3.23 Electrical, Optical, and Magnetic Properties of Materials Fall 2007

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## 3.23 Fall 2007 – Lecture 2 THINK OUTSIDE THE BOX

3.23 Electronic, Optical and Magnetic Properties of Materials - Nicola Marzari (MIT, Fall 2007)

#### More practical info

- Problem sets out on Wed (and posted on Stellar), due by 5pm of the following weekend (after that 75%, after Thu 5pm 50%, after Fri 5pm 25%)
- ~11 in total, 30% of the grade
- Sometimes I mention homework it's not the "Problem Set" @ Poilvert, Bonnet

#### Homework

- Take notes
- Revise posted lecture
- Study posted or assigned material (TEXTBOOKS – do you have them ?)
- Meet with TAs or Instructor:
   Marzari Office Hours Monday 4-5 pm
   Poilvert Office Hours Tuesday 4-5pm

3.012 Fundamentals of Materials Science: Bonding - Nicola Marzari (MIT, Fall 2005)

#### Last time: Wave mechanics

- 1. Particles, fields, and forces
- 2. Dynamics from Newton to Schroedinger
- 3. De Broglie relation  $\lambda \bullet p = h$
- 4. Waves and plane waves
- 5. Harmonic oscillator

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#### Time-dependent Schrödinger's equation

(Newton's 2<sup>nd</sup> law for quantum objects)

$$-\frac{\hbar^2}{2m}\nabla^2\Psi(\vec{r},t)+V(\vec{r},t)\Psi(\vec{r},t)=i\hbar\frac{\partial\Psi(\vec{r},t)}{\partial t}$$

1925-onwards: E. Schrödinger (wave equation), W. Heisenberg (matrix formulation), P.A.M. Dirac (relativistic)

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#### Plane waves as free particles

Our free particle  $\Psi(\vec{r},t) = A \exp[i(\vec{k}\cdot\vec{r}-\omega t)]$  satisfies the wave equation:

$$-\frac{\hbar^2}{2m}\nabla^2\Psi(\vec{r},t) = i\hbar\frac{\partial\Psi(\vec{r},t)}{\partial t} \quad \text{(provided } E = \hbar\omega = \frac{p^2}{2m} = \frac{\hbar^2k^2}{2m}\text{)}$$

#### Stationary Schrödinger's Equation (I)

$$-\frac{\hbar^{2}}{2m}\nabla^{2}\Psi(\vec{r},t)+V(\vec{r},t)\Psi(\vec{r},t)=i\hbar\frac{\partial\Psi(\vec{r},t)}{\partial t}$$

$$=\frac{\hbar^{2}}{2m}\nabla^{2}\Psi(\vec{r},t)+V(\vec{r},t)\Psi(\vec{r},t)=i\hbar\frac{\partial\Psi(\vec{r},t)}{\partial t}$$

$$-\frac{\hbar^{2}}{2m}\nabla^{2}(\varphi f)+V(\vec{r})\varphi f=i\hbar\frac{\partial(\varphi f)}{\partial t}$$

$$-\frac{\hbar^{2}}{2m}f\nabla^{2}\varphi+V\varphi f=i\hbar\frac{\partial f}{\partial t}$$

$$-\frac{42}{2m}\frac{7}{9} + V = i\frac{1}{5}i\frac{1}{5}i\frac{1}{5} = \frac{1}{5}i\frac$$

#### Stationary Schrödinger's Equation (II)

$$\left[ -\frac{\hbar^2}{2m} \nabla^2 + V(\vec{r}) \right] \varphi(\vec{r}) = E \varphi(\vec{r})$$

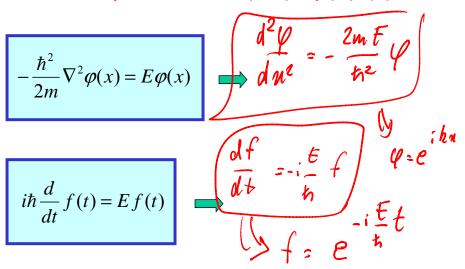
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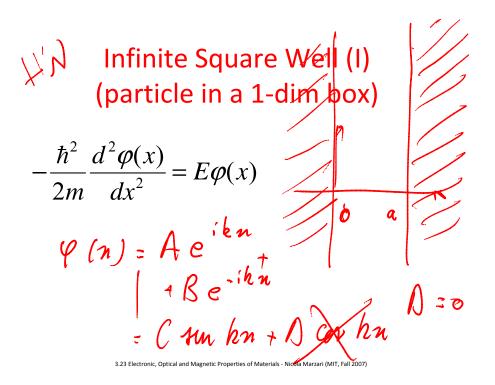
#### Stationary Schrödinger's Equation (III)

$$\left[ -\frac{\hbar^2}{2m} \nabla^2 + V(\vec{r}) \right] \varphi(\vec{r}) = E\varphi(\vec{r})$$

- 1. It's not proven it's postulated, and it is confirmed experimentally
- It's an "eigenvalue" equation: it has a solution only for certain values (discrete, or continuum intervals) of E
- 3. For those eigenvalues, the solution ("eigenstate", or "eigenfunction") is the complete descriptor of the electron in its equilibrium ground state, in a potenitial V(r).
- 4. As with all differential equations, boundary conditions must be specified
- 5. Square modulus of the wavefunction = probability of finding an electron

#### Free particle: $\Psi(x,t)=\varphi(x)f(t)$





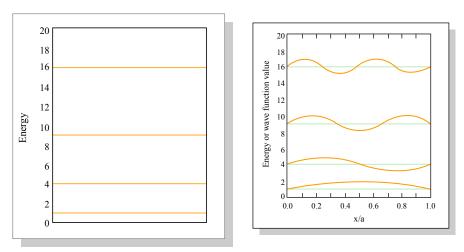
#### Infinite Square Well (II)

Chinka = 0

$$ka = hTT \qquad h = 0, +1, +2, \dots$$

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#### Infinite Square Well (III)



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#### The power of carrots

• β-carotene

Images removed due to copyright restrictions. Please see any spectrum of beta carotene, such as http://www.chm.bris.ac.uk/motm/carotene/beta-carotene\_colourings.html

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#### **Physical Observables from Wavefunctions**

• Eigenvalue equation:

$$\left[ -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V(x) \right] \varphi(x) = E\varphi(x)$$

• Expectation values for the operator (energy)

$$E = \int \varphi^*(x) \left[ -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x) \right] \varphi(x) dx \qquad F = \frac{\hbar^2}{\delta m} \left[ \frac{\hbar^2}{a^2} \right]$$

#### Particle in a 2-dim box

$$-\frac{\hbar^{2}}{2m}\left(\frac{\partial^{2}}{\partial x^{2}} + \frac{\partial^{2}}{\partial y^{2}}\right)\varphi(x, y) = E\varphi(x, y)$$

$$\varphi(x, y) = \chi(x) \chi(y)$$

$$-\frac{\hbar^{2}}{2m}\chi \frac{\partial^{2}\chi}{\partial x} - \frac{\hbar^{2}\chi}{2m}\chi \frac{\partial^{2}\chi}{\partial y^{2}} = \chi(x) \chi(y)$$

$$-\frac{\hbar^{2}}{2m}\chi \frac{\partial^{2}\chi}{\partial x} - \frac{\hbar^{2}\chi}{2m}\chi \frac{\partial^{2}\chi}{\partial y^{2}} = \chi(x) \chi(y)$$

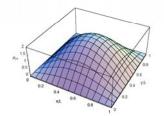
$$-\frac{\hbar^{2}}{2m}\chi \frac{\partial^{2}\chi}{\partial x} = \chi(x) \chi(x)$$

$$-\frac{\hbar^{2}}{2m}\chi \frac{\partial^{2}\chi}{\partial x} = \chi(x)$$

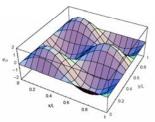
$$-\frac{\hbar^{2}}{$$

#### Particle in a 2-dim box

$$\varphi(x, y) = C \sin\left(\frac{l\pi x}{a}\right) \sin\left(\frac{m\pi y}{b}\right)$$



$$E = \frac{h^2}{8m} \left( \frac{l^2}{a^2} + \frac{m^2}{b^2} \right)$$

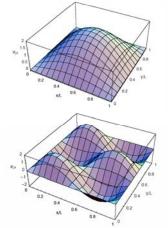


#### Particle in a 3-dim box

$$-\frac{\hbar^2}{2m}\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}\right)\varphi(x, y, z) = E\,\varphi(x, y, z)$$

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# Particle in a 3-dim box: *Farbe* defect in halides (e<sup>-</sup> bound to a negative ion vacancy)



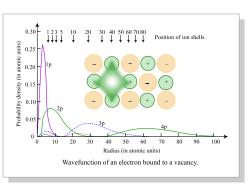


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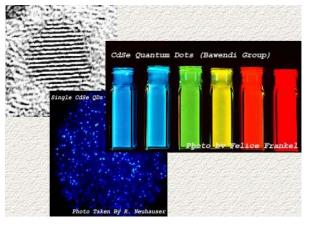
#### From Carl Zeiss to MIT...

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Avakian, P., and Smakula, A. "Color Centers in Cesium Halide Single Crystals."

Physical Review 120 (December 1960): 2007.

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#### Light absorption/emission



Courtesy M. Bawendi and Felice Frankel. Used with permission.

MIT Research: Bawendi, Mayes, Stellacci

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#### Metal Surfaces (I)

$$\left[ -\frac{\hbar^{2}}{2m} \frac{d^{2}}{dx^{2}} + V(x) \right] \varphi(x) = E \varphi(x)$$

$$V(x) \uparrow \qquad \qquad \downarrow 2 \quad \uparrow 2 \quad \downarrow 2 \quad$$

Figure by MIT OpenCourseWare.

#### Metal Surfaces (II)

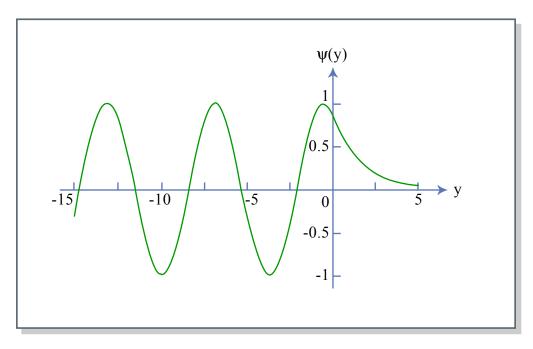


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#### **Scanning Tunnelling Microscopy**

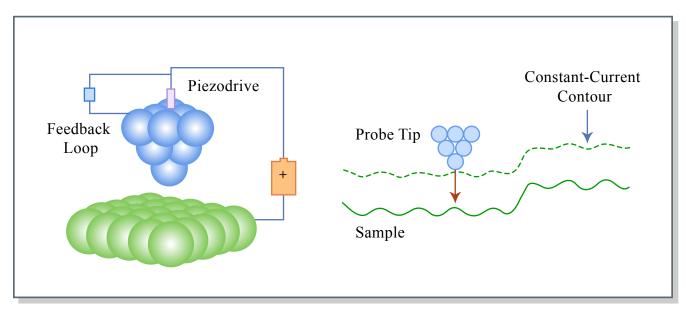


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#### Scanning Tunnelling Microscopy, cont.

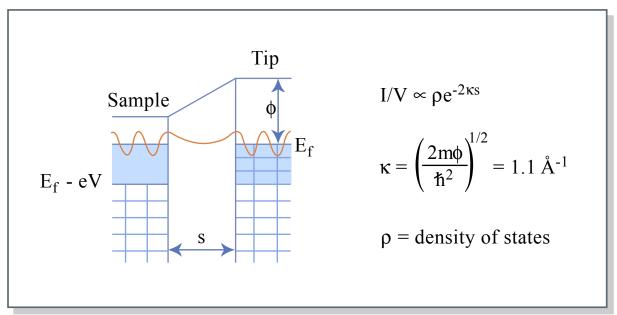
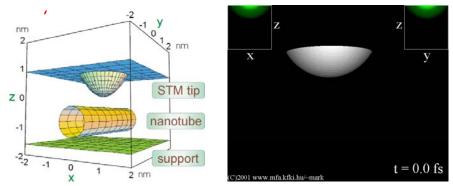


Figure by MIT OpenCourseWare.

### Wavepacket tunnelling through a nanotube



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http://newton.phy.bme.hu/education/schrd/index.html

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http://www.quantum-physics.polytechnique.fr