3.23 Electrical, Optical, and Magnetic Properties of Materials Fall 2007

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3.23 Fall 2007 – Lecture 9 BAND STRUCTURE

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Last time

- 1. Hamiltonian in a periodic potential, translation operators
- 2. Bloch's theorem (from common eigenstates of H, T_R)
- 3. n, k quantum numbers
- 4. Born-von Karman boundary conditions
- 5. Explicity proof of Bloch's theorem

Study

- Chap. 3 Singleton

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Bloch Theorem (in two equiv forms)

$$\Psi_{n\vec{k}}(\vec{r}) = \exp(i\vec{k} \cdot \vec{r}) u_{n\vec{k}}(\vec{r})$$

$$\Psi_{n\vec{k}}(\vec{r} + \vec{R}) = \exp(i\vec{k} \cdot \vec{R}) \Psi_{n\vec{k}}(\vec{r})$$

$$\psi_{n\vec{k}}(\vec{r} + \vec{R}) = \exp(i\vec{k} \cdot \vec{R}) \Psi_{n\vec{k}}(\vec{r})$$

$$\psi_{n\vec{k}}(\vec{r} + \vec{R}) = \exp(i\vec{k} \cdot \vec{R}) \Psi_{n\vec{k}}(\vec{r})$$

Hamiltonian in the Bloch representation

Explicit solution for the Bloch orbitals
$$\left(\frac{\hbar^{2}(q-G')^{2}}{2m}-E\right)C_{q-G'}+\sum_{G'}V_{G''-G'}C_{q-G''}=0$$

$$\sqrt{(r)}=\sum_{G}V_{G}e^{i\vec{G}\cdot\vec{r}}$$

$$e^{i\vec{G}\cdot\vec{r}}$$

$$\sqrt{(r)}e^{i\vec{G}\cdot\vec{r}}$$

$$\sqrt{(r)}e^{i\vec{G}\cdot\vec{r}}$$

Explicit solution for the Bloch orbitals

$$\left(\frac{\hbar^2 \left(q - G' \right)^2}{2m} - E \right) C_{q-G'} + \sum_{G'} V_{G'-G'} C_{q-G'} = 0$$

$$\left(\frac{\hbar^2 \left(q - G' \right)^2}{2m} (q - 2G)^2 \quad V_{-G} \quad V_{-2G} \quad V_{-3G} \quad V_{-4G} \right)$$

$$V_G \quad \frac{\hbar^2}{2m} (q - G)^2 \quad V_{-G} \quad V_{-2G} \quad V_{-3G} \quad V_{-3G}$$

$$V_{2G} \quad V_G \quad \frac{\hbar^2}{2m} (q)^2 \quad V_{-G} \quad V_{-2G} \quad V_{-2G} \quad C_{q-G} \quad C_{q-G} \quad C_{q-G} \quad C_{q-G} \quad C_{q-G} \quad C_{q-G} \quad C_{q+G} \quad C_{$$

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Free, free, set me free

$$\begin{pmatrix}
\frac{\hbar^{2}}{2m}(q-2G)^{2} & V & V_{2G} & V_{3G} & V_{4G} \\
V_{G} & \frac{\hbar^{2}}{2m}(q-G)^{2} & V & V_{2G} & V_{2G} \\
V_{2G} & V_{G} & \frac{\hbar^{2}}{2m}(q)^{2} & V_{G} & V_{G} \\
V_{3G} & V_{2O} & V_{G} & \frac{\hbar^{2}}{2m}(q+G)^{2} & V_{G} \\
V_{4G} & V_{3G} & V_{2G} & V_{G} & \frac{\hbar^{2}}{2m}(q+G)^{2}
\end{pmatrix}$$

$$V_{4G} = V_{3G} & V_{2G} & V_{G} & \frac{\hbar^{2}}{2m}(q+2G)^{2}$$

$$C_{q+G} \\
C_{q+G} \\
C_{q+G} \\
C_{q+2G} \\
C_$$

Free electron dispersions, 1-d

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Band Structures: Free Electron Gas, Silicon

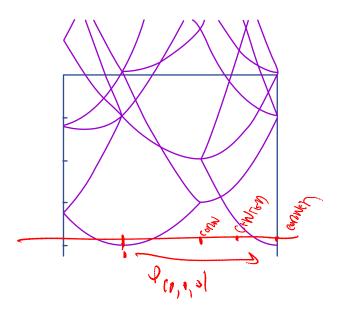


Figure by MIT OpenCourseWare.

Band Structures: Free Electron Gas, Silicon

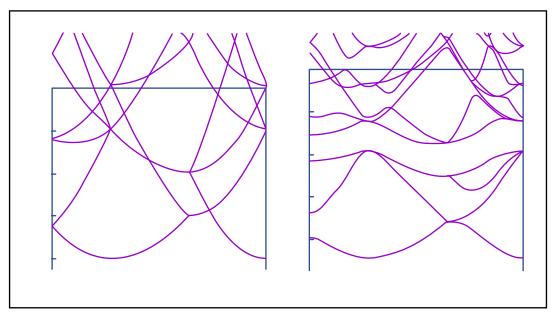
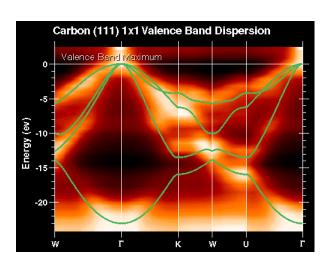


Figure by MIT OpenCourseWare.

Band Structure of Diamond



Courtesy L. J. Terminello and F. J. Himpsel. Used with permission.

$$\begin{pmatrix} \frac{\hbar^2}{2m}(q-2G)^2 & V_{-G} & V_{-G} & V_{-G} & V_{-G} \\ V_{G} & \frac{\hbar^2}{2m}(q-G)^2 & V_{-G} & V_{-G} & V_{-G} \\ V_{G'} & V_{G} & \frac{\hbar^2}{2m}(q)^2 & V_{-G} & V_{-G} \\ V_{G'} & V_{G} & \frac{\hbar^2}{2m}(q+G)^2 & V_{-G} & C_{q+G} \\ V_{+G} & V_{G} & \frac{\hbar^2}{2m}(q+2G)^2 \end{pmatrix} = \begin{pmatrix} C_{q+2G} \\ C_{q+G} \\ C_{q+G} \\ C_{q+2G} \end{pmatrix}$$

$$q = \frac{17}{4}, -\frac{11}{4} = \frac{6}{2}, -\frac{6}{2}$$

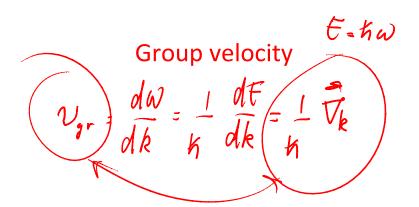
$$\begin{pmatrix} k_{2}^{2} & (\frac{6}{2})^{2} & \sqrt{6} \\ \sqrt{6} & \frac{k_{2}^{2}}{2m} & (\frac{6}{2})^{2} \end{pmatrix} \begin{pmatrix} c_{\frac{6}{2}} \\ c_{\frac{7}{2}} & c_{\frac{7}{2}} \end{pmatrix} = 6 \begin{pmatrix} c_{\frac{6}{2}} \\ c_{-\frac{6}{2}} & c_{\frac{7}{2}} \end{pmatrix}$$

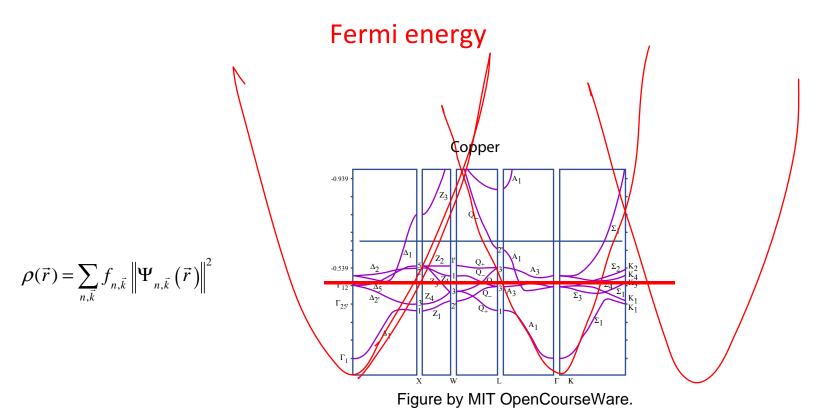
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Band Edge

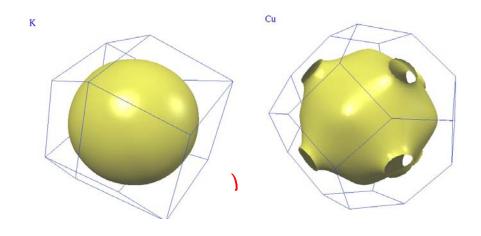
$$\frac{h^2}{2m} \left(\frac{9}{2} \right)^2 \pm V_G$$

$\Psi_{nk}(r)$ is not a momentum eigenstate





The Fermi surface



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Images from the Fermi Surface Database. Used with permission. Please see: http://www.phys.ufl.edu/fermisurface/jpg/K.jpg, http://www.phys.ufl.edu/fermisurface/jpg/Cu.jpg.

D (VRML) Fermi Surface Database

http://www.phys.ufl.e

