

# Solving PDEs

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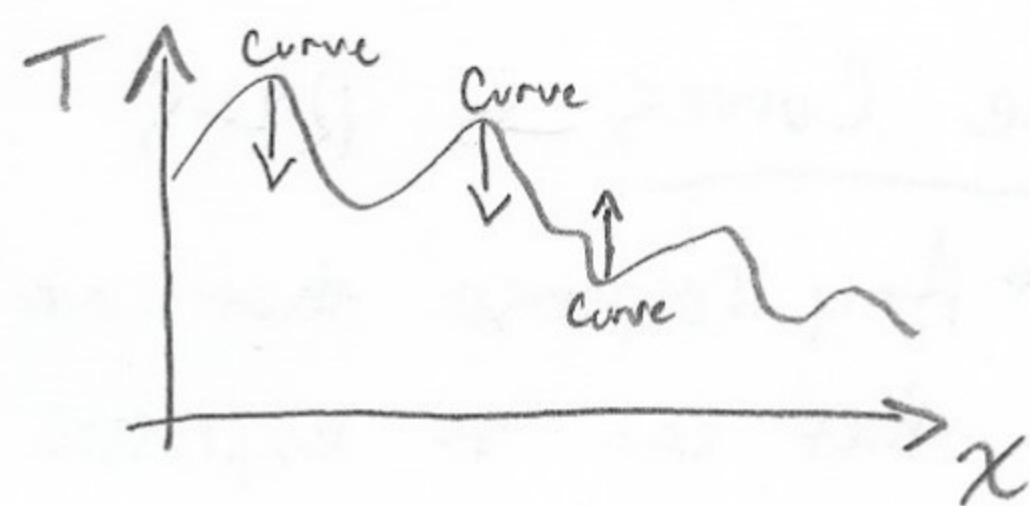
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## - The Heat Equation (Heat)

Heat:  $\frac{\partial T}{\partial t}(x, t) = \alpha \frac{\partial^2 T}{\partial x^2}(x, t)$

↑  
Change in temp.  
as time progresses

↑  
rate of change of the change in temp.  
for every step in space.  $\alpha$  a const.



E.g. How does the temp. distribution change as time progresses?

\* → "Where there's curve in space, there's curve in time."

To solve: Remember: When you're solving a PDE, you're really finding a function that will satisfy several (3) constraints:

1. The PDE,  $\frac{\partial T}{\partial t}(x, t) = \alpha \frac{\partial^2 T}{\partial x^2}(x, t)$

2. Boundary condition (BC):

$$\frac{\partial T}{\partial x}(0, t) = \frac{\partial T}{\partial x}(L, t) = 0 \quad \text{for all } t > 0,$$

this is saying that the BC's must be a constant value,  
as a const. produces a first derivative = 0

Remember: Elements change according to their neighbor, elements on either end of a 1D rod have no "neighbor" to their left/right, so they will tend toward their only neighbor w/in the body. Also, the slope (derivative) at a given point is proportional to the rate of heat flow at that point.

↘ slope = -2.1  
→ Heat flow

To model the restriction that no heat flows at either end, the slope must be zero for  $t > 0$ .  
Physical intuition: We expect  $T \rightarrow 0$  as  $t \rightarrow \infty$



### 3. Initial Condition

- Where are we starting?

#### ① Sine Curves → RHS

↳ Any response that we encounter in practice, is a function (signal) that can be expressed as a sum of sine waves.

↳ Easy to work with.

@  $t=0$   $\frac{\partial T}{\partial x}(x,0) = \cos(x)$  ,  $\frac{\partial^2 T}{\partial x^2}(x,0) = -\sin(x)$

#### Exponentials → LHS

↳ Each point (node) changes its temp at a rate proportional to the point itself; and, at every time step everything scales down by some factor.

\* ↳ When the rate of some value changes is proportional to that value itself, this should immediately tell us that this is an exponential.

\* Keeping these two core concepts in mind (Sines and exponentials), we can assume a solution:

$$T(x,t) = \sin(x)e^{-\alpha t}$$

↳ Sinusoid in space, exponential decay in time.

Check that this solution "fits" the Heat Equation

$$\frac{\partial T}{\partial t} = \alpha \frac{\partial^2 T}{\partial x^2}$$

$$T(x,t) = \sin(x)e^{-\alpha t}$$

$$\frac{\partial}{\partial t}$$

$$\frac{\partial}{\partial x}$$

$$\cos(x)e^{-\alpha t}$$

$$\frac{\partial}{\partial x}$$

$$-\alpha \sin(x)e^{-\alpha t}$$

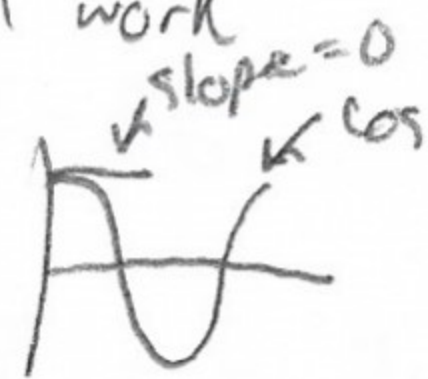
$$= \alpha [-\sin(x)e^{-\alpha t}]$$

→ Both sides are indeed equal and therefore "fit" the Heat Equation.



② BCs Now that we have a function that makes the PDE true, we can move on to the next constraint to satisfy: BCs  
 Question we have to answer: How to apply the BCs that make the slope of either end equal to zero.

↳ Cosine wave starts at a local maximum, so this will work  
 $\cos(x)e^{-\alpha t} \xrightarrow{\frac{\partial}{\partial x}} -\sin(x)e^{-\alpha t} \xrightarrow{\frac{\partial}{\partial x}} -\cos(x)e^{-\alpha t}$   
 for the left side...



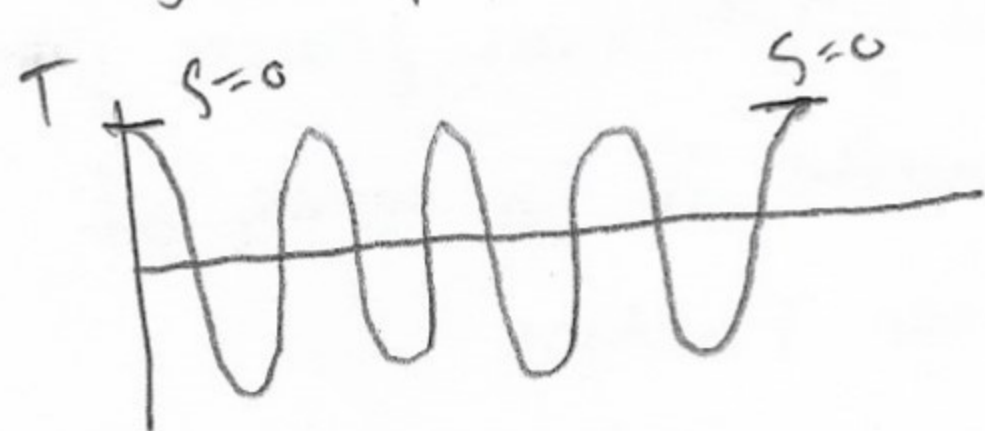
What about the right??

↳ we adjust the frequency,  $\omega$   
 $\cos(\omega \cdot x)$

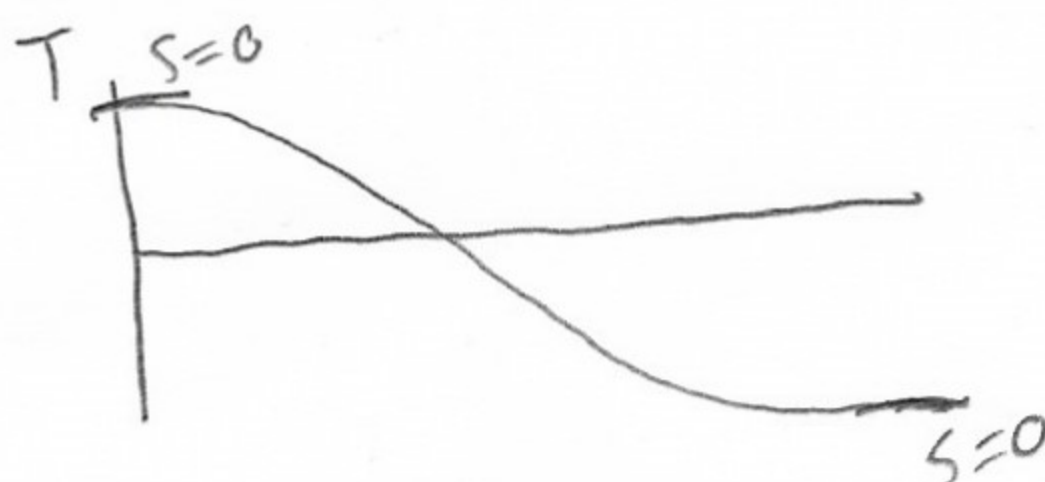
[Recall: The frequency is simply a measure of "how often" a wave moves across a certain distance.]

by adjusting  $\omega$ , we can 'force' the function to end at a maximum, making the slope = 0  
 (or min)

↑ High freq. curves



↓ Low freq. curves



because the frequency affects the "sharpness" of the curvature, so does the derivative of the function.

$$\cos(\omega \cdot x) \xrightarrow{\frac{\partial}{\partial x}} -\omega \sin(\omega \cdot x) \xrightarrow{\frac{\partial}{\partial x}} -\omega^2 \cos(\omega \cdot x)$$

chain rule...

Because we are adding an  $\omega$  to the RHS, to balance the equation we must also add it to the left - LHS (exponential decay part).

↳ Our revised solution becomes:  $T(x, t) = \cos(\omega \cdot x)e^{-\alpha \omega^2 t}$

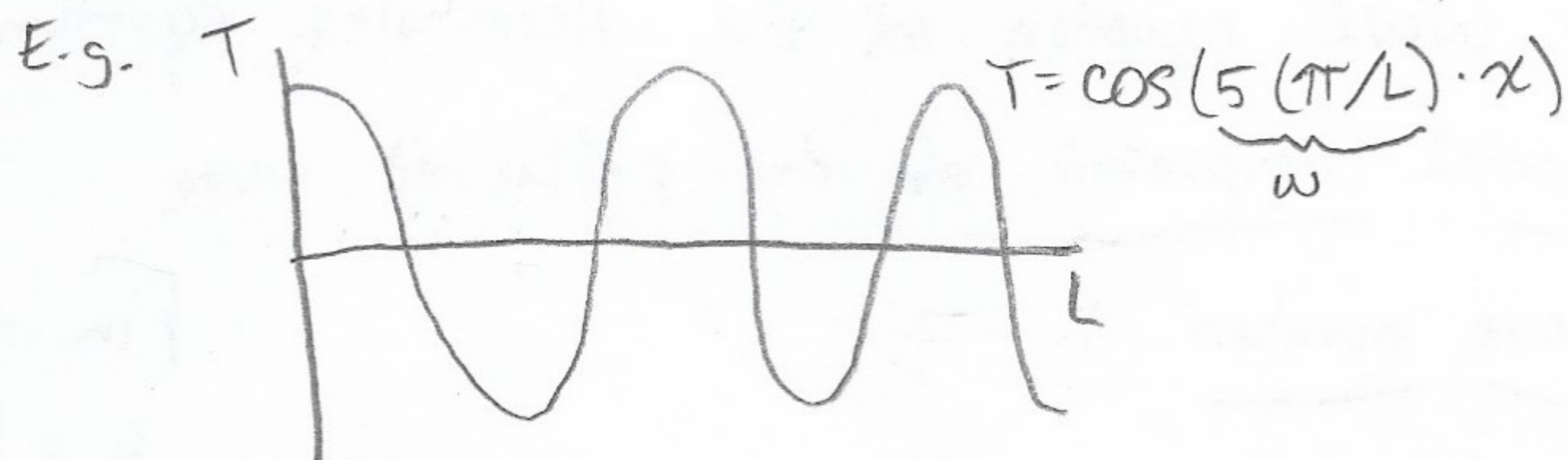
Physical intuition: A temp. function with sharper curves decays much quicker than one w/out as many.  
 → It does so quadratically.



②, cont. The lowest frequency to make both ends  
 \* have a slope of 0, if  $x_{\max} = L$ ,  $\omega = \pi/L$

↳ when  $x = L$ , the input into the cosine expression becomes  $\pi \rightarrow$  this is half the period of a cosine wave

↳ to find other omega values that satisfy, we move through harmonics to find the base # frequencies.



\* ↳ this gives us an infinite family of functions that satisfy both the PDE and the BCs.

I.e. there are multiple functions that fit the Heat Equation PDE and the BCs.

### ③ Linearity

Remember this: Everything stems from this basic observation:

↳ A function that looks like a sine curve in space and an exponential decay in time fits this equation, relating second derivatives in space with first derivatives in time.

$$\rightarrow \frac{\partial T}{\partial t}(x, t) = \alpha \frac{\partial^2 T}{\partial x^2}(x, t)$$



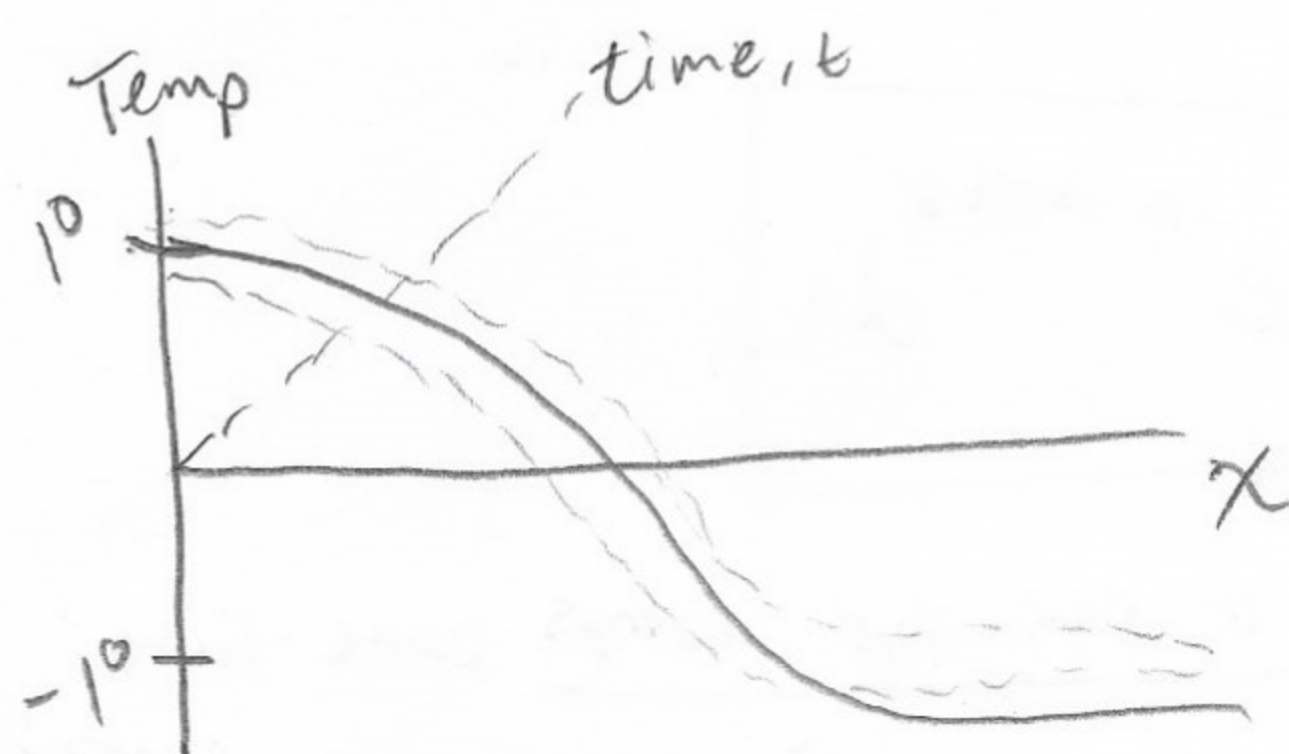
↳ A linear solution simply means, that if we know two solutions, we add them up and the sum becomes a new solution.

↳ Cosine waves - we take the family of solutions, scale them, and add them to create a custom solution - which is a combination of cosine waves.

(boldly)

Fourier asked: How can I represent a function as a sum of sine waves?

Heat equation solutions: - take these, associated with the cosine waves, and add them all up - you will arrive at an exact solution, which describes how the function (temp. distribution + decay) will evolve over time.



$$T(x, t) = \frac{4}{\pi} \left( \frac{\cos(1\pi x)}{1} e^{-\alpha 1^2 t} - \frac{\cos(3\pi x)}{3} e^{-\alpha 3^2 t} + \frac{\cos(5\pi x)}{5} e^{-\alpha 5^2 t} - \dots \right)$$

\* The key challenge is to find the coefficients.

↳ Pairs of vectors rotating in opposite directions.

↳ This is the broader (and what makes computations 'easier') case - rotating vectors composed of a real and imaginary part.



→ The heart + soul of the Fourier Series is the complex exponential,

Frequency  $\omega = 1$ ;  $\underbrace{e^{it}}_{\substack{\uparrow \\ \text{1 cycle/sec rotation.}}} ; \underbrace{e^{2\pi i t}}_{\substack{\uparrow \\ \text{Formula}}}$

\* The knobs that we get to adjust, are both the size and direction of the vectors, by multiplying each one by some constant.  $C_n$ .

$\omega = 1$   
Formula:  $C_1 e^{w 2\pi i t} = \underbrace{0.30 e^{(\pi/4)i}}_{C_1} e^{1 \cdot 2\pi i t}$

\* General formula for  $C_n = \int_0^1 f(t) e^{-n \cdot 2\pi i t} dt$

Recall the step function when two rods of dissimilar temps come into contact.  
(see plot on pg. 5)

↳ To find the coefficients, we need to compute this integral:  
(over and over again until our function is refined)

$$C_n = \int_0^1 \text{step}(t) e^{-n \cdot 2\pi i t} dt = \int_0^{0.5} 1 \cdot e^{-n \cdot 2\pi i t} dt + \int_{0.5}^1 -1 \cdot e^{-n \cdot 2\pi i t} dt$$