  $Z$ , and  $X'$ ,  $Y'$ ,  $Z'$ ,  $d\frac{dx}{dy}$  and  $d\frac{dy}{dx}$  represent displacements which differ only by a rotation, we must have

$$X_y = Y_x, \quad (375)$$

and for similar reasons,

$$Y_z = Z_y, \quad Z_x = X_z. \quad (376)$$

The six quantities  $X_x$ ,  $Y_y$ ,  $Z_z$ ,  $X_y$  or  $Y_x$ ,  $Y_z$  or  $Z_y$ , and  $Z_x$  or  $X_z$  are called the *rectangular components of stress*, the three first being the *longitudinal stresses* and the three last the *shearing stresses*. The mechanical conditions of internal equilibrium for a solid under the influence of gravity may therefore be expressed by the equations

$$\left. \begin{aligned} \frac{dX_x}{dx} + \frac{dX_y}{dy} + \frac{dX_z}{dz} &= 0, \\ \frac{dY_x}{dx} + \frac{dY_y}{dy} + \frac{dY_z}{dz} &= 0, \\ \frac{dZ_x}{dx} + \frac{dZ_y}{dy} + \frac{dZ_z}{dz} &= g' I, \end{aligned} \right\} \quad (377)$$

where  $I'$  denotes the density of the element to which the other symbols relate. Equations (375), (376) are rather to be regarded as expressing necessary relations (when  $X_x, \dots, Z_z$  are regarded as internal forces determined by the state of strain of the solid) than as expressing conditions of equilibrium. They will hold true of a solid which is not in equilibrium,—of one, for example, through which vibrations are propagated,—which is not the case with equations (377).

Equation (373) expresses the mechanical conditions of equilibrium for a surface of discontinuity within the solid. If we set the coefficients of  $\delta x, \delta y, \delta z$ , separately equal to zero we obtain

$$\left. \begin{aligned} (\alpha' X_x + \beta' X_y + \gamma' X_z)_1 + (\alpha' X_x + \beta' X_y + \gamma' X_z)_2 &= 0, \\ (\alpha' Y_x + \beta' Y_y + \gamma' Y_z)_1 + (\alpha' Y_x + \beta' Y_y + \gamma' Y_z)_2 &= 0, \\ (\alpha' Z_x + \beta' Z_y + \gamma' Z_z)_1 + (\alpha' Z_x + \beta' Z_y + \gamma' Z_z)_2 &= 0. \end{aligned} \right\} \quad (378)$$

Now when the  $\alpha', \beta', \gamma'$  represent the direction-cosines of the normal in the state of reference on the positive side of any surface within the solid, an expression of the form

$$\alpha' X_x + \beta' X_y + \gamma' X_z \quad (379)$$

represents the component parallel to  $X$  of the force exerted upon the surface in the strained state by the matter on the positive side per unit of area measured in the state of reference. This is evident from the consideration that in estimating the force upon any surface we may substitute for the given surface a broken one consisting of elements for each of which either  $x'$  or  $y'$  or  $z'$  is constant. Applied to a surface bounding a solid, or any portion of a solid which may not be continuous with the rest, when the normal is drawn outward as usual, the same expression taken negatively represents the component parallel to  $X$  of the force exerted upon the surface (per unit of surface measured in the state of reference) by the interior of the solid, or of the portion considered. Equations (378) therefore express the condition that the force exerted upon the surface of discontinuity by the matter on one side and determined by its state of strain shall be equal and opposite to that exerted by the matter on the other side. Since

$$(\alpha')_1 = -(\alpha')_2, \quad (\beta')_1 = -(\beta')_2, \quad (\gamma')_1 = -(\gamma')_2,$$

we may also write

$$\left. \begin{aligned} \alpha'(X_x)_1 + \beta'(X_y)_1 + \gamma'(X_z)_1 &= \alpha'(X_x)_2 + \beta'(X_y)_2 + \gamma'(X_z)_2, \\ \text{etc.,} \end{aligned} \right\} \quad (380)$$

where the signs of  $\alpha'$ ,  $\beta'$ ,  $\gamma'$  may be determined by the normal on either side of the surface of discontinuity.

Equation (371) expresses the mechanical condition of equilibrium for a surface where the solid meets a fluid. It involves the separate equations

$$\left. \begin{aligned} \alpha' X_x + \beta' Y_x + \gamma' Z_x &= -\alpha p \frac{Ds}{Ds'}, \\ \alpha' Y_x + \beta' Y_y + \gamma' Z_y &= -\beta p \frac{Ds}{Ds'}, \\ \alpha' Z_x + \beta' Z_y + \gamma' Z_z &= -\gamma p \frac{Ds}{Ds'}, \end{aligned} \right\} \quad (381)$$

the fraction  $\frac{Ds}{Ds'}$  denoting the ratio of the areas of the same element of the surface in the strained and unstrained states of the solid. These equations evidently express that the force exerted by the interior of the solid upon an element of its surface, and determined by the strain of the solid, must be normal to the surface and equal (but acting in the opposite direction) to the pressure exerted by the fluid upon the same element of surface.

If we wish to replace  $\alpha$  and  $Ds$  by  $\alpha'$ ,  $\beta'$ ,  $\gamma'$ , and the quantities which express the strain of the element, we may make use of the following considerations. The product  $\alpha Ds$  is the projection of the

element  $Ds$  on the  $Y$ - $Z$  plane. Now since the ratio  $\frac{Ds}{Ds'}$  is independent of the form of the element, we may suppose that it has any convenient form. Let it be bounded by the three surfaces  $x' = \text{const.}$ ,  $y' = \text{const.}$ ,  $z' = \text{const.}$ , and let the parts of each of these surfaces included by the two others with the surface of the body be denoted by  $L$ ,  $M$ , and  $N$ , or by  $L'$ ,  $M'$  and  $N'$ , according as we have reference to the strained or unstrained state of the body. The areas of  $L'$ ,  $M'$ , and  $N'$  are evidently  $\alpha' Ds'$ ,  $\beta' Ds'$ , and  $\gamma' Ds'$ ; and the sum of the projections of  $L$ ,  $M$  and  $N$  upon any plane is equal to the projection of  $Ds$  upon that plane, since  $L$ ,  $M$ , and  $N$  with  $Ds$  include a solid figure. (In propositions of this kind the *sides* of surfaces must be distinguished. If the normal to  $Ds$  falls outward from the small solid figure, the normals to  $L$ ,  $M$ , and  $N$  must fall inward, and vice

*versa*). Now  $L'$  is a right-angled triangle of which the perpendicular sides may be called  $dy'$  and  $dz'$ . The projection of  $L$  on the  $Y\text{-}Z$  plane will be a triangle, the angular points of which are determined by the co-ordinates

$$y, z; \quad y + \frac{dy}{dy'} dy', \quad z + \frac{dz}{dy'} dy'; \quad y + \frac{dy}{dz'} dz', \quad z + \frac{dz}{dz'} dz';$$

the area of such a triangle is

$$\cdot \frac{1}{2} \left( \frac{dy}{dy'} \frac{dz}{dz'} - \frac{dz}{dy'} \frac{dy}{dz'} \right) dy' dz',$$

or, since  $\frac{1}{2} dy' dz'$  represents the area of  $L'$ ,

$$\left( \frac{dy}{dy'} \frac{dz}{dz'} - \frac{dz}{dy'} \frac{dy}{dz'} \right) \alpha' Ds'.$$

(That this expression has the proper sign will appear if we suppose for the moment that the strain vanishes.) The areas of the projections of  $M$  and  $N$  upon the same plane will be obtained by changing  $y', z'$  and  $\alpha'$  in this expression into  $z', x'$ , and  $\beta'$ , and into  $x', y'$ , and  $\gamma'$ . The sum of the three expressions may be substituted for  $\alpha' Ds$  in (381).

We shall hereafter use  $\Sigma'$  to denote the sum of the three terms obtained by rotary substitutions of quantities relating to the axes  $X', Y', Z'$ , (i. e., by changing  $x', y', z'$  into  $y', z', x'$ , and into  $z', x', y'$ , with similar changes in regard to  $\alpha', \beta', \gamma'$ , and other quantities relating to these axes,) and  $\Sigma$  to denote the sum of the three terms obtained by similar rotary changes of quantities relating to the axes  $X, Y, Z$ . This is only an extension of our previous use of these symbols.

With this understanding, equations (381) may be reduced to the form

$$\left. \begin{aligned} \Sigma' (\alpha' X_x) + p \Sigma' \left\{ \alpha' \left( \frac{dy}{dy'} \frac{dz}{dz'} - \frac{dz}{dy'} \frac{dy}{dz'} \right) \right\} = 0, \\ \text{etc.} \end{aligned} \right\} \quad (382)$$

The formula (372) expresses the additional condition of equilibrium which relates to the dissolving of the solid, or its growth without discontinuity. If the solid consists entirely of substances which are actual components of the fluid, and there are no passive resistances which impede the formation or dissolving of the solid,  $\delta V'$  may have either positive or negative values, and we must have

$$\varepsilon_v - t \eta_v + p v_v = \Sigma_1 (\mu_1 \Gamma_1'). \quad (383)$$

But if some of the components of the solid are only possible com-

ponents (see page 117) of the fluid,  $\delta N'$  is incapable of positive values, as the quantity of the solid cannot be increased, and it is sufficient for equilibrium that

$$\varepsilon_v - t \eta_v + p_v \leq \Sigma_1 (\mu_1 \Gamma'_1). \quad (384)$$

To express condition (383) in a form independent of the state of reference, we may use  $\varepsilon_v$ ,  $\eta_v$ ,  $\Gamma'_1$ , etc., to denote the densities of energy, of entropy, and of the several component substances in the *variable* state of the solid. We shall obtain, on dividing the equation by  $v_v$ ,

$$\varepsilon_v - t \eta_v + p = \Sigma_1 (\mu_1 \Gamma'_1). \quad (385)$$

It will be remembered that the summation relates to the several components of the solid. If the solid is of uniform composition throughout, or if we only care to consider the contact of the solid and the fluid at a single point, we may treat the solid as composed of a single substance. If we use  $\mu_1$  to denote the potential for this substance in the fluid, and  $\Gamma$  to denote the density of the solid in the variable state, ( $\Gamma'$ , as before denoting its density in the state of reference,) we shall have

$$\varepsilon_v - t \eta_v + p v_v = \mu_1 \Gamma', \quad (386)$$

and

$$\varepsilon_v - t \eta_v + p = \mu_1 \Gamma. \quad (387)$$

To fix our ideas in discussing this condition, let us apply it to the case of a solid body which is homogeneous in nature and in state of strain. If we denote by  $\varepsilon$ ,  $\eta$ ,  $v$ , and  $m$ , its energy, entropy, volume, and mass, we have

$$\varepsilon - t \eta + p v = \mu_1 m. \quad (388)$$

Now the mechanical conditions of equilibrium for the surface where a solid meets a fluid require that the traction upon the surface determined by the state of strain of the solid shall be normal to the surface. This condition is always satisfied with respect to three surfaces at right angles to one another. In proving this well known proposition, we shall lose nothing in generality, if we make the state of reference, which is arbitrary, coincident with the state under discussion, the axes to which these states are referred being also coincident. We shall then have, for the normal component of the traction per unit of surface across any surface for which the direction-cosines of the normal are  $\alpha$ ,  $\beta$ ,  $\gamma$ , [compare (379), and for the notation  $X_x$ , etc., page 349,]

$$\begin{aligned} S = & \alpha(\alpha X_x + \beta Y_x + \gamma Z_x) \\ & + \beta(\alpha X_y + \beta Y_y + \gamma Z_y) \\ & + \gamma(\alpha X_z + \beta Y_z + \gamma Z_z), \end{aligned}$$

or, by (375), (376),

$$\begin{aligned} S = & \alpha^2 X_x + \beta^2 Y_y + \gamma^2 Z_z \\ & + 2\alpha\beta X_y + 2\beta\gamma Y_z + 2\gamma\alpha Z_x. \end{aligned} \quad (389)$$

We may also choose any convenient directions for the co-ordinate axes. Let us suppose that the direction of the axis of  $X$  is so chosen that the value of  $S$  for the surface perpendicular to this axis is as great as for any other surface, and that the direction of the axis of  $Y$  (supposed at right angles to  $X$ ) is such that the value of  $S$  for the surface perpendicular to it is as great as for any other surface passing through the axis of  $X$ . Then, if we write  $\frac{dS}{d\alpha}, \frac{dS}{d\beta}, \frac{dS}{d\gamma}$  for the differential coefficients derived from the last equation by treating  $\alpha, \beta$ , and  $\gamma$  as *independent* variables,

$$\frac{dS}{d\alpha} d\alpha + \frac{dS}{d\beta} d\beta + \frac{dS}{d\gamma} d\gamma = 0,$$

when  $\alpha d\alpha + \beta d\beta + \gamma d\gamma = 0$ ,

and  $\alpha = 1, \beta = 0, \gamma = 0$ .

That is,  $\frac{dS}{d\beta} = 0$ , and  $\frac{dS}{d\gamma} = 0$ ,

when  $\alpha = 1, \beta = 0, \gamma = 0$ .

Hence,  $X_y = 0$ , and  $Z_x = 0$ . (390)

Moreover,  $\frac{dS}{d\beta} d\beta + \frac{dS}{d\gamma} d\gamma = 0$ ,

when  $\alpha = 0, d\alpha = 0$ ,

$$\beta d\beta + \gamma d\gamma = 0,$$

and  $\beta = 1, \gamma = 0$ .

Hence  $Y_z = 0$ . (391)

Therefore, when the co-ordinate axes have the supposed directions, which are called the *principal axes of stress*, the rectangular components of the traction across any surface ( $\alpha, \beta, \gamma$ ) are by (379)

$$\alpha X_x, \beta Y_y, \gamma Z_z. \quad (392)$$

Hence, the traction across any surface will be normal to that surface,—

(1), when the surface is perpendicular to a principal axis of stress;

(2), if two of the *principal tractions*  $X_x$ ,  $Y_y$ ,  $Z_z$  are equal, when the surface is perpendicular to the plane containing the two corresponding axes, (in this case the traction across any such surface is equal to the common value of the two principal tractions);

(3), if the principal tractions are all equal, the traction is normal and constant for all surfaces.

It will be observed that in the second and third cases the position of the principal axes of stress are partially or wholly indeterminate, (so that these cases may be regarded as included in the first,) but the values of the principal tractions are always determinate, although not always different.

If, therefore, a solid which is homogeneous in nature and in state of strain is bounded by six surfaces perpendicular to the principal axes of strain, the mechanical conditions of equilibrium for these surfaces may be satisfied by the contact of fluids having the proper pressures, [see (381),] which will in general be different for the different pairs of opposite sides, and may be denoted by  $p'$ ,  $p''$ ,  $p'''$ . (These pressures are equal to the principal tractions of the solid taken negatively.) It will then be necessary for equilibrium with respect to the tendency of the solid to dissolve that the potential for the substance of the solid in the fluids shall have values  $\mu_1'$ ,  $\mu_1''$ ,  $\mu_1'''$  determined by the equations

$$\varepsilon - t\eta + p'v = \mu_1'm, \quad (393)$$

$$\varepsilon - t\eta + p''v = \mu_1''m, \quad (394)$$

$$\varepsilon - t\eta + p'''v = \mu_1'''m. \quad (395)$$

These values, it will be observed, are entirely determined by the nature and state of the solid, and their differences are equal to the differences of the corresponding pressures divided by the density of the solid.

It may be interesting to compare one of these potentials, as  $\mu_1'$ , with the potential (for the same substance) in a fluid of the same temperature  $t$  and pressure  $p'$  which would be in equilibrium with the same solid subjected on all sides to the uniform pressure  $p'$ . If we write  $[\varepsilon]_{p'}$ ,  $[\eta]_{p'}$ ,  $[v]_{p'}$ , and  $[\mu_1]_{p'}$  for the values which  $\varepsilon$ ,  $\eta$ ,  $v$ , and  $\mu_1$  would receive on this supposition, we shall have

$$[\varepsilon]_{p'} - t[\eta]_{p'} + p'[v]_{p'} = [\mu_1]_{p'}m. \quad (396)$$

Subtracting this from (393), we obtain

$$\varepsilon - [\varepsilon]_{p'} - t\eta + t[\eta]_{p'} + p'v - p'[v]_{p'} = \mu_1m - [\mu_1]_{p'}m. \quad (397)$$

Now it follows immediately from the definitions of energy and entropy

that the first four terms of this equation represent the work spent upon the solid in bringing it from the state of hydrostatic stress to the other state without change of temperature, and  $p' v - p'[v]_h$  evidently denotes the work done in displacing a fluid of pressure  $p'$  surrounding the solid during the operation. Therefore, the first number of the equation represents the total work done in bringing the solid *when surrounded by a fluid of pressure  $p'$*  from the state of hydrostatic stress  $p'$  to the state of stress  $p', p'', p'''$ . This quantity is necessarily positive, except of course in the limiting case when  $p' = p'' = p'''$ . If the quantity of matter of the solid body be unity, the increase of the potential in the fluid on the side of the solid on which the pressure remains constant, which will be necessary to maintain equilibrium, is equal to the work done as above described. Hence,  $\mu_1'$  is greater than  $[\mu_1]_{p'}$ , and for similar reasons,  $\mu_1''$  is greater than the value of the potential which would be necessary for equilibrium if the solid were subjected to the uniform pressure  $p''$ , and  $\mu_1'''$  greater than that which would be necessary for equilibrium if the solid were subjected to the uniform pressure  $p'''$ . That is, (if we adapt our language to what we may regard as the most general case, viz., that in which the fluids contain the substance of the solid but are not wholly composed of that substance,) the fluids in equilibrium with the solid are all supersaturated with respect to the substance of the solid, except when the solid is in a state of hydrostatic stress; so that if there were present in any one of these fluids any small fragment of the same kind of solid subject to the hydrostatic pressure of the fluid, such a fragment would tend to increase. Even when no such fragment is present, although there must be perfect equilibrium so far as concerns the tendency of the solid to dissolve or to increase by the accretion of similarly strained matter, yet the presence of the solid which is subject to the distorting stresses, will doubtless facilitate the commencement of the formation of a solid of hydrostatic stress upon its surface, to the same extent, perhaps, in the case of an amorphous body, as if it were itself subject only to hydrostatic stress. This may sometimes, or perhaps generally, make it a necessary condition of equilibrium in cases of contact between a fluid and an amorphous solid which can be formed out of it that the solid at the surface where it meets the fluid shall be sensibly in a state of hydrostatic stress.

But in the case of a crystalline solid, subjected to distorting stresses and in contact with solutions satisfying the conditions deduced above, although crystals of hydrostatic stress would doubtless commence to

form upon its surface (if the distorting stresses and consequent supersaturation of the fluid should be carried too far), before they would commence to be formed within the fluid or on the surface of most other bodies, yet within certain limits the relations expressed by equations (393)–(395) must admit of realization, especially when the solutions are such as can be easily supersaturated.\*

It may be interesting to compare the variations of  $\rho$ , the pressure in the fluid which determines in part the stresses and the state of strain of the solid, with other variations of the stresses or strains in the solid, with respect to the relation expressed by equation (388). To examine this point with complete generality, we may proceed in the following manner.

Let us consider so much of the solid as has in the state of reference the form of a cube, the edges of which are equal to unity, and parallel to the co-ordinate axes. We may suppose this body to be homogeneous in nature and in state of strain both in its state of reference and in its variable state. (This involves no loss of generality, since we may make the unit of length as small as we choose.) Let the fluid meet the solid on one or both of the surfaces for which  $Z'$  is constant. We may suppose these surfaces to remain perpendicular to the axis of  $Z$  in the variable state of the solid, and the edges in which  $y'$  and  $z'$  are both constant to remain parallel to the axis of  $X$ . It will be observed that these suppositions only fix the position of the strained body relatively to the co-ordinate axes, and do not in any way limit its state of strain.

It follows from the suppositions which we have made that

$$\frac{dz}{dx'} = \text{const.} = 0, \quad \frac{dz}{dy'} = \text{const.} = 0, \quad \frac{dy}{dx'} = \text{const.} = 0; \quad (398)$$

and

$$X_{zz} = 0, \quad Y_{zz} = 0, \quad Z_{zz} = -p \frac{dx}{dx'} \frac{dy}{dy'} \quad (399)$$

Hence, by (355),

$$d\varepsilon_{yy} = t d\eta_{yy} + X_{yy} t \frac{dx}{dx'} + X_{yy} t \frac{dx}{dy'} + Y_{yy} t \frac{dy}{dy'} - p \frac{dx}{dx'} \frac{dy}{dy'} d \frac{dz}{dz'}. \quad (400)$$

Again, by (388),

\* The effect of distorting stresses in a solid on the phenomena of crystallization and liquefaction, as well as the effect of change of hydrostatic pressure common to the solid and liquid, was first described by Professor James Thomson. See *Trans. R. S. Edin.*, vol. xvi, p. 575: and *Proc. Roy. Soc.*, vol. xi, p. 473, or *Phil. Mag.*, s. 4, vol. xxiv, p. 395.

$$d\varepsilon = t d\eta + \eta dt - p dv - v dp + m d\mu_1. \quad (401)$$

Now the suppositions which have been made require that

$$v = \frac{dx}{dx'} \frac{dy}{dy'} \frac{dz}{dz'}, \quad (402)$$

and

$$dv = \frac{dy}{dy'} \frac{dz}{dz'} d \frac{dx}{dx'} + \frac{dz}{dz'} \frac{dx}{dx'} d \frac{dy}{dy'} + \frac{dx}{dx'} \frac{dy}{dy'} d \frac{dz}{dz'}. \quad (403)$$

Combining equations (400), (401), and (403), and observing that  $\varepsilon_v$  and  $\eta_v$  are equivalent to  $\varepsilon$  and  $\eta$ , we obtain

$$\begin{aligned} & \eta dt - v dp + m d\mu_1 \\ &= \left( X_x + p \frac{dy}{dy'} \frac{dz}{dz'} \right) d \frac{dx}{dx'} + X_y d \frac{dx}{dy'} + \left( Y_y + p \frac{dz}{dz'} \frac{dx}{dx'} \right) d \frac{dy}{dy'}. \end{aligned} \quad (404)$$

The reader will observe that when the solid is subjected on all sides to the uniform normal pressure  $p$ , the coefficients of the differentials in the second member of this equation will vanish. For the expression  $\frac{dy}{dy'} \frac{dz}{dz'}$  represents the projection on the  $Y$ - $Z$  plane of a side of the parallelopiped for which  $x'$  is constant, and multiplied by  $p$  it will be equal to the component parallel to the axis of  $X$  of the total pressure across this side, i. e., it will be equal to  $X_x$ , taken negatively.

The case is similar with respect to the coefficient of  $d \frac{dy}{dy'}$ ; and  $X_y$  evidently denotes a force tangential to the surface on which it acts. It will also be observed, that if we regard the forces acting upon the sides of the solid parallelopiped as composed of the hydrostatic pressure  $p$  together with addition forces, the work done in any infinitesimal variation of the state of strain of the solid by these additional forces will be represented by the second member of the equation.

We will first consider the case in which the fluid is identical in substance with the solid. We have then, by equation (97), for a mass of the fluid equal to that of the solid,

$$\eta_F dt - v_F dp + m d\mu_1 = 0, \quad (405)$$

$\eta_F$  and  $v_F$  denoting the entropy and volume of the fluid. By subtraction we obtain

$$\begin{aligned} & -(\eta_F - \eta) dt + (v_F - v) dp \\ &= \left( X_x + p \frac{dy}{dy'} \frac{dz}{dz'} \right) d \frac{dx}{dx'} + X_y d \frac{dx}{dy'} + \left( Y_y + p \frac{dz}{dz'} \frac{dx}{dx'} \right) d \frac{dy}{dy'}. \end{aligned} \quad (406)$$

Now if the quantities  $\frac{dx}{dx'}, \frac{dy}{dy'}, \frac{dz}{dz'}$  remain constant, we shall have for the relation between the variations of temperature and pressure which is necessary for the preservation of equilibrium

$$\frac{dt}{dp} = \frac{v_F - v}{\eta_F - \eta} = t \frac{v_F - v}{Q}, \quad (407)$$

where  $Q$  denotes the heat which would be absorbed if the solid body should pass into the fluid state without change of temperature or pressure. This equation is similar to (131), which applies to bodies subject to hydrostatic pressure. But the value of  $\frac{dt}{dp}$  will not generally be the same as if the solid were subject on all sides to the uniform normal pressure  $p$ ; for the quantities  $v$  and  $\eta$  (and therefore  $Q$ ) will in general have different values. But when the pressures on all sides are normal and equal, the value of  $\frac{dt}{dp}$  will be the same, whether we consider the pressure when varied as still normal and equal on all sides, or consider the quantities  $\frac{dx}{dx'}, \frac{dx}{dy'}, \frac{dy}{dy'}$  as constant.

But if we wish to know how the temperature is affected if the pressure between the solid and fluid remains constant, but the strain of the solid is varied in any way consistent with this supposition, the differential coefficients of  $t$  with respect to the quantities which express the strain are indicated by equation (406). These differential coefficients all vanish, when the pressures on all sides are normal and equal, but the differential coefficient  $\frac{dt}{dp}$ , when  $\frac{dx}{dx'}, \frac{dx}{dy'}, \frac{dy}{dy'}$  are constant, or when the pressures on all sides are normal and equal, vanishes only when the density of the fluid is equal to that of the solid.

The case is nearly the same when the fluid is not identical in substance with the solid, if we suppose the composition of the fluid to remain unchanged. We have necessarily with respect to the fluid

$$d\mu_1 = \left( \frac{d\mu_1}{dt} \right)_{p, m}^{(F)} dt + \left( \frac{d\mu_1}{dp} \right)_{t, m}^{(F)} dp,^* \quad (408)$$

where the index (F) is used to indicate that the expression to which it is affixed relates to the fluid. But by equation (92)

$$\left( \frac{d\mu_1}{dt} \right)_{p, m}^{(F)} = - \left( \frac{d\eta}{dm_1} \right)_{t, p, m}^{(F)}, \text{ and } \left( \frac{d\mu_1}{dp} \right)_{t, m}^{(F)} = \left( \frac{dv}{dm_1} \right)_{t, p, m}^{(F)}. \quad (409)$$

Substituting these values in the preceding equation, transposing terms, and multiplying by  $m$ , we obtain

\* A suffixed  $m$  stands here, as elsewhere in this paper, for all the symbols  $m_1, m_2$ , etc., except such as may occur in the differential coefficient.

$$m \left( \frac{d\eta}{dm_1} \right)_{t, p, m}^{(F)} dt - m \left( \frac{dv}{dm_1} \right)_{t, p, m}^{(F)} dp + m d\mu_1 = 0. \quad (410)$$

By subtracting this equation from (404) we may obtain an equation similar to (406), except that in place of  $\eta_F$  and  $v_F$  we shall have the expressions

$$m \left( \frac{d\eta}{dm_1} \right)_{t, p, m}^{(F)} \quad \text{and} \quad m \left( \frac{dv}{dm_1} \right)_{t, p, m}^{(F)}$$

The discussion of equation (406) will therefore apply *mutatis mutandis* to this case.

We may also wish to find the variations in the composition of the fluid which will be necessary for equilibrium when the pressure  $p$  or the quantities  $\frac{dx}{dx'}, \frac{dy}{dy'}, \frac{dz}{dz'}$  are varied, the temperature remaining constant. If we know the value for the fluid of the quantity represented by  $\zeta$  on page 142 in terms of  $t, p$ , and the quantities of the several components  $m_1, m_2, m_3$ , etc., the first of which relates to the substance of which the solid is formed, we can easily find the value of  $\mu_1$  in terms of the same variables. Now in considering variations in the composition of the solid, it will be sufficient if we make all but one of the components variable. We may therefore give to  $m_1$  a constant value, and making  $t$  also constant, we shall have

$$d\mu_1 = \left( \frac{d\mu_1}{dp} \right)_{t, m}^{(F)} dp + \left( \frac{d\mu_1}{dm_2} \right)_{t, p, m}^{(F)} dm_2 + \left( \frac{d\mu_1}{dm_3} \right)_{t, p, m}^{(F)} dm_3 + \text{etc.}$$

Substituting this value in equation (404), and cancelling the term containing  $dt$ , we obtain

$$\begin{aligned} & \left\{ m \left( \frac{d\mu_1}{dp} \right)_{t, m}^{(F)} - v \right\} dp + m \left( \frac{d\mu_1}{dm_2} \right)_{t, p, m}^{(F)} dm_2 \\ & + m \left( \frac{d\mu_1}{dm_3} \right)_{t, p, m}^{(F)} dm_3 + \text{etc.} = \left( X_{x'} + p \frac{dy}{dy'} \frac{dz}{dz'} \right) d \frac{dx}{dx'} \\ & + X_{y'} d \frac{dx}{dy'} + \left( Y_{y'} + p \frac{dz}{dz'} \frac{dx}{dx'} \right) d \frac{dy}{dy'} \end{aligned} \quad (411)$$

This equation shows the variation in the quantity of any one of the components of the fluid (other than the substance which forms the solid) which will balance a variation of  $p$ , or of  $\frac{dx}{dx'}, \frac{dy}{dy'}, \frac{dz}{dz'}$  with respect to the tendency of the solid to dissolve.

*Fundamental Equations for Solids.*

The principles developed in the preceding pages show that the solution of problems relating to the equilibrium of a solid, or at least their reduction to purely analytical processes, may be made to depend upon our knowledge of the composition and density of the solid at every point in some particular state, which we have called the state of reference, and of the relation existing between the quantities which have been represented by  $\varepsilon_{vr}$ ,  $\eta_{vr}$ ,  $\frac{dx}{dx'}$ ,  $\frac{dy}{dy'}$ ,  $\dots$ ,  $\frac{dz}{dz'}$ ,  $x'$ ,  $y'$ , and  $z'$ . When the solid is in contact with fluids, a certain knowledge of the properties of the fluids is also requisite, but only such as is necessary for the solution of problems relating to the equilibrium of fluids among themselves.

If in any state of which a solid is capable, it is homogeneous in its nature and in its state of strain, we may choose this state as the state of reference, and the relation between  $\varepsilon_{vr}$ ,  $\eta_{vr}$ ,  $\frac{dx}{dx'}$ ,  $\dots$ ,  $\frac{dz}{dz'}$ , will be independent of  $x'$ ,  $y'$ ,  $z'$ . But it is not always possible, even in the case of bodies which are homogeneous in nature, to bring all the elements simultaneously into the same state of strain. It would not be possible, for example, in the case of a Prince Rupert's drop.

If, however, we know the relation between  $\varepsilon_{vr}$ ,  $\eta_{vr}$ ,  $\frac{dx}{dx'}$ ,  $\dots$ ,  $\frac{dz}{dz'}$ , for any kind of homogeneous solid, with respect to any given state of reference, we may derive from it a similar relation with respect to any other state as a state of reference. For if  $x'$ ,  $y'$ ,  $z'$  denote the co-ordinates of points of the solid in the first state of reference, and  $x''$ ,  $y''$ ,  $z''$  the co-ordinates of the same points in the second state of reference, we shall have necessarily

$$\frac{dx}{dx'} = \frac{dx}{dx''} \frac{dx''}{dx'} + \frac{dx}{dy''} \frac{dy''}{dx'} + \frac{dx}{dz''} \frac{dz''}{dx'}, \text{ etc. (nine equations), } (412)$$

and if we write  $R$  for the volume of an element in the state  $(x'', y'', z'')$  divided by its volume in the state  $(x', y', z')$ , we shall have

$$R = \begin{vmatrix} \frac{dx''}{dx'} & \frac{dx''}{dy'} & \frac{dx''}{dz'} \\ \frac{dy''}{dx'} & \frac{dy''}{dy'} & \frac{dy''}{dz'} \\ \frac{dz''}{dx'} & \frac{dz''}{dy'} & \frac{dz''}{dz'} \end{vmatrix}, \quad (413)$$

$$\varepsilon_{vn} = R \varepsilon_{v'n}, \quad \eta_{vn} = R \eta_{v'n}. \quad (414)$$

If, then, we have ascertained by experiment the value of  $\varepsilon_v$ , in terms of  $\eta_{vi}$ ,  $\frac{dx}{dx'}$ ,  $\dots$   $\frac{dz}{dz'}$ , and the quantities which express the composition of the body, by the substitution of the values given in (412)–(414), we shall obtain  $\varepsilon_{vn}$  in terms of  $\eta_{vn}$ ,  $\frac{dx}{dx''}$ ,  $\dots$   $\frac{dz''}{dz''}$ ,  $\frac{dx''}{dx'}$ ,  $\dots$   $\frac{dz'}{dz'}$ , and the quantities which express the composition of the body.

We may apply this to the elements of a body which may be variable from point to point in composition and state of strain in a given state of reference ( $x'', y'', z''$ ), and if the body is fully described in that state of reference, both in respect to its composition and to the displacement which it would be necessary to give to a homogeneous solid of the same composition, for which  $\varepsilon_v$  is known in terms of  $\eta_{vi}$ ,  $\frac{dx}{dx'}$ ,  $\dots$   $\frac{dz}{dz'}$ , and the quantities which express its composition, to bring it from the state of reference ( $x', y', z'$ ) into a similar and similarly situated state of strain with that of the element of the non-homogeneous body, we may evidently regard  $\frac{dx''}{dx'}$ ,  $\dots$   $\frac{dz''}{dz'}$  as known for each element of the body, that is, as known in terms of  $x'', y'', z''$ . We shall then have  $\varepsilon_{vn}$  in terms of  $\eta_{vn}$ ,  $\frac{dx}{dx''}$ ,  $\dots$   $\frac{dz}{dz''}$ ,  $x'', y'', z''$ ; and since the composition of the body is known in terms of  $x'', y'', z''$ , and the density, if not given directly, can be determined from the density of the homogeneous body in its state of reference ( $x', y', z'$ ), this is sufficient for determining the equilibrium of any given state of the non-homogeneous solid.

An equation, therefore, which expresses for any kind of solid, and with reference to any determined state of reference, the relation between the quantities denoted by  $\varepsilon_{vi}$ ,  $\eta_{vi}$ ,  $\frac{dx}{dx'}$ ,  $\dots$   $\frac{dz}{dz'}$ , involving also the quantities which express the composition of the body, when that is capable of continuous variation, or any other equation from which the same relations may be deduced, may be called a *fundamental equation* for that kind of solid. It will be observed that the sense in which this term is here used, is entirely analogous to that in which we have already applied the term to fluids and solids which are subject only to hydrostatic pressure.

When the fundamental equation between  $\varepsilon_{vi}$ ,  $\eta_{vi}$ ,  $\frac{dx}{dx'}$ ,  $\dots$   $\frac{dz}{dz'}$  is

known, we may obtain by differentiation the values of  $t$ ,  $X_x$ , . . .  $Z_z$ , in terms of the former quantities, which will give eleven independent relations between the twenty-one quantities

$$\varepsilon_v, \eta_v, \frac{dx}{dx'}, \dots \frac{dz}{dz'}, t, X_x, \dots Z_z, \quad (415)$$

which are all that exist, since ten of these quantities are independent. All these equations may also involve variables which express the composition of the body, when that is capable of continuous variation.

If we use the symbol  $\psi_v$ , to denote the value of  $\psi$  (as defined on pages 144, 145) for any element of a solid divided by the volume of the element in the state of reference, we shall have

$$\psi_v = \varepsilon_v - t \eta_v. \quad (416)$$

The equation (356) may be reduced to the form

$$\delta\psi_v = -\eta_v \delta t + \Sigma \Sigma' \left( X_x \delta \frac{dx}{dx'} \right). \quad (417)$$

Therefore, if we know the value of  $\psi_v$ , in terms of the variables  $t$ ,  $\frac{dx}{dx'}$ , . . .  $\frac{dz}{dz'}$ , together with those which express the composition of the body, we may obtain by differentiation the values of  $\eta_v$ ,  $X_x$ , . . .  $Z_z$ , in terms of the same variables. This will make eleven independent relations between the same quantities as before, except that we shall have  $\psi_v$ , instead of  $\varepsilon_v$ . Or if we eliminate  $\psi_v$ , by means of equation (416), we shall obtain eleven independent equations between the quantities in (415) and those which express the composition of the body. An equation, therefore, which determines the value of  $\psi_v$ , as a function of the quantities  $t$ ,  $\frac{dx}{dx'}$ , . . .  $\frac{dz}{dz'}$ , and the quantities which express the composition of the body when it is capable of continuous variation, is a fundamental equation for the kind of solid to which it relates.

In the discussion of the conditions of equilibrium of a solid, we might have started with the principle that it is necessary and sufficient for equilibrium that the temperature shall be uniform throughout the whole mass in question, and that the variation of the force-function ( $-\psi$ ) of the same mass shall be null or negative for any variation in the state of the mass not affecting its temperature. We might have assumed that the value of  $\psi$  for any same element of the solid is a

function of the temperature and the state of strain, so that for constant temperature we might write

$$\delta\psi_v = \Sigma \Sigma' \left( X_{x_i} \delta \frac{dx}{dx'} \right),$$

the quantities  $X_{x_1}, \dots, X_{x_n}$  being defined by this equation. This would be only a formal change in the definition of  $X_{x_1}, \dots, X_{x_n}$ , and would not affect their values, for this equation holds true of  $X_{x_1}, \dots, X_{x_n}$  as defined by equation (355). With such data, by transformations similar to those which we have employed, we might obtain similar results.\* It is evident that the only difference in the equations would be that  $\psi_v$  would take the place of  $\varepsilon_v$ , and that the terms relating to entropy would be wanting. Such a method is evidently preferable with respect to the directness with which the results are obtained. The method of this paper shows more distinctly the rôle of *energy* and *entropy* in the theory of equilibrium, and can be extended more naturally to those dynamical problems in which motions take place under the condition of constancy of entropy of the elements of a solid (as when vibrations are propagated through a solid), just as the other method can be more naturally extended to dynamical problems in which the temperature is constant. (See note on page 145.)

We have already had occasion to remark that the state of strain of any element considered without reference to directions in space is capable of only six independent variations. Hence, it must be possible to express the state of strain of an element by six functions of  $\frac{dx}{dx'}, \dots, \frac{dz}{dz'}$ , which are independent of the position of the element.

For these quantities we may choose the squares of the ratios of elongation of lines parallel to the three co-ordinate axes in the state of reference, and the products of the ratios of elongation for each pair of these lines multiplied by the cosine of the angle which they include in the variable state of the solid. If we denote these quantities by  $A, B, C, a, b, c$ , we shall have

\* For an example of this method, see Thomson and Tait's *Natural Philosophy*, vol. i, p. 705. With regard to the general theory of elastic solids, compare also Thomson's Memoir "On the Thermo-elastic and Thermo-magnetic Properties of Matter" in the *Quarterly Journal of Mathematics*, vol. i, p. 57 (1855), and Green's memoirs on the propagation, reflection, and refraction of light in the *Transactions of the Cambridge Philosophical Society*, vol. vii.

$$A = \Sigma \left( \frac{dx}{dx'} \right)^2, \quad B = \Sigma \left( \frac{dx}{dy'} \right)^2, \quad C = \Sigma \left( \frac{dx}{dz'} \right)^2, \quad (418)$$

$$a = \Sigma \left( \frac{dx}{dy'} \frac{dx}{dz'} \right), \quad b = \Sigma \left( \frac{dx}{dz'} \frac{dx}{dx'} \right), \quad c = \Sigma \left( \frac{dx}{dx'} \frac{dx}{dy'} \right). \quad (419)$$

The determination of the fundamental equation for a solid is thus reduced to the determination of the relation between  $\varepsilon_{v\prime}$ ,  $\eta_{v\prime}$ ,  $A$ ,  $B$ ,  $C$ ,  $a$ ,  $b$ ,  $c$ , or of the relation between  $\psi_{v\prime}$ ,  $t$ ,  $A$ ,  $B$ ,  $C$ ,  $a$ ,  $b$ ,  $c$ .

In the case of isotropic solids, the state of strain of an element, so far as it can affect the relation of  $\varepsilon_{v\prime}$  and  $\eta_{v\prime}$ , or of  $\psi_{v\prime}$  and  $t$ , is capable of only three independent variations. This appears most distinctly as a consequence of the proposition that for any given strain of an element there are three lines in the element which are at right angles to one another both in its unstrained and in its strained state. If the unstrained element is isotropic, the ratios of elongation for these three lines must with  $\eta_{v\prime}$  determine the value  $\varepsilon_{v\prime}$ , or with  $t$  determine the value of  $\psi_{v\prime}$ .

To demonstrate the existence of such lines, which are called the *principal axes of strain*, and to find the relations of the elongations of such lines to the quantities  $\frac{dx}{dx'}$ ,  $\dots$ ,  $\frac{dz}{dz'}$ , we may proceed as follows. The ratio of elongation  $r$  of any line of which  $\alpha'$ ,  $\beta'$ ,  $\gamma'$  are the direction-cosines in the state of reference is evidently given by the equation

$$\begin{aligned} r^2 &= \left( \frac{dx}{dx'} \alpha' + \frac{dx}{dy'} \beta' + \frac{dx}{dz'} \gamma' \right)^2 \\ &\quad + \left( \frac{dy}{dx'} \alpha' + \frac{dy}{dy'} \beta' + \frac{dy}{dz'} \gamma' \right)^2 \\ &\quad + \left( \frac{dz}{dx'} \alpha' + \frac{dz}{dy'} \beta' + \frac{dz}{dz'} \gamma' \right)^2. \end{aligned} \quad (420)$$

Now the proposition to be established is evidently equivalent to this—that it is always possible to give such directions to the two systems of rectangular axes  $X'$ ,  $Y'$ ,  $Z'$ , and  $X$ ,  $Y$ ,  $Z$ , that

$$\left. \begin{aligned} \frac{dx}{dy'} &= 0, & \frac{dx}{dz'} &= 0, & \frac{dy}{dz'} &= 0, \\ \frac{dy}{dx'} &= 0, & \frac{dz}{dx'} &= 0, & \frac{dz}{dy'} &= 0. \end{aligned} \right\} \quad (421)$$

We may choose a line in the element for which the value of  $r$  is at least as great as for any other, and make the axes of  $X$  and  $X'$  parallel to this line in the strained and unstrained states respectively.

Then

$$\frac{dy}{dx'} = 0, \quad \frac{dz}{dx'} = 0. \quad (422)$$

Moreover, if we write  $\frac{d(r^2)}{d\alpha'}, \frac{d(r^2)}{d\beta'}, \frac{d(r^2)}{d\gamma'}$  for the differential coefficients obtained from (420) by treating  $\alpha', \beta', \gamma'$  as *independent* variables,

$$\frac{d(r^2)}{d\alpha'} d\alpha' + \frac{d(r^2)}{d\beta'} d\beta' + \frac{d(r^2)}{d\gamma'} d\gamma' = 0,$$

when

$$\alpha' d\alpha' + \beta' d\beta' + \gamma' d\gamma' = 0,$$

and

$$\alpha' = 1, \quad \beta' = 0, \quad \gamma' = 0.$$

That is,

$$\frac{d(r^2)}{d\beta'} = 0, \quad \text{and} \quad \frac{d(r^2)}{d\gamma'} = 0,$$

when

$$\alpha' = 1, \quad \beta' = 0, \quad \gamma' = 0.$$

Hence,

$$\frac{dx}{dy'} = 0, \quad \frac{dx}{dz'} = 0. \quad (423)$$

Therefore a line of the element which in the unstrained state is perpendicular to  $X'$  is perpendicular to  $X$  in the strained state. Of all such lines we may choose one for which the value of  $r$  is at least as great as for any other, and make the axes of  $X'$  and  $X$  parallel to this line in the unstrained and in the strained state respectively.

Then

$$\frac{dz}{dy'} = 0; \quad (424)$$

and it may easily be shown by reasoning similar to that which has just been employed that

$$\frac{dy}{dz'} = 0. \quad (425)$$

Lines parallel to the axes of  $X'$ ,  $Y'$ , and  $Z'$  in the unstrained body will therefore be parallel to  $X$ ,  $Y$ , and  $Z$  in the strained body, and the ratios of elongation for such lines will be

$$\frac{dx}{dx'}, \quad \frac{dy}{dy'}, \quad \frac{dz}{dz'}.$$

These lines have the common property of a stationary value of the ratio of elongation for varying directions of the line. This appears from the form to which the general value of  $r^2$  is reduced by the positions of the co-ordinate axes, viz.,

$$r^2 = \left(\frac{dx}{dx'}\right)^2 \alpha'^2 + \left(\frac{dy}{dy'}\right)^2 \beta'^2 + \left(\frac{dz}{dz'}\right)^2 \gamma'^2.$$

Having thus proved the existence of lines, with reference to any particular strain, which have the properties mentioned, let us proceed to find the relations between the ratios of elongation for these lines (the *principal axes of strain*) and the quantities  $\frac{dx}{dx'}, \dots, \frac{dz}{dz'}$  under the most general supposition with respect to the position of the co-ordinate axes.

For any principal axis of strain we have

$$\frac{d(r^2)}{d\alpha'} d\alpha' + \frac{d(r^2)}{d\beta'} d\beta' + \frac{d(r^2)}{d\gamma'} d\gamma' = 0,$$

when

$$\alpha' d\alpha' + \beta' d\beta' + \gamma' d\gamma' = 0,$$

the differential coefficients in the first of these equations being determined from (420) as before. Therefore,

$$\frac{1}{\alpha'} \frac{d(r^2)}{d\alpha'} = \frac{1}{\beta'} \frac{d(r^2)}{d\beta'} = \frac{1}{\gamma'} \frac{d(r^2)}{d\gamma'}. \quad (426)$$

From (420) we obtain directly

$$\frac{\alpha'}{2} \frac{d(r^2)}{d\alpha'} + \frac{\beta'}{2} \frac{d(r^2)}{d\beta'} + \frac{\gamma'}{2} \frac{d(r^2)}{d\gamma'} = r^2. \quad (427)$$

From the two last equations, in virtue of the necessary relation  $\alpha'^2 + \beta'^2 + \gamma'^2 = 1$ , we obtain

$$\frac{1}{2} \frac{d(r^2)}{d\alpha'} = \alpha' r^2, \quad \frac{1}{2} \frac{d(r^2)}{d\beta'} = \beta' r^2, \quad \frac{1}{2} \frac{d(r^2)}{d\gamma'} = \gamma' r^2, \quad (428)$$

or, if we substitute the values of the differential coefficients taken from (420),

$$\left. \begin{aligned} \alpha' \Sigma \left( \frac{dx}{dx'} \right)^2 + \beta' \Sigma \left( \frac{dx}{dx'} \frac{dx}{dy'} \right) + \gamma' \Sigma \left( \frac{dx}{dx'} \frac{dx}{dz'} \right) &= \alpha' r^2, \\ \alpha' \Sigma \left( \frac{dx}{dy'} \frac{dx}{dx'} \right) + \beta' \Sigma \left( \frac{dx}{dy'} \right)^2 + \gamma' \Sigma \left( \frac{dx}{dy'} \frac{dx}{dz'} \right) &= \beta' r^2, \\ \alpha' \Sigma \left( \frac{dx}{dz'} \frac{dx}{dx'} \right) + \beta' \Sigma \left( \frac{dx}{dz'} \frac{dx}{dy'} \right) + \gamma' \Sigma \left( \frac{dx}{dz'} \right)^2 &= \gamma' r^2. \end{aligned} \right\} \quad (429)$$

If we eliminate  $\alpha'$ ,  $\beta'$ ,  $\gamma'$  from these equations, we may write the result in the form,

$$\left| \begin{array}{ccc} \Sigma \left( \frac{dx}{dx'} \right)^2 - r^2 & \Sigma \left( \frac{dx}{dx'} \frac{dx}{dy'} \right) & \Sigma \left( \frac{dx}{dx'} \frac{dx}{dz'} \right) \\ \Sigma \left( \frac{dx}{dy'} \frac{dx}{dx'} \right) & \Sigma \left( \frac{dx}{dy'} \right)^2 - r^2 & \Sigma \left( \frac{dx}{dy'} \frac{dx}{dz'} \right) \\ \Sigma \left( \frac{dx}{dz'} \frac{dx}{dx'} \right) & \Sigma \left( \frac{dx}{dz'} \frac{dx}{dy'} \right) & \Sigma \left( \frac{dx}{dz'} \right)^2 - r^2 \end{array} \right| = 0. \quad (430)$$

We may write

$$-r^6 + Er^4 - Fr^2 + G = 0. \quad (431)$$

Then

$$E = \Sigma' \Sigma \left( \frac{dx}{dx'} \right)^2. \quad (432)$$

Also\*

$$\begin{aligned} F &= \Sigma' \left\{ \Sigma \left( \frac{dx}{dx'} \right)^2 \Sigma \left( \frac{dx}{dy'} \right)^2 - \Sigma \left( \frac{dx}{dx'} \frac{dx}{dy'} \right) \Sigma \left( \frac{dx}{dx'} \frac{dx}{dy'} \right) \right\} \\ &= \Sigma' \Sigma \left\{ \left( \frac{dx}{dx'} \right)^2 \Sigma \left( \frac{dx}{dy'} \right)^2 - \frac{dx}{dx'} \frac{dx}{dy'} \Sigma \left( \frac{dx}{dx'} \frac{dx}{dy'} \right) \right\} = \\ &= \Sigma' \Sigma \left[ \left( \frac{dx}{dx'} \right)^2 \left( \frac{dy}{dy'} \right)^2 + \left( \frac{dx}{dx'} \right)^2 \left( \frac{dz}{dy'} \right)^2 - \frac{dx}{dx'} \frac{dx}{dy'} \frac{dy}{dx'} \frac{dy}{dy'} - \frac{dx}{dx'} \frac{dx}{dy'} \frac{dz}{dx'} \frac{dz}{dy'} \right] \\ &= \Sigma' \Sigma \left\{ \left( \frac{dx}{dx'} \right)^2 \left( \frac{dy}{dy'} \right)^2 + \left( \frac{dy}{dx'} \right)^2 \left( \frac{dx}{dy'} \right)^2 - 2 \frac{dx}{dx'} \frac{dx}{dy'} \frac{dy}{dx'} \frac{dy}{dy'} \right\} \\ &= \Sigma' \Sigma \left( \frac{dx}{dx'} \frac{dy}{dy'} - \frac{dy}{dx'} \frac{dx}{dy'} \right)^2. \end{aligned} \quad (433)$$

This may also be written

$$F = \Sigma' \Sigma \begin{vmatrix} \frac{dx}{dx'} & \frac{dx}{dy'} \\ \frac{dy}{dx'} & \frac{dy}{dy'} \end{vmatrix}^2. \quad (434)$$

In the reduction of the value of  $G$ , it will be convenient to use the symbol  $\Sigma_{3+3}$  to denote the sum of the six terms formed by changing  $x, y, z$ , into  $y, z, x$ ;  $z, x, y$ ;  $x, z, y$ ;  $y, x, z$ ; and  $z, y, x$ ; and the symbol  $\Sigma_{3-3}$  in the same sense except that the last three terms are to be taken negatively; also to use  $\Sigma'_{3-3}$  in a similar sense with respect to  $x', y', z'$ ; and to use  $x', y', z'$  as equivalent to  $x, y, z$ , except that they are not to be affected by the sign of summation. With this understanding we may write

$$G = \Sigma'_{3-3} \left\{ \Sigma \left( \frac{dx}{dx'} \frac{dx}{dx'} \right) \Sigma \left( \frac{dx}{dy'} \frac{dx}{dy'} \right) \Sigma \left( \frac{dx}{dz'} \frac{dx}{dz'} \right) \right\}. \quad (435)$$

\* The values of  $F$  and  $G$  given in equations (434) and (438), which are here deduced at length, may be derived from inspection of equation (430) by means of the usual theorems relating to the multiplication of determinants. See Salmon's *Lessons Introductory to the Modern Higher Algebra*, 2d Ed., Lesson III; or Baltzer's *Theorie und Anwendung der Determinanten*, § 5.

In expanding the product of the three sums, we may cancel on account of the sign  $\Sigma'$  the terms which do not contain all the three expressions  $dx$ ,  $dy$ , and  $dz$ . Hence we may write

$$\begin{aligned} G &= \sum'_{3-3} \sum_{3+3} \left( \frac{dx}{dx'} \frac{dx}{dy'} \frac{dy}{dy'} \frac{dy}{dz'} \frac{dz}{dz'} \frac{dz}{dx'} \right) \\ &= \sum_{3+3} \left\{ \frac{dx}{dx'} \frac{dy}{dy'} \frac{dz}{dz'} \sum'_{3-3} \left( \frac{dx}{dx'} \frac{dy}{dy'} \frac{dz}{dz'} \right) \right\} \\ &= \sum_{3-3} \left( \frac{dx}{dx'} \frac{dy}{dy'} \frac{dz}{dz'} \right) \sum'_{3-3} \left( \frac{dx}{dx'} \frac{dy}{dy'} \frac{dz}{dz'} \right). \end{aligned} \quad (436)$$

Or, if we set

$$H = \begin{vmatrix} dx & dx & dx \\ \frac{dx}{dx'} & \frac{dy}{dy'} & \frac{dz}{dz'} \\ dy & dy & dy \\ \frac{dy}{dx'} & \frac{dy}{dy'} & \frac{dy}{dz'} \\ dz & dz & dz \\ \frac{dz}{dx'} & \frac{dy}{dy'} & \frac{dz}{dz'} \end{vmatrix}, \quad (437)$$

we shall have

$$G = H^2. \quad (438)$$

It will be observed that  $F$  represents the sum of the squares of the nine minors which can be formed from the determinant in (437), and that  $E$  represents the sum of the squares of the nine constituents of the same determinant.

Now we know by the theory of equations that equation (431) will be satisfied in general by three different values of  $r^2$ , which we may denote by  $r_1^2$ ,  $r_2^2$ ,  $r_3^2$ , and which must represent the squares of the ratios of elongation for the three principal axes of strain; also that  $E$ ,  $F$ ,  $G$ , are symmetrical functions of  $r_1^2$ ,  $r_2^2$ ,  $r_3^2$ , viz.,

$$\begin{aligned} E &= r_1^2 + r_2^2 + r_3^2, \quad F = r_1^2 r_2^2 + r_2^2 r_3^2 + r_3^2 r_1^2, \\ G &= r_1^2 r_2^2 r_3^2. \end{aligned} \quad \{ \quad (439)$$

Hence, although it is possible to solve equation (431) by the use of trigonometrical functions, it will be more simple to regard  $\varepsilon_v$  as a function of  $\eta_v$ , and the quantities  $E$ ,  $F$ ,  $G$  (or  $H$ ), which we have expressed in terms of  $\frac{dx}{dx'}, \dots, \frac{dz}{dz'}$ . Since  $\varepsilon_v$  is a single-valued function of  $\eta_v$ , and  $r_1^2$ ,  $r_2^2$ ,  $r_3^2$  (with respect to all the changes of which the body is capable), and a symmetrical function with respect to  $r_1^2$ ,  $r_2^2$ ,  $r_3^2$ , and since  $r_1^2$ ,  $r_2^2$ ,  $r_3^2$  are *collectively* determined without ambiguity by the values of  $E$ ,  $F$ , and  $H$ , the quantity  $\varepsilon_v$  must be a

single-valued function of  $\eta_v$ ,  $E$ ,  $F$ , and  $H$ . The determination of the fundamental equation for isotropic bodies is therefore reduced to the determination of this function, or (as appears from similar considerations) the determination of  $\psi_v$  as a function of  $t$ ,  $E$ ,  $F$ , and  $H$ .

It appears from equations (439) that  $E$  represents the sum of the squares of the ratios of elongation for the principal axes of strain, that  $F$  represents the sum of the squares of the ratios of enlargement for the three surfaces determined by these axes, and that  $G$  represents the square of the ratio of enlargement of volume. Again, equation (432) shows that  $E$  represents the sum of the squares of the ratios of elongation for lines parallel to  $X'$ ,  $Y'$ , and  $Z'$ ; equation (434) shows that  $F$  represents the sum of the squares of the ratios of enlargement for surfaces parallel to the planes  $X'-Y'$ ,  $Y'-Z'$ ,  $Z'-X'$ ; and equation (438), like (439), shows that  $G$  represents the square of the ratio of enlargement of volume. Since the position of the co-ordinate axes is arbitrary, it follows that the sum of the squares of the ratios of elongation or enlargement of three lines or surfaces which in the unstrained state are at right angles to one another, is otherwise independent of the direction of the lines or surfaces. Hence,  $\frac{1}{3}E$  and  $\frac{1}{3}F$  are the mean squares of the ratios of linear elongation and of superficial enlargement, for all possible directions in the unstrained solid.

There is not only a practical advantage in regarding the strain as determined by  $E$ ,  $F$ , and  $H$ , instead of  $E$ ,  $F$ , and  $G$ , because  $H$  is more simply expressed in terms of  $\frac{dx}{dx'}$ ,  $\dots$ ,  $\frac{dz}{dz'}$ , but there is also a certain theoretical advantage on the side of  $E$ ,  $F$ ,  $H$ . If the systems of co-ordinate axes  $X$ ,  $Y$ ,  $Z$ , and  $X'$ ,  $Y'$ ,  $Z'$ , are either identical or such as are capable of superposition, which it will always be convenient to suppose, the determinant  $H$  will always have a positive value for any strain of which a body can be capable. But it is possible to give to  $x$ ,  $y$ ,  $z$  such values as functions of  $x'$ ,  $y'$ ,  $z'$  that  $H$  shall have a negative value. For example, we may make

$$x = x', \quad y = y', \quad z = -z'. \quad (440)$$

This will give  $H = -1$ , while

$$x = x', \quad y = y', \quad z = z' \quad (441)$$

will give  $H = 1$ . Both (440) and (441) give  $G = 1$ . Now although such a change in the position of the particles of a body as is represented by (440) cannot take place while the body remains solid, yet

a method of representing strains may be considered incomplete, which confuses the cases represented by (440) and (441).

We may avoid all such confusion by using  $E$ ,  $F$ , and  $H$  to represent a strain. Let us consider an element of the body strained which in the state  $(x', y', z')$  is a cube with its edges parallel to the axes of  $X'$ ,  $Y'$ ,  $Z'$ , and call the edges  $dx'$ ,  $dy'$ ,  $dz$  according to the axes to which they are parallel, and consider the ends of the edges as positive for which the values of  $x'$ ,  $y'$ , or  $z'$  are the greater. Whatever may be the nature of the parallelopiped in the state  $(x, y, z)$  which corresponds to the cube  $dx', dy', dz'$  and is determined by the quantities  $\frac{dx}{dx'}$ ,  $\dots \frac{dz}{dz'}$ , it may always be brought by continuous changes to the form of a cube and to a position in which the edges  $dx'$ ,  $dy'$  shall be parallel to the axes of  $X$  and  $Y$ , the positive ends of the edges toward the positive directions of the axes, and this may be done without giving the volume of the parallelopiped the value zero, and therefore without changing the sign of  $H$ . Now two cases are possible;—the positive end of the edge  $dz'$  may be turned toward the positive or toward the negative direction of the axis of  $Z$ . In the first case,  $H$  is evidently positive; in the second, negative. The determinant  $H$  will therefore be positive or negative,—we may say, if we choose, that the volume will be positive or negative,—according as the element can or cannot be brought from the state  $(x, y, z)$  to the state  $(x', y', z')$  by continuous changes without giving its volume the value zero.

If we now recur to the consideration of the principal axes of strain and the principal ratios of elongation  $r_1$ ,  $r_2$ ,  $r_3$ , and denote by  $U_1$ ,  $U_2$ ,  $U_3$  and  $U'_1$ ,  $U'_2$ ,  $U'_3$  the principal axes of strain in the strained and unstrained element respectively, it is evident that the sign of  $r_1$ , for example, depends upon the direction in  $U_1$  which we regard as corresponding to a given direction in  $U'_1$ . If we choose to associate directions in these axes so that  $r_1$ ,  $r_2$ ,  $r_3$  shall all be positive, the positive or negative value of  $H$  will determine whether the system of axes  $U_1$ ,  $U_2$ ,  $U_3$  is or is not capable of superposition upon the system  $U'_1$ ,  $U'_2$ ,  $U'_3$  so that corresponding directions in the axes shall coincide. Or, if we prefer to associate directions in the two systems of axes, so that they shall be capable of superposition, corresponding directions coinciding, the positive or negative value of  $H$  will determine whether an even or an odd number of the quantities  $r_1$ ,  $r_2$ ,  $r_3$  are negative. In this case we may write

$$r_1 r_2 r_3 = H = \begin{vmatrix} \frac{dx}{dx'} & \frac{dx}{dy'} & \frac{dx}{dz'} \\ \frac{dy}{dx'} & \frac{dy}{dy'} & \frac{dy}{dz'} \\ \frac{dz}{dx'} & \frac{dz}{dy'} & \frac{dz}{dz'} \end{vmatrix}. \quad (442)$$

It will be observed that to change the signs of two of the quantities  $r_1, r_2, r_3$  is simply to give a certain rotation to the body without changing its state of strain.

Whichever supposition we make with respect to the axes  $U_1, U_2, U_3$ , it is evident that the state of strain is completely determined by the values  $E, F$ , and  $H$ , not only when we limit ourselves to the consideration of such strains as are consistent with the idea of solidity, but also when we regard any values of  $\frac{dx}{dx'}, \dots, \frac{dz}{dz'}$  as possible.

*Approximative Formulae.*—For many purposes the value of  $\varepsilon_v$ , for an isotropic solid may be represented with sufficient accuracy by the formula

$$\varepsilon_v = i' + e' E + f' F + h' H, \quad (443)$$

where  $i', e', f'$ , and  $h'$  denote functions of  $\eta_v$ ; or the value of  $\psi_v$ , by the formula

$$\psi_v = i + e E + f F + h H, \quad (444)$$

where  $i, e, f$ , and  $h$  denote functions of  $t$ . Let us first consider the second of these formulae. Since  $E, F$ , and  $H$  are symmetrical functions of  $r_1, r_2, r_3$ , if  $\psi_v$  is any function of  $t, E, F, H$ , we must have

$$\left. \begin{aligned} \frac{d\psi_v}{dr_1} &= \frac{d\psi_v}{dr_2} = \frac{d\psi_v}{dr_3}, \\ \frac{d^2\psi_v}{dr_1^2} &= \frac{d^2\psi_v}{dr_2^2} = \frac{d^2\psi_v}{dr_3^2}, \\ \frac{d^2\psi_v}{dr_1 dr_2} &= \frac{d^2\psi_v}{dr_2 dr_3} = \frac{d^2\psi_v}{dr_3 dr_1}, \end{aligned} \right\} \quad (445)$$

whenever  $r_1 = r_2 = r_3$ . Now  $i, e, f$ , and  $h$  may be determined (as functions of  $t$ ) so as to give to

$$\psi_v, \quad \frac{d\psi_v}{dr_1}, \quad \frac{d^2\psi_v}{dr_1^2}, \quad \frac{d^2\psi_v}{dr_1 dr_2}$$

their proper values at every temperature for some isotropic state of strain, which may be determined by any desired condition. We shall suppose that they are determined so as to give the proper

values to  $\psi_{v_i}$ , etc., when the stresses in the solid vanish. If we denote by  $r_0$  the common value of  $r_1, r_2, r_3$  which will make the stresses vanish at any given temperature, and imagine the true value of  $\psi_{v_i}$ , and also the value given by equation (444) to be expressed in terms of the ascending powers of

$$r_1 - r_0, \quad r_2 - r_0, \quad r_3 - r_0, \quad (446)$$

it is evident that the expressions will coincide as far as the terms of the second degree *inclusive*. That is, the errors of the values of  $\psi_{v_i}$  given by equation (444) are of the same order of magnitude as the cubes of the above differences. The errors of the values of

$$\frac{d\psi_{v_i}}{dr_1}, \quad \frac{d\psi_{v_i}}{dr_2}, \quad \frac{d\psi_{v_i}}{dr_3}$$

will be of the same order of magnitude as the squares of the same differences. Therefore, since

$$\frac{d\psi_{v_i}}{dx'} = \frac{d\psi_{v_i}}{dr_1} \frac{dr_1}{dx'} + \frac{d\psi_{v_i}}{dr_2} \frac{dr_2}{dx'} + \frac{d\psi_{v_i}}{dr_3} \frac{dr_3}{dx'} \quad (447)$$

whether we regard the true value of  $\psi_{v_i}$  or the value given by equation (444), and since the error in (444) does not affect the values of

$$\frac{dr_1}{dx'}, \quad \frac{dr_2}{dx'}, \quad \frac{dr_3}{dx'}$$

which we may regard as determined by equations (431), (432), (434), (437) and (438), the errors in the values of  $X_x$ , derived from (444) will be of the same order of magnitude as the squares of the differences in (446). The same will be true with respect to  $X_y$ ,  $X_z$ ,  $Y_x$ , etc., etc.

It will be interesting to see how the quantities  $e$ ,  $f$ , and  $h$  are related to those which most simply represent the elastic properties of isotropic solids. If we denote by  $V$  and  $R$  the *elasticity of volume* and the *rigidity*\* (both determined under the condition of constant temperature and for states of vanishing stress), we shall have as definitions:

$$V = -v \left( \frac{dp}{dv} \right)_t, \quad \text{when } v = r_0^3 v', \quad (448)$$

\* See Thomson and Tait's *Natural Philosophy*, vol. i, p. 711.

where  $\rho$  denotes a uniform pressure to which the solid is subjected,  $v$  its volume, and  $v'$  its volume in the state of reference; and

$$\left. \begin{aligned} R &= \frac{dX_{v'}}{d\frac{dx}{dy'}} = -\frac{d^2\psi_{v'}}{\left(d\frac{dx}{dy'}\right)^2}, \\ \text{when } &\frac{dx}{dy'} = \frac{dy}{dy'} = \frac{dz}{dz'} = r_0, \\ \text{and } &\frac{dx}{dy'} = \frac{dx}{dz'} = \frac{dy}{dz'} = \frac{dy}{dx'} = \frac{dz}{dx'} = \frac{dz}{dy'} = 0. \end{aligned} \right\} \quad (449)$$

Now when the solid is subject to uniform pressure on all sides, if we consider so much of it as has the volume unity in the state of reference, we shall have

$$r_1 = r_2 = r_3 = v^{\frac{1}{3}}, \quad (450)$$

and by (444) and (439),

$$\psi_{v'} = i + 3e v^{\frac{2}{3}} + 3f v^{\frac{4}{3}} + h v. \quad (451)$$

Hence, by equation (88), since  $\psi_v$  is equivalent to  $\psi$ ,

$$-\rho = \left(\frac{d\psi}{dv}\right)_t = 2e v^{-\frac{1}{3}} + 4f v^{\frac{1}{3}} + h, \quad (452)$$

$$-v \left(\frac{dp}{dv}\right)_t = -\frac{2}{3}e v^{-\frac{4}{3}} + \frac{4}{3}f v^{\frac{1}{3}}; \quad (453)$$

and by (448),

$$V = -\frac{2}{3} \frac{e}{r_0} + \frac{4}{3} f r_0. \quad (454)$$

To obtain the value of  $R$  in accordance with the definition (449), we may suppose the values of  $E$ ,  $F$ , and  $H$  given by equations (432), (434), and (437) to be substituted in equation (444). This will give for the value of  $R$

$$R = 2e + 4f r_0^2. \quad (455)$$

Moreover, since  $\rho$  must vanish in (452) when  $v = r_0^3$ , we have

$$2e + 4f r_0^2 + h r_0 = 0. \quad (456)$$

From the three last equations may be obtained the values of  $e$ ,  $f$ ,  $h$ , in terms of  $r_0$ ,  $V$ , and  $R$ ; viz.,

$$e = \frac{R + 3r_0 V}{4}, \quad f = \frac{R + 3r_0 V}{8r_0^2}, \quad h = -\frac{R}{r_0}. \quad (457)$$

The quantity  $r_0$ , like  $R$  and  $V$ , is a function of the temperature, the differential coefficient  $\frac{d \log r_0}{dt}$  representing the rate of linear expansion of the solid when without stress.

It will not be necessary to discuss equation (443) at length, as the case is entirely analogous to that which has just been treated. [It must be remembered that  $\eta_{Vi}$ , in the discussion of (443) will take the place everywhere of the temperature in the discussion of (444).] If we denote by  $V'$  and  $R'$  the *elasticity of volume* and the *rigidity*, both determined under the condition of *constant entropy*, (i. e., of *no transmission of heat*), and for states of vanishing stress, we shall have the equations:

$$V' = -\frac{2e'}{3r_0} + \frac{4}{3}f'r_0, \quad (458)$$

$$R' = 2e' + 4f'r_0^2, \quad (459)$$

$$2e' + 4f'r_0^2 + h'r_0 = 0. \quad (460)$$

Whence

$$e' = \frac{R' - 3r_0V'}{4}, \quad f' = \frac{R' + 3r_0V'}{8r_0^2}, \quad h' = -\frac{R'}{r_0}. \quad (461)$$

In these equations  $r_0$ ,  $R'$ , and  $V'$  are to be regarded as functions of the quantity  $\eta_{Vi}$ .

If we wish to change from one state of reference to another (also isotropic), the changes required in the fundamental equation are easily made. If  $a$  denotes the length of any line of the solid in the second state of reference divided by its length in the first, it is evident that when we change from the first state of reference to the second the values of the symbols  $\varepsilon_{Vi}$ ,  $\eta_{Vi}$ ,  $\psi_{Vi}$ ,  $H$  are divided by  $a^3$ , that of  $E$  by  $a^2$ , and that of  $F$  by  $a^4$ . In making the change of the state of reference, we must therefore substitute in the fundamental equation of the form (444)  $a^3\psi_{Vi}$ ,  $a^2E$ ,  $a^4F$ ,  $a^3H$  for  $\psi_{Vi}$ ,  $E$ ,  $F$ , and  $H$ , respectively. In the fundamental equation of the form (443), we must make the analogous substitutions, and also substitute  $a^3\eta_{Vi}$  for  $\eta_{Vi}$ . [It will be remembered that  $i'$ ,  $e'$ ,  $f'$ , and  $h'$  represent functions of  $\eta_{Vi}$ , and that it is only when their values in terms of  $\eta_{Vi}$  are substituted, that equation (443) becomes a fundamental equation.]

### *Concerning Solids which absorb Fluids.*

There are certain bodies which are solid with respect to some of their components, while they have other components which are fluid. In the following discussion, we shall suppose both the solidity and the fluidity to be perfect, so far as any properties are concerned which can affect the conditions of equilibrium,—i. e., we shall suppose that the solid matter of the body is entirely free from plasticity, and that there are no passive resistances to the motion of the fluid

components except such as vanish with the velocity of the motion,—leaving it to be determined by experiment how far and in what cases these suppositions are realized.

It is evident that equation (356) must hold true with regard to such a body, when the quantities of the fluid components contained in a given element of the solid remain constant. Let  $\Gamma'_a$ ,  $\Gamma'_b$ , etc., denote the quantities of the several fluid components contained in an element of the body divided by the volume of the element in the state of reference, or, in other words, let these symbols denote the densities which the several fluid components would have, if the body should be brought to the state of reference while the matter contained in each element remained unchanged. We may then say that equation (356) will hold true, when  $\Gamma'_a$ ,  $\Gamma'_b$ , etc., are constant. The complete value of the differential of  $\varepsilon_v$ , will therefore be given by an equation of the form

$$d\varepsilon_v = t d\eta_v + \Sigma \Sigma' \left( X_x d\frac{dx}{dx'} \right) + L_a d\Gamma'_a + L_b d\Gamma'_b + \text{etc.} \quad (462)$$

Now when the body is in a state of hydrostatic stress, the term in this equation containing the signs of summation will reduce to  $-p dv_v$ , ( $v_v$ , denoting, as elsewhere, the volume of the element divided by its volume in the state of reference). For in this case

$$X_x = -p \left( \frac{dy}{dy'} \frac{dz}{dz'} - \frac{dz}{dy'} \frac{dy}{dz'} \right), \quad (463)$$

$$\begin{aligned} \Sigma \Sigma' \left( X_x d\frac{dx}{dx'} \right) &= -p \Sigma \Sigma' \left\{ \left( \frac{dy}{dy'} \frac{dz}{dz'} - \frac{dz}{dy'} \frac{dy}{dz'} \right) d\frac{dx}{dx'} \right\} \\ &= -p d \begin{vmatrix} \frac{dx}{dx'} & \frac{dx}{dy'} & \frac{dx}{dz'} \\ \frac{dy}{dx'} & \frac{dy}{dy'} & \frac{dy}{dz'} \\ \frac{dz}{dx'} & \frac{dz}{dy'} & \frac{dz}{dz'} \end{vmatrix} \\ &= -p dv_v. \end{aligned} \quad (464)$$

We have, therefore, for a state of hydrostatic stress,

$$d\varepsilon_v = t d\eta_v - p dv_v + L_a d\Gamma'_a + L_b d\Gamma'_b + \text{etc.}, \quad (465)$$

and multiplying by the volume of the element in the state of reference, which we may regard as constant,

$$d\varepsilon = t d\eta - p dv + L_a dm_a + L_b dm_b + \text{etc.}, \quad (466)$$

where  $\varepsilon$ ,  $\eta$ ,  $v$ ,  $m_a$ ,  $m_b$ , etc., denote the energy, entropy, and volume of the element, and the quantities of its several fluid components. It is evident that the equation will also hold true, if these symbols are understood as relating to a homogeneous body of finite size. The only limitation with respect to the variations is that the element or body to which the symbols relate shall always contain the same solid matter. The varied state may be one of hydrostatic stress or otherwise.

But when the body is in a state of hydrostatic stress, and the solid matter is considered invariable, we have by equation (12)

$$d\varepsilon = t d\eta - p dv + \mu_a dm_a + \mu_b dm_b + \text{etc.} \quad (467)$$

It should be remembered that the equation cited occurs in a discussion which relates only to bodies of hydrostatic stress, so that the varied state as well as the initial is there regarded as one of hydrostatic stress. But a comparison of the two last equations shows that the last will hold true without any such limitation, and moreover, that the quantities  $L_a$ ,  $L_b$ , etc., when determined for a state of hydrostatic stress, are equal to the potentials  $\mu_a$ ,  $\mu_b$ , etc.

Since we have hitherto used the term *potential* solely with reference to bodies of hydrostatic stress, we may apply this term as we choose with regard to other bodies. We may therefore call the quantities  $L_a$ ,  $L_b$ , etc., the *potentials* for the several fluid components in the body considered, whether the state of the body is one of hydrostatic stress or not, since this use of the term involves only an extension of its former definition. It will also be convenient to use our ordinary symbol for a potential to represent these quantities. Equation (462) may then be written

$$d\varepsilon_v = t d\eta_v + \sum \left( X_x, d \frac{dx}{dx'} \right) + \mu_a d\Gamma'_a + \mu_b d\Gamma'_b + \text{etc.} \quad (468)$$

This equation holds true of solids having fluid components without any limitation with respect to the initial state or to the variations, except that the solid matter to which the symbols relate shall remain the same.

In regard to the conditions of equilibrium for a body of this kind, it is evident in the first place that if we make  $\Gamma'_a$ ,  $\Gamma'_b$ , etc., constant, we shall obtain from the general criterion of equilibrium all the conditions which we have obtained for ordinary solids, and which are expressed by the formulae (364), (374), (380), (382)–(384). The quantities  $\Gamma'_1$ ,  $\Gamma'_2$ , etc., in the last two formulae include of course

those which have just been represented by  $\Gamma'_a$ ,  $\Gamma'_b$ , etc., and which relate to the fluid components of the body, as well as the corresponding quantities relating to its solid components. Again, if we suppose the solid matter of the body to remain without variation in quantity or position, it will easily appear that the potentials for the substances which form the fluid components of the solid body must satisfy the same conditions in the solid body and in the fluids in contact with it, as in the case of entirely fluid masses. See eqs. (22).

The above conditions must however be slightly modified in order to make them sufficient for equilibrium. It is evident that if the solid is dissolved at its surface, the fluid components which are set free may be absorbed by the solid as well as by the fluid mass, and in like manner if the quantity of the solid is increased, the fluid components of the new portion may be taken from the previously existing solid mass. Hence, whenever the *solid* components of the solid body are actual components of the fluid mass, (whether the case is the same with the *fluid* components of the solid body or not,) an equation of the form (383) must be satisfied, in which the potentials  $\mu_a$ ,  $\mu_b$ , etc., contained implicitly in the second member of the equation are determined from the solid body. Also if the *solid* components of the solid body are all possible but not all actual components of the fluid mass, a condition of the form (384) must be satisfied, the values of the potentials in the second member being determined as in the preceding case.

The quantities

$$t, \quad X_{x'}, \dots, Z_{z'}, \quad \mu_a, \quad \mu_b, \quad \text{etc.}, \quad (469)$$

being differential coefficients of  $\varepsilon_v$ , with respect to the variables

$$\eta_{v'}, \quad \frac{dx}{dx'}, \dots, \frac{dz}{dz'}, \quad \Gamma'_a, \quad \Gamma'_b, \quad \text{etc.}, \quad (470)$$

will of course satisfy the necessary relations

$$\frac{\frac{dt}{dx}}{\frac{d\eta_{v'}}{dx'}} = \frac{dX_{x'}}{d\eta_{v'}}, \quad \text{etc.} \quad (471)$$

This result may be generalized as follows. Not only is the second member of equation (468) a complete differential in its present form, but it will remain such if we transfer the sign of differentiation ( $d$ ) from one factor to the other of any term (the sum indicated by the symbol  $\geq \Sigma'$  is here supposed to be expanded into nine terms), and at the same time change the sign of the term from + to -. For to

substitute  $- \eta_v dt$  for  $t d\eta_v$ , for example, is equivalent to subtracting the complete differential  $d(t\eta_v)$ . Therefore, if we consider the quantities in (469) and (470) which occur in any same term in equation (468) as forming a pair, we may choose as independent variables either quantity of each pair, and the differential coefficient of the remaining quantity of any pair with respect to the independent variable of another pair will be equal to the differential coefficient of the remaining quantity of the second pair with respect to the independent variable of the first, taken positively, if the independent variables of these pairs are both affected by the sign  $d$  in equation (468), or are neither thus affected, but otherwise taken negatively. Thus

$$\left( \frac{dX_{x'}}{dI_a'} \right)_{dx} = \left( \frac{d\mu_a}{d\overline{dx}} \right)_{\Gamma_a'}, \quad \left( \frac{dX_{x'}}{d\mu_a} \right)_{dx} = - \left( \frac{dI_a'}{d\overline{dx}} \right)_{\mu_a}, \quad (472)$$

$$\left( \frac{d\overline{dx}}{d\mu_a} \right)_{X_{x'}} = \left( \frac{dI_a'}{dX_{x'}} \right)_{\mu_a}, \quad \left( \frac{d\overline{dx}}{dI_a'} \right)_{X_{x'}} = - \left( \frac{d\mu_a}{dX_{x'}} \right)_{\Gamma_a'}, \quad (473)$$

where in addition to the quantities indicated by the suffixes, the following are to be considered as constant: either  $t$  or  $\eta_v$ , either  $X_y$ , or  $\frac{dx}{dy'}$ , . . . either  $Z_z$ , or  $\frac{dz}{dz'}$ , either  $\mu_b$  or  $I_b'$ , etc.

It will be observed that when the temperature is constant the conditions  $\mu_a = \text{const.}$ ,  $\mu_b = \text{const.}$  represent the physical condition of a body in contact with a fluid of which the phase does not vary, and which contains the components to which the potentials relate. Also that when  $I_a'$ ,  $I_b'$ , etc., are constant, the heat absorbed by the body in any infinitesimal change of condition per unit of volume measured in the state of reference is represented by  $t d\eta_v$ . If we denote this quantity by  $dQ_v$ , and use the suffix  $q$  to denote the condition of no transmission of heat, we may write

$$\left( \frac{d \log t}{d \overline{dx}} \right)_q = \left( \frac{dX_{x'}}{dQ_v} \right)_{\overline{dx}}, \quad \left( \frac{d \log t}{dX_{x'}} \right)_q = - \left( \frac{d\overline{dx}}{dQ_v} \right)_{X_{x'}}, \quad (474)$$

$$\left( \frac{dQ_v}{dX_{x'}} \right)_t = \left( \frac{d\overline{dx}}{d \log t} \right)_{X_{x'}}, \quad \left( \frac{dQ_v}{d\overline{dx}} \right)_t = - \left( \frac{dX_{x'}}{d \log t} \right)_{\overline{dx}}, \quad (475)$$

where  $I_a'$ ,  $I_b'$ , etc., must be regarded as constant in all the equations, and either  $X_y$ , or  $\frac{dx}{dy'}$ , . . . either  $Z_z$ , or  $\frac{dz}{dz'}$ , in each equation.

INFLUENCE OF SURFACES OF DISCONTINUITY UPON THE EQUILIBRIUM  
OF HETEROGENEOUS MASSES.—THEORY OF CAPILLARITY.

We have hitherto supposed, in treating of heterogeneous masses in contact, that they might be considered as separated by mathematical surfaces, each mass being unaffected by the vicinity of the others, so that it might be homogeneous quite up to the separating surfaces both with respect to the density of each of its various components and also with respect to the densities of energy and entropy. That such is not rigorously the case is evident from the consideration that if it were so with respect to the densities of the components it could not be so in general with respect to the density of energy, as the sphere of molecular action is not infinitely small. But we know from observation that it is only within very small distances of such a surface that any mass is sensibly affected by its vicinity,—a natural consequence of the exceedingly small sphere of sensible molecular action,—and this fact renders possible a simple method of taking account of the variations in the densities of the component substances and of energy and entropy, which occur in the vicinity of surfaces of discontinuity. We may use this term, for the sake of brevity, without implying that the discontinuity is absolute, or that the term distinguishes any surface with mathematical precision. It may be taken to denote the non-homogeneous film which separates homogeneous or nearly homogeneous masses.

Let us consider such a surface of discontinuity in a fluid mass which is in equilibrium and uninfluenced by gravity. For the precise measurement of the quantities with which we have to do, it will be convenient to be able to refer to a geometrical surface, which shall be sensibly coincident with the physical surface of discontinuity, but shall have a precisely determined position. For this end, let us take some point in or very near to the physical surface of discontinuity, and imagine a geometrical surface to pass through this point and all other points which are similarly situated with respect to the condition of the adjacent matter. Let this geometrical surface be called the *dividing surface*, and designated by the symbol S. It will be observed that the position of this surface is as yet to a certain extent arbitrary, but that the directions of its normals are already everywhere determined, since all the surfaces which can be formed in the manner described are evidently parallel to one another. Let us also imagine a closed surface cutting the surface S and including a part of the homogeneous mass on each side. We will so far limit the

form of this closed surface as to suppose that on each side of  $S$ , as far as there is any want of perfect homogeneity in the fluid masses, the closed surface is such as may be generated by a moving normal to  $S$ . Let the portion of  $S$  which is included by the closed surface be denoted by  $\mathbf{s}$ , and the area of this portion by  $s$ . Moreover, let the mass contained within the closed surface be divided into three parts by two surfaces, one on each side of  $S$ , and very near to that surface, although at such distance as to lie entirely beyond the influence of the discontinuity in its vicinity. Let us call the part which contains the surface  $\mathbf{s}$  (with the physical surface of discontinuity)  $M$ , and the homogeneous parts  $M'$  and  $M''$ , and distinguish by  $\epsilon, \epsilon', \epsilon'', \eta, \eta', \eta''$ ,  $m_1, m'_1, m''_1, m_2, m'_2, m''_2$ , etc., the energies and entropies of these masses, and the quantities which they contain of their various components.

It is necessary, however, to define more precisely what is to be understood in cases like the present by the energy of masses which are only separated from other masses by imaginary surfaces. A part of the total energy which belongs to the matter in the vicinity of the separating surface, relates to pairs of particles which are on different sides of the surface, and such energy is not in the nature of things referable to either mass by itself. Yet, to avoid the necessity of taking separate account of such energy, it will often be convenient to include it in the energies which we refer to the separate masses. When there is no break in the homogeneity at the surface, it is natural to treat the energy as distributed with a uniform density. This is essentially the case with the initial state of the system which we are considering, for it has been divided by surfaces passing in general through homogeneous masses. The only exception—that of the surface which cuts at right angles the non-homogeneous film—(apart from the consideration that without any important loss of generality we may regard the part of this surface within the film as very small compared with the other surfaces) is rather apparent than real, as there is no change in the state of the matter *in the direction perpendicular to this surface*. But in the variations to be considered in the state of the system, it will not be convenient to limit ourselves to such as do not create any discontinuity at the surfaces bounding the masses  $M, M', M''$ : we must therefore determine how we will estimate the energies of the masses in case of such infinitesimal discontinuities as may be supposed to arise. Now the energy of each mass will be most easily estimated by neglecting the discontinuity, i. e., if we estimate the energy on the supposition that

beyond the bounding surface the phase is identical with that within the surface. This will evidently be allowable, if it does not affect the total amount of energy. To show that it does not affect this quantity, we have only to observe that, if the energy of the mass on one side of a surface where there is an infinitesimal discontinuity of phase is greater as determined by this rule than if determined by any other (suitable) rule, the energy of the mass on the other side must be less by the same amount when determined by the first rule than when determined by the second, since the discontinuity relative to the second mass is equal but opposite in character to the discontinuity relative to the first.

If the entropy of the mass which occupies any one of the spaces considered is not in the nature of things determined without reference to the surrounding masses, we may suppose a similar method to be applied to the estimation of entropy.

With this understanding, let us return to the consideration of the equilibrium of the three masses  $M$ ,  $M'$ , and  $M''$ . We shall suppose that there are no limitations to the possible variations of the system due to any want of perfect mobility of the components by means of which we express the composition of the masses, and that these components are independent, i. e., that no one of them can be formed out of the others.

With regard to the mass  $M$ , which includes the surface of discontinuity, it is necessary for its internal equilibrium that when its boundaries are considered constant, and when we consider only *reversible* variations (i. e., those of which the opposite are also possible), the variation of its energy should vanish with the variations of its entropy and of the quantities of its various components. For changes within this mass will not affect the energy or the entropy of the surrounding masses (when these quantities are estimated on the principle which we have adopted), and it may therefore be treated as an isolated system. For fixed boundaries of the mass  $M$ , and for reversible variations, we may therefore write

$$\delta\varepsilon = A_0 \delta\eta + A_1 \delta m_1 + A_2 \delta m_2 + \text{etc.}, \quad (476)$$

where  $A_0$ ,  $A_1$ ,  $A_2$ , etc., are quantities determined by the initial (unvaried) condition of the system. It is evident that  $A_0$  is the temperature of the lamelliform mass to which the equation relates, or the *temperature at the surface of discontinuity*. By comparison of this equation with (12) it will be seen that the definition of  $A_1$ ,  $A_2$ , etc., is entirely analogous to that of the potentials in homo-

geneous masses, although the mass to which the former quantities relate is not homogeneous, while in our previous definition of potentials, only homogeneous masses were considered. By a natural extension of the term *potential*, we may call the quantities  $A_1$ ,  $A_2$ , etc., the *potentials at the surface of discontinuity*. This designation will be farther justified by the fact, which will appear hereafter, that the value of these quantities is independent of the thickness of the lamina ( $M$ ) to which they relate. If we employ our ordinary symbols for temperature and potentials, we may write

$$\delta\varepsilon = t\delta\eta + \mu_1\delta m_1 + \mu_2\delta m_2 + \text{etc.} \quad (477)$$

If we substitute  $\geq$  for  $=$  in this equation, the formula will hold true of all variations whether reversible or not;\* for if the variation of energy could have a value less than that of the second member of the equation, there must be variation in the condition of  $M$  in which its energy is diminished without change of its entropy or of the quantities of its various components.

It is important, however, to observe that for any given values of  $\delta\eta$ ,  $\delta m_1$ ,  $\delta m_2$ , etc., while there *may* be possible variations of the nature and state of  $M$  for which the value of  $\delta\varepsilon$  is greater than that of the second member of (477), there *must* always be possible variations for which the value of  $\delta\varepsilon$  is equal to that of the second member.

\* To illustrate the difference between variations which are reversible, and those which are not, we may conceive of two entirely different substances meeting in equilibrium at a mathematical surface without being at all mixed. We may also conceive of them as mixed in a thin film about the surface where they meet, and then the amount of mixture is capable of variation both by increase and by diminution. But when they are absolutely unmixed, the amount of mixture can be increased, but is incapable of diminution, and it is then consistent with equilibrium that the value of  $\delta\varepsilon$  (for a variation of the system in which the substances commence to mix) should be greater than the second member of (477). It is not necessary to determine whether precisely such cases actually occur; but it would not be legitimate to overlook the possible occurrence of cases in which variations may be possible while the opposite variations are not.

It will be observed that the sense in which the term *reversible* is here used is entirely different from that in which it is frequently used in treatises on thermodynamics, where a process by which a system is brought from a state A to a state B is called reversible, to signify that the system may also be brought from the state B to the state A through the same series of intermediate states taken in the reverse order by means of external agencies of the opposite character. The variation of a system from a state A to a state B (supposed to differ infinitely little from the first) is here called reversible when the system is capable of another state  $B'$  which bears the same relation to the state A that A bears to B.

It will be convenient to have a notation which will enable us to express this by an equation. Let  $\delta\varepsilon$  denote the smallest value (i. e., the value nearest to  $-\infty$ ) of  $\delta\varepsilon$  consistent with given values of the other variations, then

$$\delta\varepsilon = t \delta\eta + \mu_1 \delta m_1 + \mu_2 \delta m_2 + \text{etc.} \quad (478)$$

For the internal equilibrium of the whole mass which consists of the parts M, M', M'', it is necessary that

$$\delta\varepsilon + \delta\varepsilon' + \delta\varepsilon'' \geq 0 \quad (479)$$

for all variations which do not affect the enclosing surface or the total entropy or the total quantity of any of the various components. If we also regard the surfaces separating M, M', and M'' as invariable, we may derive from this condition, by equations (478) and (12), the following as a *necessary* condition of equilibrium:

$$\begin{aligned} & t \delta\eta + \mu_1 \delta m_1 + \mu_2 \delta m_2 + \text{etc.} \\ & + t' \delta\eta' + \mu_1' \delta m_1' + \mu_2' \delta m_2' + \text{etc.} \\ & + t'' \delta\eta'' + \mu_1'' \delta m_1'' + \mu_2'' \delta m_2'' + \text{etc.} \geq 0, \end{aligned} \quad (480)$$

the variations being subject to the equations of conditions

$$\left. \begin{aligned} \delta\eta + \delta\eta' + \delta\eta'' &= 0, \\ \delta m_1 + \delta m_1' + \delta m_1'' &= 0, \\ \delta m_2 + \delta m_2' + \delta m_2'' &= 0, \\ \text{etc.} & \end{aligned} \right\} \quad (481)$$

It may also be the case that some of the quantities  $\delta m_1'$ ,  $\delta m_1''$ ,  $\delta m_2'$ ,  $\delta m_2''$ , etc., are incapable of negative values or can only have the value zero. This will be the case when the substances to which these quantities relate are not actual or possible components of M' or M''. (See page 117.) To satisfy the above condition it is necessary and sufficient that

$$t = t' = t'', \quad (482)$$

$$\mu_1' \delta m_1' \geq \mu_1 \delta m_1', \quad \mu_2' \delta m_2' \geq \mu_2 \delta m_2', \quad \text{etc.,} \quad (483)$$

$$\mu_1'' \delta m_1'' \geq \mu_1 \delta m_1'', \quad \mu_2'' \delta m_2'' \geq \mu_2 \delta m_2'', \quad \text{etc.} \quad (484)$$

It will be observed that, if the substance to which  $\mu_1$ , for instance, relates is an actual component of each of the homogeneous masses, we shall have  $\mu_1 = \mu_1' = \mu_1''$ . If it is an actual component of the first only of these masses, we shall have  $\mu_1 = \mu_1'$ . If it is also a possible component of the second homogeneous mass, we shall also have  $\mu_1 \geq \mu_1''$ . If this substance occurs only at the surface of dis-

continuity, the value of the potential  $\mu_1$  will not be determined by any equation, but cannot be greater than the potential for the same substance in either of the homogeneous masses in which it may be a possible component.

It appears, therefore, that the particular conditions of equilibrium *relating to temperature and the potentials* which we have before obtained by neglecting the influence of the surfaces of discontinuity (pp. 119, 120, 128) are not invalidated by the influence of such discontinuity in their application to homogeneous parts of the system bounded like  $M'$  and  $M''$  by imaginary surfaces lying within the limits of homogeneity,—a condition which may be fulfilled by surfaces very near to the surfaces of discontinuity. It appears also that similar conditions will apply to the non-homogeneous films like  $M'$ , which separate such homogeneous masses. The properties of such films, which are of course different from those of homogeneous masses, require our farther attention.

The volume occupied by the mass  $M$  is divided by the surface  $S$  into two parts, which we will call  $v'''$  and  $v''''$ ,  $v'''$  lying next to  $M'$ , and  $v''''$  to  $M''$ . Let us imagine these volumes filled by masses having throughout the same temperature, pressure and potentials, and the same densities of energy and entropy, and of the various components, as the masses  $M'$  and  $M''$  respectively. We shall then have, by equation (12), if we regard the volumes as constant,

$$\delta\varepsilon''' = t' \delta\eta''' + \mu_1' \delta m_1''' + \mu_2' \delta m_2''' + \text{etc.}, \quad (485)$$

$$\delta\varepsilon'''' = t'' \delta\eta'''' + \mu_1'' \delta m_1'''' + \mu_2'' \delta m_2'''' + \text{etc.}; \quad (486)$$

whence, by (482)–(484), we have for reversible variations

$$\delta\varepsilon''' = t \delta\eta''' + \mu_1 \delta m_1''' + \mu_2 \delta m_2''' + \text{etc.}, \quad (487)$$

$$\delta\varepsilon'''' = t \delta\eta'''' + \mu_1 \delta m_1'''' + \mu_2 \delta m_2'''' + \text{etc.} \quad (488)$$

From these equations and (477), we have for reversible variations

$$\begin{aligned} \delta(\varepsilon - \varepsilon''' - \varepsilon''') &= t \delta(\eta - \eta''' - \eta''') \\ &+ \mu_1 \delta(m_1 - m_1''' - m_1''') + \mu_2 \delta(m_2 - m_2''' - m_2''') + \text{etc.} \end{aligned} \quad (489)$$

Or, if we set\*

$$\varepsilon^s = \varepsilon - \varepsilon''' - \varepsilon''', \quad \eta^s = \eta - \eta''' - \eta''', \quad (490)$$

$$m_1^s = m_1 - m_1''' - m_1''', \quad m_2^s = m_2 - m_2''' - m_2''', \quad \text{etc.,} \quad (491)$$

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\* It will be understood that the  $s$  here used is not an algebraic exponent, but is only intended as a distinguishing mark. The Roman letter  $S$  has not been used to denote any quantity.

we may write

$$\delta \varepsilon^s = t \delta \eta^s + \mu_1 \delta m_1^s + \mu_2 \delta m_2^s + \text{etc.} \quad (492)$$

This is true of reversible variations in which the surfaces which have been considered are fixed. It will be observed that  $\varepsilon^s$  denotes the excess of the energy of the actual mass which occupies the total volume which we have considered over that energy which it would have, if on each side of the surface  $S$  the density of energy had the same uniform value quite up to that surface which it has at a sensible distance from it; and that  $\eta^s$ ,  $m_1^s$ ,  $m_2^s$ , etc., have analogous significations. It will be convenient, and need not be a source of any misconception, to call  $\varepsilon^s$  and  $\eta^s$  the energy and entropy of *the surface* (or the *superficial* energy and entropy),  $\frac{\varepsilon^s}{s}$  and  $\frac{\eta^s}{s}$  the *superficial densities* of energy and entropy,  $\frac{m_1^s}{s}$ ,  $\frac{m_2^s}{s}$ , etc., the *superficial densities* of the several components.

Now these quantities ( $\varepsilon^s$ ,  $\eta^s$ ,  $m_1^s$ , etc.) are determined partly by the state of the physical system which we are considering, and partly by the various imaginary surfaces by means of which these quantities have been defined. The position of these surfaces, it will be remembered, has been regarded as fixed in the variation of the system. It is evident, however, that the form of that portion of these surfaces, which lies in the region of homogeneity on either side of the surface of discontinuity cannot affect the values of these quantities. To obtain the complete value of  $\delta \varepsilon^s$  for reversible variations, we have therefore only to regard variations in the position and form of the limited surface  $s$ , as this determines all of the surfaces in question lying within the region of non-homogeneity. Let us first suppose the form of  $s$  to remain unvaried and only its position in space to vary, either by translation or rotation. No change in (492) will be necessary to make it valid in this case. For the equation is valid if  $s$  remains fixed and the material system is varied in position; also, if the material system and  $s$  are both varied in position, while their relative position remains unchanged. Therefore, it will be valid if the surface alone varies its position.

But if the form of  $s$  be varied, we must add to the second member (492) terms which shall represent the value of

$$\delta \varepsilon^s - t \delta \eta^s - \mu_1 \delta m_1^s - \mu_2 \delta m_2^s - \text{etc.}$$

due to such variation in the form of  $s$ . If we suppose  $s$  to be suffi-

ciently small to be considered uniform throughout in its curvatures and in respect to the state of the surrounding matter, the value of the above expression will be determined by the variation of its area  $\delta s$  and the variations of its principal curvatures  $\delta c_1$  and  $\delta c_2$ , and we may write

$$\begin{aligned}\delta \varepsilon^s = & t \delta \eta^s + \mu_1 \delta m_1^s + \mu_2 \delta m_2^s + \text{etc.} \\ & + \sigma \delta s + C_1 \delta c_1 + C_2 \delta c_2,\end{aligned}\quad (493)$$

or

$$\begin{aligned}\delta \varepsilon^s = & t \delta \eta^s + \mu_1 \delta m_1^s + \mu_2 \delta m_2^s + \text{etc.} \\ & + \sigma \delta s + \frac{1}{2}(C_1 + C_2) \delta(c_1 + c_2) + \frac{1}{2}(C_1 - C_2) \delta(c_1 - c_2),\end{aligned}\quad (494)$$

$\sigma$ ,  $C_1$ , and  $C_2$  denoting quantities which are determined by the initial state of the system and position and form of  $\mathbf{s}$ . The above is the complete value of the variation of  $\varepsilon^s$  for reversible variations of the system. But it is always possible to give such a position to the surface  $\mathbf{s}$  that  $C_1 + C_2$  shall vanish.

To show this, it will be convenient to write the equation in the longer form [see (490), (491)]

$$\begin{aligned}\delta \varepsilon = & t \delta \eta - \mu_1 \delta m_1 - \mu_2 \delta m_2 - \text{etc.} \\ = & \delta \varepsilon''' + t \delta \eta''' + \mu_1 \delta m_1''' + \mu_2 \delta m_2''' + \text{etc.} \\ = & \delta \varepsilon''' + t \delta \eta''' + \mu_1 \delta m_1''' + \mu_2 \delta m_2''' + \text{etc.} \\ = & \sigma \delta s + \frac{1}{2}(C_1 + C_2) \delta(c_1 + c_2) + \frac{1}{2}(C_1 - C_2) \delta(c_1 - c_2),\end{aligned}\quad (495)$$

i. e., by (482)–(484) and (12),

$$\begin{aligned}\delta \varepsilon = & t \delta \eta - \mu_1 \delta m_1 - \mu_2 \delta m_2 - \text{etc.} + p' \delta v''' + p'' \delta v''' \\ = & \sigma \delta s + \frac{1}{2}(C_1 + C_2) \delta(c_1 + c_2) + \frac{1}{2}(C_1 - C_2) \delta(c_1 - c_2).\end{aligned}\quad (496)$$

From this equation it appears in the first place that the pressure is the same in the two homogeneous masses separated by a plane surface of discontinuity. For let us imagine the material system to remain unchanged, while the plane surface  $\mathbf{s}$  without change of area or of form moves in the direction of its normal. As this does not affect the boundaries of the mass  $M$ ,

$$\delta \varepsilon - t \delta \eta - \mu_1 \delta m_1 - \mu_2 \delta m_2 - \text{etc.} = 0.$$

Also  $\delta s = 0$ ,  $\delta(c_1 + c_2) = 0$ ,  $\delta(c_1 - c_2) = 0$ , and  $\delta v''' = -\delta v'''$ . Hence  $p' = p''$ , when the surface of discontinuity is plane.

Let us now examine the effect of different positions of the surface  $\mathbf{s}$  in the same material system upon the value of  $C_1 + C_2$ , supposing at first that in the initial state of the system the surface of discontinuity is plane. Let us give the surface  $\mathbf{s}$  some particular position. In the

initial state of the system this surface will of course be plane like the physical surface of discontinuity, to which it is parallel. In the varied state of the system, let it become a portion of a spherical surface having positive curvature; and at sensible distances from this surface let the matter be homogeneous and with the same phases as in the initial state of the system; also at and about the surface let the state of the matter so far as possible be the same as at and about the plane surface in the initial state of the system. (Such a variation in the system may evidently take place negatively as well as positively, as the surface may be curved toward either side. But whether such a variation is consistent with the maintenance of equilibrium is of no consequence, since in the preceding equations only the initial state is supposed to be one of equilibrium.) Let the surface  $\mathbf{s}$ , placed as supposed, whether in the initial or the varied state of the surface, be distinguished by the symbol  $\mathbf{s}'$ . Without changing either the initial or the varied state of the material system, let us make another supposition with respect to the imaginary surface  $\mathbf{s}$ . In the unvaried system let it be parallel to its former position but removed from it a distance  $\lambda$  on the side on which lie the centers of positive curvature. In the varied state of the system, let it be spherical and concentric with  $\mathbf{s}'$ , and separated from it by the same distance  $\lambda$ . It will of course lie on the same side of  $\mathbf{s}'$  as in the unvaried system. Let the surface  $\mathbf{s}$ , placed in accordance with this second supposition, be distinguished by the symbol  $\mathbf{s}''$ . Both in the initial and the varied state, let the perimeters of  $\mathbf{s}'$  and  $\mathbf{s}''$  be traced by a common normal. Now the value of

$$\delta\varepsilon - t\delta\eta - \mu_1\delta m_1 - \mu_2\delta m_2 - \text{etc.}$$

in equation (496) is not affected by the position of  $\mathbf{s}$ , being determined simply by the body  $M$ : the same is true  $p'\delta v''' + p''\delta v''''$  or  $p'\delta(v''' + v''''), v''' + v''''$  being the volume of  $M$ . Therefore the second member of (496) will have the same value whether the expressions relate to  $\mathbf{s}'$  or  $\mathbf{s}''$ . Moreover,  $\delta(c_1 - c_2) = 0$  both for  $\mathbf{s}'$  and  $\mathbf{s}''$ . If we distinguish the quantities determined for  $\mathbf{s}'$  and for  $\mathbf{s}''$  by the marks ' and ", we may therefore write

$$\sigma'\delta s' + \frac{1}{2}(C_1' + C_2')\delta(c_1' + c_2') = \sigma''\delta s'' + \frac{1}{2}(C_1'' + C_2'')\delta(c_1'' + c_2'').$$

Now if we make  $\delta s'' = 0$ ,

we shall have by geometrical necessity

$$\delta s' = s\lambda\delta(c_1'' + c_2'').$$

Hence

$$\sigma' s \lambda \delta(c_1'' + c_2'') + \frac{1}{2} (C_1' + C_2') \delta(c_1' + c_2') = \frac{1}{2} (C_2'' + C_2'') \delta(c_1'' + c_2'').$$

$$\text{But } \delta(c_1' + c_2') = \delta(c_1'' + c_2'').$$

$$\text{Therefore, } C_1' + C_2' + 2 \sigma' s \lambda = C_1'' + C_2''.$$

This equation shows that we may give a positive or negative value to  $C_1'' + C_2''$  by placing  $\mathbf{s}''$  a sufficient distance on one or on the other side of  $\mathbf{s}'$ . Since this is true when the (unvaried) surface is plane, it must also be true when the surface is nearly plane. And for this purpose a surface may be regarded as nearly plane, when the radii of curvature are very large in proportion to the thickness of the non-homogeneous film. This is the case when the radii of curvature have any sensible size. In general, therefore, whether the surface of discontinuity is plane or curved it is possible to place the surface  $\mathbf{s}$  so that  $C_1 + C_2$  in equation (494) shall vanish.

Now we may easily convince ourselves by equation (493) that if  $\mathbf{s}$  is placed within the non-homogeneous film, and  $s = 1$ , the quantity  $\sigma$  is of the same order of magnitude as the values of  $\varepsilon^s$ ,  $\eta^s$ ,  $m_1^s$ ,  $m_2^s$ , etc., while the values of  $C_1$  and  $C_2$  are of the same order of magnitude as the changes in the values of the former quantities caused by increasing the curvature of  $\mathbf{s}$  by unity. Hence, on account of the thinness of the non-homogeneous film, since it can be very little affected by such a change of curvature in  $\mathbf{s}$ , the values of  $C_1$  and  $C_2$  must in general be very small relatively to  $\sigma$ . And hence, if  $\mathbf{s}'$  be placed within the non-homogeneous film, the value of  $\lambda$  which will make  $C_1'' + C_2''$  vanish must be very small (of the same order of magnitude as the thickness of the non-homogeneous film). The position of  $\mathbf{s}$ , therefore, which will make  $C_1 + C_2$  in (494) vanish, will in general be sensibly coincident with the physical surface of discontinuity.

We shall hereafter suppose, when the contrary is not distinctly indicated that the surface  $\mathbf{s}$ , in the unvaried state of the system, has such a position as to make  $C_1 + C_2 = 0$ . It will be remembered that the surface  $\mathbf{s}$  is a part of a larger surface  $S$ , which we have called the *dividing surface*, and which is coextensive with the physical surface of discontinuity. We may suppose that the position of the dividing surface is everywhere determined by similar considerations. This is evidently consistent with the suppositions made on page 380 with regard to this surface.

We may therefore cancel the term

$$\frac{1}{2} (C_1 + C_2) \delta (c_1 + c_2)$$

in (494). In regard to the following term, it will be observed that  $C_1$  must necessarily be equal to  $C_2$ , when  $c_1 = c_2$ , which is the case when the surface of discontinuity is plane. Now on account of the thinness of the non-homogeneous film, we may always regard it as composed of parts which are approximately plane. Therefore, without danger of sensible error, we may also cancel the term

$$\frac{1}{2} (C_1 - C_2) \delta (c_1 - c_2).$$

Equation (494) is thus reduced to the form

$$\delta \varepsilon^s = t \delta \eta^s + \sigma \delta s + \mu_1 \delta m_1^s + \mu_2 \delta m_2^s + \text{etc.} \quad (497)$$

We may regard this as the complete value of  $\delta \varepsilon^s$ , for all reversible variations in the state of the system supposed initially in equilibrium, when the dividing surface has its initial position determined in the manner described.

The above equation is of fundamental importance in the theory of capillarity. It expresses a relation with regard to surfaces of discontinuity analogous to that expressed by equation (12) with regard to homogeneous masses. From the two equations may be directly deduced the conditions of equilibrium of heterogeneous masses in contact, subject or not to the action of gravity, without disregard of the influence of the surfaces of discontinuity. The general problem, including the action of gravity, we shall take up hereafter: at present we shall only consider, as hitherto, a small part of a surface of discontinuity with a part of the homogeneous mass on either side, in order to deduce the additional condition which may be found when we take account of the motion of the dividing surface.

We suppose as before that the mass especially considered is bounded by a surface of which all that lies in the region of non-homogeneity is such as may be traced by a moving normal to the dividing surface. But instead of dividing the mass as before into four parts, it will be sufficient to regard it as divided into two parts by the dividing surface. The energy, entropy, etc., of these parts, estimated on the supposition that its nature (including density of energy, etc.) is uniform quite up to the dividing surface, will be denoted by  $\varepsilon'$ ,  $\eta'$ , etc.,  $\varepsilon''$ ,  $\eta''$ , etc. Then the total energy will be  $\varepsilon^s + \varepsilon' + \varepsilon''$ , and the general condition of internal equilibrium will be that

$$\delta \varepsilon^s + \delta \varepsilon' + \delta \varepsilon'' \geq 0, \quad (498)$$

when the bounding surface is fixed, and the total entropy and total quantities of the various components are constant. We may suppose  $\eta^s, \eta', \eta'', m_1^s, m_1', m_1'', m_2^s, m_2', m_2'',$  etc., to be all constant. Then by (497) and (12) the condition reduces to

$$\sigma \delta s - p' \delta v' - p'' \delta v'' = 0. \quad (499)$$

(We may set  $=$  for  $\geq$ , since changes in the position of the dividing surface can evidently take place in either of two opposite directions.) This equation has evidently the same form as if a membrane without rigidity and having a tension  $\sigma$ , uniform in all directions, existed at the dividing surface. Hence, the particular position which we have chosen for this surface may be called the surface of tension, and  $\sigma$  the superficial tension. If all parts of the dividing surface move a uniform normal distance  $\delta N$ , we shall have

$$\delta s = (c_1 + c_2) s \delta N, \quad \delta v' = s \delta N, \quad \delta v'' = -s \delta N;$$

$$\text{whence} \quad \sigma (c_1 + c_2) = p' - p'', \quad (500)$$

the curvatures being positive when their centers lie on the side to which  $p'$  relates. This is the condition which takes the place of that of equality of pressure (see pp. 119, 128) for heterogeneous fluid masses in contact, when we take account of the influence of the surfaces of discontinuity. We have already seen that the conditions relating to temperature and the potentials are not affected by these surfaces.

#### *Fundamental Equations for Surfaces of Discontinuity.*

In equation (497) the initial state of the system is supposed to be one of equilibrium. The only limitation with respect to the varied state is that the variation shall be reversible, i. e., that an opposite variation shall be possible. Let us now confine our attention to variations in which the system remains in equilibrium. To distinguish this case, we may use the character  $d$  instead  $\delta$ , and write

$$ds = t dh^s + \sigma ds + \mu_1 dm_1^s + \mu_2 dm_2^s + \text{etc.} \quad (501)$$

Both the states considered being states of equilibrium, the limitation with respect to the reversibility of the variations may be neglected, since the variations will always be reversible in at least one of the states considered.

If we integrate this equation, supposing the area  $s$  to increase from zero to any finite value  $s$ , while the material system to a part of which the equation relates remains without change, we obtain

$$s^s = t \eta^s + \sigma s + \mu_1 m_1^s + \mu_2 m_2^s + \text{etc.}, \quad (502)$$

which may be applied to any portion of any surface of discontinuity (in equilibrium) which is of the same nature throughout, or throughout which the values of  $t$ ,  $\sigma$ ,  $\mu_1$ ,  $\mu_2$ , etc. are constant.

If we differentiate this equation, regarding all the quantities as variable, and compare the result with (501), we obtain

$$\eta^s dt + s d\sigma + m_1^s d\mu_1 + m_2^s d\mu_2 + \text{etc.} = 0. \quad (503)$$

If we denote the *superficial densities* of energy, of entropy, and of the several component substances (see page 386) by  $\varepsilon_s$ ,  $\eta_s$ ,  $\Gamma_1$ ,  $\Gamma_2$ , etc., we have

$$\varepsilon_s = \frac{\varepsilon^s}{s}, \quad \eta_s = \frac{\eta^s}{s}, \quad (504)$$

$$\Gamma_1 = \frac{m_1^s}{s}, \quad \Gamma_2 = \frac{m_2^s}{s}, \quad \text{etc.}, \quad (505)$$

and the preceding equations may be reduced to the form:—

$$d\varepsilon_s = t d\eta_s + \mu_1 d\Gamma_1 + \mu_2 d\Gamma_2 + \text{etc.}, \quad (506)$$

$$\varepsilon_s = t \eta_s + \sigma + \mu_1 \Gamma_1 + \mu_2 \Gamma_2 + \text{etc.}, \quad (507)$$

$$d\sigma = -\eta_s dt - \Gamma_1 d\mu_1 - \Gamma_2 d\mu_2 - \text{etc.} \quad (508)$$

Now the contact of the two homogeneous masses does not impose any restriction upon the variations of phase of either, except that the temperature and the potentials for actual components shall have the same value in both. [See (482)–(484) and (500).] For however the values of the pressures in the homogeneous masses may vary (on account of arbitrary variations of the temperature and potentials), and however the superficial tension may vary, equation (500) may always be satisfied by giving the proper curvature to the surface of tension, so long, at least, as the difference of pressures is not great. Moreover, if any of the potentials  $\mu_1$ ,  $\mu_2$ , etc. relate to substances which are found only at the surface of discontinuity, their values may be varied by varying the superficial densities of those substances. The values of  $t$ ,  $\mu_1$ ,  $\mu_2$ , etc. are therefore independently variable, and it appears from equation (508) that  $\sigma$  is a function of these quantities. If the form of this function is known, we may derive from it by differentiation  $n+1$  equations ( $n$  denoting the total number of component substances) giving the values of  $\eta_s$ ,  $\Gamma_1$ ,  $\Gamma_2$ , etc. in terms of the variables just mentioned. This will give us, with (507),  $n+3$  independent equations between the  $2n+4$  quantities which occur in that equation. These are all that exist, since  $n+1$

of these quantities are independently variable. Or, we may consider that we have  $n+3$  independent equations between the  $2n+5$  quantities occurring in equation (502), of which  $n+2$  are independently variable.

An equation, therefore, between

$$\sigma, \quad t, \quad \mu_1, \quad \mu_2, \quad \text{etc.,} \quad (509)$$

may be called a fundamental equation for the surface of discontinuity. An equation between

$$\varepsilon^s, \quad \eta^s, \quad s, \quad m_1^s \quad m_2^s, \quad \text{etc.,} \quad (510)$$

or between  $\varepsilon_s, \quad \eta_s, \quad \Gamma_1, \quad \Gamma_2, \quad \text{etc.,}$  (511)

may also be called a fundamental equation in the same sense. For it is evident from (501) that an equation may be regarded as subsisting between the variables (510), and if this equation be known, since  $n+2$  of the variables may be regarded as independent (viz.,  $n+1$  for the  $n+1$  variations in the nature of the surface of discontinuity, and one for the area of the surface considered), we may obtain by differentiation and comparison with (501),  $n+2$  additional equations between the  $2n+5$  quantities occurring in (502). Equation (506) shows that equivalent relations can be deduced from an equation between the variables (511). It is moreover quite evident that an equation between the variables (510) must be reducible to the form of an equation between the ratios of these variables, and therefore to an equation between the variables (511).

The same designation may be applied to any equation from which, by differentiation and the aid only of general principles and relations,  $n+3$  independent relations between the same  $2n+5$  quantities may be obtained.

If we set  $\psi^s = \varepsilon^s - t \eta^s,$  (512)

we obtain by differentiation and comparison with (501)

$$d\psi^s = -\eta^s dt + \sigma ds + \mu_1 dm_1^s + \mu_2 dm_2^s + \text{etc.} \quad (513)$$

An equation, therefore, between  $\psi^s, t, s, m_1^s, m_2^s, \text{etc.,}$  is a fundamental equation, and is to be regarded as entirely equivalent to either of the other fundamental equations which have been mentioned.

The reader will not fail to notice the analogy between these fundamental equations, which relate to surfaces of discontinuity, and those relating to homogeneous masses, which have been described on pages 140–144.

*On the Experimental Determination of Fundamental Equations for Surfaces of Discontinuity.*

When all the substances which are found at a surface of discontinuity are components of one or the other of the homogeneous masses, the potentials  $\mu_1, \mu_2$ , etc., as well as the temperature, may be determined from these homogeneous masses.\* The tension  $\sigma$  may be determined by means of the relation (500). But our measurements are practically confined to cases in which the difference of the pressures in the homogeneous masses is small; for with increasing differences of pressure the radii of curvature soon become too small for measurement. Therefore, although the equation  $p' = p''$  (which is equivalent to an equation between  $t, \mu_1, \mu_2$ , etc., since  $p'$  and  $p''$  are both functions of these variables) may not be exactly satisfied in cases in which it is convenient to measure the tension, yet this equation is so nearly satisfied in all the measurements of tension which we can make, that we must regard such measurements as simply establishing the values of  $\sigma$  for values of  $t, \mu_1, \mu_2$ , etc., which satisfy the equation  $p' = p''$ , but not as sufficient to establish the rate of change in the value of  $\sigma$  for variations of  $t, \mu_1, \mu_2$ , etc., which are inconsistent with the equation  $p' = p''$ .

To show this more distinctly, let  $t, \mu_2, m_3$ , etc. remain constant, then by (508) and (98)

$$d\sigma = -\Gamma_1 d\mu_1,$$

$$dp' = \gamma_1' d\mu_1,$$

$$dp'' = \gamma_1'' d\mu_1,$$

$\gamma_1'$  and  $\gamma_1''$  denoting the densities  $\frac{m_1'}{v'}$  and  $\frac{m_1''}{v''}$ . Hence,

$$dp' - dp'' = (\gamma_1' - \gamma_1'') d\mu_1,$$

and  $\Gamma_1 d(p' - p'') = (\gamma_1'' - \gamma_1') d\sigma$ .

But by (500)

$$(c_1 + c_2) d\sigma + \sigma d(c_1 + c_2) = d(p' - p'').$$

Therefore,

$$\Gamma_1 (c_1 + c_2) d\sigma + \Gamma_1 \sigma d(c_1 + c_2) = (\gamma_1'' - \gamma_1') d\sigma,$$

$$\text{or } \{\gamma_1'' - \gamma_1' - \Gamma_1 (c_1 + c_2)\} d\sigma = \Gamma_1 \sigma d(c_1 + c_2).$$

\* It is here supposed that the thermodynamic properties of the homogeneous masses have already been investigated, and that the fundamental equations of these masses may be regarded as known.

Now  $\Gamma_1(c_1 + c_2)$  will generally be very small compared with  $\gamma_1'' - \gamma_1'$ . Neglecting the former term, we have

$$\frac{d\sigma}{\sigma} = \frac{\Gamma_1}{\gamma_1'' - \gamma_1'} d(c_1 + c_2).$$

To integrate this equation, we may regard  $\Gamma_1, \gamma_1', \gamma_1''$  as constant. This will give, as an approximate value,

$$\log \frac{\sigma}{\sigma'} = \frac{\Gamma_1}{\gamma_1'' - \gamma_1'} (c_1 + c_2),$$

$\sigma'$  denoting the value of  $\sigma$  when the surface is plane. From this it appears that when the radii of curvature have any sensible magnitude, the value of  $\sigma$  will be sensibly the same as when the surface is plane and the temperature and all the potentials except one have the same values, unless the component for which the potential has not the same value has very nearly the same density in the two homogeneous masses, in which case, the condition under which the variations take place is nearly equivalent to the condition that the pressures shall remain equal.

Accordingly, we cannot in general expect to determine the superficial density  $\Gamma_1$  from its value  $-\left(\frac{d\sigma}{d\mu_1}\right)_{t, \mu}^*$  by measurements of superficial tensions. The case will be the same with  $\Gamma_2, \Gamma_3$ , etc., and also with  $\eta_s$ , the superficial density of entropy.

The quantities  $\epsilon_s, \eta_s, \Gamma_1, \Gamma_2$ , etc. are evidently too small in general to admit of direct measurement. When one of the components, however, is found only at the surface of discontinuity, it may be more easy to measure its superficial density than its potential. But except in this case, which is of secondary interest, it will generally be easy to determine  $\sigma$  in terms of  $t, \mu_1, \mu_2$ , etc., with considerable accuracy for plane surfaces, and extremely difficult or impossible to determine the fundamental equation more completely.

#### *Fundamental Equations for Plane Surfaces of Discontinuity.*

An equation giving  $\sigma$  in terms of  $t, \mu_1, \mu_2$ , etc., which will hold true only so long as the surface of discontinuity is plane, may be called a fundamental equation for a plane surface of discontinuity. It will be interesting to see precisely what results can be obtained from such an equation, especially with respect to the energy and entropy

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\* The suffixed  $\mu$  is used to denote that all the potentials except that occurring in the denominator of the differential coefficient are to be regarded as constant.

and the quantities of the component substances in the vicinity of the surface of discontinuity.

These results can be exhibited in a more simple form, if we deviate to a certain extent from the method which we have been following. The particular position adopted for the dividing surface (which determines the superficial densities) was chosen in order to make the term  $\frac{1}{2}(C_1+C_2)\delta(c_1+c_2)$  in (494) vanish. But when the curvature of the surface is not supposed to vary, such a position of the dividing surface is not necessary for the simplification of the formula. It is evident that equation (501) will hold true for plane surfaces (supposed to remain such) without reference to the position of the dividing surfaces, except that it shall be parallel to the surface of discontinuity. We are therefore at liberty to choose such a position for the dividing surface as may for any purpose be convenient.

None of the equations (502)–(513), which are either derived from (501), or serve to define new symbols, will be affected by such a change in the position of the dividing surface. But the expressions  $\varepsilon^s$ ,  $\eta^s$ ,  $m_1^s$ ,  $m_2^s$ , etc., as also  $\varepsilon_s$ ,  $\eta_s$ ,  $\Gamma_1$ ,  $\Gamma_2$ , etc. and  $\psi^s$ , will of course have different values when the position of that surface is changed. The quantity  $\sigma$ , however, which we may regard as defined by equations (501), or, if we choose, by (502) or (507), will not be affected in value by such a change. For if the dividing surface be moved a distance  $\lambda$  measured normally and toward the side to which  $v''$  relates, the quantities

$$\varepsilon_s, \quad \eta_s, \quad \Gamma_1, \quad \Gamma_2, \quad \text{etc.},$$

will evidently receive the respective increments

$\lambda(\varepsilon_v'' - \varepsilon_v')$ ,  $\lambda(\eta_v'' - \eta_v')$ ,  $\lambda(\gamma_1'' - \gamma_1')$ ,  $\lambda(\gamma_2'' - \gamma_2')$ , etc.,  $\varepsilon_v'$ ,  $\varepsilon_v''$ ,  $\eta_v'$ ,  $\eta_v''$  denoting the densities of energy and entropy in the two homogeneous masses. Hence, by equation (507),  $\sigma$  will receive the increment

$$\lambda(\varepsilon_v'' - \varepsilon_v') - t\lambda(\eta_v'' - \eta_v') - \mu_1\lambda(\gamma_1'' - \gamma_1') - \mu_2\lambda(\gamma_2'' - \gamma_2') - \text{etc.}$$

But by (93)

$$\begin{aligned} -\rho'' &= \varepsilon_v'' - t\eta_v'' - \mu_1\gamma_1'' - \mu_2\gamma_2'' - \text{etc.}, \\ -\rho' &= \varepsilon_v' - t\eta_v' - \mu_1\gamma_1' - \mu_2\gamma_2' - \text{etc.} \end{aligned}$$

Therefore, since  $\rho' = \rho''$ , the increment in the value of  $\sigma$  is zero. The value of  $\sigma$  is therefore independent of the position of the dividing surface, when this surface is plane. But when we call this quantity the superficial tension, we must remember that it will not have

its characteristic properties as a tension with reference to any arbitrary surfacee. Considered as a tension, its position is in the surface which we have called the surface of tension, and, strictly speaking, nowhere else. The positions of the dividing surfacee, however, which we shall consider, will not vary from the surface of tension sufficiently to make this distinction of any practical importance.

It is generally possible to place the dividing surface so that the total quantity of any desired component in the vicinity of the surface of discontinuity shall be the same as if the density of that component were uniform on each side quite up to the dividing surface. In other words, we may place the dividing surface so as to make any one of the quantities  $\Gamma_1$ ,  $\Gamma_2$ , etc., vanish. The only exception is with regard to a component which has the same density in the two homogeneous masses. With regard to a component which has very nearly the same density in the two masses such a location of the dividing surfacee might be objectionable, as the dividing surface might fail to coincide sensibly with the physical surface of discontinuity. Let us suppose that  $\gamma'_1$  is not equal (nor very nearly equal) to  $\gamma''_1$ , and that the dividing surface is so placed as to make  $\Gamma_1 = 0$ . Then equation (508) reduces to

$$d\sigma = -\eta_{s(1)} dt - \Gamma_{2(1)} d\mu_2 - \Gamma_{3(1)} d\mu_3 - \text{etc.}, \quad (514)$$

where the symbols  $\eta_{s(1)}$ ,  $\Gamma_{2(1)}$ , etc., are used for greater distinctness to denote the values of  $\eta_s$ ,  $\Gamma_2$ , etc., as determined by a dividing surface placed so that  $\Gamma_1 = 0$ . Now we may consider all the differentials in the second member of this equation as independent, without violating the condition that the surface shall remain plane, i. e., that  $dp' = dp''$ . This appears at once from the values of  $dp'$  and  $dp''$  given by equation (98). Moreover, as has already been observed, when the fundamental equations of the two homogeneous masses are known, the equation  $p' = p''$  affords a relation between the quantities  $t$ ,  $\mu_1$ ,  $\mu_2$ , etc. Hence, when the value of  $\sigma$  is also known for plane surfacees in terms of  $t$ ,  $\mu_1$ ,  $\mu_2$ , etc., we can eliminate  $\mu_1$  from this expression by means of the relation derived from the equality of pressures, and obtain the value of  $\sigma$  for plane surfacees in terms of  $t$ ,  $\mu_2$ ,  $\mu_3$ , etc. From this, by differentiation, we may obtain directly the values of  $\eta_{s(1)}$ ,  $\Gamma_{2(1)}$ ,  $\Gamma_{3(1)}$ , etc., in terms of  $t$ ,  $\mu_2$ ,  $\mu_3$ , etc. This would be a convenient form of the fundamental equation. But, if the elimination of  $p'$ ,  $p''$ , and  $\mu_1$  from the finite equations presents algebraic difficulties, we can in all cases easily eliminate  $dp'$ ,  $dp''$ ,  $d\mu_1$  from the corresponding differential equations and thus obtain a

differential equation from which the values of  $\eta_{S(1)}$ ,  $\Gamma_{2(1)}$ ,  $\Gamma_{3(1)}$ , etc. in terms of  $t$ ,  $\mu_1$ ,  $\mu_2$ , etc., may be at once obtained by comparison with (514).\*

\* If liquid mercury meets the mixed vapors of water and mercury in a plane surface, and we use  $\mu_1$  and  $\mu_2$  to denote the potentials of mercury and water respectively, and place the dividing surface so that  $\Gamma_1=0$ , i. e., so that the total quantity of mercury is the same as if the liquid mercury reached this surface on one side and the mercury vapor on the other without change of density on either side, then  $\Gamma_{2(1)}$  will represent the amount of water in the vicinity of this surface, per unit of surface, above that which there would be, if the water-vapor just reached the surface without change of density, and this quantity (which we may call the quantity of water condensed upon the surface of the mercury) will be determined by the equation

$$\Gamma_{2(1)} = -\frac{d\sigma}{d\mu_2}.$$

(In this differential coefficient as well as the following, the temperature is supposed to remain constant and the surface of discontinuity plane. Practically, the latter condition may be regarded as fulfilled in the case of any ordinary curvatures.)

If the pressure in the mixed vapors conforms to the law of Dalton (see pp. 215, 218), we shall have for constant temperature

$$dp_2 = \gamma_2 d\mu_2,$$

where  $p_2$  denotes the part of the pressure in the vapor due to the water-vapor, and  $\gamma_2$  the density of the water-vapor. Hence we obtain

$$\Gamma_{2(1)} = -\gamma_2 \frac{d\sigma}{dp_2}.$$

For temperatures below 100° centigrade, this will certainly be accurate, since the pressure due to the vapor of mercury may be neglected.

The value of  $\sigma$  for  $p_2=0$  and the temperature of 20° centigrade must be nearly the same as the superficial tension of mercury in contact with air, or 55.03 grammes per linear metre according to Quincke (Pogg. Ann., Bd. 139, p. 27). The value of  $\sigma$  at the same temperature, when the condensed water begins to have the properties of water in mass, will be equal to the sum of the superficial tensions of mercury in contact with water and of water in contact with its own vapor. This will be, according to the same authority, 42.58 + 8.25, or 50.83 grammes per metre, if we neglect the difference of the tensions of water with its vapor and water with air. As  $p_2$ , therefore, increases from zero to 236400 grammes per square metre (when water begins to be condensed *in mass*),  $\sigma$  diminishes from about 55.03 to about 50.83 grammes per linear metre. If the general course of the values of  $\sigma$  for intermediate values of  $p_2$  were determined by experiment, we could easily form an approximate estimate of the values of the superficial density  $\Gamma_{2(1)}$  for different pressures less than that of saturated vapor. It will be observed that the determination of the superficial density does not by any means depend upon inappreciable differences of superficial tension. The greatest difficulty in the determination would doubtless be that of distinguishing between the diminution of superficial tension due to the water and that due to other substances which might accidentally be present. Such determinations are of considerable practical importance on account of the use of mercury in measurements of the specific gravity of vapors.

The same physical relations may of course be deduced without giving up the use of the surface of tension as a dividing surface, but the formulæ which express them will be less simple. If we make  $t, \mu_3, \mu_4$ , etc. constant, we have by (98) and (508)

$$\begin{aligned} dp' &= \gamma_1' d\mu_1 + \gamma_2' d\mu_2, \\ dp'' &= \gamma_1'' d\mu_1 + \gamma_2'' d\mu_2, \\ d\sigma &= -\Gamma_1 d\mu_1 - \Gamma_2 d\mu_2, \end{aligned}$$

where we may suppose  $\Gamma_1$  and  $\Gamma_2$  to be determined with reference to the surface of tension. Then, if  $dp' = dp''$ ,

$$(\gamma_1' - \gamma_1'') d\mu_1 + (\gamma_2' - \gamma_2'') d\mu_2 = 0,$$

and

$$d\sigma = \Gamma_1 \frac{\gamma_2' - \gamma_2''}{\gamma_1' - \gamma_1''} d\mu_2 - \Gamma_2 d\mu_2.$$

That is,

$$\left( \frac{d\sigma}{d\mu_2} \right)_{p' = p'', t, \mu_3, \mu_4, \text{etc.}} = -\Gamma_2 + \Gamma_1 \frac{\gamma_2' - \gamma_2''}{\gamma_1' - \gamma_1''}. \quad (515)$$

The reader will observe that  $\frac{\Gamma_1}{\gamma_1' - \gamma_1''}$  represents the distance between the surface of tension and that dividing surface which would make  $\Gamma_1 = 0$ ; the second member of the last equation is therefore equivalent to  $-\Gamma_{2(1)}$ .

If any component substance has the same density in the two homogeneous masses separated by a plane surface of discontinuity, the value of the superficial density for that component is independent of the position of the dividing surface. In this case alone we may derive the value of the superficial density of a component with reference to the surface of tension from the fundamental equation for plane surfaces alone. Thus in the last equation, when  $\gamma_2' = \gamma_2''$ , the second member will reduce to  $-\Gamma_2$ . It will be observed that to make  $p' = p'', t, \mu_3, \mu_4$ , etc. constant is in this case equivalent to making  $t, \mu_1, \mu_3, \mu_4$ , etc. constant.

Substantially the same is true of the superficial density of entropy or of energy, when either of these has the same density in the two homogeneous masses.\*

\* With respect to questions which concern only the *form* of surfaces of discontinuity, such precision as we have employed in regard to the position of the dividing surface is evidently quite unnecessary. This precision has not been used for the sake of the mechanical part of the problem, which does not require the surface to be defined with greater nicety than we can employ in our observations, but in order to give

*Concerning the Stability of Surfaces of Discontinuity.*

We shall first consider the stability of a film separating homogeneous masses with respect to changes in its nature, while its position and the nature of the homogeneous masses are not altered. For this purpose, it will be convenient to suppose that the homogeneous masses are very large, and thoroughly stable with respect to the possible formation of any different homogeneous masses out of their components, and that the surface of discontinuity is plane and uniform.

Let us distinguish the quantities which relate to the actual components of one or both of the homogeneous masses by the suffixes  $a$ ,  $b$ , etc., and those which relate to components which are found only at the surface of discontinuity by the suffixes  $s$ ,  $t$ , etc., and consider the variation of the energy of the whole system in consequence of a given change in the nature of a small part of the surface of discontinuity, while the entropy of the whole system and the total quantities of the several components remain constant, as well as the volume of each of the homogeneous masses, as determined by the surface of tension. This small part of the surface of discontinuity in its changed state is supposed to be still uniform in nature, and such as may subsist in equilibrium between the given homogeneous masses, which will evidently not be sensibly altered in nature or thermodynamic state. The remainder of the surface of discontinuity is also supposed to remain uniform, and on account of its infinitely greater size to be infinitely less altered in its nature than the first part. Let  $\Delta\epsilon^s$  denote the increment of the superficial energy of this first part,  $\Delta\eta^s$ ,  $\Delta m_a^s$ ,  $\Delta m_b^s$ , etc.,  $\Delta m_g^s$ ,  $\Delta m_h^s$ , etc., the increments of its superficial

determinate values to the superficial densities of energy, entropy, and the component substances, which quantities, as has been seen, play an important part in the relations between the tension of a surface of discontinuity, and the composition of the masses which it separates.

The product  $\sigma s$  of the superficial tension and the area of the surface, may be regarded as the *available energy* due to the surface in a system in which the temperature and the potentials  $\mu_1$ ,  $\mu_2$ , etc.—or the differences of these potentials and the gravitational potential (see page 208) when the system is subject to gravity—are maintained sensibly constant. The value of  $\sigma$ , as well as that of  $s$ , is sensibly independent of the precise position which we may assign to the dividing surface (so long as this is sensibly coincident with the surface of discontinuity), but  $\epsilon_s$ , the *superficial density of energy*, as the term is used in this paper, like the superficial densities of entropy and of the component substances, requires a more precise localization of the dividing surface.

entropy and of the quantities of the components which we regard as belonging to the surface. The increments of entropy and of the various components which the rest of the system receive will be expressed by

$$-\Delta\eta^S, -\Delta m_a^S, -\Delta m_b^S, \text{ etc.}, -\Delta m_g^S, -\Delta m_h^S, \text{ etc.},$$

and the consequent increment of energy will be by (12) and (501)

$$-t\Delta\eta^S + \mu_a\Delta m_a^S + \mu_b\Delta m_b^S + \text{etc.} - \mu_g\Delta m_g^S - \mu_h\Delta m_h^S + \text{etc.}$$

Hence the total increment of energy in the whole system will be

$$\left. \begin{aligned} \Delta\epsilon^S - t\Delta\eta^S + \mu_a\Delta m_a^S + \mu_b\Delta m_b^S + \text{etc.} \\ - \mu_g\Delta m_g^S - \mu_h\Delta m_h^S + \text{etc.} \end{aligned} \right\} \quad (516)$$

If the value of this expression is necessarily positive, for finite changes as well as infinitesimal in the nature of the part of the film to which  $\Delta\epsilon^S$ , etc. relate,\* the increment of energy of the whole system will be positive for any possible changes in the nature of the film, and the film will be stable, at least with respect to changes in its nature, as distinguished from its position. For, if we write

$$\Delta\epsilon^S, \Delta\eta^S, \Delta m_a^S, \Delta m_b^S, \text{ etc.}, \Delta m_g^S, \Delta m_h^S, \text{ etc.}$$

for the energy, etc. of any element of the surface of discontinuity, we have from the supposition just made

$$\begin{aligned} \Delta D\epsilon^S - t\Delta D\eta^S + \mu_a\Delta Dm_a^S + \mu_b\Delta Dm_b^S + \text{etc.} \\ - \mu_g\Delta Dm_g^S - \mu_h\Delta Dm_h^S + \text{etc.} > 0; \end{aligned} \quad (517)$$

and integrating for the whole surface, since

$$\Delta f'Dm_g^S = 0, \Delta f'Dm_h^S = 0, \text{ etc.},$$

we have

$$\Delta f'D\epsilon^S - t\Delta f'D\eta^S + \mu_a\Delta f'Dm_a^S + \mu_b\Delta f'Dm_b^S + \text{etc.} > 0. \quad (518)$$

Now  $\Delta f'D\eta^S$  is the increment of the entropy of the whole surface, and  $-\Delta f'D\eta^S$  is therefore the increment of the entropy of the two homogeneous masses. In like manner,  $-\Delta f'Dm_a^S$ ,  $-\Delta f'Dm_b^S$ , etc. are the increments of the quantities of the components in these masses. The expression

$$-t\Delta f'D\eta^S + \mu_a\Delta f'Dm_a^S + \mu_b\Delta f'Dm_b^S + \text{etc.}$$

\* In the case of infinitesimal changes in the nature of the film, the sign  $\Delta$  must be interpreted, as elsewhere in this paper, without neglect of infinitesimals of the higher orders. Otherwise, by equation (501), the above expression would have the value zero.

denotes therefore, according to equation (12), the increment of energy of the two homogeneous masses, and since  $\Delta f D\epsilon^s$  denotes the increment of energy of the surface, the above condition expresses that the increment of the total energy of the system is positive. That we have only considered the possible formation of such films as are capable of existing in equilibrium between the given homogeneous masses can not invalidate the conclusion in regard to the stability of the film, for in considering whether any state of the system will have less energy than the given state, we need only consider the state of least energy, which is necessarily one of equilibrium.

If the expression (516) is capable of a negative value for an infinitesimal change in the nature of the part of the film to which the symbols relate, the film is obviously unstable.

If the expression is capable of a negative value, but only for finite and not for infinitesimal changes in the nature of this part of the film, the film is *practically unstable*,\* i. e., if such a change were made in a small part of the film, the disturbance would tend to increase. But it might be necessary that the initial disturbance should also have a finite magnitude in respect to the extent of surface in which it occurs; for we cannot suppose that the thermodynamic relations of an infinitesimal part of a surface of discontinuity are independent of the adjacent parts. On the other hand, the changes which we have been considering are such that every part of the film remains in equilibrium with the homogeneous masses on each side; and if the energy of the system can be diminished by a finite change satisfying this condition, it may perhaps be capable of diminution by an infinitesimal change which does not satisfy the same condition. We must therefore leave it undetermined whether the film, which in this case is practically unstable, is or is not unstable in the strict mathematical sense of the term.

Let us consider more particularly the condition of practical stability, in which we need not distinguish between finite and infinitesimal changes. To determine whether the expression (516) is capable of a negative value, we need only consider the least value of which it is capable. Let us write it in the fuller form

$$\begin{aligned} \epsilon^{su} - \epsilon^{st} - t(\eta^{su} - \eta^{st}) - \mu_a(m_a^{su} - m_a^{st}) - \mu_b(m_b^{su} - m_b^{st}) - \text{etc.}, \\ - \mu_g'(m_g^{su} - m_g^{st}) - \mu_h'(m_h^{su} - m_h^{st}) - \text{etc.}, \end{aligned} \quad \left. \right\} \quad (519)$$

where the single and double accents distinguish the quantities which

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\* With respect to the sense in which this term is used, compare page 133.

relate to the first and second states of the film, the letters without accents denoting those quantities which have the same value in both states. The differential of this expression when the quantities distinguished by double accents are alone considered variable, and the area of the surface is constant, will reduce by (501) to the form

$$(\mu_g'' - \mu_g') dm_g'' + (\mu_h'' - \mu_h') dm_h'' + \text{etc.}$$

To make this incapable of a negative value, we must have

$$\mu_g'' = \mu_g', \quad \text{unless} \quad m_g'' = 0,$$

$$\mu_h'' = \mu_h', \quad \text{unless} \quad m_h'' = 0.$$

In virtue of these relations and by equation (502), the expression (519), i. e., (516), will reduce to

$$\sigma'' s - \sigma' s,$$

which will be positive or negative according as

$$\sigma'' - \sigma' \quad (520)$$

is positive or negative.

That is, if the tension of the film is less than that of any other film which can exist between the same homogeneous masses (which has therefore the same values of  $t$ ,  $\mu_a$ ,  $\mu_b$ , etc.), and which moreover has the same values of the potentials  $\mu_g$ ,  $\mu_h$ , etc., so far as it contains the substances to which these relate, then the first film will be stable. But the film will be practically unstable, if any other such film has a less tension. [Compare the expression (141), by which the practical stability of homogeneous masses is tested.]

It is, however, evidently necessary for the stability of the surface of discontinuity with respect to *deformation*, that the value of the superficial tension should be positive. Moreover, since we have by (502) for the surface of discontinuity

$$\varepsilon^S - t \eta^S - \mu_a m_a^S - \mu_b m_b^S - \text{etc.} - \mu_g m_g^S - \mu_h m_h^S - \text{etc.} = \sigma s,$$

and by (93) for the two homogeneous masses

$$\varepsilon' - t \eta' + p v' - \mu_a m_a' - \mu_b m_b' - \text{etc.} = 0,$$

$$\varepsilon'' - t \eta'' + p v'' - \mu_a m_a'' - \mu_b m_b'' - \text{etc.} = 0,$$

if we denote by

$$\varepsilon, \quad \eta, \quad v, \quad m_a, \quad m_b, \quad \text{etc.}, \quad m_g, \quad m_h, \quad \text{etc.},$$

the total energy, etc. of a composite mass consisting of two such homogeneous masses divided by such a surface of discontinuity, we shall have by addition of these equations

$$\varepsilon - t\eta + p\nu - \mu_a m_a - \mu_b m_b - \text{etc.} - \mu_g m_g - \mu_h m_h - \text{etc.} = \sigma s.$$

Now if the value of  $\sigma$  is negative, the value of the first member of this equation will decrease as  $s$  increases, and may therefore be decreased by making the mass to consist of thin alternate strata of the two kinds of homogeneous masses which we are considering. There will be no limit to the decrease which is thus possible with a given value of  $v$ , so long as the equation is applicable, i. e., so long as the strata have the properties of similar bodies in mass. But it may easily be shown (as in a similar case on pages 131, 132) that when the values of

$$t, \quad p, \quad \mu_a, \quad \mu_b, \quad \text{etc.}, \quad \mu_g, \quad \mu_h, \quad \text{etc.}$$

are regarded as fixed, being determined by the surface of discontinuity in question, and the values of

$$\varepsilon, \quad \eta, \quad m_a, \quad m_b, \quad \text{etc.}, \quad m_g, \quad m_h, \quad \text{etc.}$$

are variable and may be determined by any body having the given volume  $v$ , the first member of this equation cannot have an infinite negative value, and must therefore have a least possible value, which will be negative, if any value is negative, that is, if  $\sigma$  is negative.

The body determining  $\varepsilon$ ,  $\eta$ , etc. which will give this least value to this expression will evidently be sensibly homogeneous. With respect to the formation of such a body, the system consisting of the two homogeneous masses and the surface of discontinuity with the negative tension is by (53) (see also page 133) at least practically unstable, if the surface of discontinuity is very large, so that it can afford the requisite material without sensible alteration of the values of the potentials. (This limitation disappears, if all the component substances are found in the homogeneous masses.) Therefore, in a system satisfying the conditions of practical stability with respect to the possible formation of all kinds of homogeneous masses, negative tensions of the surfaces of discontinuity are necessarily excluded.

Let us now consider the condition which we obtain by applying (516) to infinitesimal changes. The expression may be expanded as before to the form (519), and then reduced by equation (502) to the form

$$s(\sigma'' - \sigma') + m_g^{s''}(\mu_g'' - \mu_g') + m_h^{s''}(\mu_h'' - \mu_h') + \text{etc.}$$

That the value of this expression shall be positive when the quantities are determined by two films which differ infinitely little is a necessary condition of the stability of the film to which the single

accents relate. But if one film is stable, the other will in general be so too, and the distinction between the films with respect to stability is of importance only at the limits of stability. If all films for all values of  $\mu_g$ ,  $\mu_h$ , etc. are stable, or all within certain limits, it is evident that the value of the expression must be positive when the quantities are determined by any two infinitesimally different films within the same limits. For such collective determinations of stability the condition may be written

$$-s \Delta\sigma - m_g^s \Delta\mu_g - m_h^s \Delta\mu_h - \text{etc.} > 0,$$

or

$$\Delta\sigma < -I_g \Delta\mu_g - I_h \Delta\mu_h - \text{etc.} \quad (521)$$

On comparison of this formula with (508), it appears that within the limits of stability the second and higher differential coefficients of the tension considered as a function of the potentials for the substances which are found only at the surface of discontinuity (the potentials for the substances found in the homogeneous masses and the temperature being regarded as constant) satisfy the conditions which would make the tension a maximum if the necessary conditions relative to the first differential coefficients were fulfilled.

In the foregoing discussion of stability, the surface of discontinuity is supposed plane. In this case, as the tension is supposed positive, there can be no tendency to a change of form of the surface. We now pass to the consideration of changes consisting in or connected with motion and change of form of the surface of tension, which we shall at first suppose to be and to remain spherical and uniform throughout.

In order that the equilibrium of a spherical mass entirely surrounded by an indefinitely large mass of different nature shall be neutral with respect to changes in the value of  $r$ , the radius of the sphere, it is evidently necessary that equation (500), which in this may be written

$$2\sigma = r(p' - p''), \quad (522)$$

as well as the other conditions of equilibrium, shall continue to hold true for varying values of  $r$ . Hence, for a state of equilibrium which is on the limit between stability and instability, it is necessary that the equation

$$2d\sigma = (p' - p'')dr + rdp'$$

shall be satisfied, when the relations between  $d\sigma$ ,  $dp'$ , and  $dr$  are determined from the fundamental equations on the supposition that

the conditions of equilibrium relating to temperature and the potentials remain satisfied. (The differential coefficients in the equations which follow are to be determined on this supposition.) Moreover, if

$$r \frac{dp'}{dr} < 2 \frac{d\sigma}{dr} - p' + p'', \quad (523)$$

i. e., if the pressure of the interior mass increases less rapidly (or decreases more rapidly) with increasing radius than is necessary to preserve neutral equilibrium, the equilibrium is stable. But if

$$r \frac{dp'}{dr} > 2 \frac{d\sigma}{dr} - p' + p'', \quad (524)$$

the equilibrium is unstable. In the remaining case, when

$$r \frac{dp'}{dr} = 2 \frac{d\sigma}{dr} - p' + p'', \quad (525)$$

farther conditions are of course necessary to determine absolutely whether the equilibrium is stable or unstable, but in general the equilibrium will be stable in respect to change in one direction and unstable in respect to change in the opposite direction, and is therefore to be considered unstable. In general, therefore, we may call (523) the condition of stability.

When the interior mass and the surface of discontinuity are formed entirely of substances which are components of the external mass,  $p'$  and  $\sigma$  cannot vary and condition (524) being satisfied the equilibrium is unstable.

But if either the interior homogeneous mass or the surface of discontinuity contains substances which are not components of the enveloping mass, the equilibrium may be stable. If there is but one such substance, and we denote its densities and potential by  $\gamma'_1$ ,  $\Gamma_1$ , and  $\mu_1$ , the condition of stability (523) will reduce to the form

$$\left( r \frac{dp'}{d\mu_1} - 2 \frac{d\sigma}{d\mu_1} \right) \frac{d\mu_1}{dr} < p'' - p',$$

or, by (98) and (508),

$$(r \gamma'_1 + 2 \Gamma_1) \frac{d\mu_1}{dr} < p'' - p'. \quad (526)$$

In these equations and in all which follow in the discussion of this case, the temperature and the potentials  $\mu_2$ ,  $\mu_3$ , etc. are to be regarded as constant. But

$$\gamma'_1 v' + \Gamma_1 s,$$

which represents the total quantity of the component specified by the suffix, must be constant. It is evidently equal to

$$\frac{4}{3}\pi r^3 \gamma_1' + 4\pi r^2 \Gamma_1.$$

Dividing by  $4\pi$  and differentiating, we obtain

$$(r^2 \gamma_1' + 2r \Gamma_1) dr + \frac{1}{3}r^3 d\gamma_1' + r^2 d\Gamma_1 = 0,$$

or, since  $\gamma_1'$  and  $\Gamma_1$  are functions of  $\mu_1$ ,

$$(r \gamma_1' + 2 \Gamma_1) dr + \left( \frac{r^2}{3} \frac{d\gamma_1'}{d\mu_1} + r \frac{d\Gamma_1}{d\mu_1} \right) d\mu_1 = 0. \quad (527)$$

By means of this equation, the condition of stability is brought to the form

$$\frac{(r \gamma_1' + 2 \Gamma_1)^2}{\frac{r^2}{3} \frac{d\gamma_1'}{d\mu_1} + r \frac{d\Gamma_1}{d\mu_1}} > p' - p''. \quad (528)$$

If we eliminate  $r$  by equation (522), we have

$$\frac{\left( \frac{\gamma_1'}{p' - p''} + \frac{\Gamma_1}{\sigma} \right)^2}{\frac{1}{3(p' - p'')} \frac{d\gamma_1'}{d\mu_1} + \frac{1}{2\sigma} \frac{d\Gamma_1}{d\mu_1}} > 1. \quad (529)$$

If  $p'$  and  $\sigma$  are known in terms of  $t, \mu_1, \mu_2$ , etc., we may express the first member of this condition in terms of the same variables and  $p''$ . This will enable us to determine, for any given state of the external mass, the values of  $\mu_1$  which will make the equilibrium stable or unstable.

If the component to which  $\gamma_1'$  and  $\Gamma_1$  relate is found only at the surface of discontinuity, the condition of stability reduces to

$$\frac{\Gamma_1^2}{\sigma} \frac{d\mu_1}{d\Gamma_1} > \frac{1}{2}. \quad (530)$$

Since  $\Gamma_1 = -\frac{d\sigma}{d\mu_1}$ ,

we may also write

$$\frac{\Gamma_1}{\sigma} \frac{d\sigma}{d\Gamma_1} < -\frac{1}{2}, \quad \text{or} \quad \frac{d \log \sigma}{d \log \Gamma_1} < -\frac{1}{2}. \quad (531)$$

Again, if  $\Gamma_1 = 0$  and  $\frac{d\Gamma_1}{d\mu_1} = 0$ , the condition of stability reduces to

$$\frac{\frac{3}{2} \gamma_1'^2}{p' - p''} \frac{d\mu_1}{d\gamma_1'} > 1. \quad (532)$$

Since  $\gamma_1' = \frac{dp'}{d\mu_1}$ ,

we may also write

$$\frac{\gamma_1'}{p' - p''} \frac{dp'}{d\gamma_1'} > \frac{1}{3}, \quad \text{or} \quad \frac{d \log (p' - p'')}{d \log \gamma_1'} > \frac{1}{3}. \quad (533)$$

When  $r$  is large, this will be a close approximation for any values of  $\Gamma_1$ , unless  $\gamma_1'$  is very small. The two special conditions (531) and (533) might be derived from very elementary considerations.

Similar conditions of stability may be found when there are more substances than one in the inner mass or the surface of discontinuity, which are not components of the enveloping mass. In this case, we have instead of (526) a condition of the form

$$(r \gamma_1' + 2 \Gamma_1) \frac{d\mu_1}{dr} + (r \gamma_2' + 2 \Gamma_2) \frac{d\mu_2}{dr} + \text{etc.} < p'' - p', \quad (534)$$

from which  $\frac{d\mu_1}{dr}$ ,  $\frac{d\mu_2}{dr}$ , etc. may be eliminated by means of equations derived from the conditions that

$$\gamma_1' v' + \Gamma_1 s, \quad \gamma_2' v' + \Gamma_2 s, \quad \text{etc.}$$

must be constant.

Nearly the same method may be applied to the following problem. Two different homogeneous fluids are separated by a diaphragm having a circular orifice, their volumes being invariable except by the motion of the surface of discontinuity, which adheres to the edge of the orifice:—to determine the stability or instability of this surface when in equilibrium.

The condition of stability derived from (522) may in this case be written

$$r \frac{d(p' - p'')}{dv'} < 2 \frac{d\sigma}{dv'} - (p' - p'') \frac{dr}{dv'}, \quad (535)$$

where the quantities relating to the concave side of the surface of tension are distinguished by a single accent.

If both the masses are infinitely large, or if one which contains all the components of the system is infinitely large,  $p' - p''$  and  $\sigma$  will be constant, and the condition reduces to

$$\frac{dr}{dv'} < 0.$$

The equilibrium will therefore be stable or unstable according as the surface of tension is less or greater than a hemisphere.

To return to the general problem:—if we denote by  $x$  the part of the axis of the circular orifice intercepted between the center of the orifice and the surface of tension, by  $R$  the radius of the orifice, and by  $V'$  the value of  $v'$  when the surface of tension is plane, we shall have the geometrical relations

$$R^2 = 2 r x - x^2,$$

$$\begin{aligned} \text{and} \quad v' &= V' + \frac{2}{3} \pi r^2 x - \frac{1}{3} \pi R^2 (r - x) \\ &= V' + \pi r x^2 - \frac{1}{3} \pi x^3. \end{aligned}$$

By differentiation we obtain

$$(r - x) dx + x dr = 0,$$

and  $dv' = \pi x^2 dr + (2\pi r x - \pi x^2) dx;$

whence  $(r - x) dv' = -\pi r x^2 dr.$  (536)

By means of this relation, the condition of stability may be reduced to the form

$$\frac{dp'}{dv'} - \frac{dp''}{dv'} - \frac{2}{r} \frac{d\sigma}{dv'} < (p' - p'') \frac{r - x}{\pi r^2 x^2}. \quad (537)$$

Let us now suppose that the temperature and all the potentials except one,  $\mu_1$ , are to be regarded as constant. This will be the case when one of the homogeneous masses is very large and contains all the components of the system except one, or when both these masses are very large and there is a single substance at the surface of discontinuity which is not a component of either; also when the whole system contains but a single component, and is exposed to a constant temperature at its surface. Condition (537) will reduce by (98) and (508) to the form

$$\left(\gamma_1' - \gamma_1'' + \frac{2\Gamma_1}{r}\right) \frac{d\mu_1}{dv'} < (p' - p'') \frac{r - x}{\pi r^2 x^2}. \quad (538)$$

But  $\gamma_1' v' + \gamma_1'' v'' + \Gamma_1 s$

(the total quantity of the component specified by the suffix) must be constant; therefore, since

$$dv'' = -dv', \text{ and } ds = \frac{2}{r} dv',$$

$$\left(v' \frac{d\gamma_1'}{d\mu_1} + v'' \frac{d\gamma_1''}{d\mu_1} + s \frac{d\Gamma_1}{d\mu_1}\right) d\mu_1 + \left(\gamma_1' - \gamma_1'' + \frac{2\Gamma_1}{r}\right) dv' = 0. \quad (539)$$

By this equation, the condition of stability is brought to the form

$$\frac{\left(\gamma_1' - \gamma_1'' + \frac{2\Gamma_1}{r}\right)^2}{v' \frac{d\gamma_1'}{d\mu_1} + v'' \frac{d\gamma_1''}{d\mu_1} + s \frac{d\Gamma_1}{d\mu_1}} > (p' - p'') \frac{x - r}{\pi x^2 r^2}. \quad (540)$$

When the substance specified by the suffix is a component of either of the homogeneous masses, the terms  $\frac{2\Gamma_1}{r}$  and  $s \frac{d\Gamma_1}{d\mu_1}$  may generally be neglected. When it is not a component of either, the terms  $\gamma_1'$ ,  $\gamma_1''$ ,  $v' \frac{d\gamma_1'}{d\mu_1}$ ,  $v'' \frac{d\gamma_1''}{d\mu_1}$  may of course be cancelled, but we must not apply the formula to cases in which the substance spreads over the diaphragm separating the homogeneous masses.

In the cases just discussed, the problem of the stability of certain surfaces of tension has been solved by considering the case of neutral equilibrium,—a condition of neutral equilibrium affording the equation of the limit of stability. This method probably leads as directly as any to the result, when that consists in the determination of the value of a certain quantity at the limit of stability, or of the relation which exists at that limit between certain quantities specifying the state of the system. But problems of a more general character may require a more general treatment.

Let it be required to ascertain the stability or instability of a fluid system in a given state of equilibrium with respect to motion of the surfaces of tension and accompanying changes. It is supposed that the conditions of internal stability for the separate homogeneous masses are satisfied, as well as those conditions of stability for the surfaces of discontinuity which relate to small portions of these surfaces with the adjacent masses. (The conditions of stability which are here supposed to be satisfied have been already discussed in part and will be farther discussed hereafter.) The fundamental equations for all the masses and surfaces occurring in the system are supposed to be known. In applying the general criteria of stability which are given on page 110, we encounter the following difficulty.

The question of the stability of the system is to be determined by the consideration of states of the system which are slightly varied from that of which the stability is in question. These varied states of the system are not in general states of equilibrium, and the relations expressed by the fundamental equations may not hold true of them. More than this,—if we attempt to describe a varied state of the system by varied values of the quantities which describe the initial state, if these varied values are such as are inconsistent with equilibrium, they may fail to determine with precision any state of the system. Thus, when the phases of two contiguous homogeneous masses are specified, if these phases are such as satisfy all the conditions of equilibrium, the nature of the surface of discontinuity (if without additional components) is entirely determined; but if the phases do not satisfy all the conditions of equilibrium, the nature of the surface of discontinuity is not only undetermined, but incapable of determination by specified values of such quantities as we have employed to express the nature of surfaces of discontinuity in equilibrium. For example, if the temperatures in contiguous homogeneous masses are different, we cannot specify the thermal state of the surface of discontinuity by assigning to it any particular temperature. It would be

necessary to give the law by which the temperature passes over from one value to the other. And if this were given, we could make no use of it in the determination of other quantities, unless the rate of change of the temperature were so gradual, that at every point we could regard the thermodynamic state as unaffected by the change of temperature in its vicinity. It is true that we are also ignorant in respect to surfaces of discontinuity *in equilibrium* of the law of change of those quantities which are different in the two phases in contact, such as the densities of the components, but this, although unknown to us, is entirely determined by the nature of the phases in contact, so that no vagueness is occasioned in the definition of any of the quantities which we have occasion to use with reference to such surfaces of discontinuity.

It may be observed that we have established certain differential equations, especially (497), in which only the initial state is necessarily one of equilibrium. Such equations may be regarded as establishing certain properties of states bordering upon those of equilibrium. But these are properties which hold true only when we disregard quantities proportional to the square of those which express the degree of variation of the system from equilibrium. Such equations are therefore sufficient for the determination of the conditions of equilibrium, but not sufficient for the determination of the conditions of stability.

We may, however, use the following method to decide the question of stability in such a case as has been described.

Beside the real system of which the stability is in question, it will be convenient to conceive of another system, to which we shall attribute in its initial state the same homogeneous masses and surfaces of discontinuity which belong to the real system. We shall also suppose that the homogeneous masses and surfaces of discontinuity of this system, which we may call the imaginary system, have the same fundamental equations as those of the real system. But the imaginary system is to differ from the real in that the variations of its state are limited to such as do not violate the conditions of equilibrium relating to temperature and the potentials, and that the fundamental equations of the surfaces of discontinuity hold true for these varied states, although the condition of equilibrium expressed by equation (500) may not be satisfied.

Before proceeding farther, we must decide whether we are to examine the question of stability under the condition of a constant external temperature, or under the condition of no transmission of

heat to or from external bodies, and in general, to what external influences we are to regard the system as subject. It will be convenient to suppose that the exterior of the system is fixed, and that neither matter nor heat can be transmitted through it. Other cases may easily be reduced to this, or treated in a manner entirely analogous.

Now if the real system in the given state is unstable, there must be some slightly varied state in which the energy is less, but the entropy and the quantities of the components the same as in the given state, and the exterior of the system unvaried. But it may easily be shown that the given state of the system may be made stable by constraining the surfaces of discontinuity to pass through certain fixed lines situated in the unvaried surfaces. Hence, if the surfaces of discontinuity are constrained to pass through corresponding fixed lines in the surfaces of discontinuity belonging to the varied state just mentioned, there must be a state of stable equilibrium for the system thus constrained which will differ infinitely little from the given state of the system, the stability of which is in question, and will have the same entropy, quantities of components and exterior, but less energy. The imaginary system will have a similar state, since the real and imaginary systems do not differ in respect to those states which satisfy all the conditions of equilibrium for each surface of discontinuity. That is, the imaginary system has a state, differing infinitely little from the given state, and with the same entropy, quantities of components, and exterior, but with less energy.

Conversely, if the imaginary system has such a state as that just described, the real system will also have such a state. This may be shown by fixing certain lines in the surfaces of discontinuity of the imaginary system in its state of less energy and then making the energy a minimum under the conditions. The state thus determined will satisfy all the conditions of equilibrium for each surface of discontinuity, and the real system will therefore have a corresponding state, in which the entropy, quantities of components, and exterior will be the same as in the given state, but the energy less.

We may therefore determine whether the given system is or is not unstable, by applying the general criterion of instability (7) to the imaginary system.

If the system is not unstable, the equilibrium is either neutral or stable. Of course we can determine which of these is the case by reference to the imaginary system, since this determination depends upon states of equilibrium, in regard to which the real and imaginary

systems do not differ. We may therefore determine whether the equilibrium of the given system is stable, neutral, or unstable, by applying the criteria (3)–(7) to the imaginary system.

The result which we have obtained may be expressed as follows:—In applying to a fluid system which is in equilibrium, and of which all the small parts taken separately are stable, the criteria of stable, neutral, and unstable equilibrium, we may regard the system as under constraint to satisfy the conditions of equilibrium relating to temperature and the potentials, and as satisfying the relations expressed by the fundamental equations for masses and surfaces, even when the condition of equilibrium relating to pressure [equation (500)] is not satisfied.

It follows immediately from this principle, in connection with equations (501) and (86), that in a stable system each surface of tension must be a surface of minimum area for constant values of the volumes which it divides, when the other surfaces bounding these volumes and the perimeter of the surface of tension are regarded as fixed; that in a system in neutral equilibrium each surface of tension will have as small an area as it can receive by any slight variations under the same limitations; and that in seeking the remaining conditions of stable or neutral equilibrium, when these are satisfied, it is only necessary to consider such varied surfaces of tension as have similar properties with reference to the varied volumes and perimeters.

We may illustrate the method which has been described by applying it to a problem but slightly different from one already (pp. 408, 409) discussed by a different method. It is required to determine the conditions of stability for a system in equilibrium, consisting of two different homogeneous masses meeting at a surface of discontinuity, the perimeter of which is invariable, as well as the exterior of the whole system, which is also impermeable to heat.

To determine what is necessary for stability in addition to the condition of minimum area for the surface of tension, we need only consider those varied surfaces of tension which satisfy the same condition. We may therefore regard the surfaces of tension as determined by  $v'$ , the volume of one of the homogeneous masses. But the state of the system would evidently be completely determined by the position of the surface of tension and the temperature and potentials, if the entropy and the quantities of the components were variable; and therefore, since the entropy and the quantities of the components are constant, the state of the system must be completely determined by the position of the surface of tension. We may therefore regard

all the quantities relating to the system as functions of  $v'$ , and the condition of stability may be written

$$\frac{d\varepsilon}{dv'} dv' + \frac{1}{2} \frac{d^2\varepsilon}{dv'^2} dv'^2 + \text{etc.} > 0,$$

where  $\varepsilon$  denotes the total energy of the system. Now the conditions of equilibrium require that

$$\frac{d\varepsilon}{dv'} = 0,$$

Hence, the general condition of stability is that

$$\frac{d^2\varepsilon}{dv'^2} > 0. \quad (541)$$

Now if we write  $\varepsilon'$ ,  $\varepsilon''$ ,  $\varepsilon^s$  for the energies of the two masses and of the surface, we have by (86) and (501), since the total entropy and the total quantities of the several components are constant,

$$d\varepsilon = d\varepsilon' + d\varepsilon'' + d\varepsilon^s = -p' dv' - p'' dv'' + \sigma ds,$$

or, since  $dv'' = -dv'$ ,

$$\frac{d\varepsilon}{dv'} = -p' + p'' + \sigma \frac{ds}{dv'}. \quad (542)$$

Hence,

$$\frac{d^2\varepsilon}{dv'^2} = -\frac{dp'}{dv'} + \frac{dp''}{dv'} + \frac{d\sigma}{dv'} \frac{ds}{dv'} + \sigma \frac{d^2s}{dv'^2}, \quad (543)$$

and the condition of stability may be written

$$\sigma \frac{d^2s}{dv'^2} > \frac{dp'}{dv'} - \frac{dp''}{dv'} - \frac{d\sigma}{dv'} \frac{ds}{dv'}. \quad (544)$$

If we now simplify the problem by supposing, as in the similar case on page 409, that we may disregard the variations of the temperature and of all the potentials except one, the condition will reduce to

$$\sigma \frac{d^2s}{dv'^2} > \left( \gamma_1' - \gamma_1'' + \Gamma_1 \frac{ds}{dv'} \right) \frac{d\mu_1}{dv'}. \quad (545)$$

The total quantity of the substance indicated by the suffix  $_1$  is

$$\gamma_1' v' + \gamma_1'' v'' + \Gamma_1 s.$$

Making this constant, we have

$$\left( \gamma_1' - \gamma_1'' + \Gamma_1 \frac{ds}{dv'} \right) dv' + \left( v' \frac{d\gamma_1'}{d\mu_1} + v'' \frac{d\gamma_1''}{d\mu_1} + s \frac{d\Gamma_1}{d\mu_1} \right) d\mu_1 = 0. \quad (546)$$

The condition of equilibrium is thus reduced to the form

$$\sigma \frac{d^2s}{dv'^2} > - \frac{\left( \gamma_1' - \gamma_1'' + \Gamma_1 \frac{ds}{dv'} \right)^2}{v' \frac{d\gamma_1'}{d\mu_1} + v'' \frac{d\gamma_1''}{d\mu_1} + s \frac{d\Gamma_1}{d\mu_1}}, \quad (547)$$

where  $\frac{ds}{dv'}$  and  $\frac{d^2s}{dv'^2}$  are to be determined from the form of the surface of tension by purely geometrical considerations, and the other differential coefficients are to be determined from the fundamental equations of the homogeneous masses and the surface of discontinuity. Condition (540) may be easily deduced from this as a particular case.

The condition of stability with reference to motion of surfaces of discontinuity admits of a very simple expression when we can treat the temperature and potentials as constant. This will be the case when one or more of the homogeneous masses, containing together all the component substances, may be considered as indefinitely large, the surfaces of discontinuity being finite. For if we write  $\Sigma \Delta \varepsilon$  for the sum of the variations of the energies of the several homogeneous masses, and  $\Sigma \Delta \varepsilon^s$  for the sum of the variations of the energies of the several surfaces of discontinuity, the condition of stability may be written

$$\Sigma \Delta \varepsilon + \Sigma \Delta \varepsilon^s > 0, \quad (548)$$

the total entropy and the total quantities of the several components being constant. The variations to be considered are infinitesimal, but the character  $\Delta$  signifies, as elsewhere in this paper, that the expression is to be interpreted without neglect of infinitesimals of the higher orders. Since the temperature and potentials are sensibly constant, the same will be true of the pressures and surface-tensions, and by integration of (86) and (501) we may obtain for any homogeneous mass

$$\Delta \varepsilon = t \Delta \eta - p \Delta v + \mu_1 \Delta m_1 + \mu_2 \Delta m_2 + \text{etc.},$$

and for any surface of discontinuity

$$\Delta \varepsilon^s = t \Delta \eta^s + \sigma \Delta s + \mu_1^s \Delta m_1^s + \mu_2^s \Delta m_2^s + \text{etc.}$$

These equations will hold true of finite differences, when  $t, p, \sigma, \mu_1, \mu_2$ , etc. are constant, and will therefore hold true of infinitesimal differences, under the same limitations, without neglect of the infinitesimals of the higher orders. By substitution of these values, the condition of stability will reduce to the form

$$- \Sigma(p \Delta v) + \Sigma(\sigma \Delta s) > 0,$$

$$\text{or} \quad \Sigma(p \Delta v) - \Sigma(\sigma \Delta s) < 0. \quad (549)$$

That is, the sum of the products of the volumes of the masses by their pressures diminished by the sum of the products of the areas of the surfaces of discontinuity by their tensions must be a maximum. This is a purely geometrical condition, since the pressures and ten-

sions are constant. This condition is of interest, because it is always sufficient for stability with reference to motion of surfaces of discontinuity. For any system may be reduced to the kind described by putting certain parts of the system in communication (by means of fine tubes if necessary) with large masses of the proper temperatures and potentials. This may be done without introducing any new movable surfaces of discontinuity. The condition (549) when applied to the altered system is therefore the same as when applied to the original system. But it is sufficient for the stability of the altered system, and therefore sufficient for its stability if we diminish its freedom by breaking the connection between the original system and the additional parts, and therefore sufficient for the stability of the original system.

*On the Possibility of the Formation of a Fluid of different Phase within any Homogeneous Fluid.*

The study of surfaces of discontinuity throws considerable light upon the subject of the stability of such homogeneous fluid masses as have a less pressure than others formed of the same components (or some of them) and having the same temperature and the same potentials for their actual components.\*

In considering this subject, we must first of all inquire how far our method of treating surfaces of discontinuity is applicable to cases in which the radii of curvature of the surfaces are of insensible magnitude. That it should not be applied to such cases without limitation is evident from the consideration that we have neglected the term  $\frac{1}{2}(C_1 - C_2)\delta(c_1 - c_2)$  in equation (494) on account of the magnitude of the radii of curvature compared with the thickness of the non-homogeneous film. (See page 390). When, however, only spherical masses are considered, this term will always disappear, since  $C_1$  and  $C_2$  will necessarily be equal.

Again, the surfaces of discontinuity have been regarded as separating homogeneous masses. But we may easily conceive that a globular mass (surrounded by a large homogeneous mass of different nature) may be so small that no part of it will be homogeneous, and that even at its center the matter cannot be regarded as having any phase of matter *in mass*. This, however, will cause no difficulty, if we regard the phase of the interior mass as determined by the same

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\* See page 161, where the term stable is used (as indicated on page 159) in a less strict sense than in the discussion which here follows.

relations to the exterior mass as in other cases. Beside the phase of the exterior mass, there will always be another phase having the same temperature and potentials, but of the general nature of the small globule which is surrounded by that mass and in equilibrium with it. This phase is completely determined by the system considered, and in general entirely stable and perfectly capable of realization in mass, although not such that the exterior mass could exist in contact with it at a plane surface. This is the phase which we are to attribute to the mass which we conceive as existing within the dividing surface.\*

With this understanding with regard to the phase of the fictitious interior mass, there will be no ambiguity in the meaning of any of the symbols which we have employed, when applied to cases in which the surface of discontinuity is spherical, however small the radius may be. Nor will the demonstration of the general theorems require any material modification. The dividing surface, which determines the value of  $\varepsilon^s$ ,  $\eta^s$ ,  $m_1^s$ ,  $m_2^s$ , etc., is as in other cases to be placed so as to make the term  $\frac{1}{2}(C_1 + C_2)\delta(c_1 + c_2)$  in equation (494) vanish, i. e., so as to make equation (497) valid. It has been shown on pages 387–389 that when thus placed it will sensibly coincide with the physical surface of discontinuity, when this consists of a non-homogeneous film separating homogeneous masses, and having radii of curvature which are large compared with its thickness. But in regard to globular masses too small for this theorem to have any application, it will be worth while to examine how far we may be certain that the radius of the dividing surface will have a real and positive value, since it is only then that our method will have any natural application.

The value of the radius of the dividing surface, supposed spherical, of any globule in equilibrium with a surrounding homogeneous fluid may be most easily obtained by eliminating  $\sigma$  from equations (500) and (502), which have been derived from (497), and contain the radius implicitly. If we write  $r$  for this radius, equation (500) may be written

$$2\sigma = (\rho' - \rho'')r, \quad (550)$$

the single and double accents referring respectively to the interior and exterior masses. If we write  $[\varepsilon]$ ,  $[\eta]$ ,  $[m_1]$ ,  $[m_2]$ , etc. for the

\* For example, in applying our formulae to a microscopic globule of water in steam, by the density or pressure of the interior mass we should understand, not the actual density or pressure at the center of the globule, but the density of liquid water (in large quantities) which has the temperature and potential of the steam.

excess of the total energy, entropy, etc. in and about the globular mass above what would be in the same space if it were uniformly filled with matter of the phase of the exterior mass, we shall have necessarily with reference to the whole dividing surface

$$\varepsilon^s = [\varepsilon] - v' (\varepsilon_v' - \varepsilon_v''), \quad \eta^s = [\eta] - v' (\eta_v' - \eta_v''),$$

$m_1^s = [m_1] - v' (\gamma_1' - \gamma_1'')$ ,  $m_2^s = [m_2] - v' (\gamma_2' - \gamma_2'')$ , etc., where  $\varepsilon_v'$ ,  $\varepsilon_v''$ ,  $\eta_v'$ ,  $\eta_v''$ ,  $\gamma_1'$ ,  $\gamma_1''$ , etc. denote, in accordance with our usage elsewhere, the volume-densities of energy, of entropy, and of the various components, in the two homogeneous masses. We may thus obtain from equation (502).

$$\begin{aligned} \sigma s &= [\varepsilon] - v' (\varepsilon_v' - \varepsilon_v'') - t[\eta] + t v' (\eta_v' - \eta_v'') \\ &\quad - \mu_1 [m_1] + \mu_1 v' (\gamma_1' - \gamma_1'') - \mu_2 [m_2] + \mu_2 v' (\gamma_2' - \gamma_2'') - \text{etc.} \end{aligned} \quad (551)$$

But by (93),

$$\begin{aligned} p' &= -\varepsilon_v' + t \eta_v' + \mu_1 \gamma_1' + \mu_2 \gamma_2' + \text{etc.}, \\ p'' &= -\varepsilon_v'' + t \eta_v'' + \mu_1 \gamma_1'' + \mu_2 \gamma_2'' + \text{etc.} \end{aligned}$$

Let us also write for brevity

$$W = [\varepsilon] - t[\eta] - \mu_1 [m_1] - \mu_2 [m_2] - \text{etc.} \quad (552)$$

(It will be observed that the value of  $W$  is entirely determined by the nature of the physical system considered, and that the notion of the dividing surface does not in any way enter into its definition.) We shall then have

$$\sigma s = W + v' (p' - p''), \quad (553)$$

or, substituting for  $s$  and  $v'$  their values in terms of  $r$ ,

$$4\pi r^2 \sigma = W + \frac{4}{3}\pi r^3 (p' - p''), \quad (554)$$

and eliminating  $\sigma$  by (550),

$$\frac{2}{3}\pi r^3 (p' - p'') = W, \quad (555)$$

$$r = \left( \frac{3W}{2\pi(p' - p'')} \right)^{\frac{1}{3}}. \quad (556)$$

If we eliminate  $r$  instead  $\sigma$ , we have

$$\frac{16\pi\sigma^3}{3(p' - p'')^2} = W, \quad (557)$$

$$\sigma = \left( \frac{3W(p' - p'')^2}{16\pi} \right)^{\frac{1}{3}}. \quad (558)$$

Now, if we first suppose the difference of the pressures in the homogeneous masses to be very small, so that the surface of discontinuity is nearly plane, since without any important loss of generality

we may regard  $\sigma$  as positive (for if  $\sigma$  is not positive when  $p'=p''$ , the surface when plane would not be stable in regard to position, as it certainly is, in every actual case, when the proper conditions are fulfilled with respect to its perimeter), we see by (550) that the pressure in the interior mass must be the greater; i. e., we may regard  $\sigma$ ,  $p'-p''$ , and  $r$  as all positive. By (555), the value of  $W$  will also be positive. But it is evident from equation (552), which defines  $W$ , that the value of this quantity is necessarily real, in any possible case of equilibrium, and can only become infinite when  $r$  becomes infinite and  $p'=p''$ . Hence, by (556) and (558), as  $p'-p''$  increases from very small values,  $W$ ,  $r$ , and  $\sigma$  have single, real, and positive values until they simultaneously reach the value zero. Within this limit, our method is evidently applicable; beyond this limit, if such exist, it will hardly be profitable to seek to interpret the equations. But it must be remembered that the vanishing of the radius of the somewhat arbitrarily determined *dividing surface* may not necessarily involve the vanishing of the physical heterogeneity. It is evident, however, (see pp. 387–389,) that the globule must become insensible in magnitude before  $r$  can vanish.

It may easily be shown that the quantity denoted by  $W$  is the work which would be required to form (by a reversible process) the heterogeneous globule in the interior of a very large mass having initially the uniform phase of the exterior mass. For this work is equal to the increment of energy of the system when the globule is formed without change of the entropy or volume of the whole system or of the quantities of the several components. Now  $[\eta]$ ,  $[m_1]$ ,  $[m_2]$ , etc. denote the increments of entropy and of the components in the space where the globule is formed. Hence these quantities with the negative sign will be equal to the increments of entropy and of the components in the rest of the system. And hence, by equation (86),

$$-t[\eta] - \mu_1[m_1] - \mu_2[m_2] - \text{etc.}$$

will denote the increment of energy in all the system except where the globule is formed. But  $[\varepsilon]$  denotes the increment of energy in that part of the system. Therefore, by (552),  $W$  denotes the total increment of energy in the circumstances supposed, or the work required for the formation of the globule.

The conclusions which may be drawn from these considerations with respect to the stability of the homogeneous mass of the pressure  $p''$  (supposed less than  $p'$ , the pressure belonging to a different phase of the same temperature and potentials) are very obvious.

Within those limits within which the method used has been justified, the mass in question must be regarded as in strictness stable with respect to the growth of a globule of the kind considered, since  $W$ , the work required for the formation of such a globule of a certain size (viz., that which would be in equilibrium with the surrounding mass), will always be positive. Nor can smaller globules be formed, for they can neither be in equilibrium with the surrounding mass, being too small, nor grow to the size of that to which  $W$  relates. If, however, by any external agency such a globular mass (of the size necessary for equilibrium) were formed, the equilibrium has already (page 406) been shown to be unstable, and with the least excess in size, the interior mass would tend to increase without limit except that depending on the magnitude of the exterior mass. We may therefore regard the quantity  $W$  as affording a kind of measure of the *stability* of the phase to which  $p''$  relates. In equation (557) the value of  $W$  is given in terms of  $\sigma$  and  $p'-p''$ . If the three fundamental equations which give  $\sigma$ ,  $p'$ , and  $p''$  in terms of the temperature and the potentials were known, we might regard the stability ( $W$ ) as known in terms of the same variables. It will be observed that when  $p'=p''$  the value of  $W$  is infinite. If  $p'-p''$  increases without greater changes of the phases than are necessary for such increase,  $W$  will vary at first very nearly inversely as the square of  $p'-p''$ . If  $p'-p''$  continues to increase, it may perhaps occur that  $W$  reaches the value zero; but until this occurs the phase is certainly stable with respect to the kind of change considered. Another kind of change is conceivable, which initially is small in degree but may be great in its extent in space. Stability in this respect or *stability in respect to continuous changes of phase* has already been discussed (see page 162), and its limits determined. These limits depend entirely upon the fundamental equation of the homogeneous mass of which the stability is in question. But with respect to the kind of changes here considered, which are initially small in extent but great in degree, it does not appear how we can fix the limits of stability with the same precision. But it is safe to say that if there is such a limit it must be at or beyond the limit at which  $\sigma$  vanishes. This latter limit is determined entirely by the fundamental equation of the surface of discontinuity between the phase of which the stability is in question and that of which the possible formation is in question. We have already seen that when  $\sigma$  vanishes, the radius of the dividing surface and the work  $W$  vanish with it. If the fault in the homogeneity of the mass vanishes at the same time, (it evidently

cannot vanish sooner,) the phase becomes unstable at this limit. But if the fault in the homogeneity of the physical mass does not vanish with  $r$ ,  $\sigma$  and  $W$ ,—and no sufficient reason appears why this should not be considered as the general case,—although the amount of work necessary to upset the equilibrium of the phase is infinitesimal, this is not enough to make the phase unstable. It appears therefore that  $W$  is a somewhat one-sided measure of stability.

It must be remembered in this connection that the fundamental equation of a surface of discontinuity can hardly be regarded as capable of experimental determination, except for plane surfaces, (see pp. 394, 395,) although the relation for spherical surfaces is in the nature of things entirely determined, at least so far as the phases are separately capable of existence. Yet the foregoing discussion yields the following practical results. It has been shown that the real stability of a phase extends in general beyond that limit (discussed on pages 160, 161), which may be called the limit of practical stability, at which the phase can exist in contact with another at a plane surface, and a formula has been deduced to express the degree of stability in such cases as measured by the amount of work necessary to upset the equilibrium of the phase when supposed to extend indefinitely in space. It has also been shown to be entirely consistent with the principles established that this stability should have limits, and the manner in which the general equations would accommodate themselves to this case has been pointed out.

By equation (553), which may be written

$$W = \sigma s - (\rho' - \rho'') v', \quad (559)$$

we see that the work  $W$  consists of two parts, of which one is always positive, and is expressed by the product of the superficial tension and the area of the surface of tension, and the other is always negative, and is numerically equal to the product of the difference of pressure by the volume of the interior mass. We may regard the first part as expressing the work spent in forming the surface of tension, and the second part the work gained in forming the interior mass.\*

\* To make the physical significance of the above more clear, we may suppose the two processes to be performed separately in the following manner. We may suppose a large mass of the same phase as that which has the volume  $v'$  to exist initially in the interior of the other. Of course, it must be surrounded by a resisting envelop, on account of the difference of the pressures. We may, however, suppose this envelop permeable to all the component substances, although not of such properties that a mass can form on the exterior like that within. We may allow the

Moreover, the second of these quantities, if we neglect its sign, is always equal to two-thirds of the first, as appears from equation (550) and the geometrical relation  $v' = \frac{1}{3}rs$ . We may therefore write

$$W = \frac{1}{3} \sigma s = \frac{1}{2} (p' - p'') v'. \quad (560)$$

*On the Possible Formation at the Surface where two different Homogeneous Fluids meet of a Fluid of different Phase from either.*

Let A, B, and C be three different fluid phases of matter, which satisfy all the conditions necessary for equilibrium when they meet at plane surfaces. The components of A and B may be the same or different, but C must have no components except such as belong to A or B. Let us suppose masses of the phases A and B to be separated by a very thin sheet of the phase C. This sheet will not necessarily be plane, but the sum of its principal curvatures must be zero. We may treat such a system as consisting simply of masses of the phases A and B with a certain surface of discontinuity, for in our previous discussion there has been nothing to limit the thickness or the nature of the film separating homogeneous masses, except that its thickness has generally been supposed to be small in comparison with its radii of curvature. The value of the superficial tension for such a film will be  $\sigma_{AC} + \sigma_{BC}$ , if we denote by these symbols the tensions of the surfaces of contact of the phases A and C, and B and C, respectively. This not only appears from evident mechanical considerations, but may also be easily verified by equations (502) and (93), the first of which may be regarded as defining the quantity  $\sigma$ . This value will not be affected by diminishing the thickness of the film, until the

envelop to yield to the internal pressure until its contents are increased by  $v'$  without materially affecting its superficial area. If this be done sufficiently slowly, the phase of the mass within will remain constant. (See page 139.) A homogeneous mass of the volume  $v'$  and of the desired phase has thus been produced, and the work gained is evidently  $(p' - p'')v'$ .

Let us suppose that a small aperture is now opened and closed in the envelop so as to let out exactly the volume  $v'$  of the mass within, the envelop being pressed inwards in another place so as to diminish its contents by this amount. During the extrusion of the drop and until the orifice is entirely closed, the surface of the drop must adhere to the edge of the orifice, but not elsewhere to the outside surface of the envelop. The work done in forming the surface of the drop will evidently be  $\sigma s$  or  $\frac{3}{2}(p' - p'')v'$ . Of this work, the amount  $(p' - p'')v'$  will be expended in pressing the envelop inward, and the rest in opening and closing the orifice. Both the opening and the closing will be resisted by the capillary tension. If the orifice is circular, it must have, when widest open, the radius determined by equation (550).

limit is reached at which the interior of the film ceases to have the properties of matter in mass. Now if  $\sigma_{AC} + \sigma_{BC}$  is greater than  $\sigma_{AB}$ , the tension of the ordinary surface between A and B, such a film will be at least practically unstable. (See page 403.) We cannot suppose that  $\sigma_{AB} > \sigma_{AC} + \sigma_{BC}$ , for this would make the ordinary surface between A and B unstable and difficult to realize. If  $\sigma_{AB} = \sigma_{AC} + \sigma_{BC}$ , we may assume, in general, that this relation is not accidental, and that the ordinary surface of contact for A and B is of the kind which we have described.

Let us now suppose the phases A and B to vary, so as still to satisfy the conditions of equilibrium at plane contact, but so that the pressure of the phase C determined by the temperature and potentials of A and B shall become less than the pressure of A and B. A system consisting of the phases A and B will be entirely stable with respect to the formation of any phase like C. (The case is not quite identical with that considered on page 161, since the system in question contains two different phases, but the principles involved are entirely the same.)

With respect to variations of the phases A and B in the opposite direction we must consider two cases separately. It will be convenient to denote the pressures of the three phases by  $p_A$ ,  $p_B$ ,  $p_C$ , and to regard these quantities as functions of the temperature and potentials.

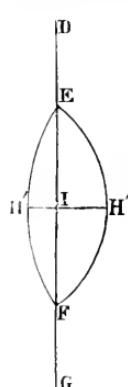
If  $\sigma_{AB} = \sigma_{AC} + \sigma_{BC}$  for values of the temperature and potentials which make  $p_A = p_B = p_C$ , it will not be possible to alter the temperature and potentials at the surface of contact of the phases A and B so that  $p_A \neq p_B$ , and  $p_C > p_A$ , for the relation of the temperature and potentials necessary for the equality of the three pressures will be preserved by the increase of the mass of the phase C. Such variations of the phases A and B might be brought about in separate masses, but if these were brought into contact, there would be an immediate formation of a mass of the phase C, with reduction of the phases of the adjacent masses to such as satisfy the conditions of equilibrium with that phase.

But if  $\sigma_{AB} < \sigma_{AC} + \sigma_{BC}$ , we can vary the temperature and potentials so that  $p_A = p_B$ , and  $p_C > p_A$ , and it will not be possible for a sheet of the phase of C to form *immediately*, i. e., while the pressure of C is sensibly equal to that of A and B; for mechanical work equal to  $\sigma_{AC} + \sigma_{BC} - \sigma_{AB}$  per unit of surface might be obtained by bringing the system into its original condition, and therefore produced without any external expenditure, unless it be that of heat at the temperature of the system, which is evidently incapable of producing the work.

The stability of the system in respect to such a change must therefore extend beyond the point where the pressure of C commences to be less than that of A and B. We arrive at the same result if we use the expression (520) as a test of stability. Since this expression has a finite positive value when the pressures of the phases are all equal, the ordinary surface of discontinuity must be stable, and it must require a finite change in the circumstances of the case to make it become unstable.\*

In the preceding paragraph it is shown that the surface of contact of phases A and B is stable under certain circumstances, with respect to the formation of a thin sheet of the phase C. To complete the demonstration of the stability of the surface with respect to the formation of the phase C, it is necessary to show that this phase cannot be formed at the surface in lentiform masses. This is the more necessary, since it is in this manner, if at all, that the phase is likely to be formed, for an incipient sheet of phase C would evidently be unstable when  $\sigma_{AB} < \sigma_{AC} + \sigma_{BC}$ , and would immediately break up into lentiform masses.

It will be convenient to consider first a lentiform mass of phase C



in equilibrium between masses of phases A and B which meet in a plane surface. Let figure 10 represent a section of such a system through the centers of the spherical surfaces, the mass of phase A lying on the left of D E H' F G, and that of phase B on the right of D E H'' F G. Let the line joining the centers cut the spherical surfaces in II' and II'', and the plane of the surface of contact of A and B in I. Let the radii of E H' F and E H'' F be denoted by  $r'$ ,  $r''$ , and the segments I H', I H'' by  $x'$ ,  $x''$ . Also let I E, the radius of the circle in which the spherical surfaces intersect, be denoted by  $R$ . By a suitable application of the general condition of equilibrium we

may easily obtain the equation

$$\sigma_{AC} \frac{r' - x'}{r'} + \sigma_{BC} \frac{r'' - x''}{r''} = \sigma_{AB}, \quad (561)$$

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\* It is true that such a case as we are now considering is formally excluded in the discussion referred to, which relates to a plane surface, and in which the system is supposed thoroughly stable with respect to the possible formation of any different homogeneous masses. Yet the reader will easily convince himself that the criterion (520) is perfectly valid in this case with respect to the possible formation of a thin sheet of the phase C, which, as we have seen, may be treated simply as a different kind of surface of discontinuity.

which signifies that the components parallel to EF of the tension  $\sigma_{AC}$  and  $\sigma_{BC}$  are together equal to  $\sigma_{AB}$ . If we denote by  $W$  the amount of work which must be expended in order to form such a lentiform mass as we are considering between masses of indefinite extent having the phases A and B, we may write

$$W = M - N, \quad (562)$$

where  $M$  denotes the work expended in replacing the surface between A and B by the surfaces between A and C and B and C, and  $N$  denotes the work gained in replacing the masses of phases A and B by the mass of phase C. Then

$$M = \sigma_{AC} s_{AC} + \sigma_{BC} s_{BC} - \sigma_{AB} s_{AB}, \quad (563)$$

where  $s_{AC}$ ,  $s_{BC}$ ,  $s_{AB}$  denote the areas of the three surfaces concerned; and

$$N = V' (p_c - p_A) + V'' (p_c - p_B), \quad (564)$$

where  $V'$  and  $V''$  denote the volumes of the masses of the phases A and B which are replaced. Now by (500),

$$p_c - p_A = \frac{2\sigma_{AC}}{r'}, \quad \text{and} \quad p_c - p_B = \frac{2\sigma_{BC}}{r''}. \quad (565)$$

We have also the geometrical relations

$$\left. \begin{aligned} V' &= \frac{2}{3} \pi r'^2 x' - \frac{1}{3} \pi R^2 (r' - x'), \\ V'' &= \frac{2}{3} \pi r''^2 x'' - \frac{1}{3} \pi R^2 (r'' - x''). \end{aligned} \right\} \quad (566)$$

By substitution we obtain

$$\begin{aligned} N &= \frac{4}{3} \pi \sigma_{AC} r' x' - \frac{2}{3} \pi R^2 \sigma_{AC} \frac{r' - x'}{r'} \\ &\quad + \frac{4}{3} \pi \sigma_{BC} r'' x'' - \frac{2}{3} \pi R^2 \sigma_{BC} \frac{r'' - x''}{r''}, \end{aligned} \quad (567)$$

and by (561),

$$N = \frac{4}{3} \pi \sigma_{AC} r' x' + \frac{4}{3} \pi \sigma_{BC} r'' x'' - \frac{2}{3} \pi R^2 \sigma_{AB}. \quad (568)$$

Since

$$2\pi r' x' = s_{AC}, \quad 2\pi r'' x'' = s_{BC}, \quad \pi R^2 = s_{AB},$$

we may write

$$N = \frac{2}{3} (\sigma_{AC} s_{AC} + \sigma_{BC} s_{BC} - \sigma_{AB} s_{AB}). \quad (569)$$

(The reader will observe that the ratio of  $M$  and  $N$  is the same as that of the corresponding quantities in the case of the spherical mass treated on pages 416-422.) We have therefore

$$W = \frac{1}{3} (\sigma_{AC} s_{AC} + \sigma_{BC} s_{BC} - \sigma_{AB} s_{AB}). \quad (570)$$

This value is positive so long as

$$\sigma_{AC} + \sigma_{BC} > \sigma_{AB},$$

since  $s_{AC} > s_{AB}$ , and  $s_{BC} > s_{AB}$ .

But at the limit, when

$$\sigma_{AC} + \sigma_{BC} = \sigma_{AB},$$

we see by (561) that

$$s_{AC} = s_{AB}, \text{ and } s_{BC} = s_{AB},$$

and therefore

$$W = 0.$$

It should however be observed that in the immediate vicinity of the circle in which the three surfaces of discontinuity intersect, the physical state of each of these surfaces must be affected by the vicinity of the others. We cannot, therefore, rely upon the formula (570) except when the dimensions of the lentiform mass are of sensible magnitude.

We may conclude that after we pass the limit at which  $p_C$  becomes greater than  $p_A$  and  $p_B$  (supposed equal) lentiform masses of phase C will not be formed until either  $\sigma_{AB} = \sigma_{AC} + \sigma_{BC}$ , or  $p_C - p_A$  becomes so great that the lentiform mass which would be in equilibrium is one of insensible magnitude. [The diminution of the radii with increasing values of  $p_C - p_A$  is indicated by equation (565).] Hence, no mass of phase C will be formed until one of these limits is reached. Although the demonstration relates to a *plane* surface between A and B, the result must be applicable whenever the radii of curvature have a sensible magnitude, since the effect of such curvature may be disregarded when the lentiform mass is of sufficiently small.

The equilibrium of the lentiform mass of phase C is easily proved to be unstable, so that the quantity  $W$  affords a kind of measure of the stability of plane surfaces of contact of the phases A and B.\*

\* If we represent phases by the position of points in such a manner that coexistent phases (in the sense in which the term is used on page 152) are represented by the same point, and allow ourselves, for brevity, to speak of the phases as having the positions of the points by which they are represented, we may say that three coexistent phases are situated where three series of pairs of coexistent phases meet or intersect. If the three phases are all fluid, or when the effects of solidity may be disregarded, two cases are to be distinguished. Either the three series of coexistent phases all intersect,—this is when each of the three surface-tensions is less than the sum of the two others,—or one of the series terminates where the two others intersect,—this is where one surface tension is equal to the sum of the others. The series of coexistent phases will be represented by lines or surfaces, according as the phases have one or two independently variable components. Similar relations exist when the number of components is greater, except that they are not capable of geometrical representation without some limitation, as that of constant temperature or pressure or certain constant potentials.

Essentially the same principles apply to the more general problem in which the phases A and B have moderately different pressures, so that their surfaces of contact must be curved, but the radii of curvature have a sensible magnitude.

In order that a thin film of the phase C may be in equilibrium between masses of the phases A and B, the following equations must be satisfied—

$$\sigma_{AC}(c_1 + c_2) = p_A - p_C,$$

$$\sigma_{BC}(c_1 + c_2) = p_C - p_B,$$

where  $c_1$  and  $c_2$  denote the principal curvatures of the film, the centers of positive curvature lying in the mass having the phase A. Eliminating  $c_1 + c_2$ , we have

$$\sigma_{BC}(p_A - p_C) = \sigma_{AC}(p_C - p_B),$$

or 
$$p_C = \frac{\sigma_{BC}p_A + \sigma_{AC}p_B}{\sigma_{BC} + \sigma_{AC}}. \quad (571)$$

It is evident that if  $p_C$  has a value greater than that determined by this equation, such a film will develop into a larger mass; if  $p_C$  has a less value, such a film will tend to diminish. Hence, when

$$p_C < \frac{\sigma_{BC}p_A + \sigma_{AC}p_B}{\sigma_{BC} + \sigma_{AC}}, \quad (572)$$

the phases A and B have a stable surface of contact.

Again, if more than one kind of surface of discontinuity is possible between A and B, for any given values of the temperature and potentials, it will be impossible for that having the greater tension to displace the other, at the temperature and with the potentials considered. Hence, when  $p_C$  has the value determined by equation (571), and consequently  $\sigma_{AC} + \sigma_{BC}$  is one value of the tension for the surface between A and B, it is impossible that the ordinary tension of the surface  $\sigma_{AB}$  should be greater than this. If  $\sigma_{AB} = \sigma_{AC} + \sigma_{BC}$ , when equation (571) is satisfied, we may presume that a thin film of the phase C actually exists at the surface between A and B, and that a variation of the phases such as would make  $p_C$  greater than the second number of (571) cannot be brought about at that surface, as it would be prevented by the formation of a larger mass of the phase C. But if  $\sigma_{AB} < \sigma_{AC} + \sigma_{BC}$  when equation (571) is satisfied, this equation does not mark the limit of the stability of the surface between A and B, for the temperature or potentials must receive a finite

change before the film of phase C, or (as we shall see in the following paragraph) a lentiform mass of that phase, can be formed.

The work which must be expended in order to form on the surface between indefinitely large masses of phases A and B a lentiform mass of phase C in equilibrium, may evidently be represented by the formula

$$W = \sigma_{AC} S_{AC} + \sigma_{BC} S_{BC} - \sigma_{AB} S_{AB} \\ - p_c V_c + p_A V_A + p_B V_B, \quad (573)$$

where  $S_{AC}$ ,  $S_{BC}$  denote the areas of the surfaces formed between A and C, and B and C,  $S_{AB}$  the diminution of the area of the surface between A and B,  $V_c$  the volume formed of the phase C, and  $V_A$ ,  $V_B$  the diminution of the volumes of the phases A and B. Let us now suppose  $\sigma_{AC}$ ,  $\sigma_{BC}$ ,  $\sigma_{AB}$ ,  $p_A$ ,  $p_B$  to remain constant and the external boundary of the surface between A and B to remain fixed, while  $p_c$  increases and the surfaces of tension receive such alterations as are necessary for equilibrium. It is not necessary that this should be physically possible in the actual system; we may suppose the changes to take place, for the sake of argument, although involving changes in the fundamental equations of the masses and surfaces considered. Then, regarding  $W$  simply as an abbreviation for the second member of the preceding equation, we have

$$dW = \sigma_{AC} dS_{AC} + \sigma_{BC} dS_{BC} - \sigma_{AB} dS_{AB} \\ - p_c dV_c + p_A dV_A + p_B dV_B - V_c dp_c. \quad (574)$$

But the conditions of equilibrium require that

$$\sigma_{AC} dS_{AC} + \sigma_{BC} dS_{BC} - \sigma_{AB} dS_{AB} \\ - p_c dV_c + p_A dV_A + p_B dV_B = 0. \quad (575)$$

Hence,

$$dW = -V_c dp_c. \quad (576)$$

Now it is evident that  $V_c$  will diminish as  $p_c$  increases. Let us integrate the last equation supposing  $p_c$  to increase from its original value until  $V_c$  vanishes. This will give

$$W'' - W' = \text{a negative quantity}, \quad (577)$$

where  $W'$  and  $W''$  denote the initial and final values of  $W$ . But  $W=0$ . Hence  $W'$  is positive. But this is the value of  $W$  in the original system containing the lentiform mass, and expresses the work necessary to form the mass between the phases A and B. It is therefore impossible that such a mass should form on a surface be-

tween these phases. We must however observe the same limitation as in the less general case already discussed,—that  $p_C - p_A$ ,  $p_C - p_B$  must not be so great that the dimensions of the lentiform mass are of insensible magnitude. It may also be observed that the value of these differences may be so small that there will not be room on the surface between the masses of phases A and B for a mass of phase C sufficiently large for equilibrium. In this case we may consider a mass of phase C which is in equilibrium upon the surface between A and B in virtue of a *constraint* applied to the line in which the three surfaces of discontinuity intersect, which will not allow this line to become longer, although not preventing it from becoming shorter. We may prove that the value of  $W$  is positive by such an integration as we have used before.

#### *Substitution of Pressures for Potentials in Fundamental Equations for Surfaces.*

The fundamental equation of a surface which gives the value of the tension in terms of the temperature and potentials seems best adapted to the purposes of theoretical discussion, especially when the number of components is large or undetermined. But the experimental determination of the fundamental equations, or the application of any result indicated by theory to actual cases, will be facilitated by the use of other quantities in place of the potentials, which shall be capable of more direct measurement, and of which the numerical expression (when the necessary measurements have been made) shall depend upon less complex considerations. The numerical value of a potential depends not only upon the system of units employed, but also upon the arbitrary constants involved in the definition of the energy and entropy of the substance to which the potential relates, or, it may be, of the elementary substances of which that substance is formed. (See page 152.) This fact and the want of means of direct measurement may give a certain vagueness to the idea of the potentials, and render the equations which involve them less fitted to give a clear idea of physical relations.

Now the fundamental equation of each of the homogeneous masses which are separated by any surface of discontinuity affords a relation between the pressure in that mass and the temperature and potentials. We are therefore able to eliminate one or two potentials from the fundamental equation of a surface by introducing the pressures in the adjacent masses. Again, when one of these masses is a gas

mixture which satisfies Dalton's law as given on page 215, the potential for each simple gas may be expressed in terms of the temperature and the partial pressure belonging to that gas. By the introduction of these partial pressures we may eliminate as many potentials from the fundamental equation of the surface as there are simple gases in the gas-mixture.

An equation obtained by such substitutions may be regarded as a fundamental equation for the surface of discontinuity to which it relates, for when the fundamental equations of the adjacent masses are known, the equation in question is evidently equivalent to an equation between the tension, temperature, and potentials, and we must regard the knowledge of the properties of the adjacent masses as an indispensable preliminary, or an essential part, of a complete knowledge of any surface of discontinuity. It is evident, however, that from these fundamental equations involving pressures instead of potentials we cannot obtain by differentiation (without the use of the fundamental equations of the homogeneous masses) precisely the same relations as by the differentiation of the equations between the tensions, temperatures, and potentials. It will be interesting to inquire, at least in the more important cases, what relations may be obtained by differentiation from the fundamental equations just described alone.

If there is but one component, the fundamental equations of the two homogeneous masses afford one relation more than is necessary for the elimination of the potential. It may be convenient to regard the tension as a function of the temperature and the difference of the pressures. Now we have by (508) and (98)

$$d\sigma = -\eta_s dt - \Gamma d\mu_1,$$

$$d(p' - p'') = (\eta_v' - \eta_v'') dt + (\gamma' - \gamma'') d\mu_1.$$

Hence we derive the equation

$$d\sigma = -\left(\eta_s - \frac{\Gamma}{\gamma' - \gamma''} (\eta_v' - \eta_v'')\right) dt - \frac{\Gamma}{\gamma' - \gamma''} d(p' - p''), \quad (578)$$

which indicates the differential coefficients of  $\sigma$  with respect to  $t$  and  $p' - p''$ . For surfaces which may be regarded as nearly plane, it is

evident that  $\frac{\Gamma}{\gamma' - \gamma''}$  represents the distance from the surface of tension to a dividing surface located so as to make the superficial density of the single component vanish, (being positive, when the

latter surface is on the side specified by the double accents,) and that the coefficient of  $dt$  (without the negative sign) represents the superficial density of entropy as determined by the latter dividing surface, i. e., the quantity denoted by  $\eta_{s_1}$ , on page 397.

When there are two components, neither of which is confined to the surface of discontinuity, we may regard the tension as a function of the temperature and the pressures in the two homogeneous masses. The values of the differential coefficients of the tension with respect to these variables may be represented in a simple form if we choose such substances for the components that in the particular state considered each mass shall consist of a single component. This will always be possible when the composition of the two masses is not identical, and will evidently not affect the values of the differential coefficients. We then have

$$\begin{aligned} d\sigma &= -\eta_s dt - \Gamma'_s d\mu_i - \Gamma''_s d\mu_u, \\ dp' &= \eta_v' dt + \gamma' d\mu_i, \\ dp'' &= \eta_v'' dt + \gamma'' d\mu_u, \end{aligned}$$

where the marks, and  $_u$  are used instead of the usual  $_1$  and  $_2$  to indicate the identity of the component specified with the substance of the homogeneous masses specified by  $'$  and  $''$ . Eliminating  $d\mu_i$  and  $d\mu_u$  we obtain

$$d\sigma = -\left(\eta_s - \frac{\Gamma'_s}{\gamma'} \eta_v' - \frac{\Gamma''_s}{\gamma''} \eta_v''\right) dt - \frac{\Gamma'_s}{\gamma'} dp' - \frac{\Gamma''_s}{\gamma''} dp''. \quad (579)$$

We may generally neglect the difference of  $p'$  and  $p''$ , and write

$$d\sigma = -\left(\eta_s - \frac{\Gamma'_s}{\gamma'} \eta_v' - \frac{\Gamma''_s}{\gamma''} \eta_v''\right) dt - \left(\frac{\Gamma'_s}{\gamma'} + \frac{\Gamma''_s}{\gamma''}\right) dp. \quad (580)$$

The equation thus modified is strictly to be regarded as the equation for a plane surface. It is evident that  $\frac{\Gamma'_s}{\gamma'}$  and  $\frac{\Gamma''_s}{\gamma''}$  represent the distances from the surface of tension of the two surfaces of which one would make  $\Gamma'_s$  vanish, and the other  $\Gamma''_s$ , that  $\frac{\Gamma'_s}{\gamma'} + \frac{\Gamma''_s}{\gamma''}$  represents the distance between these two surfaces, or the *diminution of volume* due to a unit of the surface of discontinuity, and that the coefficient of  $dt$  (without the negative sign) represents the excess of entropy in a system consisting of a unit of the surface of discontinuity with a part of each of the adjacent masses above that which the same matter would have if it existed in two homogeneous masses of the same phases but without any surface of discontinuity.

(A mass thus existing without any surface of discontinuity must of course be entirely surrounded by matter of the same phase.)\*

The form in which the values of  $\left(\frac{d\sigma}{dt}\right)_p$  and  $\left(\frac{d\sigma}{dp}\right)_t$  are given in equation (580) is adapted to give a clear idea of the relations of these quantities to the particular state of the system for which they are to be determined, but not to show how they vary with the state of the system. For this purpose it will be convenient to have the values of these differential coefficients expressed with reference to ordinary components. Let these be specified as usual by  $\gamma_1$  and  $\gamma_2$ . If we eliminate  $d\mu_1$  and  $d\mu_2$  from the equations

$$\begin{aligned} -d\sigma &= \eta_s dt + \gamma'_1 d\mu_1 + \gamma'_2 d\mu_2, \\ dp &= \eta_v' dt + \gamma'_1 d\mu_1 + \gamma'_2 d\mu_2, \\ dp &= \eta_v'' dt + \gamma''_1 d\mu_1 + \gamma''_2 d\mu_2, \end{aligned}$$


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\* If we set

$$V = -\frac{\Gamma'_1}{\gamma'} - \frac{\Gamma'_2}{\gamma''}, \quad (a)$$

$$H_s = \eta_s - \frac{\Gamma'_1}{\gamma'} \eta_v' - \frac{\Gamma''_2}{\gamma''} \eta_v'', \quad (b)$$

and in like manner

$$E_s = \varepsilon_s - \frac{\Gamma'_1}{\gamma'} \varepsilon_v' - \frac{\Gamma''_2}{\gamma''} \varepsilon_v'', \quad (c)$$

we may easily obtain, by means of equations (93) and (507),

$$E_s = t H_s + \sigma - p V. \quad (d)$$

Now equation (580) may be written

$$d\sigma = -H_s dt + V dp. \quad (e)$$

Differentiating (d), and comparing the result with (e), we obtain

$$dE_s = t dH_s - p dV. \quad (f)$$

The quantities  $E_s$  and  $H_s$  might be called the superficial densities of energy and entropy quite as properly as those which we denote by  $\varepsilon_s$  and  $\eta_s$ . In fact, when the composition of both of the homogeneous masses is invariable, the quantities  $E_s$  and  $H_s$  are much more simple in their definition than  $\varepsilon_s$  and  $\eta_s$ , and would probably be more naturally suggested by the terms *superficial density of energy* and *of entropy*. It would also be natural in this case to regard the quantities of the homogeneous masses as determined by the total quantities of matter, and not by the surface of tension or any other dividing surface. But such a nomenclature and method could not readily be extended so as to treat cases of more than two components with entire generality.

In the treatment of surfaces of discontinuity in this paper, the definitions and nomenclature which have been adopted will be strictly adhered to. The object of this note is to suggest to the reader how a different method might be used in some cases with advantage, and to show the precise relations between the quantities which are used in this paper and others which might be confounded with them, and which may be made more prominent when the subject is treated differently.

we obtain

$$d\sigma = \frac{B}{A} dt + \frac{C}{A} dp, \quad (581)$$

where

$$A = \gamma_1'' \gamma_2' - \gamma_1' \gamma_2'', \quad (582)$$

$$B = \begin{vmatrix} \eta_s & \Gamma_1 & \Gamma_2 \\ \eta_v' & \gamma_1' & \gamma_2' \\ \eta_v'' & \gamma_1'' & \gamma_2'' \end{vmatrix}, \quad (583)$$

$$C = \Gamma_1 (\gamma_2'' - \gamma_2') + \Gamma_2 (\gamma_1' - \gamma_1''). \quad (584)$$

It will be observed that  $A$  vanishes when the composition of the two homogeneous masses is identical, while  $B$  and  $C$  do not, in general, and that the value of  $A$  is negative or positive according as the mass specified by ' contains the component specified by <sub>1</sub> in a greater or less proportion than the other mass. Hence, the values both of  $\left(\frac{d\sigma}{dt}\right)_p$  and of  $\left(\frac{d\sigma}{dp}\right)_t$  become infinite when the difference in the composition of the masses vanishes, and change sign when the greater proportion of a component passes from one mass to the other. This might be inferred from the statements on page 155 respecting coexistent phases which are identical in composition, from which it appears that when two coexistent phases have nearly the same composition, a small variation of the temperature or pressure of the coexistent phases will cause a relatively very great variation in the composition of the phases. The same relations are indicated by the graphical method represented in figure 6 on page 184.

With regard to gas-mixtures which conform to Dalton's law, we shall only consider the fundamental equation for plane surfaces, and shall suppose that there is not more than one component in the liquid which does not appear in the gas-mixture. We have already seen that in limiting the fundamental equation to plane surfaces we can get rid of one potential by choosing such a dividing surface that the superficial density of one of the components vanishes. Let this be done with respect to the component peculiar to the liquid, if such there is; if there is no such component, let it be done with respect to one of the gaseous components. Let the remaining potentials be eliminated by means of the fundamental equations of the simple gases. We may thus obtain an equation between the superficial tension, the temperature, and the several pressures of the simple gases in the gas-mixture or all but one of these pressures. Now, if we eliminate  $d\mu_2$ ,  $d\mu_3$ , etc. from the equations

$$\begin{aligned} d\sigma &= -\eta_{s(1)} dt - \Gamma_{2(1)} d\mu_2 + \Gamma_{3(1)} d\mu_3 + \text{etc.}, \\ dp_2 &= \eta_{v2} dt + \gamma_2 d\mu_2, \\ dp_3 &= \eta_{v3} dt + \gamma_3 d\mu_3, \\ &\text{etc.}, \end{aligned}$$

where the suffix <sub>1</sub> relates to the component of which the surface-density has been made to vanish, and  $\gamma_2, \gamma_3$ , etc. denote the densities of the gases specified in the gas mixture, and  $p_2, p_3$ , etc.,  $\eta_{v2}, \eta_{v3}$ , etc. the pressures and the densities of entropy due to these several gases, we obtain

$$\begin{aligned} d\sigma &= -\left(\eta_{s(1)} - \frac{\Gamma_{2(1)}}{\gamma_2} \eta_{v2} - \frac{\Gamma_{3(1)}}{\gamma_3} \eta_{v3} - \text{etc.}\right) dt \\ &\quad - \frac{\Gamma_{2(1)}}{\gamma_2} dp_2 - \frac{\Gamma_{3(1)}}{\gamma_3} dp_3 - \text{etc.} \quad (585) \end{aligned}$$

This equation affords values of the differential coefficients of  $\sigma$  with respect to  $t, p_2, p_3$ , etc., which may be set equal to those obtained by differentiating the equation between these variables.

#### *Thermal and Mechanical Relations pertaining to the Extension of a Surface of Discontinuity.*

The fundamental equation of a surface of discontinuity with one or two component substances, beside its statical applications, is of use to determine the heat absorbed when the surface is extended under certain conditions.

Let us first consider the case in which there is only a single component substance. We may treat the surface as plane, and place the dividing surface so that the surface density of the single component vanishes. (See page 397.) If we suppose the area of the surface to be increased by unity without change of temperature or of the quantities of liquid and vapor, the entropy of the whole will be increased by  $\eta_{s(1)}$ . Therefore, if we denote by  $Q$  the quantity of heat which must be added to satisfy the conditions, we shall have

$$Q = t \eta_{s(1)}, \quad (586)$$

and by (514),

$$Q = -t \frac{d\sigma}{dt} = -\frac{d\sigma}{d \log t}, \quad (587)$$

It will be observed that the condition of constant quantities of liquid and vapor as determined by the dividing surface which we have adopted is equivalent to the condition that the total volume shall remain constant.

Again, if the surface is extended without application of heat, while the pressure in the liquid and vapor remains constant, the temperature will evidently be maintained constant by condensation of the vapor. If we denote by  $M$  the mass of vapor condensed per unit of surface formed, and by  $\eta_M''$  and  $\eta_M'$  the entropies of the liquid and vapor per unit of mass, the condition of no addition of heat will require that

$$M(\eta_M'' - \eta_M') = \eta_{S(1)}. \quad (588)$$

The increase of the volume of liquid will be

$$\gamma' \frac{\eta_{S(1)}}{(\eta_M'' - \eta_M')}, \quad (589)$$

and the diminution of the volume of vapor

$$\gamma'' \frac{\eta_{S(1)}}{(\eta_M'' - \eta_M')} \cdot \quad (590)$$

Hence, for the work done (per unit of surface formed) by the external bodies which maintain the pressure, we shall have

$$W = \frac{P \eta_{S(1)}}{\eta_M'' - \eta_M'} \left( \frac{1}{\gamma''} - \frac{1}{\gamma'} \right). \quad (591)$$

and, by (514) and (131),

$$W = -P \frac{d\sigma}{dt} \frac{dt}{dp} = -P \frac{d\sigma}{dp} = -\frac{d\sigma}{d \log p}. \quad (592)$$

The work expended directly in extending the film will of course be equal to  $\sigma$ .

Let us now consider the case in which there are two component substances, neither of which is confined to the surface. Since we cannot make the superficial density of both these substances vanish by any dividing surface, it will be best to regard the surface of tension as the dividing surface. We may, however, simplify the formula by choosing such substances for components that each homogeneous mass shall consist of a single component. Quantities relating to these components will be distinguished as on page 431. If the surface is extended until its area is increased by unity, while heat is added at the surface so as to keep the temperature constant, and the pressure of the homogeneous masses is also kept constant, the phase of these masses will necessarily remain unchanged, but the quantity of one will be diminished by  $\Gamma_v'$ , and that of the other by  $\Gamma_v''$ . Their

entropies will therefore be diminished by  $\frac{\Gamma_v'}{\gamma'} \eta_v'$  and  $\frac{\Gamma_v''}{\gamma''} \eta_v''$ , respect-

ively. Hence, since the surface receives the increment of entropy  $\eta_s$ , the total quantity of entropy will be increased by

$$\eta_s = \frac{\Gamma_1}{\gamma'} \eta_{v'} + \frac{\Gamma_2}{\gamma''} \eta_{v''},$$

which by equation (580) is equal to

$$- \left( \frac{d\sigma}{dt} \right)_p.$$

Therefore, for the quantity of heat  $Q$  imparted to the surface, we shall have

$$Q = - t \left( \frac{d\sigma}{dt} \right)_p = - \left( \frac{d\sigma}{d \log t} \right)_p. \quad (593)$$

We must notice the difference between this formula and (587). In (593) the quantity of heat  $Q$  is determined by the condition that the temperature and pressures shall remain constant. In (587) these conditions are equivalent and insufficient to determine the quantity of heat. The additional condition by which  $Q$  is determined may be most simply expressed by saying that the total volume must remain constant. Again, the differential coefficient in (593) is defined by considering  $p$  as constant; in the differential coefficient in (587)  $p$  cannot be considered as constant, and no condition is necessary to give the expression a definite value. Yet, notwithstanding the difference of the two cases, it is quite possible to give a single demonstration which shall be applicable to both. This may be done by considering a cycle of operations after the method employed by Sir William Thomson, who first pointed out these relations.\*

The diminution of volume (per unit of surface formed) will be

$$V = \frac{\Gamma_1}{\gamma'} + \frac{\Gamma_2}{\gamma''} = - \left( \frac{d\sigma}{dp} \right)_t; \quad (594)$$

and the work done (per unit of surface formed) by the external bodies which maintain the pressure constant will be

$$W = - p \left( \frac{d\sigma}{dp} \right)_t = - \left( \frac{d\sigma}{d \log p} \right)_t \quad (595)$$

Compare equation (592).

The values of  $Q$  and  $W$  may also be expressed in terms of quantities relating to the ordinary components. By substitution in (593) and (595) of the values of the differential coefficients which are given by (581), we obtain

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\* See *Proc. Roy. Soc.*, vol. ix, p. 255, (June, 1858); or *Phil. Mag.*, Ser. 4, vol. xvii, p. 61.

$$Q = -t \frac{B}{A}, \quad W = -p \frac{C}{A}. \quad (596)$$

where  $A$ ,  $B$ , and  $C$  represent the expressions indicated by (582)–(584). It will be observed that the values of  $Q$  and  $W$  are in general infinite for the surface of discontinuity between coexistent phases which differ infinitesimally in composition, and change sign with the quantity  $A$ . When the phases are absolutely identical in composition, it is not in general possible to counteract the effect of extension of the surface of discontinuity by any supply of heat. For the matter at the surface will not in general have the same composition as the homogeneous masses, and the matter required for the increased surface cannot be obtained from these masses without altering their phase. The infinite values of  $Q$  and  $W$  are explained by the fact that when the phases are nearly identical in composition, the extension of the surface of discontinuity is accompanied by the vaporization or condensation of a very large mass, according as the liquid or the vapor is the richer in that component which is necessary for the formation of the surface of discontinuity.

If, instead of considering the amount of heat necessary to keep the phases from altering while the surface of discontinuity is extended, we consider the variation of temperature caused by the extension of the surface while the pressures remain constant, it appears that this variation of temperature changes sign with  $\gamma_1''\gamma_2' - \gamma_1'\gamma_2''$ , but vanishes with this quantity, i. e., vanishes when the composition of the phases becomes the same. This may be inferred from the statements on page 155, or from a consideration of the figure on page 184. When the composition of the homogeneous masses is initially absolutely identical, the effect on the temperature of a finite extension or contraction of the surface of discontinuity will be the same,—either of the two will lower or raise the temperature according as the temperature is a maximum or minimum for constant pressure.

The effect of the extension of a surface of discontinuity which is most easily verified by experiment is the effect upon the tension before complete equilibrium has been reestablished throughout the adjacent masses. A fresh surface between coexistent phases may be regarded in this connection as an extreme case of a recently extended surface. When sufficient time has elapsed after the extension of a surface originally in equilibrium between coexistent phases, the superficial tension will evidently have sensibly its original value, unless there are substances at the surface which are either not found

at all in the adjacent masses, or are found only in quantities comparable to those in which they exist at the surface. But a surface newly formed or extended may have a very different tension.

This will not be the case, however, when there is only a single component substance, since all the processes necessary for equilibrium are confined to a film of insensible thickness, and will require no appreciable time for their completion.

When there are two components, neither of which is confined to the surface of discontinuity, the reestablishment of equilibrium after the extension of the surface does not necessitate any processes reaching into the interior of the masses except the transmission of heat between the surface of discontinuity and the interior of the masses. It appears from equation (593) that if the tension of the surface diminishes with a rise of temperature, heat must be supplied to the surface to maintain the temperature uniform when the surface is extended, i. e., the effect of extending the surface is to cool it; but if the tension of any surface increases with the temperature, the effect of extending the surface will be to raise its temperature. In either case, it will be observed, the immediate effect of extending the surface is to increase its tension. A contraction of the surface will of course have the opposite effect. But the time necessary for the reestablishment of sensible thermal equilibrium after extension or contraction of the surface must in most cases be very short.

In regard to the formation or extension of a surface between two coexistent phases of more than two components, there are two extreme cases which it is desirable to notice. When the superficial density of each of the components is exceeding small compared with its density in either of the homogeneous masses, the matter (as well as the heat) necessary for the formation or extension of the normal surface can be taken from the immediate vicinity of the surface without sensibly changing the properties of the masses from which it is taken. But if any one of these superficial densities has a considerable value, while the density of the same component is very small in each of the homogeneous masses, both absolutely and relatively to the densities of the other components, the matter necessary for the formation or extension of the normal surface must come from a considerable distance. Especially if we consider that a small difference of density of such a component in one of the homogeneous masses will probably make a considerable difference in the value of the corresponding potential [see eq. (217)], and that a small difference in the value of the potential will make a considerable difference in the ten-

sion [see eq. (508)], it will be evident that in this case a considerable time will be necessary after the formation of a fresh surface or the extension of an old one for the reëstablishment of the normal value of the superficial tension. In intermediate cases, the reëstablishment of the normal tension will take place with different degrees of rapidity.

But whatever the number of component substances, provided that it is greater than one, and whether the reëstablishment of equilibrium is slow or rapid, extension of the surface will generally produce increase and contraction decrease of the tension. It would evidently be inconsistent with stability that the opposite effects should be produced. In general, therefore, a fresh surface between coexistent phases has a greater tension than an old one.\* By the use of fresh surfaces, in experiments in capillarity, we may sometimes avoid the effect of minute quantities of foreign substances, which may be present without our knowledge or desire, in the fluids which meet at the surface investigated.

When the establishment of equilibrium is rapid, the variation of the tension from its normal value will be manifested especially during the extension or contraction of the surface, the phenomenon resembling that of viscosity, except that the variations of tension arising from variations in the densities at and about the surface will be the same in all directions, while the variations of tension due to any property of the surface really analogous to viscosity would be greatest in the direction of the most rapid extension.

We may here notice the different action of traees in the homogeneous masses of those substances which increase the tension and of those which diminish it. When the volume-densities of a component are very small, its surface-density may have a considerable positive value, but can only have a very minute negative one.† For the value when negative cannot exceed (numerically) the product of the greater volume-density by the thickness of the non-homogeneous

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\* When, however, homogeneous masses which have not coexistent phases are brought into contact, the superficial tension may increase with the course of time. The superficial tension of a drop of alcohol and water placed in a large room will increase as the potential for alcohol is equalized throughout the room, and is diminished in the vicinity of the surface of discontinuity.

† It is here supposed that we have chosen for components such substances as are incapable of resolution into other components which are independently variable in the homogeneous masses. In a mixture of alcohol and water, for example, the components must be pure alcohol and pure water.

film. Each of these quantities is exceedingly small. The surface-density when positive is of the same order of magnitude as the thickness of the non-homogeneous film, but is not necessarily small compared with other surface-densities because the volume-densities of the same substance in the adjacent masses are small. Now the potential of a substance which forms a very small part of a homogeneous mass certainly increases, and probably very rapidly, as the proportion of that component is increased. [See (171) and (217).] The pressure, temperature, and the other potentials, will not be sensibly affected. [See (98).] But the effect on the tension of this increase of the potential will be proportional to the surface-density, and will be to diminish the tension when the surface-density is positive. [See (508).] It is therefore quite possible that a very small trace of a substance in the homogeneous masses should greatly diminish the tension, but not possible that such a trace should greatly increase it.\*

#### *Impermeable Films.*

We have so far supposed, in treating of surfaces of discontinuity, that they afford no obstacle to the passage of any of the component substances from either of the homogeneous masses to the other. The case, however, must be considered, in which there is a film of matter at the surface of discontinuity which is impermeable to some or all of

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\* From the experiments of M. E. Duclaux, (*Annales de Chimie et de Physique*, Ser. 4, vol. xxi, p. 383,) it appears that one per cent. of alcohol in water will diminish the superficial tension to .933, the value for pure water being unity. The experiments do not extend to pure alcohol, but the difference of the tensions for mixtures of alcohol and water containing 10 and 20 per cent. water is comparatively small, the tensions being .322 and .336 respectively.

According to the same authority (page 427 of the volume cited), one 3200th part of Castile soap will reduce the superficial tension of water by one-fourth; one 800th part of soap by one-half. These determinations, as well as those relating to alcohol and water, are made by the method of drops, the weight of the drops of different liquids (from the same pipette) being regarded as proportional to their superficial tensions.

M. Athanase Dupré has determined the superficial tensions of solutions of soap by different methods. A statical method gives for one part of common soap in 5000 of water a superficial tension about one-half as great as for pure water, but if the tension be measured on a jet close to the orifice, the value (for the same solution) is sensibly identical with that of pure water. He explains these different values of the superficial tension of the same solution as well as the great effect on the superficial tension which a very small quantity of soap or other trifling impurity may produce, by the tendency of the soap or other substance to form a film on the surface of the liquid. (See *Annales de Chimie et de Physique*, Ser. 4, vol. vii, p. 409, and vol. ix, p. 379.)

the components of the contiguous masses. Such may be the case, for example, when a film of oil is spread on a surface of water, even when the film is too thin to exhibit the properties of the oil in mass. In such cases, if there is communication between the contiguous masses through other parts of the system to which they belong, such that the components in question can pass freely from one mass to the other, the impossibility of a direct passage through the film may be regarded as an immaterial circumstance, so far as states of equilibrium are concerned, and our formulae will require no change. But when there is no such indirect communication, the potential for any component for which the film is impermeable may have entirely different values on opposite sides of the film, and the case evidently requires a modification of our usual method.

A single consideration will suggest the proper treatment of such cases. If a certain component which is found on both sides of a film cannot pass from either side to the other, the fact that the part of the component which is on one side is the same kind of matter with the part on the other side may be disregarded. All the general relations must hold true, which would hold if they were really different substances. We may therefore write  $\mu_1$  for the potential of the component on one side of the film, and  $\mu_2$  for the potential of the same substance (to be treated as if it were a different substance) on the other side;  $m_1^s$  for the excess of the quantity of the substance on the first side of the film above the quantity which would be on that side of the dividing surface (whether this is determined by the surface of tension or otherwise) if the density of the substance were the same near the dividing surface as at a distance, and  $m_2^s$  for a similar quantity relating to the other side of the film and dividing surface. On the same principle, we may use  $\Gamma_1$  and  $\Gamma_2$  to denote the values of  $m_1^s$  and  $m_2^s$  per unit of surface, and  $m_1', m_2'', \gamma_1', \gamma_2''$  to denote the quantities of the substance and its densities in the two homogeneous masses.

With such a notation, which may be extended to cases in which the film is impermeable to any number of components, the equations relating to the surface and the contiguous masses will evidently have the same form as if the substances specified by the different suffixes were all really different. The superficial tension will be a function of  $\mu_1$  and  $\mu_2$ , with the temperature and the potentials for the other components, and  $-\Gamma_1$ ,  $-\Gamma_2$  will be equal to its differential coefficients with respect to  $\mu_1$  and  $\mu_2$ . In a word, all the general relations which have been demonstrated may be applied to this case, if

we remember always to treat the component as a different substance according as it is found on one side or the other of the impermeable film.

When there is free passage for the component specified by the suffixes  $_1$  and  $_2$  through other parts of the system, (or through any flaws in the film,) we shall have in case of equilibrium  $\mu_1 = \mu_2$ . If we wish to obtain the fundamental equation for the surface when satisfying this condition, without reference to other possible states of the surface, we may set a single symbol for  $\mu_1$  and  $\mu_2$  in the more general form of the fundamental equation. Cases may occur of an impermeability which is not absolute, but which renders the transmission of some of the components exceedingly slow. In such cases, it may be necessary to distinguish at least two different fundamental equations, one relating to a state of approximate equilibrium which may be quickly established, and another relating to the ultimate state of complete equilibrium. The former may be derived from the latter by such substitutions as that just indicated.

*The Conditions of Internal Equilibrium for a System of Heterogeneous Fluid Masses without neglect of the Influence of the Surfaces of Discontinuity or of Gravity.*

Let us now seek the complete value of the variation of the energy of a system of heterogeneous fluid masses, in which the influence of gravity and of the surfaces of discontinuity shall be included, and deduce from it the conditions of internal equilibrium for such a system. In accordance with the method which has been developed, the intrinsic energy, (*i. e.*, the part of the energy which is independent of gravity,) the entropy, and the quantities of the several components must each be divided into two parts, one of which we regard as belonging to the surfaces which divide approximately homogeneous masses, and the other as belonging to these masses. The elements of intrinsic energy, entropy, etc., relating to an element of surface  $Ds$  will be denoted by  $D\varepsilon^s$ ,  $D\eta^s$ ,  $Dm_1^s$ ,  $Dm_2^s$ , etc., and those relating to an element of volume  $Dv$ , by  $D\varepsilon^v$ ,  $D\eta^v$ ,  $Dm_1^v$ ,  $Dm_2^v$ , etc. We shall also use  $Dm^s$  or  $\Gamma Ds$  and  $Dm^v$  or  $\gamma Dv$  to denote the total quantities of matter relating to the elements  $Ds$  and  $Dv$  respectively. That is,

$$Dm^s = \Gamma Ds = Dm_1^s + Dm_2^s + \text{etc.}, \quad (597)$$

$$Dm^v = \gamma Dv = Dm_1^v + Dm_2^v + \text{etc.} \quad (598)$$

The part of the energy which is due to gravity must also be divided

into two parts, one of which relates to the elements  $Dm^s$ , and the other to the elements  $Dm^v$ . The complete value of the variation of the energy of the system will be represented by the expression

$$\delta \int D\varepsilon^v + \delta \int D\varepsilon^s + \delta \int g z Dm^v + \delta \int g z Dm^s, \quad (599)$$

in which  $g$  denotes the force of gravity, and  $z$  the height of the element above a fixed horizontal plane.

It will be convenient to limit ourselves at first to the consideration of reversible variations. This will exclude the formation of new masses or surfaces. We may therefore regard any infinitesimal variation in the state of the system as consisting of infinitesimal variations of the quantities relating to its several elements, and bring the sign of variation in the preceding formula after the sign of integration. If we then substitute for  $\delta D\varepsilon^v$ ,  $\delta D\varepsilon^s$ ,  $\delta Dm^v$ ,  $\delta Dm^s$ , the values given by equations (13), (497), (597), (598), we shall have for the condition of equilibrium with respect to reversible variations of the internal state of the system

$$\begin{aligned} & \int t \delta D\eta^v - \int p \delta Dv + \int \mu_1 \delta Dm_1^v + \int \mu_2 \delta Dm_2^v + \text{etc.} \\ & + \int t \delta D\eta^s + \int \sigma \delta Ds + \int \mu_1 \delta Dm_1^s + \int \mu_2 \delta Dm_2^s + \text{etc.} \\ & + \int g \delta z Dm^v + \int g z \delta Dm_1^v + \int g z \delta Dm_2^v + \text{etc.} \\ & + \int g \delta z Dm^s + \int g z \delta Dm_1^s + \int g z \delta Dm_2^s + \text{etc.} = 0, \end{aligned} \quad (600)$$

Since equation (497) relates to surfaces of discontinuity which are initially in equilibrium, it might seem that this condition, although always necessary for equilibrium, may not always be sufficient. It is evident, however, from the form of the condition, that it includes the particular conditions of equilibrium relating to every possible deformation of the system, or reversible variation in the distribution of entropy or of the several components. It therefore includes all the relations between the different parts of the system which are necessary for equilibrium, so far as reversible variations are concerned. (The necessary relations between the various quantities relating to each element of the masses and surfaces are expressed by the fundamental equation of the mass or surface concerned, or may be immediately derived from it. See pp. 140–144 and 391–393.)

The variations in (600) are subject to the conditions which arise from the nature of the system and from the supposition that the changes in the system are not such as to affect external bodies. This supposition is necessary, unless we are to consider the variations in the state of the external bodies, and is evidently allowable in seeking the conditions of equilibrium which relate to the interior of the sys-

tem.\* But before we consider the equations of condition in detail, we may divide the condition of equilibrium (600) into the three conditions

$$\int t \delta D\eta^v + \int t \delta D\eta^s = 0, \quad (601)$$

$$-\int p \delta Dv + \int \sigma \delta Ds + \int g \delta z Dm^v + \int g \delta z Dm^s = 0, \quad (602)$$

$$\begin{aligned} & \int \mu_1 \delta Dm_1^v + \int \mu_1 \delta Dm_1^s + \int g z \delta Dm_1^v + \int g z \delta Dm_1^s \\ & + \int \mu_2 \delta Dm_2^v + \int \mu_2 \delta Dm_2^s + \int g z \delta Dm_2^v + \int g z \delta Dm_2^s \\ & + \text{etc.} = 0. \end{aligned} \quad (603)$$

For the variations which occur in any one of the three are evidently independent of those which occur in the other two, and the equations of condition will relate to one or another of these conditions separately.

The variations in condition (601) are subject to the condition that the entropy of the whole system shall remain constant. This may be expressed by the equation

$$\int \delta D\eta^v + \int \delta D\eta^s = 0. \quad (604)$$

To satisfy the condition thus limited it is necessary and sufficient that

$$t = \text{const.} \quad (605)$$

throughout the whole system, which is the condition of thermal equilibrium.

The conditions of mechanical equilibrium, or those that relate to the possible deformation of the system, are contained in (602), which may also be written

$$-\int p \delta Dr + \int \sigma \delta Ds + \int g \gamma \delta z Dv + \int g \Gamma \delta z Ds = 0. \quad (606)$$

It will be observed that this condition has the same form as if the different fluids were separated by heavy and elastic membranes without rigidity and having at every point a tension uniform in all directions in the plane of the surface. The variations in this formula,

\* We have sometimes given a physical expression to a supposition of this kind, in problems in which the peculiar condition of matter in the vicinity of surfaces of discontinuity was to be neglected, by regarding the system as surrounded by a rigid and impermeable envelop. But the more exact treatment which we are now to give the problem of equilibrium would require us to take account of the influence of the envelop on the immediately adjacent matter. Since this involves the consideration of surfaces of discontinuity between solids and fluids, and we wish to limit ourselves at present to the consideration of the equilibrium of fluid masses, we shall give up the conception of an impermeable envelop, and regard the system as bounded simply by a imaginary surface, which is not a surface of discontinuity. The variations of the system must be such as do not deform the surface, nor affect the matter external to it.

beside their necessary geometrical relations, are subject to the conditions that the external surface of the system, and the lines in which the surfaces of discontinuity meet it, are fixed. The formula may be reduced by any of the usual methods, so as to give the particular conditions of mechanical equilibrium. Perhaps the following method will lead as directly as any to the desired result.

It will be observed the quantities affected by  $\delta$  in (606) relate exclusively to the position and size of the elements of volume and surface into which the system is divided, and that the variations  $\delta p$  and  $\delta\sigma$  do not enter into the formula either explicitly or implicitly. The equations of condition which concern this formula also relate exclusively to the variations of the system of geometrical elements, and do not contain either  $\delta p$  or  $\delta\sigma$ . Hence, in determining whether the first member of the formula has the value zero for every possible variation of the system of geometrical elements, we may assign to  $\delta p$  and  $\delta\sigma$  any values whatever, which may simplify the solution of the problem, without inquiring whether such values are physically possible.

Now when the system is in its initial state, the pressure  $p$ , in each of the parts into which the system is divided by the surfaces of tension, is a function of the co-ordinates which determine the position of the element  $Dv$ , to which the pressure relates. In the varied state of the system, the element  $Dv$  will in general have a different position. Let the variation  $\delta p$  be determined solely by the change in position of the element  $Dv$ . This may be expressed by the equation

$$\delta p = \frac{dp}{dx} \delta x + \frac{dp}{dy} \delta y + \frac{dp}{dz} \delta z, \quad (607)$$

in which  $\frac{dp}{dx}$ ,  $\frac{dp}{dy}$ ,  $\frac{dp}{dz}$  are determined by the function mentioned, and  $\delta x$ ,  $\delta y$ ,  $\delta z$  by the variation of the position of the element  $Dv$ .

Again, in the initial state of the system the tension  $\sigma$ , in each of the different surfaces of discontinuity, is a function of two co-ordinates,  $\omega_1$ ,  $\omega_2$ , which determine the position of the element  $Ds$ . In the varied state of the system, this element will in general have a different position. The change of position may be resolved into a component lying in the surface and another normal to it. Let the variation  $\delta\sigma$  be determined solely by the first of these components of the motion of  $Ds$ . This may be expressed by the equation

$$\delta\sigma = \frac{d\sigma}{d\omega_1} \delta\omega_1 + \frac{d\sigma}{d\omega_2} \delta\omega_2, \quad (608)$$

in which  $\frac{d\sigma}{d\omega_1}$ ,  $\frac{d\sigma}{d\omega_2}$  are determined by the function mentioned, and  $\delta\omega_1$ ,  $\delta\omega_2$ , by the component of the motion of  $Ds$  which lies in the plane of the surface.

With this understanding, which is also to apply to  $\delta p$  and  $\delta\sigma$  when contained implicitly in any expression, we shall proceed to the reduction of the condition (606).

With respect to any one of the volumes into which the system is divided by the surfaces of discontinuity, we may write

$$\int p \delta Dv = \delta \int p Dv - \int \delta p Dv.$$

But it is evident that

$$\delta \int p Dv = \int p \delta N Ds,$$

where the second integral relates to the surfaces of discontinuity bounding the volume considered, and  $\delta N$  denotes the normal component of the motion of an element of the surface, measured outward. Hence,

$$\int p \delta Dv = \int p \delta N Ds - \int \delta p Dv.$$

Since this equation is true of each separate volume into which the system is divided, we may write for the whole system

$$\int p \delta Dv = \int (p' - p'') \delta N Ds - \int \delta p Dv, \quad (609)$$

where  $p'$  and  $p''$  denote the pressures on opposite sides of the element  $Ds$ , and  $\delta N$  is measured toward the side specified by double accents.

Again, for each of the surfaces of discontinuity, taken separately,

$$\int \sigma \delta Ds = \delta \int \sigma Ds - \int \delta \sigma Ds,$$

and

$$\delta \int \sigma Ds = \int \sigma (c_1 + c_2) \delta N Ds + \int \sigma \delta T Dl,$$

where  $c_1$  and  $c_2$  denote the principal curvatures of the surface, (positive, when the centers are on the side opposite to that toward which  $\delta N$  is measured,)  $Dl$  an element of the perimeter of the surface, and  $\delta T$  the component of the motion of this element which lies in the plane of the surface and is perpendicular to the perimeter, (positive, when it extends the surface). Hence we have for the whole system

$$\int \sigma \delta Ds = \int \sigma (c_1 + c_2) \delta N Ds + \int \Sigma (\sigma \delta T) Dl - \int \delta \sigma Ds, \quad (610)$$

where the integration of the elements  $Dl$  extends to all the lines in which the surfaces of discontinuity meet, and the symbol  $\Sigma$  denotes a summation with respect to the several surfaces which meet in such a line.

By equations (609) and (610), the general condition of mechanical equilibrium is reduced to the form

$$-\int(p' - p'') \delta N Ds + \int \delta p Dv + \int \sigma(c_1 + c_2) \delta N Ds \\ + \int \Sigma(\sigma \delta T) Dl - \int \delta \sigma Ds + \int g \gamma \delta z Dv + \int g l \delta z Ds = 0.$$

Arranging and combining terms, we have

$$\int(g \gamma \delta z + \delta p) Dv \\ + \int[(p'' - p') \delta N + \sigma(c_1 + c_2) \delta N + g l \delta z - \delta \sigma] Ds \\ + \int \Sigma(\sigma \delta T) Dl = 0. \quad (611)$$

To satisfy this condition, it is evidently necessary that the coefficients of  $Dv$ ,  $Ds$ , and  $Dl$  shall vanish throughout the system.

In order that the coefficient of  $Dv$  shall vanish, it is necessary and sufficient that, in each of the masses into which the system is divided by the surfaces of tension,  $p$  shall be a function of  $z$  alone, such that

$$\frac{dp}{dz} = -g \gamma. \quad (612)$$

In order that the coefficient of  $Ds$  shall vanish in all cases, it is necessary and sufficient that it shall vanish for normal and for tangential movements of the surface. For normal movements we may write

$$\delta \sigma = 0, \quad \text{and} \quad \delta z = \cos \vartheta \delta N,$$

where  $\vartheta$  denotes the angle which the normal makes with a vertical line. The first condition therefore gives the equation

$$p' - p'' = \sigma(c_1 + c_2) + g l \cos \vartheta, \quad (613)$$

which must hold true at every point in every surface of discontinuity. The condition with respect to tangential movements shows that in each surface of tension  $\sigma$  is a function of  $z$  alone, such that

$$\frac{d\sigma}{dz} = g l. \quad (614)$$

In order that the coefficient of  $Dl$  in (611) shall vanish, we must have, for every point in every line in which surfaces of discontinuity meet, and for any infinitesimal displacement of the line,

$$\Sigma(\sigma \delta T) = 0. \quad (615)$$

This condition evidently expresses the same relations between the tensions of the surfaces meeting in the line and the directions of perpendiculars to the line drawn in the planes of the various surfaces, which hold for the magnitudes and directions of forces in equilibrium in a plane.

In condition (603), the variations which relate to any component are to be regarded as having the value zero in any part of the system in

which that substance is not an actual component.\* The same is true with respect to the equations of condition, which are of the form

$$\left. \begin{array}{l} f' \delta Dm_1^v + f' \delta Dm_1^s = 0, \\ f' \delta Dm_2^v + f' \delta Dm_2^s = 0, \\ \text{etc.} \end{array} \right\} \quad (616)$$

(It is here supposed that the various components are independent, *i.e.*, that none can be formed out of others, and that the parts of the system in which any component actually occurs are not entirely separated by parts in which it does not occur.) To satisfy the condition (603), subject to these equations of condition, it is necessary and sufficient that the conditions

$$\left. \begin{array}{l} \mu_1 + gz = M_1, \\ \mu_2 + gz = M_2, \\ \text{etc.,} \end{array} \right\} \quad (617)$$

( $M_1$ ,  $M_2$ , etc. denoting constants,) shall each hold true in those parts of the system in which the substance specified is an actual component. We may here add the condition of equilibrium relative to the possible absorption of any substance (to be specified by the suffix  $\alpha$ ) by parts of the system of which it is not an actual component, *viz.*, that the expression  $\mu_\alpha + gz$  must not have a less value in such parts of the system than in a contiguous part in which the substance is an actual component.

From equation (613) with (605) and (617) we may easily obtain the differential equation of a surface of tension (in the geometrical sense of the term), when  $p'$ ,  $p''$ , and  $\sigma$  are known in terms of the temperature and potentials. For  $c_1 + c_2$  and  $\mathfrak{L}$  may be expressed in terms of the first and second differential coefficients of  $z$  with respect to the horizontal co-ordinates, and  $p'$ ,  $p''$ ,  $\sigma$ , and  $\Gamma$  in terms of the temperature and potentials. But the temperature is constant, and for each of the potentials we may substitute  $-gz$  increased by a constant. We thus obtain an equation in which the only variables are  $z$  and its first and second differential coefficients with respect to the horizontal co-ordinates. But it will rarely be necessary to use so exact a method. Within moderate differences of level, we may regard  $p'$ ,  $p''$ , and  $\sigma$  as constant. We may then integrate the equation [derived from (612)]

$$d(p' - p'') = g(\gamma'' - \gamma') dz,$$

\* The term *actual component* has been defined for homogeneous masses on page 117, and the definition may be extended to surfaces of discontinuity. It will be observed that if a substance is an actual component of either of the masses separated by a surface of discontinuity, it must be regarded as an actual component for that surface, as well as when it occurs at the surface but not in either of the contiguous masses.

which will give

$$p' - p'' = g(\gamma'' - \gamma')z, \quad (618)$$

where  $z$  is to be measured from the horizontal plane for which  $p' = p''$ . Substituting this value in (613), and neglecting the term containing  $T$ , we have

$$c_1 + c_2 = \frac{g(\gamma'' - \gamma')}{\sigma} z, \quad (619)$$

where the coefficient of  $z$  is to be regarded as constant. Now the value of  $z$  cannot be very large, in any surface of sensible dimensions, unless  $\gamma'' - \gamma'$  is very small. We may therefore consider this equation as practically exact, unless the densities of the contiguous masses are very nearly equal. If we substitute for the sum of the curvatures its value in terms of the differential coefficients of  $z$  with respect to the horizontal rectangular co-ordinates,  $x$  and  $y$ , we have

$$\frac{\left(1 + \frac{dz^2}{dy^2}\right) \frac{d^2z}{dx^2} - 2 \frac{dz}{dx} \frac{dz}{dy} \frac{d^2z}{dxdy} + \left(1 + \frac{dz^2}{dx^2}\right) \frac{d^2z}{dy^2}}{\left(1 + \frac{dz^2}{dx^2} + \frac{dz^2}{dy^2}\right)^{\frac{3}{2}}} = \frac{g(\gamma'' - \gamma')}{\sigma} z. \quad (620)$$

With regard to the sign of the root in the denominator of the fraction, it is to be observed that, if we always take the positive value of the root, the value of the whole fraction will be positive or negative according as the greater concavity is turned upward or downward. But we wish the value of the fraction to be positive when the greater concavity is turned toward the mass specified by a single accent. We should therefore take the positive or negative value of the root according as this mass is above or below the surface.

The particular conditions of equilibrium which are given in the last paragraph but one may be regarded in general as the conditions of chemical equilibrium between the different parts of the system, since they relate to the separate components.\* But such a designation is not entirely appropriate unless the number of components is greater than one. In no case are the conditions of mechanical equilibrium entirely independent of those which relate to temperature and the potentials. For the conditions (612) and (614) may be regarded as consequences of (605) and (617) in virtue of the necessary relations (98) and (508).†

\* Concerning another kind of conditions of chemical equilibrium, which relate to the molecular arrangement of the components, and not to their sensible distribution in space, see pages 197–203.

† Compare page 206, where a similar problem is treated without regard to the influence of the surfaces of discontinuity.

The mechanical conditions of equilibrium, however, have an especial importance, since we may always regard them as satisfied in any liquid (and not decidedly viscous) mass in which no sensible motions are observable. In such a mass, when isolated, the attainment of mechanical equilibrium will take place very soon; thermal and chemical equilibrium will follow more slowly. The thermal equilibrium will generally require less time for its approximate attainment than the chemical; but the processes by which the latter is produced will generally cause certain inequalities of temperature until a state of complete equilibrium is reached.

When a surface of discontinuity has more components than one which do not occur in the contiguous masses, the adjustment of the potentials for these components in accordance with equations (617) may take place very slowly, or not at all, for want of sufficient mobility in the components of the surface. But when this surface has only one component which does not occur in the contiguous masses, and the temperature and potentials in these masses satisfy the conditions of equilibrium, the potential for the component peculiar to the surface will very quickly conform to the law expressed in (617), since this is a necessary consequence of the condition of mechanical equilibrium (614) in connection with the conditions relating to temperature and the potentials which we have supposed to be satisfied. The necessary distribution of the substance peculiar to the surface will be brought about by expansions and contractions of the surface. If the surface meets a third mass containing this component and no other which is foreign to the masses divided by the surface, the potential for this component in the surface will of course be determined by that in the mass which it meets.

The particular conditions of mechanical equilibrium (612)–(615), which may be regarded as expressing the relations which must subsist between contiguous portions of a fluid system in a state of mechanical equilibrium, are serviceable in determining whether a given system is or is not in such a state. But the mechanical theorems which relate to finite parts of the system, although they may be deduced from these conditions by integration, may generally be more easily obtained by a suitable application of the general condition of mechanical equilibrium (606), or by the application of ordinary mechanical principles to the system regarded as subject to the forces indicated by this equation.

It will be observed that the conditions of equilibrium relating to temperature and the potentials are not affected by the surfaces of

discontinuity. [Compare (228) and (234).] \* Since a phase cannot vary continuously without variations of the temperature or the potentials, it follows from these conditions that the phase at any point in a fluid system which has the same independently variable components throughout, and is in equilibrium under the influence of gravity, must be one of a certain number of phases which are completely determined by the phase at any given point and the difference of level of the two points considered. If the phases throughout the fluid system satisfy the general condition of practical stability for phases existing in large masses (viz., that the pressure shall be the least consistent with the temperature and potentials), they will be entirely determined by the phase at any given point and the differences of level. (Compare page 210, where the subject is treated without regard to the influence of the surfaces of discontinuity.)

*Conditions of equilibrium relating to irreversible changes.*—The conditions of equilibrium relating to the absorption by any part of the system of substances which are not actual components of that part have been given on page 448. Those relating to the formation of new masses and surfaces are included in the conditions of stability relating to such changes, and are not always distinguishable from them. They are evidently independent of the action of gravity. We have already discussed the conditions of stability with respect to the formation of new fluid masses within a homogeneous fluid and at the surface when two such masses meet (see pages 416–429), as well as the condition relating to the possibility of a change in the nature of a surface of discontinuity. (See pages 400–403, where the surface considered is plane, but the result may easily be extended to curved surfaces.) We shall hereafter consider, in some of the more important cases, the conditions of stability with respect to the formation of new masses and surfaces which are peculiar to lines in which several surfaces of discontinuity meet, and points in which several such lines meet.

*Conditions of stability relating to the whole system.*—Beside the conditions of stability relating to very small parts of a system, which are substantially independent of the action of gravity, and are discussed elsewhere, there are other conditions, which relate to the

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\* If the fluid system is divided into separate masses by solid diaphragms which are permeable to all the components of the fluids independently, the conditions of equilibrium of the fluids relating to temperature and the potentials will not be affected. (Compare page 139.) The propositions which follow in the above paragraph may be extended to this case.

whole system or to considerable parts of it. To determine the question of the stability of a given fluid system under the influence of gravity, when all the conditions of equilibrium are satisfied as well as those conditions of stability which relate to small parts of the system taken separately, we may use the method described on page 413, the demonstration of which (pages 411, 412) will not require any essential modification on account of gravity.

When the variations of temperature and of the quantities  $M_1$ ,  $M_2$ , etc. [see (617)] involved in the changes considered are so small that they may be neglected, the condition of stability takes a very simple form, as we have already seen to be the case with respect to a system uninfluenced by gravity. (See page 415.)

We have to consider a varied state of the system in which the total entropy and the total quantities of the various components are unchanged, and all variations vanish at the exterior of the system,—in which, moreover, the conditions of equilibrium relating to temperature and the potentials are satisfied, and the relations expressed by the fundamental equations of the masses and surfaces are to be regarded as satisfied, although the state of the system is not one of complete equilibrium. Let us imagine the state of the system to vary continuously in the course of time in accordance with these conditions and use the symbol  $d$  to denote the simultaneous changes which take place at any instant. If we denote the total energy of the system by  $E$ , the value of  $dE$  may be expanded like that of  $\delta E$  in (599) and (600), and then reduced (since the values of  $t$ ,  $\mu_1 + gz$ ,  $\mu_2 + gz$ , etc. are uniform throughout the system, and the total entropy and total quantities of the several components are constant) to the form

$$\begin{aligned} dE &= -\int p \, dDv + \int g \, dz \, Dm^v + \int \sigma \, dDs + \int g \, dz \, Dm^s \\ &= -\int p \, dDv + \int g \gamma \, dz \, Dv + \int \sigma \, dDs + \int g \Gamma \, dz \, Ds, \end{aligned} \quad (621)$$

where the integrations relate to the elements expressed by the symbol  $D$ . The value of  $p$  at any point in any of the various masses, and that of  $\sigma$  at any point in any of the various surfaces of discontinuity are entirely determined by the temperature and potentials at the point considered. If the variations of  $t$  and  $M_1$ ,  $M_2$ , etc. are to be neglected, the variations of  $p$  and  $\sigma$  will be determined solely by the change in position of the point considered. Therefore, by (612) and (614),

$$dp = -g \gamma \, dz, \quad d\sigma = g \Gamma \, dz;$$

and

$$\begin{aligned} dE &= -f p \, dDv - f dp \, Dv + f \sigma \, dDs + f d\sigma \, Ds \\ &= -d f p \, Dv + d f \sigma \, Ds. \end{aligned} \quad (622)$$

If we now integrate with respect to  $d$ , commencing at the given state of the system, we obtain

$$\Delta E = -\Delta f p \, Dv + \Delta f \sigma \, Ds, \quad (623)$$

where  $\Delta$  denotes the value of a quantity in a varied state of the system diminished by its value in the given state. This is true for finite variations, and is therefore true for infinitesimal variations without neglect of the infinitesimals of the higher orders. The condition of stability is therefore that

$$\Delta f p \, Dv - \Delta f \sigma \, Ds < 0, \quad (624)$$

or that the quantity

$$f p \, Dv - f \sigma \, Ds \quad (625)$$

has a maximum value, the values of  $p$  and  $\sigma$ , for each different mass or surface, being regarded as determined functions of  $z$ . (In ordinary cases  $\sigma$  may be regarded as constant in each surface of discontinuity, and  $p$  as a linear function of  $z$  in each different mass.) It may easily be shown (compare page 416) that this condition is always sufficient for stability with reference to motion of surfaces of discontinuity, even when the variations of  $t$ ,  $M_1$ ,  $M_2$ , etc. cannot be neglected in the determination of the necessary condition of stability with respect to such changes.

#### *On the Possibility of the Formation of a New Surface of Discontinuity where several Surfaces of Discontinuity meet.*

When more than three surfaces of discontinuity between homogeneous masses meet along a line, we may conceive of a new surface being formed between any two of the masses which do not meet in a surface in the original state of the system. The condition of stability with respect to the formation of such a surface may be easily obtained by the consideration of the limit between stability and instability, as exemplified by a system which is in equilibrium when a very small surface of the kind is formed.

To fix our ideas, let us suppose that there are four homogeneous masses A, B, C, and D, which meet one another in four surfaces, which we may call A-B, B-C, C-D, and D-A, these surfaces all meeting along a line L. This is indicated in figure 11 by a section of the sur-

faces cutting the line L at right angles at a point O. In an infinitesimal variation of the state of the system, we may conceive of a small surface being formed between A and C (to be called A-C), so that the section of the surfaces of discontinuity by the same plane takes the form indicated in figure 12. Let us suppose that

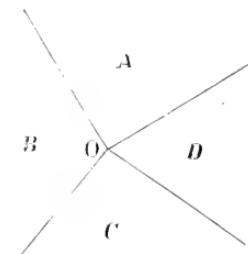


FIG. 11.

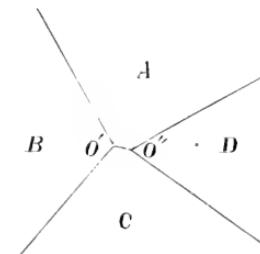


FIG. 12.

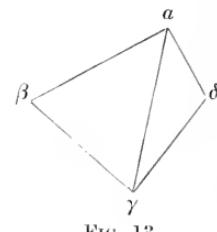


FIG. 13.

the condition of equilibrium (615) is satisfied both for the line L in which the surfaces of discontinuity meet in the original state of the system, and for the two such lines (which we may call L' and L'') in the varied state of the system, at least at the points O' and O'' where they are cut by the plane of section. We may therefore form a quadrilateral of which the sides  $\alpha\beta$ ,  $\beta\gamma$ ,  $\gamma\delta$ ,  $\delta\alpha$  are equal in numerical value to the tensions of the several surfaces A-B, B-C, C-D, D-A, and are parallel to the normals to these surfaces at the point O in the original state of the system. In like manner, for the varied state of the system we can construct two triangles having similar relations to the surfaces of discontinuity meeting at O' and O''. But the directions of the normals to the surfaces A-B and B-C at O' and to C-D and D-A at O'' in the varied state of the system differ infinitely little from the directions of the corresponding normals at O in the initial state. We may therefore regard  $\alpha\beta$ ,  $\beta\gamma$  as two sides of the triangle representing the surfaces meeting at O', and  $\gamma\delta$ ,  $\delta\alpha$  as two sides of the triangle representing the surfaces meeting at O''. Therefore, if we join  $\alpha\gamma$ , this line will represent the direction of the normal to the surface A-C, and the value of its tension. If the tension of a surface between such masses as A and C had been greater than that represented by  $\alpha\gamma$ , it is evident that the initial state of the system of surfaces (represented in figure 11) would have been stable with respect to the possible formation of any such surface. If the tension had been less, the state of the system would have been at least practically unstable. To determine whether it is unstable in the strict sense of the term, or whether or not it is prop-

erly to be regarded as in equilibrium, would require a more refined analysis than we have used.\*

The result which we have obtained may be generalized as follows. When more than three surfaces of discontinuity in a fluid system meet in equilibrium along a line, with respect to the surfaces and masses immediately adjacent to any point of this line we may form a polygon of which the angular points shall correspond in order to the different masses separated by the surfaces of discontinuity, and the sides to these surfaces, each side being perpendicular to the corresponding surface, and equal to its tension. With respect to the formation of new surfaces of discontinuity in the vicinity of the point especially considered, the system is stable, if every diagonal of the polygon is less, and practically unstable, if any diagonal is greater, than the tension which would belong to the surface of discontinuity between the corresponding masses. In the limiting case, when the diagonal is exactly equal to the tension of the corresponding surface, the system may often be determined to be unstable by the application of the principle enunciated to an adjacent point of the line in which the surfaces of discontinuity meet. But when, in the polygons constructed for all points of the line, no diagonal is in any case greater

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\* We may here remark that a nearer approximation in the theory of equilibrium and stability might be attained, by taking special account, in our general equations, of the lines in which surfaces of discontinuity meet. These lines might be treated in a manner entirely analogous to that in which we have treated surfaces of discontinuity. We might recognize linear densities of energy, of entropy, and of the several substances which occur about the line, also a certain linear tension. With respect to these quantities and the temperature and potentials, relations would hold analogous to those which have been demonstrated for surfaces of discontinuity. (See pp. 391–393.) If the sum of the tensions of the lines  $L'$  and  $L''$ , mentioned above, is greater than the tension of the line  $L$ , this line will be in strictness stable (although practically unstable) with respect to the formation of a surface between A and C, when the tension of such a surface is a little less than that represented by the diagonal  $ay$ .

The different use of the term *practically unstable* in different parts of this paper need not create confusion, since the general meaning of the term is in all cases the same. A system is called practically unstable when a very small (not necessarily indefinitely small) disturbance or variation in its condition will produce a considerable change. In the former part of this paper, in which the influence of surfaces of discontinuity was neglected, a system was regarded as practically unstable when such a result would be produced by a disturbance of the same order of magnitude as the quantities relating to surfaces of discontinuity which were neglected. But where surfaces of discontinuity are considered, a system is not regarded as practically unstable, unless the disturbance which will produce such a result is very small compared with the quantities relating to surfaces of discontinuity of any appreciable magnitude.

than the tension of the corresponding surface, but a certain diagonal is equal to the tension in the polygons constructed for a finite portion of the line, farther investigations are necessary to determine the stability of the system. For this purpose, the method described on page 413 is evidently applicable.

A similar proposition may be enunciated in many cases with respect to a point about which the angular space is divided into solid angles by surfaces of discontinuity. If these surfaces are in equilibrium, we can always form a closed solid figure without re-entrant angles of which the angular points shall correspond to the several masses, the edges to the surfaces of discontinuity, and the sides to the lines in which these edges meet, the edges being perpendicular to the corresponding surfaces, and equal to their tensions, and the sides being perpendicular to the corresponding lines. Now if the solid angles in the physical system are such as may be subtended by the sides and bases of a triangular prism enclosing the vertical point, or can be derived from such by deformation, the figure representing the tensions will have the form of two triangular pyramids on opposite sides of the same base, and the system will be stable or practically unstable with respect to the formation of a surface between the masses which only meet in a point, according as the tension of a surface between such masses is greater or less than the diagonal joining the corresponding angular points of the solid representing the tensions. This will easily appear on consideration of the case in which a very small surface between the masses would be in equilibrium.

*The Conditions of Stability for Fluids relating to the Formation  
of a New Phase at a Line in which Three Surfaces of  
Discontinuity meet.*

With regard to the formation of new phases there are particular conditions of stability which relate to lines in which several surfaces of discontinuity meet. We may limit ourselves to the case in which there are three such surfaces, this being the only one of frequent occurrence, and may treat them as meeting in a straight line. It will be convenient to commence by considering the equilibrium of a system in which such a line is replaced by a filament of a different phase.

Let us suppose that three homogeneous fluid masses, A, B, and C, are separated by cylindrical (or plane) surfaces, A-B, B-C, C-A, which at first meet in a straight line, each of the surface-tensions  $\sigma_{AB}$ ,  $\sigma_{BC}$ ,  $\sigma_{CA}$  being less than the sum of the other two. Let us suppose that the

system is then modified by the introduction of a fourth fluid mass D, which is placed between A, B, and C, and is separated from them by cylindrical surfaces D-A, D-B, D-C meeting A-B, B-C, and C-A in straight lines. The general form of the surfaces is shown by figure 14, in which the full lines represent a section perpendicular to all the surfaces. The system thus modified is to be in equilibrium, as well as the original system, the position of the surfaces A-B, B-C, C-A being unchanged. That the last condition is consistent with equilibrium will appear from the following mechanical considerations.

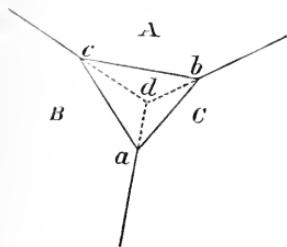


FIG. 14.

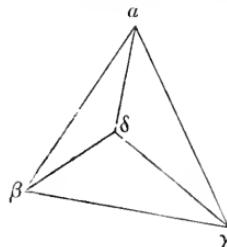


FIG. 15.

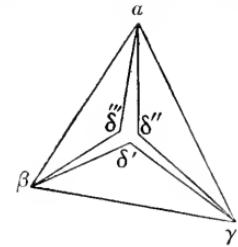


FIG. 16.

Let  $v_D$  denote the volume of the mass D per unit of length or the area of the curvilinear triangle  $a b c$ . Equilibrium is evidently possible for any values of the surface-tensions (if only  $\sigma_{AB}$ ,  $\sigma_{BC}$ ,  $\sigma_{CA}$  satisfy the condition mentioned above, and the tensions of the three surfaces meeting at each of the edges of D satisfy a similar condition) with any value (not too large) of  $v_D$ , if the edges of D are constrained to remain in the original surfaces A-B, B-C, and C-A, or these surfaces extended, if necessary, without change of curvature. (In certain cases one of the surfaces D-A, D-B, D-C may disappear and D will be bounded by only two cylindrical surfaces.) We may therefore regard the system as maintained in equilibrium by forces applied to the edges of D and acting at right angles to A-B, B-C, C-A. The same forces would keep the system in equilibrium if D were rigid. They must therefore have a zero resultant, since the nature of the mass D is immaterial when it is rigid, and no forces external to the system would be required to keep a corresponding part of the original system in equilibrium. But it is evident from the points of application and directions of these forces that they cannot have a zero resultant unless each force is zero. We may therefore introduce a fourth mass D without disturbing the parts which remain of the surfaces A-B, B-C, C-D.

It will be observed that all the angles at  $a$ ,  $b$ ,  $c$ , and  $d$  in figure 14 are entirely determined by the six surface-tensions  $\sigma_{AB}$ ,  $\sigma_{BC}$ ,  $\sigma_{CA}$ ,  $\sigma_{DA}$ ,  $\sigma_{DB}$ ,  $\sigma_{DC}$ . [See (615).] The angles may be derived from the tensions

by the following construction, which will also indicate some important properties. If we form a triangle  $\alpha\beta\gamma$  (figure 15 or 16) having sides equal to  $\sigma_{AB}$ ,  $\sigma_{BC}$ ,  $\sigma_{CA}$ , the angles of the triangle will be supplements of the angles at  $d$ . To fix our ideas, we may suppose the sides of the triangle to be perpendicular to the surfaces at  $d$ . Upon  $\beta\gamma$  we may then construct (as in figure 16) a triangle  $\beta\gamma\delta'$  having sides equal to  $\sigma_{BC}$ ,  $\sigma_{DC}$ ,  $\sigma_{DB}$ , upon  $\gamma\alpha$  a triangle  $\gamma\alpha\delta''$  having sides equal to  $\sigma_{CA}$ ,  $\sigma_{DA}$ ,  $\sigma_{DC}$ , and upon  $\alpha\beta$  a triangle  $\alpha\beta\delta'''$  having sides equal to  $\sigma_{AB}$ ,  $\sigma_{DB}$ ,  $\sigma_{DA}$ . These triangles are to be on the same sides of the lines  $\beta\gamma$ ,  $\gamma\alpha$ ,  $\alpha\beta$ , respectively, as the triangle  $\alpha\beta\gamma$ . The angles of these triangles will be supplements of the angles of the surfaces of discontinuity at  $a$ ,  $b$ , and  $c$ . Thus  $\beta\gamma\delta' = d\alpha b$ , and  $\alpha\gamma\delta'' = d\beta a$ . Now if  $\delta'$  and  $\delta''$  fall together in a single point  $\delta$  within the triangle  $\alpha\beta\gamma$ ,  $\delta'''$  will fall in the same point, as in figure 15. In this case we shall have  $\beta\gamma\delta + \alpha\gamma\delta = \alpha\gamma\beta$ , and the three angles of the curvilinear triangle  $adb$  will be together equal to two right angles. The same will be true of the three angles of each of the triangles  $bdc$ ,  $cda$ , and hence of the three angles of the triangle  $abc$ . But if  $\delta'$ ,  $\delta''$ ,  $\delta'''$  do not fall together in the same point within the triangle  $\alpha\beta\gamma$ , it is either possible to bring these points to coincide within the triangle by increasing some or all of the tensions  $\sigma_{DA}$ ,  $\sigma_{DB}$ ,  $\sigma_{DC}$ , or to effect the same result by diminishing some or all of these tensions. (This will easily appear when one of the points  $\delta'$ ,  $\delta''$ ,  $\delta'''$  falls within the triangle, if we let the two tensions which determine this point remain constant, and the third tension vary. When all the points  $\delta'$ ,  $\delta''$ ,  $\delta'''$  fall without the triangle  $\alpha\beta\gamma$ , we may suppose the greatest of the tensions  $\sigma_{DA}$ ,  $\sigma_{DB}$ ,  $\sigma_{DC}$ —the two greatest, when these are equal, and all three when they all are equal—to diminish until one of the points  $\delta'$ ,  $\delta''$ ,  $\delta'''$  is brought within the triangle  $\alpha\beta\gamma$ .) In the first case we may say that the tensions of the new surfaces are too small to be represented by the distances of an internal point from the vertices of the triangle representing the tensions of the original surfaces (or, for brevity, that they are too small to be represented as in figure 15); in the second case we may say that they are too great to be thus represented. In the first case, the sum of the angles in each of the triangles  $adb$ ,  $bdc$ ,  $cda$  is less than two right angles (compare figures 14 and 16): in the second case, each pair of the triangles  $\alpha\beta\delta'''$ ,  $\beta\gamma\delta''$ ,  $\gamma\alpha\delta''$  will overlap, at least when the tensions  $\sigma_{DA}$ ,  $\sigma_{DB}$ ,  $\sigma_{DC}$  are only a little too great to be represented as in figure 15, and the sum of the angles of each of the triangles  $adb$ ,  $bdc$ ,  $cda$  will be greater than two right angles.

Let us denote by  $v_A$ ,  $v_B$ ,  $v_C$  the portions of  $v_D$  which were originally occupied by the masses A, B, C, respectively, by  $s_{DA}$ ,  $s_{DB}$ ,  $s_{DC}$ , the areas of the surfaces specified per unit of length of the mass D, and by  $s_{AB}$ ,  $s_{BC}$ ,  $s_{CA}$ , the areas of the surfaces specified which were replaced by the mass D per unit of its length. In numerical value,  $v_A$ ,  $v_B$ ,  $v_C$  will be equal to the areas of the curvilinear triangles  $bcd$ ,  $cad$ ,  $abd$ ; and  $s_{DA}$ ,  $s_{DB}$ ,  $s_{DC}$ ,  $s_{AB}$ ,  $s_{BC}$ ,  $s_{CA}$  to the lengths of the lines  $bc$ ,  $ca$ ,  $ab$ ,  $cd$ ,  $ad$ ,  $bd$ . Also let

$$W_s = \sigma_{DA} s_{DA} + \sigma_{DB} s_{DB} + \sigma_{DC} s_{DC} - \sigma_{AB} s_{AB} - \sigma_{BC} s_{BC} - \sigma_{CA} s_{CA}, \quad (626)$$

and  $W_v = p_D v_D - p_A v_A - p_B v_B - p_C v_C. \quad (627)$

The general condition of mechanical equilibrium for a system of homogeneous masses not influenced by gravity, when the exterior of the whole system is fixed, may be written

$$\Sigma (\sigma \delta s) - \Sigma (p \delta v) = 0. \quad (628)$$

[See (606).] If we apply this both to the original system consisting of the masses A, B, and C, and to the system modified by the introduction of the mass D, and take the difference of the results, supposing the deformation of the system to be the same in each case, we shall have

$$\begin{aligned} \sigma_{DA} \delta s_{DA} + \sigma_{DB} \delta s_{DB} + \sigma_{DC} \delta s_{DC} - \sigma_{AB} \delta s_{AB} - \sigma_{BC} \delta s_{BC} \\ - \sigma_{CA} \delta s_{CA} - p_D \delta v_D + p_A \delta v_A + p_B \delta v_B + p_C \delta v_C = 0. \end{aligned} \quad (629)$$

In view of this relation, if we differentiate (626) and (627) regarding all quantities except the pressures as variable, we obtain

$$\begin{aligned} dW_s - dW_v = s_{DA} d\sigma_{DA} + s_{DB} d\sigma_{DB} + s_{DC} d\sigma_{DC} \\ - s_{AB} d\sigma_{AB} - s_{BC} d\sigma_{BC} - s_{CA} d\sigma_{CA}. \end{aligned} \quad (630)$$

Let us now suppose the system to vary in size, remaining always similar to itself in form, and that the tensions diminish in the same ratio as lines, while the pressures remain constant. Such changes will evidently not impair the equilibrium. Since all the quantities  $s_{DA}$ ,  $\sigma_{DA}$ ,  $s_{DB}$ ,  $\sigma_{DB}$ , etc. vary in the same ratio,

$$s_{DA} d\sigma_{DA} = \frac{1}{2} d(\sigma_{DA} s_{DA}), \quad s_{DB} d\sigma_{DB} = \frac{1}{2} d(\sigma_{DB} s_{DB}), \quad \text{etc.} \quad (631)$$

We have therefore by integration of (630)

$$W_s - W_v = \frac{1}{2} (\sigma_{DA} s_{DA} + \sigma_{DB} s_{DB} + \sigma_{DC} s_{DC} - \sigma_{AB} s_{AB} - \sigma_{BC} s_{BC} - \sigma_{CA} s_{CA}), \quad (632)$$

whence, by (626),

$$W_s = 2 W_v, \quad (633)$$

The condition of stability for the system when the pressures and tensions are regarded as constant, and the position of the surfaces

A-B, B-C, C-A as fixed, is that  $W_s - W_v$  shall be a minimum under the same conditions. [See (549).] Now for any constant values of the tensions and of  $p_A, p_B, p_C$ , we may make  $v_D$  so small that when it varies, the system remaining in equilibrium, (which will in general require a variation of  $p_D$ ), we may neglect the curvatures of the lines  $a, b, c$ , and regard the figure  $a b c d$  as remaining similar to itself. For the *total curvature* (*i. e.*, the curvature measured in degrees) of each of the lines  $a b, b c, c a$  may be regarded as constant, being equal to the constant difference of the sum of the angles of one of the curvilinear triangles  $a b d, b d c, c d a$  and two right angles. Therefore, when  $v_D$  is very small, and the system is so deformed that equilibrium would be preserved if  $p_D$  had the proper variation, but this pressure as well as the others and all the tensions remain constant,  $W_s$  will vary as the lines in the figure  $a b c d$ , and  $W_v$  as the square of these lines. Therefore, for such deformations,

$$W_v \propto W_s^2.$$

This shows that the system cannot be stable for constant pressures and tensions when  $v_D$  is small and  $W_v$  is positive, since  $W_s - W_v$  will not be a minimum. It also shows that the system is stable when  $W_v$  is negative. For, to determine whether  $W_s - W_v$  is a minimum for constant values of the pressures and tensions, it will evidently be sufficient to consider such varied forms of the system as give the least value to  $W_s - W_v$  for any value of  $v_D$  in connection with the constant pressures and tensions. And it may easily be shown that such forms of the system are those which would preserve equilibrium if  $p_D$  had the proper value.

These results will enable us to determine the most important questions relating to the stability of a line along which three homogeneous fluids A, B, C meet, with respect to the formation of a different fluid D. The components of D must of course be such as are found in the surrounding bodies. We shall regard  $p_D$  and  $\sigma_{DA}, \sigma_{DB}, \sigma_{DC}$  as determined by that phase of D which satisfies the conditions of equilibrium with the other bodies relating to temperature and the potentials. These quantities are therefore determinable, by means of the fundamental equations of the mass D and of the surfaces D-A, D-B, D-C, from the temperature and potentials of the given system.

Let us first consider the case in which the tensions, thus determined, can be represented as in figure 15, and  $p_D$  has a value consistent with the equilibrium of a small mass such as we have been considering. It appears from the preceding discussion that when  $v_D$  is

sufficiently small the figure  $a b c d$  may be regarded as rectilinear, and that its angles are entirely determined by its tensions. Hence the ratios of  $v_A$ ,  $v_B$ ,  $v_C$ ,  $v_D$ , for sufficiently small values of  $v_D$ , are determined by the tensions alone, and for convenience in calculating these ratios, we may suppose  $p_A$ ,  $p_B$ ,  $p_C$  to be equal, which will make the figure  $a b c d$  absolutely rectilinear, and make  $p_D$  equal to the other pressures, since it is supposed that this quantity has the value necessary for equilibrium. We may obtain a simple expression for the ratios of  $v_A$ ,  $v_B$ ,  $v_C$ ,  $v_D$  in terms of the tensions in the following manner. We shall write  $[D B C]$ ,  $[D C A]$ , etc., to denote the areas of triangles having sides equal to the tensions of the surfaces between the masses specified.

$$\begin{aligned} v_A : v_B &:: \text{triangle } b d e : \text{triangle } a d c \\ &:: bc \sin bed : ac \sin aed \\ &:: \sin bac \sin bed : \sin abc \sin aed \\ &:: \sin \gamma \delta \beta \sin \delta \alpha \beta : \sin \gamma \delta \alpha \sin \delta \beta \alpha \\ &:: \sin \gamma \delta \beta \delta \beta : \sin \gamma \delta \alpha \delta \alpha \\ &:: \text{triangle } \gamma \delta \beta : \text{triangle } \gamma \delta \alpha \\ &:: [D B C] : [D C A]. \end{aligned}$$

Hence,

$$v_A : v_B : v_C : v_D :: [D B C] : [D C A] : [D A B] : [A B C], \quad (634)$$

where

$\frac{1}{4}\sqrt{[(\sigma_{AB} + \sigma_{BC} + \sigma_{CA})(\sigma_{AB} + \sigma_{BC} - \sigma_{CA})(\sigma_{BC} + \sigma_{CA} - \sigma_{AB})(\sigma_{CA} + \sigma_{AB} - \sigma_{BC})]}$  may be written for  $[A B C]$ , and analogous expressions for the other symbols, the sign  $\sqrt{\phantom{x}}$  denoting the positive root of the necessarily positive expression which follows. This proportion will hold true in any case of equilibrium, when the tensions satisfy the condition mentioned and  $v_D$  is sufficiently small. Now if  $p_A = p_B = p_C$ ,  $p_D$  will have the same value, and we shall have by (627)  $W_v = 0$ , and by (633)  $W_s = 0$ . But when  $v_D$  is very small, the value of  $W_s$  is entirely determined by the tensions and  $v_D$ . Therefore, whenever the tensions satisfy the condition supposed, and  $v_D$  is very small (whether  $p_A$ ,  $p_B$ ,  $p_C$  are equal or unequal,)

$$0 = W_s = W_v = p_D v_D - p_A v_A - p_B v_B - p_C v_C, \quad (635)$$

which with (634) gives

$$p_D = \frac{[D B C] p_A + [D C A] p_B + [D A B] p_C}{[D B C] + [D C A] + [D A B]}. \quad (636)$$

Since this is the only value of  $p_D$  for which equilibrium is possible when

the tensions satisfy the condition supposed and  $v_D$  is small, it follows that when  $p_D$  has a less value, the line where the fluids A, B, C meet is stable with respect to the formation of the fluid D. When  $p_D$  has a greater value, if such a line can exist at all, it must be at least practically unstable, *i. e.*, if only a very small mass of the fluid D should be formed it would tend to increase.

Let us next consider the case in which the tensions of the new surfaces are too small to be represented as in figure 15. If the pressures and tensions are consistent with equilibrium for any very small value of  $v_D$ , the angles of each of the curvilinear triangles  $adb$ ,  $bde$ ,  $cda$  will be together less than two right angles, and the lines  $ab$ ,  $bc$ ,  $ca$ , will be convex toward the mass D. For given values of the pressures and tensions, it will be easy to determine the magnitude of  $v_D$ . For the tensions will give the total curvatures (in degrees) of the lines  $ab$ ,  $bc$ ,  $ca$ ; and the pressures will give the radii of curvature. These lines are thus completely determined. In order that  $v_D$  shall be very small it is evidently necessary that  $p_D$  shall be less than the other pressures. Yet if the tensions of the new surfaces are only a very little too small to be represented as in figure 15,  $v_D$  may be quite small when the value of  $p_D$  is only a little less than that given by equation (636). In any case, when the tensions of the new surfaces are too small to be represented as in figure 15, and  $v_D$  is small,  $W_v$  is negative, and the equilibrium of the mass D is stable. Moreover,  $W_s - W_v$ , which represents the work necessary to form the mass D with its surfaces in place of the other masses and surfaces, is negative.

With respect to the stability of a line in which the surfaces A-B, B-C, C-A meet, when the tensions of the new surfaces are too small to be represented as in figure 15, we first observe that when the pressures and tensions are such as to make  $v_D$  moderately small but not so small as to be neglected, [this will be when  $p_D$  is somewhat smaller than the second member of (636),—more or less smaller according as the tensions differ more or less from such as are represented in figure 15,] the equilibrium of such a line as that supposed (if it is capable of existing at all) is at least practically unstable. For greater values of  $p_D$  (with the same values of the other pressures and the tensions) the same will be true. For somewhat smaller values of  $p_D$ , the mass of the phase D which will be formed will be so small, that we may neglect this mass and regard the surfaces A-B, B-C, C-A as meeting in a line in stable equilibrium. For still smaller values of  $p_D$ , we may likewise regard the surfaces A-B, B-C, C-A as capable

of meeting in stable equilibrium. It may be observed that when  $v_D$ , as determined by our equations, becomes quite insensible, the conception of a small mass D having the properties deducible from our equations ceases to be accurate, since the matter in the vicinity of a line where these surfaces of discontinuity meet must be in a peculiar state of equilibrium not recognized by our equations.\* But this cannot affect the validity of our conclusion with respect to the stability of the line in question.

The case remains to be considered in which the tensions of the new surfaces are too great to be represented as in figure 15. Let us suppose that they are not very much too great to be thus represented. When the pressures are such as to make  $v_D$  moderately small (in case of equilibrium) but not so small that the mass D to which it relates ceases to have the properties of matter in mass, [this will be when  $p_D$  is somewhat greater than the second member of (636),—more or less greater according as the tensions differ more or less from such as are represented in figure 15,] the line where the surfaces A-B, B-C, C-A meet will be in stable equilibrium with respect to the formation of such a mass as we have considered, since  $W_s - W_v$  will be positive. The same will be true for less values of  $p_D$ . For greater values of  $p_D$ , the value of  $W_s - W_v$ , which measures the stability with respect to the kind of change considered, diminishes. It does not vanish, according to our equations, for finite values of  $p_D$ . But these equations are not to be trusted beyond the limit at which the mass D ceases to be of sensible magnitude.

But when the tensions are such as we now suppose, we must also consider the possible formation of a mass D within a closed figure in which the surfaces D-A, D-B, D-C meet together (with the surfaces A-B, B-C, C-A) in two opposite points. If such a figure is to be in equilibrium, the six tensions must be such as can be represented by

\* See note on page 455. We may here add that the linear tension there mentioned may have a negative value. This would be the case with respect to a line in which three surfaces of discontinuity are regarded as meeting, but where nevertheless there really exists in stable equilibrium a filament of different phase from the three surrounding masses. The value of the linear tension for the supposed line, would be nearly equal to the value of  $W_s - W_v$  for the actually existing filament. (For the exact value of the linear tension it would be necessary to add the sum of the linear tensions of the three edges of the filament.) We may regard two soap-bubbles adhering together as an example of this case. The reader will easily convince himself that in an exact treatment of the equilibrium of such a double bubble we must recognize a certain negative tension in the line of intersection of the three surfaces of discontinuity.

the six distances of four points in space (see page 455),—a condition which evidently agrees with the supposition which we have made. If we denote by  $w_v$  the work gained in forming the mass D (of such size and form as to be in equilibrium) in place of the other masses, and by  $w_s$  the work expended in forming the new surfaces in place of the old, it may easily be shown by a method similar to that used on page 459 that  $w_s = \frac{3}{2}w_v$ . From this we obtain  $w_s - w_v = \frac{1}{2}w_v$ . This is evidently positive when  $p_D$  is greater than the other pressures. But it diminishes with increase of  $p_D$ , as easily appears from the equivalent expression  $\frac{1}{3}w_s$ . Hence the line of intersection of the surfaces of discontinuity A-B, B-C, C-A is stable for values of  $p_D$  greater than the other pressures (and therefore for all values of  $p_D$ ) so long as our method is to be regarded as accurate, which will be so long as the mass D which would be in equilibrium has a sensible size.

In certain cases in which the tensions of the new surfaces are much too large to be represented as in figure 15, the reasoning of the two last paragraphs will cease to be applicable. These are cases in which the six tensions cannot be represented by the sides of a tetrahedron. It is not necessary to discuss these cases, which are distinguished by the different shape which the mass D would take if it should be formed, since it is evident that they can constitute no exception to the results which we have obtained. For an increase of the values of  $\sigma_{DA}$ ,  $\sigma_{DB}$ ,  $\sigma_{DC}$  cannot favor the formation of D, and hence cannot impair the stability of the line considered, as deduced from our equations. Nor can an increase of these tensions essentially affect the fact that the stability thus demonstrated may fail to be realized when  $p_D$  is considerably greater than the other pressures, since the *a priori* demonstration of the stability of any one of the surfaces A-B, B-C, C-A, taken singly, is subject to the limitation mentioned. (See page 426.)

*The Condition of Stability for Fluids relating to the Formation  
of a New Phase at a Point where the Vertices of  
Four Different Masses meet.*

Let four different fluid masses A, B, C, D meet about a point, so as to form the six surfaces of discontinuity A-B, B-C, C-A, D-A, D-B, D-C, which meet in the four lines A-B-C, B-C-D, C-D-A, D-A-B, these lines meeting in the vertical point. Let us suppose the system stable in other respects, and consider the conditions of stability for the vertical point with respect to the possible formation of a different fluid mass E.

If the system can be in equilibrium when the vertical point has been replaced by a mass E against which the four masses A, B, C, D abut, being truncated at their vertices, it is evident that E will have four vertices, at each of which six surfaces of discontinuity meet. (Thus at one vertex there will be the surfaces formed by A, B, C, and E.) The tensions of each set of six surfaces (like those of the six surfaces formed by A, B, C, and D) must therefore be such that they can be represented by the six edges of a tetrahedron. When the tensions do not satisfy these relations, there will be no particular condition of stability for the point about which A, B, C, and D meet, since if a mass E should be formed, it would distribute itself along some of the lines or surfaces which meet at the vertical point, and it is therefore sufficient to consider the stability of these lines and surfaces. We shall suppose that the relations mentioned are satisfied.

If we denote by  $W_v$  the work gained in forming the mass E (of such size and form as to be in equilibrium) in place of the portions of the other masses which are suppressed, and by  $W_s$  the work expended in forming the new surfaces in place of the old, it may easily be shown by a method similar to that used on page 459 that

$$W_s = \frac{3}{2} W_v, \quad (637)$$

whence

$$W_s - W_v = \frac{1}{2} W_v; \quad (638)$$

also, that when the volume E is small, the equilibrium of E will be stable or unstable according as  $W_s$  and  $W_v$  are negative or positive.

A critical relation for the tensions is that which makes equilibrium possible for the system of the five masses A, B, C, D, E, when all the surfaces are plane. The ten tensions may then be represented in magnitude and direction by the ten distances of five points in space  $\alpha, \beta, \gamma, \delta, \varepsilon$ , viz., the tension of A-B and the direction of its normal by the line  $\alpha\beta$ , etc. The point  $\varepsilon$  will lie within the tetrahedron formed by the other points. If we write  $v_E$  for the volume of E, and  $v_A, v_B, v_C, v_D$  for the volumes of the parts of the other masses which are suppressed to make room for E, we have evidently

$$W_v = p_E v_E - p_A v_A - p_B v_B - p_C v_C - p_D v_D. \quad (639)$$

Hence, when all the surfaces are plane,  $W_v = 0$ , and  $W_s = 0$ . Now equilibrium is always possible for a given small value of  $v_E$  with any given values of the tensions and of  $p_A, p_B, p_C, p_D$ . When the tensions satisfy the critical relation,  $W_s = 0$ , if  $p_A = p_B = p_C = p_D$ . But when  $v_E$  is small and constant, the value of  $W_s$  must be independent of  $p_A, p_B, p_C, p_D$ , since the angles of the surfaces are determined by the tensions and their curvatures may be neglected. Hence,  $W_s = 0$ , and

$v_E = 0$ , when the critical relation is satisfied and  $v_E$  small. This gives

$$\rho_E = \frac{v_A p_A + v_B p_B + v_C p_C + v_D p_D}{v_E}. \quad (640)$$

In calculating the ratios of  $v_A, v_B, v_C, v_D, v_E$ , we may suppose all the surfaces to be plane. Then E will have the form of a tetrahedron, the vertices of which may be called a, b, c, d, (each vertex being named after the mass which is not found there,) and  $v_A, v_B, v_C, v_D$  will be the volumes of the tetrahedra into which it may be divided by planes passing through its edges and an interior point e. The volumes of these tetrahedra are proportional to those of the five tetrahedra of the figure  $\alpha\beta\gamma\delta\varepsilon$ , as will easily appear if we recollect that the line ab is common to the surfaces C-D, D-E, E-C, and therefore perpendicular to the surface common to the lines  $\gamma\delta, \delta\varepsilon, \varepsilon\gamma$ , i.e., to the surface  $\gamma\delta\varepsilon$ , and so in other cases, (it will be observed that  $\gamma, \delta$ , and  $\varepsilon$  are the letters which do not correspond to a or b); also that the surface abc is the surface D-E and therefore perpendicular to  $\delta\varepsilon$ , etc. Let tetr abed, trian abc, etc. denote the volume of the tetrahedron or the area of the triangle specified,  $\sin(ab, bc)$ ,  $\sin(abc, dbc)$ ,  $\sin(abc, ad)$ , etc. the sines of the angles made by the lines and surfacees specified, and [B C D E], [C D E A], etc., the volumes of tetrahedra having edges equal to the tensions of the surfaces between the masses specified. Then, since we may express the volume of a tetrahedron either by  $\frac{1}{3}$  of the product of one side, an edge leading to the opposite vertex, and the sine of the angle which these make, or by  $\frac{2}{3}$  of the product of two sides divided by the common edge and multiplied by the sine of the included angle,

$$v_A : v_B :: \text{tetr } bced : \text{tetr } acde$$

$$:: bc \sin(bc, ede) : ae \sin(ac, ede)$$

$$:: \sin(ba, ac) \sin(bc, ede) : \sin(ab, bc) \sin(ac, ede)$$

$$:: \sin(\gamma\delta\varepsilon, \beta\delta\varepsilon) \sin(\alpha\delta\varepsilon, \alpha\beta) : \sin(\gamma\delta\varepsilon, \alpha\delta\varepsilon) \sin(\beta\delta\varepsilon, \alpha\beta)$$

$$:: \frac{\text{tetr } \gamma\beta\delta\varepsilon \text{ tetr } \beta\alpha\delta\varepsilon}{\text{trian } \beta\delta\varepsilon} : \frac{\text{tetr } \gamma\alpha\delta\varepsilon \text{ tetr } \alpha\beta\delta\varepsilon}{\text{trian } \alpha\delta\varepsilon} \quad \text{trian } \alpha\delta\varepsilon \quad \text{trian } \beta\delta\varepsilon$$

$$:: \text{tetr } \gamma\beta\delta\varepsilon : \text{tetr } \gamma\alpha\delta\varepsilon$$

$$:: [B C D E] : [C D E A].$$

Hence,

$$v_A : v_B : v_C : v_D :: [B C D E] : [C D E A] : [D E A B] : [E A B C], \quad (641)$$

and (640) may be written

$$\rho_E = \frac{[B C D E]p_A + [C D E A]p_B + [D E A B]p_C + [E A B C]p_D}{[B C D E] + [C D E A] + [D E A B] + [E A B C]}. \quad (642)$$

If the value of  $p_E$  is less than this, when the tensions satisfy the critical relation, the point where vertices of the masses A, B, C, D meet is stable with respect to the formation of any mass of the nature of E. But if the value of  $p_E$  is greater, either the masses A, B, C, D cannot meet at a point in equilibrium, or the equilibrium will be at least practically unstable.

When the tensions of the new surfaces are too small to satisfy the critical relation with the other tensions, these surfaces will be convex toward E; when their tensions are too great for that relation, the surfaces will be concave toward E. In the first case,  $W_v$  is negative, and the equilibrium of the five masses A, B, C, D, E is stable, but the equilibrium of the four masses A, B, C, D meeting at a point is impossible or at least practically unstable. This is subject to the limitation that when  $p_E$  is sufficiently small the mass E which will form will be so small that it may be neglected. This will only be the case when  $p_E$  is smaller—in general considerably smaller—than the second number of (642). In the second case, the equilibrium of the five masses A, B, C, D, E will be unstable, but the equilibrium of the four masses A, B, C, D will be stable unless  $v_E$  (calculated for the case of the five masses) is of insensible magnitude. This will only be the case when  $p_E$  is greater—in general considerably greater—than the second member of (642).

#### Liquid Films.

When a fluid exists in the form of a thin film between other fluids, the great inequality of its extension in different directions will give rise to certain peculiar properties, even when its thickness is sufficient for its interior to have the properties of matter in mass. The frequent occurrence of such films, and the remarkable properties which they exhibit, entitle them to particular consideration. To fix our ideas, we shall suppose that the film is liquid and that the contiguous fluids are gaseous. The reader will observe our results are not dependent, so far as their general character is concerned, upon this supposition.

Let us imagine the film to be divided by surfaces perpendicular to its sides into small portions of which all the dimensions are of the same order of magnitude as the thickness of the film,—such portions to be called *elements of the film*,—it is evident that far less time will in general be required for the attainment of approximate equilibrium between the different parts of any such element and the other fluids which are immediately contiguous, than for the attainment of equi-

librium between all the different elements of the film. There will accordingly be a time, commencing shortly after the formation of the film, in which its separate elements may be regarded as satisfying the conditions of internal equilibrium, and of equilibrium with the contiguous gases, while they may not satisfy all the conditions of equilibrium with each other. It is when the changes due to this want of complete equilibrium take place so slowly that the film appears to be at rest, except so far as it accommodates itself to any change in the external conditions to which it is subjected, that the characteristic properties of the film are most striking and most sharply defined.

Let us therefore consider the properties which will belong to a film sufficiently thick for its interior to have the properties of matter in mass, in virtue of the approximate equilibrium of all its elements taken separately, when the matter contained in each element is regarded as invariable, with the exception of certain substances which are components of the contiguous gas-masses and have their potentials thereby determined. The occurrence of a film which precisely satisfies these conditions may be exceptional, but the discussion of this somewhat ideal case will enable us to understand the principal laws which determine the behavior of liquid films in general.

Let us first consider the properties which will belong to each element of the film under the conditions mentioned. Let us suppose the element extended, while the temperature and the potentials which are determined by the contiguous gas-masses are unchanged. If the film has no components except those of which the potentials are maintained constant, there will be no variation of tension in its surfaces. The same will be true when the film has only one component of which the potential is not maintained constant, provided that this is a component of the interior of the film and not of its surface alone. If we regard the thickness of the film as determined by *dividing surfaces* which make the surface-density of this component vanish, the thickness will vary inversely as the area of the element of the film, but no change will be produced in the nature or the tension of its surfaces. If, however, the single component of which the potential is not maintained constant is confined to the surfaces of the film, an extension of the element will generally produce a decrease in the potential of this component, and an increase of tension. This will certainly be true in those cases in which the component shows a tendency to distribute itself with a uniform superficial density.

When the film has two or more components of which the potentials are not maintained constant by the contiguous gas masses, they will not in general exist in the same proportion in the interior of the film as on its surfaces, but those components which diminish the tensions will be found in greater proportion on the surfaces. When the film is extended, there will therefore not be enough of these substances to keep up the same volume- and surface-densities as before, and the deficiency will cause a certain increase of tension. The value of the *elasticity of the film*, (*i. e.*, the infinitesimal increase of the united tensions of its surfaces divided by the infinitesimal increase of area in a unit of surface), may be calculated from the quantities which specify the nature of the film, when the fundamental equations of the interior mass, of the contiguous gas-masses, and of the two surfaces of discontinuity are known. We may illustrate this by a simple example.

Let us suppose that the two surfaces of a plane film are entirely alike, that the contiguous gas-masses are identical in phase, and that they determine the potentials of all the components of the film except two. Let us call these components  $S_1$  and  $S_2$ , the latter denoting that which occurs in greater proportion on the surface than in the interior of the film. Let us denote by  $\gamma_1$  and  $\gamma_2$  the densities of these components in the interior of the film, by  $\lambda$  the thickness of the film determined by such dividing surfaces as make the surface-density of  $S_1$  vanish (see page 397), by  $\Gamma_{2(1)}$  the surface-density of the other component as determined by the same surfaces, by  $\sigma$  and  $s$  the tension and area of one of these surfaces, and by  $E$  the elasticity of the film when extended under the supposition that the total quantities of  $S_1$  and  $S_2$  in the part of the film extended are invariable, as also the temperature and the potentials of the other components. From the definition of  $E$  we have

$$2 d\sigma = E \frac{ds}{s}, \quad (643)$$

and from the conditions of the extension of the film

$$\frac{ds}{s} = - \frac{d(\lambda \gamma_1)}{\lambda \gamma_1} = - \frac{d(\lambda \gamma_2 + 2 \Gamma_{2(1)})}{\lambda \gamma_2 + 2 \Gamma_{2(1)}}. \quad (644)$$

Hence we obtain

$$\lambda \gamma_1 \frac{ds}{s} = - \gamma_1 d\lambda - \lambda d\gamma_1,$$

$$(\lambda \gamma_2 + 2 \Gamma_{2(1)}) \frac{ds}{s} = - \gamma_2 d\lambda - \lambda d\gamma_2 - 2 d\Gamma_{2(1)};$$

and eliminating  $d\lambda$ ,

$$2\gamma_1 \Gamma_{2(1)} \frac{ds}{s} = -\lambda \gamma_1 d\gamma_2 + \lambda \gamma_2 d\gamma_1 - 2\gamma_1 d\Gamma_{2(1)}, \quad (645)$$

If we set  $r = \frac{\gamma_2}{\gamma_1}$ , (646)

we have  $dr = \frac{\gamma_1 d\gamma_2 - \gamma_2 d\gamma_1}{\gamma_1^2}$ , (647)

and  $2\Gamma_{2(1)} \frac{ds}{s} = -\lambda \gamma_1 dr - 2d\Gamma_{2(1)}$ . (648)

With this equation we may eliminate  $ds$  from (643). We may also eliminate  $d\sigma$  by the necessary relation [see (514)]

$$d\sigma = -\Gamma_{2(1)} d\mu_2.$$

This will give

$$4\Gamma_{2(1)}^2 d\mu_2 = E(\lambda \gamma_1 dr + 2d\Gamma_{2(1)}), \quad (649)$$

or

$$\frac{4\Gamma_{2(1)}^2}{E} = \lambda \gamma_1 \frac{dr}{d\mu_2} + 2 \frac{d\Gamma_{2(1)}}{d\mu_2}, \quad (650)$$

where the differential coefficients are to be determined on the conditions that the temperature and all the potentials except  $\mu_1$  and  $\mu_2$  are constant, and that the pressure in the interior of the film shall remain equal to that in the contiguous gas-masses. The latter condition may be expressed by the equation

$$(\gamma_1 - \gamma_1') d\mu_1 + (\gamma_2 - \gamma_2') d\mu_2 = 0, \quad (651)$$

in which  $\gamma_1'$  and  $\gamma_2'$  denote the densities of  $S_1$  and  $S_2$  in the contiguous gas-masses. [See (98).] When the tension of the surfaces of the film and the pressures in its interior and in the contiguous gas-masses are known in terms of the temperature and potentials, equation (650) will give the value of  $E$  in terms of the same variables together with  $\lambda$ .

If we write  $G_1$  and  $G_2$  for the total quantities of  $S_1$  and  $S_2$  per unit of area of the film, we have

$$G_1 = \lambda \gamma_1, \quad (652)$$

$$G_2 = \lambda \gamma_2 + 2\Gamma_{2(1)}, \quad (653)$$

Therefore,

$$G_2 = G_1 r + 2\Gamma_{2(1)},$$

$$\left( \frac{dG_2}{d\mu_2} \right)_{G_1} = \lambda \gamma_1 \frac{dr}{d\mu_2} + 2 \frac{d\Gamma_{2(1)}}{d\mu_2}, \quad (654)$$

where the differential coefficients in the second member are to be determined as in (650), and that in the first member with the additional condition that  $G_1$  is constant. Therefore,

$$\frac{4 \Gamma_{2(1)}^2}{E} = \left( \frac{dG_2}{d\mu_2} \right)_{G_1},$$

and

$$E = 4 \Gamma_{2(1)}^2 \left( \frac{d\mu_2}{dG_2} \right)_{G_1}, \quad (655)$$

the last differential coefficient being determined by the same conditions as that in the preceding equation. It will be observed that the value of  $E$  will be positive in any ordinary case.

These equations give the elasticity of any element of the film when the temperature and the potentials for the substances which are found in the contiguous gas-masses are regarded as constant, and the potentials for the other components,  $\mu_1$  and  $\mu_2$ , have had time to equalize themselves throughout the element considered. The increase of tension immediately after a rapid extension will be greater than that given by these equations.

The existence of this elasticity, which has thus been established from *a priori* considerations, is clearly indicated by the phenomena which liquid films present. Yet it is not to be demonstrated simply by comparing the tensions of films of different thickness, even when they are made from the same liquid, for difference of thickness does not necessarily involve any difference of tension. When the phases within the films as well as without are the same, and the surfaces of the films are also the same, there will be no difference of tension. Nor will the tension of the same film be altered, if a part of the interior drains away in the course of time, without affecting the surfaces. In case the thickness of the film is reduced by evaporation, the tension may be either increased or diminished. (The evaporation of the substance  $S_1$ , in the case we have just considered, would diminish the tension.) Yet it may easily be shown that extension increases the tension of a film and contraction diminishes it. When a plane film is held vertically, the tension of the upper portions must evidently be greater than that of the lower. The tensions in every part of the film may be reduced to equality by turning it into a horizontal position. By restoring the original position we may restore the original tensions, or nearly so. It is evident that the same element of the film is capable of supporting very unequal tensions. Nor can this be always attributed to viscosity of the film. For in many cases, if we hold the film nearly horizontal, and elevate first one side and then another, the lighter portions of the film will dart from one side to the other, so as to show a very striking mobility in the film. The differences of tension which cause these rapid movements are only a very

small fraction of the difference of tension in the upper and lower portions of the film when held vertically.

If we account for the power of an element of the film to support an increase of tension by viscosity, it will be necessary to suppose that the viscosity offers a resistance to a deformation of the film in which its surface is enlarged and its thickness diminished, which is enormously great in comparison with the resistance to a deformation in which the film is extended in the direction of one tangent and contracted in the direction of another, while its thickness and the areas of its surfaces remain constant. This is not to be readily admitted as a physical explanation, although to a certain extent the phenomena resemble those which would be caused by such a singular viscosity. (See page 439.) The only natural explanation of the phenomena is that the extension of an element of the film, which is the immediate result of an increase of external force applied to its perimeter, causes an increase of its tension, by which it is brought into true equilibrium with the external forces.

The phenomena to which we have referred are such as are apparent to a very cursory observation. In the following experiment, which is described by M. Plateau,\* an increased tension is manifested in a film while contracting after a previous extension. The warmth of a finger brought near to a bubble of soap-water with glycerine, which is thin enough to show colors, causes a spot to appear indicating a diminution of thickness. When the finger is removed, the spot returns to its original color. This indicates a contraction, which would be resisted by any viscosity of the film, and can only be due to an excess of tension in the portion stretched on the return of its original temperature.

We have so far supposed that the film is thick enough for its interior to have the properties of matter in mass. Its properties are then entirely determined by those of the three phases and the two surfaces of discontinuity. From these we can also determine, in part at least, the properties of a film at the limit at which the interior ceases to have the properties of matter in mass. The elasticity of the film, which increases with its thinness, cannot of course vanish at that limit, so that the film cannot become unstable with respect to extension and contraction of its elements immediately after passing that limit.

Yet a certain kind of instability will probably arise, which we may

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\* "Statique expérimentale et théorique des liquides soumis aux seules forces moléculaires," vol. i, p. 294.

here notice, although it relates to changes in which the condition of the invariability of the quantities of certain components in an element of the film is not satisfied. With respect to variations in the distribution of its components, a film will in general be stable, when its interior has the properties of matter in mass, with the single exception of variations affecting its thickness without any change of phase or of the nature of the surfaces. With respect to this kind of change, which may be brought about by a current in the interior of the film, the equilibrium is neutral. But when the interior ceases to have the properties of matter in mass, it is to be supposed that the equilibrium will generally become unstable in this respect. For it is not likely that the neutral equilibrium will be unaffected by such a change of circumstances, and since the film certainly becomes unstable when it is sufficiently reduced in thickness, it is most natural to suppose that the first effect of diminishing the thickness will be in the direction of instability rather than in that of stability. (We are here considering liquid films between gaseous masses. In certain other cases, the opposite supposition might be more natural, as in respect to a film of water between mercury and air, which would certainly become stable when sufficiently reduced in thickness.)

Let us now return to our former suppositions—that the film is thick enough for the interior to have the properties of matter in mass, and that the matter in each element is invariable, except with respect to those substances which have their potentials determined by the contiguous gas-masses—and consider what conditions are necessary for equilibrium in such a case.

In consequence of the supposed equilibrium of its several elements, such a film may be treated as a simple surface of discontinuity between the contiguous gas-masses (which may be similar or different), whenever its radius of curvature is very large in comparison with its thickness,—a condition which we shall always suppose to be fulfilled. With respect to the film considered in this light, the mechanical conditions of equilibrium will always be satisfied, or very nearly so, as soon as a state of approximate rest is attained, except in those cases in which the film exhibits a decided viscosity. That is, the relations (613), (614), (615) will hold true, when by  $\sigma$  we understand the tension of the film regarded as a simple surface of discontinuity (this is equivalent to the sum of the tensions of the two surfaces of the film), and by  $I'$  its mass per unit of area diminished by the mass of gas which would occupy the same space if the film should be suppressed and the gases should meet at its surface of tension. This

*surface of tension of the film* will evidently divide the distance between the surfaces of tension for the two surfaces of the film taken separately, in the inverse ratio of their tensions. For practical purposes, we may regard  $\Gamma$  simply as the mass of the film per unit of area. It will be observed that the terms containing  $\Gamma$  in (613) and (614) are not to be neglected in our present application of these equations.

But the mechanical conditions of equilibrium for the film regarded as an approximately homogeneous mass in the form of a thin sheet bounded by two surfaces of discontinuity are not necessarily satisfied when the film is in a state of apparent rest. In fact, these conditions cannot be satisfied (in any place where the force of gravity has an appreciable intensity) unless the film is horizontal. For the pressure in the interior of the film cannot satisfy simultaneously condition (612), which requires it to vary rapidly with the height  $z$ , and condition (613) applied separately to the different surfaces, which makes it a certain mean between the pressures in the adjacent gas-masses. Nor can these conditions be deduced from the general condition of mechanical equilibrium (606) or (611), without supposing that the interior of the film is free to move independently of the surfaces, which is contrary to what we have supposed.

Moreover, the potentials of the various components of the film will not in general satisfy conditions (617), and cannot (when the temperature is uniform) unless the film is horizontal. For if these conditions were satisfied, equation (612) would follow as a consequence. (See page 449.)

We may here remark that such a film as we are considering cannot form any exception to the principle indicated on page 450,—that when a surface of discontinuity which satisfies the conditions of mechanical equilibrium has only one component which is not found in the contiguous masses, and these masses satisfy all the conditions of equilibrium, the potential for the component mentioned must satisfy the law expressed in (617), as a consequence of the condition of mechanical equilibrium (614). Therefore, as we have just seen that it is impossible that all the potentials in a liquid film which is not horizontal should conform to (617) when the temperature is uniform, it follows that if a liquid film exhibits any persistence which is not due to viscosity, or to a horizontal position, or to differences of temperature, it must have more than one component of which the potential is not determined by the contiguous gas-masses in accordance with (617).

The difficulties of the quantitative experimental verification of the properties which have been described would be very great, even in cases in which the conditions we have imagined were entirely fulfilled. Yet the general effect of any divergence from these conditions will be easily perceived, and when allowance is made for such divergence, the general behavior of liquid films will be seen to agree with the requirements of theory.

The formation of a liquid film takes place most symmetrically when a bubble of air rises to the top of a mass of the liquid. The motion of the liquid, as it is displaced by the bubble, is evidently such as to stretch the two surfaces in which the liquid meets the air, where these surfaces approach one another. This will cause an increase of tension, which will tend to restrain the extension of the surfaces. The extent to which this effect is produced will vary with the nature of the liquid. Let us suppose that the case is one in which the liquid contains one or more components which, although constituting but a very small part of its mass, greatly reduce its tension. Such components will exist in excess on the surfaces of the liquid. In this case the restraint upon the extension of the surfaces will be considerable, and as the bubble of air rises above the general level of the liquid, the motion of the latter will consist largely of a running out from between the two surfaces. But this running out of the liquid will be greatly retarded by its viscosity as soon as it is reduced to the thickness of a film, and the effect of the extension of the surfaces in increasing their tension will become greater and more permanent as the quantity of liquid diminishes which is available for supplying the substances which go to form the increased surfaces.

We may form a rough estimate of the amount of motion which is possible for the interior of a liquid film, relatively to its exterior, by calculating the descent of water between parallel vertical planes at which the motion of the water is reduced to zero. If we use the coefficient of viscosity as determined by Helmholtz and Piotrowski,\* we obtain

$$V = 581 D^2, \quad (656)$$

where  $V$  denotes the mean velocity of the water (*i. e.*, that velocity

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\* Sitzungsberichte der Wiener Akademie, (mathemat.-naturwiss. Classe), B. xl, S. 607. The calculation of formula (656) and that of the factor ( $\frac{8}{3}$ ) applied to the formula of Poiseuille, to adapt it to a current between plane surfaces, have been made by means of the general equations of the motion of a viscous liquid as given in the memoir referred to.

which, if it were uniform throughout the whole space between the fixed planes, would give the same discharge of water as the actual variable velocity) expressed in millimetres per second, and  $D$  denotes the distance in millimetres between the fixed planes, which is supposed to be very small in proportion to their other dimensions. This is for the temperature of  $24.5^{\circ}$  C. For the same temperature, the experiments of Poisenille\* give

$$V = 337 D^2$$

for the descent of water in long capillary tubes, which is equivalent to

$$V = 899 D^2 \quad (657)$$

for descent between parallel planes. The numerical coefficient in this equation differs considerably from that in (656), which is derived from experiments of an entirely different nature, but we may at least conclude that in a film of a liquid which has a viscosity and specific gravity not very different from those of water at the temperature mentioned the mean velocity of the interior relatively to the surfaces will not probably exceed  $1000 D^2$ . This is a velocity of  $.1^{\text{mm}}$  per second for a thickness of  $.01^{\text{mm}}$ ,  $.06^{\text{mm}}$  per minute for a thickness of  $.001$  (which corresponds to the red of the fifth order in a film of water), and  $.036^{\text{mm}}$  per hour for a thickness of  $.0001^{\text{mm}}$  (which corresponds to the white of the first order). Such an internal current is evidently consistent with great persistence of the film, especially in those cases in which the film can exist in a state of the greatest tenuity. On the other hand, the above equations give so large a value of  $V$  for thicknesses of  $1^{\text{mm}}$  or  $.1^{\text{mm}}$ , that the film can evidently be formed without carrying up any great weight of liquid, and any such thicknesses as these can have only a momentary existence.

A little consideration will show that the phenomenon is essentially of the same nature when films are formed in any other way, as by dipping a ring or the mouth of a cup in the liquid and then withdrawing it. When the film is formed in the mouth of a pipe, it may sometimes be extended so as to form a large bubble. Since the elasticity (*i. e.*, the increase of the tension with extension) is greater in the thinner parts, the thicker parts will be most extended, and the effect of this process (so far as it is not modified by gravity) will be to diminish the ratio of the greatest to the least thickness of the film. During this extension, as well as at other times, the increased elasticity due to imperfect communication of heat, etc., will serve to protect the bubble from fracture by shocks received from the air or the

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\* *Ibid.*, p. 653; or *Mémoires des Savants Étrangers*, vol. ix, p. 532.

pipe. If the bubble is now laid upon a suitable support, the condition (613) will be realized almost instantly. The bubble will then tend toward conformity with condition (614), the lighter portions rising to the top, more or less slowly, according to the viscosity of the film. The resulting difference of thickness between the upper and the lower parts of the bubble is due partly to the greater tension to which the upper parts are subject, and partly to a difference in the matter of which they are composed. When the film has only two components of which the potentials are not determined by the contiguous atmosphere, the laws which govern the arrangement of the elements of the film may be very simply expressed. If we call these components  $S_1$  and  $S_2$ , the latter denoting (as on page 469) that which exists in excess at the surface, one element of the film will tend toward the same level with another, or a higher, or a lower level, according as the quantity of  $S_2$  bears the same ratio to the quantity of  $S_1$  in the first element as in the second, or a greater, or a less ratio.

When a film, however formed, satisfies both the conditions (613) and (614), its thickness being sufficient for its interior to have the properties of matter in mass, the interior will still be subject to the slow current which we have already described, if it is truly fluid, however great its viscosity may be. It seems probable, however, that this process is often totally arrested by a certain gelatinous consistency of the mass in question, in virtue of which, although practically fluid in its behavior with reference to ordinary stresses, it may have the properties of a solid with respect to such very small stresses as those which are caused by gravity in the interior of a very thin film which satisfies the conditions (613) and (614).

However this may be, there is another cause which is often more potent in producing changes in a film, when the conditions just mentioned are approximately satisfied, than the action of gravity on its interior. This will be seen if we turn our attention to the edge where the film is terminated. At such an edge we generally find a liquid mass, continuous in phase with the interior of the film, which is bounded by concave surfaces, and in which the pressure is therefore less than in the interior of the film. This liquid mass therefore exerts a strong suction upon the interior of the film, by which its thickness is rapidly reduced. This effect is best seen when a film which has been formed in a ring is held in a vertical position. Unless the film is very viscous, its diminished thickness near the edge causes a rapid upward current on each side, while the central portion slowly

descends. Also at the bottom of the film, where the edge is nearly horizontal, portions which have become thinned escape from their position of unstable equilibrium beneath heavier portions, and pass upwards, traversing the central portion of the film until they find a position of stable equilibrium. By these processes, the whole film is rapidly reduced in thickness.

The energy of the suction which produces these effects may be inferred from the following considerations. The pressure in the slender liquid mass which encircles the film is of course variable, being greater in the lower portions than in the upper, but it is everywhere less than the pressure of the atmosphere. Let us take a point where the pressure is less than that of the atmosphere by an amount represented by a column of the liquid one centimetre in height. (It is probable that much greater differences of pressure occur.) At a point near by in the interior of the film the pressure is that of the atmosphere. Now if the difference of pressure of these two points were distributed uniformly through the space of one centimetre, the intensity of its action would be exactly equal to that of gravity. But since the change of pressure must take place very suddenly (in a small fraction of a millimetre), its effect in producing a current in a limited space must be enormously great compared with that of gravity.

Since the process just described is connected with the descent of the liquid in the mass encircling the film, we may regard it as another example of the downward tendency of the interior of the film. There is a third way in which this descent may take place, when the principal component of the interior is volatile, viz., through the air. Thus, in the case of a film of soap-water, if we suppose the atmosphere to be of such humidity that the potential for water at a level mid-way between the top and bottom of the film has the same value in the atmosphere as in the film, it may easily be shown that evaporation will take place in the upper portions and condensation in the lower. These processes, if the atmosphere were otherwise undisturbed, would occasion currents of diffusion and other currents, the general effect of which would be to carry the moisture downward. Such a precise adjustment would be hardly attainable, and the processes described would not be so rapid as to have a practical importance.

But when the potential for water in the atmosphere differs considerably from that in the film, as in the case of a film of soap-water in a dry atmosphere, or a film of soap-water with glycerine in a moist

atmosphere, the effect of evaporation or condensation is not to be neglected. In the first case, the diminution of the thickness of the film will be accelerated, in the second, retarded. In the case of the film containing glycerine, it should be observed that the water condensed cannot in all respects replace the fluid carried down by the internal current but that the two processes together will tend to wash out the glycerine from the film.

But when a component which greatly diminishes the tension of the film, although forming but a small fraction of its mass, (therefore existing in excess at the surface,) is volatile, the effect of evaporation and condensation may be considerable, even when the mean value of the potential for that component is the same in the film as in the surrounding atmosphere. To illustrate this, let us take the simple case of two components  $S_1$  and  $S_2$ , as before. (See page 469.) It appears from equation (508) that the potentials must vary in the film with the height  $z$ , since the tension does, and from (98) that these variations must (very nearly) satisfy the relation

$$\gamma_1 \frac{d\mu_1}{dz} + \gamma_2 \frac{d\mu_2}{dz} = 0, \quad (658)$$

$\gamma_1$  and  $\gamma_2$  denoting the densities of  $S_1$  and  $S_2$  in the interior of the film. The variation of the potential of  $S_2$  as we pass from one level to another is therefore as much more rapid than that of  $S_1$ , as its density in the interior of the film is less. If then the resistances restraining the evaporation, transmission through the atmosphere, and condensation of the two substances are the same, these processes will go on much more rapidly with respect to  $S_2$ . It will be observed that the values of  $\frac{d\mu_1}{dz}$  and  $\frac{d\mu_2}{dz}$  will have opposite signs,

the tendency of  $S_1$  being to pass down through the atmosphere, and that of  $S_2$  to pass up. Moreover, it may easily be shown that the evaporation or condensation of  $S_2$  will produce a very much greater effect than the evaporation or condensation of the same quantity of  $S_1$ . These effects are really of the same kind. For if condensation of  $S_2$  takes place at the top of the film, it will cause a diminution of tension, and thus occasion an extension of this part of the film, by which its thickness will be reduced, as it would be by evaporation of  $S_1$ . We may infer that it is a general condition of the persistence of liquid films, that the substance which causes the diminution of tension in the upper parts of the film must not be volatile.

But apart from any action of the atmosphere, we have seen that a

film which is truly fluid in its interior is in general subject to a continual diminution of thickness by the internal currents due to gravity and the suction at its edge. Sooner or later, the interior will somewhere cease to have the properties of matter in mass. The film will then probably become unstable with respect to a flux of the interior (see page 473), the thinnest parts tending to become still more thin (apart from any external cause) very much as if there were an attraction between the surfaces of the film, insensible at greater distances, but becoming sensible when the thickness of the film is sufficiently reduced. We should expect this to determine the rupture of the film, and such is doubtless the case with most liquids. In a film of soap-water, however, the rupture does not take place, and the processes which go on can be watched. It is apparent even to a very superficial observation that a film of which the tint is approaching the black exhibits a remarkable instability. The continuous change of tint is interrupted by the breaking out and rapid extension of black spots. That in the formation of these bright spots a separation of different substances takes place, and not simply an extension of a part of the film, is shown by the fact that the film is made thicker at the edge of these spots.

This is very distinctly seen in a plane vertical film, when a single black spot breaks out and spreads rapidly over a considerable area which was before of a nearly uniform tint approaching the black. The edge of the black spot as it spreads is marked as it were by a string of bright beads, which unite together on touching, and thus becoming larger, glide down across the bands of color below. Under favorable circumstances, there is often quite a shower of these bright spots. They are evidently small spots very much thicker—apparently many times thicker—than the part of the film out of which they are formed. Now if the formation of the black spots were due to a simple extension of the film, it is evident that no such appearance would be presented. The thickening of the edge of the film cannot be accounted for by *contraction*. For an extension of the upper portion of the film and contraction of the lower and thicker portion, with descent of the intervening portions, would be far less resisted by viscosity, and far more favored by gravity than such extensions and contractions as would produce the appearances described. But the rapid formation of a thin spot by an internal current would cause an accumulation at the edge of the spot of the material forming the interior of the film, and necessitate a thickening of the film in that place.

That which is most difficult to account for in the formation of the black spots is the arrest of the process by which the film grows thinner. It seems most natural to account for this, *if possible*, by passive resistance to motion due to a very viscous or gelatinous condition of the film. For it does not seem likely that the film, after becoming unstable by the flux of matter from its interior, would become stable (without the support of such resistance) by a continuance of the same process. On the other hand, gelatinous properties are very marked in soap-water which contains somewhat more soap than is best for the formation of films, and it is entirely natural that, even when such properties are wanting in the interior of a mass or thick film of a liquid, they may still exist in the immediate vicinity of the surface (where we know that the soap or some of its components exists in excess), or throughout a film which is so thin that the interior has ceased to have the properties of matter in mass.\* But these considerations do not amount to any *a priori* probability of an arrest of the tendency toward an internal current between adjacent elements of a black spot which may differ slightly in thickness, in time to prevent rupture of the film. For, in a thick film, the increase of the tension with the extension, which is necessary for its stability with respect to extension, is connected with an excess of the soap (or of some of its components) at the surface as compared with the interior of the film. With respect to the black spots, although the interior has ceased to have the properties of matter in mass, and any quantitative determinations derived from the surfaces of a mass of the liquid will not be applicable, it is natural to account for the stability with reference to extension by supposing that the same general difference of composition still exists. If therefore we account for the arrest of internal currents by the increasing density of soap or some of its components in the interior of the film, we must still suppose that the characteristic difference of composition in the interior and surface of the film has not been obliterated.

The preceding discussion relates to liquid films between masses of gas. Similar considerations will apply to liquid films between other liquids or between a liquid and a gas, and to films of gas between

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\*The experiments of M. Plateau (chapter VII of the work already cited) show that this is the case to a very remarkable degree with respect to a solution of saponine. With respect to soap-water, however, they do not indicate any greater superficial viscosity than belongs to pure water. But the resistance to an internal current, such as we are considering, is not necessarily measured by the resistance to such motions as those of the experiments referred to.

masses of liquid. The latter may be formed by gently depositing a liquid drop upon the surface of a mass of the same or a different liquid. This may be done (with suitable liquids) so that the continuity of the air separating the liquid drop and mass is not broken, but a film of air is formed, which, if the liquids are similar, is a counterpart of the liquid film which is formed by a bubble of air rising to the top of a mass of the liquid. (If the bubble has the same volume as the drop, the films will have precisely the same form, as well as the rest of the surfaces which bound the bubble and the drop.) Sometimes, when the weight and momentum of the drop carry it through the surface of the mass on which it falls, it appears surrounded by a complete spherical film of air, which is the counterpart on a small scale of a soap-bubble hovering in air.\* Since, however, the substance to which the necessary differences of tension in the film are mainly due is a component of the liquid masses on each side of the air film, the necessary differences of the potential of this substance cannot be permanently maintained, and these films have little persistence compared with films of soap-water in air. In this respect, the ease of these air-films is analogous to that of liquid films in an atmosphere containing substances by which their tension is greatly reduced. Compare page 479.

#### *Surfaces of Discontinuity between Solids and Fluids.*

We have hitherto treated of surfaces of discontinuity on the supposition that the contiguous masses are fluid. This is by far the most simple case for any rigorous treatment, since the masses are necessarily isotropic both in nature and in their state of strain. In this case, moreover, the mobility of the masses allows a satisfactory experimental verification of the mechanical conditions of equilibrium. On the other hand, the rigidity of solids is in general so great, that any tendency of the surfaces of discontinuity to variation in area or form may be neglected in comparison with the forces which are produced in the interior of the solids by any sensible strains, so that it is not generally necessary to take account of the surfaces of discontinuity in determining the state of strain of solid masses. But we must take account of the nature of the surfaces of discontinuity

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\* These spherical air-films are easily formed in soap-water. They are distinguishable from ordinary air-bubbles by their general behavior and by their appearance. The two concentric spherical surfaces are distinctly seen, the diameter of one appearing to be about three-quarters as large as that of the other. This is of course an optical illusion, depending upon the index of refraction of the liquid.

between solids and fluids with reference to the tendency toward solidification or dissolution at such surfaces, and also with reference to the tendencies of different fluids to spread over the surfaces of solids.

Let us therefore consider a surface of discontinuity between a fluid and a solid, the latter being either isotropic or of a continuous crystalline structure, and subject to any kind of stress compatible with a state of mechanical equilibrium with the fluid. We shall not exclude the case in which substances foreign to the contiguous masses are present in small quantities at the surface of discontinuity, but we shall suppose that the nature of this surface (*i. e.*, of the non-homogeneous film between the approximately homogeneous masses), is entirely determined by the nature and state of the masses which it separates, and the quantities of the foreign substances which may be present. The notions of the *dividing surface*, and of the *superficial densities* of energy, entropy, and the several components, which we have used with respect to surfaces of discontinuity between fluids (see pages 380 and 386), will evidently apply without modification to the present case. We shall use the suffix <sub>(1)</sub> with reference to the substance of the solid, and shall suppose the dividing surface to be determined so as to make the superficial density of this substance vanish. The superficial densities of energy, of entropy, and of the other component substances may then be denoted by our usual symbols (see page 397),

$$\varepsilon_{S(1)}, \quad \eta_{S(1)}, \quad \Gamma_{2(1)}, \quad \Gamma_{3(1)}, \quad \text{etc.}$$

Let the quantity  $\sigma$  be defined by the equation

$$\sigma = \varepsilon_{S(1)} - t \eta_{S(1)} - \mu_2 \Gamma_{2(1)} - \mu_3 \Gamma_{3(1)} - \text{etc.}, \quad (659)$$

in which  $t$  denotes the temperature, and  $\mu_2, \mu_3$ , etc. the potentials for the substances specified at the surface of discontinuity.

As in the case of two fluid masses, (see page 421,) we may regard  $\sigma$  as expressing the work spent in forming a unit of the surface of discontinuity—under certain conditions, which we need not here specify—but it cannot properly be regarded as expressing the tension of the surface. The latter quantity depends upon the work spent in stretching the surface, while the quantity  $\sigma$  depends upon the work spent in forming the surface. With respect to perfectly fluid masses, these processes are not distinguishable, unless the surface of discontinuity has components which are not found in the contiguous masses, and even in this case, (since the surface must be supposed to be formed out of matter supplied at the same potentials which belong to the matter in the surface,) the work spent in increasing the surface infinitesi-

mally by stretching is identical with that which must be spent in forming an equal infinitesimal amount of new surface. But when one of the masses is solid, and its states of strain are to be distinguished, there is no such equivalence between the stretching of the surface and the forming of new surface.\*

With these preliminary notions, we now proceed to discuss the condition of equilibrium which relates to the dissolving of a solid at the surface where it meets a fluid, when the thermal and mechanical conditions of equilibrium are satisfied. It will be necessary for us to consider the case of isotropic and of crystallized bodies separately, since in the former the value of  $\sigma$  is independent of the direction of the surface, except so far as it may be influenced by the state of strain of the solid, while in the latter the value of  $\sigma$  varies greatly with the direction of the surface with respect to the axes of crystallization, and in such a manner as to have a large number of sharply defined minima.† This may be inferred from the phenomena which crystalline bodies present, as will appear more distinctly in the following discussion. Accordingly, while a variation in the direction of an

\* This will appear more distinctly if we consider a particular case. Let us consider a thin plane sheet of a crystal in a vacuum (which may be regarded as a limiting case of a very attenuated fluid), and let us suppose that the two surfaces of the sheet are alike. By applying the proper forces to the edges of the sheet, we can make all stress vanish in its interior. The *tensions* of the two surfaces, are in equilibrium with these forces, and are measured by them. But the tensions of the surfaces, thus determined, may evidently have different values in different directions, and are entirely different from the quantity which we denote by  $\sigma$ , which represents the work required to form a unit of the surface by any reversible process, and is not connected with any idea of direction.

In certain cases, however, it appears probable that the values of  $\sigma$  and of the superficial tension will not greatly differ. This is especially true of the numerous bodies which, although generally (and for many purposes properly) regarded as solids, are really very viscous fluids. Even when a body exhibits no fluid properties at its actual temperature, if its surface has been formed at a higher temperature, at which the body was fluid, and the change from the fluid to the solid state has been by insensible gradations, we may suppose that the value of  $\sigma$  coincided with the superficial tension until the body was decidedly solid, and that they will only differ so far as they may be differently affected by subsequent variations of temperature and of the stresses applied to the solid. Moreover, when an amorphous solid is in a state of equilibrium with a solvent, although it may have no fluid properties in its interior, it seems not improbable that the particles at its surface, which have a greater degree of mobility, may so arrange themselves that the value of  $\sigma$  will coincide with the superficial tension, as in the case of fluids.

† The differential coefficients of  $\sigma$  with respect to the direction-cosines of the surface appear to be discontinuous functions of the latter quantities.

element of the surface may be neglected (with respect to its effect on the value of  $\sigma$ ) in the case of isotropic solids, it is quite otherwise with crystals. Also, while the surfaces of equilibrium between fluids and soluble isotropic solids are without discontinuities of direction, being in general curved, a crystal in a state of equilibrium with a fluid in which it can dissolve is bounded in general by a broken surface consisting of sensibly plane portions.

For isotropic solids, the conditions of equilibrium may be deduced as follows. If we suppose that the solid is unchanged, except that an infinitesimal portion is dissolved at the surface where it meets the fluid, and that the fluid is considerable in quantity and remains homogeneous, the increment of energy in the vicinity of the surface will be represented by the expression

$$\int [\varepsilon_v' - \varepsilon_v'' + (c_1 + c_2) \varepsilon_{s(1)}] \delta N Ds$$

where  $Ds$  denotes an element of the surface,  $\delta N$  the variation in its position (measured normally, and regarded as *negative* when the solid is dissolved),  $c_1$  and  $c_2$  its principal curvatures (positive when their centers lie on the same side as the solid),  $\varepsilon_{s(1)}$ , the surface-density of energy,  $\varepsilon_v'$  and  $\varepsilon_v''$  the volume-densities of energy in the solid and fluid respectively, and the sign of integration relates to the elements  $Ds$ . In like manner, the increments of entropy and of the quantities of the several components in the vicinity of the surface will be

$$\begin{aligned} & \int [\eta_v' - \eta_v'' + (c_1 + c_2) \eta_{s(1)}] \delta N Ds, \\ & \int [\gamma_1' - \gamma_1''] \delta N Ds, \\ & \int [-\gamma_2'' + (c_1 + c_2) \Gamma_{s(1)}] \delta N Ds, \\ & \text{etc.} \end{aligned}$$

The entropy and the matter of different kinds represented by these expressions we may suppose to be derived from the fluid mass. These expressions, therefore, with a change of sign, will represent the increments of entropy and of the quantities of the components in the whole space occupied by the fluid except that which is immediately contiguous to the solid. Since this space may be regarded as constant, the increment of energy in this space may be obtained [according to equation (12)] by multiplying the above expression relating to entropy by  $-t$ , and those relating to the components by  $-\mu_1'', -\mu_2$ , etc.,\* and taking the sum. If to this

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\* The potential  $\mu_1''$  is marked by double accents in order to indicate that its value is to be determined in the fluid mass, and to distinguish it from the potential  $\mu'$ .

we add the above expression for the increment of energy near the surface, we obtain the increment of energy for the whole system. Now by (93) we have

$$p'' = -\varepsilon_v'' + t\eta_v'' + \mu_1''\gamma_1'' + \mu_2''\gamma_2'' + \text{etc.}$$

By this equation and (659), our expression for the total increment of energy in the system may be reduced to the form

$$\gamma'[\varepsilon_v' - t\eta_v' - \mu_1'\gamma_1' + p'' + (c_1 + c_2)\sigma]\delta NDs. \quad (660)$$

In order that this shall vanish for any values of  $\delta N$ , it is necessary that the coefficient of  $\delta NDs$  shall vanish. This gives for the condition of equilibrium

$$\mu_1'' = \frac{\varepsilon_v' - t\eta_v' + p'' + (c_1 + c_2)\sigma}{\gamma_1'}. \quad (661)$$

This equation is identical with (387), with the exception of the term containing  $\sigma$ , which vanishes when the surface is plane.\*

We may also observe that when the solid has no stresses except an isotropic pressure, if the quantity represented by  $\sigma$  is equal to the true tension of the surface,  $p'' + (c_1 + c_2)\sigma$  will represent the pressure in the interior of the solid, and the second member of the equation will represent [see equation (93)] the value of the potential in the solid for the substance of which it consists. In this case, therefore, the equation reduces to

$$\mu_1'' = \mu_1',$$

that is, it expresses the equality of the potentials for the substance of the solid in the two masses—the same condition which would subsist if both masses were fluid.

Moreover, the compressibility of all solids is so small that, although  $\sigma$  may not represent the true tension of the surface, nor  $p'' + (c_1 + c_2)\sigma$  the true pressure in the solid when its stresses are isotropic, the quantities  $\varepsilon_v'$  and  $\eta_v'$  if calculated for the pressure  $p'' + (c_1 + c_2)\sigma$  with the actual temperature will have sensibly the same values as if calculated for the true pressure of the solid. Hence, the second member

relating to the solid mass (when this is in a state of isotropic stress), which, as we shall see, may not always have the same value. The other potentials  $\mu_2$ , etc., have the same values as in (659), and consist of two classes, one of which relates to substances which are components of the fluid mass, (these might be marked by the double accents,) and the other relates to substances found only at the surface of discontinuity. The expressions to be multiplied by the potentials of this latter class all have the value zero.

\* In equation (387), the density of the solid is denoted by  $\Gamma$ , which is therefore equivalent to  $\gamma_1'$  in (661).

of equation (661), when the stresses of the solid are sensibly isotropic, is sensibly equal to the potential of the same body at the same temperature but with the pressure  $p'' + (c_1 + c_2) \sigma$ , and the condition of equilibrium with respect to dissolving for a solid of isotropic stresses may be expressed with sufficient accuracy by saying that the potential for the substance of the solid in the fluid must have this value. In like manner, when the solid is not in a state of isotropic stress, the difference of the two pressures in question will not sensibly affect the values of  $\varepsilon_v'$  and  $\eta_v'$ , and the value of the second member of the equation may be calculated as if  $p'' + (c_1 + c_2) \sigma$  represented the true pressure in the solid in the direction of the normal to the surface. Therefore, if we had taken for granted that the quantity  $\sigma$  represents the tension of a surface between a solid and a fluid, as it does when both masses are fluid, this assumption would not have led us into any practical error in determining the value of the potential  $\mu_1''$  which is necessary for equilibrium. On the other hand, if in the case of any amorphous body the value of  $\sigma$  differs notably from the true surface-tension, the latter quantity substituted for  $\sigma$  in (661) will make the second member of the equation equal to the true value of  $\mu_1'$ , when the stresses are isotropic, but this will not be equal to the value of  $\mu_1''$  in case of equilibrium, unless  $c_1 + c_2 = 0$ .

When the stresses in the solid are not isotropic, equation (661) may be regarded as expressing the condition of equilibrium with respect to the dissolving of the solid, and is to be distinguished from the condition of equilibrium with respect to an increase of solid matter, since the new matter would doubtless be deposited in a state of isotropic stress. (The case would of course be different with crystalline bodies, which are not considered here.) The value of  $\mu_1''$  necessary for equilibrium with respect to the formation of new matter is a little less than that necessary for equilibrium with respect to the dissolving of the solid. In regard to the actual behavior of the solid and fluid, all that the theory enables us to predict with certainty is that the solid will not dissolve if the value of the potential  $\mu_1''$  is greater than that given by the equation for the solid with its distorting stresses, and that new matter will not be formed if the value of  $\mu_1''$  is less than the same equation would give for the case of the solid with isotropic stresses.\* It seems probable, however, that

\* The possibility that the new solid matter might differ in composition from the original solid is here left out of account. This point has been discussed on pages 134-137, but without reference to the state of strain of the solid or the influence of the curvature of the surface of discontinuity. The statement made above may be

if the fluid in contact with the solid is not renewed, the system will generally find a state of equilibrium in which the outermost portion of the solid will be in a state of isotropic stress. If at first the solid should dissolve, this would supersaturate the fluid, perhaps until a state is reached satisfying the condition of equilibrium with the stressed solid, and then, if not before, a deposition of solid matter in a state of isotropic stress would be likely to commence and go on until the fluid is reduced to a state of equilibrium with this new solid matter.

The action of gravity will not affect the nature of the condition of equilibrium for any single point at which the fluid meets the solid, but it will cause the values of  $p''$  and  $\mu_1''$  in (661) to vary according to the laws expressed by (612) and (617). If we suppose that the outer part of the solid is in a state of isotropic stress, which is the most important case, since it is the only one in which the equilibrium is in every sense stable, we have seen that the condition (661) is at least sensibly equivalent to this:—that the potential for the substance of the solid which would belong to the solid mass at the temperature  $t$  and the pressure  $p'' + (c_1 + c_2)\sigma$  must be equal to  $\mu_1''$ . Or, if we denote by  $(p')$  the pressure belonging to solid with the temperature  $t$  and the potential equal to  $\mu_1''$ , the condition may be expressed in the form

$$(p') = p'' + (c_1 + c_2)\sigma. \quad (662)$$

Now if we write  $\gamma''$  for the total density of the fluid, we have by (612)

$$dp'' = -g\gamma'' dz.$$

By (98)  $d(p') = \gamma_1' d\mu_1'',$

and by (617)  $d\mu_1'' = -g dz;$

whence  $d(p') = -g\gamma_1' dz.$

Accordingly we have

$$d(p') - dp'' = g(\gamma'' - \gamma_1') dz,$$

and

$$(p') - p'' = g(\gamma'' - \gamma_1') z,$$

$z$  being measured from the horizontal plane for which  $(p') = p''$ .

Substituting this value in (662), we obtain

$$c_1 + c_2 = \frac{g(\gamma'' - \gamma_1')}{\sigma} z, \quad (663)$$

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generalized so as to hold true of the formation of new solid matter of any kind on the surface as follows:—that new solid matter of any kind will not be formed upon the surface (with more than insensible thickness), if the second member of (661) calculated for such new matter is greater than the potential in the fluid for such matter.

precisely as if both masses were fluid, and  $\sigma$  denoted the tension of their common surface, and  $(p')$  the true pressure in the mass specified. [Compare (619).]

The obstacles to an exact experimental realization of these relations are very great, principally from the want of absolute uniformity in the internal structure of amorphous solids, and on account of the passive resistances to the processes which are necessary to bring about a state satisfying the conditions of theoretical equilibrium, but it may be easy to verify the general tendency toward diminution of surface, which is implied in the foregoing equations.\*

Let us apply the same method to the case in which the solid is a crystal. The surface between the solid and fluid will now consist of plane portions, the directions of which may be regarded as invari-

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\* It seems probable that a tendency of this kind plays an important part in some of the phenomena which have been observed with respect to the freezing together of pieces of ice. (See especially Professor Faraday's "Note on Regelation" in the *Proceedings of the Royal Society*, vol. x, p. 440; or in the *Philosophical Magazine*, 4th ser., vol. xxi, p. 146.) Although this is a body of crystalline structure, and the action which takes place is doubtless influenced to a certain extent by the directions of the axes of crystallization, yet, since the phenomena have not been observed to depend upon the orientation of the pieces of ice, we may conclude that the effect, so far as its general character is concerned, is such as might take place with an isotropic body. In other words, for the purposes of a general explanation of the phenomena we may neglect the differences in the values of  $\sigma_{IW}$  (the suffixes are used to indicate that the symbol relates to the surface between ice and water) for different orientations of the axes of crystallization, and also neglect the influence of the surface of discontinuity with respect to crystalline structure, which must be formed by the freezing together of the two masses of ice when the axes of crystallization in the two masses are not similarly directed. In reality, this surface—or the necessity of the formation of such a surface if the pieces of ice freeze together—must exert an influence adverse to their union, measured by a quantity  $\sigma_{II}$ , which is determined for this surface by the same principles as when one of two contiguous masses is fluid, and varies with the orientations of the two systems of crystallographic axes relatively to each other and to the surface. But under the circumstances of the experiment, since we may neglect the possibility of the two systems of axes having precisely the same directions, this influence is probably of a tolerably constant character, and is evidently not sufficient to alter the general nature of the result. In order wholly to prevent the tendency of pieces of ice to freeze together, when meeting in water with curved surfaces and without pressure, it would be necessary that  $\sigma_{II} \geq 2\sigma_{IW}$ , except so far as the case is modified by passive resistances to change, and by the inequality in the values of  $\sigma_{II}$  and  $\sigma_{IW}$  for different directions of the axes of crystallization.

It will be observed that this view of the phenomena is in harmony with the opinion of Professor Faraday. With respect to the union of pieces of ice as an indirect consequence of pressure, see page 198 of volume xi of the *Proceedings of the Royal Society*; or the *Philosophical Magazine*, 4th ser., vol. xxiii, p. 407.

able. If the crystal grows on one side a distance  $\delta N$ , without other change, the increment of energy in the vicinity of the surface will be

$$(\varepsilon_v' - \varepsilon_v'') s \delta N + \Sigma' (\varepsilon_{s(1)}' l' \operatorname{cosec} \omega' - \varepsilon_{s(1)}' l' \operatorname{cot} \omega') \delta N,$$

where  $\varepsilon_v'$  and  $\varepsilon_v''$  denote the volume-densities of energy in the crystal and fluid respectively,  $s$  the area of the side on which the crystal grows,  $\varepsilon_{s(1)}$  the surface-density of energy on that side,  $\varepsilon_{s(1)}'$  the surface-density of energy on an adjacent side,  $\omega'$  the external angle of these two sides,  $l'$  their common edge, and the symbol  $\Sigma'$  a summation with respect to the different sides adjacent to the first. The increments of entropy and of the quantities of the several components will be represented by analogous formulæ, and if we deduce as on pages 485, 486 the expression for the increase of energy in the whole system due to the growth of the crystal without change of the total entropy or volume, and set this expression equal to zero, we shall obtain for the condition of equilibrium

$$(\varepsilon_v' - t \eta_v' - \mu_1'' \gamma_1' + p'') s \delta N + \Sigma' (\sigma' l' \operatorname{cosec} \omega' - \sigma l' \operatorname{cot} \omega') \delta N = 0, \quad (664)$$

where  $\sigma$  and  $\sigma'$  relate respectively to the same sides as  $\varepsilon_{s(1)}$  and  $\varepsilon_{s(1)}'$  in the preceding formula. This gives

$$\mu_1'' = \frac{\varepsilon_v' - t \eta_v' + p''}{\gamma_1'} + \frac{\Sigma' (\sigma' l' \operatorname{cosec} \omega' - \sigma l' \operatorname{cot} \omega')}{s \gamma_1'}. \quad (665)$$

It will be observed that unless the side especially considered is small or narrow, we may neglect the second fraction in this equation, which will then give the same value of  $\mu_1''$  as eqn. (387), or as equation (661) applied to a plane surface.

Since a similar equation must hold true with respect to every other side of the crystal of which the equilibrium is not affected by meeting some other body, the condition of equilibrium for the crystalline form (when unaffected by gravity) is that the expression

$$\frac{\Sigma' (\sigma' l' \operatorname{cosec} \omega' - \sigma l' \operatorname{cot} \omega')}{s} \quad (666)$$

shall have the same value for each side of the crystal. (By the value of this expression for any side of the crystal is meant its value when  $\sigma$  and  $s$  are determined by that side and the other quantities by the surrounding sides in succession in connection with the first side.) This condition will not be affected by a change in the size of a crystal while its proportions remain the same. But the tendencies of similar crystals toward the form required by this condition, as measured by the inequalities in the composition or the temperature of

the surrounding fluid which would counterbalance them, will be inversely as the linear dimensions of the crystals, as appears from the preceding equation.

If we write  $v$  for the volume of a crystal, and  $\Sigma(\sigma s)$  for the sum of the areas of all its sides multiplied each by the corresponding value of  $\sigma$ , the numerator and denominator of the fraction (666), multiplied each by  $\delta N$ , may be represented by  $\delta\Sigma(\sigma s)$  and  $\delta v$  respectively. The value of the fraction is therefore equal to that of the differential coefficient

$$\frac{d\Sigma(\sigma s)}{dv}$$

as determined by the displacement of a particular side while the other sides are fixed. The condition of equilibrium for the form of a crystal (when the influence of gravity may be neglected) is that the value of this differential coefficient must be independent of the particular side which is supposed to be displaced. For a constant volume of the crystal,  $\Sigma(\sigma s)$  has therefore a minimum value when the condition of equilibrium is satisfied, as may easily be proved more directly.

When there are no foreign substances at the surfaces of the crystal, and the surrounding fluid is indefinitely extended, the quantity  $\Sigma(\sigma s)$  represents the work required to form the surfaces of the crystal, and the coefficient of  $s \delta N$  in (664) with its sign reversed represents the work gained in forming a mass of volume unity like the crystal but regarded as without surfaces. We may denote the work required to form the crystal by

$$W_s - W_v,$$

$W_s$  denoting the work required to form the surfaces [*i. e.*,  $\Sigma(\sigma s)$ ], and  $W_v$  the work gained in forming the mass as distinguished from the surfaces. Equation (664) may then be written

$$-\delta W_v + \Sigma(\sigma \delta s) = 0. \quad (667)$$

Now (664) would evidently continue to hold true if the crystal were diminished in size, remaining similar to itself in form and in nature, if the values of  $\sigma$  in all the sides were supposed to diminish in the same ratio as the linear dimensions of the crystal. The variation of  $W_s$  would then be determined by the relation

$$dW_s = d\Sigma(\sigma s) = \frac{3}{2} \Sigma(\sigma ds),$$

and that of  $W_v$  by (667). Hence,

$$dW_s = \frac{3}{2} dW_v,$$

and, since  $W_s$  and  $W_v$  vanish together,

$$W_s = \frac{3}{2} W_v, \\ W_s - W_v = \frac{1}{3} W_s = \frac{1}{2} W_v, \quad (668)$$

—the same relation which we have before seen to subsist with respect to a spherical mass of fluid as well as in other cases. (See pages 421, 425, 465.)

The equilibrium of the crystal is unstable with respect to variations in size when the surrounding fluid is indefinitely extended, but it may be made stable by limiting the quantity of the fluid.

To take account of the influence of gravity, we must give to  $\mu_1''$  and  $\rho''$  in (665) their average values in the side considered. These coincide (when the fluid is in a state of internal equilibrium) with their values at the center of gravity of the side. The values of  $\gamma_1'$ ,  $\varepsilon_v'$ ,  $\eta_v'$  may be regarded as constant, so far as the influence of gravity is concerned. Now since by (612) and (617)

$$dp'' = -g \rho'' dz,$$

and

$$d\mu_1'' = -g dz,$$

we have

$$d(\gamma_1' \mu_1'' - p'') = g(\rho'' - \gamma_1') dz.$$

Comparing (664), we see that the upper or the lower faces of the crystal will have the greater tendency to grow, (other things being equal,) according as the crystal is lighter or heavier than the fluid. When the densities of the two masses are equal, the effect of gravity on the form of the crystal may be neglected.

In the preceding paragraph the fluid is regarded as in a state of internal equilibrium. If we suppose the composition and temperature of the fluid to be uniform, the condition which will make the effect of gravity vanish will be that

$$\frac{d(\gamma_1' \mu_1'' - p'')}{dz} = 0,$$

when the value of the differential coefficient is determined in accordance with this supposition. This condition reduces to

$$\left( \frac{d\mu_1}{dp} \right)_{t, m}'' = \frac{1}{\gamma_1'}, *$$

which, by equation (92), is equivalent to

$$\left( \frac{dv}{dm_1} \right)_{t, p, m}'' = \frac{1}{\gamma_1'}. \quad (669)$$

\* A suffixed  $m$  is used to represent all the symbols  $m_1$ ,  $m_2$ , etc., except such as may occur in the differential coefficient.

The tendency of a crystal to grow will be greater in the upper or lower parts of the fluid, according as the growth of a crystal at constant temperature and pressure will produce expansion or contraction.

Again, we may suppose the composition of the fluid and its entropy per unit of mass to be uniform. The temperature will then vary with the pressure, that is, with  $z$ . We may also suppose the temperature of different crystals or different parts of the same crystal to be determined by the fluid in contact with them. These conditions express a state which may perhaps be realized when the fluid is gently stirred. Owing to the differences of temperature we cannot regard  $\varepsilon_v'$  and  $\eta_v'$  in (664) as constant, but we may regard their variations as subject to the relation  $d\varepsilon_v' = t d\eta_v'$ . Therefore, if we make  $\eta_v' = 0$  for the mean temperature of the fluid, (which involves no real loss of generality,) we may treat  $\varepsilon_v' - t \eta_v'$  as constant. We shall then have for the condition that the effect of gravity shall vanish—

$$\frac{d(\gamma_1' \mu_1'' - p'')}{dz} = 0,$$

which signifies in the present case that

$$\left( \frac{d\mu_1}{dp} \right)_{\eta, m}'' = \frac{1}{\gamma_1'},$$

or, by (90),

$$\left( \frac{dv}{dm_1} \right)_{\eta, p, m}'' = \frac{1}{\gamma_1'} \quad (670)$$

Since the entropy of the crystal is zero, this equation expresses that the dissolving of a small crystal in a considerable quantity of the fluid will produce neither expansion nor contraction, when the pressure is maintained constant and no heat is supplied or taken away.

The manner in which crystals actually grow or dissolve is often principally determined by other differences of phase in the surrounding fluid than those which have been considered in the preceding paragraph. This is especially the case when the crystal is growing or dissolving rapidly. When the great mass of the fluid is considerably supersaturated, the action of the crystal keeps the part immediately contiguous to it nearer the state of exact saturation. The farthest projecting parts of the crystal will therefore be most exposed to the action of the supersaturated fluid, and will grow most rapidly. The same parts of a crystal will dissolve most rapidly in a fluid considerably below saturation.\*

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\* See O. Lehmann "Ueber das Wachsthum der Krystalle," *Zeitschrift für Krystallographie und Mineralogie*, Bd. i, S. 453; or the review of the paper in Wiedemann's *Beiblätter*, Bd. ii, S. 1.

But even when the fluid is supersaturated only so much as is necessary in order that the crystal shall grow at all, it is not to be expected that the form in which  $\Sigma(\sigma s)$  has a minimum value (or such a modification of that form as may be due to gravity or to the influence of the body supporting the crystal) will always be the ultimate result. For we cannot imagine a body of the internal structure and external form of a crystal to grow or dissolve by an entirely continuous process, or by a process in the same sense continuous as condensation or evaporation between a liquid and gas, or the corresponding processes between an amorphous solid and a fluid. The process is rather to be regarded as periodic, and the formula (664) cannot properly represent the true value of the quantities intended unless  $\delta N$  is equal to the distance between two successive layers of molecules in the crystal, or a multiple of that distance. Since this can hardly be treated as an infinitesimal, we can only conclude with certainty that sensible changes cannot take place for which the expression (664) would have a positive value.\*

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\* That it is necessary that certain relations shall be precisely satisfied in order that equilibrium may subsist between a liquid and gas with respect to evaporation, is explained (see Clausius "Ueber die Art der Bewegung, welche wir Wärme nennen," *Pogg. Ann.*, Bd. c, S. 353; or *Abhand. über die mech. Wärmetheorie*, XIV,) by supposing that a passage of individual molecules from the one mass to the other is continually taking place, so that the slightest circumstance may give the preponderance to the passage of matter in either direction. The same supposition may be applied, at least in many cases, to the equilibrium between amorphous solids and fluids. Also in the case of crystals in equilibrium with fluids, there may be a passage of individual molecules from one mass to the other, so as to cause insensible fluctuations in the mass of the solid. If these fluctuations are such as to cause the occasional deposit or removal of a whole layer of particles, the least cause would be sufficient to make the probability of one kind of change prevail over that of the other, and it would be necessary for equilibrium that the theoretical conditions deduced above should be precisely satisfied. But this supposition seems quite improbable, except with respect to a very small side.

The following view of the molecular state of a crystal when in equilibrium with respect to growth or dissolution appears as probable as any. Since the molecules at the corners and edges of a perfect crystal would be less firmly held in their places than those in the middle of a side, we may suppose that when the condition of theoretical equilibrium (665) is satisfied several of the outermost layers of molecules on each side of the crystal are incomplete toward the edges. The boundaries of these imperfect layers probably fluctuate, as individual molecules attach themselves to the crystal or detach themselves, but not so that a layer is entirely removed (on any side of considerable size), to be restored again simply by the irregularities of the motions of the individual molecules. Single molecules or small groups of molecules may indeed attach themselves to the side of the crystal but they will speedily be dislodged, and if any molecules are thrown out from the middle of a surface, these deficiencies

Let us now examine the special condition of equilibrium which relates to a line at which three different masses meet, when one or more of these masses is solid. If we apply the method of page 685 to a system containing such a line, it is evident that we shall obtain in the expression corresponding to (660), beside the integral relating to the surfaces, a term of the form

$$\int \Sigma(\sigma \delta T) Dl$$

to be interpreted as the similar term in (611), except so far as the definition of  $\sigma$  has been modified in its extension to solid masses. In order that this term shall be incapable of a negative value it is neces-

will also soon be made good; nor will the frequency of these occurrences be such as greatly to affect the general smoothness of the surfaces, except near the edges where the surfaces fall off somewhat, as before described. Now a continued growth on any side of a crystal is impossible unless new layers can be formed. This will require a value of  $\mu_1''$  which may exceed that given by equation (665) by a finite quantity. Since the difficulty in the formation of a new layer is at or near the commencement of the formation, the necessary value of  $\mu_1''$  may be independent of the area of the side, except when the side is very small. The value of  $\mu_1''$  which is necessary for the growth of the crystal will however be different for different kinds of surfaces, and probably will generally be greatest for the surfaces for which  $\sigma$  is least.

On the whole, it seems not improbable that the form of very minute crystals in equilibrium with solvents is principally determined by equation (665), (*i. e.*, by the condition that  $\Sigma(\sigma s)$  shall be a minimum for the volume of the crystal except so far as the case is modified by gravity or the contact of other bodies,) but as they grow larger (in a solvent no more supersaturated than is necessary to make them grow at all), the deposition of new matter on the different surfaces will be determined more by the nature (orientation) of the surfaces and less by their size and relations to the surrounding surfaces. As a final result, a large crystal, thus formed, will generally be bounded by those surfaces alone on which the deposit of new matter takes place least readily, with small, perhaps insensible truncations. If one kind of surfaces satisfying this condition cannot form a closed figure, the crystal will be bounded by two or three kinds of surfaces determined by the same condition. The kinds of surface thus determined will probably generally be those for which  $\sigma$  has the least values. But the relative development of the different kinds of sides, even if unmodified by gravity or the contact of other bodies, will not be such as to make  $\Sigma(\sigma s)$  a minimum. The growth of the crystal will finally be confined to sides of a single kind.

It does not appear that any part of the operation of removing a layer of molecules presents any especial difficulty so marked as that of commencing a new layer; yet the values of  $\mu_1''$  which will just allow the different stages of the process to go on must be slightly different, and therefore, for the continued dissolving of the crystal the value of  $\mu_1''$  must be less (by a finite quantity) than that given by equation (665). It seems probable that this would be especially true of those sides for which  $\sigma$  has the least values. The effect of dissolving a crystal (even when it is done as slowly as possible) is therefore to produce a form which probably differs from that of theoretical equilibrium in a direction opposite to that of a growing crystal.

sary that at every point of the line

$$\Sigma(\sigma \delta T) \geq 0 \quad (671)$$

for any *possible* displacement of the line. Those displacements are to be regarded as possible which are not prevented by the solidity of the masses, when the interior of every solid mass is regarded as incapable of motion. At the surfaces between solid and fluid masses, the processes of solidification and dissolution will be possible in some cases, and impossible in others.

The simplest case is when two masses are fluid and the third is solid and insoluble. Let us denote the solid by S, the fluids by A and B, and the angles filled by these fluids by  $\alpha$  and  $\beta$  respectively. If the surface of the solid is continuous at the line where it meets the two fluids, the condition of equilibrium reduces to

$$\sigma_{AB} \cos \alpha = \sigma_{BS} - \sigma_{AS}. \quad (672)$$

If the line where these masses meet is at an edge of the solid, the condition of equilibrium is that

$$\begin{aligned} \sigma_{AB} \cos \alpha &\leq \sigma_{BS} - \sigma_{AS}, \\ \text{and} \quad \sigma_{AB} \cos \beta &\leq \sigma_{AS} - \sigma_{BS}; \end{aligned} \quad \left. \right\} \quad (673)$$

which reduces to the preceding when  $\alpha + \beta = \pi$ . Since the displacement of the line can take place by a purely mechanical process, this condition is capable of a more satisfactory experimental verification than those conditions which relate to processes of solidification and dissolution. Yet the frictional resistance to a displacement of the line is enormously greater than in the case of three fluids, since the relative displacements of contiguous portions of matter are enormously greater. Moreover, foreign substances adhering to the solid are not easily displaced, and cannot be distributed by extensions and contractions of the surface of discontinuity, as in the case of fluid masses. Hence, the distribution of such substances is arbitrary to a greater extent than in the case of fluid masses, (in which a single foreign substance in any surface of discontinuity is uniformly distributed, and a greater number are at least so distributed as to make the tension of the surface uniform,) and the presence of these substances will modify the conditions of equilibrium in a more irregular manner.

If one or more of three surfaces of discontinuity which meet in a line divides an amorphous solid from a fluid in which it is soluble, such a surface is to be regarded as movable, and the particular conditions involved in (671) will be accordingly modified. If the soluble solid is a crystal, the case will properly be treated by the method used on page 490. The condition of equilibrium relating to the line

will not in this case be entirely separable from those relating to the adjacent surfaces, since a displacement of the line will involve a displacement of the whole side of the crystal which is terminated at this line. But the expression for the total increment of energy in the system due to any internal changes not involving any variation in the total entropy or volume will consist of two parts, of which one relates to the properties of the masses of the system, and the other may be expressed in the form

$$\delta \Sigma(\sigma s),$$

the summation relating to all the surfaces of discontinuity. This indicates the same tendency toward changes diminishing the value of  $\Sigma(\sigma s)$ , which appears in other cases.\*

*General Relations.*—For any constant state of strain of the surface of the solid, we may write

$$d\varepsilon_{s(1)} = t d\eta_{s(1)} + \mu_2 d\Gamma_{2(1)} + \mu_3 d\Gamma_{3(1)} + \text{etc.}, \quad (674)$$

since this relation is implied in the definition of the quantities involved. From this and (659) we obtain

$$d\sigma = -\eta_{s(1)} dt - \Gamma_{2(1)} d\mu_2 - \Gamma_{3(1)} d\mu_3 - \text{etc.}, \quad (675)$$

which is subject, in strictness, to the same limitation—that the state of strain of the surface of the solid remains the same. But this limitation may in most cases be neglected. (If the quantity  $\sigma$  represented the true tension of the surface, as in the case of a surface between fluids, the limitation would be wholly unnecessary.)

*Another method and notation.*—We have so far supposed that we have to do with a non-homogeneous film of matter between two homogeneous (or very nearly homogeneous) masses, and that the nature and state of this film is in all respects determined by the

\* The freezing together of wool and ice may be mentioned here. The fact that a fiber of wool which remains in contact with a block of ice under water will become attached to it seems to be strictly analogous to the fact that if a solid body be brought into such a position that it just touches the free surface of water, the water will generally rise up about the point of contact so as to touch the solid over a surface of some extent. The condition of the latter phenomenon is

$$\sigma_{SA} + \sigma_{WA} > \sigma_{SW},$$

where the suffixes  $s$ ,  $A$ , and  $w$  refer to the solid, to air, and to water, respectively. In like manner, the condition for the freezing of the ice to the wool, if we neglect the æolotropic properties of the ice, is

$$\sigma_{SW} + \sigma_{IW} > \sigma_{SI},$$

where the suffixes  $s$ ,  $w$ , and  $i$  relate to wool, to water, and to ice, respectively. See *Proc. Roy. Soc.*, vol. x, p. 447; or *Phil. Mag.*, 4th ser., vol. xxi, p. 151.

nature and state of these masses together with the quantities of the foreign substances which may be present in the film. (See page 483.) Problems relating to processes of solidification and dissolution seem hardly capable of a satisfactory solution, except on this supposition, which appears in general allowable with respect to the surfaces produced by these processes. But in considering the equilibrium of fluids at the surface of an unchangeable solid, such a limitation is neither necessary nor convenient. The following method of treating the subject will be found more simple and at the same time more general.

Let us suppose the superficial density of energy to be determined by the excess of energy in the vicinity of the surface over that which would belong to the solid, if (with the same temperature and state of strain) it were bounded by a vacuum in place of the fluid, and to the fluid, if it extended with a uniform volume-density of energy just up to the surface of the solid, or, if in any case this does not sufficiently define a surface, to a surface determined in some definite way by the exterior particles of the solid. Let us use the symbol ( $\varepsilon_s$ ) to denote the superficial energy *thus defined*. Let us suppose a superficial density of entropy to be determined in a manner entirely analogous, and be denoted by ( $\eta_s$ ). In like manner also, for all the components of the fluid, and for all foreign fluid substances which may be present at the surface, let the superficial densities be determined, and denoted by ( $\Gamma_2$ ), ( $\Gamma_3$ ), etc. These *superficial densities of the fluid components* relate solely to the matter which is fluid or movable. All matter which is immovably attached to the solid mass is to be regarded as a part of the same. Moreover, let  $\varsigma$  be defined by the equation

$$\varsigma = (\varepsilon_s) - t(\eta_s) - \mu_2(\Gamma_2) - \mu_3(\Gamma_3) - \text{etc.} \quad (676)$$

These quantities will satisfy the following general relations—

$$d(\varepsilon_s) = t d(\eta_s) + \mu_2 d(\Gamma_2) + \mu_3 d(\Gamma_3) + \text{etc.,} \quad (677)$$

$$d\varsigma = -(\eta_s) dt - (\Gamma_2) d\mu_2 - (\Gamma_3) d\mu_3 - \text{etc.} \quad (678)$$

In strictness, these relations are subject to the same limitation as (674) and (675). But this limitation may generally be neglected. In fact, the values of  $\varsigma$ , ( $\varepsilon_s$ ), etc. must in general be much less affected by variations in the state of strain of the surface of the solid than those of  $\sigma$ ,  $\varepsilon_{s(1)}$ , etc.

The quantity  $\varsigma$  evidently represents the tendency to contraction in that portion of the surface of the fluid which is in contact with the solid. It may be called the *superficial tension of the fluid in contact with the solid*. Its value may be either positive or negative.

It will be observed that for the same solid surface and for the same temperature but for different fluids the values of  $\sigma$  (in all cases to which the definition of this quantity is applicable) will differ from those of  $\varsigma$  by a constant, viz., the value of  $\sigma$  for the solid surface in a vacuum.

For the condition of equilibrium of two different fluids at a line on the surface of the solid, we may easily obtain

$$\sigma_{AB} \cos \alpha = \varsigma_{BS} - \varsigma_{AS}, \quad (679)$$

the suffixes, etc., being used as in (672), and the condition being subject to the same modification when the fluids meet at an edge of the solid.

It must also be regarded as a condition of theoretical equilibrium at the line considered, [subject, like (679), to limitation on account of passive resistances to motion,] that if there are any foreign substances in the surfaces A-S and B-S, the potentials for these substances shall have the same value on both sides of the line; or, if any such substance is found only on one side of the line, that the potential for that substance must not have a less value on the other side; and that the potentials for the components of the mass A, for example, must have the same values in the surface B-C as in the mass A, or, if they are not actual components of the surface B-C, a value not less than in A. Hence, we cannot determine the difference of the surface-tensions of two fluids in contact with the same solid, by bringing them together upon the surface of the solid, unless these conditions are satisfied, as well as those which are necessary to prevent the mixing of the fluid masses.

The investigation on pages 442–448 of the conditions of equilibrium for a fluid system under the influence of gravity may easily be extended to the case in which the system is bounded by or includes solid masses, when these can be treated as rigid and incapable of dissolution. The general condition of mechanical equilibrium would be of the form

$$-\int p \delta Dv + \int g \gamma \delta z Dv + \int \sigma \delta Ds + \int g \Gamma \delta z Ds + \int g \delta z Dm + \int \varsigma \delta Ds + \int g (\Gamma') \delta z Ds = 0, \quad (680)$$

where the first four integrals relate to the fluid masses and the surfaces which divide them, and have the same signification as in equation (606), the fifth integral relates to the movable solid masses, and the sixth and seventh to the surfaces between the solids and fluids, ( $\Gamma'$ ) denoting the sum of the quantities ( $\Gamma_2$ ), ( $\Gamma_3$ ), etc. It should be observed that at the surface where a fluid meets a solid

$\delta z$  and  $\delta z$ , which indicate respectively the displacements of the solid and the fluid, may have different values, but the components of these displacements which are normal to the surface must be equal.

From this equation, among other particular conditions of equilibrium, we may derive the following—

$$d\varsigma = g(\Gamma) dz, \quad (681)$$

[compare (614),] which expresses the law governing the distribution of a thin fluid film on the surface of a solid, when there are no passive resistances to its motion.

By applying equation (680) to the case of a vertical cylindrical tube containing two different fluids, we may easily obtain the well-known theorem that the product of the perimeter of the internal surface by the difference  $\varsigma' - \varsigma''$  of the superficial tensions of the upper and lower fluids in contact with the tube is equal to the excess of weight of the matter in the tube above that which would be there, if the boundary between the fluids were in the horizontal plane at which their pressures would be equal. In this theorem, we may either include or exclude the weight of a film of fluid matter adhering to the tube. The proposition is usually applied to the column of fluid *in mass* between the horizontal plane for which  $p' = p''$  and the actual boundary between the two fluids. The superficial tensions  $\varsigma'$  and  $\varsigma''$  are then to be measured in the vicinity of this column. But we may also include the weight of a film adhering to the internal surface of the tube. For example, in the case of water in equilibrium with its own vapor in a tube, the weight of all the water-substance in the tube above the plane  $p' = p''$ , diminished by that of the water-vapor which would fill the same space, is equal to the perimeter multiplied by the difference in the values of  $\varsigma$  at the top of the tube and at the plane  $p' = p''$ . If the height of the tube is infinite, the value of  $\varsigma$  at the top vanishes, and the weight of the film of water adhering to the tube and of the mass of liquid water above the plane  $p' = p''$  diminished by the weight of vapor which would fill the same space is equal in numerical value but of opposite sign to the product of the perimeter of the internal surface of the tube multiplied by  $\varsigma''$ , the superficial tension of liquid water in contact with the tube at the pressure at which the water and its vapor would be in equilibrium at a plane surface. In this sense, the total weight of water which can be supported by the tube per unit of the perimeter of its surface is directly measured by the value of  $-\varsigma$  for water in contact with the tube.

**MODIFICATION OF THE CONDITIONS OF EQUILIBRIUM BY ELECTRO-MOTIVE FORCE.—THEORY OF A PERFECT ELECTRO-CHEMICAL APPARATUS.**

We know by experience that in certain fluids (electrolytic conductors) there is a connection between the fluxes of the component substances and that of electricity. The quantitative relation between these fluxes may be expressed by an equation of the form

$$De = \frac{Dm_a}{\alpha_a} + \frac{Dm_b}{\alpha_b} + \text{etc.} - \frac{Dm_g}{\alpha_g} - \frac{Dm_h}{\alpha_h} - \text{etc.}, \quad (682)$$

where  $De$ ,  $Dm_a$ , etc. denote the infinitesimal quantities of electricity and of the components of the fluid which pass simultaneously through any same surface, which may be either at rest or in motion, and  $\alpha_a$ ,  $\alpha_b$ , etc.,  $\alpha_g$ ,  $\alpha_h$ , etc. denote positive constants. We may evidently regard  $Dm_a$ ,  $Dm_b$ , etc.,  $Dm_g$ ,  $Dm_h$ , etc., as independent of one another. For, if they were not so, one or more could be expressed in terms of the others, and we could reduce the equation to a shorter form in which all the terms of this kind would be independent.

Since the motion of the fluid as a whole will not involve any electrical current, the densities of the components specified by the suffixes must satisfy the relation

$$\frac{\gamma_a}{\alpha_a} + \frac{\gamma_b}{\alpha_b} + \text{etc.} = \frac{\gamma_g}{\alpha_g} + \frac{\gamma_h}{\alpha_h} + \text{etc.} \quad (683)$$

These densities, therefore, are not independently variable, like the densities of the components which we have employed in other cases.

We may account for the relation (682) by supposing that electricity (positive or negative) is inseparably attached to the different kinds of molecules, so long as they remain in the interior of the fluid, in such a way that the quantities  $\alpha_a$ ,  $\alpha_b$ , etc. of the substances specified are each charged with a unit of positive electricity, and the quantities  $\alpha_g$ ,  $\alpha_h$ , etc. of the substances specified by these suffixes are each charged with a unit of negative electricity. The relation (683) is accounted for by the fact that the constants  $\alpha_a$ ,  $\alpha_g$ , etc. are so small that the electrical charge of any sensible portion of the fluid varying sensibly from the law expressed in (683) would be enormously great, so that the formation of such a mass would be resisted by a very great force.

It will be observed that the choice of the substances which we regard as the components of the fluid is to some extent arbitrary, and that the same physical relations may be expressed by different equa-

tions of the form (682), in which the fluxes are expressed with reference to different sets of components. If the components chosen are such as represent what we believe to be the actual molecular constitution of the fluid, those of which the fluxes appear in the equation of the form (682) are called the *ions*, and the constants of the equation are called their *electro-chemical equivalents*. For our present purpose, which has nothing to do with any theories of molecular constitution, we may choose such a set of components as may be convenient, and call those *ions*, of which the fluxes appear in the equation of the form (682), without farther limitation.

Now, since the fluxes of the independently variable components of an electrolytic fluid do not necessitate any electrical currents, all the conditions of equilibrium which relate to the movements of these components will be the same as if the fluid were incapable of the electrolytic process. Therefore all the conditions of equilibrium which we have found without reference to electrical considerations, will apply to an electrolytic fluid and its independently variable components. But we have still to seek the remaining conditions of equilibrium, which relate to the possibility of electrolytic conduction.

For simplicity, we shall suppose that the fluid is without internal surfaces of discontinuity (and therefore homogeneous except so far as it may be slightly affected by gravity), and that it meets metallic conductors (*electrodes*) in different parts of its surface, being otherwise bounded by non-conductors. The only electrical currents which it is necessary to consider are those which enter the electrolyte at one electrode and leave it at another.

If all the conditions of equilibrium are fulfilled in a given state of the system, except those which relate to changes involving a flux of electricity, and we imagine the state of the system to be varied by the passage from one electrode to another of the quantity of electricity  $\delta e$  accompanied by the quantity  $\delta m_a$  of the component specified, without any flux of the other components or any variation in the total entropy, the total variation of energy in the system will be represented by the expression

$$(V'' - V') \delta e + (\mu_a'' - \mu_a') \delta m_a + (T' - T'') \delta m_a,$$

in which  $V'$ ,  $V''$  denote the electrical potentials in pieces of the same kind of metal connected with the two electrodes,  $T'$ ,  $T''$ , the gravitational potentials at the two electrodes, and  $\mu_a'$ ,  $\mu_a''$ , the intrinsic potentials for the substance specified. The first term represents the increment of the potential energy of electricity, the second the incre-

ment of the intrinsic energy of the ponderable matter, and the third the increment of the energy due to gravitation.\* But by (682)

$$\delta m_a = \alpha_a \delta e$$

It is therefore necessary for equilibrium that

$$V'' - V' + \alpha_a (\mu_a'' - \mu_a' - F'' + F') = 0. \quad (684)$$

To extend this relation to all the electrodes we may write

$$\begin{aligned} V' + \alpha_a (\mu_a' - F') &= V'' + \alpha_a (\mu_a'' - F'') \\ &= V''' + \alpha_a (\mu_a''' - F''') = \text{etc.} \end{aligned} \quad (685)$$

For each of the other cations (specified by <sub>b</sub> etc.) there will be a similar condition, and for each of the anions a condition of the form

$$\begin{aligned} V' - \alpha_g (\mu_g' - F') &= V'' - \alpha_g (\mu_g'' - F'') \\ &= V''' - \alpha_g (\mu_g''' - F''') = \text{etc.} \end{aligned} \quad (686)$$

When the effect of gravity may be neglected, and there are but two electrodes, as in a galvanic or electrolytic cell, we have for any cation

$$V'' - V' = \alpha_a (\mu_a' - \mu_a''), \quad (687)$$

and for any anion

$$V'' - V' = \alpha_g (\mu_g'' - \mu_g'), \quad (688)$$

where  $V'' - V'$  denotes the electromotive force of the combination. That is:—

*When all the conditions of equilibrium are fulfilled in a galvanic or electrolytic cell, the electromotive force is equal to the difference in the values of the potential for any ion or apparent ion at the surfaces of the electrodes multiplied by the electro-chemical equivalent of that ion, the greater potential of an anion being at the same electrode as the greater electrical potential, and the reverse being true of a cation.*

Let us apply this principle to different cases.

(I.) If the ion is an independently variable component of an electrode, or by itself constitutes an electrode, the potential for the ion (in any case of equilibrium which does not depend upon passive resistances to change) will have the same value within the electrode as on its surface, and will be determined by the composition of the electrode with its temperature and pressure. This might be illustrated by a cell with electrodes of mercury containing certain quantities of zinc in solution (or with one such electrode and the other of pure

\* It is here supposed that the gravitational potential may be regarded as constant for each electrode. When this is not the case, the expression may be applied to small parts of the electrodes taken separately.

zinc) and an electrolytic fluid containing a salt of zinc, but not capable of dissolving the mercury.\* We may regard a cell in which hydrogen acts as an ion between electrodes of palladium charged with hydrogen as another illustration of the same principle, but the solidity of the electrodes and the consequent resistance to the diffusion of the hydrogen within them (a process which cannot be assisted by conveotive currents as in a liquid mass) present considerable obstacles to the experimental verification of the relation.

(II.) Sometimes the ion is soluble (as an independently variable component) in the electrolytic fluid. Of course its condition in the fluid when thus dissolved must be entirely different from its condition when acting on an ion, in which case its quantity is not independently variable, as we have already seen. Its diffusion in the fluid in this state of solution is not necessarily connected with any electrical current, and in other relations its properties may be entirely changed. In any discussion of the internal properties of the fluid (with respect to its fundamental equation, for example,) it would be necessary to treat it as a different substance. (See page 117.) But if the process by which the charge of electricity passes into the electrode, and the ion is dissolved in the electrolyte is *reversible*, we may evidently regard the potentials for the substance of the ion in (687) or (688) as relating to the substance thus dissolved in the electrolyte. In case of absolute equilibrium, the density of the substance thus dissolved would of course be uniform throughout the fluid, (since it can move independently of any electrical current,) so that by the strict application of our principle we only obtain the somewhat barren result, that if any of the ions are soluble in the fluid without their electrical charges, the electromotive force must vanish in any case of absolute equilibrium not dependent upon passive resistances. Nevertheless, cases in which the ion is thus dissolved in the electrolytic fluid only to a very small extent, and its passage from one electrode to the other by ordinary diffusion is extremely slow, may be regarded as approximating to the case in which it is incapable of diffusion. In such cases, we may regard the relations (687), (688) as approximately valid, although the condition of equilibrium

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\* If the electrolytic fluid dissolved the mercury as well as the zinc, equilibrium could only subsist when the electromotive force is zero, and the composition of the electrodes identical. For when the electrodes are formed of the two metals in different proportions, that which has the greater potential for zinc will have the less potential for mercury. [See equation (98).] This is inconsistent with equilibrium, according to the principle mentioned above, if both metals can act as cations.

relating to the diffusion of the dissolved ion is not satisfied. This may be the case with hydrogen and oxygen as ions (or apparent ions) between electrodes of platinum in some of its forms.

(III.) The ion may appear in mass at the electrode. If it be a conductor of electricity, it may be regarded as forming an electrode, as soon as the deposit has become thick enough to have the properties of matter in mass. The case therefore will not be different from that first considered. When the ion is a non-conductor, a continuous thick deposit on the electrode would of course prevent the possibility of an electrical current. But the case in which the ion being a non-conductor is disengaged in masses contiguous to the electrode but not entirely covering it, is an important one. It may be illustrated by hydrogen appearing in bubbles at a cathode. In case of perfect equilibrium, independent of passive resistances, the potential of the ion in (687) or (688) may be determined in such a mass. Yet the circumstances are quite unfavorable for the establishment of perfect equilibrium, unless the ion is to some extent absorbed by the electrode or electrolytic fluid, or the electrode is fluid. For if the ion must pass *immediately* into the non-conducting mass, while the electricity passes into the electrode, it is evident that the only possible terminus of an electrolytic current is at the line where the electrode, the non-conducting mass, and the electrolytic fluid meet, so that the electrolytic process is necessarily greatly retarded, and an approximate ceasing of the current cannot be regarded as evidence that a state of approximate equilibrium has been reached. But even a slight degree of solubility of the ion in the electrolytic fluid or in the electrode may greatly diminish the resistance to the electrolytic process, and help toward producing that state of complete equilibrium which is supposed in the theorem we are discussing. And the mobility of the surface of a liquid electrode may act in the same way. When the ion is absorbed by the electrode, or by the electrolytic fluid, the case of course comes under the heads which we have already considered, yet the fact that the ion is set free in mass is important, since it is in such a mass that the determination of the value of the potential will generally be most easily made.

(IV.) When the ion is not absorbed either by the electrode or by the electrolytic fluid, and is not set free in mass, it may still be deposited on the surface of the electrode. Although this can take place only to a limited extent (without forming a body having the properties of matter in mass), yet the electro-chemical equivalents of all substances are so small that a very considerable flux of electricity

may take place before the deposit will have the properties of matter in mass. Even when the ion appears in mass, or is absorbed by the electrode or electrolytic fluid, the non-homogeneous film between the electrolytic fluid and the electrode may contain an additional portion of it. Whether the ion is confined to the surface of the electrode or not, we may regard this as one of the cases in which we have to recognize a certain superficial density of substances at surfaces of discontinuity, the general theory of which we have already considered.

The deposit of the ion will affect the superficial tension of the electrode if it is liquid, or the closely related quantity which we have denoted by the same symbol  $\sigma$  (see pages 482–500) if the electrode is solid. The effect can of course be best observed in the case of a liquid electrode. But whether the electrodes are liquid or solid, if the external electromotive force  $V' - V''$  applied to an electrolytic combination is varied, when it is too weak to produce a lasting current, and the electrodes are thereby brought into a new state of polarization, in which they make equilibrium with the altered value of the electromotive force, without change in the nature of the electrodes or of the electrolytic fluid, then by (508) or (675)

$$d\sigma' = -\Gamma_a' d\mu_a',$$

$$d\sigma'' = -\Gamma_a'' d\mu_a'';$$

and by (687),

$$d(V' - V'') = -\alpha_a(d\mu_a' - d\mu_a'').$$

Hence

$$d(V' - V'') = \frac{\alpha_a}{\Gamma_a'} d\sigma' - \frac{\alpha_a}{\Gamma_a''} d\sigma''. \quad (689)$$

If we suppose that the state of polarization of only one of the electrodes is affected (as will be the case when its surface is very small compared with that of the other), we have

$$d\sigma' = \frac{\Gamma_a'}{\alpha_a} d(V' - V''). \quad (690)$$

The superficial tension of one of the electrodes is then a function of the electromotive force.

This principle has been applied by M. Lippmann to the construction of the electrometer which bears his name.\* In applying equations (689) and (690) to dilute sulphuric acid between electrodes of mercury, as in a Lippmann's electrometer, we may suppose that the

\* See his memoir: "Relations entre les phénomènes électriques et capillaires," *Annales de Chimie et de Physique*, 5e série, t. v, p. 494.

suffix refers to hydrogen. It will be most convenient to suppose the *dividing surface* to be so placed as to make the surface-density of mercury zero. (See page 397.) The matter which exists in excess or deficiency at the surface may then be expressed by the surface-densities of sulphuric acid, of water, and of hydrogen. The value of the last may be determined from equation (690). According to M. Lippmann's determinations, it is negative when the surface is in its natural state (i. e., the state to which it tends when no external electromotive force is applied), since  $\sigma'$  increases with  $V'' - V'$ . When  $V'' - V'$  is equal to nine-tenths of the electromotive force of a Daniell's cell, the electrode to which  $V''$  relates remaining in its natural state, the tension  $\sigma'$  of the surface of the other electrode has a maximum value, and there is no excess or deficiency of hydrogen at that surface. This is the condition toward which a surface tends when it is extended while no flux of electricity takes place. The flux of electricity per unit of new surface formed, which will maintain a surface

in a constant condition while it is extended, is represented by  $\frac{I_a'}{\alpha_a}$  in numerical value, and its direction, when  $I_a''$  is negative, is from the mercury into the acid.

We have so far supposed, in the main, that there are no passive resistances to change, except such as vanish with the rapidity of the processes which they resist. The actual condition of things with respect to passive resistances appears to be nearly as follows. There does not appear to be any passive resistance to the electrolytic process by which an ion is transferred from one electrode to another, except such as vanishes with the rapidity of the process. For, in any case of equilibrium, the smallest variation of the externally applied electromotive force appears to be sufficient to cause a (temporary) electrolytic current. But the case is not the same with respect to the molecular changes by which the ion passes into new combinations or relations, as when it enters into the mass of the electrodes, or separates itself in mass, or is dissolved (no longer with the properties of an ion) in the electrolytic fluid. In virtue of the passive resistance to these processes, the external electromotive force may often vary within wide limits, without creating any current by which the ion is transferred from one of the masses considered to the other. In other words, the value of  $V' - V''$  may often differ greatly from that obtained from (687) or (688) when we determine the values of the potentials for the ion as in cases I, II, and III. We may, however, regard these equations as entirely valid, when the potentials for the

ions are determined at the surface of the electrodes with reference to the ion in the condition in which it is brought there or taken away by an electrolytic current, without any attendant irreversible processes. But in a complete discussion of the properties of the surface of an electrode it may be necessary to distinguish (both in respect to surface-densities and to potentials) between the substance of the ion in this condition and the same substance in other conditions into which it cannot pass (directly) without irreversible processes. No such distinction, however, is necessary when the substance of the ion can pass at the surface of the electrode by reversible processes from any one of the conditions in which it appears to any other.

The formulae (687), (688) afford as many equations as there are ions. These, however, amount to only one independent equation additional to those which relate to the independently variable components of the electrolytic fluid. This appears from the consideration that a flux of any cation may be combined with a flux of any anion in the same direction so as to involve no electrical current, and that this may be regarded as the flux of an independently variable component of the electrolytic fluid.

#### *General Properties of a Perfect Electro-chemical Apparatus.*

When an electrical current passes through a galvanic or electrolytic cell, the state of the cell is altered. If no changes take place in the cell except during the passage of the current, and all changes which accompany the current can be reversed by reversing the current, the cell may be called a perfect electro-chemical apparatus. The electromotive force of the cell may be determined by the equations which have just been given. But some of the general relations to which such an apparatus is subject may be conveniently stated in a form in which the ions are not explicitly mentioned.

In the most general case, we may regard the cell as subject to external action of four different kinds. (1) The supply of electricity at one electrode and the withdrawal of the same quantity at the other. (2) The supply or withdrawal of a certain quantity of heat. (3) The action of gravity. (4) The motion of the surfaces enclosing the apparatus, as when its volume is increased by the liberation of gases.

The increase of the energy in the cell is necessarily equal to that which it receives from external sources. We may express this by the equation

$$d\varepsilon = (V' - V'') de + dQ + dW_e + dW_p, \quad (691)$$

in which  $d\varepsilon$  denotes the increment of the intrinsic energy of the cell,  $de$  the quantity of electricity which passes through it,  $V'$  and  $V''$  the electrical potentials in masses of the same kind of metal connected with the anode and cathode respectively,  $dQ$  the heat received from external bodies,  $dW_g$  the work done by gravity, and  $dW_p$  the work done by the pressures which act on the external surface of the apparatus.

The conditions under which we suppose the processes to take place are such that the increase of the entropy of the apparatus is equal to the entropy which it receives from external sources. The only external source of entropy is the heat which is communicated to the cell by the surrounding bodies. If we write  $d\eta$  for the increment of entropy in the cell, and  $t$  for the temperature, we have

$$d\eta = \frac{dQ}{t}. \quad (692)$$

Eliminating  $dQ$ , we obtain

$$d\varepsilon = (V' - V'') de + t d\eta + dW_g + dW_p, \quad (693)$$

or

$$V'' - V' = -\frac{d\varepsilon}{de} + t \frac{d\eta}{de} + \frac{dW_g}{de} + \frac{dW_p}{de}. \quad (694)$$

It is worth while to notice that if we give up the condition of the reversibility of the processes, so that the cell is no longer supposed to be a perfect electro-chemical apparatus, the relation (691) will still subsist. But, if we still suppose, for simplicity, that all parts of the cell have the same temperature, which is *necessarily* the case with a perfect electro-chemical apparatus, we shall have, instead of (692),

$$d\eta \geq \frac{dQ}{t}, \quad (695)$$

and instead of (693), (694)

$$(V'' - V') de \leq -d\varepsilon + t d\eta + dW_g + dW_p. \quad (696)$$

The values of the several terms of the second member of (694), for a given cell, will vary with the external influences to which the cell is subjected. If the cell is enclosed (with the products of electrolysis) in a rigid envelop, the last term will vanish. The term relating to gravity is generally to be neglected. If no heat is supplied or withdrawn, the term containing  $d\eta$  will vanish. But in the calculation of the electromotive force, which is the most important application of the equation, it is generally more convenient to suppose that the temperature remains constant.

The quantities expressed by the terms containing  $dQ$  and  $d\eta$  in (691), (693), (694), and (696) are frequently neglected in the consideration of cells of which the temperature is supposed to remain constant. In other words, it is frequently assumed that neither heat nor cold is produced by the passage of an electrical current through a perfect electro-chemical combination (except that heat which may be indefinitely diminished by increasing the time in which a given quantity of electricity passes), and that only heat can be produced in any cell, unless it be by processes of a secondary nature, which are not immediately or necessarily connected with the process of electrolysis.

It does not appear that this assumption is justified by any sufficient reason. In fact, it is easy to find a case in which the electromotive force is determined entirely by the term  $t \frac{d\eta}{de}$  in (694), all the other terms in the second member of the equation vanishing. This is true of a Grove's gas battery charged with hydrogen and nitrogen. In this case, the hydrogen passes over to the nitrogen,—a process which does not alter the energy of the cell, when maintained at a constant temperature. The work done by external pressures is evidently nothing, and that done by gravity is (or may be) nothing. Yet an electrical current is produced. The work done (or which may be done) by the current outside of the cell is the equivalent of the work (or of a part of the work) which might be gained by allowing the gases to mix in other ways. This is equal, as has been shown by Lord Rayleigh,\* to the work which may be gained by allowing each gas separately to expand at constant temperature from its initial volume to the volume occupied by the two gases together. The same work is equal, as appears from equations (278), (279) on page 217, (see also page 220,) to the increase of the entropy of the system multiplied by the temperature.

It is possible to vary the construction of the cell in such a way that nitrogen or other neutral gas will not be necessary. Let the cell consist of a U-shaped tube of sufficient height, and have pure hydrogen at each pole under very unequal pressures (as of one and two atmospheres respectively) which are maintained constant by properly weighted pistons, sliding in the arms of the tube. The difference of the pressures in the gas-masses at the two electrodes must of course be balanced by the difference in the height of the two columns of acidulated water. It will hardly be doubted that such an apparatus

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\* *Philosophical Magazine*, vol. xlix, p. 311.

would have an electromotive force acting in the direction of a current which would carry the hydrogen from the denser to the rarer mass. Certainly the gas could not be carried in the opposite direction by an external electromotive force without the expenditure of as much (electromotive) work as is equal to the mechanical work necessary to pump the gas from the one arm of the tube to the other. And if by any modification of the metallic electrodes (which remain unchanged by the passage of electricity) we could reduce the passive resistances to zero, so that the hydrogen could be carried reversibly from one mass to the other without finite variation of the electromotive force, the only possible value of the electromotive force would be represented by the expression  $t \frac{dy}{de}$ , as a very close approximation. It will be observed

that, although gravity plays an essential part in a cell of this kind by maintaining the difference of pressure in the masses of hydrogen, the electromotive force cannot possibly be ascribed to gravity, since the work done by gravity, when hydrogen passes from the denser to the rarer mass, is negative.

Again, it is entirely improbable that the electrical currents caused by differences in the concentration of solutions of salts, (as in a cell containing sulphate of zinc between zinc electrodes, or sulphate of copper between copper electrodes, the solution of the salt being of unequal strength at the two electrodes,) which have recently been investigated theoretically and experimentally by M.M. Helmholtz and Moser,\* are confined to cases in which the mixture of solutions of different degrees of concentration will produce heat. Yet in cases in which the mixture of more and less concentrated solutions is not attended with evolution or absorption of heat, the electromotive force must vanish in a cell of the kind considered, if it is determined simply by the diminution of energy in the cell. And when the mixture produces cold, the same rule would make any electromotive force impossible except in the direction which would tend to increase the difference of concentration. Such conclusions as would be quite irreconcileable with the theory of the phenomena given by Professor Helmholtz.

A more striking example of the necessity of taking account of the variations of entropy in the cell in *a priori* determinations of electromotive force is afforded by electrodes of zinc and mercury in a solution of sulphate of zinc. Since heat is absorbed when zinc is dissolved

\* *Annalen der Physik und Chemie, Neue Folge*, Band iii, February, 1878.

in mercury,\* the energy of the cell is increased by a transfer of zinc to the mercury, when the temperature is maintained constant. Yet in this combination, the electromotive force acts in the direction of the current producing such a transfer.† The couple presents certain anomalies when a considerable quantity of zinc is united with the mercury. The electromotive force changes its direction, so that this case is usually cited as an illustration of the principle that the electromotive force is in the direction of the current which diminishes the energy of the cell; *i. e.*, which produces or allows those changes which are accompanied by evolution of heat when they take place directly. But whatever may be the cause of the electromotive force which has been observed acting in the direction from the amalgam through the electrolyte to the zinc (a force which according to the determinations of M. Gaugain is only one twenty-fifth part of that which acts in the reverse direction when pure mercury takes the place of the amalgam), these anomalies can hardly affect the general conclusions with which alone we are here concerned. If the electrodes of a cell are pure zinc and an amalgam containing zinc not in excess of the amount which the mercury will dissolve at the temperature of the experiment without losing its fluidity, and if the only change (other than thermal) accompanying a current is a transfer of zinc from one electrode to the other,—conditions which may not have been satisfied in all the experiments recorded, but which it is allowable to suppose in a theoretical discussion, and which certainly will not be regarded as inconsistent with the fact that heat is absorbed when zinc is dissolved in mercury,—it is impossible that the electromotive force should be in the direction of a current transferring zinc from the amalgam to the electrode of pure zinc. For, since the zinc eliminated from the amalgam by the electrolytic process might be re-dissolved directly, such a direction of the electromotive force would involve the possibility of obtaining an indefinite amount of electromotive work, and therefore of mechanical work, without other expenditure than that of heat at the constant temperature of the cell.

None of the cases which we have been considering involve combinations by definite proportions, and, except in the case of the cell with electrodes of mercury and zinc, the electromotive forces are quite small. It may perhaps be thought that with respect to those cells in which combinations take place by definite proportions the electromotive force may be calculated with substantial accuracy from

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\* J. Regnault, *Comptes Rendus*, t. li, p. 778.

† Gaugain, *Comptes Rendus*, t. xlvi, p. 430.

the diminution of the energy, without regarding the variation of entropy. But the phenomena of chemical combination do not in general seem to indicate any possibility of obtaining from the combination of substances by any process whatever an amount of mechanical work which is equivalent to the heat produced by the direct union of the substances.

A kilogramme of hydrogen, for example, combining by combustion under the pressure of the atmosphere with eight kilogrammes of oxygen to form liquid water, yields an amount of heat which may be represented in round numbers by 34000 calories.\* We may suppose that the gases are taken at the temperature of 0° C., and that the water is reduced to the same temperature. *But this heat cannot be obtained at any temperature desired.* A very high temperature has the effect of preventing to a greater or less extent, the combination of the elements. Thus, according to M. Sainte-Claire Deville,† the temperature obtained by the combustion of hydrogen and oxygen cannot much if at all exceed 2500° C., which implies that less than one-half of the hydrogen and oxygen present combine at that temperature. This relates to combustion under the pressure of the atmosphere. According to the determinations of Professor Bunsen‡ in regard to combustion in a confined space, only one-third of a mixture of hydrogen and oxygen will form a chemical compound at the temperature of 2850° C. and a pressure of ten atmospheres, and only a little more than one-half when the temperature is reduced by the addition of nitrogen to 2024° C., and the pressure to about three atmospheres *exclusive* of the part due to the nitrogen.

Now 10 calories at 2500° C. are to be regarded as reversibly convertible into one calorie at 4° C. together with the mechanical work representing the energy of 9 calories. If, therefore, all the 34000 calories obtainable from the union of hydrogen and oxygen under atmospheric pressure could be obtained at the temperature of 2500° C., and no higher, we should estimate the electromotive work performed in a perfect electro-chemical apparatus in which these elements are combined or separated at ordinary temperatures and under atmospheric pressure as representing nine-tenths of the 34000 calories, and the heat evolved or absorbed in the apparatus as representing one-tenth of the 34000 calories.§ This, of course, would give an electromotive

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\* See Rühlmann's *Handbuch der mechanischen Wärmetheorie*, Bd. ii, p. 290.

† *Comptes Rendus*, t. lvi, p. 199; and t. lxiv, 67.

‡ *Pogg. Ann.*, Bd. exxxi (1867), p. 161.

§ These numbers are not subject to correction for the pressure of the atmosphere, since the 34000 calories relate to combustion under the same pressure.

force exactly nine-tenths as great as is obtained on the supposition that all the 34000 calories are convertible into electromotive or mechanical work. But, according to all indications, the estimate 2500° C. (for the temperature at which we may regard all the heat of combustion as obtainable) is far too high,\* and we must regard the theoretical value of the electromotive force necessary to electrolyze water as considerably less than nine-tenths of the value obtained on the supposition that it is necessary for the electromotive agent to supply all the energy necessary for the process.

The case is essentially the same with respect to the electrolysis of hydrochloric acid, which is probably a more typical example of the process than the electrolysis of water. The phenomenon of dissociation is equally marked, and occurs at a much lower temperature, more than half of the gas being dissociated at 1400° C.† And the heat which is obtained by the combination of hydrochloric acid gas with water, especially with water which already contains a considerable quantity of the acid, is probably only to be obtained at temperatures comparatively low. This indicates that the theoretical value of the electromotive force necessary to electrolyze this acid (*i. e.*, the electromotive force which would be necessary in a reversible electrochemical apparatus), must be very much less than that which could perform in electromotive work the equivalent of all the heat evolved in the combination of hydrogen, chlorine and water to form the liquid submitted to electrolysis. This presumption, based upon the phenomena exhibited in the direct combination of the substances, is corroborated by the experiments of M. Favre, who has observed an absorption of heat in the cell in which this acid was electrolyzed.‡ The

\* Unless the received ideas concerning the behavior of gases at high temperatures are quite erroneous, it is possible to indicate the general character of a process (involving at most only such difficulties as are neglected in theoretical discussions) by which water may be converted into separate masses of hydrogen and oxygen without other expenditure than that of an amount of heat equal to the difference of energy of the matter in the two states and supplied at a temperature far below 2500° C. The essential parts of the process would be (1) vaporizing the water and heating it to a temperature at which a considerable part will be dissociated, (2) the partial separation of the hydrogen and oxygen by filtration, and (3) the cooling of both gaseous masses until the vapor they contain is condensed. A little calculation will show that in a continuous process all the heat obtained in the operation of cooling the products of filtration could be utilized in heating fresh water.

† Sainte-Claire Deville, *Comptes Rendus*, t. Ixiv, p. 67.

‡ See *Mémoires des Savants Étrangers*, Sér. 2, t. xxv, No. 1, p. 142; or *Comptes Rendus*, t. lxxiii, p. 973. The figures obtained by M. Favre will be given hereafter, in connection with others of the same nature.

electromotive work expended must therefore have been less than the increase of energy in the cell.

In both cases of composition in definite proportions which we have considered, the compound has more entropy than its elements, and the difference is by no means inconsiderable. This appears to be the rule rather than the exception with respect to compounds which have less energy than their elements. Yet it would be rash to assert that it is an invariable rule. And when one substance is substituted for another in a compound, we may expect great diversity in the relations of energy and entropy.

In some cases, there is a striking correspondence between the electromotive force of a cell and the rate of diminution of its energy per unit of electricity transmitted, the temperature remaining constant. A Daniell's cell is a notable example of this correspondence. It may perhaps be regarded as a very significant case, since of all cells in common use, it has the most constant electromotive force, and most nearly approaches the condition of reversibility. If we apply our previous notation [compare (691)] with the substitution of finite for infinitesimal differences to the determinations of M. Favre,\* estimating energy in calories, we have for each equivalent (32.6 kilograms) of zinc dissolved

$$(V'' - V') \Delta e = 24327^{\text{cal}}, \quad \Delta \varepsilon = -25394^{\text{cal}}, \quad \Delta Q = -1067^{\text{cal}}.$$

It will be observed that the electromotive work performed by the cell is about four per cent. less than the diminution of energy in the cell.† The value of  $\Delta Q$ , which, when negative, represents the heat evolved in the cell when the external resistance of the circuit is very great, was determined by direct measurement, and does not appear to have been corrected for the resistance of the cell. This correction would diminish the value of  $-\Delta Q$ , and increase that of  $(V'' - V') \Delta e$ , which was obtained by subtracting  $-\Delta Q$  from  $-\Delta \varepsilon$ .

It appears that under certain conditions neither heat nor cold is produced in a Grove's cell. For M. Favre has found that with different degrees of concentration of the nitric acid sometimes heat and sometimes cold is produced.‡ When neither is produced, of course

\* See *Mém. Savants Étrang.*, loc. cit., p. 90; or *Comptes Rendus*, vol. lxix, p. 35, where the numbers are slightly different.

† A comparison of the experiments of different physicists has in some cases given a much closer correspondence. See Wiedemann's *Galvanismus*, etc., 2<sup>te</sup> Auflage, Bd. ii, §§ 1117, 1118.

‡ *Mém. Savants Étrang.*, loc. cit., p. 93; or *Comptes Rendus*, t. lxix, p. 37, and t. lxxiii, p. 893.

the electromotive force of the cell is exactly equal to its diminution of energy per unit of electricity transmitted. But such a coincidence is far less significant than the fact that an absorption of heat has been observed. With acid containing about seven equivalents of water ( $\text{HNO}_6 + 7\text{HO}$ ), M. Favre has found

$$(V'' - V') \Delta e = 46781^{\text{cal}}, \quad \Delta \varepsilon = -41824^{\text{cal}}, \quad \Delta Q = 4957^{\text{cal}};$$

and with acid containing about one equivalent of water ( $\text{HNO}_6 + \text{HO}$ ),

$$(V'' - V') \Delta e = 49847^{\text{cal}}, \quad \Delta \varepsilon = -52714^{\text{cal}}, \quad \Delta Q = -2867^{\text{cal}}.$$

In the first example, it will be observed that the quantity of heat absorbed in the cell is not small, and that the electromotive force is nearly one-eighth greater than can be accounted for by the diminution of energy in the cell.

This absorption of heat in the cell he has observed in other cases, in which the chemical processes are much more simple.

For electrodes of cadmium and platinum in hydrochloric acid his experiments give\*

$$(V'' - V') \Delta e = 9256^{\text{cal}}, \quad \Delta \varepsilon = -8258^{\text{cal}}, \\ \Delta W_p = -290^{\text{cal}}, \quad \Delta Q = 1288^{\text{cal}}.$$

In this case the electromotive force is nearly one-sixth greater than can be accounted for by the diminution of energy in the cell with the work done against the pressure of the atmosphere.

For electrodes of zinc and platinum in the same acid one series of experiments gives†

$$(V'' - V') \Delta e = 16950^{\text{cal}}, \quad \Delta \varepsilon = -16189^{\text{cal}}, \\ \Delta W_p = -290^{\text{cal}}, \quad \Delta Q = 1051^{\text{cal}};$$

and a later series,‡

$$(V'' - V') \Delta e = 16738^{\text{cal}}, \quad \Delta \varepsilon = -17702^{\text{cal}}, \\ \Delta W_p = -290^{\text{cal}}, \quad \Delta Q = -674^{\text{cal}}.$$

In the electrolysis of hydrochloric acid in a cell with a porous partition, he has found§

\* *Comptes Rendus*, t. lxviii, p. 1305. The total heat obtained in the whole circuit (including the cell) when all the electromotive work is turned into heat, was ascertained by direct experiment. This quantity, 7968 calories, is evidently represented by  $(V'' - V') \Delta e - \Delta Q$ , also by  $-\Delta \varepsilon + \Delta W_p$ . [See (691).] The value of  $(V'' - V') \Delta e$  is obtained by adding  $\Delta Q$ , and that of  $-\Delta \varepsilon$  by adding  $-\Delta W_p$ , which is easily estimated, being determined by the evolution of one kilogramme of hydrogen.

† *Ibid.*

‡ *Mém. Savants Étrang.*, loc. cit., p. 145.

§ *Ibid.*, p. 142.

$$(V' - V'') \Delta e = 34825 \text{ cal.} \quad \Delta Q = 2113 \text{ cal.}$$

whence

$$\Delta \varepsilon - \Delta W_p = 36938.$$

We cannot assign a precise value to  $\Delta W_p$ , since the quantity of chlorine which was evolved in the form of gas is not stated. But the value of  $-\Delta W_p$  must lie between  $290 \text{ cal.}$  and  $580 \text{ cal.}$ , probably nearer to the former.

The great difference in the results of the two series of experiments relating to electrodes of zinc and platinum in hydrochloric acid is most naturally explained by supposing some difference in the conditions of the experiment, as in the concentration of the acid, or in the extent to which the substitution of zinc for hydrogen took place.\* That which it is important for us to observe in all these cases is that there are conditions under which heat is absorbed in a galvanic or electrolytic cell, so that the galvanic cell has a greater electromotive force than can be accounted for by the diminution of its energy, and the operation of electrolysis requires a less electromotive force than would be calculated from the increase of energy in the cell,—especially when the work done against the pressure of the atmosphere is taken into account.

It should be noticed that in all these experiments the quantity represented by  $\Delta Q$  (which is the critical quantity with respect to the point at issue) was determined by direct measurement of the heat absorbed or evolved by the cell when placed alone in a calorimeter. The resistance of the circuit was made so great by a rheostat placed outside of the calorimeter that the resistance of the cell was regarded as insignificant in comparison, and no correction appears to have been made in any case for this resistance. With exception of the error due to this circumstance, which would in all cases diminish the heat absorbed in the cell (or increase the heat evolved), the probable error of  $\Delta Q$  must be very small in comparison with that of  $(V' - V'') \Delta e$ , or with that of  $\Delta \varepsilon$ , which were in general determined by the compar-

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\* It should perhaps be stated that in his extended memoir published in 1877 in the *Mémoires des Savants Étrangers*, in which he has presumably collected those results of his experiments which he regards as most important and most accurate, M. Favre does not mention the absorption of heat in a cell of this kind, or in the similar cell in which cadmium takes the place of zinc. This may be taken to indicate a decided preference for the later experiments which showed an evolution of heat. Whatever the ground of this preference may have been, it can hardly destroy the significance of the absorption of heat, which was a matter of direct observation in repeated experiments. See *Comptes Rendus*, t. lxviii, p. 1305.

ison of different calorimetical measurements, involving very much greater quantities of heat.

In considering the numbers which have been cited, we should remember that when hydrogen is evolved as gas the process is in general very far from reversible. In a perfect electrochemical apparatus, the same changes in the cell would yield a much greater amount of electromotive work, or absorb a much less amount. In either case, the value of  $\Delta Q$  would be much greater than in the imperfect apparatus, the difference being measured perhaps by thousands of calories.\*

It often occurs in a galvanic or electrolytic cell that an ion which is set free at one of the electrodes appears in part as gas, and is in part absorbed by the electrolytic fluid, and in part absorbed by the electrode. In such cases, a slight variation in the circumstances, which would not sensibly affect the electromotive force, would cause all of the ion to be disposed of in one of the three ways mentioned, if the current were sufficiently weak. This would make a considerable

\* Except in the case of the Grove's cell, in which the reactions are quite complicated, the absorption of heat is most marked in the electrolysis of hydrochloric acid. The latter case is interesting, since the experiments confirm the presumption afforded by the behavior of the substances in other circumstances. (See page 514.) In addition to the circumstances mentioned above tending to diminish the observed absorption of heat, the following, which are peculiar to this case, should be noticed.

The electrolysis was performed in a cell with a porous partition, in order to prevent the chlorine and hydrogen dissolved in the liquid from coming in contact with each other. It had appeared in a previous series of experiments (*Mém. Savants Étrang.*, loc. cit., p. 131; or *Comptes Rendus*, t. Ixvi, p. 1231,) that a very considerable amount of heat might be produced by the chemical union of the gases in solution. In a cell without partition, instead of an absorption, an evolution of heat took place, which sometimes exceeded 5000 calories. If, therefore, the partition did not perfectly perform its office, this could only cause a diminution in the value of  $\Delta Q$ .

A large part at least of the chlorine appears to have been absorbed by the electrolytic fluid. It is probable that a slight difference in the circumstances of the experiment—a diminution of pressure, for example,—might have caused the greater part of the chlorine to be evolved as gas, without essentially affecting the electromotive force. The solution of chlorine in water presents some anomalies, and may be attended with complex reactions, but it appears to be always attended with a very considerable evolution of heat. (See Berthelot, *Comptes Rendus*, t. lxxvi, p. 1514.) If we regard the evolution of the chlorine in the form of gas as the normal process, we may suppose that the absorption of heat in the cell was greatly diminished by the retention of the chlorine in solution.

Under certain circumstances, oxygen is evolved in the electrolysis of dilute hydrochloric acid. It does not appear that this took place to any considerable extent in the experiments which we are considering. But so far as it may have occurred, we may regard it as a case of the electrolysis of water. The significance of the fact of the absorption of heat is not thereby affected.

difference in the variation of energy in the cell, and the electromotive force cannot certainly be calculated from the variation of energy alone in all these cases. The correction due to the work performed against the pressure of the atmosphere when the ion is set free as gas will not help us in reconciling these differences. It will appear on consideration that this correction will in general increase the discordance in the values of the electromotive force. Nor does it distinctly appear which of these cases is to be regarded as normal and which are to be rejected as involving secondary processes.\*

If in any case secondary processes are excluded, we should expect it to be when the ion is identical in substance with the electrode upon which it is deposited, or from which it passes into the electrolyte. But even in this case we do not escape the difficulty of the different forms in which the substance may appear. If the temperature of the experiment is at the melting point of a metal which forms the ion and the electrode, a slight variation of temperature will cause the ion to be deposited in the solid or in the liquid state, or, if the current is in the opposite direction, to be taken up from a solid or from a liquid body. Since this will make a considerable difference in the variation of energy, we obtain different values for the electromotive force above and below the melting point of the metal, unless we also take account of the variations of entropy. Experiment does not indicate the existence of any such difference,† and when we take account of variations of entropy, as in equation (694), it is apparent

that there ought not to be any, the terms  $\frac{d\epsilon}{de}$  and  $t \frac{d\eta}{de}$  being both

\* It will be observed that in using the formulae (694) and (696) we do not have to make any distinction between *primary* and *secondary* processes. The only limitation to the generality of these formulae depends upon the *reversibility* of the processes, and this limitation does not apply to (696).

† M. Raoult has experimented with a galvanic element having an electrode of bismuth in contact with phosphoric acid containing phosphate of bismuth in solution. (See *Comptes Rendus*, t. lxviii. p. 643.) Since this metal absorbs in melting 12.64 calories per kilogramme or 885 calories per equivalent ( $70^{\text{kil}}$ ), while a Daniell's cell yields about 24000 calories of electromotive work per equivalent of metal, the solid or liquid state of the bismuth ought to make a difference of electromotive force represented by .037 of a Daniell's cell, if the electromotive force depended simply upon the energy of the cell. But in M. Raoult's experiments no sudden change of electromotive force was manifested at the moment when the bismuth changed its state of aggregation. In fact, a change of temperature in the electrode from about fifteen degrees above to about fifteen degrees below the temperature of fusion only occasioned a variation of electromotive force equal to .002 of a Daniell's cell.

Experiments upon lead and tin gave similar results.

affected by the same difference, viz., the heat of fusion of an electrochemical equivalent of the metal. In fact, if such a difference existed, it would be easy to devise arrangements by which the heat yielded by a metal in passing from the liquid to the solid state could be transformed into electromotive work (and therefore into mechanical work) without other expenditure.

The foregoing examples will be sufficient, it is believed, to show the necessity of regarding other considerations in determining the electromotive force of a galvanic or electrolytic cell than the variation of its energy alone (when its temperature is supposed to remain constant), or corrected only for the work which may be done by external pressures or by gravity. But the relations expressed by (693), (694), and (696) may be put in a briefer form.

If we set, as on page 144,

$$\psi = \varepsilon - t \eta,$$

we have, for any constant temperature,

$$d\psi = d\varepsilon - t d\eta;$$

and for any perfect electrochemical apparatus, the temperature of which is maintained constant,

$$V'' - V' = - \frac{d\psi}{de} + \frac{dW_g}{de} + \frac{dW_p}{de}; \quad (697)$$

and for any cell whatever, when the temperature is maintained uniform and constant,

$$(V'' - V') de \leq - d\psi + dW_g + dW_p. \quad (698)$$

In a cell of any ordinary dimensions, the work done by gravity, as well as the inequalities of pressure in different parts of the cell may be neglected. If the pressure as well as the temperature is maintained uniform and constant, and we set, as on page 147,

$$\zeta = \varepsilon - t \eta + p v,$$

where  $p$  denotes the pressure in the cell, and  $v$  its total volume (including the products of electrolysis), we have

$$d\zeta = d\varepsilon - t d\eta + p dv,$$

and for a perfect electro-chemical apparatus,

$$V'' - V' = - \frac{d\zeta}{de}, \quad (699)$$

or for any cell,

$$(V'' - V') de \leq - d\zeta. \quad (700)$$

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## ERRATA TO VOL. III.

- Page 1, line 13, for Hagerman, read Hagenman.  
Page 5, line 30, for *Vetumnus*, read *Vertumnus*.  
Page 9, last line, for 1873, read 1874.  
Page 11, line 31, for *virticillata*, read *verticillato*.  
Page 13, line 34, for *cappilare*, read *capillare*.  
Page 23, line 4, for *Smythella*, read *Somythella*.  
Page 28, line 19, for **Caridon**, read **Caridion**.  
Page 35, line 3, for **Scapellum**, read **Scalpellum**.  
Page 58, line 14, for branches, read branchlets.  
Page 60, line 12, for Plate X, read IX.  
Page 167, formula (168), for  $m_1$ , read  $\mu_1$ .  
Page 167, formula (169), for  $m_1, \dots m_{n-1}$ , read  $\mu_1, \dots \mu_{n-1}$ .  
Page 239, formula (333), for  $\frac{C}{t}$  read  $\frac{C}{u_2 t}$ .  
Page 295, note ‡, 1st line, for Ueber de Gliedmassen, read Ueber das Skelett der Gliedmassen.  
Page 304, note, for prominal, read proximal.  
Page 356, last line but two, for crystalline solid, read solid of continuous crystalline structure.  
Page 385, line 13, for  $M'$ , read  $M$ .  
Pages 391, 394, 395, 400, in headings, after *Discontinuity*, add *between Fluid Masses*.  
Page 403, line 16, after any other film, add of the same components.  
Page 405, line 29, after this, add ease.  
Page 432, line 15 of foot-note, for  $H$ , read  $H_s$ .

In figure 40, plate XVI, there is a series of ovals around one-half of the real double points. There should be added to the curve, as represented, a like series of ovals around each of the remaining real double points.

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