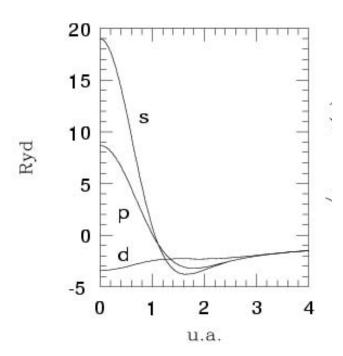
3.320: Lecture 9 (Mar 3 2005) SV((ESS AND FAILURE

Photos of soccer players removed for copyright reasons.

Pseudopotentials



Occupation, Eigenvalue

$$(-\frac{1}{2}\nabla^2 + v_{eff}) \psi_j = \varepsilon_j \psi_j$$

$$\left(-\frac{1}{2}\nabla^2 + \upsilon_{eff}^{(ps)}\right)\psi_j^{(ps)} = \varepsilon_j\psi_j^{(ps)}$$

Figure by MIT OCW.

After Pehlke, Eckhard. Lecture on "The Plane-Wave Pseudopotential Method."

Bloch Theorem

$$\psi_{n\vec{k}}(\vec{r}) \quad u_{n\vec{k}}(\vec{r}) \exp(i\vec{k}\cdot\vec{r})$$

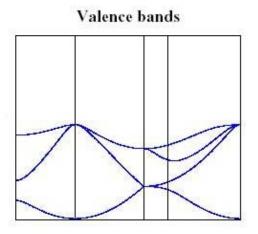
periodic u is expanded in planewaves, labeled according to the reciprocal lattice vectors

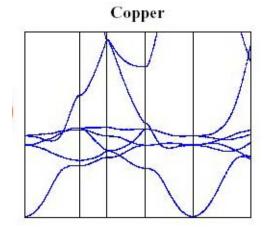
$$u_{n\vec{k}}(\vec{r}) \sum_{\vec{G}} c_{n\vec{k}}^{\vec{G}} \exp(i \vec{G} \cdot \vec{F})$$

Other possibilities - many

- Gaussian basis sets (Hartree-Fock codes)
- Real space representations
- LCAO (Linear combination of atomic orbitals)
- LMTO (Linear muffin-tin orbitals), LAPW (Linearized augmented plane waves), PAW (Projector-augmented wave)

Brillouin Zone integrations





- Sampling at one point (the best Baldereschi point, or the simplest Gamma point)
- Sampling at regular meshes (Monkhorst-Pack grids)
- For metallic systems, integration of the discontinuity is improved introducing a fictitious electronic temperature

Iterations to Selfconsistency

- Construct the external potential (array of non-local psp)
- Choose the plane-wave basis set cutoff, k-point sampling

- Pick a trial electronic density
- Construct the Hamiltonian operator: Hartree and exchange-correlation
- Solve Kohn-Sham equations for the given Hamiltonian (e.g. by diagonalization)
- Calculate the new charge density
- Iterate

Self-consistent ground state

• Iterative diagonalizations (Davidson, Lancsoz, non-scf conjugate gradients + mixing strategy)

$$\hat{H}^{[n(\vec{r})]}\psi(\vec{r}) \quad E\psi(\vec{r})$$

• Direct minimization (Car-Parrinello, conjugate gradients)

Matrix Diagonalization

$$\hat{H}^{[n(\vec{r})]}\psi(\vec{r}) \quad E\psi(\vec{r}) \qquad \left|\psi\right\rangle \quad \sum_{n=1,k} c_n \left|\varphi_n\right\rangle$$

$$\sum_{n=1,k} c_n \langle \varphi_m | \hat{H} | \varphi_{\overline{n}} \rangle \quad Ec_m$$

$$\begin{pmatrix}
H_{11} & \dots & H_{1k} \\
= \cdot & & \cdot \\
= \cdot & = \cdot \\
= \cdot & \vdots \\
= \cdot & \cdot \\
= \cdot & = \cdot \\
= \cdot & \cdot \\
= \cdot & \vdots \\
= \cdot & \cdot \\
C_k
\end{pmatrix} = \begin{pmatrix} c_1 \\ c_1 \\ c_2 \\ \vdots \\ c_k \end{pmatrix}$$

Constrained non-linear minimization

$$\begin{split} E[\{\psi_i\}] &= \sum_{i=1}^N -\frac{1}{2} \int \psi_i^{\star}(\mathbf{r}) \nabla^2 \psi_i(\mathbf{r}) \ d\mathbf{r} + E_H[n(\mathbf{r})] + \\ &+ E_{xc}[n(\mathbf{r})] + \int v_{ext}(\mathbf{r}) n(\mathbf{r}) \ d\mathbf{r} \end{split}$$

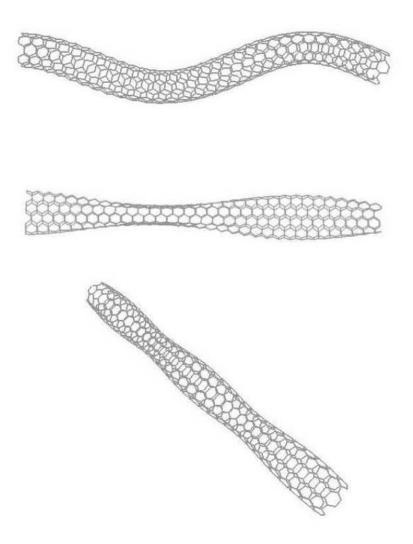
$$E_{H}[n(\vec{r})] \frac{1}{2} \iint \frac{n(\vec{r}_{1})n(\vec{r}_{2})}{|\vec{r}_{1} - \vec{r}_{2}|} d\vec{r}_{1} d\vec{r}_{2}$$

Constrained non-linear minimization

Practice and Outlook

- Techniques and applications (from k-points to DNA and superconductivity)
- Beyond GGA: WDA, TDDFT, QMC
- Connection to approximate methods

Phonons



Phonon Dispersions (Linear Response Theory)

- Exp. neutron scattering data (J.L. Warren et al, Phys. Rev. 158, 805 1967)
- Our GGA results

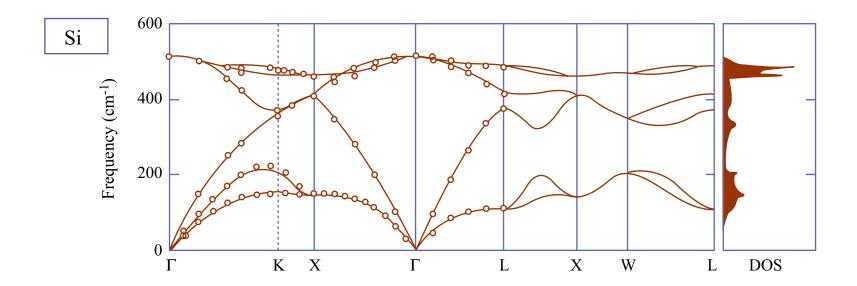
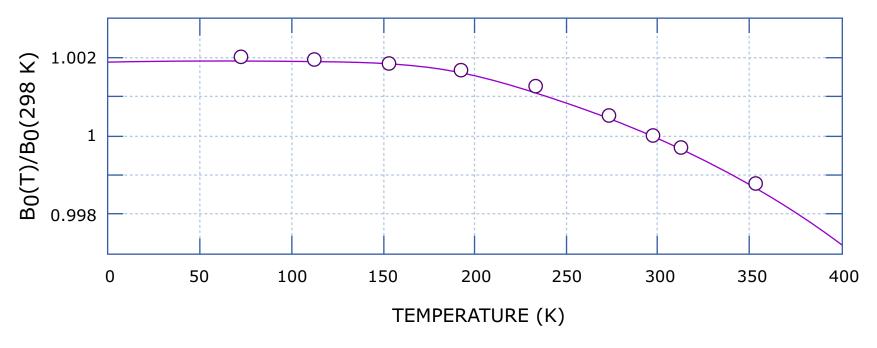


Figure by MIT OCW.

Thermodynamical Properties (Vibrational Free Energy)

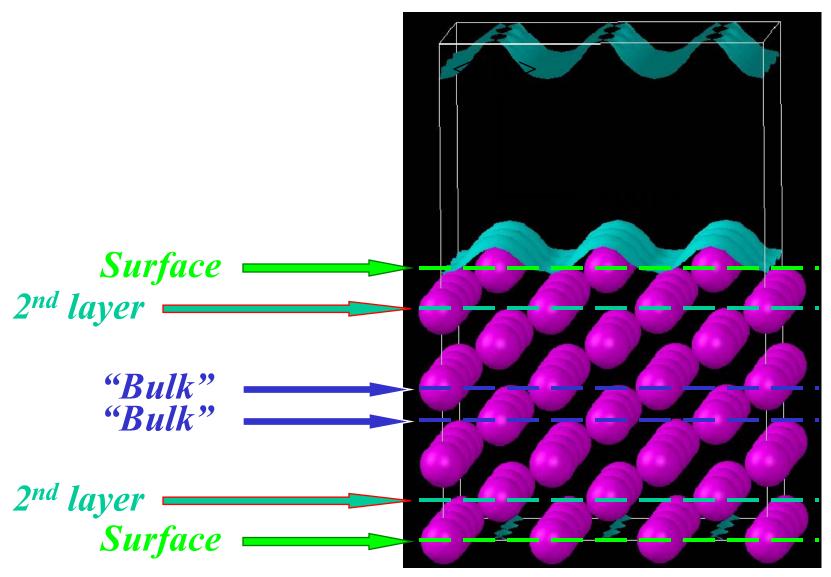


- Our ab-initio results
- Experimental data (McSkimin, H. J., et. al. J Appl Phys 43 (1972): 2944.

Figure by MIT OCW.

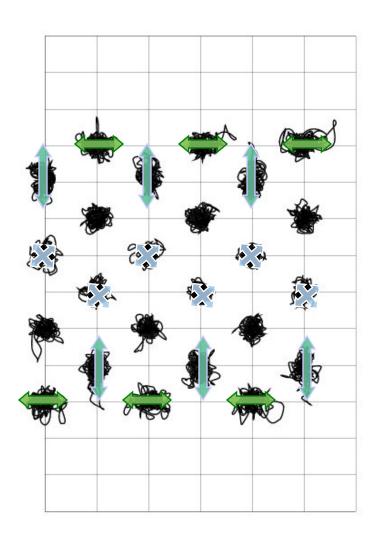
Ab-initio results compared to experimental data of H.J. McSkimin et al., J. Appl. Phys. 43, 2944 -1972.

Dynamics in Al(110)



Mar 3 2005 3.320 Atomistic Modeling of Materials -- Gerbrand Ceder and Nicola Marzari

Al(110) Mean Square Displacements



Computational LEED

Graph removed for copyright reasons.

Interlayer relaxation vs. Temperature (K).

Source: Marzari, N., et al. "Thermal contraction and disordering of the

Al(110) surface." Physical Review Letters 82, no.16 (1999): pp.3296-9.

References (theory)

- W. Koch, M. C. Holthausen, A Chemist's Guide to Density Functional Theory
- R. G. Parr, W. Yang, Density-Functional Theory of Atoms and Molecules
- W. Kohn, Nobel lecture
- F. Jensen, Introduction to Computational Chemistry
- J. M. Thijssen, Computational Physics
- B. H. Bransden and C. J. Joachim, *Physics of Atoms and Molecules*
- K. Burke: *The ABC of DFT*, http://dft.rutgers.edu/kieron/beta/

References (practice)

- Payne, Teter, Allan, Arias, Joannopoulos, *Review of Modern Physics* **64**, 1045 (1992).
- Lecture notes from

http://www.FHI-Berlin.MPG.DE/th/Meetings/FHImd2001/program.html

- (L3 Pehlke, L2 Kratzer, L4 Fuchs)
- Hartree-Fock for solids, Dovesi *et al.*, *Physica Status Solidi (b)* **217**, 63 (2000).

Software

- Gaussian (http://www.gaussian.com) (\$\$) (chemistry, Hartree-Fock, DFT, correlated approaches)
- Gamess-UK (http://www.cse.clrc.ac.uk/qcg/gamess-uk/) (\$) (chemistry, Hartree-Fock, DFT, correlated approaches)
- Materials Studio/Cerius (http://www.accelrys.com) (\$\$) (DFT, planewave, ultrasoft)
- Crystal (http://www.chimifm.unito.it/teorica/crystal) (\$) (Hartree-Fock)
- VASP (http://cms.mpi.univie.ac.at/vasp) (\$) (DFT, planewave, ultrasoft, PAW)
- ESPRESSO (http://www.pwscf.org) (free) (DFT, planewave, ultrasoft, linear-response theory, Car-Parrinello)
- ABINIT (http://www.abinit.org) (free) (DFT, planewave, linear-response theory, GW)
- CPMD (http://www.cpmd.org) (free) (DFT, planewave, Car-Parrinello, time-dependent DFT)
- CASINO (http://www.tcm.phy.cam.ac.uk/~mdt26/cqmc.html) (free) (Quantum Monte Carlo)

3.320/SMA5107: Lecture 9b (3/3/05)

Applications and Performance of DFT methods



Spin Polarization: The Local Spin Density Approximation (LSDA)

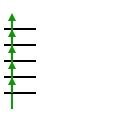
- Electrons have spin +(-) $1/2 \mu_B$
- Spin is treated as a scalar quantity (this is approximate, as relation to angular momentum makes it a vector quantity)
- •Two spin states often referred to as "up" and "down" ↓
- Up-Up interaction is different from Up-Down

Exchange Effects: Refresher

Hund's rule: A consequence of Pauling exclusion principle

Even in solids where energy levels are split, can get

parallel filling if splitting is not too large







Why is this important?

Filling of different orbitals may give the atom different chemical properties

Materials can carry a magnetic moment

$$\rho_{\uparrow} \neq \rho_{\downarrow}$$

Spins in Density Functional Theory

In principle $E_{xc}[\rho]$ "knows" about this effect, but in practice it is poorly approximated since only total charge density is variable

But in practice, need to help LDA along ...

Solution: Treat up and down densities separately

Local Spin Density (LSD) = LDA with different charge density for up and down electrons

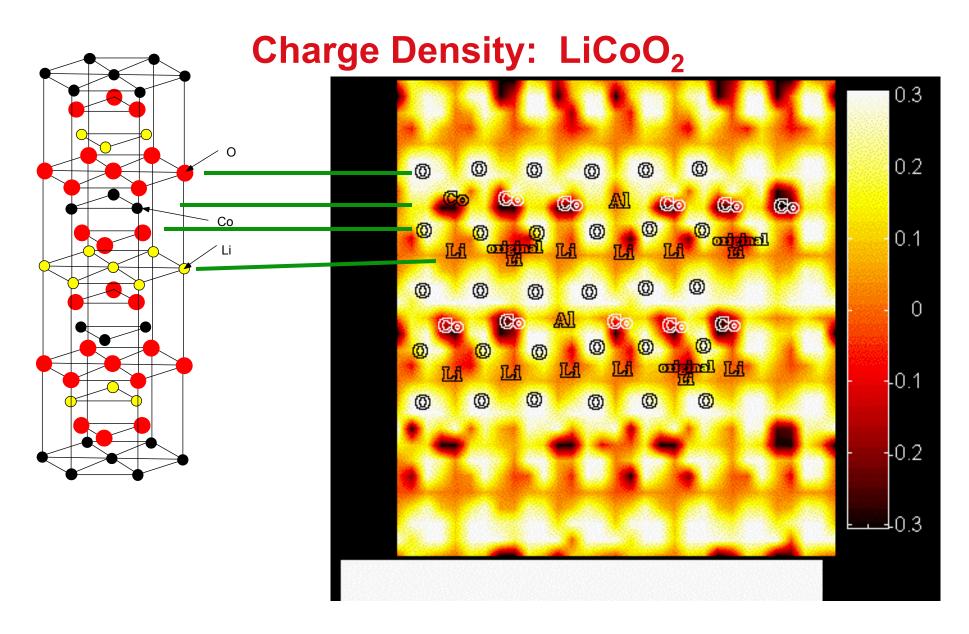
$$v_{eff}^{\uparrow} = v(r) + \int \frac{\rho(r')}{|r - r'|} dr' + \frac{\partial E_{xc} \left[\rho^{\uparrow}, \rho^{\downarrow}\right]}{\partial \rho^{\uparrow}}$$

$$v_{eff}^{\downarrow} = v(r) + \int \frac{\rho(r')}{|r - r'|} dr' + \frac{\partial E_{xc} \left[\rho^{\uparrow}, \rho^{\downarrow}\right]}{\partial \rho^{\downarrow}}$$

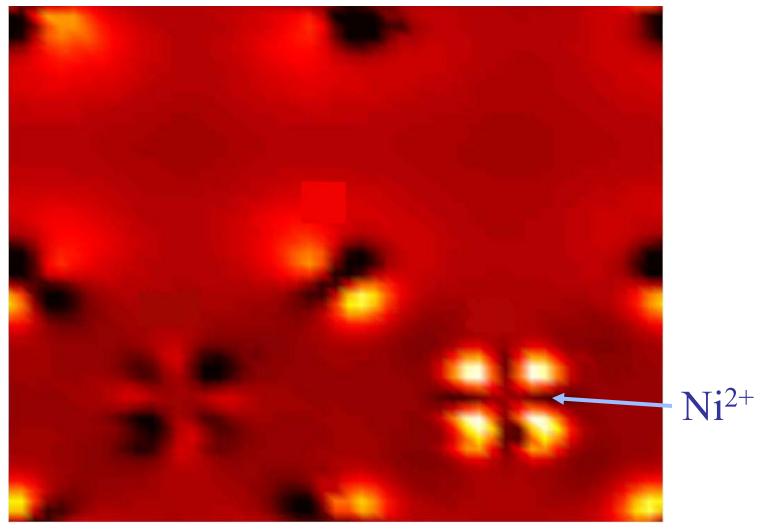
Similar to restricted and unrestricted Hartree Fock

Regular LDA can not capture exchange effect well since it is non-local

A spin-polarized version of GGA exists as well



Spin Density: filters out unpolarized ions (e.g. O^{2-})



How good is the output from LDA/GGA?

Computational quantum mechanics provides very basic information: Energy, charge density, band structure, optimized atomic positions, etc.

How do I make a faster car from this?

Photo of sports car removed for copyright reasons.

First evaluate accuracy of basic information -> then understands how that propagates into higher order models

Energies of Atoms

energy in Ry	Li	Ве	В	С	N	0	F	Ne
GGA: PW91	214.928	229.296	249.24	275.561	2108.926	2149.997	2199.433	2257.893
LDA	214.668	228.892	248.686	274.849	2108.045	2148.939	2198.189	2256.455
Expt.	214.958	229.334	249.308	275.688	2109.174	2150.126	2199.45	2257.856
energy in Ry	Na	Mg	Al	Si	Р	S	CI	Ar
GGA: PW91	2324.514	2400.12	2484.686	2578.669	2682.386	2796.152	2920.278	21055.077
LDA	2322.867	2398.265	2482.618	2576.384	2679.88	2793.419	2917.313	21051.876
Expt.	2324.49	2400.086	2484.672	2578.696	2682.764	2796.2	2920.298	21055.098
energy in Ry	K	Са	Ga	Ge	As	Se	Br	Kr
GGA: PW91	21199.825	21355.144	23850.018	24154.2	24471.917	24803.334	25148.619	25507.943
LDA	21196.383	21351.466	23843.66	24147.583	24465.038	24796.191	25141.209	25500.263

LDA underestimates stability of atom, GGA is closer

Small Molecules

Binding Energy

	exp (ev)	LDA	GGA	HF
H2	-4.753	-4.913	-4.540	-3.64
LiH	-2.509	-2.648	-2.322	
O2	-5.230	-7.595	-6.237	-1.28
H20	-10.078	-11.567	-10.165	
F2	-1.66	-3.32		1.37

Binding energy too high in LDA, GGA is closer but sometimes bound to weak. Pure Hartree Fock without corrections is terrible.

Lattice Parameters in Solids

	exp	LDA	Δ	GGA	Δ
Si	5.427	5.4	-0.50%	5.49	1.16%
Ge	5.65	5.62	-0.53%	5.74	1.59%
GaAs	5.65	5.62	-0.53%	5.73	1.42%
Al	4.03	3.98	-1.31%	4.09	1.57%
Cu	3.60	3.52	-2.35%	3.62	0.44%
Ag	4.07	4.00	-1.69%	4.17	2.47%
Та	3.30	3.26	-1.12%	3.32	0.80%
W	3.16	3.14	-0.67%	3.18	0.67%
Pt	3.91	3.90	-0.41%	3.97	1.49%
Au	4.06	4.05	-0.13%	4.16	2.48%

LDA tends to "overbind", GGA "underbinds" GGA error more variable

Bulk Modulus in Solids (in GPA)

	exp	LDA	Δ	GGA	Δ
Si	99	96	-3.03%	83	-16.16%
Ge	77	78	1.30%	61	-20.78%
GaAs	76	74	-2.63%	65	-14.47%
Al	77	84	9.09%	73	-5.19%
Cu	138	192	39.13%	151	9.42%
Ag	102	139	36.27%	85	-16.67%
Та	193	224	16.06%	197	2.07%
W	310	337	8.71%	307	-0.97%
Pt	283	307	8.48%	246	-13.07%
Au	172	198	15.12%	142	-17.44%

LDA tends to be too stiff. GGA too soft

Oxides

	exp	LDA	Δ	GGA	Δ
MgO	4.21	4.17	-0.95%		
TiO2 (a)	4.59	4.548	-0.92%	4.623	0.72%
TiO2 (c)	2.958	2.944	-0.47%	2.987	0.98%
Al2O2	5.128	5.091	-0.72%	5.185	1.11%
BaTiO3	4	3.94	-1.50%		
PbTiO3	3.9	3.833	-1.72%	3.891	-0.23%
SnO2	4.737	4.637	-2.11%		
β-MnO2 (a)	4.404	4.346	-1.32%	4.444	0.91%
b-MnO2 (c)	2.876	2.81	-2.29%	2.891	0.52%

Summary of Geometry Prediction

LDA under-predicts bond lengths (always?)

GGA error is less systematic though over-prediction is common.

errors are in many cases < 1%, for transition metal oxides < 5%

Predicting Structure: The Energy Scales

Accuracy required to predict physical behavior is astonishing

V Atomic energy: -1894.074 Ry

Fcc V : -1894.7325 Ry

Bcc V : -1894.7125 Ry

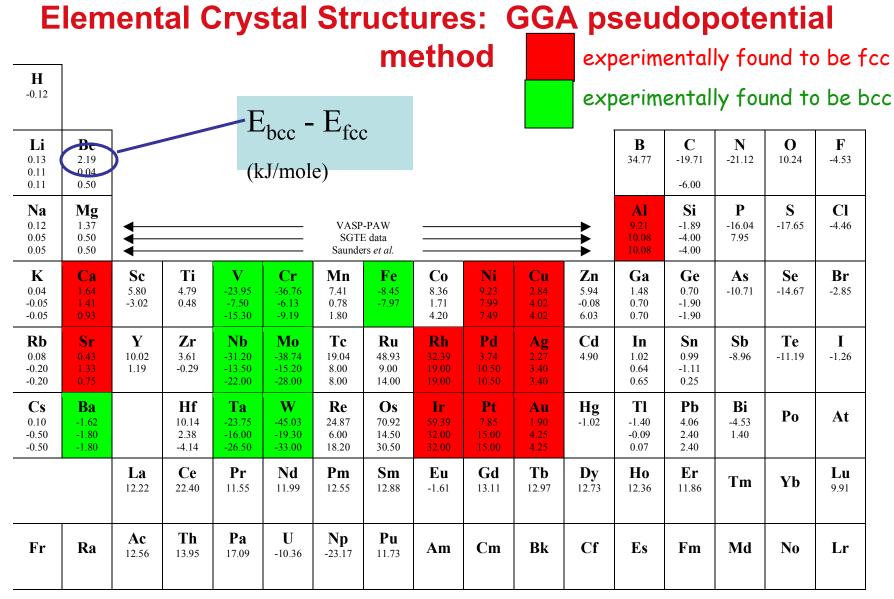
Cohesive energy is 0.638 Ry (0.03% of total E)

Fcc/bcc difference is 0.02 Ry (0.001% of total E)

Mixing energies are also order 10⁻⁶ fraction of total E

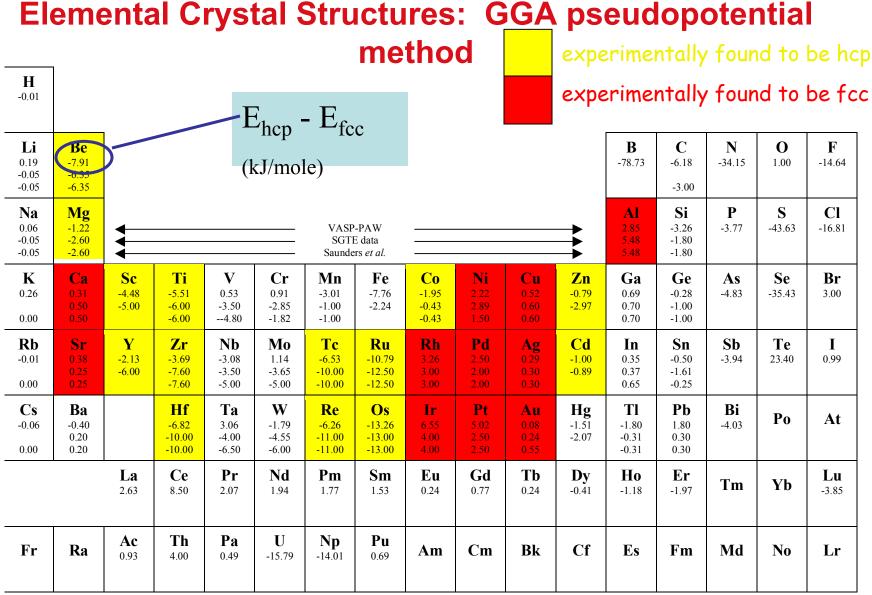
How can we ever get physical behavior correct?

Large cancellation of errors!



data taken from:

Y. Wang,*a S. Curtarolo, et al. *Ab Initio Lattice Stability in Comparison with CALPHAD Lattice Stability*, Computer Coupling of Phase Diagrams and Thermochemistry (Calphad) Vol. 28, Issue 1, March 2004, Pages 79-90.



data taken from:

Y. Wang,*a S. Curtarolo, et al. *Ab Initio Lattice Stability in Comparison with CALPHAD Lattice Stability*, Computer Coupling of Phase Diagram and Thermochemistry (Calphad) Vol. 28, Issue 1, March 2004, Pages 79-90.

Summary: Comparing Energy of Structures

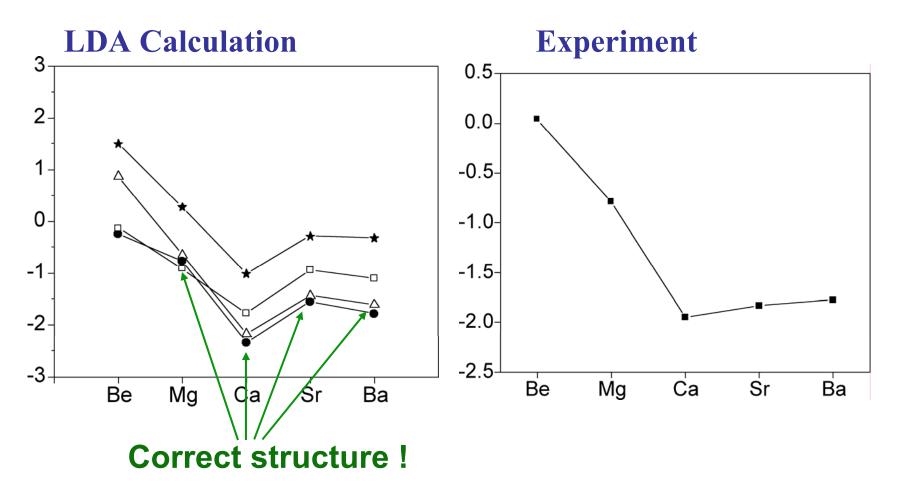
For most elements, both LDA and GGA predict the correct structure for a material (as far as we know)

Notable exceptions: Fe in LDA; materials with substantial electron correlation effects (e.g. Pu)

High Throughput studies are now possible.

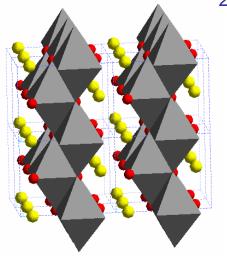
Metal Hydrides

Hydride formation energy: $M + H_2 \rightarrow MH_2$

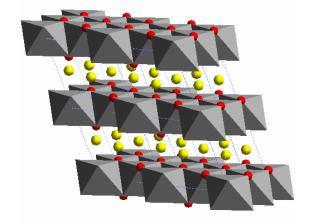


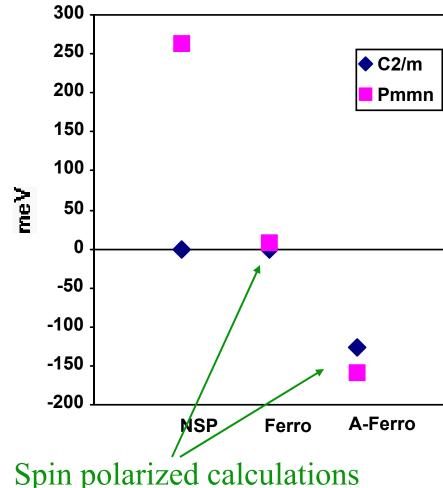
Spin Polarization can be Important (in particular for transition metal compounds

Pmmn: structure of LiMnO₂

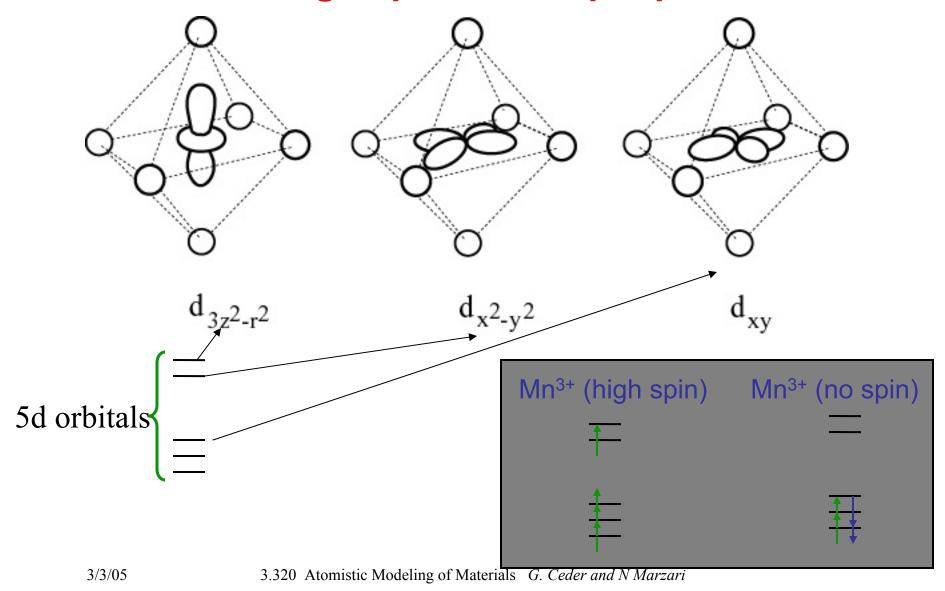


C2/m: other common structure





Orbital filling depends on spin polarization



Reaction Energies

Reaction	Exp (eV)	LDA (eV)	Δ
Li (bcc) + Al (fcc) -> LiAl (compound)	-0.2457	-0.2234	9.08%
0.5 Cu(fcc) + 0.5 Au(fcc) -> CuAu	-0.053	-0.0193	63.58%
Li(bcc) + CoO2 -> LiCoO2	-4.25	-3.75	11.76%

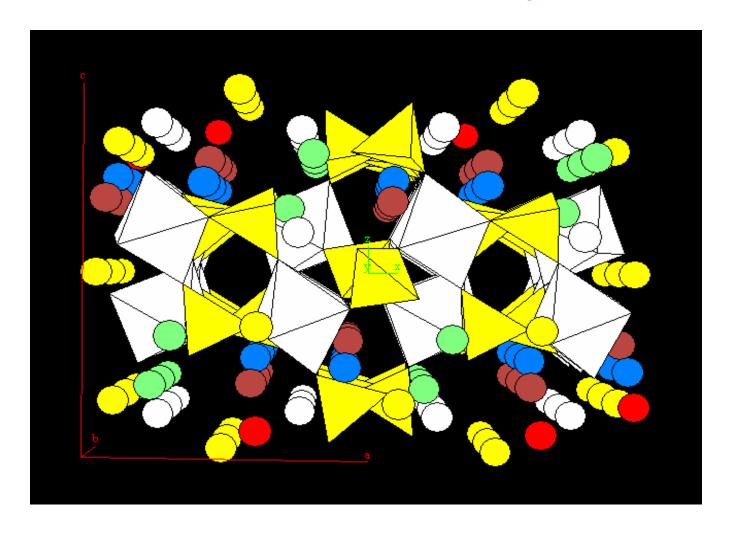
Redox Reactions can be more Problematic

FePO₄ + Li -> LiFePO₄ GGA Exp
2.8 eV 3.5 eV
MnO₂ + Li -> LiMn₂O₄ 3.6 eV 4.1 eV

$$V_2(PO_4)_3$$
 + Li -> LiV₂(PO₄)₃ 3.3 eV 4.6 eV

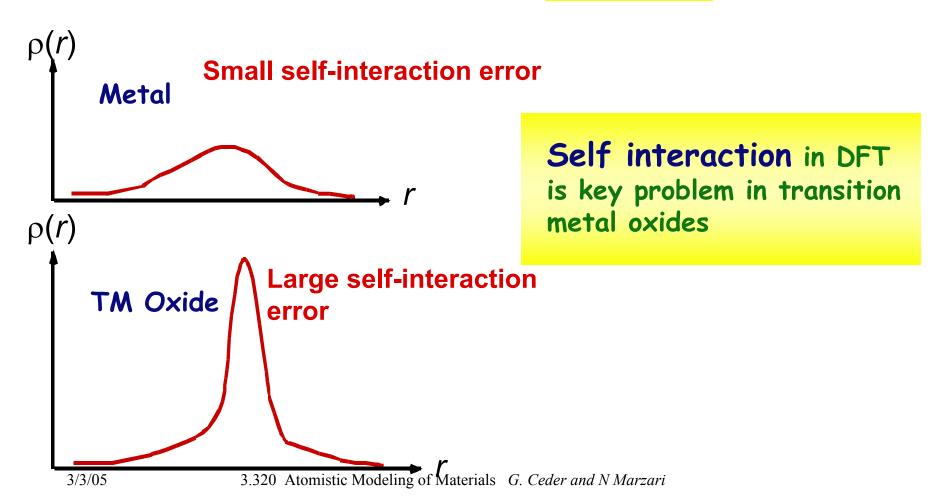
All these reactions involve the transfer of an electron from a delocalized state in Li metal to a localized state in the transition metal oxide (phosphate)

Structure of LiV₂(PO₄)₃

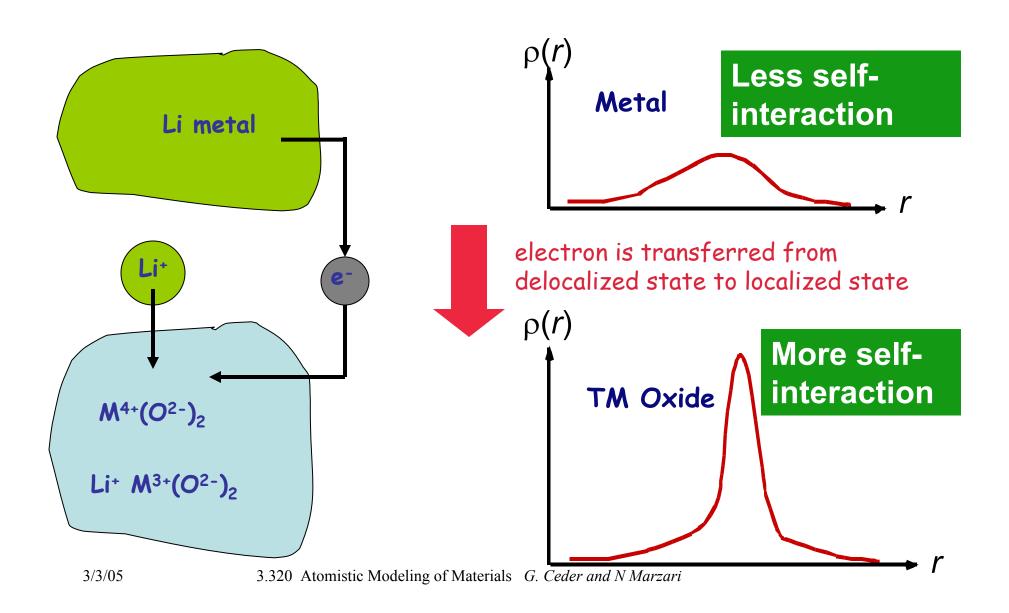


In standard DFT an electron interacts with the effective potential generated by all the electrons (including itself)

$$H = \sum_{i} H_{i} = \sum_{i=1}^{N_{e}} \nabla_{i}^{2} + \sum_{i=1}^{N_{e}} V_{nuclear}(r_{i}) + \sum_{i=1}^{N_{e}} V_{effective}(r_{i})$$



Li insertion into cathode transfers electron from a delocalized to a localized state



Summary (LDA)

Lattice constants: 1-3% too small

Cohesive Energies: 5-20% too strongly bound

Bulk Modulus: 5-20% (largest errors for late TM)

Bandgaps: too small

GGA gives better cohesive energies. Effect on lattice parameters is more random. GGA important for magnetic systems.