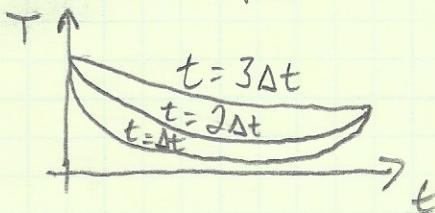


What is a parabolic equation?

See Chapra p. 815

↳ Determines how an unknown varies in both space and time.

↳ It's called parabolic because the function has the form of a parabola, the curvature changes as we step in time.



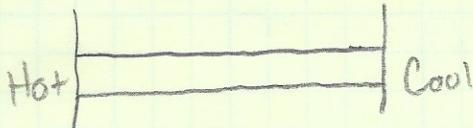
↳ Manifested by the presence of both spatial and temporal derivatives in the heat conduction equation.

↳ The solution will consist of a series of spatial distributions corresponding to the state of the element at various times.

### 30.1 The Heat Conduction Equation

see Chapra p.840

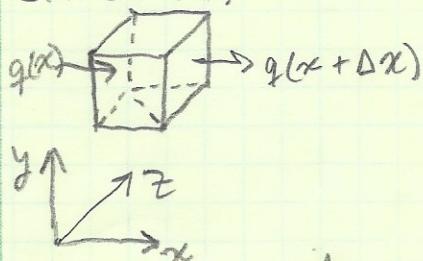
Ex. 1-D case; A thin rod insulated at all points except the ends.



Derivation: Conservation of heat principle is used to develop a heat balance for the differential element:

$$( \text{rate of heat in} ) = ( \text{rate of heat out} ) + ( \text{rate of heat stored} ) \quad (1)$$

Dif. element; or input - outputs = stored



where  $g(x)$  is the heat input in the  $x$ -direction, to obtain the rate we multiply by the area of the face on which it enters and the duration (time)

Here the heat output is a function of the change in direction in which the heat traveled

$$\therefore g(x) \Delta y \Delta z \Delta t - g(x+\Delta x) \Delta y \Delta z \Delta t = \Delta x \Delta y \Delta z \rho C \Delta T, \quad (3)$$

Area duration      Volume

$\rho$  = density of material

$C$  = heat capacity of the material [J/(kg·°C)]

$\Delta T$  = change in temp. across the body in the  $x$ -direction

Dividing (3) by the volume of the element and  $\Delta t$  gives:

$$\frac{g(x) - g(x + \Delta x)}{\Delta x} = \rho C \frac{\Delta T}{\Delta t}$$

the spatial rate  
of change of heat input

↑  
change in temp.  
after a change  
in time.

- OR -

a change in heat input  
after a step in space

(4)

Now, to make this more general, we encode the limit to the ratios after smaller and smaller changes to the input with  $\partial (\delta)$ :

$$-\frac{\partial g}{\partial x} = \rho C \frac{\partial T}{\partial t}$$
(5)

The link between flux and temperature is provided by Fourier's law of heat conduction:

$$q_i = -K f C \frac{\partial T}{\partial i}, \text{ where } \quad \text{see Chapra, p. 821}$$
(29.4)

$q_i$  = heat flux in the direction of the  $i$  dimension  
 $K$  (flow of energy per unit of area per unit of time)

$K$  = coefficient of thermal diffusivity

Taking the derivative of (29.4) to substitute it into (5) gives:

$$\frac{\partial q}{\partial i} = -K f C \frac{\partial^2 T}{\partial i^2}, \text{ then (6) into (5)}$$
(6)

$K \frac{\partial^2 T}{\partial x^2} = \frac{\partial T}{\partial t} \iff \text{Heat conduction equation}$

(30.1)

↳ Contains a second derivative in space and a first derivative in time.

Summary: The rate at which temp. at a point changes w/r respect to time is equal to:

- The second derivative of that point w/r respect to space.

↳ "Where there's curvature in space, there's change in time."

→ The LHS is telling us how the temp. distribution curves as we move across space.

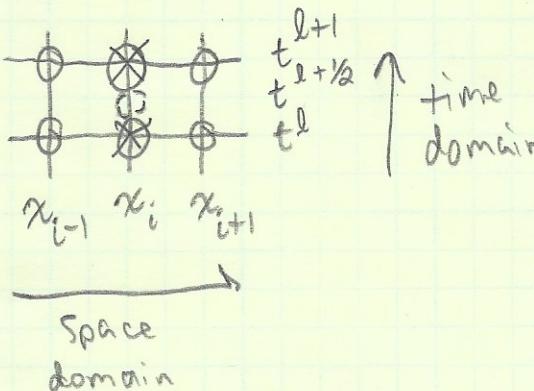
Solving Parabolic PDEs (e.g. heat conduction eq.) using the finite difference Crank-Nicolson method:

See Chaps., p. 849

↳ this method is second-order accurate in both space and time

Consider the grid of computational molecules for the Crank-Nicolson method:

Fig. 30.9



X grid point involved in time difference

O grid point involved in space difference

the temporal first derivative is approx. at  $t^{l+1/2}$  by

$$\frac{\partial T}{\partial t} \approx \frac{T_i^{l+1} - T_i^l}{\Delta t} \quad (30.12)$$

the spatial second derivative is determined at the midpoint by averaging the difference approx. at  $t^l$  and  $t^{l+1}$  of the time increment

$$\frac{\partial^2 T}{\partial x^2} \approx \frac{1}{2} \left[ \frac{T_{i+1}^l - 2T_i^l + T_{i-1}^l}{(\Delta x)^2} + \frac{T_{i+1}^{l+1} - 2T_i^{l+1} + T_{i-1}^{l+1}}{(\Delta x)^2} \right] \quad (30.13)$$

Because we're taking the average of two computational molecules

approx. the 2nd derivative of Temp. w/ respect to the step in space at the beginning of the time increment  $t^l$

same, but at the end of the time increment  $t^{l+1}$

Sub. (30.12) and (30.13) into (30.1) =

$$\frac{K}{2} \left[ \frac{T_{i+1}^l - 2T_i^l + T_{i-1}^l}{(\Delta x)^2} + \frac{T_{i+1}^{l+1} - 2T_i^{l+1} + T_{i-1}^{l+1}}{(\Delta x)^2} \right] = \frac{T_i^{l+1} - T_i^l}{\Delta t}$$

Collect terms:

$$-\lambda T_{i-1}^{l+1} + 2(1+\lambda)T_i^{l+1} - \lambda T_{i+1}^{l+1} = \lambda T_{i-1}^l + 2(1-\lambda)T_i^l + \lambda T_{i+1}^l \quad (30.14)$$

where,  $\lambda = \frac{K \Delta t}{(\Delta x)^2}$ ,  $\lambda$  is known as a (computational) stability criterion

Boundary conditions:

$T_0^{l+1} = f_0(t^{l+1})$ , first ~~interior~~ node, this is a function describing how the boundary temp. changes w/time. why is this plus 1 ??  $\rightarrow$  it doesn't make much sense, why would we use this superscript; why not  $f_0(t^l)$ ? Oh, I think it's related to express that BC temp. changes w/time.

$T_{m+1}^{l+1} = f_{m+1}(t^{l+1})$ , last ~~interior~~ node  $= 50^\circ C$ ?

(30.14.1) and (30.14.2) are sub. into (30.14) to obtain the eqs. for the first and last nodes, respectively.

$\hookrightarrow$  Generating eqs. (30.15) and (30.16), see Chapra, p.850

Evaluate (30.15) numerically for node 1 ( $i=1$ )

$$2.04174 T_1^{l+1} - 0.02087 T_2^{l+1} = 0.02087 (100^\circ C) + 2(1 - 0.02087) T_1^l + 0.02087 T_2^l + 0.02087 (100^\circ C)$$

$\hookrightarrow$  The problem statement says the BCs are fixed for all times at

$$\begin{aligned} T(0) &= 100^\circ C \\ T(L) &= 50^\circ C \end{aligned}$$

$$\therefore T_0^l = T_0^{l+1} = T_0^{l+2} = T_0^{l+n} = 100^\circ C$$

$\hookrightarrow$  and, at  $t=0$ , the temp of the nod = 0

$\hookrightarrow$  at time = 0 ( $l=0$ )

$$\Rightarrow 2.04174 T_1^l - 0.02087 T_2^l = 4.174$$

This is the first equation in the system of equations

$t = \text{time}$  Evaluate (30.15) numerically for node 1 ( $i=1$ ) 6/23/21 16:00 4/6

$i = \text{space}$

~~W~~

$$-\cancel{\lambda T_{i-1}^{l+1}} + 2(1+\lambda)T_i^{l+1} - \cancel{\lambda T_{i+1}^{l+1}} = \cancel{\lambda T_{i-1}^l} + 2(1-\lambda)T_i^l + \cancel{\lambda T_{i+1}^l} \quad (30.14)$$

Boundary Condition ( $i=0$ )

$T_0^{l+1} = f_0(t^{l+1}) \rightarrow$  this is a function describing how the boundary temp. changes w/time. (30.14.1)

Derive the equation for the first interior node ( $i=1$ ):

$$-\cancel{\lambda T_0^{l+1}} + 2(1+\lambda)T_1^{l+1} - \cancel{\lambda T_2^{l+1}} = \cancel{\lambda T_0^l} + 2(1-\lambda)T_1^l + \cancel{\lambda T_2^l}$$


Recall: The problem statement says the BCs are fixed for all times at  $T(0) = 100^\circ\text{C}$   $T(L) = 50^\circ\text{C}$

$$\therefore T_0^l = T_0^{l+1} = T_0^{l+2} = T_0^{l+n} = 100^\circ\text{C}$$

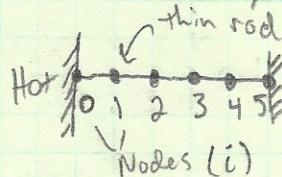
And, at  $t=0$ , the temp of the rod = 0 (Initial Condition)

$$\Rightarrow 2(1+\lambda)T_1^{l+1} - \cancel{\lambda T_2^{l+1}} = \cancel{\lambda T_0^l} + 2(1-\lambda)T_1^l + \cancel{\lambda T_2^l} + \cancel{\lambda T_0^l} + \cancel{\lambda T_0^{l+1}}$$

$$= \lambda(T_0^l + T_0^{l+1})$$

@  $t=n$  ( $l=1$ )

$$2(1+\lambda)T_1^{l+1} - \cancel{\lambda T_2^{l+1}} = \lambda f_0(t^l) + 2(1-\lambda)T_1^l + \cancel{\lambda T_2^l} + \lambda f_0(t^{l+1})$$



Let's setup the linear system of eqs.

for 6 nodes, 5 elements (incl. BCs) at  $t=t_1$ , ( $l=1$ ),  
since  $L=10$ ,  $\Delta x=2 \Rightarrow N_{el}=10/2=5$ ; % elements

Note: Because the BCs are fixed for all time, we only "care" about  
the temp. distribution w/in the rod, i.e. at nodes 1:4

Using Eq. (30.14), evaluate  $T_i^1$  @  $l=t_1=0.15$ , for  $i=1:4$

~~$i=1: -\lambda T_0^2 + 2(1+\lambda)T_1^2 - \lambda T_2^2 = \lambda T_0^1 + 2(1-\lambda)T_1^1 + \lambda T_2^1$~~

Bringing  $T_0$  and  $T_2$  to LHS,  $T_0$  to RHS, collect terms...

~~$2(1+\lambda)T_1^2 - \lambda T_2^2 - 2(1-\lambda)T_1^1 - \lambda T_2^1 = \lambda T_0^1 + \lambda T_0^2$~~ 

Left node, BC

~~$i=2: -\lambda T_1^2 + 2(1+\lambda)T_2^2 - \lambda T_3^2 = \lambda T_1^1 + 2(1-\lambda)T_2^1 + \lambda T_3^1$~~

~~$2(1+\lambda)T_2^2 - \lambda T_3^2 - 2(1-\lambda)T_2^1 - \lambda T_3^1 = \lambda T_1^1 + \lambda T_1^2$~~ 

Left node

~~$i=3: -\lambda T_2^2 + 2(1+\lambda)T_3^2 - \lambda T_4^2 = \lambda T_2^1 + 2(1-\lambda)T_3^1 + \lambda T_4^1$~~

~~$2(1+\lambda)T_3^2 - \lambda T_4^2 - 2(1-\lambda)T_3^1 - \lambda T_4^1 = \lambda T_2^1 + \lambda T_2^2$~~

~~$i=4: -\lambda T_3^2 + 2(1+\lambda)T_4^2 - \lambda T_5^2 = \lambda T_3^1 + 2(1-\lambda)T_4^1 + \lambda T_5^1$~~

~~$-\lambda T_3^2 + 2(1+\lambda)T_4^2 - \lambda T_3^1 - 2(1-\lambda)T_4^1 = \lambda T_5^1 + \lambda T_5^2$~~ 

Right node, BC

@  $t=1$  (time =  $t'$ ), system of eqs.:

$$\begin{array}{l}
 i=1: 2(1+\lambda)T_1^2 - \cancel{\lambda T_2^2} - 2(1-\lambda)T_1^1 - \cancel{\lambda T_2^1} = \cancel{\lambda T_2^1} + \cancel{\lambda T_0^2} \\
 i=2: 2(1+\lambda)T_2^2 - \cancel{\lambda T_3^2} - 2(1-\lambda)T_2^1 - \cancel{\lambda T_3^1} = \cancel{\lambda T_3^1} + \cancel{\lambda T_2^2} \\
 i=3: 2(1+\lambda)T_3^2 - \cancel{\lambda T_4^2} - 2(1-\lambda)T_3^1 - \cancel{\lambda T_4^1} = \cancel{\lambda T_4^1} + \cancel{\lambda T_3^2} \\
 i=4: -\cancel{\lambda T_3^2} + 2(1+\lambda)T_4^2 - \cancel{\lambda T_3^1} - 2(1-\lambda)T_4^1 = \cancel{\lambda T_5^1} + \cancel{\lambda T_5^2}
 \end{array}
 \quad \left. \begin{array}{c} BC \\ \text{Unknowns} \end{array} \right| \begin{array}{c} 64 \\ 2 \\ 2 \\ 0 \end{array}$$

↳ 6 unknowns, 4 eqs. ??

↳ the problem is that I'm starting my analysis at  $t'$ , when I should be starting at  $t^0$ ; starting at  $t^0$  would make the internal nodes = 0, since this is the initial condition.

→ Okay, so let's do that.

$$\begin{array}{l}
 @ t=0 \text{ (time = } t^0\text{)}, \text{ system of eqs. :} \\
 \begin{array}{l}
 i=1: 2(1+\lambda)T_1^1 - \cancel{\lambda T_2^1} - 2(1-\lambda)T_1^0 - \cancel{\lambda T_2^0} = \cancel{\lambda T_0^1} + \cancel{\lambda T_0^0} \\
 i=2: 2(1+\lambda)T_2^1 - \cancel{\lambda T_3^1} - 2(1-\lambda)T_2^0 - \cancel{\lambda T_3^0} = \cancel{\lambda T_1^1} + \cancel{\lambda T_1^0} \\
 i=3: -\cancel{\lambda T_2^1} + 2(1+\lambda)T_3^1 - \cancel{\lambda T_3^0} = 0 \\
 i=3: 2(1+\lambda)T_3^1 - \cancel{\lambda T_4^1} - 2(1-\lambda)T_3^0 - \cancel{\lambda T_4^0} = \cancel{\lambda T_2^1} + \cancel{\lambda T_2^0} \\
 \qquad - \cancel{\lambda T_2^1} + 2(1+\lambda)T_4^1 - \cancel{\lambda T_4^0} = 0 \\
 i=4: -\cancel{\lambda T_3^1} + 2(1+\lambda)T_4^1 - \cancel{\lambda T_3^0} - 2(1-\lambda)T_4^0 = \cancel{\lambda T_5^1} + \cancel{\lambda T_5^0}
 \end{array}
 \quad \left. \begin{array}{c} IC \\ IC \\ BC \\ \text{Unknowns} \end{array} \right| \begin{array}{c} 2 \\ 1 \\ 1 \\ 4 \end{array}
 \end{array}$$

↳ 4 unknowns, 4 eqs.

↳ So now I see where the tridiagonal pattern is coming from.

↳ But this pattern will not hold for  $t=t'$ , right?

let's re-write the system for  $t=1$  (time =  $t'$ ); noting that all terms with  $T'$  were determined from the previous system ( $t=0$ ).  $\frac{5}{6}$

$$i=1: \quad 2(1+\lambda)T_1^2 - \lambda T_2^2 = \lambda T_0' + \lambda T_0^2 + 2(1-\lambda)T_1' + \lambda T_2' \\ i=2: \quad -\lambda T_1^2 + 2(1+\lambda)T_2^2 - \lambda T_3^2 = \lambda T_1' + \underbrace{2(1-\lambda)T_2' - \lambda T_3'}_{\substack{\text{determined @ previous} \\ \text{iteration}}} \\ i=3: \quad \dots$$

$$i=4: -\lambda T_3^2 + 2(1+\lambda)T_4^2 = \lambda T_5^1 + \lambda T_5^2 + \lambda T_3^1 + 2(1-\lambda)T_4^1$$

↳ Actually, the tridiagonal pattern holds.  $\overset{BC}{\checkmark}$

↳ Okay, so essentially, the RHS vector is the heat flow from the neighboring nodes (which the quantity is known/calc. from previous loop), we'll call this vector  $\{T_{\text{neighbors}}\}$

Is the general matrix form can be expressed as follows:

$$[\lambda] \{ T_{\text{node, time}} \} = \{ T_{\text{neighbors}} \}$$

↳ Alright! I love it when things make sense!

↳ Now I need to write a MATLAB program to solve this for all time.

$$[\lambda] = \begin{bmatrix} & T_1 & T_2 & T_3 & T_4 \\ 1 & 2(1+\lambda) & -\lambda & 0 & 0 \\ 2 & -\lambda & 2(1+\lambda) & -\lambda & 0 \\ 3 & 0 & -\lambda & 2(1+\lambda) & -\lambda \\ 4 & 0 & 0 & -\lambda & 2(1+\lambda) \end{bmatrix}$$

$$\{T_{\text{node}, \text{time}}\} = \left\{ \begin{array}{l} T_1^{t+1} \\ T_2^{t+1} \\ T_3^{t+1} \\ T_4^{t+1} \end{array} \right\}$$

Unknowns  $\rightarrow$  quantities for the next moment in time

$$\{T_{\text{neighbors}}\} = \left\{ \begin{array}{l} \lambda T_0^t + \lambda T_0^{t+1} + 2(1-\lambda)T_1^t + \lambda T_2^t \\ \lambda T_1^t + 2(1-\lambda)T_2^t - \lambda T_3^t \\ \lambda T_2^t + 2(1-\lambda)T_3^t + \lambda T_4^t \\ \lambda T_3^t + 2(1-\lambda)T_4^t + \lambda T_5^t + \lambda T_5^{t+1} \end{array} \right\}$$

BCs

(RHS temp. vector)

All quantities known for the current moment in time, with the exception of BCs.

MATLAB syntax:

$$\{T_{\text{node}, \text{time}}\} = [\lambda] \setminus \{T_{\text{neighbors}}\}$$

or  
 $\{T_{\text{rhs}}\}$   
 Right hand side

6/25/21

Crank-Nicolson method

6/6

What does the solution look like?

$$\{T_{\text{time, node}}\} = [\lambda] \setminus \{T_{\text{neighbors}}\}$$

↳ This vector output is only for one time step, and all nodes.  
I need to create a matrix that will hold all time steps  
for all nodes in the system.

$$[T_{\text{solution}}] = \begin{bmatrix} \text{transpose } \{T_{t^1, \text{nodes}}\} \\ \text{transpose } \{T_{t^2, \text{nodes}}\} \\ \vdots \\ \text{transpose } \{T_{t^f, \text{nodes}}\} \end{bmatrix}$$