12.520 Lecture Notes 15

Special Cases

In general $\tau_{ij} = \lambda e_{ii} \delta_{ij} + 2\mu e_{ij}$

Plane stress: (e.g., $\tau_{zz} = 0$, $\tau_{xz} = 0$)

$$\underline{\tau} = \begin{bmatrix} \tau_{xx} & \tau_{xy} & 0 \\ \tau_{xy} & \tau_{yy} & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

$$p = -\frac{\tau_{xx} + \tau_{yy}}{3}$$

$$\underline{\tau}^{\text{dev}} = \begin{bmatrix}
\frac{2\tau_{xx} - \tau_{yy}}{3} & \tau_{xy} & 0 \\
\tau_{xy} & \frac{2\tau_{yy} - \tau_{xx}}{3} & 0 \\
0 & 0 & \frac{-\tau_{xx} - \tau_{yy}}{3}
\end{bmatrix}$$

$$\tau_{zz}^{\text{dev}} = -\frac{\tau_{xx} + \tau_{yy}}{3}$$

Solving for strains:

$$e_{xx} = \frac{1}{E} (\tau_{xx} - v\tau_{yy})$$

$$e_{yy} = \frac{1}{E} (\tau_{yy} - v\tau_{xx})$$

$$e_{zz} = -\frac{v}{E} (\tau_{xx} + \tau_{yy})$$

$$e_{xy} = \frac{\tau_{xy}}{2u}$$

Plane strain: (e.g., $e_{zz} = e_{xz} = 0 = \tau_{zz}^{\text{dev}} \implies \tau_{xx}^{\text{dev}} = -\tau_{yy}^{\text{dev}}$)

$$\underline{\tau} = \begin{bmatrix} \tau_{xx} & \tau_{xy} & 0 \\ \tau_{xy} & \tau_{yy} & 0 \\ 0 & 0 & \frac{\tau_{xx} + \tau_{yy}}{2} \end{bmatrix}$$

$$p = -\tau_{zz}$$

$$\underline{\tau}^{\text{dev}} = \begin{bmatrix}
\frac{\tau_{xx} - \tau_{yy}}{2} & \tau_{xy} & 0 \\
\tau_{xy} & \frac{\tau_{yy} - \tau_{xx}}{2} & 0 \\
0 & 0 & 0
\end{bmatrix}$$

$$e_{xx} = \frac{1+\nu}{E} \left[\tau_{xx} (1-\nu) - \nu \tau_{yy} \right]$$

$$e_{yy} = \frac{1+\nu}{E} \left[\tau_{yy} (1-\nu) - \nu \tau_{xx} \right]$$

$$e_{xy} = \frac{\tau_{xy}}{2\mu}$$

Note:

Plate stress
$$\rightarrow$$
 Plane strain
$$E \rightarrow \frac{E}{1-v^2}$$

$$v \rightarrow \frac{v}{1-v}$$

$$\frac{2\lambda\mu}{\lambda+2\mu} \leftarrow \lambda$$

Plane stress: (e.g., $\tau_{33} = 0$)

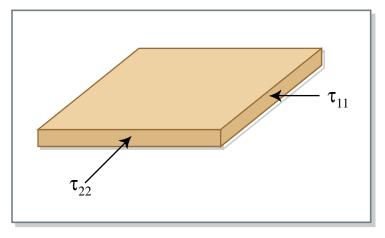


Figure 15.1 Figure by MIT OCW.

$$e_{11} = \frac{1}{E} (\tau_{11} - v\tau_{22})$$

$$e_{22} = \frac{1}{E} (\tau_{22} - v\tau_{11})$$

$$e_{33} = -\frac{v}{E} (\tau_{11} + \tau_{22})$$

Tectonic Stress – Often plane stress is useful $(3D \rightarrow 2D)$.

Assume mostly in plate stress

Assume

$$\tau_{zz} = -\rho gz \quad \text{(lithostatic)}$$

$$\tau_{xx} = -\rho gz + \Delta \tau_{xx}$$

$$\tau_{yy} = -\rho gz + \Delta \tau_{yy}$$

The non-lithostatic (as opposed to deviatoric) stress:

$$\tau_{zz}' = 0$$

$$\tau_{xx}' = \Delta \tau_{xx}$$

$$\tau_{yy}' = \Delta \tau_{yy}$$
plane stress

Assume for simplicity $\Delta \tau_{xx} = \Delta \tau_{yy} = \sigma$

$$e_{11} = e_{22} = \frac{(1-\nu)}{E}\sigma \implies \text{compression for } \sigma \text{ compression}$$
 $e_{33} = -\frac{2\nu}{E}\sigma \implies \text{extension for } \sigma \text{ compression}$

$$\text{dilation } \theta = (2-4\nu)\frac{\sigma}{E} = -\frac{\delta\rho}{\rho}$$

Look at the change in lithostatic pressure at a depth $h \to h + \delta h = h(1 - 2v\frac{\sigma}{E})$

$$\rho gh \rightarrow (\rho + \delta \rho)g(h + \delta h)$$

Plane stress, $\sigma_1 = \sigma_2 = \sigma$

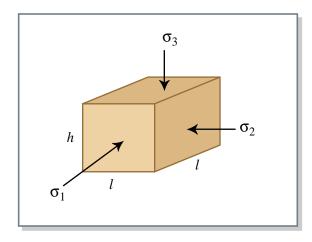


Figure 15.2 Figure by MIT OCW.

$$\varepsilon_{1} = \varepsilon_{2} = \frac{(1 - v)}{E} \sigma$$

$$\varepsilon_{3} = \frac{-2v}{E} \sigma$$

$$\theta = \varepsilon_{1} + \varepsilon_{2} + \varepsilon_{3} = \frac{(2 - 4v)}{E} \sigma = \frac{2(1 - 2v)}{E} \sigma$$

Now to conserve mass, $\frac{\delta \rho}{\rho} = -\frac{\delta V}{V} = \theta$

New lithostat:

$$\begin{split} \delta \tau_{zz} &= -(\rho + \delta \rho)g(h + \delta h) + \rho gh \\ &= (1 + \theta)g(1 + \frac{2v}{E}\sigma)\rho h - \rho gh \\ &= (\theta + \frac{2v}{E}\sigma)\rho gh \end{split}$$

$$\frac{\delta \tau_{zz}}{\sigma} = \frac{2(1-\nu)}{E} \rho g h$$

Use

$$v = 0.25$$
 $E = 100 \text{ GPa} = 1 \text{ Mbar}$
 $h = 100 \text{ km}$
 $\rho = 3 \text{ Mg/m}^3 = 3 \text{ g/cm}^3$

$$\frac{\delta \tau_{zz}}{\sigma} = \frac{2(1-\nu)}{E} \rho g h \; ; \; 5\%$$