

Computer assisted proofs in differential equations

Lecture 4 - Dec. 13, 2019

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Systems of first order differential equations

- ▶ Consider the IVP

$$\dot{x} = f(x), \quad x(0) = x_0 \in \mathbb{R}^N,$$

where $f = (f_1, \dots, f_N)$ is a polynomial nonlinearity.

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where $\alpha \in \mathbb{N}^N$ is a multi-index, $|\alpha| = \alpha_1 + \dots + \alpha_N$, and $x^\alpha = x_1^{\alpha_1} \cdots x_n^{\alpha_N}$.

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- ▶ We assume that time rescaling, if needed, is absorbed implicitly in the coefficients c_α .
- ▶ The goal is to obtain solutions $x(t) = (x_1(t), \dots, x_N(t))$ with power series expansion

$$x_n(t) = \sum_{k=0}^{\infty} (a_n)_k t^k = \sum_{k=0}^{\infty} a_{n,k} t^k, \quad n = 1, \dots, N.$$

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- ▶ The n -th component of $f(x(t))$ is of the form

$$f_n(x(t)) = \sum_{k \geq 0} (\phi_n(a))_k t^k,$$

where the k -th Taylor coefficient $(\phi_n(a))_k$ of $f_n(x(t))$ is a combination of Cauchy products involving the terms of a .

- ▶ To make $\phi_n(a)$ explicit, we use the notation

$$a^k := a * a^{k-1} = \underbrace{a * \cdots * a}_{k \text{ times}}, \quad \text{with} \quad a^1 = a.$$

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- ▶ It is also convenient to denote

$$(a^0)_k := \delta_{k0}.$$

This is natural in view of $1 = \sum_{k=0}^{\infty} \delta_{k0} t^k$. Also, we now have $a = a * a^0$.

Systems of first order differential equations

- We then have

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- We need to solve

$$\begin{aligned}\dot{x}_n(t) - f_n(x(t)) &= \sum_{k \geq 1} k a_{n,k} t^{k-1} - \sum_{k \geq 0} \phi_n(a)_k t^k \\ &= \sum_{k \geq 1} (k a_{n,k} - \phi_n(a)_{k-1}) t^{k-1} = 0.\end{aligned}$$

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- We define $F_n(a) = \{F_n(a)_k\}_{k \geq 0}$, for each $n = 1, \dots, N$, component-wise by

$$(F_n(a))_k := \begin{cases} a_{n,0} - x_{0,n} & \text{if } k = 0, \\ k a_{n,k} - \phi_n(a)_{k-1} & \text{if } k \geq 1. \end{cases}$$

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- ▶ In particular, a graphically intuitive representation of A is given by

$$A = \left[\begin{array}{cc|cc|c|cc|cc} A_{1,1}^{(N)} & 0 & A_{1,2}^{(N)} & 0 & \cdots & \cdots & A_{1,n}^{(N)} & 0 \\ 0 & \Lambda_N & 0 & 0 & \cdots & \cdots & 0 & 0 \\ \hline A_{2,1}^{(N)} & 0 & A_{2,2}^{(N)} & 0 & & & A_{2,n}^{(N)} & 0 \\ 0 & 0 & 0 & \Lambda_N & & & 0 & 0 \\ \hline \vdots & \vdots & & & \ddots & & \vdots & \vdots \\ \vdots & \vdots & & & & \ddots & \vdots & \vdots \\ \hline A_{n,1}^{(N)} & 0 & A_{n,2}^{(N)} & 0 & \cdots & \cdots & A_{n,n}^{(N)} & 0 \\ 0 & 0 & 0 & 0 & \cdots & \cdots & 0 & \Lambda_N \end{array} \right].$$

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- ▶ Where

$$(\Lambda_K a)_k = \begin{cases} 0 & \text{if } 0 \leq k \leq K, \\ \frac{1}{k} a_k & \text{if } k \geq K + 1. \end{cases}$$

Example - Lorenz system

- ▶ Consider the vector field

$$f(x) = \begin{pmatrix} \sigma(x_2 - x_1) \\ \rho x_1 - x_2 - x_1 x_3 \\ -\beta x_3 + x_1 x_2 \end{pmatrix},$$

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- ▶ Let $a = (a_1, a_2, a_3) \in (\ell_\nu^1)^3$, and

$$(\phi_j(a))_k = \begin{cases} \sigma((a_2)_k - (a_1)_k), & j = 1 \\ \rho(a_1)_k - (a_2)_k - (a_1 * a_3)_k, & j = 2 \\ -\beta(a_3)_k + (a_1 * a_2)_k, & j = 3 \end{cases}$$

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- ▶ After rescaling time by L we solve recursively

$$\begin{pmatrix} (a_1)_k \\ (a_2)_k \\ (a_3)_k \end{pmatrix} = \frac{L}{k} \begin{pmatrix} \sigma((a_2)_{k-1} - (a_1)_{k-1}) \\ \rho(a_1)_{k-1} - (a_2)_{k-1} - (a_1 * a_3)_{k-1} \\ -\beta(a_3)_{k-1} + (a_1 * a_2)_{k-1} \end{pmatrix}.$$

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$$Y_0^{(2)} = \nu \sum_{k=N}^{2N} \frac{L}{k+1} |\rho(\bar{a}_1)_k - (\bar{a}_2)_k - (\bar{a}_1 * \bar{a}_3)_k| \nu^k,$$

$$Y_0^{(3)} = \nu \sum_{k=N}^{2N} \frac{L}{k+1} |-\beta(\bar{a}_3)_k + (\bar{a}_1 * \bar{a}_2)_k| \nu^k,$$

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- ▶ We also define

$$Z(r) := \frac{\nu}{N+1} \left(\max \left\{ z_0^{(1)}, z_0^{(2)}, z_0^{(3)} \right\} + \max \left\{ z_1^{(2)}, z_1^{(3)} \right\} r \right),$$

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- ▶ where

$$z_0^{(1)} = 2L|\sigma|$$

$$z_0^{(2)} = L(|\rho| + 1 + \|\bar{a}_1\|_{1,\nu} + \|\bar{a}_3\|_{1,\nu})$$

$$z_1^{(2)} = 2L$$

$$z_0^{(3)} = L(|\beta| + \|\bar{a}_1\|_{1,\nu} + \|\bar{a}_2\|_{1,\nu})$$

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and

$$\sup_{u \in \overline{B_r(0)}} \|DT(u)\| \leq Z(r).$$

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- ▶ Hence we have the radii polynomial

$$p(r) = Z(r)r - r + Y_0$$

and we can apply the contraction mapping radii polynomial theorem.

Example - Lorenz system

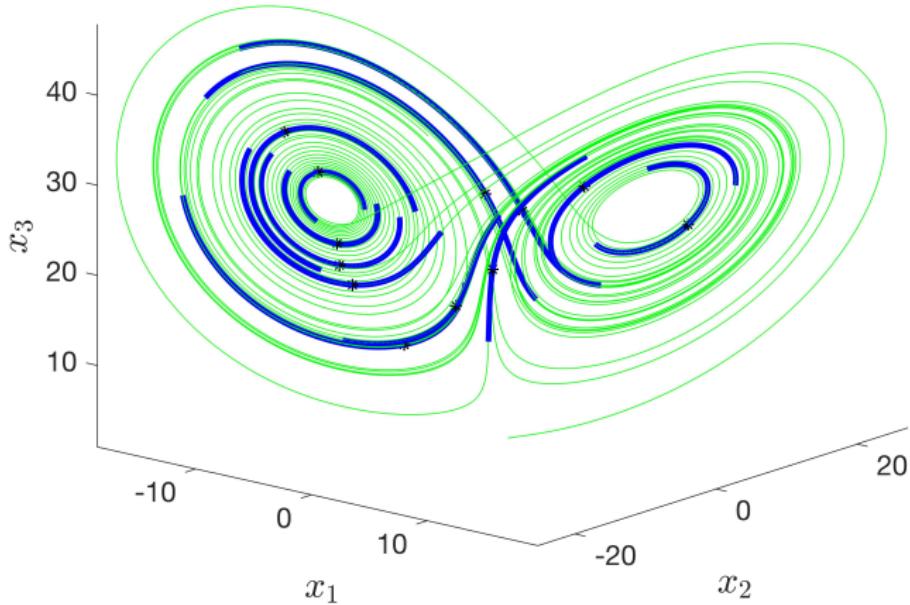


Figure: The green curve is a numerical integration of a single initial condition for 50 time units. The blue curves represent one step of Taylor integration with order $N = 200$ for different initial conditions, plotted for times $t \in (-0.2, 0.2)$.

Periodic solutions of vector fields

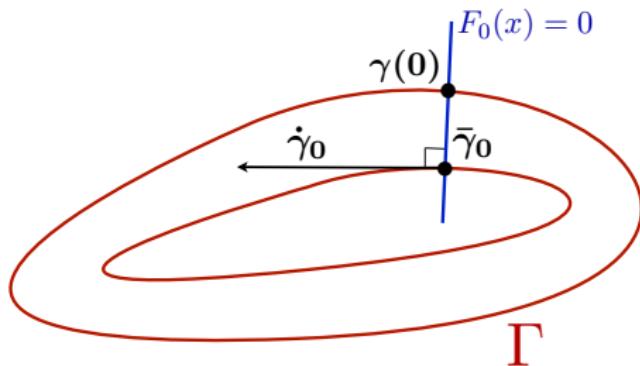
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- Given a periodic solution $\gamma(t)$ of an ODE $\dot{x} = f(x)$, then $\gamma_\theta(t) := \gamma(t + \theta)$ is also a periodic solution for all $\theta \in \mathbb{R}$.
- We make use of the so-called **Poincaré phase condition** in order to isolate $\gamma(t)$ in function space

$$\dot{\gamma}_0 \cdot (\bar{\gamma}_0 - \gamma(0)) = 0,$$

for fixed $\bar{\gamma}_0 = (\bar{\gamma}_{0,1}, \dots, \bar{\gamma}_{0,n}) \in \mathbb{R}^n$ and $\dot{\gamma}_0 = (\dot{\gamma}_{0,1}, \dots, \dot{\gamma}_{0,n}) \in \mathbb{R}^n$.



Periodic solutions of vector fields

- ▶ Let $\gamma: \mathbb{R} \rightarrow \mathbb{R}^n$ be a $\frac{2\pi}{\omega}$ -periodic solution to $\dot{x} = f(x)$, with $f: \mathbb{R}^n \rightarrow \mathbb{R}^n$.

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- ▶ For each $j = 1, \dots, n$, denote by $a_j := \{(a_j)_k\}_{k \in \mathbb{Z}}$ the sequence of Fourier coefficients of γ_j , that is

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$$f_j(\gamma(t)) = \sum_{k \in \mathbb{Z}} (\phi_j(a))_k e^{ik\omega t},$$

where the k -th Fourier coefficient $(\phi_j(a))_k$ of $f_j(\gamma(t))$ is a combination of discrete convolutions involving the terms of a .

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Periodic solutions of vector fields

- ▶ The assumption that γ is a periodic solution to the ODE is equivalent to

$$\dot{\gamma}_j(t) - f_j(\gamma(t)) = \sum_{k \in \mathbb{Z}} (ik\omega(a_j)_k - (\phi_j(a))_k) e^{ik\omega t} = 0,$$

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- ▶ In terms of the Fourier coefficients, the Poincaré phase condition is equivalent to

$$\bar{\gamma}_0 \cdot f(\bar{\gamma}_0) - \sum_{n=1}^N f_n(\bar{\gamma}_0) \sum_{k \in \mathbb{Z}} a_{n,k} = 0.$$

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- ▶ However this condition can be relaxed, by fixing $k_0 \in \mathbb{N}$, to the following

$$F_0(x) := \sum_{j=1}^n \bar{\gamma}_{0,j} \dot{\gamma}_{0,j} - \sum_{j=1}^n \sum_{|k| \leq k_0} \dot{\gamma}_{0,j}(a_j)_k = 0.$$

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- ▶ Considering $X := \mathbb{C} \times (\ell_\nu^1)^n$ with norm

$$\|x\|_X = \|(\omega, a)\|_X := \max \{|\omega|, \|a_1\|_{1,\nu}, \|a_2\|_{1,\nu}, \dots, \|a_n\|_{1,\nu}\}.$$

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- ▶ For each $j = 1, \dots, n$, let

$$F_j(x) = \{(F_j(x))_k\}_{k \in \mathbb{Z}}.$$

- ▶ We need to solve

$$F(x) = \begin{pmatrix} F_0(x) \\ F_1(x) \\ \vdots \\ F_n(x) \end{pmatrix} = 0$$

in the Banach space $X = \mathbb{C} \times (\ell_\nu^1)^n$.

Example - Lorenz system

- ▶ Consider again the Lorenz system is given by

$$\begin{cases} \dot{x}_1 = \sigma(x_2 - x_1) \\ \dot{x}_2 = \rho x_1 - x_2 - x_1 x_3 \\ \dot{x}_3 = -\beta x_3 + x_1 x_2, \end{cases}$$

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- The functional F takes the form

$$F_j(x) = \begin{cases} \sum_{j=1}^3 \bar{\gamma}_{0,j} \dot{\gamma}_{0,j} - \sum_{j=1}^3 \dot{\gamma}_{0,j} \sum_{|k| \leq k_0} (a_j)_k, & j = 0 \\ \{ik\omega(a_1)_k - (\phi_1(a))_k\}_{k \in \mathbb{Z}}, & j = 1 \\ \{ik\omega(a_2)_k - (\phi_2(a))_k\}_{k \in \mathbb{Z}}, & j = 2 \\ \{ik\omega(a_3)_k - (\phi_3(a))_k\}_{k \in \mathbb{Z}}, & j = 3, \end{cases}$$

where

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Example - Lorenz system

For the classical parameter values $\rho = 28$, $\beta = -8/3$ and $\sigma = 10$ we prove the existence of the periodic orbits below (with $\nu = 1.01$, except for the orbit AAAAAAABBABBBAAAA where $\nu = 1.009$).

orbit itinerary	period	N	\mathcal{I}	
AB	1.5587	80	$[3.0112 \times 10^{-14}$	0.0078691]
AAB	2.3059	110	$[6.0866 \times 10^{-13}$	0.0041978]
ABB	2.3059	110	$[5.1273 \times 10^{-13}$	0.0031017]
AAAB	3.0236	140	$[2.1023 \times 10^{-12}$	0.0018607]
AABB	3.0843	140	$[2.6987 \times 10^{-13}$	0.0015571]
ABBB	3.0236	140	$[2.016 \times 10^{-12}$	0.0014146]
AAAAB	3.7256	170	$[1.6132 \times 10^{-12}$	0.00089004]
AAABB	3.8203	170	$[2.6086 \times 10^{-12}$	0.00059818]
AABAB	3.8695	170	$[3.5391 \times 10^{-12}$	0.00064646]
AABBB	3.8203	170	$[2.1882 \times 10^{-12}$	0.00055587]
AAAABABB	6.0824	300	$[2.7194 \times 10^{-11}$	6.8867×10^{-5}
AABBAABB	8.957	560	$[2.6555 \times 10^{-11}$	6.4514×10^{-5}
AAAAAAABBABBBAAAA	11.9973	870	$[1.2693 \times 10^{-10}$	3.0684×10^{-5}

Example - Lorenz system

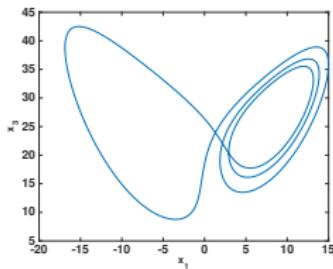
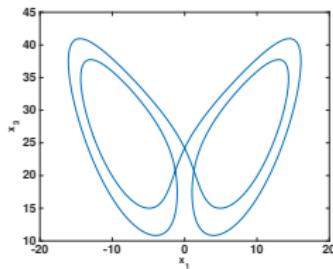
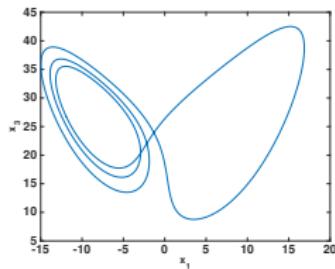
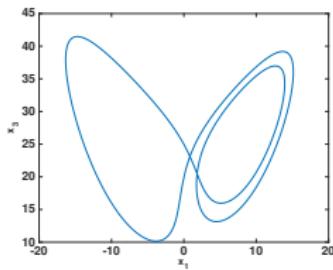
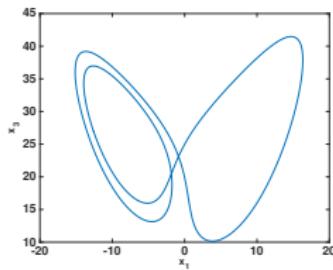
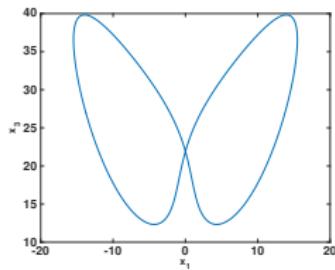


Figure: Top left: AB, Top center: AAB, Top right: ABB, Bottom left: AAAB, Bottom center: AABB, Bottom right: ABBB.

Example - Lorenz system

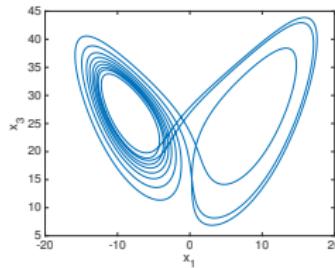
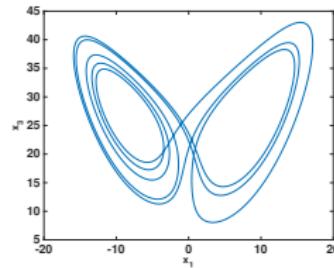
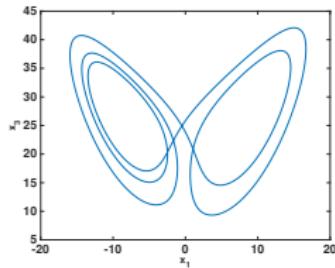
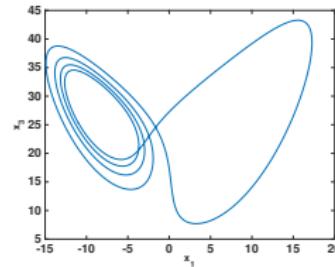
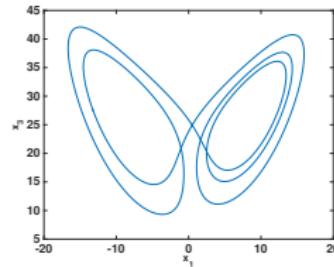
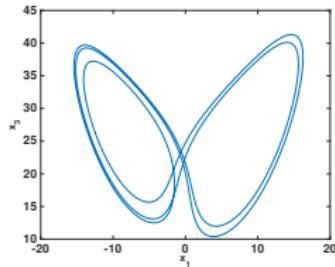


Figure: **Top left:** AABAB, **Top center:** AABB, **Top right:** AAAAB, **Bottom left:** AAABB, **Bottom center:** AAAABABB, **Bottom right:** AAAAABAAAABB.

Example - Lorenz system

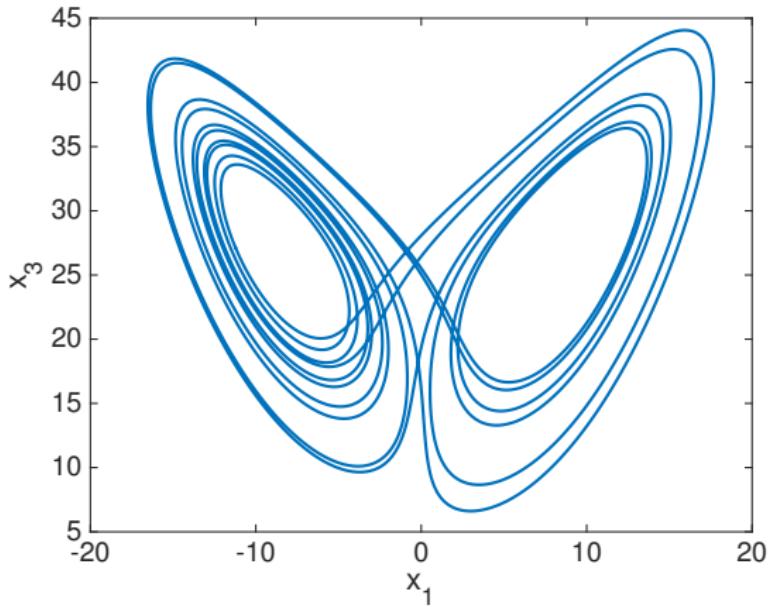


Figure: A periodic orbit AAAAABBBBAAAA with period $p \approx 11.9973$. The C^0 -error between the true solution and the numerical approximation is $r = 1.2693 \times 10^{-10}$. The proof was obtained with $N = 870$ and $\nu = 1.009$.

The Parameterization Method for (Un)Stable Manifolds

Theorem ((Un)Stable Manifold Theorem)

Consider the ODE

$$\dot{x} = f(x), \quad x \in \mathbb{R}^n$$

with $f \in C^k(\mathbb{R}^n)$, $k \in \{1, 2, 3, \dots, \infty\} \cup \{\omega\}$ for which $f(0) = 0$. Assume

$$Df(0) = \begin{bmatrix} A^u & 0 \\ 0 & A^s \end{bmatrix}$$

is hyperbolic where $A^u \in M_d(\mathbb{R})$, $A^s \in M_{n-d}(\mathbb{R})$, and the real parts of the all the eigenvalues of A^u and A^s are positive and negative, respectively. Then, there exists a neighborhood $U = U^u \times U^s \subset \mathbb{R}^d \times \mathbb{R}^{n-d}$ of the origin and C^k functions $P^u: U^u \rightarrow \mathbb{R}^n$ and $P^s: U^s \rightarrow \mathbb{R}^n$ that are tangent to $\mathbb{R}^d \times 0$ and $0 \times \mathbb{R}^{n-d}$ at the point 0, respectively, such that

$$W_{loc}^u(0) = P^u(U^u) \quad \text{and} \quad W_{loc}^s(0) = P^s(U^s).$$

The Parameterization Method

- ▶ Denote the diagonal matrix of stable eigenvalues by

$$\Lambda = \begin{bmatrix} \lambda_1 & \dots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \dots & \lambda_m \end{bmatrix}$$

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- ▶ Our goal is to construct an analytic function $P: B^m \rightarrow \mathbb{R}^n$ such that $P(B^m) = W_{loc}^s(\tilde{x})$.

The Parameterization Method

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- ▶ and

$$\varphi(t, P(\sigma)) = P(e^{\Lambda t}\sigma),$$

for all $\sigma \in B^m$.

The Parameterization Method

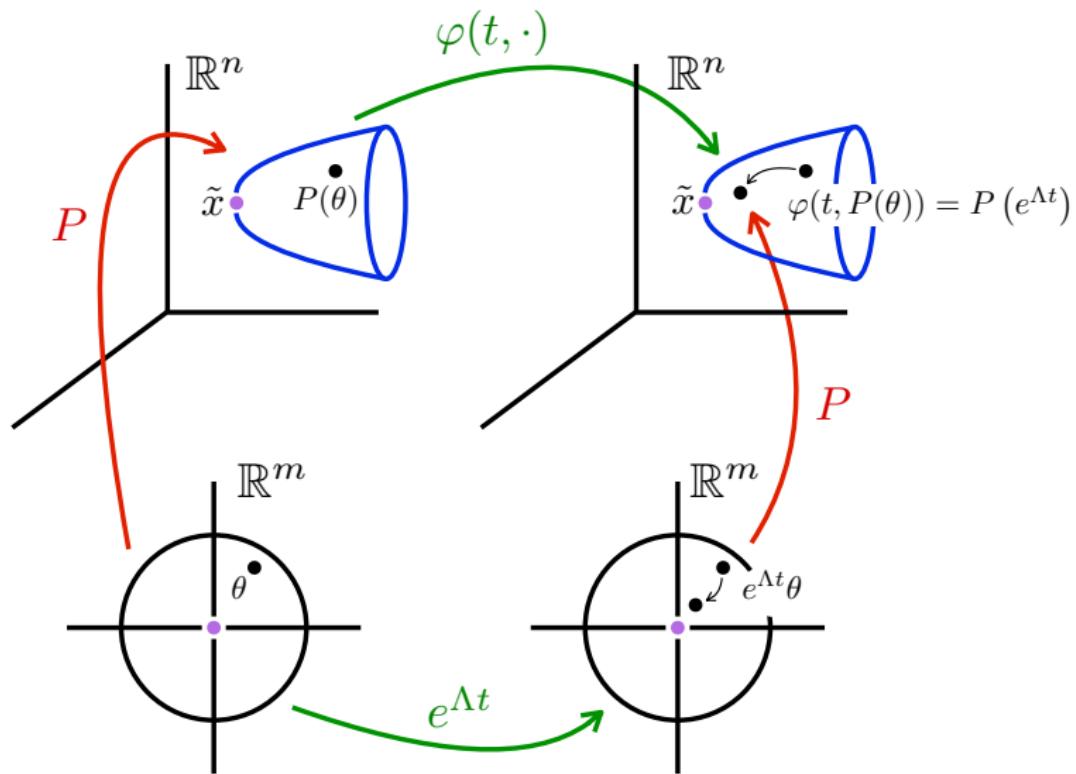


Figure: Schematic of the Parameterization Method

The Parameterization Method

Lemma (Parameterization Lemma)

Let $P: B^m \subset \mathbb{R}^m \rightarrow \mathbb{R}^n$ be a smooth function with

$$P(0) = \tilde{x} \quad \text{and} \quad DP(0) = A_0.$$

Then $P(\sigma)$ satisfies the conjugacy relationship above if and only if P is a solution of the partial differential equation

$$\lambda_1\sigma_1 \frac{\partial}{\partial\sigma_1} P(\sigma_1, \dots, \sigma_m) + \dots + \lambda_m\sigma_m \frac{\partial}{\partial\sigma_m} P(\sigma_1, \dots, \sigma_m) = f(P(\sigma_1, \dots, \sigma_m))$$

for all $\sigma = (\sigma_1, \dots, \sigma_m) \in B^m$.

The invariance equation above can be written more concisely as

$$DP(\sigma)\Lambda\sigma = f(P(\sigma)), \quad \text{for } \sigma \in B^m.$$

Example - Lorenz system

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- ▶ Imposing the first order constraints leads to

$$p_0 = \begin{pmatrix} a_0 \\ b_0 \\ c_0 \end{pmatrix} = \tilde{x}, \quad \text{and} \quad p_1 = \begin{pmatrix} a_1 \\ b_1 \\ c_1 \end{pmatrix} = \xi.$$

Example - Lorenz system

- ▶ To understand the terms with $n \geq 2$ first note that, by differentiating, we get that

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- ▶ For the Lorenz system, we see that

$$f(P(s)) = \sum_{n=0}^{\infty} \begin{pmatrix} \sigma(b_n - a_n) \\ \rho a_n - b_n - \sum_{k=0}^n a_{n-k} c_k \\ -\beta c_n + \sum_{k=0}^n a_{n-k} b_k \end{pmatrix} s^n.$$

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- Upon matching like powers we have

$$n \lambda \begin{pmatrix} a_n \\ b_n \\ c_n \end{pmatrix} = \begin{pmatrix} \sigma(b_n - a_n) \\ \rho a_n - b_n - \sum_{k=0}^n a_{n-k} c_k \\ -\beta c_n + \sum_{k=0}^n a_{n-k} b_k \end{pmatrix},$$

for all $n \geq 2$.

Example - Lorenz system

- The sums contain “hidden” terms of order n , and we write

$$\sum_{k=0}^n a_{n-k} c_k = a_0 c_n + c_0 a_n + \sum_{k=1}^{n-1} a_{n-k} c_k,$$

and

$$\sum_{k=0}^n a_{n-k} b_k = a_0 b_n + b_0 a_n + \sum_{k=1}^{n-1} a_{n-k} b_k,$$

where the sums now depend only on terms of order strictly less than n .

- Isolating terms of order n on the left hand side gives

$$\begin{pmatrix} \sigma(b_n - a_n) - n\lambda a_n \\ \rho a_n - b_n - a_0 c_n - c_0 a_n - n\lambda b_n \\ -\beta c_n + a_0 b_n + b_0 a_n - n\lambda c_n \end{pmatrix} = \begin{pmatrix} 0 \\ \sum_{k=1}^{n-1} a_{n-k} c_k \\ -\sum_{k=1}^{n-1} a_{n-k} b_k \end{pmatrix}.$$

Example - Lorenz system

- ▶ Noticing that we can rewrite the left hand side as

$$\begin{pmatrix} \sigma(b_n - a_n) - n\lambda a_n \\ \rho a_n - b_n - a_0 c_n - c_0 a_n - n\lambda b_n \\ -\beta c_n + a_0 b_n + b_0 a_n - n\lambda c_n \end{pmatrix} =$$
$$\begin{bmatrix} -\sigma - n\lambda & \sigma & 0 \\ \rho - c_0 & -1 - n\lambda & -a_0 \\ b_0 & a_0 & -\beta - n\lambda \end{bmatrix} \begin{pmatrix} a_n \\ b_n \\ c_n \end{pmatrix} = [Df(\tilde{x}) - n\lambda I] p_n.$$

- ▶ We can express the system above in matrix form as

$$[Df(\tilde{x}) - n\lambda I] p_n = R_n,$$

where

$$R_n := \begin{pmatrix} 0 \\ \sum_{k=1}^{n-1} a_{n-k} c_k \\ -\sum_{k=1}^{n-1} a_{n-k} b_k \end{pmatrix}.$$

Example - Lorenz system

- ▶ Noticing that $n\lambda < \lambda < 0$ is never an eigenvalue for $n \geq 2$, since λ is the only stable eigenvalue, we have the recursive formula

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- ▶ Hence we have the N -th order approximation for the stable manifold

$$P^N(s) = \sum_{n=0}^N p_n s^n.$$

Example - Lorenz system

- ▶ For 2-dimensional manifolds the formulas becomes

$$[Df(\tilde{x}) - (m\lambda_1 + n\lambda_2)I] p_{mn} = R_{mn},$$

where

$$R_{mn} = \begin{pmatrix} 0 \\ \sum_{k=0}^m \sum_{l=0}^n \hat{\delta}_{kl}^{mn} a_{(m-k)(n-l)} c_{kl} \\ - \sum_{k=0}^m \sum_{l=0}^n \hat{\delta}_{kl}^{mn} a_{(m-k)(n-l)} b_{kl} \end{pmatrix}.$$

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- The N -th order numerical approximation is

$$P^N(s_1, s_2) = \sum_{n=0}^N \sum_{m=0}^n p_{n-m, m} s_1^{n-m} s_2^m.$$