

# Convergence\_WaveSystem\_Staggered\_SQUARE\_squares

November 26, 2018

## 1 Staggered scheme for Wave System on square meshes

### 1.1 The Wave System on the square

We consider the following Wave system with periodic boundary conditions

$$\begin{cases} \partial_t p + c^2 \nabla \cdot \vec{q} = 0 \\ \partial_t \vec{q} + \vec{\nabla} p = 0 \end{cases}.$$

The wave system can be written in matrix form

$$\partial_t \begin{pmatrix} p \\ \vec{q} \end{pmatrix} + \begin{pmatrix} 0 & c^2 \nabla \cdot \\ \vec{\nabla} & 0 \end{pmatrix} \begin{pmatrix} p \\ \vec{q} \end{pmatrix} = \begin{pmatrix} 0 \\ \vec{0} \end{pmatrix}$$

In  $d$  space dimensions the wave system is an hyperbolic system of  $d + 1$  equations

$$\partial_t U + \sum_{i=1}^d A_i \partial_{x_i} U = 0, \quad U = {}^t(p, \vec{q})$$

where the jacobian matrix is

$$A(\vec{n}) = \sum_{i=1}^d n_i A_i = \begin{pmatrix} 0 & c^2 \vec{n} \\ \vec{n} & 0 \end{pmatrix}, \quad \vec{n} \in \mathbb{R}^d.$$

has  $d + 1$  eigenvalues  $-c, 0, \dots, 0, c$ .

On the square domain  $\Omega = [0, 1] \times [0, 1]$  we consider the initial data

$$\begin{cases} p_0(x, y) = \text{constant} \\ q_{x0}(x, y) = \sin(\pi x) \cos(\pi y) \\ q_{y0}(x, y) = -\sin(\pi y) \cos(\pi x) \end{cases}.$$

The initial data  $(p_0, q_x, q_y)$  is a stationary solution of the wave system.

### 1.2 A 2D Staggered scheme for the Wave System

In  $2D$ , the linear wave system can be written using the cartesian coordinate system  $\vec{q} = (q_x, q_y)$  as

$$\begin{cases} \partial_t p + c^2(\partial_x q_x + \partial_y q_y) = 0 \\ \partial_t q_x + \partial_x p = 0 \\ \partial_t q_y + \partial_y p = 0. \end{cases}$$

We consider a 2D rectangular grid made of  $N = n_x \times n_y$  cells.

**The cells** are indexed by two integers  $i = 1, \dots, n_x$  ( $x$ -direction), and  $j = 1, \dots, n_y$  ( $y$ -direction).

**The pressure**  $p$  is discretised at the cell centers and is indexed with integer values  $p_{i,j}, i = 1, \dots, n_x, j = 1, \dots, n_y$ .

**The horizontal component**  $q_x$  of the momentum is discretised at the vertical cell interfaces and is indexed with a half-integer followed by an integer  $q_{i-\frac{1}{2},j}, i = 1, \dots, n_x, j = 1, \dots, n_y$ .

**The vertical component**  $q_y$  of the momentum is discretised at the horizontal cell interfaces and is indexed with an integer followed by a half-integer  $q_{i,j-\frac{1}{2}}, i = 1, \dots, n_x, j = 1, \dots, n_y$ .

The discrete equations read

$$\begin{cases} \partial_t p_{i,j} + c^2 \frac{q_{i+\frac{1}{2},j} - q_{i-\frac{1}{2},j}}{\Delta x} + c^2 \frac{q_{i,j+\frac{1}{2}} - q_{i,j-\frac{1}{2}}}{\Delta y} = 0 \\ \partial_t q_{i-\frac{1}{2},j} + \frac{p_{i,j} - p_{i-1,j}}{\Delta x} = 0 \\ \partial_t q_{i,j-\frac{1}{2}} + \frac{p_{i,j} - p_{i,j-1}}{\Delta y} = 0, \end{cases}$$

for  $i = 1, \dots, n_x, j = 1, \dots, n_y$

with the notations  $p_0 = p_{n_x}, q_{n_x+\frac{1}{2},j} = q_{\frac{1}{2},j}$  and  $q_{i,n_y+\frac{1}{2}} = q_{i,\frac{1}{2}}$  at the periodic boundaries.

We are therefore led to a linear system of  $3N = 3n_x \times n_y$  ODEs to solve.

### 1.3 The 2D Staggered scheme in matrix form

Define the unknown vector of the semi-discrete system as

$$\mathcal{U} = \begin{pmatrix} \mathcal{P} \\ \mathcal{Q}_x \\ \mathcal{Q}_y \end{pmatrix}, \quad \mathcal{P} = \begin{pmatrix} p_1 \\ \vdots \\ p_N \end{pmatrix}, \quad \mathcal{Q}_x = \begin{pmatrix} q_1^x \\ \vdots \\ q_N^x \end{pmatrix}, \quad \mathcal{Q}_y = \begin{pmatrix} q_1^y \\ \vdots \\ q_N^y \end{pmatrix},$$

where the global index for the pressure unknown for any cell  $(i, j)$  is  $p_k = p_{jn_x+i}$   
the global index for the  $x$ -momentum unknown for any vertical cell interface  $(i-\frac{1}{2}, j)$  is  $q_k^x = q_{jn_x+i}$   
and global index for the the  $y$ -momentum unknown for any horizontal cell interface  $(i, j-\frac{1}{2})$  is  $q_k^y = q_{jn_x+i}$ .

With these notations, the discrete equations read for  $k = 0, \dots, N$

$$\begin{cases} \partial_t p_{jn_x+i} + c^2 \frac{q_{jn_x+(i+1)\%n_x}^x - q_{jn_x+i}^x}{\Delta x} + c^2 \frac{q_{((j+1)\%n_y)n_x+i}^y - q_{jn_x+i}^y}{\Delta y} = 0 \\ \partial_t q_{jn_x+i}^x + \frac{p_{jn_x+i} - p_{jn_x+(i-1)\%n_x}}{\Delta x} = 0 \\ \partial_t q_{jn_x+i}^y + \frac{p_{jn_x+i} - p_{((j-1)\%n_y)n_x+i}}{\Delta y} = 0 \end{cases}.$$

The discrete staggered scheme takes the matrix form

$$\partial_t \mathcal{U} + \mathcal{M} \mathcal{U} = 0,$$

with

$$\begin{aligned}\mathcal{M} &= \begin{pmatrix} 0 & c^2 \mathcal{C}_x^{2d} & c^2 \mathcal{C}_y^{2d} \\ -{}^t \mathcal{C}_x^{2d} & 0 & 0 \\ -{}^t \mathcal{C}_y^{2d} & 0 & 0 \end{pmatrix} \in \mathcal{M}_{3n_x n_y}(\mathbb{R}), \quad \mathcal{C}_x^{2d}, \mathcal{C}_y^{2d} \in \mathcal{M}_{n_x n_y}(\mathbb{R}) \\ \mathcal{C}_x^{2d} &= \frac{1}{\Delta x} \begin{pmatrix} \mathcal{C}^{1d} & 0 & 0 \\ 0 & \ddots & 0 \\ 0 & 0 & \mathcal{C}^{1d} \end{pmatrix} \in \mathcal{M}_{n_x n_y}(\mathbb{R}), \quad \mathcal{C}^{1d} \in \mathcal{M}_{n_x}(\mathbb{R}) \\ \mathcal{C}_y^{2d} &= \frac{1}{\Delta y} \begin{pmatrix} -\mathbb{I}_{n_x} & \mathbb{I}_{n_x} & 0 & 0 \\ 0 & \ddots & \ddots & 0 \\ 0 & 0 & \ddots & \mathbb{I}_{n_x} \\ \mathbb{I}_{n_x} & 0 & 0 & -\mathbb{I}_{n_x} \end{pmatrix} \in \mathcal{M}_{n_x n_y}(\mathbb{R}).\end{aligned}$$

## 1.4 Staggered scheme stability

With the new unknown variable

$$\mathcal{V} = \begin{pmatrix} \frac{1}{c} \mathcal{P} \\ \mathcal{Q}_x \\ \mathcal{Q}_y \end{pmatrix},$$

which yields the discrete system

$$\partial_t \mathcal{V} + \mathcal{M}' \mathcal{V} = 0,$$

with the antisymmetric matrix :

$$\mathcal{M}' = \begin{pmatrix} 0 & c \mathcal{C}_x^{2d} & c \mathcal{C}_y^{2d} \\ -c {}^t \mathcal{C}_x^{2d} & 0 & 0 \\ -c {}^t \mathcal{C}_y^{2d} & 0 & 0 \end{pmatrix}.$$

Hence the norm of  $\mathcal{V}$  is constant :

$$\frac{1}{2} \partial_t \|\mathcal{V}\|^2 = {}^t \mathcal{V} \partial_t \mathcal{V} = -\frac{c}{\Delta x} {}^t \mathcal{V} \begin{pmatrix} 0 & c \mathcal{C}_x^{2d} & c \mathcal{C}_y^{2d} \\ -c {}^t \mathcal{C}_x^{2d} & 0 & 0 \\ -c {}^t \mathcal{C}_y^{2d} & 0 & 0 \end{pmatrix} \mathcal{V} = 0.$$

Since

$$\|\mathcal{U}(t)\| \leq \max \left\{ 1, \frac{1}{c} \right\} \|\mathcal{V}(t)\| = \max \left\{ 1, \frac{1}{c} \right\} \|\mathcal{V}(0)\|,$$

we deduce that  $\|\mathcal{U}\|$  is bounded and the scheme is therefore stable.

## 1.5 The script

```
#Condition initiale :
```

```
#Warning : the velocity is based on cells with the principle that the x component is  
pressure_field, velocity_field = initial_conditions_wave_system(my_mesh)
```

```
#Pas de temps
```

```
dt = cfl * dx_min / c0
```

```
#Matrice des systèmes linéaires
```

```

divMat=computeDivergenceMatrix(my_mesh,nbVoisinsMax,dt,test_bc)

# Construction du vecteur inconnu
Un=cdmath.Vector(nbCells*(dim+1))
for k in range(nbCells):
    Un[k*(dim+1)+0] = pressure_field[k]
    Un[k*(dim+1)+1] =rho0*velocity_field[k,0]
    Un[k*(dim+1)+2] =rho0*velocity_field[k,1]

# Création du solveur linéaire
LS=cdmath.LinearSolver(divMat,Un,iterGMRESMax, precision, "GMRES","ILU")

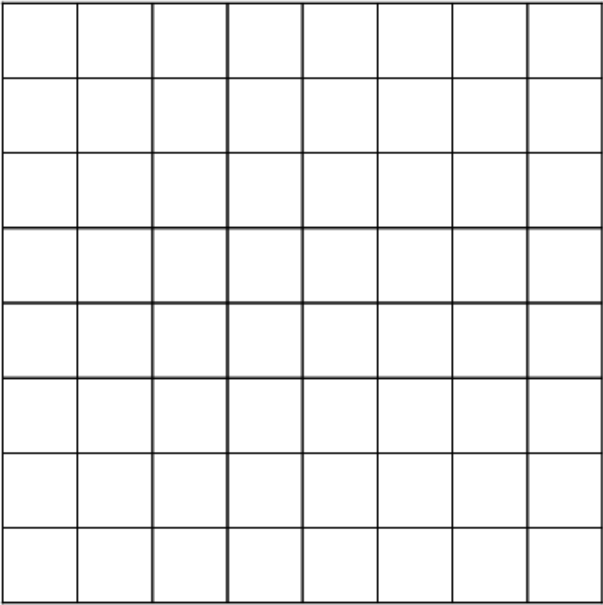
# Time loop
while (it<ntmax and time <= tmax and not isStationary):
    LS.setSndMember(Un)
    Un=LS.solve();
    Un.writeVTK

# Automatic postprocessing : save 2D picture and plot diagonal data
#=====
diag_data=VTK_routines.Extract_field_data_over_line_to_numpyArray(my_ResultField,[0
plt.legend()
plt.xlabel('Position on diagonal line')
plt.ylabel('Value on diagonal line')
if len(sys.argv) >1 :
    plt.title('Plot over diagonal line for Staggered Finite Volumes \n for Wave sys
    plt.plot(curv_abs, diag_data, label= str(nbCells)+ ' cells mesh')
    plt.savefig("FiniteVolumes2D_square_ResultField_"+str(nbCells)+ '_cells'+"_Plot

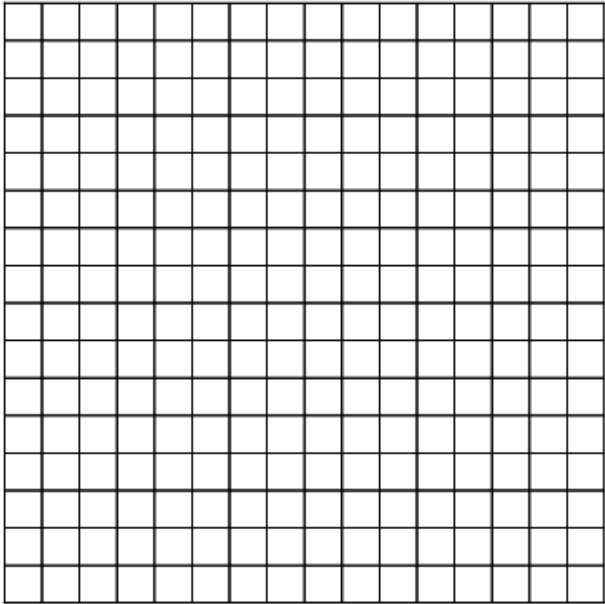
```

1.6 Numerical results

1.6.1 Cartesian meshes

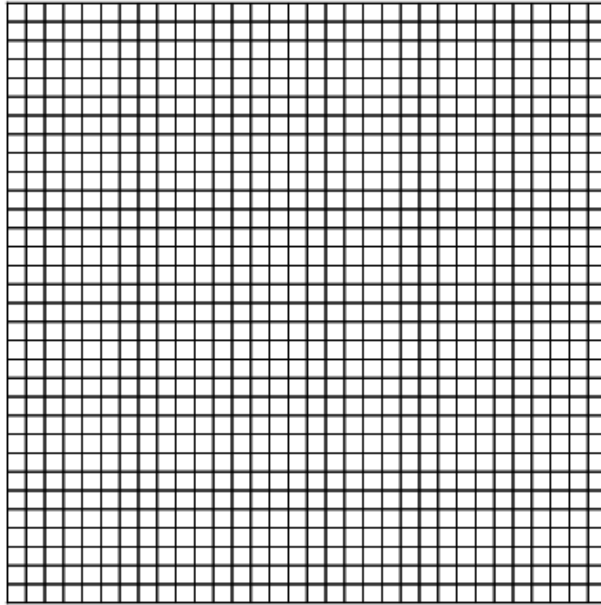


mesh 1 | mesh 2 | mesh 3 - | - - | -



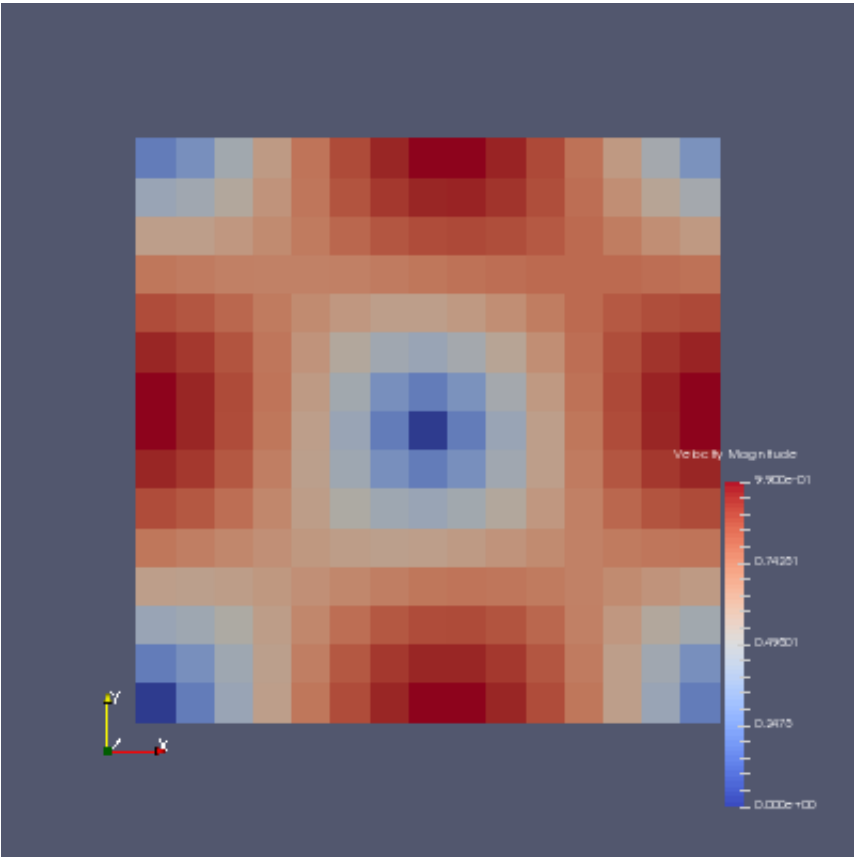
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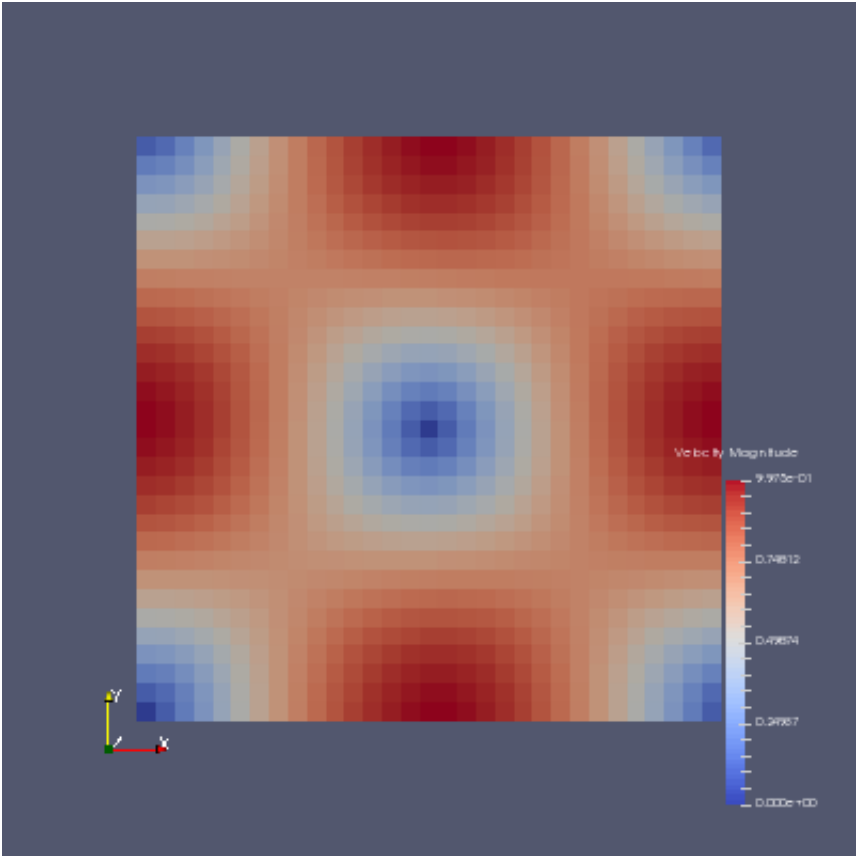




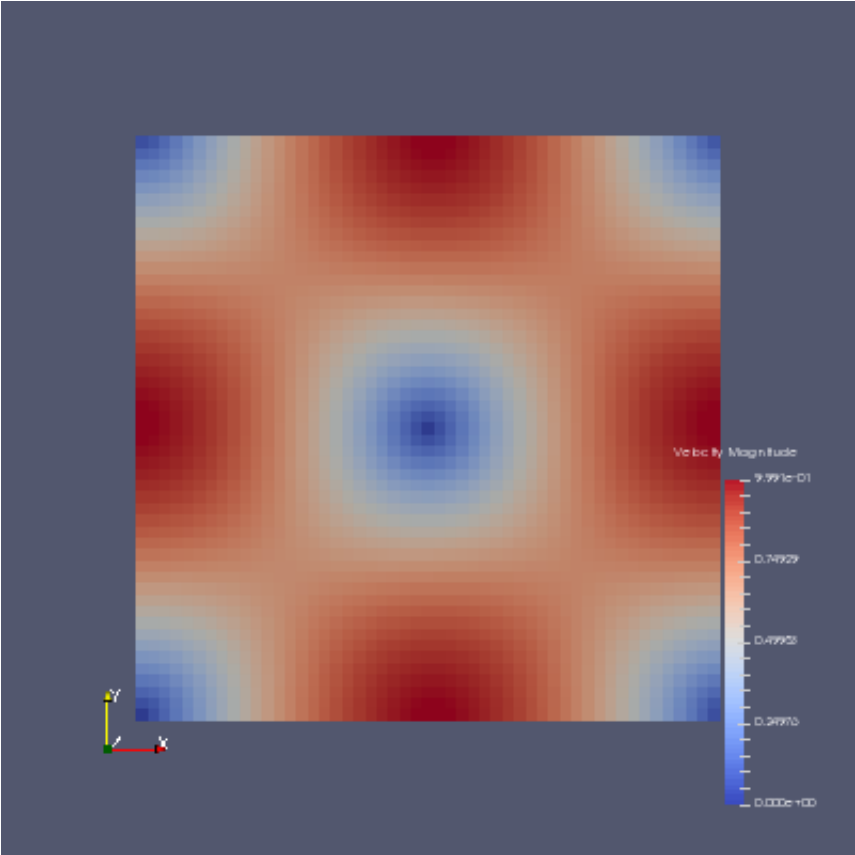
1.6.2 Velocity initial data (magnitude)



result 1 | result 2 | result 3 - | - - | -

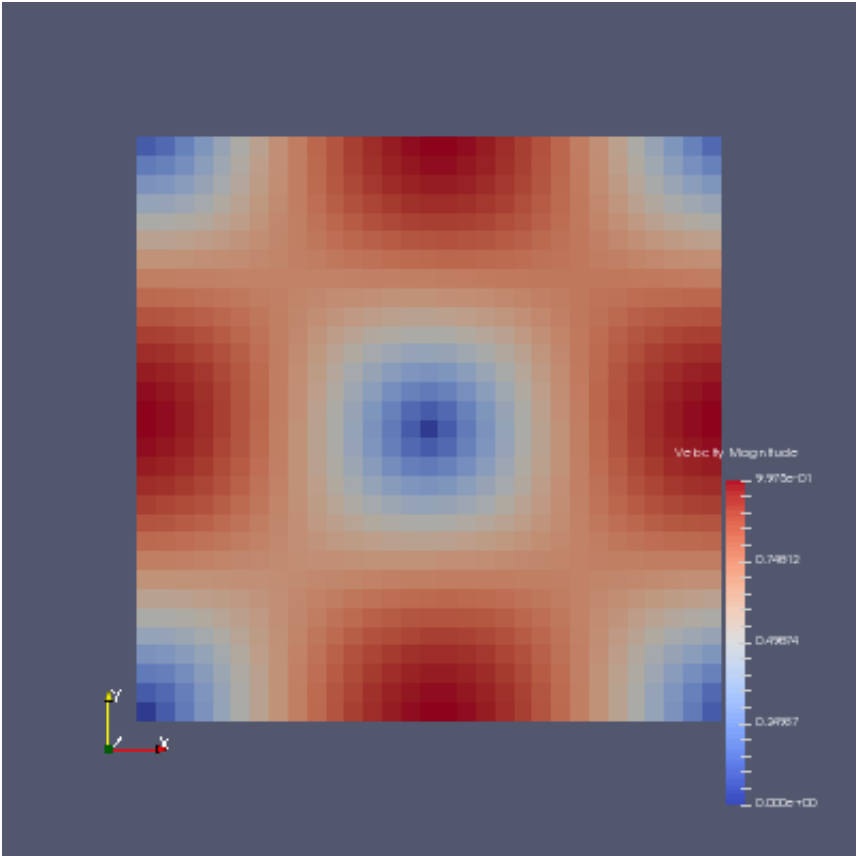
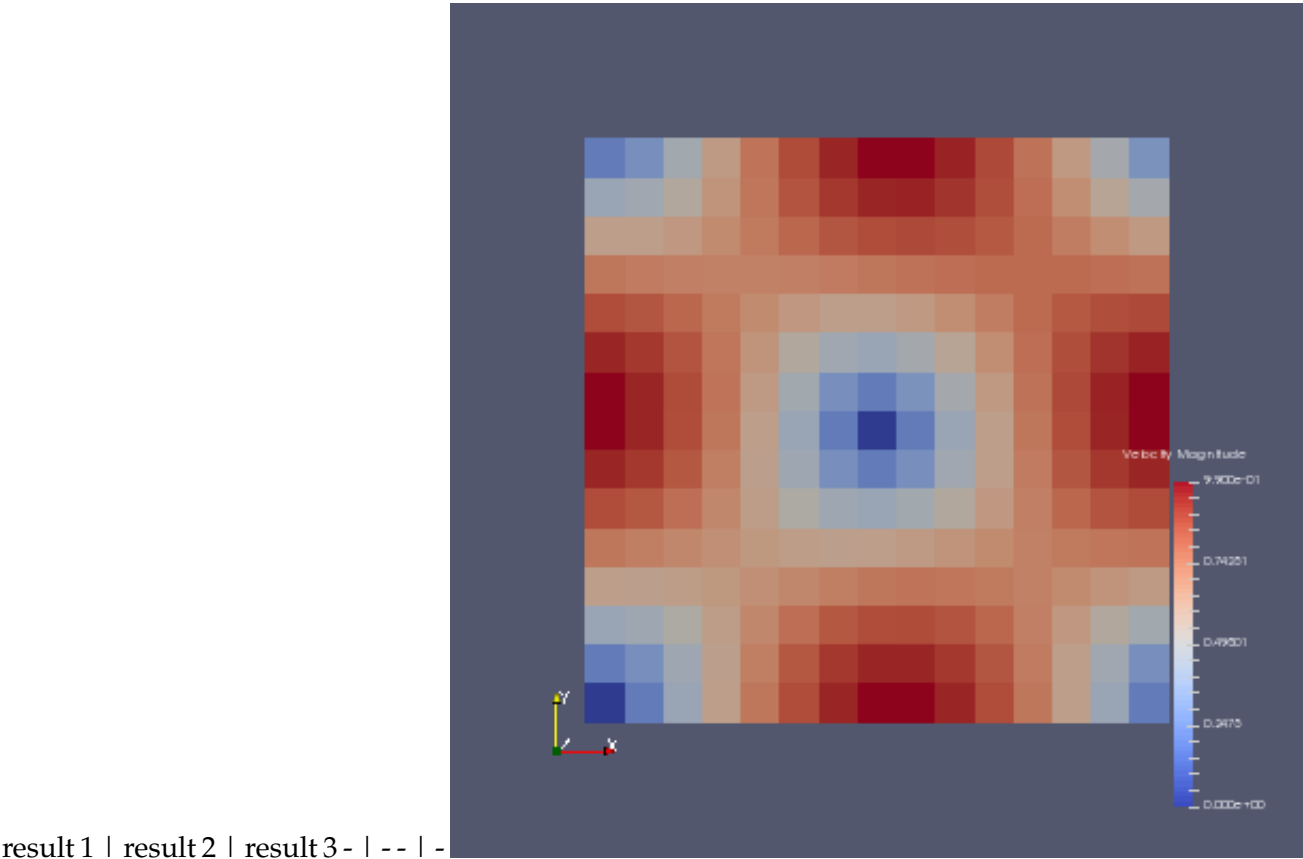


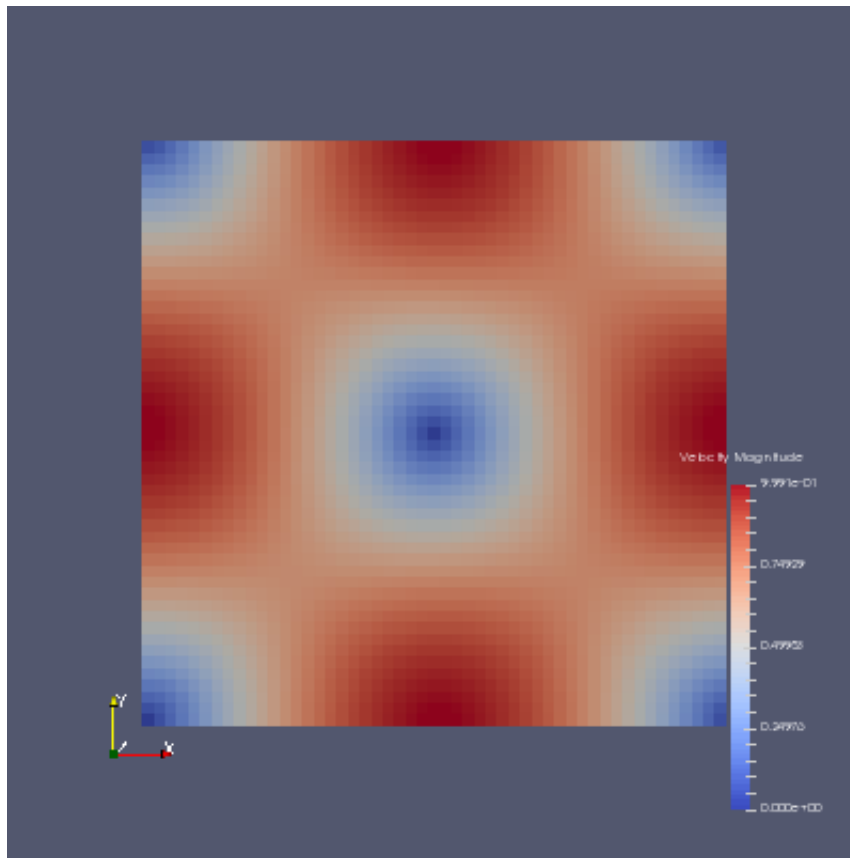






1.6.3 Stationary velocity (magnitude)





## 1.7 Convergence on stationary velocity

