$$\frac{1}{2\pi/L_X} = \frac{1}{2\pi/L_X} = \frac{(2\pi)^2}{L_X L_Y}$$

$$\#(\epsilon) = \pi \left(\frac{\epsilon}{h}\right)^{4/3} D(k)$$

$$D(\epsilon) = \frac{d\#}{d\epsilon} = \frac{4}{3} \frac{A}{(2\pi)^2} \pi \frac{\epsilon^{1/3}}{\int_{-1}^{4/3} = \frac{A}{3\pi \int_{-1}^{4/3} \epsilon^{1/3}} \epsilon^{1/3} = \frac{A}{3\pi \int_{-1}^{4/3} \epsilon^{1/3}} \epsilon^{1/3}$$

€ = b k

c)
$$E = \int_{a}^{\infty} E \left(\tilde{n} + \frac{1}{a} \right) D(E) dE = \int_{a}^{\infty} E \left(\frac{1}{e^{E/hT}} + \frac{1}{a} \right) D(E) dE$$

$$C_{A} = \frac{\partial E}{\partial T} \Big|_{A} = \int_{a}^{\infty} E \frac{e^{E/hT}}{(e^{E/kT}-1)^{2}} \frac{A}{3\pi b^{4/3}} E^{1/3} dE$$

$$= \frac{Ak}{3\pi b^{4/3}} (kT)^{4/3} \int_{a}^{\infty} \frac{x^{7/3}e^{x}}{(e^{x}-1)^{2}} dx \propto T^{4/3}$$

One could also proceed directly from E

$$E = \frac{A}{317 \, b^{9/3}} \int_{0}^{\infty} \frac{e^{9/3}}{e^{6/kT}} d\epsilon + C = \frac{A(kT)^{7/3}}{317 \, b^{9/3}} \int_{0}^{\infty} \frac{y/s}{e^{\kappa_{-1}}} d\kappa + C$$

$$C_{A} = \frac{\partial E}{\partial T} \Big|_{A} = \frac{7A \, k}{917 \, b^{9/3}} (kT) \int_{0}^{\infty} \frac{x^{9/3}}{e^{\kappa_{-1}}} d\kappa \propto T$$

d) There is no energy gap behavior because there is no energy gap. For any ht There are always oscillators with towakT.

- THERE ARE 10 SINGLE PARTICLE STATES INCLUDING SAW.

 # 3-PARTICLE STATES = # WAYS OF CHOOSING 3 FROM 10

 WHEN ORDER DOES NOT MATTER = 10x9x8/3x241 = 120
 - E=1.5D: 2 ELECTRONS IN Q, I IN BORC > 4 STATES

 E=1.5D: 2 ELECTRONS IN Q, I IN dore > 4 STATES

 -0/AT -1.50/AT

 Z = 4e +4e +----
 - C) E=4.5 D: 3 ELECTRONS IN & AND OR C > 4 STATES

 E=4 D: 2 ELECTRONS IN & AND OR C > 6 WAYS

 AND I ELECTRON IN 6 OR C > 4 WAYS

 S 6 x 4 = 24 STATES

Z = + 24 e +4 e -4.50/AT

- d) <u>k In 4</u> e) <u>k In 120</u>
 - f) C(T) -> O SINCE THE TOTAL ENERGY HAI AN

 T->00

 UPPER BOOND
 - g-6) E=0: ALL 3 BOSONS IN $Q \rightarrow 1$ STATE E=0: 2 BOSONS IN <math>Q, IIN b OR $C \rightarrow 2$ STATES $\frac{Z=1+2e}{+\cdots}$

- g-c) E=4.5D: 3 BOSONS IN d AND $|ore \rightarrow 4|$ STATES $E=4.0\Delta$: 2 BOSONS IN d AND $|ore \rightarrow 3|$ WAYS

 AND I BOSON IN b OR $c \rightarrow 2|$ WAYS $\Rightarrow 3 \times 2 = c$ STATES = -40/RT = -4.50/RT
 - g-d) 561 kln 1 =0
 - 9-f) C(T) +0 SINCE THE TOTAL ENERGY HAT AN
 - g-a) FOR INFORMATION ONLY!

THE TOTAL NUMBER OF STATES = THE # OF WAYS

OF POTTING 3 SPINIESS BOSONS IN 5 SPATIAL

STATES = # WAYS OF POTTING 3 BALLS IN

5 DIFFERENT BOXES WHEN ORDER DOES NOT

MATTER = # WAYS OF ORDERING 3 BALLS AND

4 PARTITIONS WHEN THE ORDER OF THE BALLS

DOES NOT COUNT

$$= \frac{7!}{(7-3)!} \times \frac{1}{3!} = \frac{7 \cdot c \cdot 5}{3 \cdot 2 \cdot 1} = 35$$

IN PARTICULAR THERE ARE 5 WAY! OF PUTTING ALL

3 BOSON! IN THE SAME STATE, 544 = 20 WAY! OF

PUTTING 2 IN I STATE AND I IN AMOTHER, AND

544/2! = 10 WAY! OF PUTTING EACH IN A SEMBATE STATE.

3 a)
$$\frac{e}{(1-A)} = T^{4}$$

$$\frac{e}{(1-A)} = T^{4}$$

$$\frac{e}{(1-A)} = T^{4}$$

$$\Rightarrow e = A = T^{4}$$

$$T_{H} = T^{4}$$

$$T_{C} = T^{4}$$

$$T_{H}^{Y} - (1-\lambda)T_{H}^{Y} - \lambda T^{Y} = -T_{e}^{Y} + (1-\lambda)T_{e}^{Y} + \lambda T^{Y}$$

$$\lambda T_{H}^{Y} + \lambda T_{e}^{Y} = 2\lambda T^{Y}$$

$$T^{Y} = \frac{T_{H}^{Y} + T_{e}^{Y}}{2}$$

b)
$$\sigma \chi T_{H}^{q} - \sigma \chi T^{q} = \chi \sigma (T_{H}^{q} - T^{q}) = \frac{\chi \sigma}{2} (T_{H}^{q} - T_{e}^{q})$$

$$J \text{ IN ABJENCE OF SHEET} = \sigma (T_{H}^{q} - T_{e}^{q}) = J_{o}$$

$$J_{\text{SHEET}} = \frac{\chi}{2} J_{o} = (\frac{1-\Upsilon}{2}) J_{o}$$

THE ENTROPY OF THE INEAL PARAMAGNET

DEPENDS ON H AND T ONLY THROUGH THE

RATIO Y = LEVEL SPACING & H/T

ADIABATIC => CONSTANT S => CONSTANT H/T

$$\frac{H_f}{H_i} = \frac{20}{6,000} = \frac{1}{400} \implies T_f = \frac{1}{400} T_i = \frac{1}{400} K = 2.5 m K$$

NOTE: SINCE THE SPIN SYSTEM IS USUALLY SATURATED AT THE BEGINNING OF SOCH AN EXPERIMENT, ARGUMENTS BAIFD ON THE CURIE LAW ARE NOT SUFFICIENT.

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